

Spin Polarized State in the Two-Dimensional Electron Liquid at the Surface of SrTiO_3

M. Gabay

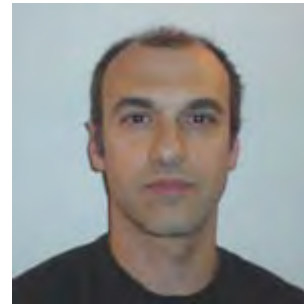


Schloß Waldthausen 06/29-07/02/2015

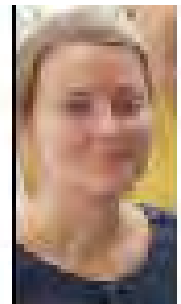
A. Santander-Syro, C. Martins

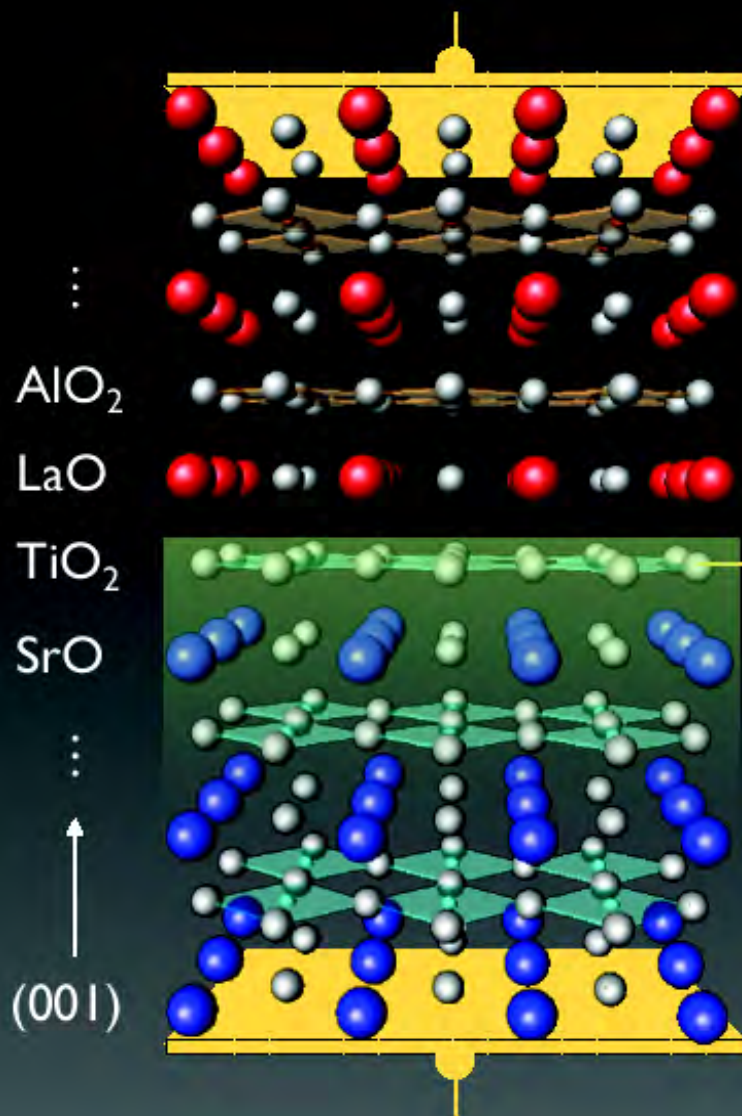


M. Rozenberg, O. Hijano Cubelos



R. Valenti, H. Jeschke. M. Altmeyer
Goethe Universität Frankfurt am Main





LaAlO_3 :

