



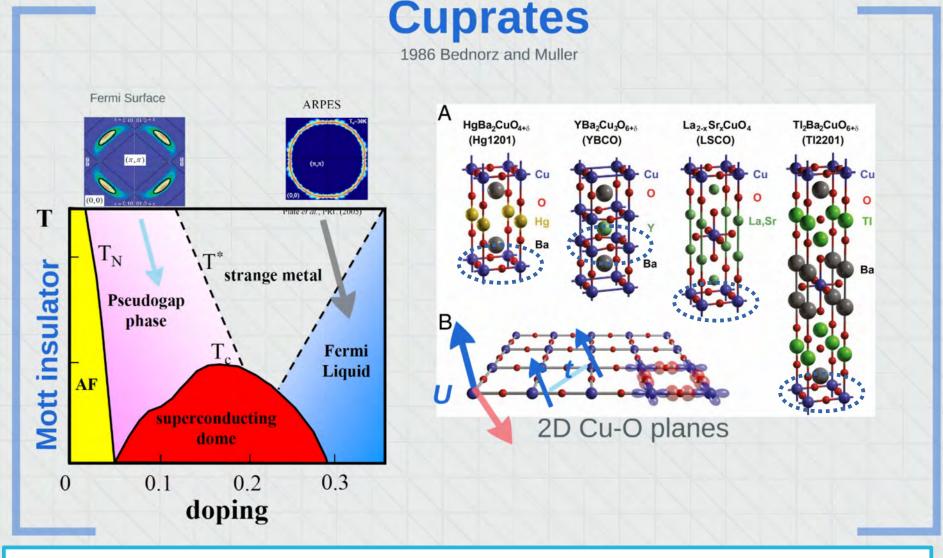


Comprendre le monde, construire l'avenir®

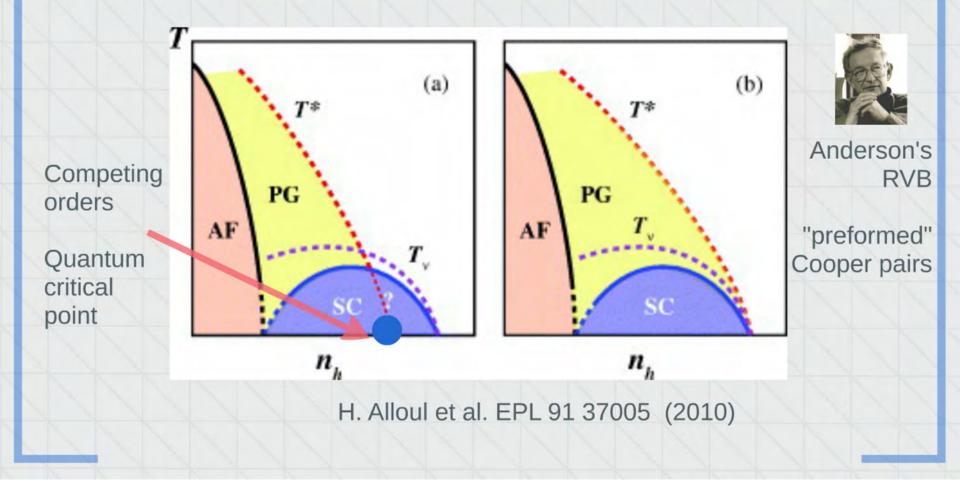
Pseudo-Gap to the Extreme

Marcello Civelli

PSEUDOGAP IS AN UNCONVENTIONAL (BAD) METAL

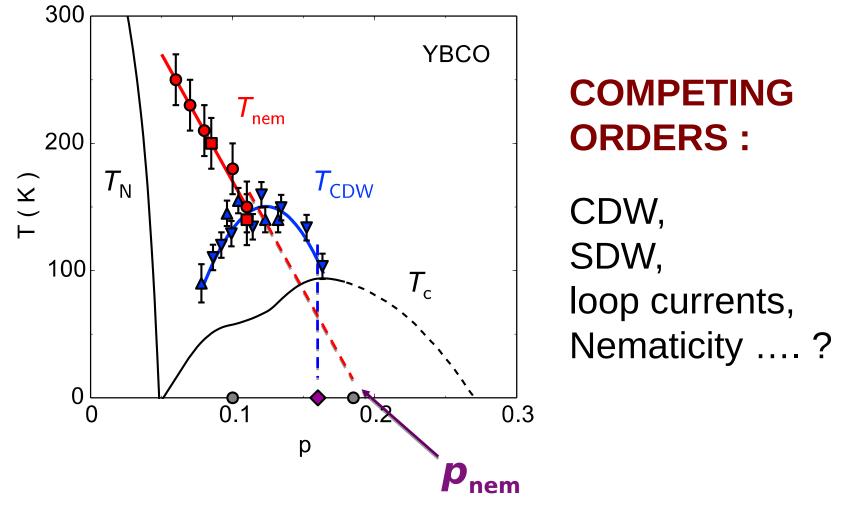


PseudoGap vs SupraGap: a long-standing open issue



IS THE HIGH-TC MECHANISM BURIED ALREADY IN THE NORMAL PG PHASE ? CAN WE UNDERSTAND PG ?

