

Tensor network algorithms and quantum frustrated magnetism

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Collaborators

Theorists

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K. Schmidt (Dortmund), J. Dorier (Lausanne)
S. Manmana, A. Honecker (Göttingen)

P. Corboz (Amsterdam)

Experimentalists

M. Takigawa, Y. Matsuda (ISSP, Tokyo)
C. Berthier, M. Horvatic (Grenoble)



Scope

- Introduction:
 - Numerical simulations in 2D
 - **Tensor network algorithms**
 - Models of **2D frustrated quantum magnetism**
- Shastry-Sutherland model:
 - Zero-field phase diagram
 - **Magnetization plateaux**
- Conclusions

Numerical simulations in 2D

- Exact diagonalizations → small clusters
- Quantum Monte Carlo → minus sign
- Degenerate perturbation theory
→ effective model still to be solved
- Series expansion → choice of initial state
- Variational Monte Carlo with Gutzwiller projected wave functions → not general
- DMRG → breaks lattice rotational symmetry

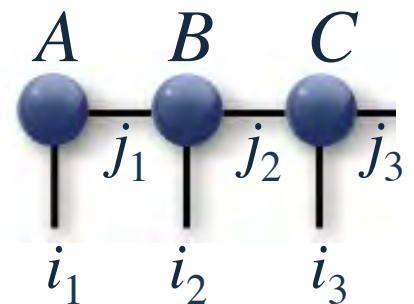
Tensor network ansatz

$$|\psi\rangle = \sum_{i_1 \dots i_N} c_{i_1 \dots i_N} |i_1\rangle \otimes \dots \otimes |i_N\rangle$$

$c_{i_1 \dots i_N} \simeq$ trace over a product of tensors

Example: Matrix product state in 1D

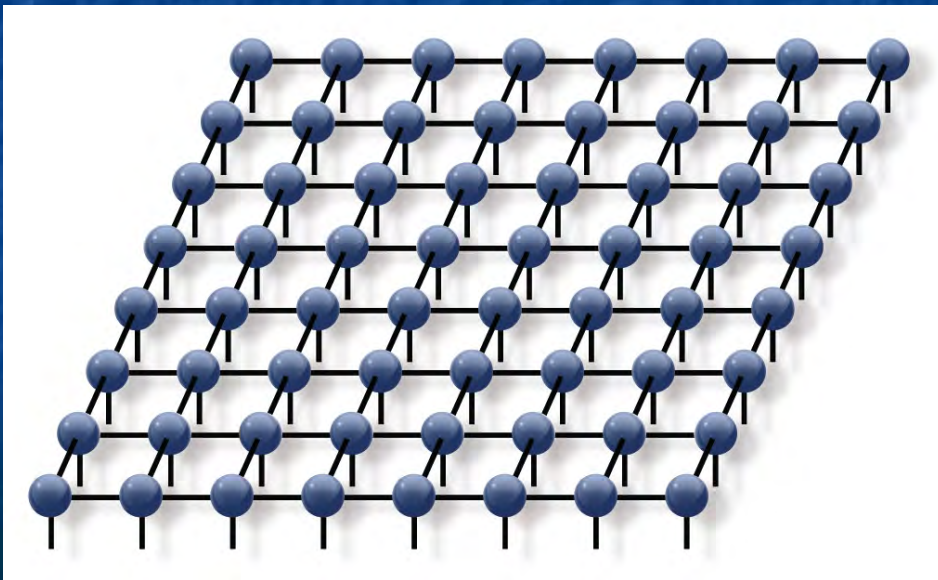
$$c_{i_1 i_2 i_3 \dots} \simeq \sum_{j_1 j_2 \dots} A_{i_1}^{j_1} B_{i_2}^{j_1 j_2} C_{i_3}^{j_2 j_3} \dots$$



Generalization to 2D

PEPS = product of entangled pair states

Verstraete and Cirac, 2004



$$A_i^{j_1 j_2 j_3 j_4} = \text{rank-5 tensor}$$

$$j_1, j_2, j_3, j_4 = 1, \dots, D$$

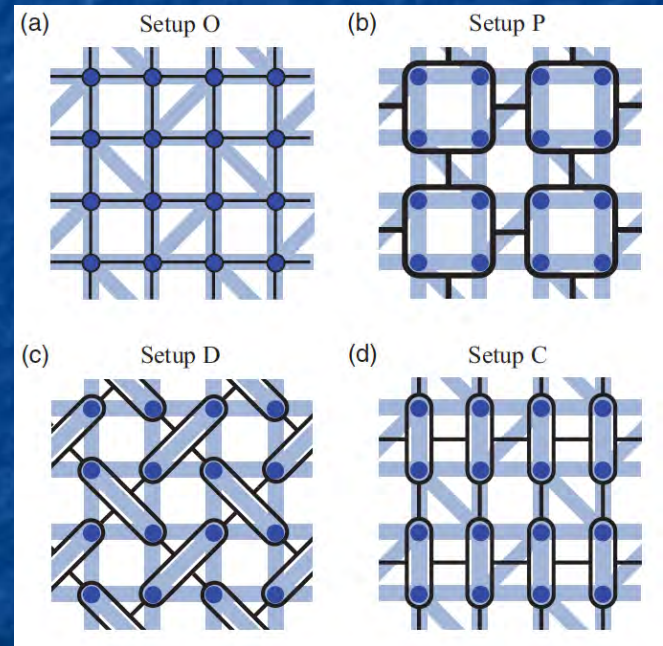
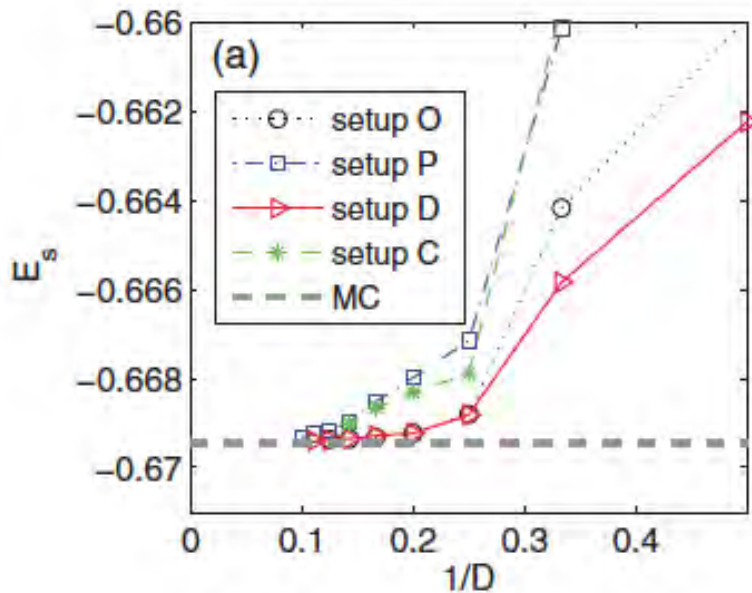
Variational approach

- **PEPS**: minimize the energy w.r.t. tensor elements
- Other schemes: renormalization (MERA,...)
- Advantage: **dim=pol(D,N)**, not exp(N)
- Why can it work?
 - reproduces the '**area law**' for the entanglement entropy in the GS of a local Hamiltonian

$$S = -\text{tr} (\rho_A \log \rho_A) \sim \partial A$$

- How large should D be? It depends...

Spin-1/2 Heisenberg model on square lattice



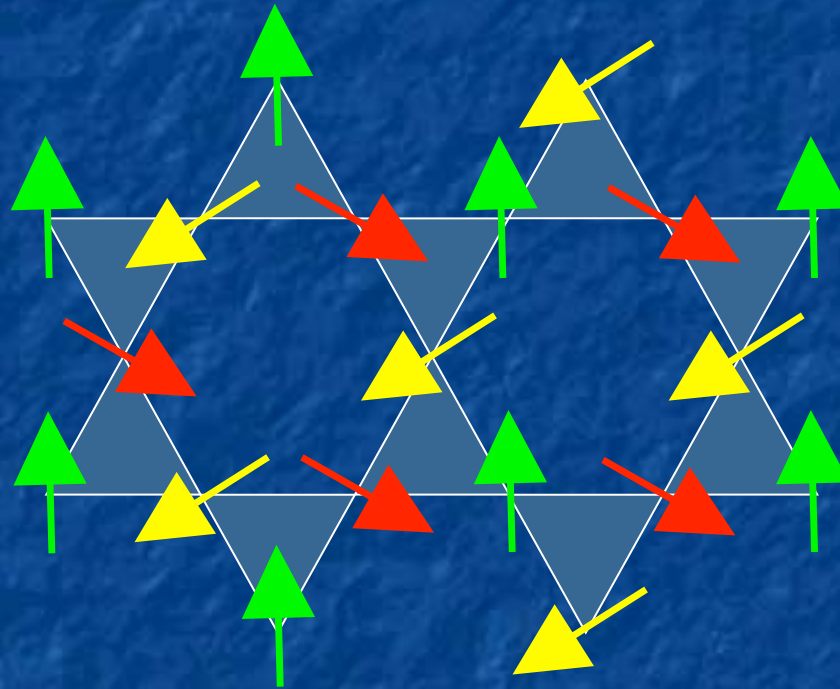
Excellent energy with $D=10$

Frustrated quantum magnetism

$$\mathcal{H} = \sum_{(i,j)} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

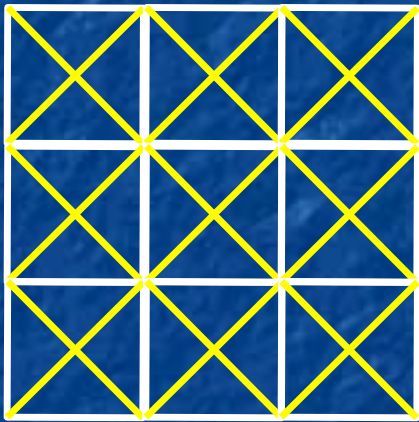
Infinite degeneracy of classical GS

Kagome lattice

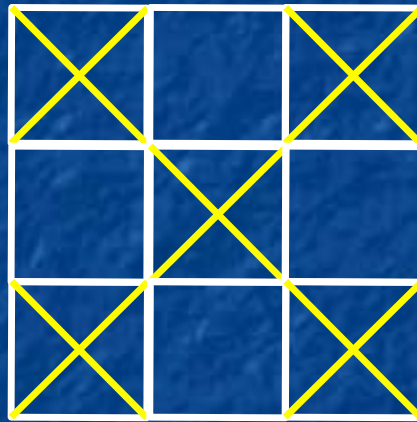


Spin liquid? Gapped? Algebraic? Valence-bond crystal?