band insulator

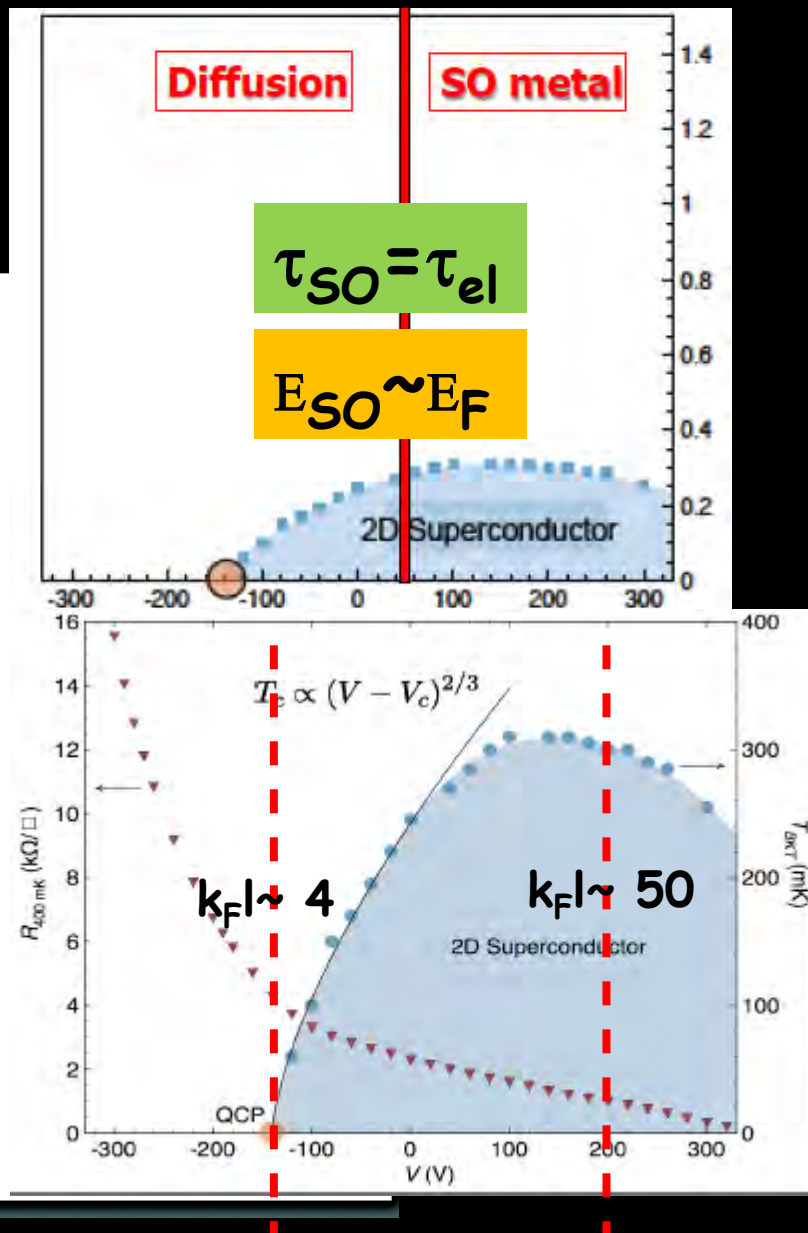
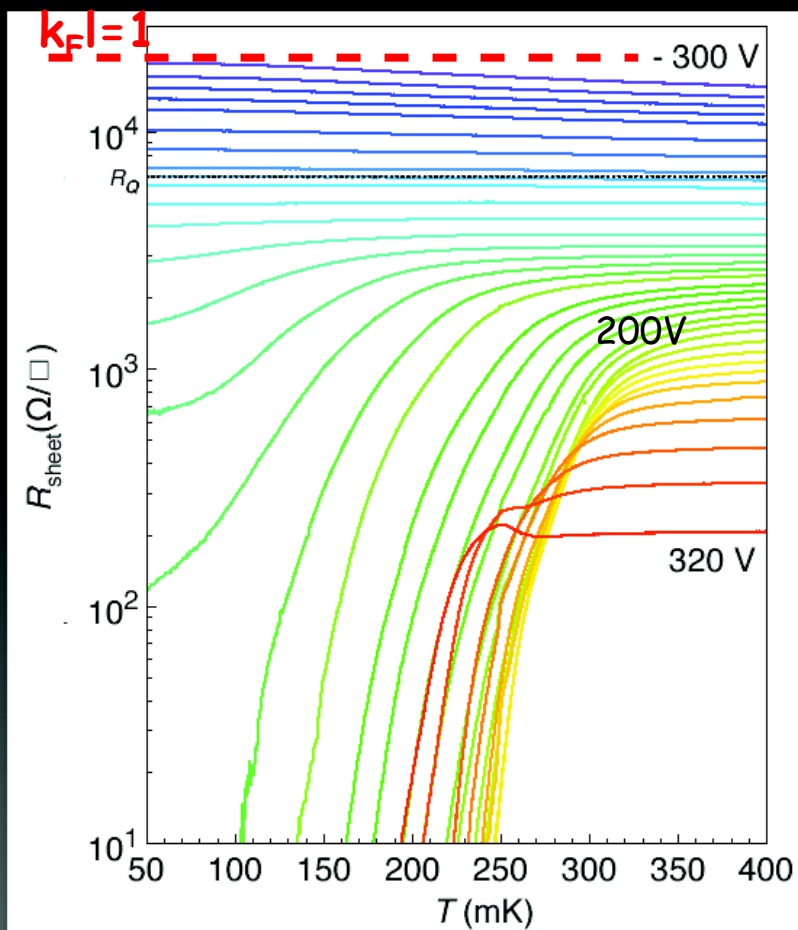
$$\Delta = 5.6 \text{ eV}, \quad \kappa = 24$$