PG phase: the longstanding search for a hidden order



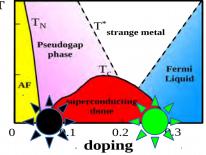
Cyr-Choinière et al., arXiv:1504.06972 (2015)

OUTLINE

Cluster DMFT :

- 1) WE DO NOT NEED A BROKEN ORDER TO EXPLAIN PG
- 2) PG vs SuperConductivity (different beasts? friends ?...):

FOCUS ON EXTREMES OF THE DOME



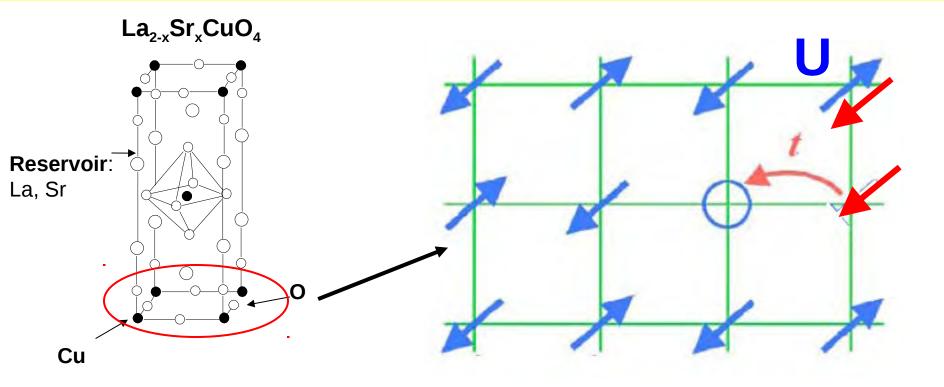


Low doping: <u>non-nodal PG</u> vs <u>nodal SC gap</u> $^{\circ}$ theory as compared to Raman experiments...

igh doping: Raman collapse of PG at a VanHove Sing.

CAN the interplay between PG and SC BOOST UP TC ?.

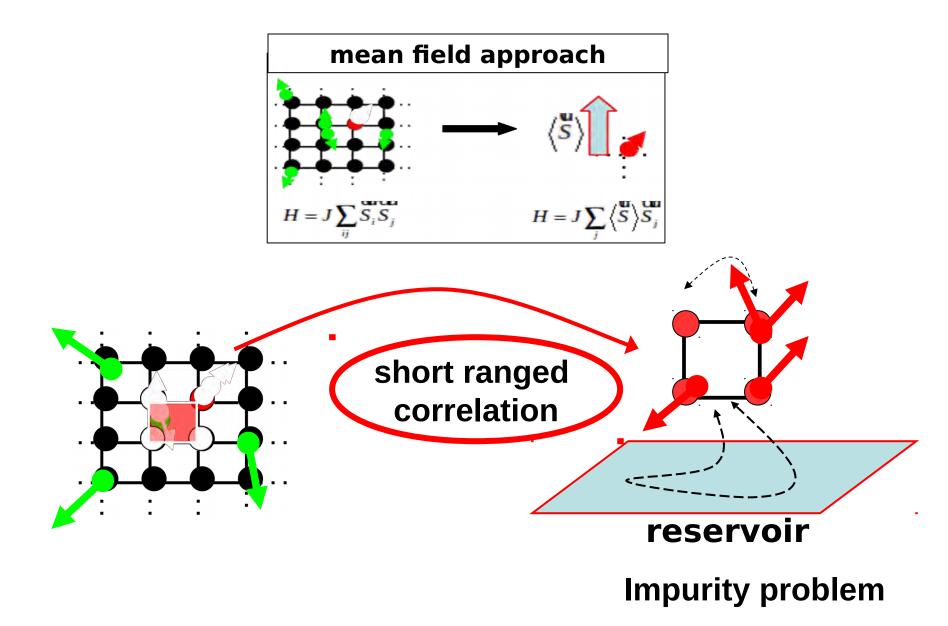
Theoretical play-ground model



2D Hubbard Model $H = -t \sum_{ij\sigma} c^{\dagger}_{i\sigma} c_{j\sigma} + U \sum_{i} n_{i} n_{i} n_{i} n_{i}$

