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## **Komplexe Algebraische Geometrie**

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### **Workshop: Komplexe Algebraische Geometrie**

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## Abstracts

### On the Quantisation of Completely Integrable Hamiltonian Systems

DUCO VAN STRATEN

(joint work with Mauricio Garay)

Classical mechanics is described by a hamiltonian function that induces a flow in a phase space. The mathematical model is that of a symplectic manifold  $M$ , where the symplectic form  $\omega$  defines an identification  $\phi$  between the cotangent bundle  $\Omega_M$  and the tangent bundle  $\Theta_M$ ; a function  $H$  on  $M$  defines a flow by integrating the hamiltonian vector field  $\phi(dH)$ , [1].

We consider the case  $M = \mathbb{C}^{2n}$  with canonical coordinates  $(p_1, \dots, p_n, q_1, \dots, q_n)$  such that  $\omega = \sum_{i=1}^n dp_i \wedge dq_i$ . The dynamics is described by the Hamilton equations

$$\dot{p}_i = -\partial H / \partial q_i, \quad \dot{q}_i = \partial H / \partial p_i$$

where the hamiltonian  $H$  is a function of the  $2n$  coordinates  $(p, q)$ . The time derivative of an arbitrary function is then given by  $\dot{F} = \{H, F\}$ , where

$$\{F, G\} = \sum_{i=1}^n \frac{\partial F}{\partial p_i} \frac{\partial G}{\partial q_i} - \frac{\partial G}{\partial p_i} \frac{\partial F}{\partial q_i}$$

is the *Poisson-bracket* of  $F$  and  $G$ .  $F$  is called a *conserved quantity* if  $\dot{F} = 0$ , or, what is the same  $F$  *Poisson commutes* with  $H$ ,  $\{F, H\} = 0$ .

In general we call  $I_1, I_2, \dots, I_n \in R := \mathbb{C}[p_1, \dots, p_n, q_1, \dots, q_n]$  which are functionally independent and with  $\{I_i, I_j\} = 0$  for all  $i, j$  a *(polynomial classical) integrable system*. Although they are rare and hard to construct, several examples are known, like the tops of Euler, Lagrange, Kovalevskaya; special cases of the Henon-Heiles system, the Calogero-Moser systems, to mention a few. In many cases the fibres of the map  $I := (I_1, \dots, I_n) : \mathbb{C}^{2n} \rightarrow \mathbb{C}^n$  are affine pieces of abelian varieties, see [6] for an overview. In algebraic geometry one encounters the integrable Hitchin system, the systems of Beauville-Mukai, which the global situation of Lagrangian fibrations on hyperkähler manifolds.

In their 1925 paper [2], Born and Jordan realised that quantum mechanics is a *non-commutative deformation of classical mechanics*: the ring  $R = \mathbb{C}[p, q]$  is replaced by the non-commutative Heisenberg algebra  $Q := \mathbb{C} \langle \hbar, p, q \rangle$  with the relation

$$pq - qp = \hbar, \quad \hbar := \frac{h}{2\pi i}, \quad (h = 6.10^{-34} \text{ Js})$$

$\hbar$  should be considered as a central element, and classical mechanics is recovered by putting  $\hbar = 0$ . Indeed, one can consider  $R$  as a quotient of  $Q$ :  $Q/\hbar Q = R$ . It

was observed by Dirac, that the Poisson-bracket is recovered from the commutator via

$$\{f, g\} := \frac{1}{\hbar}[F, G] \quad \text{mod } \hbar Q$$

**Question:** *Given a integrable system  $I_1, \dots, I_n \in R$ , do there exist  $J_1, \dots, J_n \in Q$  such that  $[J_i, J_j] = 0$  and  $J_i = I_i \quad \text{mod } \hbar$ ?*

If we can find such commuting  $J_1, \dots, J_n$ , we will say the system is *quantum completely integrable*. We have no general answer to this question, but for many integrable systems explicit quantisations are known. The quantisation of the Hitchin system plays a central role in the *geometric Langlands program* [3].