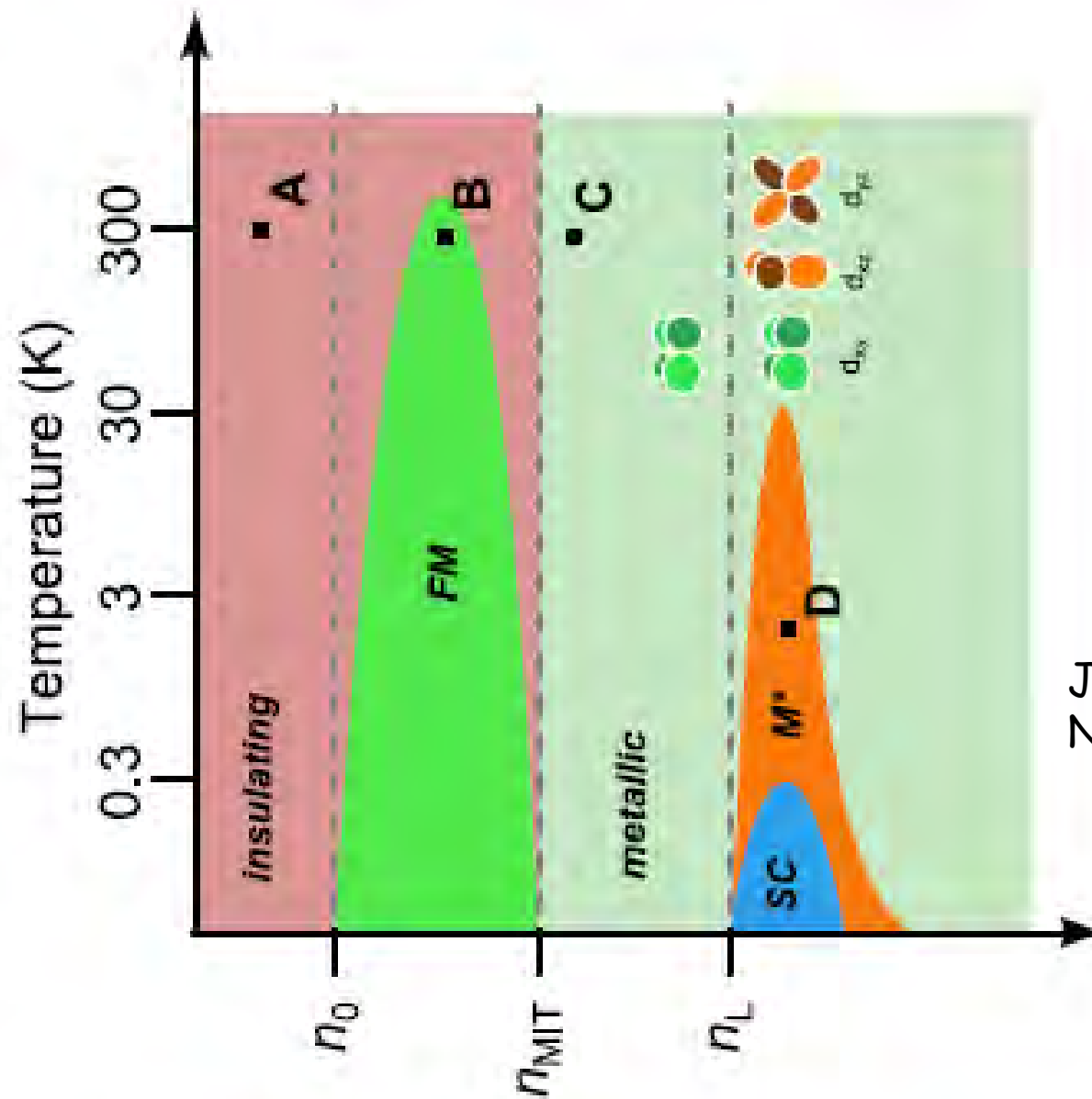
SrTiO_3 :

band insulator

$$\Delta = 3.2 \text{ eV}, \quad \kappa(300 \text{ K}) = 300$$

quantum paraelectric

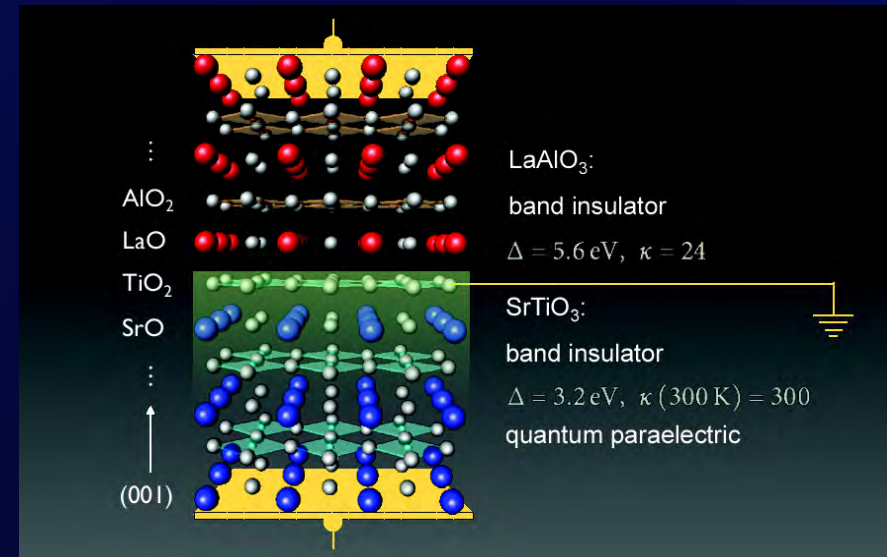
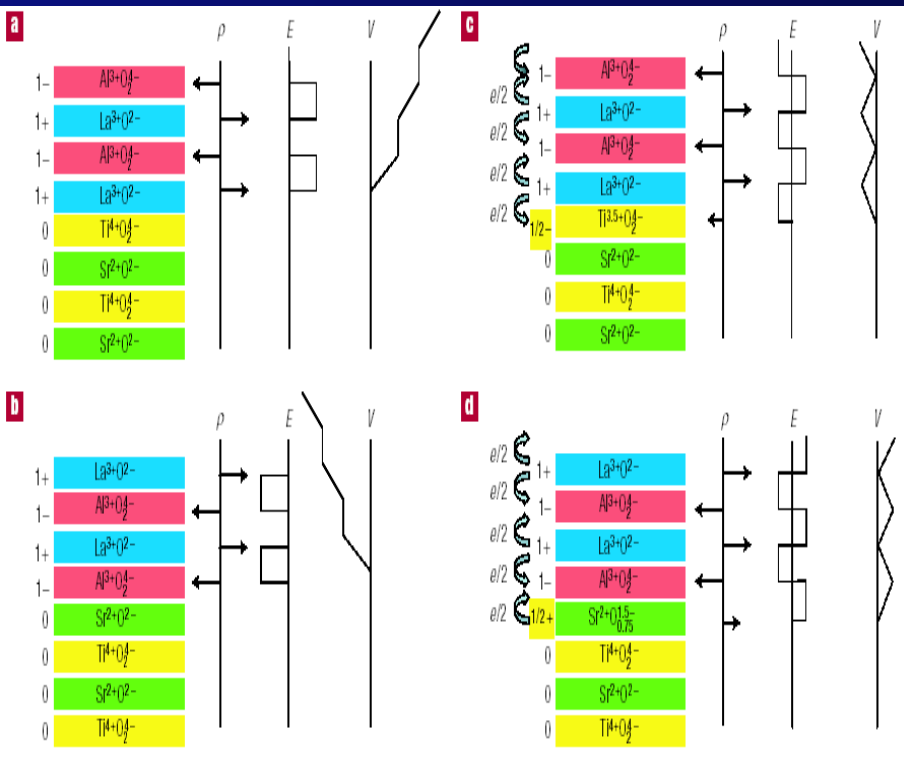