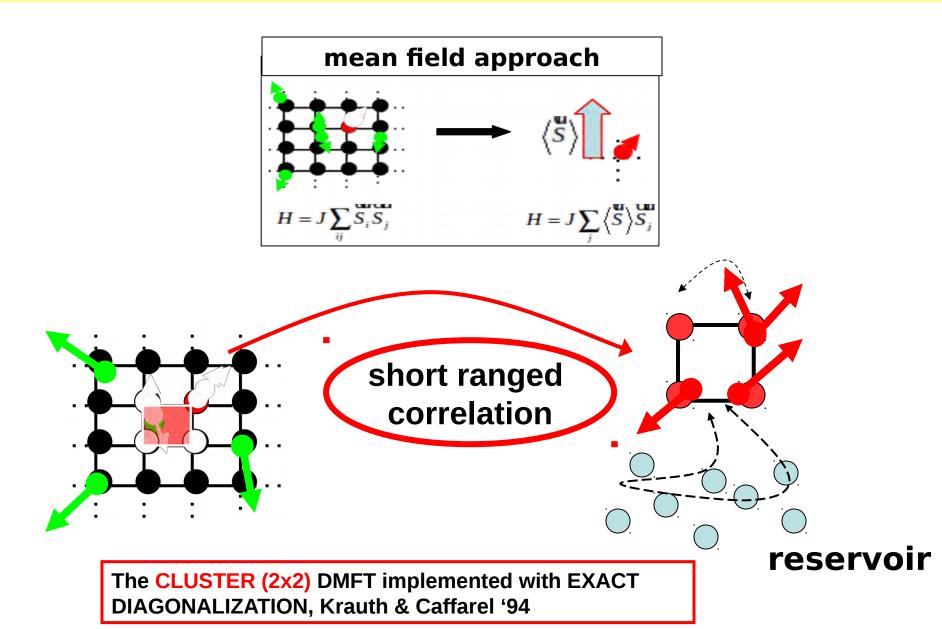
The CLUSTER (2x2) DMFT

Th. Maier et al., Rev. of Mod. Phys. 77,1027 '05; G. Kotliar et al., Rev. of Mod. Phys. 78, 865 '06.



The CLUSTER (2x2) DMFT

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The CLUSTER (2x2) DMFT

GOAL: the one-article propagator \rightarrow **SPECTRAL FUNCTION** $G(k,\omega) = \langle \langle c^{\dagger}_{-k\sigma} c_{k\sigma} \rangle \rangle (\omega) \qquad A(k,\omega) = -\Im G(k,\omega)$

$$G_{loc}(\omega) = \langle \langle c_{1\sigma}^{\dagger} c_{1\sigma} \rangle \rangle \langle \omega \rangle$$

$$G_{n.n}(\omega) = \langle \langle c_{1\sigma}^{\dagger} c_{2\sigma} \rangle \rangle \langle \omega \rangle$$

$$G_{n.n.}(\omega) = \langle \langle c_{1\sigma}^{\dagger} c_{3\sigma} \rangle \rangle \langle \omega \rangle$$

CDMFT (2x2) gives energy ω and approximate k-dependence:

$$G(k,\omega) = G_{loc} + G_{n.n} (\cos k_x + \cos k_y) + G_{n.n.n} \cos(k_x) \cos(k_y)$$

DMFT single site Delicate point !

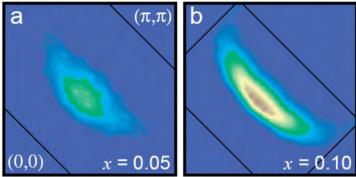
Normal state

PSEUDOGAP and Fermi arcs

Civelli et al: PRL 95, 106402 (2005)

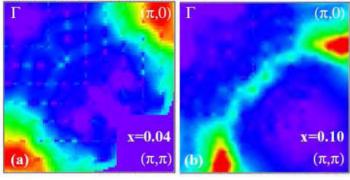
ARPES Experiments

hole doped



K.M. Shen et al., Science '04

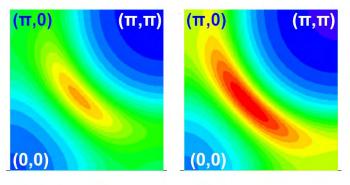
e-doped



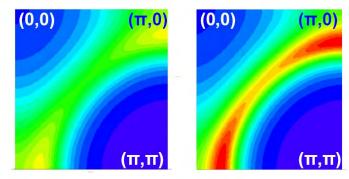
N.P. Armitage et al., PRL '02

 $A(k, \omega = 0)$ CDMFT Theory

hole doped

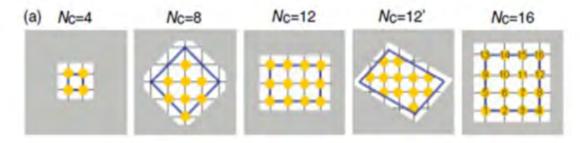


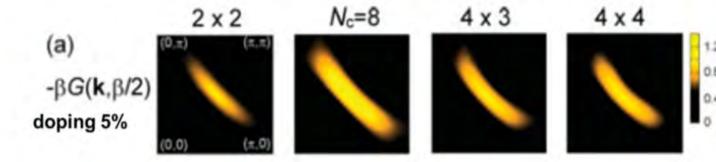
e- doped



Increasing the cluster size

Shiro Sakai et al. PRB 85, 035102 (2012) CDMFT + Continuous Time QMC





The Pseudogap is a solid feature!

