

Correlated electronic structures and spin-dynamics, some recent investigations

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Collaborators

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Overview

1.An electronic structure data-base2.DMFT implementation with examples3.Calculation of exchange parameters4.Spin-wave excitations5.Cohesion and spectra of rare-earths



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Electronic structure database

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	11 Na	12 Mg							13 A1	14 Si	15 P	16 S	17 C1	18 Ar				
	19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
	37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	.53 1	54 Xe
	55 Cs	56 Ba	*	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 A1	86 Rn
	87 Fr	88 Ra	**	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Uun	111 Uuu	112 Uub	113 Uut	114 Uuq	115 Uup	116 Uuh	117 Uus	118 Uuo
Computational Materials Science 44, 1042 (2009)			*	57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu
			**	89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr



Possible high temperature superconductors





Data-mining criterion-structure



Computational Materials Science 67 282 (2013)



Data-mining criterion-bands/bonds





Dynamical mean field theory

 $H = H_{LDA} + \sum \sum U_{\xi_1 \xi_2 \xi_3 \xi_4} c^{\dagger}_{R\xi_1} c^{\dagger}_{R\xi_2} c_{R\xi_3} c_{R\xi_4}$ $R \ \xi_1 \xi_2 \xi_3 \xi_4$

The Hubbard model is mapped into an Anderson Impurity Model

FICTITIOUS SYSTEM REPRODUCING THE DYNAMICS

The mapping is made with the condition of preserving the local Green's function and is exact in the limit of infinite nearest neighbors



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Springer Series in Solid-State Sciences 167

John XI Wils Mebarek Alouari Per Andersson Anna Delin Olek Gresson Delsky Grechmon Full-Potential Electronic Structure Method Energy and Force Calculations with Density Unicincian and Opmania Mean Feld Theory

This book covers the theory of electronic structure of materials, with special emphasis on the usage of linear muffin-tin orbitals. Methodological aspects are given in detail as are examples of the method when applied to various matrials. Different exchange and correlation functionals are described and how they are implemented within the basis of linear muffin-tin orbitals. Functionals covered are the local spin density approximation, generalised gradient approximation, self-interaction correction and dynamical mean field theory.

> Full-Potential Electronic Structure Method

Wills · Delin ·

Andersson Grechnyev



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Full-Potential Electronic Structure Method

Energy and Force Calculations with Density Functional and Dynamical Mean Field Theory

123

DMFT implementation:

1. Grechnev, Di Marco et al. PRB, 76, 35107 (2007) 2. Di Marco et al PRB 79 115111 (2009) 3. Thunström et al. PRB 79, 165104 (2009) 4. Grånäs et al. Comp. Mat. Sci. 55, 295 (2012) 5. Thunström, Di Marco et al. PRL 109, 186401 (2012).

Katsnelson+Lichtenstein



Correlated basis

$$\hat{\boldsymbol{A}}_{R} \equiv \sum_{\boldsymbol{\xi}, \boldsymbol{\xi}'} |\boldsymbol{R}, \boldsymbol{\xi}\rangle \langle \boldsymbol{R}, \boldsymbol{\xi}| \sum_{\mathbf{k}} \hat{\boldsymbol{A}}_{\mathbf{k}} |\boldsymbol{R}, \boldsymbol{\xi}'\rangle \langle \boldsymbol{R}, \boldsymbol{\xi}'|$$

Two choices of $|R,\xi\rangle,$ muffin-tin based and orthogonal: i) MT

$$\chi_{lm} = i^l Y_{lm} (a\phi_{l\nu} + b\phi'_{l\nu})$$
$$\chi_{lm} = i^l Y_{lm} n_l$$

ii) ORT

$$(\mathbf{H} - \epsilon \mathbf{O})\mathbf{x} = 0$$
$$\mathbf{O} = \mathbf{L}\mathbf{L}^{h}$$

Cholesky decomposition gives $(\mathbf{L}^{-1}\mathbf{H}\mathbf{L}^{-h} - \epsilon \mathbf{1})\mathbf{y} = 0$ $\mathbf{y} = \mathbf{L}^{h}\mathbf{x}$

Exact Diagonalization Solver

The finite size problem can be solved exactly with a direct construction of all the accessible many-body states.

N=5 electrons in K=10 orbitals:

$ \Psi_1^5 angle$	=	$ 1111100000\rangle,$
$ \Psi_2^5 angle$	=	$ 1111010000\rangle,$

 $|\Psi_M^5\rangle = |0000011111\rangle.$

<₽

M corresponds to $\binom{K}{N}$

 \mathbf{P}

Too large for standard computational resources!

Block diagonalization

up to 30 bath states!



Local correlation effects in the electronic structure of Mn doped GaAs with LDA+DMFT

Igor Di Marco

Exact Diagonalization Solver

The finite size problem can be solved exactly with a direct construction of all the accessible many-body states.