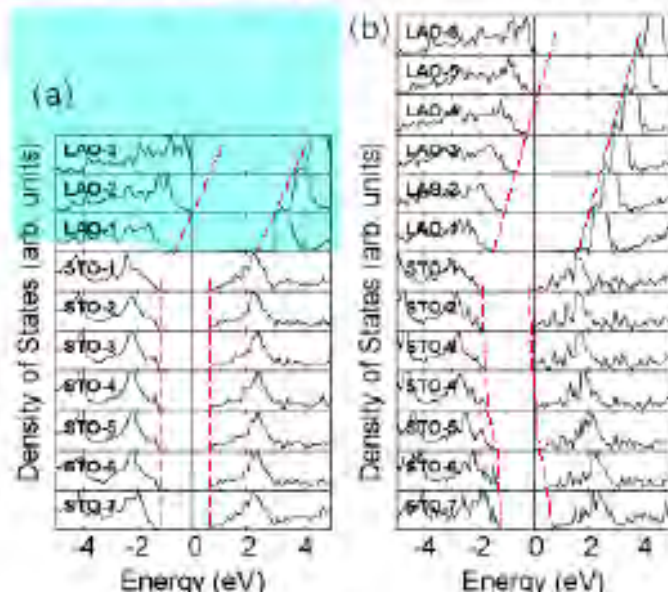
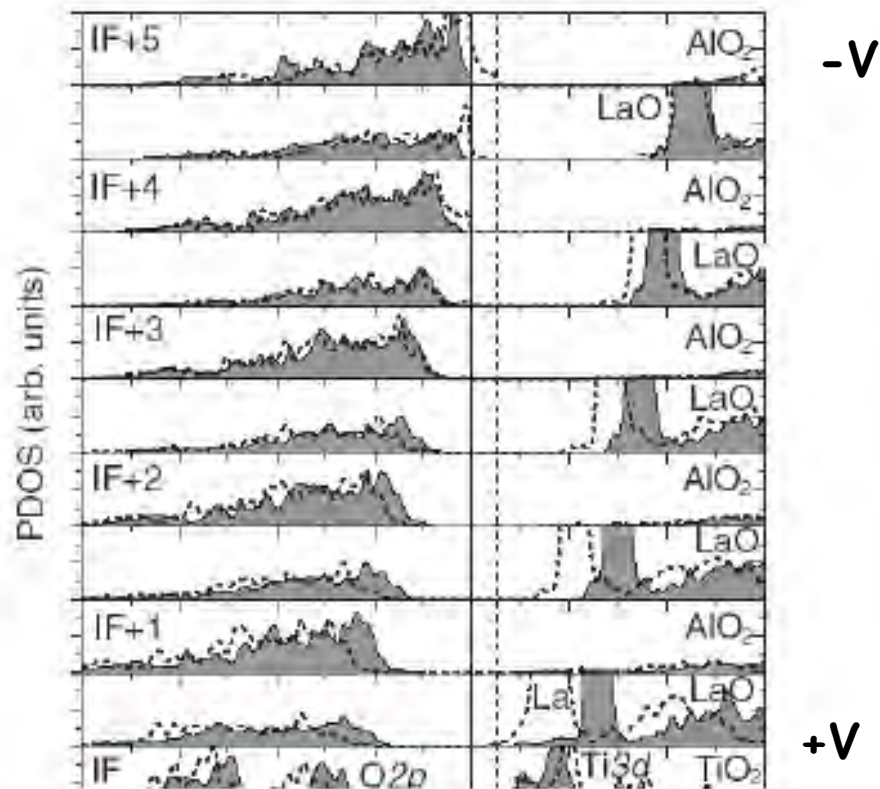