Many groups have been developing **CDMFT (DCA)** to study the **normal and SC state of the 2D Hubbard Model**

Mark Jarrell and co-woekers (DCA, large clusters) G. Kotliar and K. Haule group in Rutgers (CDMFT)

A.-M. Trambley recently in collaboration with G. Sordi and K. Haule (CDMFT CTQMC)

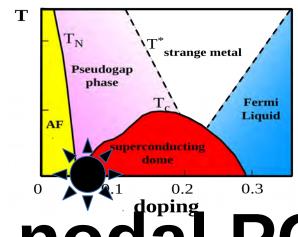
A.J. Millis and E. Gull (DCA CTQMC)

A. Georges, O. Parcollet, M. Ferrero (from small to large cluster....CTQMC and DCA)

S. Sakai and M. Imada (CDMFT, ED, CTQMC...

Everybody agrees on the gross picture of the PG

Underdoped side





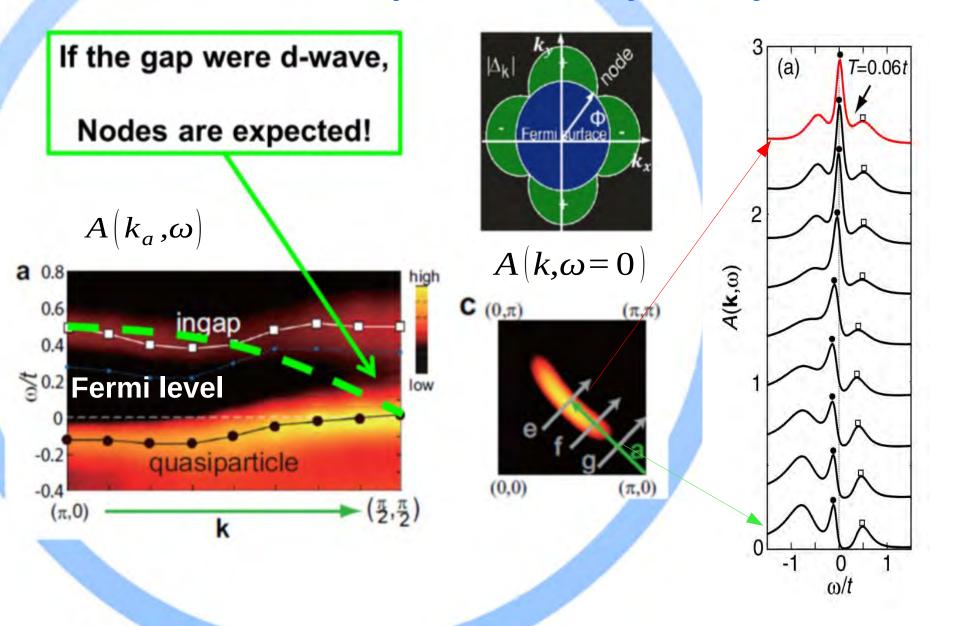
Non-nodal PG v.s. d-wave nodal SC gap

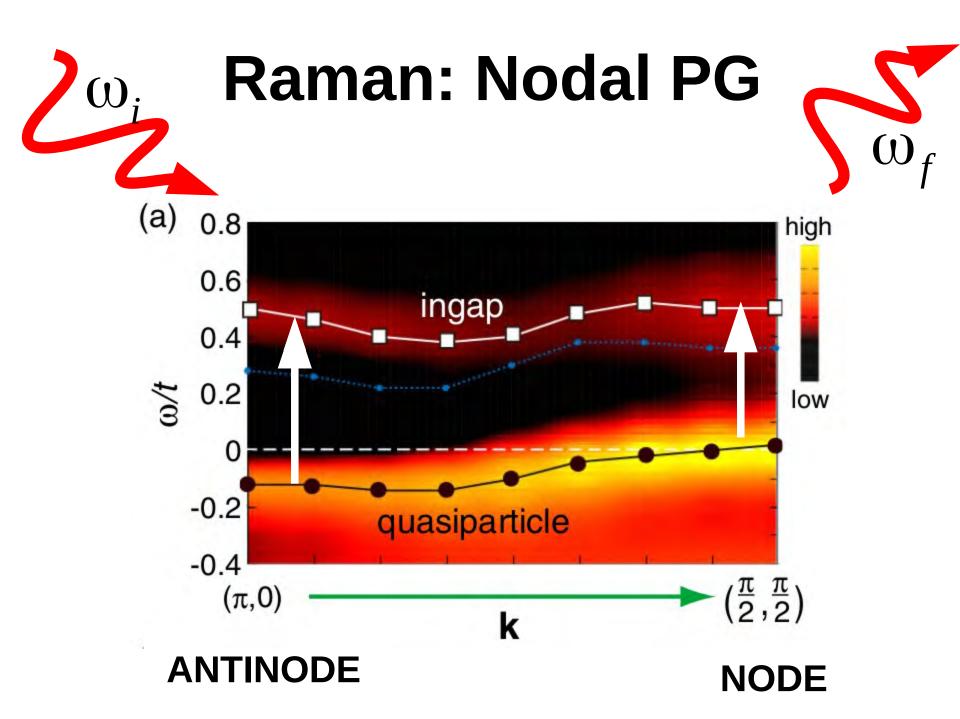


S. Sakai et al. Phys. Rev. Lett. 111, 107001 (2013) Experiments : Sacuto group Paris 7