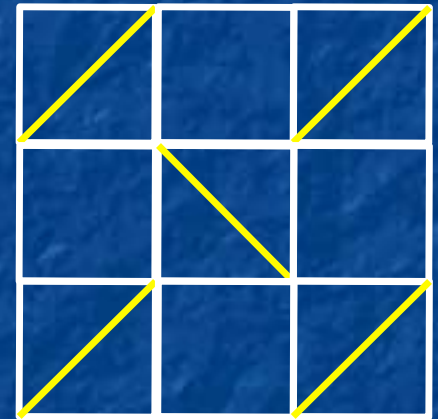
Square lattice



J_1 - J_2



Checkerboard



Shastry-Sutherland



Results on frustrated magnets

- **Kagome:** single site does not work. Needs 3 sites.
- **J_1 - J_2 on square lattice:** intermediate phase around $J_2/J_1=1/2$, but no consensus on its nature

→ Plaquette?

Yu and Kao, PRB 2012

→ Topological spin liquid?

Mezzacapo, PRB 2012; Wang, Gu, Verstraete, Wen 2012

- **Checkerboard:** intermediate plaquette phase, in agreement with previous results

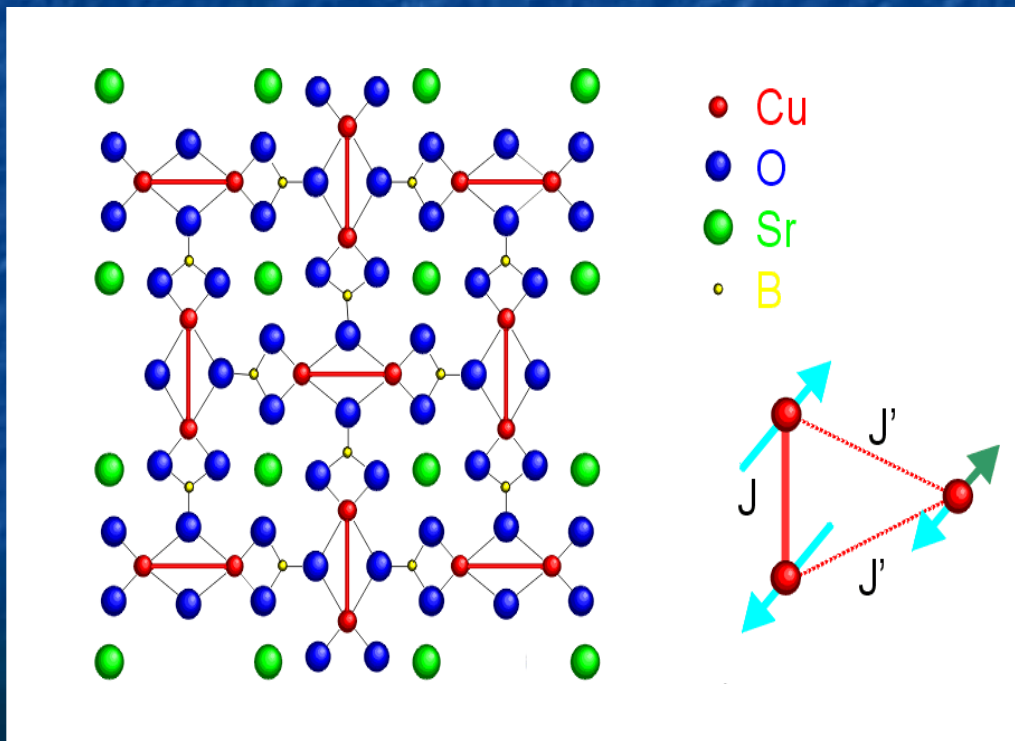
Chan, Han, Duan, PRB 2011

Shastry-Sutherland

- Isacsson and Syljuasen, PRE 2006
 - $D=2$, no intermediate phase in zero field
- Lou, Suzuki, Harada, Kawashima, arxiv 2012
 - MERA
 - intermediate plaquette phase
 - plateaux and supersolid phases in a field
- P. Corboz and FM, PRB 2013, PRL 2013 and 2014
 - plateau sequence of $\text{SrCu}_2(\text{BO}_3)_2$

SrCu₂(BO₃)₂

Smith and Keszler, JSSC 1991

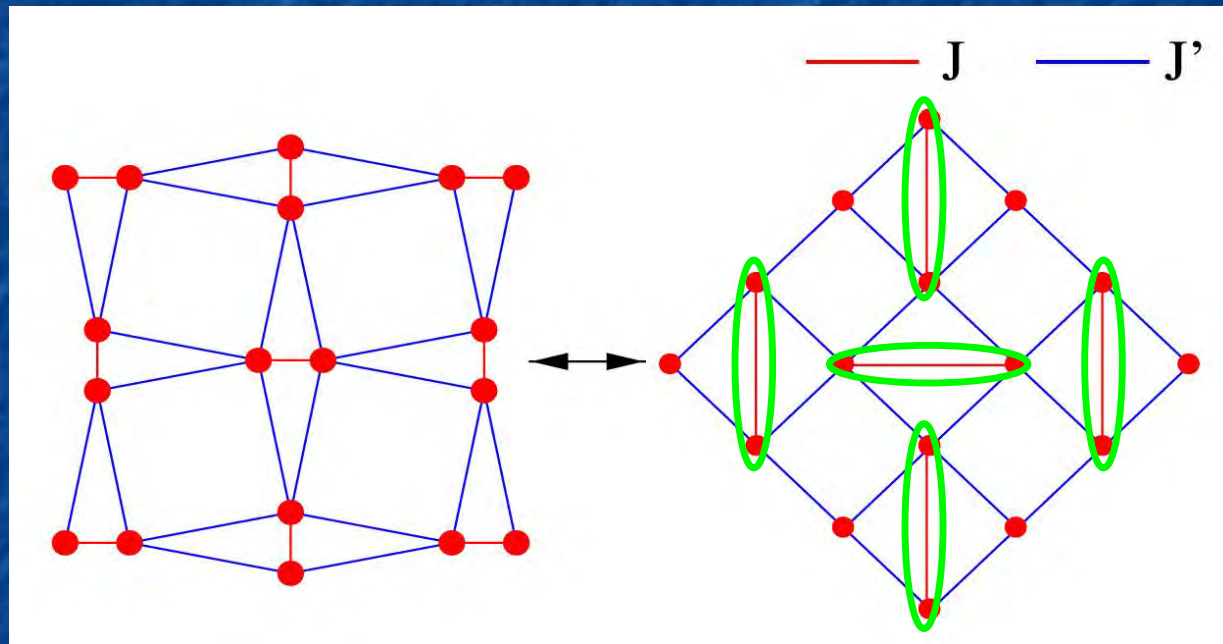


Cu²⁺ -> Spin 1/2

$J \approx 85 \text{ K}$

$J'/J \approx 0.65$

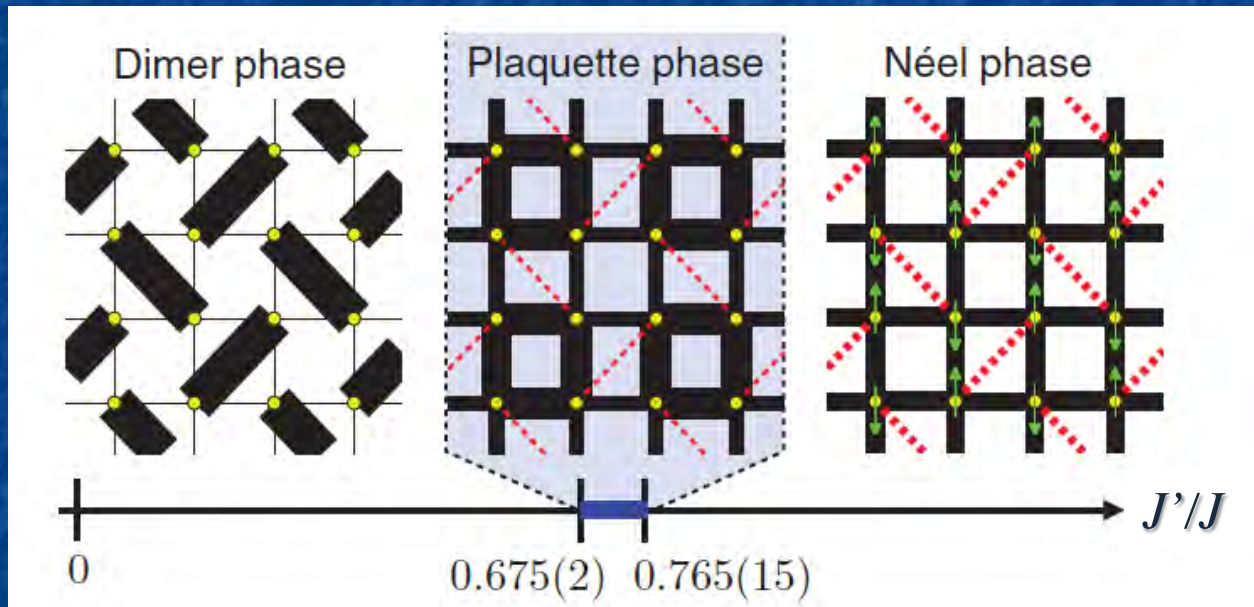
Shastry-Sutherland model



Ground-state = Product of singlets
on J -bonds if J'/J not too large

Shastry and Sutherland, '81

iPEPS zero-field phase diagram

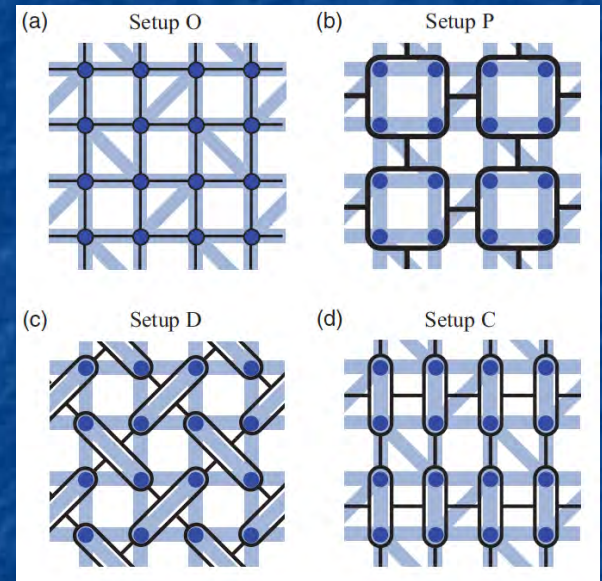
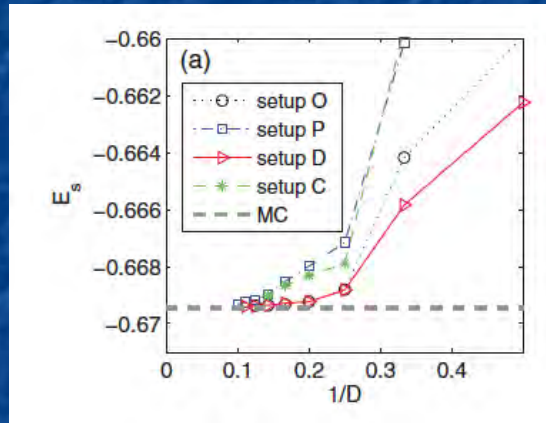