J. Levy et al.,
Nature Communications
5, 5019 (2014)

2DEGS in perovskite oxides

❖ Interface polar-non polar compounds



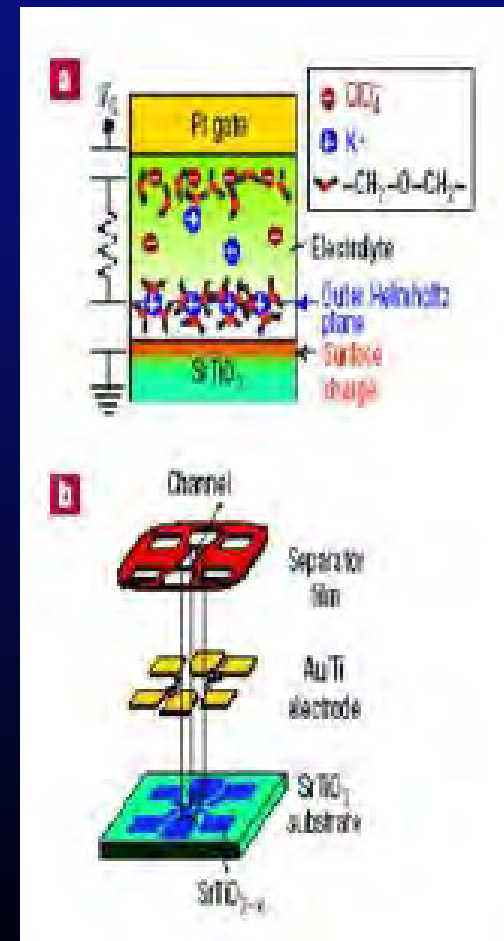


Rossitza Pentcheva¹ and Warren E. Pickett¹

PRL **102**, 107602 (2009)