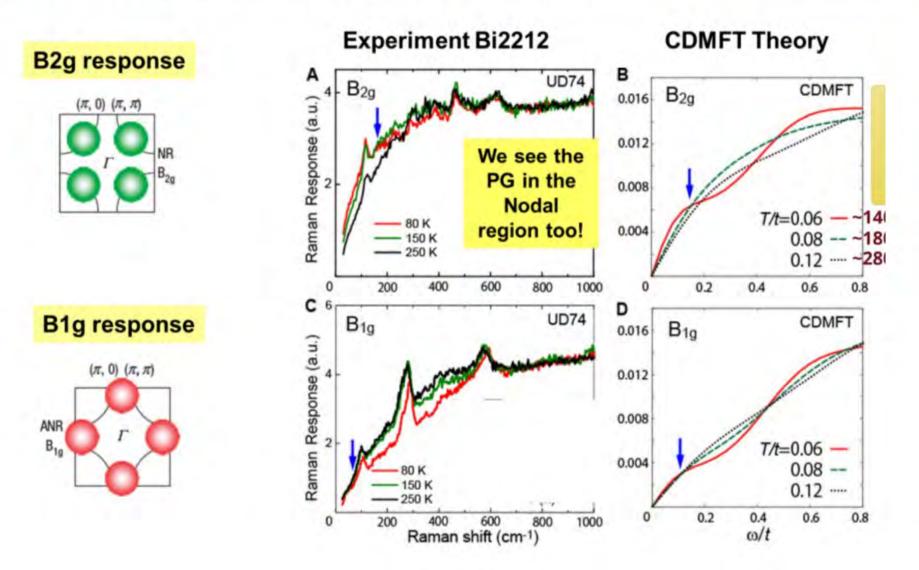


PG abd SC gap have different structures in the nodes and different particle-hole symmetry

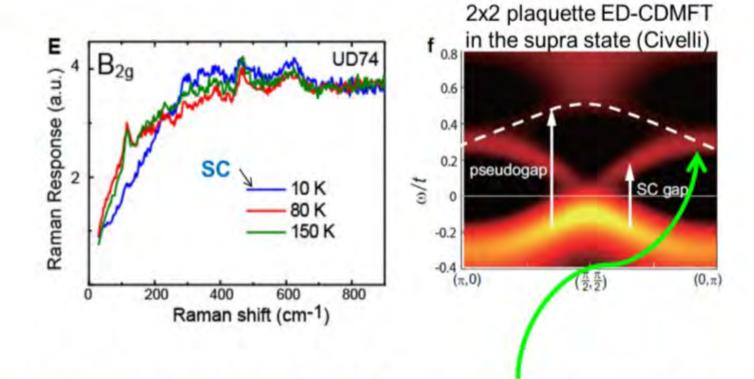




RAMAN RESPONSE : 1) Polarized light sees different regions of k-space 2) Access the ω> 0 unoccuppied spectrum



Pseudo and superconducting gaps in Raman B2g (nodes) Alain Sacuto SQUAP group ay Paris 7

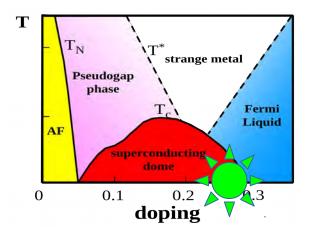


The opening of the SC d-wave gap (Bogoliubov bands formation) appears to

insert electronic states within the s-wave pseudogap at the nodes (B2g Raman)

See also recent arXiv I. Gull, Olivier Parcollet, A. J. Millis, 16-site DCA In the superconducting state

Overdoped side



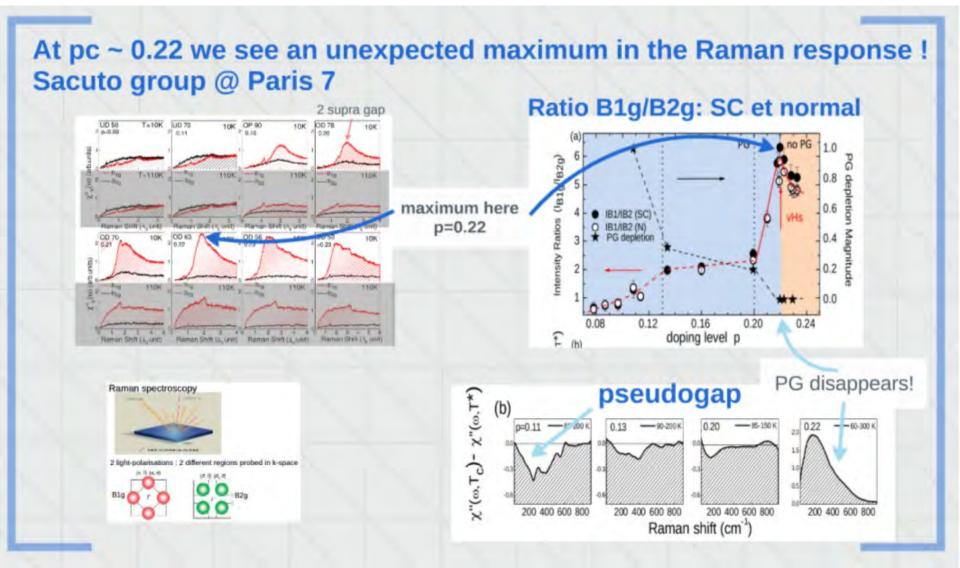
Collapse of the PG @ a Lifshitz transition