P. Corboz and FM, PRB 2013

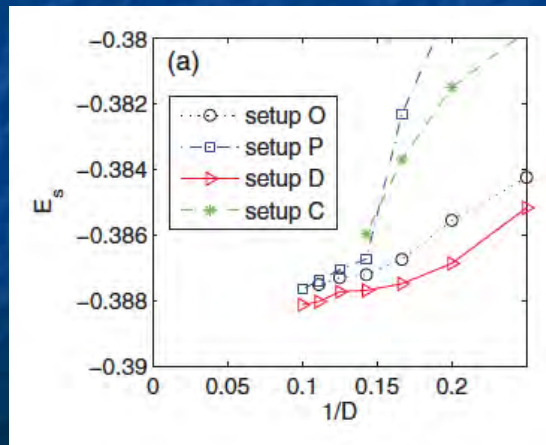
Agrees qualitatively with former studies
Koga-Kawakami 2000, Läuchli-Wessel-Sigrist 2002, etc.

Convergence with D

Neel ($J=0$)

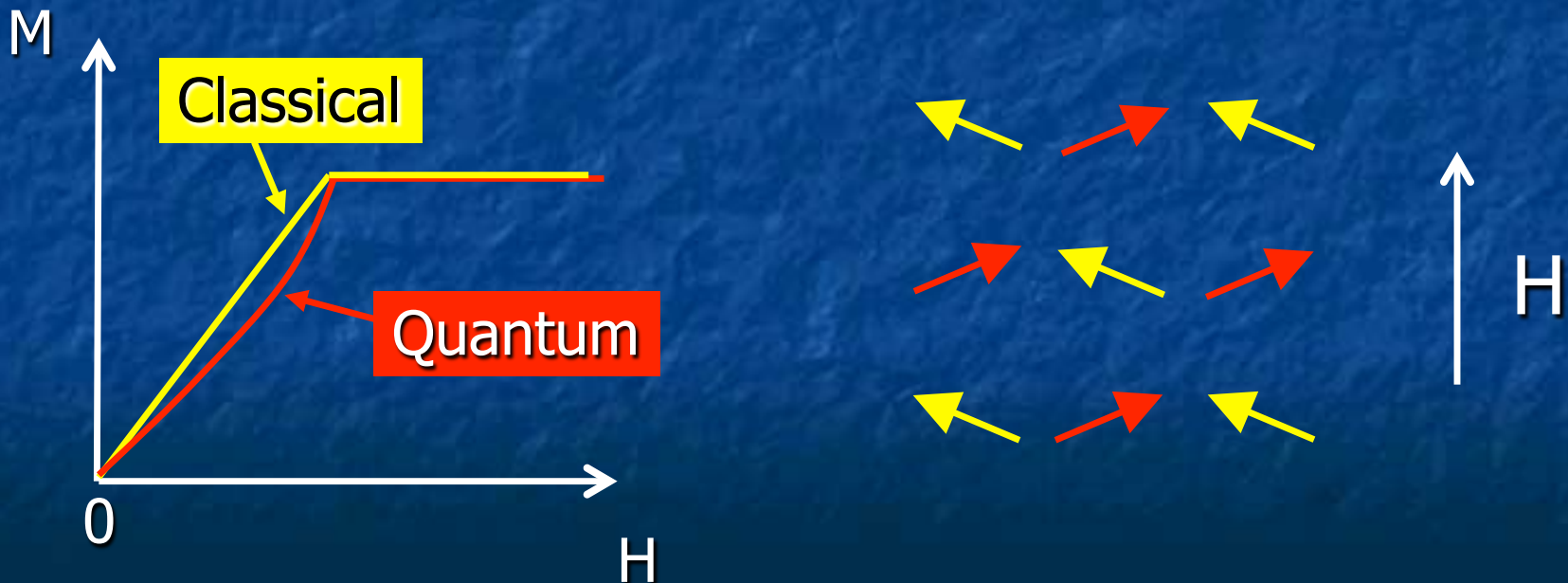


Plaquette phase



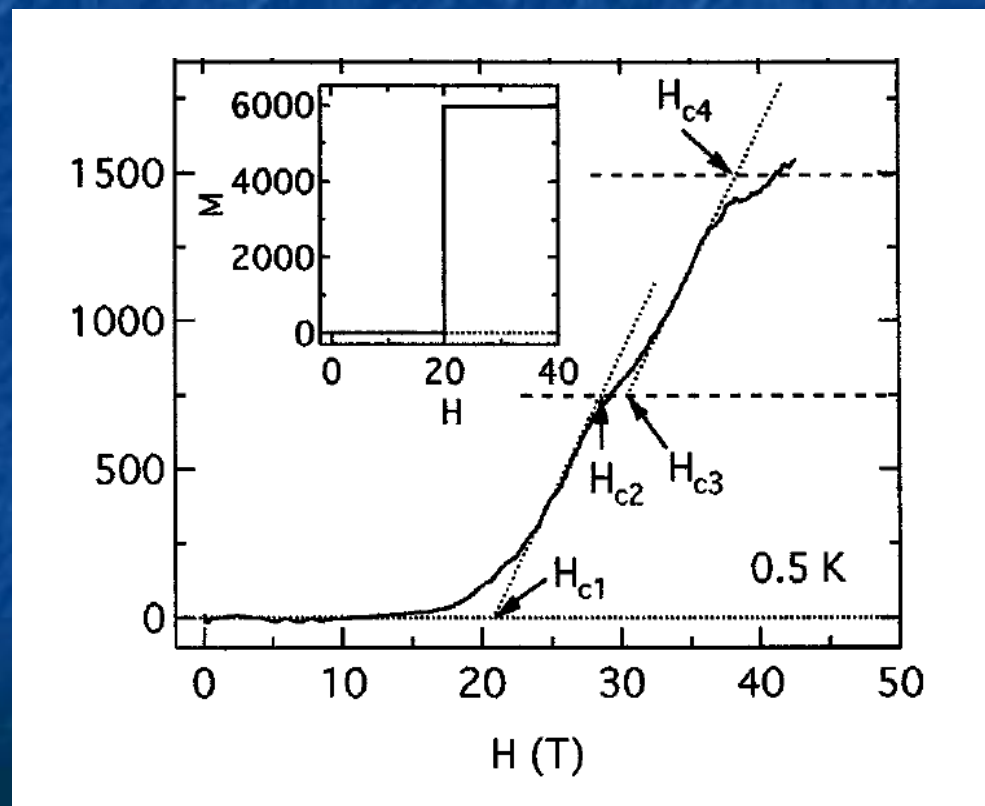
Magnetization of Néel AF

$$\mathcal{H} = J \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j - g\mu_B H \sum_i S_i^z$$



Exact Dimer Ground State and Quantized Magnetization Plateaus in the Two-Dimensional Spin System $\text{SrCu}_2(\text{BO}_3)_2$

H. Kageyama,^{1,2,*} K. Yoshimura,^{1,3,†} R. Stern,³ N. V. Mushnikov,² K. Onizuka,² M. Kato,¹ K. Kosuge,¹
C. P. Slichter,³ T. Goto,² and Y. Ueda²

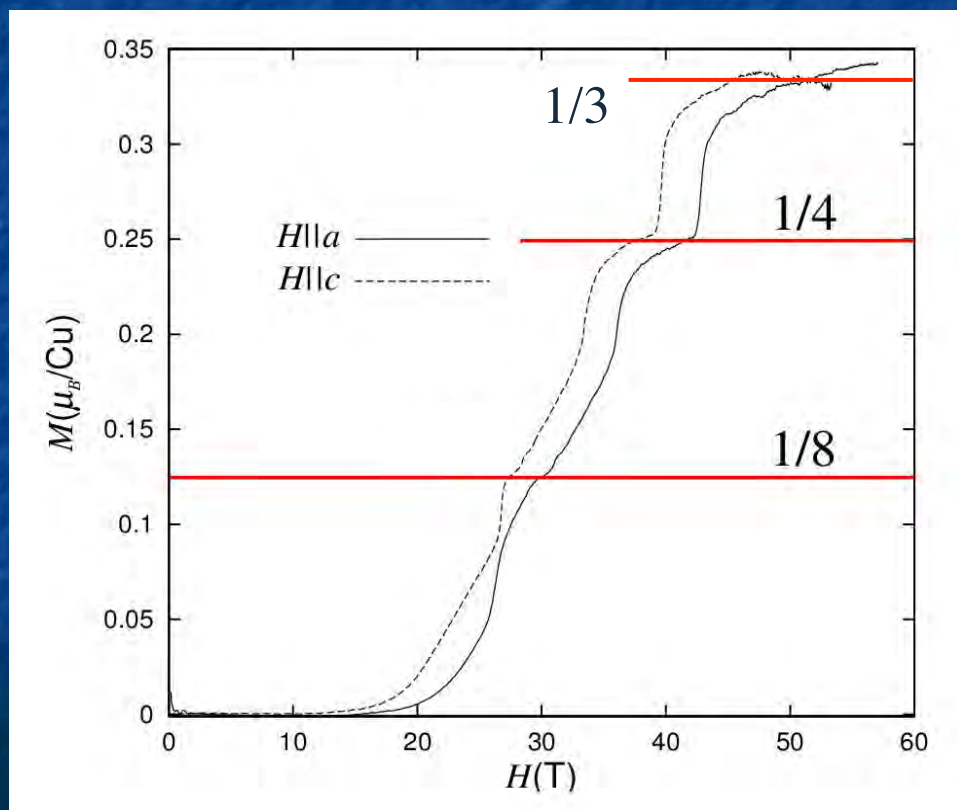


Anomalies

- $M=0$
- $M=1/8$
- $M=1/4$

1/3 Magnetization Plateau in $\text{SrCu}_2(\text{BO}_3)_2$ - Stripe Order of Excited Triplets -