N=5 electrons in K=10 orbitals:

Once the many-body states have been determined, the one-particle Green's function can be obtained through the Lehmann representation

$$G^{\rm ED}(i\omega)_{\xi_1\xi_2} = \frac{1}{Z} \sum_{\nu\mu} \frac{\langle \mu | \hat{c}_{\xi_1} | \nu \rangle \langle \nu | \hat{c}_{\xi_1}^{\dagger} | \mu \rangle}{i\omega + E_{\mu} - E_{\nu}} \left(e^{-\beta E_{\mu}} + e^{-\beta E_{\nu}} \right)$$



Igor Di Marco



Valence band of Mn-doped GaAs





P. Thunström et al. PRL 109 186401 (2012)

A two-step approach to model excitations

1. DFT calculation and extraction of the exchange parameters

$$\hat{H} = -\sum_{i \neq j} J_{ij} \cdot (\vec{e}_i \cdot \vec{e}_j)$$

2. Solution of a parameterised Heisenberg Model.

Finite-temperature magnetism. Simulations for the magnon spectra, Curie Temperature

Ab initio exchange parameters from DFT

Lichtenstein et al JMMM 67 65 (1987)

$$J_{ij} = \frac{-1}{4\pi} \int_{-\infty}^{E_f} d\varepsilon \operatorname{Tr}_L \left[\Delta_i \cdot G_{ij}^{\uparrow}(\varepsilon) \cdot \Delta_j \cdot G_{ji}^{\downarrow}(\varepsilon) \right]$$

Local exchange splitting Inter-site Green's function

 $\Delta_{i} = \begin{pmatrix} \hat{H}_{i}^{\uparrow} - \hat{H}_{i}^{\downarrow} \end{pmatrix}$ DOS(E) E

$$G_{ij}^{\sigma}(z) = \left\langle i \left| \hat{G}(z) \right| j \right\rangle = \left\langle i \left| \frac{1}{z - \hat{H}^{\sigma}} \right| j \right\rangle$$



bcc Fe: *J_{ij}*'s

LDA+U (Static Mean Field)

Dynamical Mean Field Theory (DMFT)



Hubbard U: *U*=2.3 eV $J=0.9 \ eV$





hcp Gd Ferromagnetic metal (T_c=292 K)



hcp Gd: magnons

<u>Results</u>: Atomistic spin dynamics (UppASD software)

Solving the equation of motion:

a

$$\frac{d\vec{S}_j}{dt} = \frac{1}{\hbar} \begin{bmatrix} \partial \hat{H} \\ \partial \vec{S}_j \end{bmatrix} \times \vec{S}_j$$
contains J_{ij} 's

Simulated magnon spectrum at T=78 K



Theory: Y.K. & D. Rodriguez Exp: Koehler *et al.* PRL **24** 16 (1970)

Magnetic properties of NiO



	PRB 91 125133 (2015)				
Computational setup	J_1	J_2			
LSDA	0.04	-1.58			
LSDA + DMFT	-0.003	-0.48			
LSDA + U	-0.002	-0.50			
LSDA + $U(U = 8 \text{ eV})$ (Ref. [1])	0.004/0.0	-0.53			
Exp. 1 (Ref. [2])	-0.051	-0.637			
Exp. 2 (Ref. [3])	0.051	-0.67			

LDA+U already provides a good description of the J_{ij} 's

[1] Jacobsson PRB 88 134427 (2013)

[2] Hutchings PRB 6 3447 (1972)

[3] Shanker PRB 7 5000 (1973)

NiO: J_2 as a function of U



Does $\frac{t^2}{T}$ hold?

NiO: J_2 as a function of U















Mannens dika åldrar

MELLAN 15-20 AR AR MANNEN SOM EN FIAT, LITEN MEN ETTRIS,

MELLAN 20-30 AR AR MANNEN SOM EN PORSCHE, SNABB OCH VÄLSMORP.

MELLAN 30-50 ÅR ÄR MANNEN SOM EN OPEL RECORD LOVAR MER ÄN DEN HÅLLER.

MELLAN SO-60 AR AR MANNEN SOM EN SAMMAL FORD MAN MÅSTE HÄLLA I HONOM LITE SPRIT FÖR ATT FÅ HONOM ATT STARTA.

Impact of correlations on J's

Exchange parameters from Dynamical Mean Field Theory (DMFT):

$$J_{ij} = \frac{\mathrm{T}}{4} \sum_{n} \mathrm{Tr} \Big[\hat{\Delta}_{i}(i\omega_{n}) \hat{G}_{ij}^{\dagger}(i\omega_{n}) \hat{\Delta}_{j}(i\omega_{n}) \hat{G}_{ji}^{\downarrow}(i\omega_{n}) \Big]$$

Katsnelson & Lichtenstein PRB 61 8906 (2000)

G contains self-energy: $G_{ij}(z) = \left\langle i \left| \frac{1}{z - \hat{H} - \hat{\Sigma}(z)} \right| j \right\rangle$

Exchange field becomes dynamical! $\Delta(i\omega_n) = \left(\hat{H}_i^{\uparrow} - \hat{H}_i^{\downarrow}\right) + \left(\Sigma_i^{\uparrow}(i\omega_n) - \Sigma_i^{\downarrow}(i\omega_n)\right)$