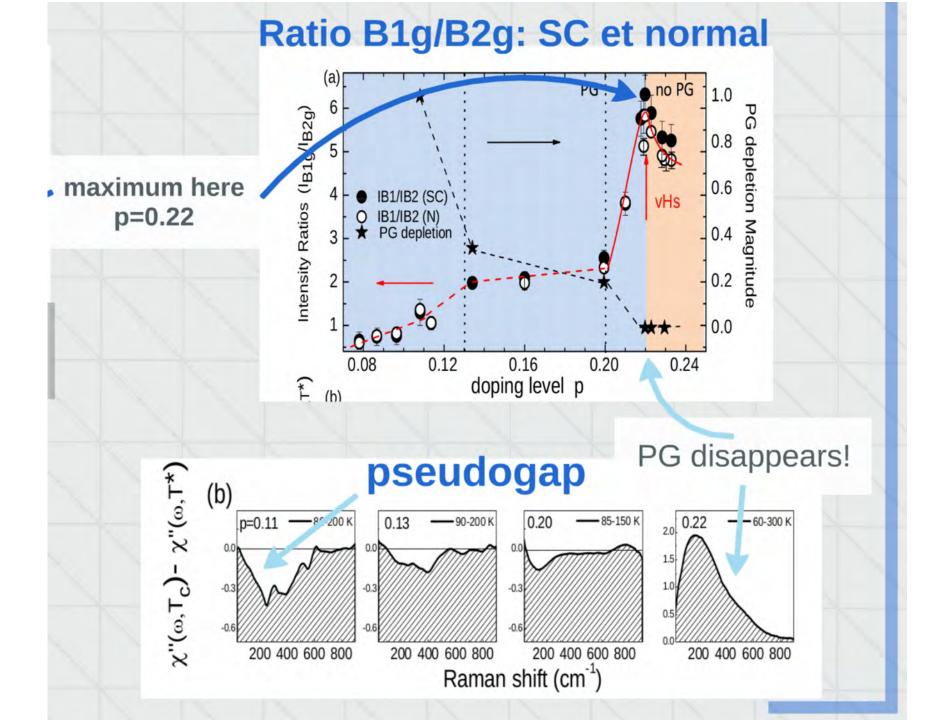




Raman experiments of the SQUAP gropu at Paris Diderot (**Benhabib**, Sacuto...) Theoretical interpretation : I. Paul, Paris Diderot Physical Review Letters 114, 147001, 2015 Experiments by the Squap team, Paris Diderot

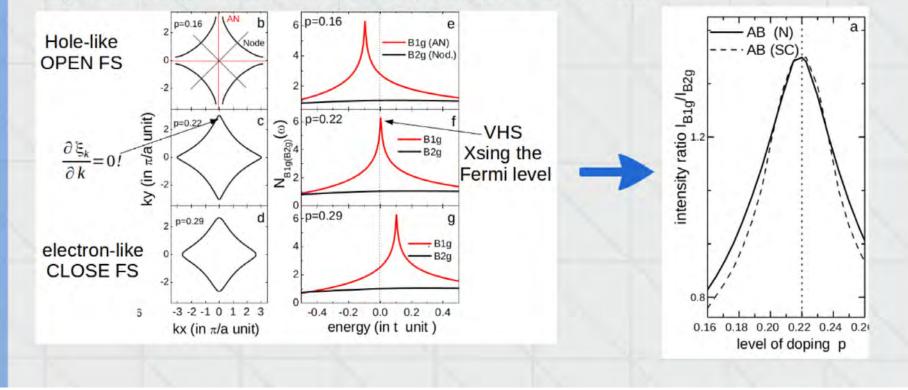
S. Benhabib et al. PRL 114, 147001, 2015





Theory: Vhs and Lifshitz transition

tight-binding Model of Bi2212, collab. with I. Paul



S. Benhabib et al. PRL 114, 147001, 2015

CONCLUSIONS

Van hove singularity can explain the maximum in the B1g reponse !!

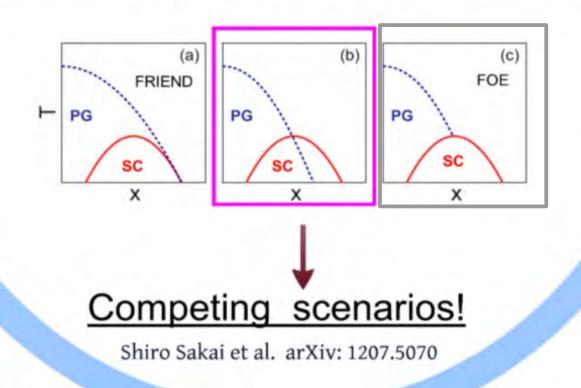
This point marks a change in topology of the Fermi surface (Lifshitz transition), no QCP needed

The SC seems not much affected by this transition (and this explains the differences among different materials...)