Kenzo ONIZUKA, Hiroshi KAGEYAMA*, Yasuo NARUMI^{1,2},
Koichi KINDO^{2,1}, Yutaka UEDA and Tsuneaki GOTO



Plateaus

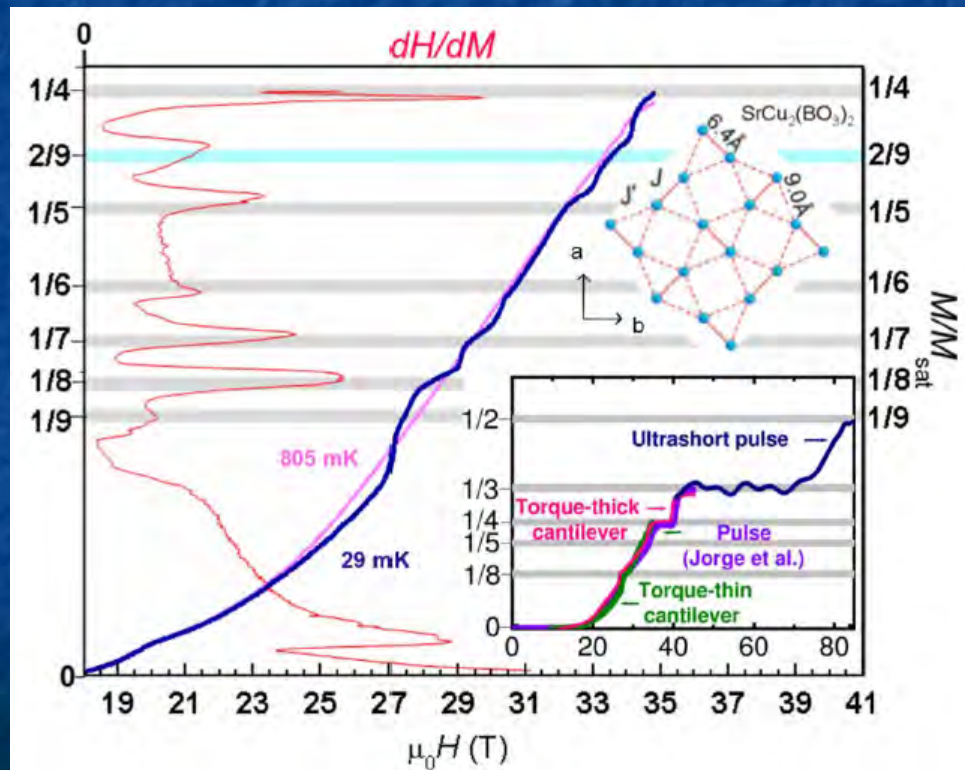
- $M=0$
- $M=1/8$
- $M=1/4$
- $M=1/3$

Fractalization drives crystalline states in a frustrated spin system

Suchitra E. Sebastian^{a,1}, N. Harrison^b, P. Sengupta^c, C. D. Batista^c, S. Francoual^b, E. Palm^d, T. Murphy^d, N. Marciano^a, H. A. Dabkowska^e, and B. D. Gaulin^e

www.pnas.org/cgi/doi/10.1073/pnas.0804320105

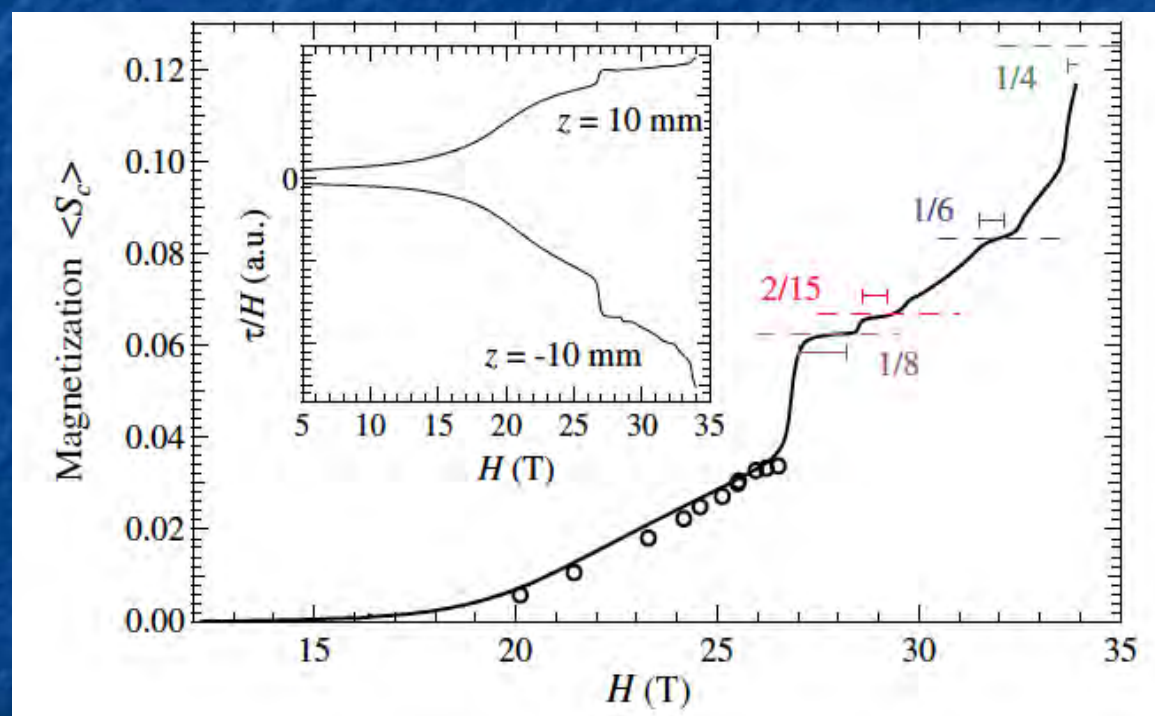
PNAS | December 23, 2008 | vol. 105 | no. 51 | 20157–20160



$1/9, 1/8, 1/7, 1/6,$
 $1/5, 2/9, 1/4$

Incomplete Devil's Staircase in the Magnetization Curve of $\text{SrCu}_2(\text{BO}_3)_2$

M. Takigawa,^{1,*} M. Horvatić,² T. Waki,³ S. Krämer,² C. Berthier,² F. Lévy-Bertrand,^{2,†} I. Sheikin,² H. Kageyama,⁴
Y. Ueda,¹ and F. Mila⁵



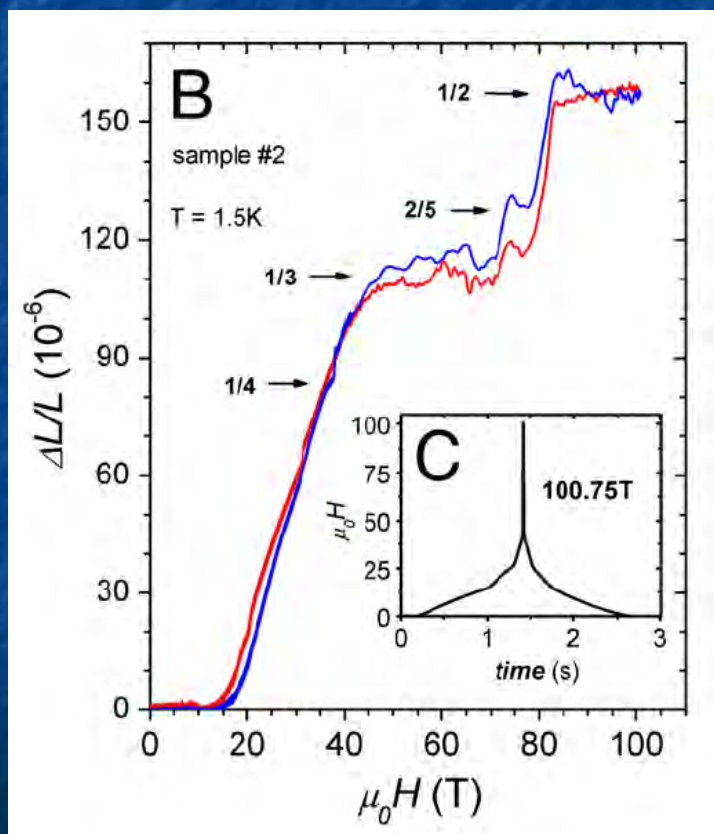
$1/8, 2/15, 1/6, 1/4, \dots$

Magnetostriction and magnetic texture to 100.75 Tesla in frustrated $\text{SrCu}_2(\text{BO}_3)_2$

Marcelo Jaime^{a,b,1}, Ramzy Daou^{c,d}, Scott A. Crooker^{a,b}, Franziska Weickert^b, Atsuko Uchida^{a,b}, Adrian E. Feiguin^e, Cristian D. Batista^f, Hanna A. Dabkowska^g, and Bruce D. Gaulin^{g,h}