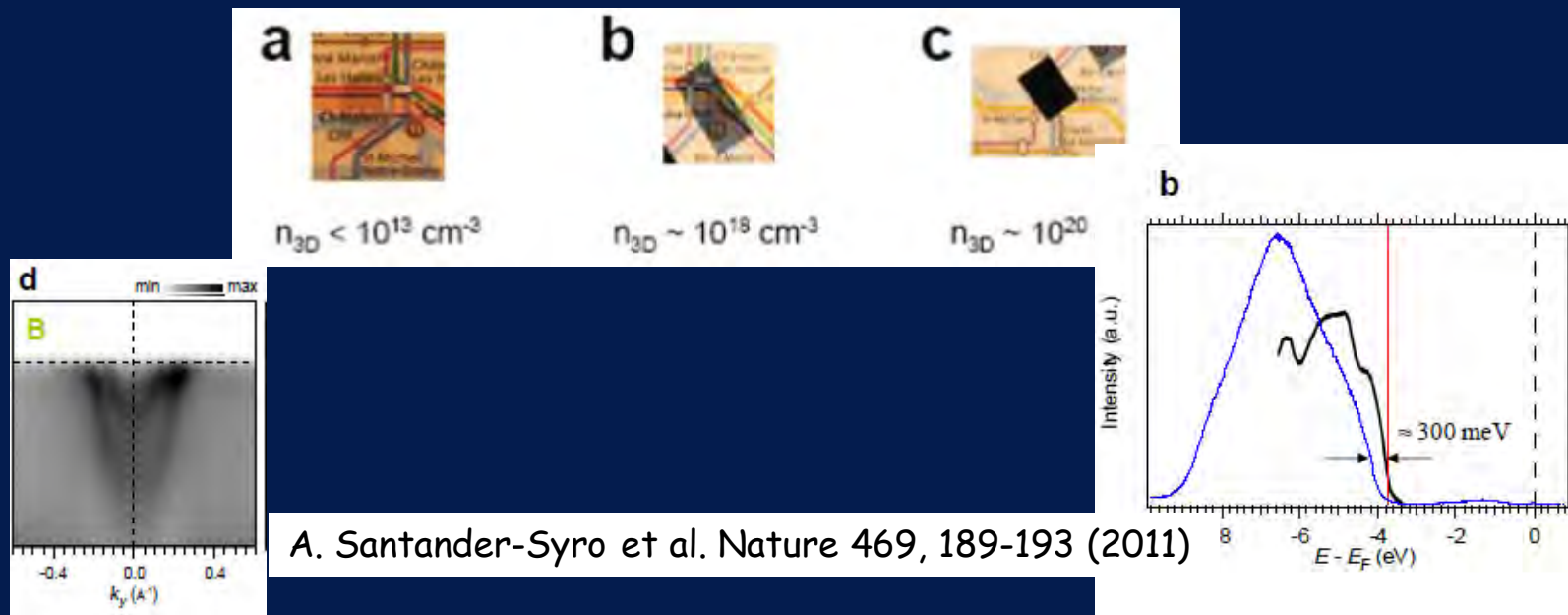
2DEGS in perovskite oxides

- ❖ Interface polar-non polar compounds
- ❖ Surface non polar material

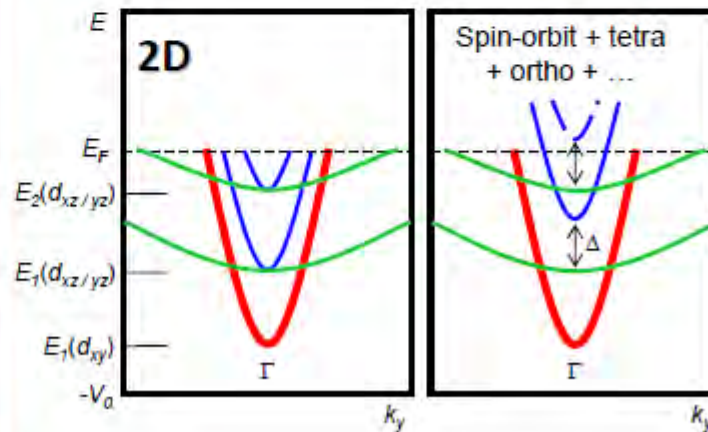
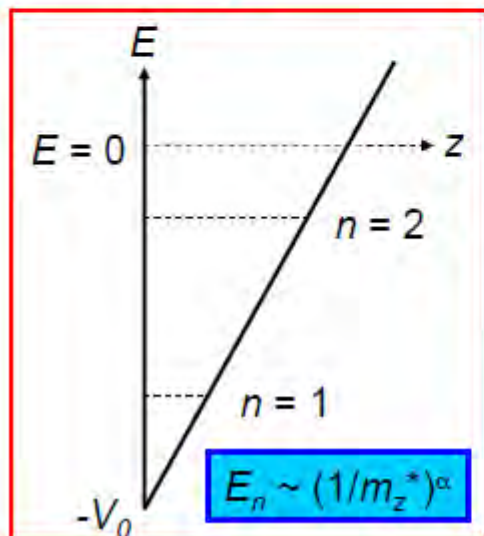
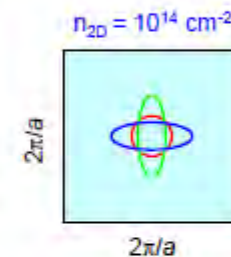
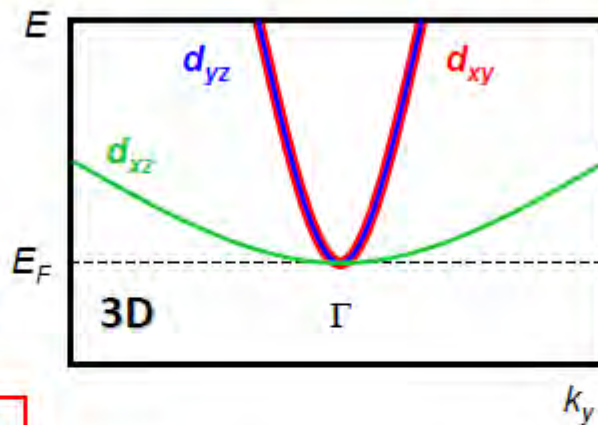
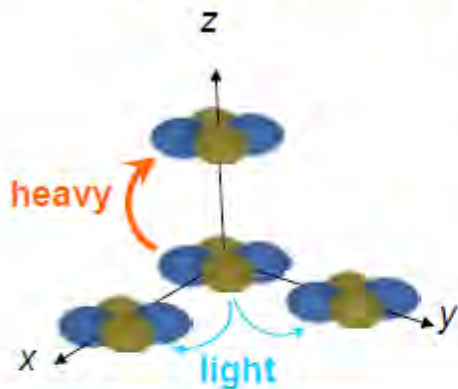


2DEGS in perovskite oxides

- ❖ Interface polar-non polar compounds
- ❖ Surface non polar material
- ❖ Or cleaved surface in high vacuum