SC and PG are different animals!! (but no QCP needed)

Conclusions: PG and SC looks having a different structure

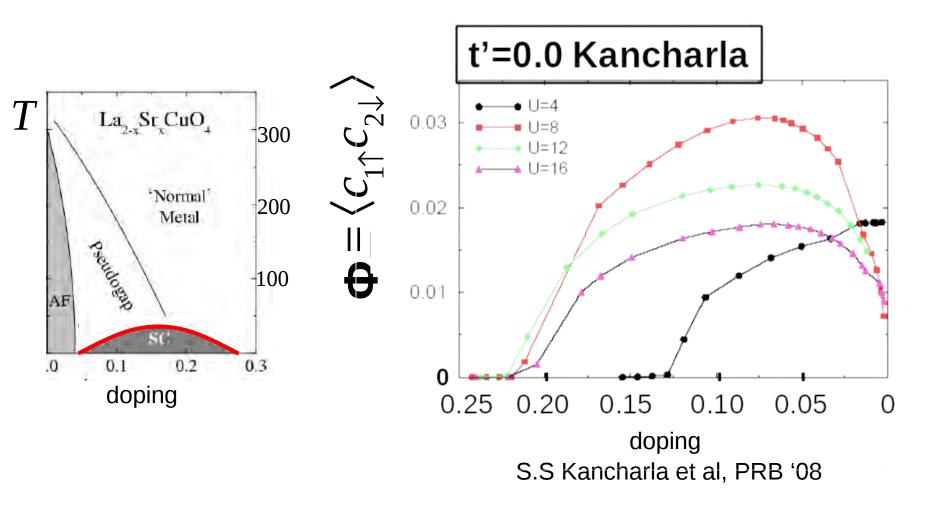
The pseudogap: friend or foe of high Tc? Norman, Pines, Kallin (2005)



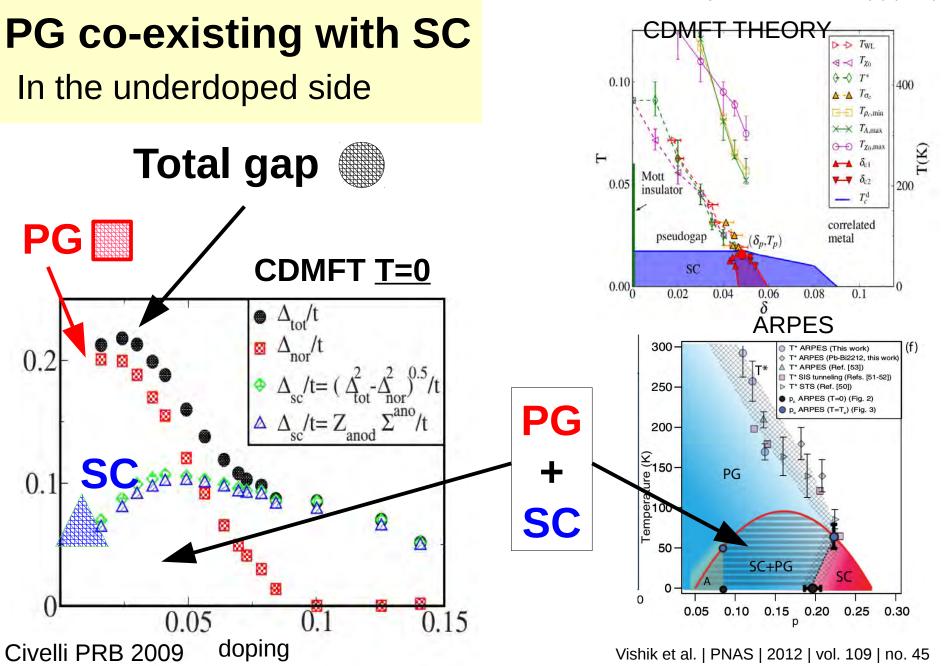
PG and Superconducting state

Superconducting d-wave state!