12404–12407 | PNAS | July 31, 2012 | vol. 109 | no. 31

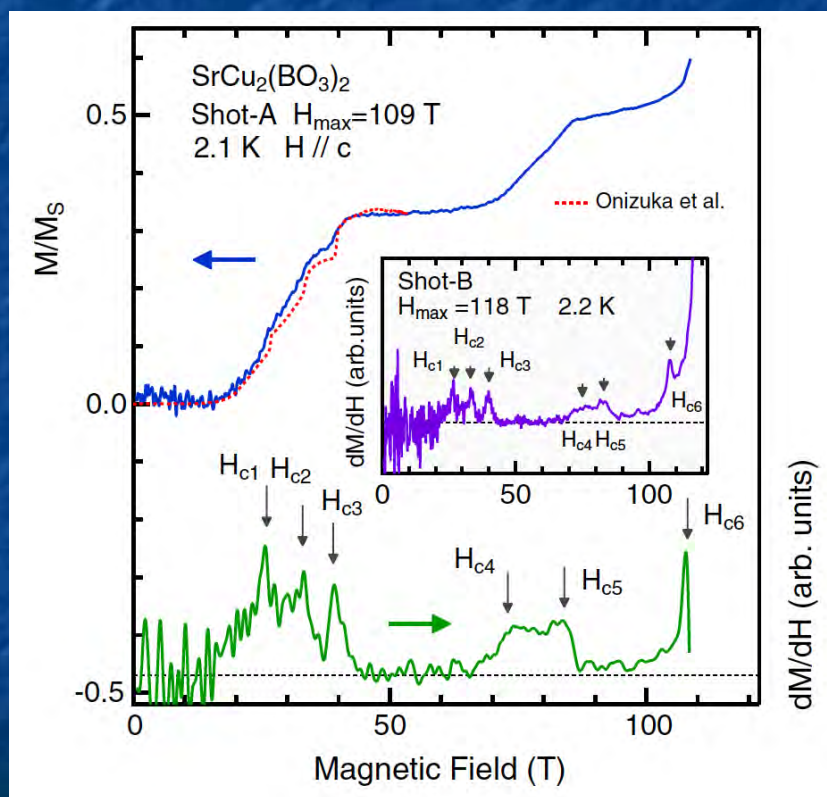
www.pnas.org/cgi/doi/10.1073/pnas.1200743109



..., 1/4, 1/3, 2/5, 1/2

Magnetization of $\text{SrCu}_2(\text{BO}_3)_2$ in Ultrahigh Magnetic Fields up to 118 T

Y.H. Matsuda,^{1,*} N. Abe,¹ S. Takeyama,¹ H. Kageyama,² P. Corboz,³ A. Honecker,^{4,5} S.R. Manmana,⁴
G.R. Foltin,⁶ K.P. Schmidt,⁶ and F. Mila⁷



...1/3, 1/2

Critical summary I

0, 1/9, 1/8, 2/15, 1/7, 1/6, 1/5, 2/9, 1/4, 1/3, 2/5, 1/2

Too many suspects!

Critical summary II

Only pulsed field

Only magnetostriction

0, ~~1/9~~, 1/8, 2/15, ~~1/7~~, 1/6, ~~1/5~~, ~~2/9~~, 1/4, 1/3, ~~2/5~~, 1/2

0, 1/8, 2/15, 1/6, 1/4, 1/3, 1/2

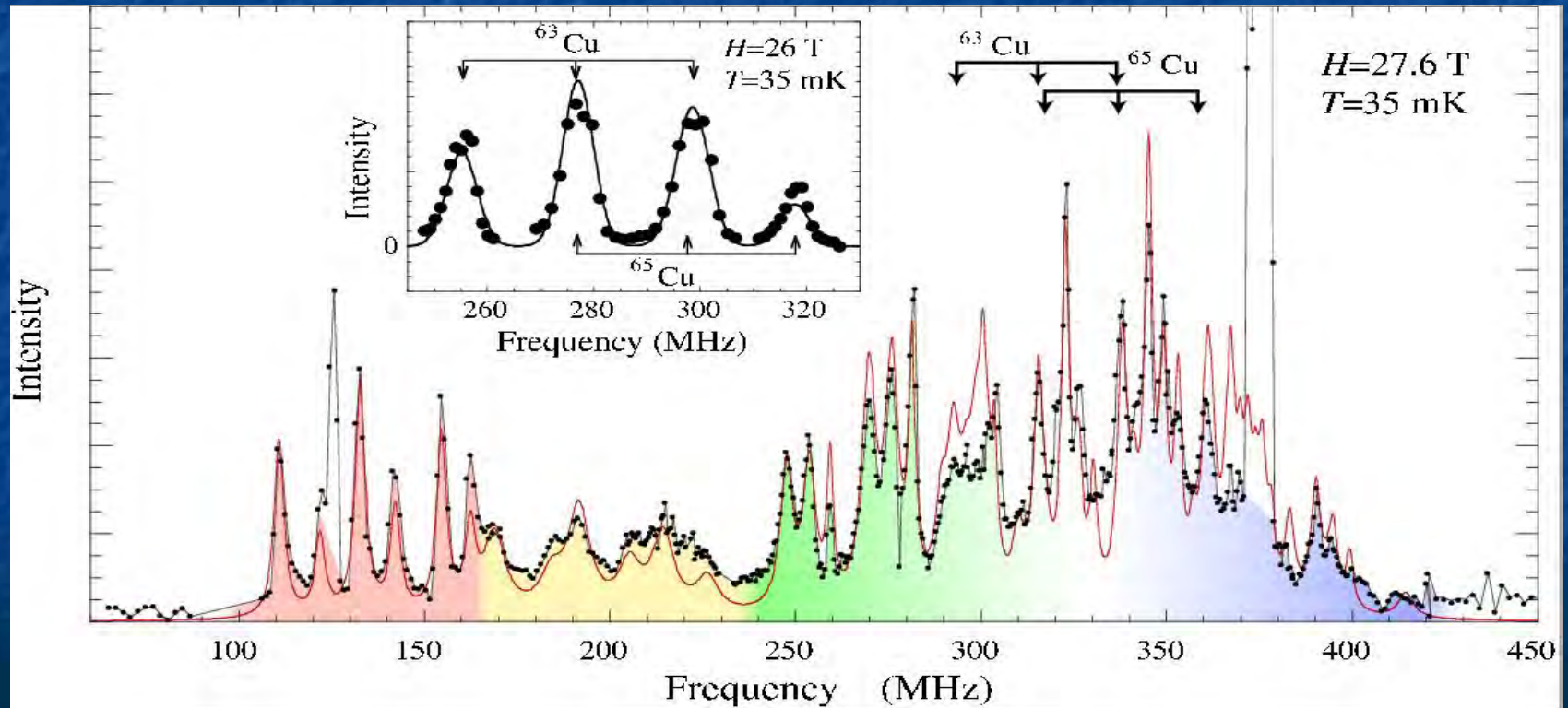
Quantum Hall Effect or broken spatial symmetry?

Magnetic Superstructure in the Two-Dimensional Quantum Antiferromagnet $\text{SrCu}_2(\text{BO}_3)_2$

K. Kodama,¹ M. Takigawa,^{1*} M. Horvatić,² C. Berthier,^{2,3}
H. Kageyama,¹ Y. Ueda,¹ S. Miyahara,^{1,4} F. Becca,⁴ F. Mila⁴

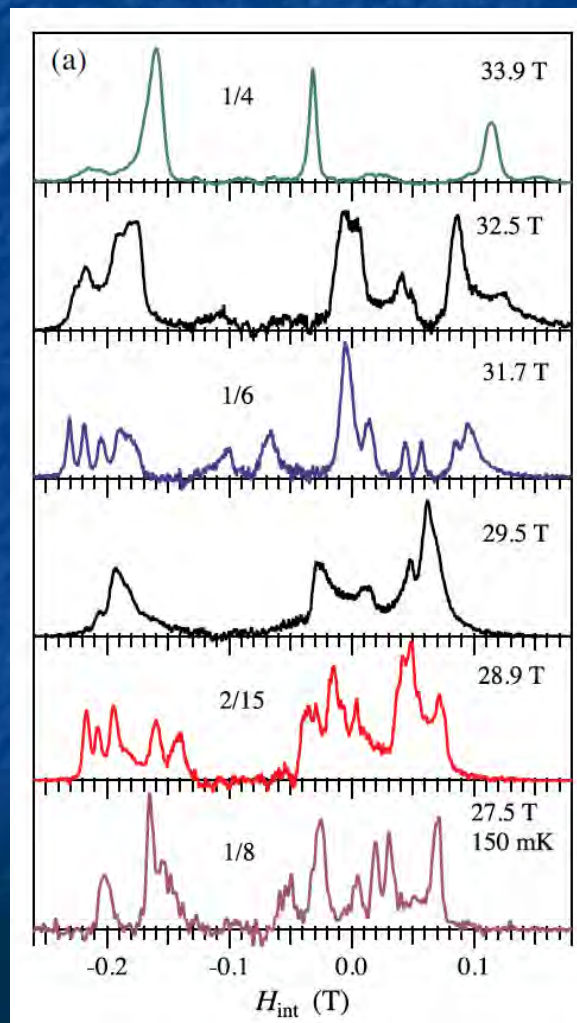
Broken symmetry
in 1/8 plateau

www.sciencemag.org SCIENCE VOL 298 11 OCTOBER 2002



Incomplete Devil's Staircase in the Magnetization Curve of $\text{SrCu}_2(\text{BO}_3)_2$