SrTiO₃: bulk vs 2D-confined electronic structure



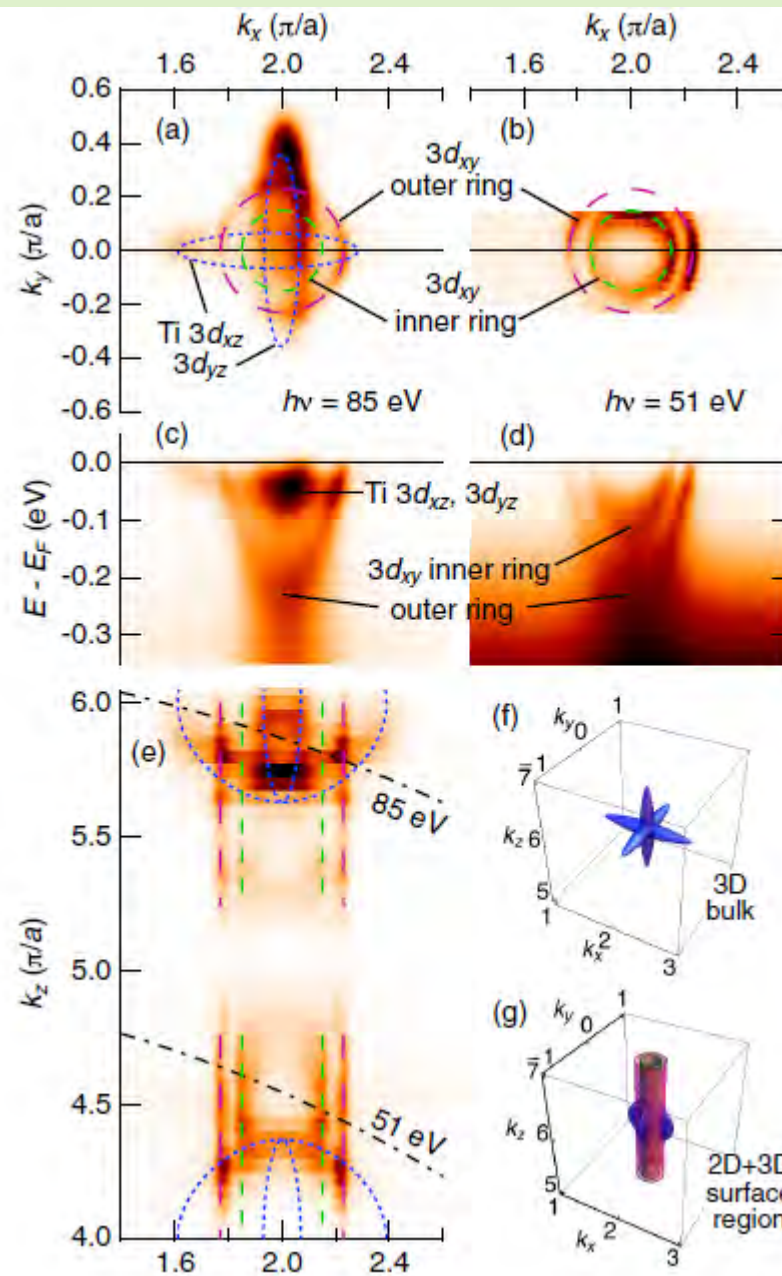
$$4 \text{ u.c.} < L < 8 \text{ u.c.}$$

$$n_{2D} \sim 0.25 \text{ e/u.c.}$$

$F \sim 100 \text{ MV/m}$

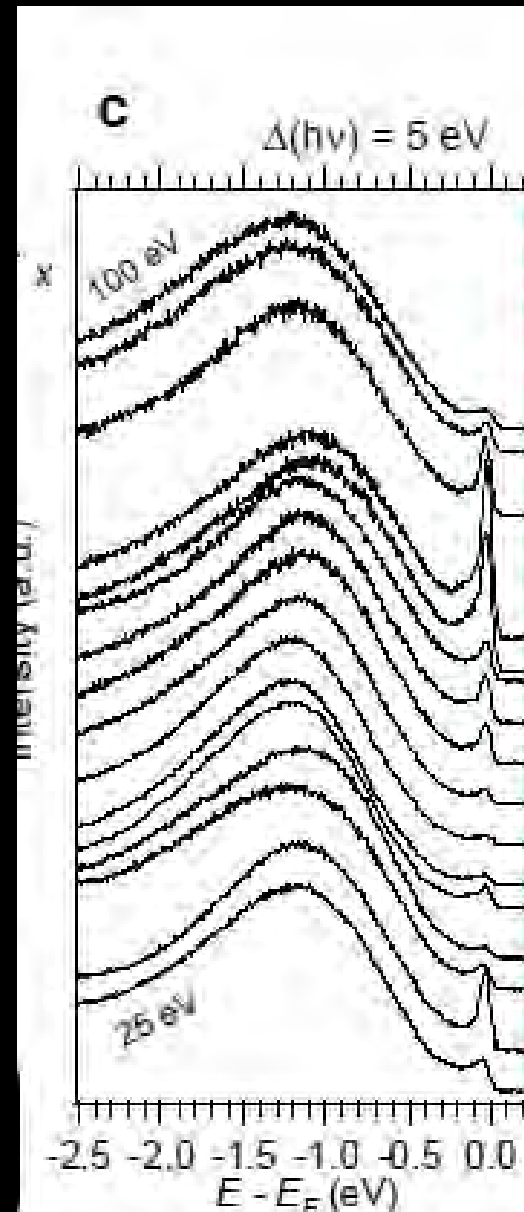
A.F. Santander Syro et al. Nature 469, 189 (2011)

Meevasana et al. Nat. Mater. 10, 114 (2011)



N.C Plumb et al

PRL 113, 086801 (2014)



A. Santander-Syro et al. Nature 469, 189-193 (2011)

Unveiling a two-dimensional electron gas with universal subbands
at the surface of SrTiO_3