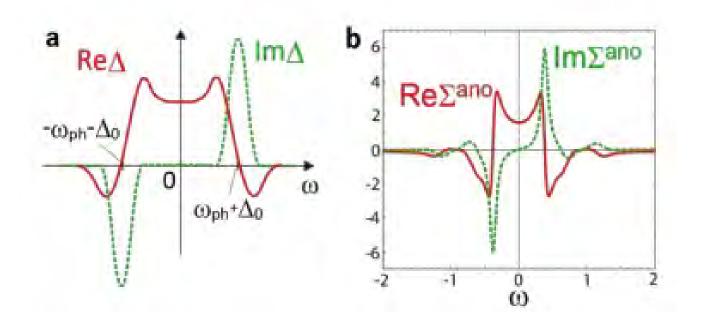
order parameter in the Hubbard Model $\phi = \langle c_1 \uparrow c_2 \downarrow \rangle$



Sordi-Tremblay-Haule CDMFt CTQM Phys. Rev. B 87, 041101(R) (2013)

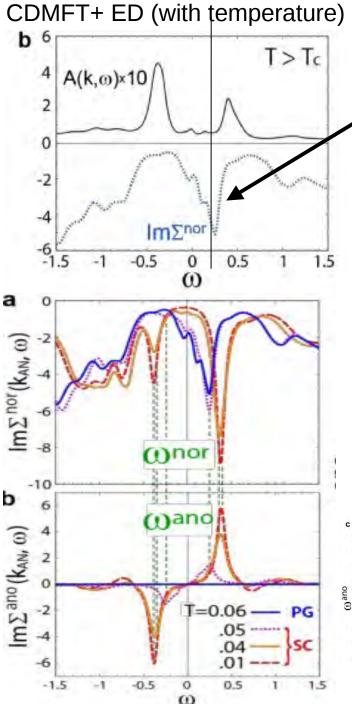


Pairing function (2x2 cdmft U=8t) Sakai, Civelli, Imada, http://arxiv.org/pdf/1411.4365.pdf





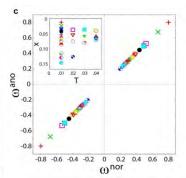
The **Peak** in ImΣano is in standard (Migdal-Eliashberg) theory associated to the boson mediating the pairing interaction $\text{Re}\Delta(\mathbf{k},\omega=0) = \frac{2}{\pi} \int_0^\infty \frac{\text{Im}\Delta(\mathbf{k},\omega')}{\omega'} d\omega'$ Phonos, spin-density waves,...



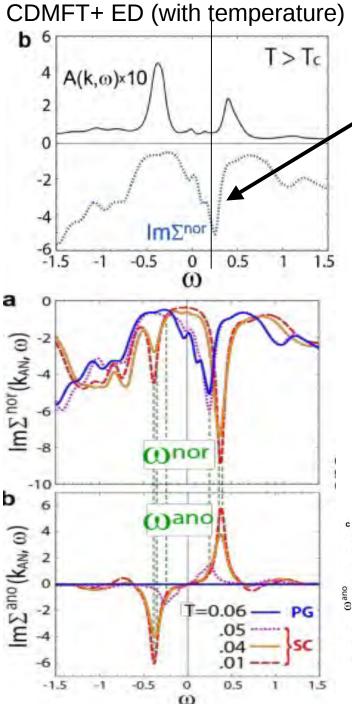
The peak in ImΣnor is responsible for the pseudogap in A(k,ω)

Peak in $Im\Sigma_{ano}$ has been seen already in CDMFT Haule and Kotliar PRB 76 2007 Maier Poilblanc Scalapino PRL 2008 Civelli PRL 2009 Kyung, Senechal, Trambley PRB 2009 Millis and Gull PRL 2013, PRB 2015 **BUT**

Sakai, Civelli, Imada arXiv :1411.4365



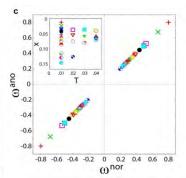
The <u>peak</u> in Im Σ nor and Im Σ ano <u>are the SAME</u>



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Sakai, Civelli, Imada arXiv :1411.4365



The <u>peak</u> in Im Σ nor and Im Σ ano <u>are the SAME</u>

Mechanism driving the PG drives the H-Tc?

$$\Re \Sigma_{ano}(k,\omega=0) = \frac{1}{\pi} \frac{\int \Im \Sigma_{ano}(k,\omega')}{\omega'} d\omega'$$

1) A peak in $Im\Sigma_{ano}$ can push up $Re\Sigma_{ano}$ (ω =0), the peak comes from the PG

$$\Delta(k,\omega) = \frac{\Sigma_{ano}(k,\omega)}{\left(1 - \frac{\Sigma_{nor}(k,\omega) - \Sigma_{nor}(k,-\omega)^{\dagger}}{2\omega}\right)}$$
$$\Delta(k,\omega) = Z(k,\omega) \Sigma_{ano}(k,\omega)$$

2) A peak in $Im\Sigma_{nor}$ implies $Z \rightarrow 0$, destroys Δ (Anderson ? Interpretation in terms of «<u>hidden fermion</u> » by M. Imada

Sakai, Civelli, <u>Imada</u> arXiv :1411.4365

CONCLUSIONS

- 1) Cluster DMFT : WE DO NOT NEED A BROKEN ORDER TO EXPLAIN PG, which comes from Mott physics (pole in Σ_{nor})
- 2) **PG** and **SC** are competing (in overdoped and underdoped sides, agreement with Raman....)
- **3)** The origin of the PG creates peaks in the pairing function which can boost Tc
- 4) At the same time PG=MOTT, quasiparticles and SC are bound to go down approaching the Mott insulator

PG friend AND foe of H-Tc