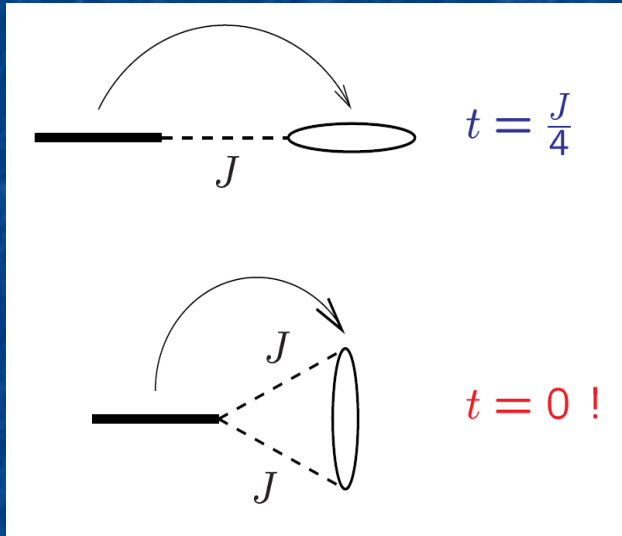
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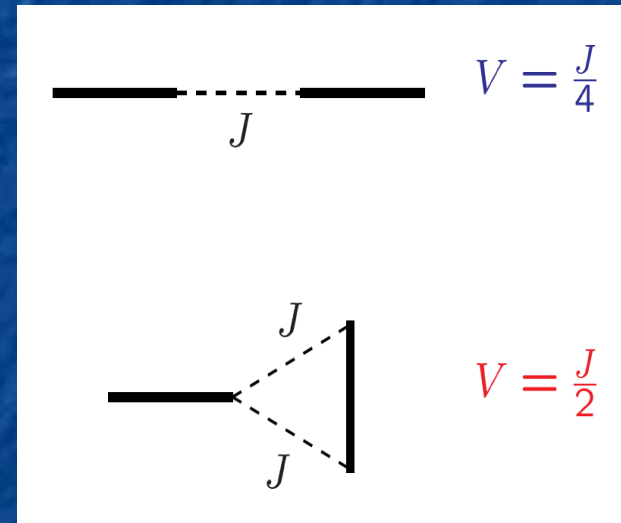
Broken symmetry in
1/8, 2/15, 1/6 and 1/4 plateaus

Effect of frustration

Triplet Hopping



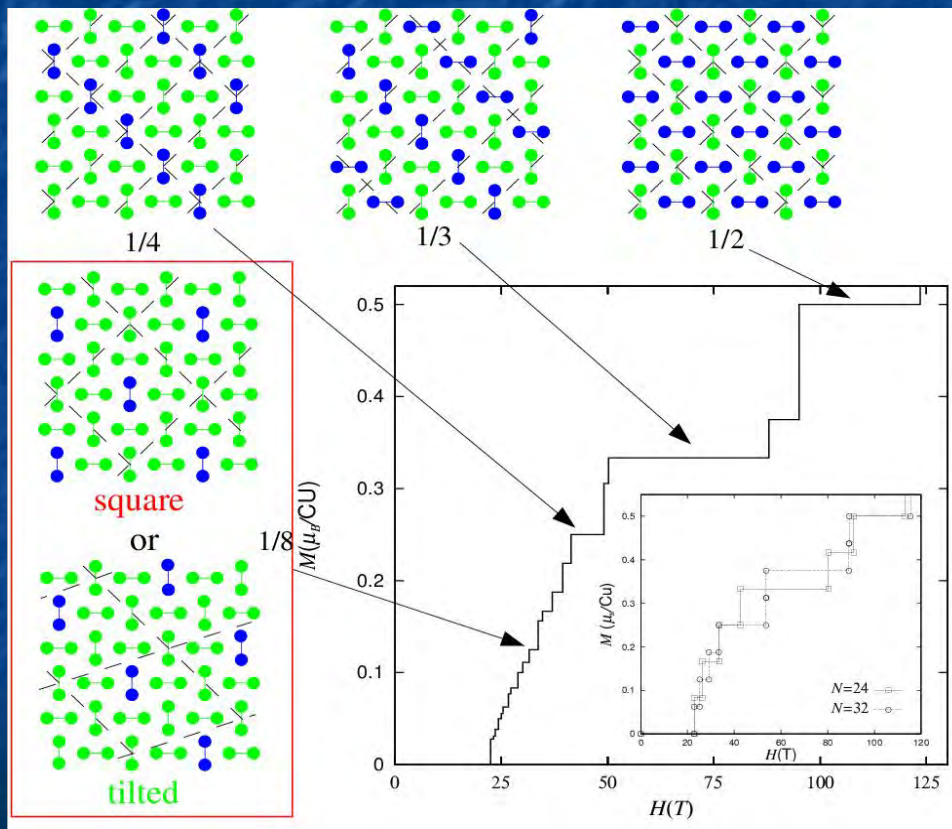
Triplet Repulsion



- Kinetic energy \ll potential energy
- Long-range repulsion
 - Crystals of triplets with high commensurability
 - **Magnetization plateaux**

Superstructures at magnetization plateaus in $\text{SrCu}_2(\text{BO}_3)_2$

Shin Miyahara and Kazuo Ueda



- Simple ansatz for long-range triplet-triplet interaction
- Many plateaus

Magnetization plateaus of the Shastry-Sutherland model for $\text{SrCu}_2(\text{BO}_3)_2$: Spin-density wave, supersolid, and bound states

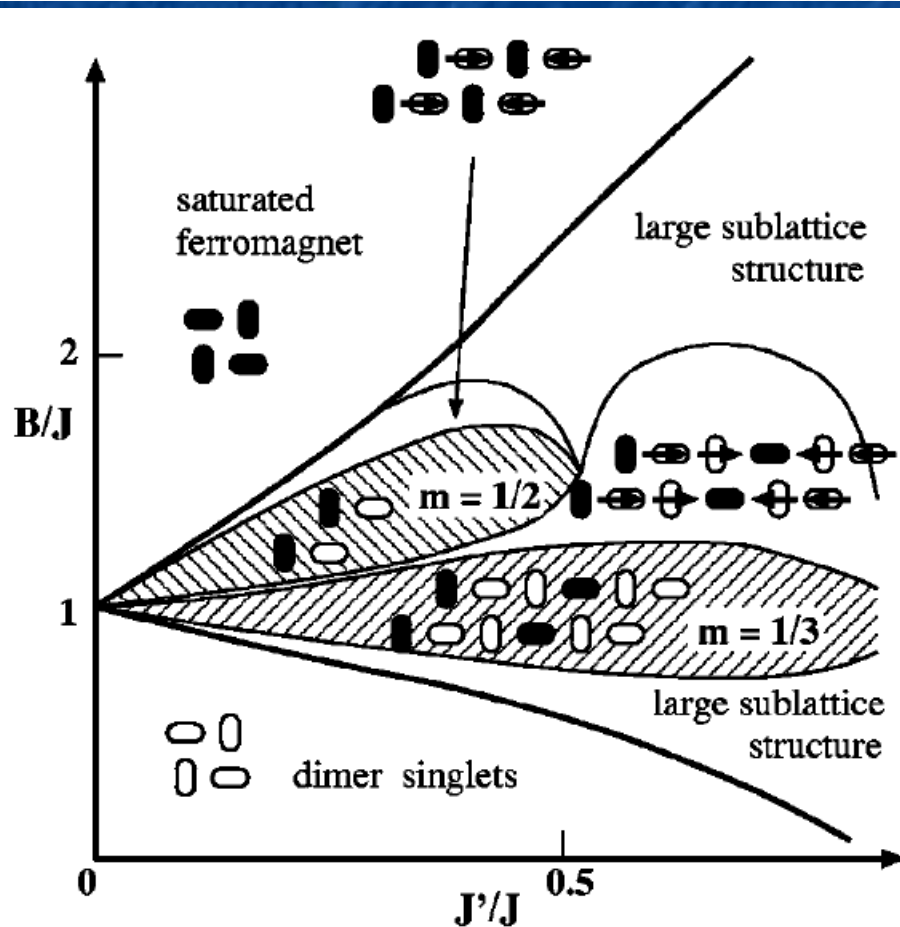
Tsutomu Momoi*

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138

Keisuke Totsuka

Department of Physics, Kyushu University, Hakozaki, Higashi-ku, Fukuoka-shi 812-8581, Japan

(Received 1 June 2000)

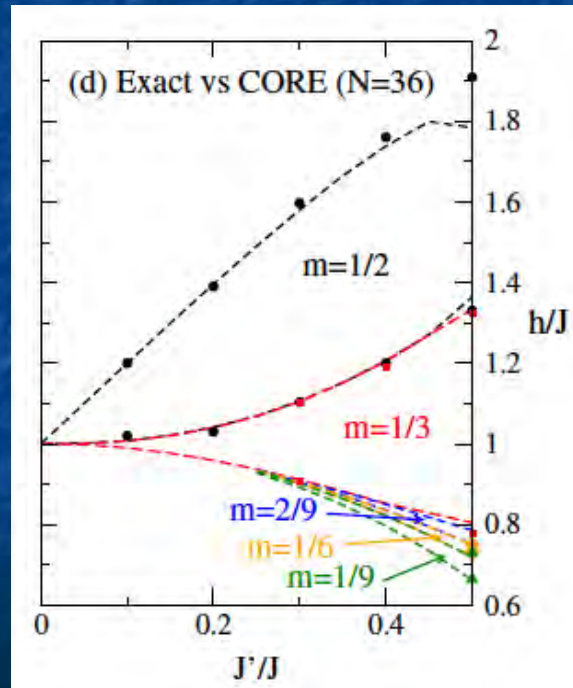
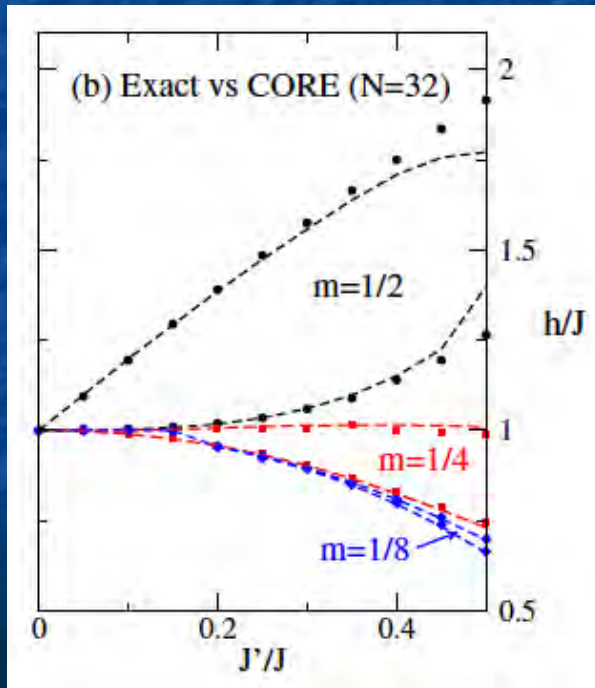


- Effective model to third order in $(J'/J)^3$
- Only short-range triplet-triplet interaction
- Only $1/3$ and $1/2$ plateaus
- Spin-supersolids