It is natural to work order by order in  $\hbar$  and put  $Q_k := Q/\hbar^k Q$  and replace  $Q$  by the completion  $\hat{Q} = \lim_{\leftarrow k} Q_k$ . We consider the polynomial ring  $A = \mathbb{C}[I_1, \dots, I_n] \xrightarrow{\iota_1} Q_1 = R$  which we try to lift  $\iota_1$  order by order to  $A \xrightarrow{\iota_2} Q_2, \dots, A \xrightarrow{\iota_k} Q_k$ . The Poisson-commutativity of the  $I_i$  is equivalent to the liftability of  $\iota_1$  to  $\iota_2$ .

Let  $\Theta_A := \text{Der}(A, A) = \bigoplus_{i=1}^n A \frac{\partial}{\partial I_i}$  and put  $C^p := R \otimes_A \wedge^p \Theta_A$ . We have  $n$  commuting derivations  $f \mapsto \{I_i, f\}$  of  $R$ , which combine to define a differential

$$\delta : C^p \longrightarrow C^{p+1}, \quad fw \mapsto \sum_{i=1}^n \{f, I_i\} \frac{\partial}{\partial I_i} \wedge w$$

**Proposition** [5]: Consider  $\iota_k : A \longrightarrow Q_k$  and a lifting to  $\iota_{k+1} : A \longrightarrow Q_{k+1}$ . Then there exists a well-defined obstruction element

$$\Xi = \Xi(\iota_k) \in H^2(C^\bullet, \delta).$$

with the following property:  $\iota_k$  can be lifted to  $\iota_{k+2} : A \longrightarrow Q_{k+2}$  by changing the lift  $\iota_{k+1}$  if and only if  $\Xi(\iota_k) = 0$ .

We put  $X = \text{Spec}(R) = \mathbb{C}^{2n}$ ,  $S = \text{Spec}(A) = \mathbb{C}^n$  and let  $I : X \longrightarrow S$  the corresponding map. There is a discriminant set  $\Sigma \subset S$ , such that the pull-back  $I' : X' \longrightarrow S' := S \setminus \Sigma$  is smooth and for  $s \in S'$  the fibre  $X_s$  is a smooth Lagrangian subvariety of  $X$ . The complex  $(C^\bullet, \delta)$  can be sheafied to a sheaf complex  $\mathcal{C}^\bullet$  on  $X$ .

**Proposition** [5]: There is a natural map of complexes

$$\rho : (\Omega_{X/S}^\bullet, d) \longrightarrow (\mathcal{C}^\bullet, \delta)$$

which is an isomorphism on  $X'$ .

As a consequence, the obstruction class  $\Xi$  induces for  $s \in S'$  an element

$$\Xi_s \in H^2(\Omega_{X_s}) = H^2(X_s, \mathbb{C})$$

If one makes reasonable assumptions on the structure of the singularities, one can show coherence of the cohomology, using the classical Kiehl-Verdier approach:

**Theorem [4]:** If  $I : X \rightarrow S$  is *pyramidal*, then  $H^i(\mathcal{C}^\bullet, \delta)$  are  $\mathcal{O}_S$ -coherent.

**Corollary:** If  $H^2(\mathcal{C}^\bullet, \delta)$  is torsion free, then the obstruction  $\Xi$  is zero if and only if  $\Xi_s = 0$  for generic  $s \in S'$ .

In fact, the modules  $H^i$  are in fact free modules in all examples we calculated.

The classical Darboux-Givental'-Weinstein theorem says that in the  $C^\infty$  context, a neighbourhood of a Lagrange submanifold  $L$  is symplectomorphic to a neighbourhood in the cotangent bundle  $T^*L$ . The same is true in our situation for  $L = X_s \subset X$ , because  $L$  is a Stein space. As a consequence of the rigidity of the Poisson structure, it seems one can construct a formal quantisation on a formal generic fibre. This *Quantum Darboux theorem* would imply the vanishing of  $\Xi_s$  for  $s$  generic. One would obtain the following corollary: If  $I : X \rightarrow S$  is pyramidal and  $H^2(\mathcal{C}^\bullet, \delta)$  is torsion free, then there  $I$  lifts to a formal quantum integrable system: we find  $J_i \in \hat{Q}$ ,  $[J_1, J_j] = 0$  and  $J_i = I_i \pmod{\hbar}$ .

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