Effective Theory of Magnetization Plateaux in the Shastry-Sutherland Lattice

A. Abendschein^{1,2} and S. Capponi^{1,2,*}

Effective model with CORE (Contractor Renormalization)
Exact diagonalization on small clusters

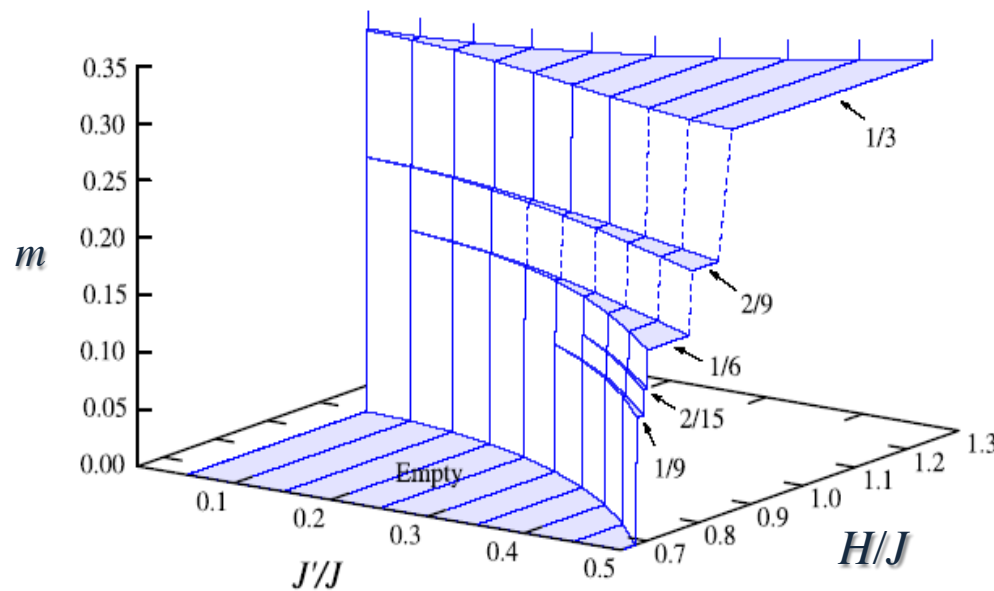


- Large finite-size effects
- No 2/15 plateau
- No information on plateau structures

Theory of Magnetization Plateaux in the Shastry-Sutherland Model

J. Dorier,¹ K. P. Schmidt,^{2,*} and F. Mila¹

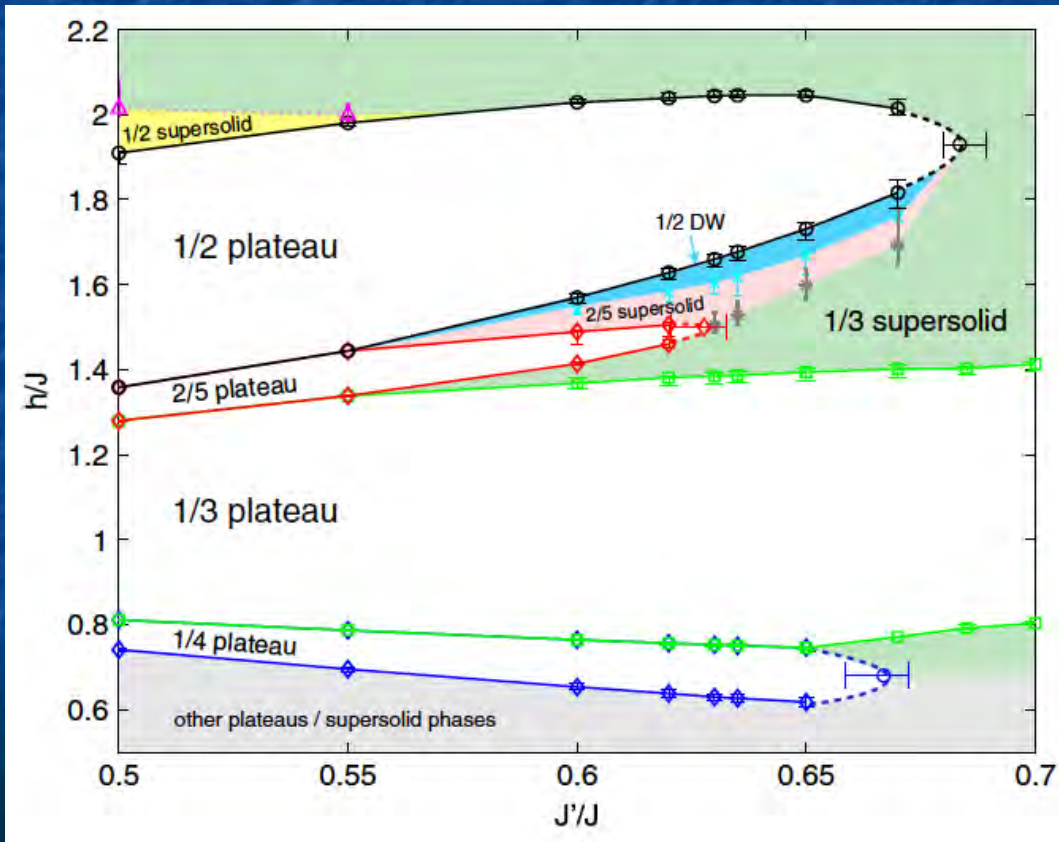
Long-range triplet-triplet interaction with high-order perturbation theory in J'/J



Triplet crystals:
not the right sequence
(at least up to $J'/J=0.5$)

Magnetization of $\text{SrCu}_2(\text{BO}_3)_2$ in Ultrahigh Magnetic Fields up to 118 T

Y.H. Matsuda,^{1,*} N. Abe,¹ S. Takeyama,¹ H. Kageyama,² P. Corboz,³ A. Honecker,^{4,5} S.R. Manmana,⁴
G.R. Foltin,⁶ K.P. Schmidt,⁶ and F. Mila⁷

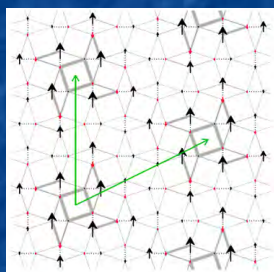
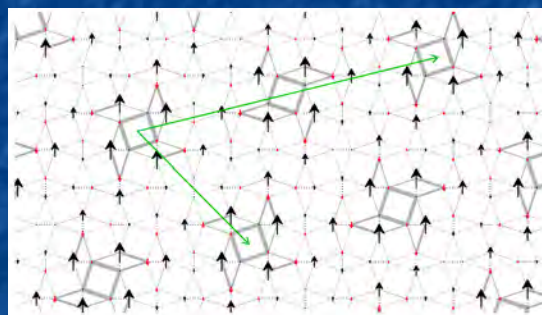
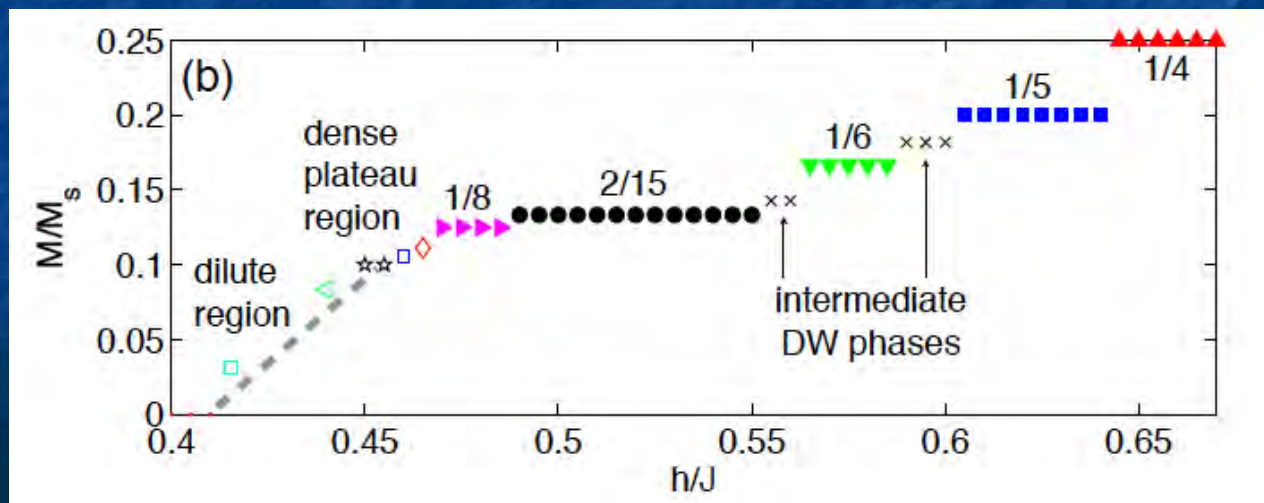
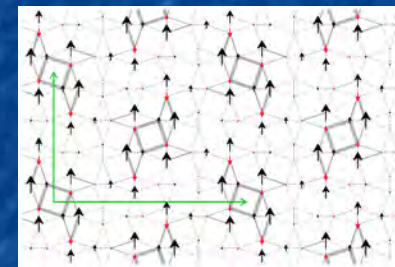


iPEPS

A variational approach
derived from quantum
information

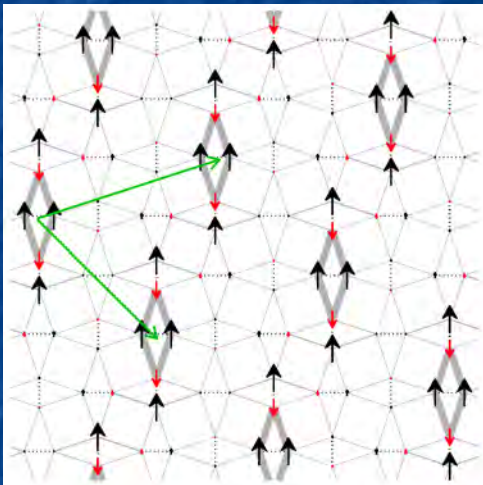
Verstraete & Cirac, 2004

Crystals of Bound States in the Magnetization Plateaus of the Shastry-Sutherland Model

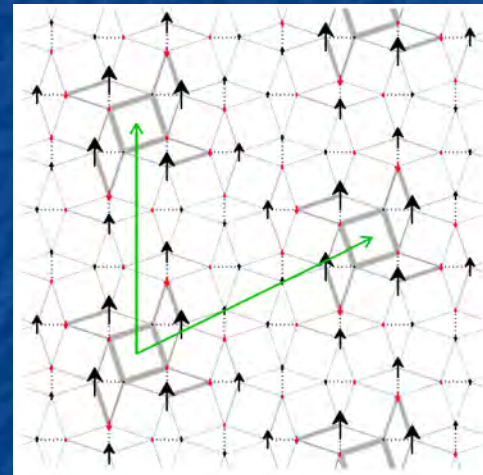
Philippe Corboz¹ and Frédéric Mila² $1/8$  $2/15$  $1/6$ 

Triplets versus bound-states

Triplet crystal



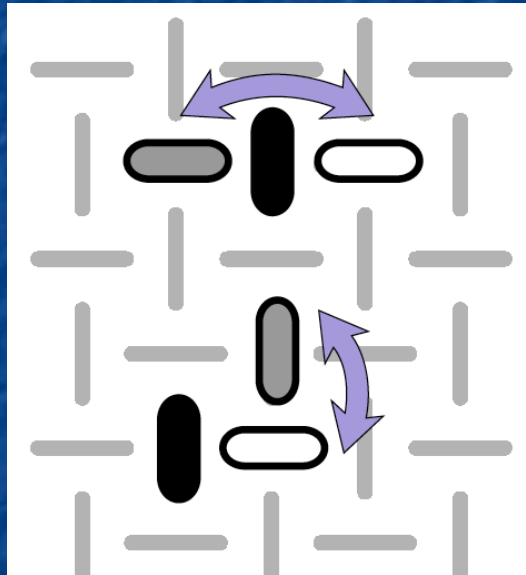
Bound state crystal



- Why are bound state crystals favoured?
- Why did it take 15 years to identify them?

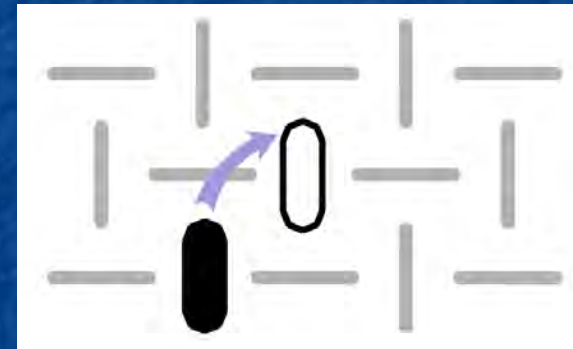
Why are bound states favoured?

Correlated hopping



$\propto (J'/J)^2 \rightarrow$ gain in kinetic energy

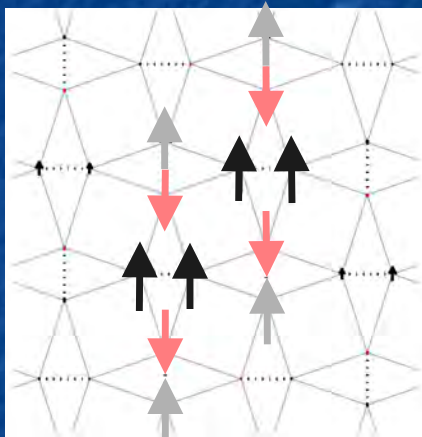
Single particle hopping



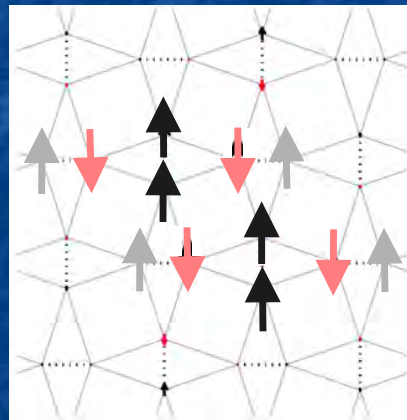
$\propto (J'/J)^6$

Boundstates as pinwheels

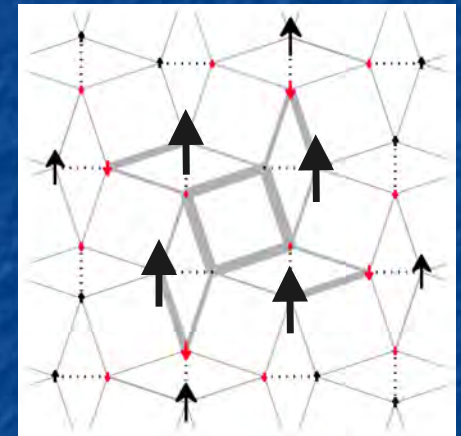
Order $(J'/J)^3$



+



=



Higher order

Magnetization plateaus of the Shastry-Sutherland model for $\text{SrCu}_2(\text{BO}_3)_2$: Spin-density wave, supersolid, and bound states

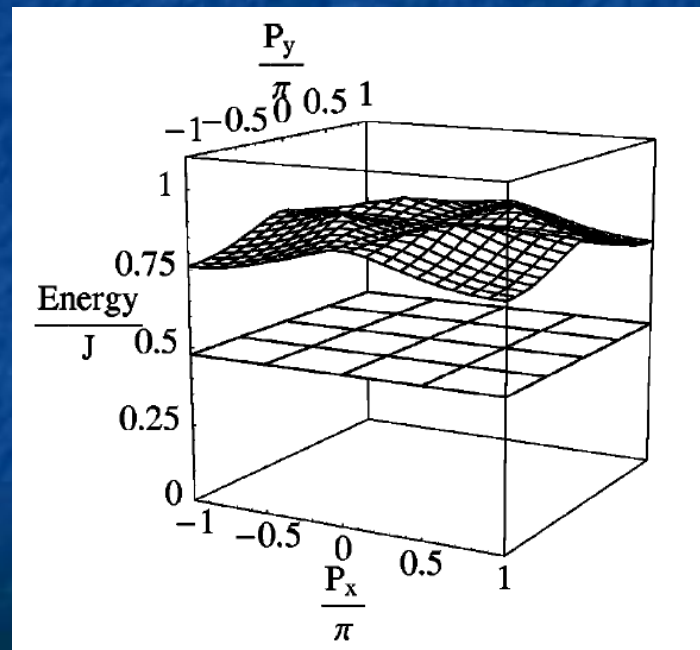
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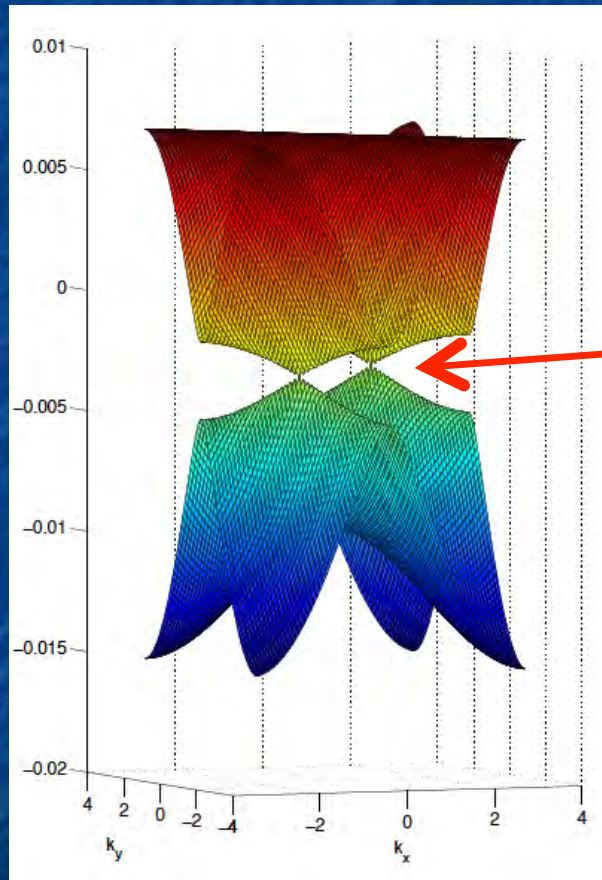
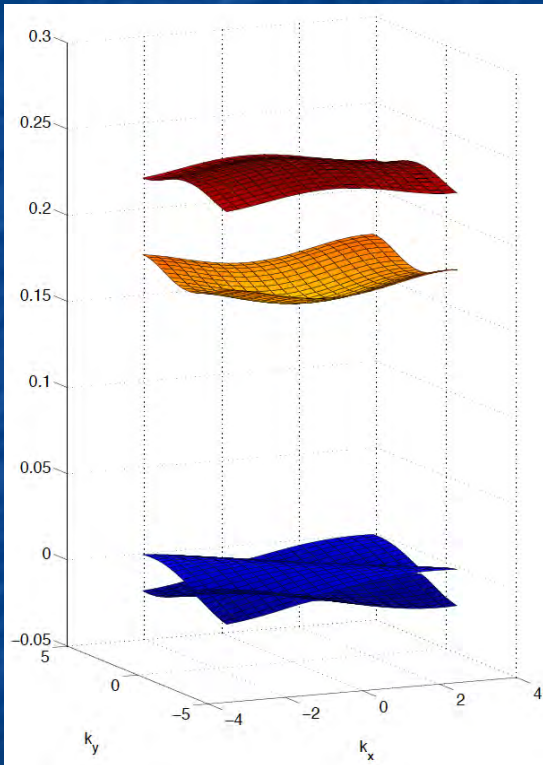
Department of Physics, Kyushu University, Hakozaki, Higashi-ku, Fukuoka-shi 812-8581, Japan

(Received 1 June 2000)



**Bound state dispersion:
minimum at the zone corner**

Localized bound state



Negative
reference
energy!

Conclusions on $\text{SrCu}_2(\text{BO}_3)_2$

- **Natural plateau sequence**
 - $1/8, 2/15, 1/6$: crystals of **bound states**
 - new predictions for their symmetry
- To be done next
 - revisit effective triplet model
 - compare predictions with NMR
 - X-ray or **neutron in high field** (>27 T)?

Conclusions

- Tensor network algorithms
 - **very promising** new method in 2D
 - encouraging results on various frustrated models
 - **solved the 15-year old puzzle of magnetization plateaux in $\text{SrCu}_2(\text{BO}_3)_2$**