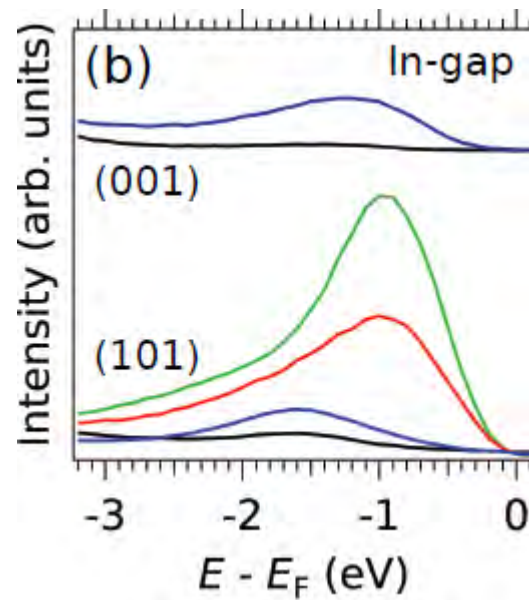


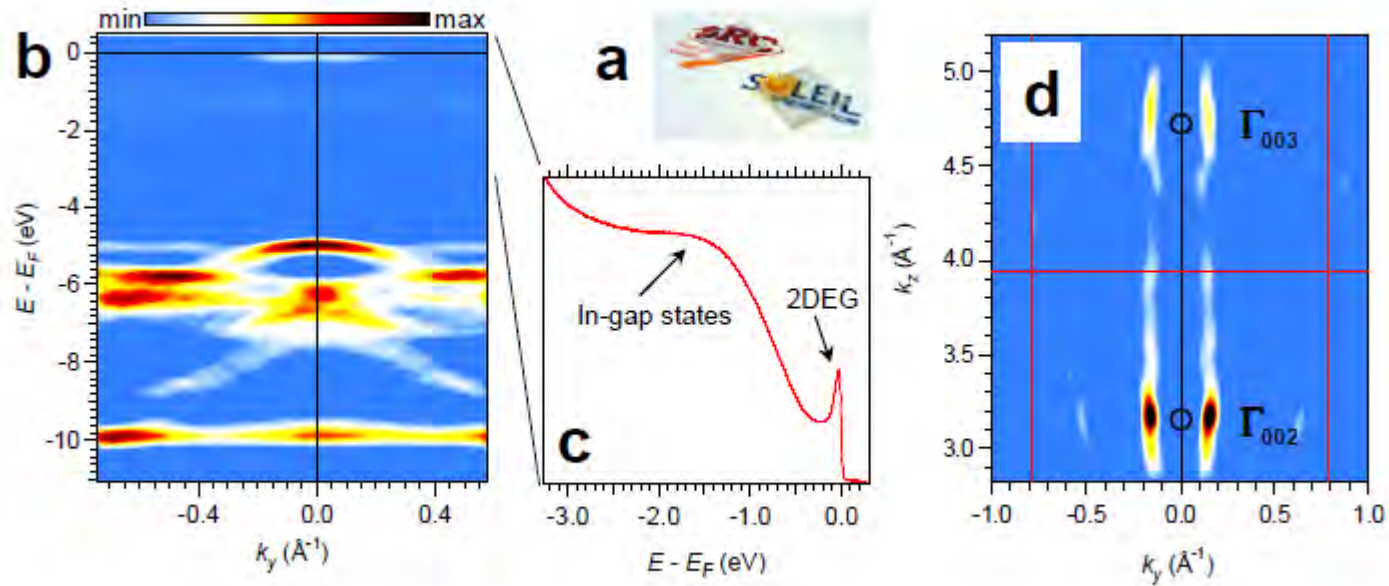
Figure 1: (Color online) Angle integrated spectra of (a) the Ti-3*p* peak, and (b) the in-gap state of the anatase TiO₂ (001) and (101) surfaces (upper and lower curves, respectively) measured at $h\nu = 100$ eV. The black curve was measured shortly after the first light exposure, and the blue, red and green curves at sequentially increasing times respectively. (c) Binding energy of the Ti⁴⁺ peak and concentration of oxygen

Light engineering of two dimensional electron gases at the (001) and (101) surfaces of TiO₂ anatase

T. C. Rödel,^{†,‡} F. Fortuna,[†] F. Bertran,[‡] M. Gabay,[¶] M. J. Rozenberg,[¶]

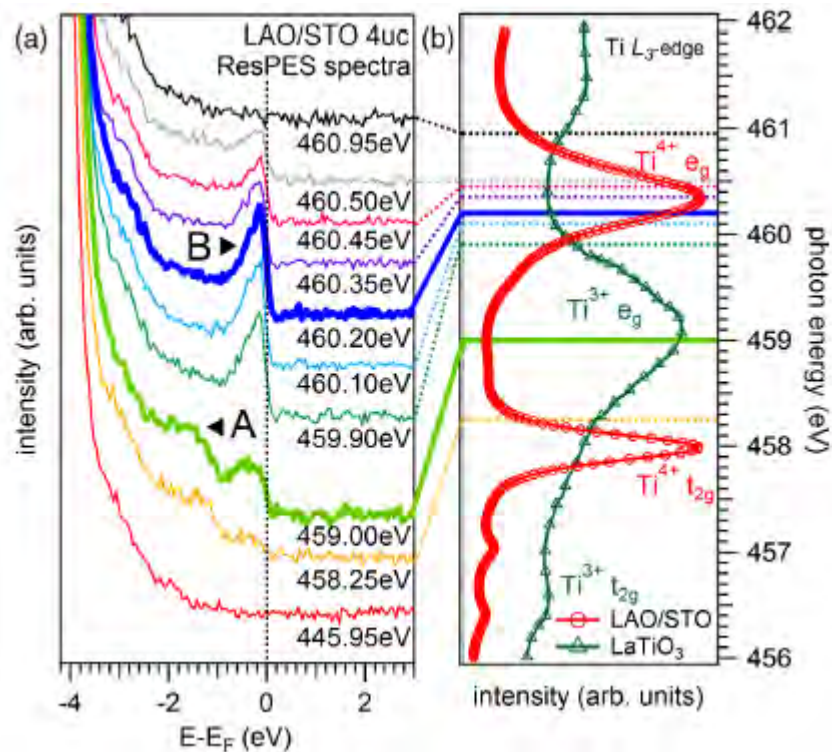
A. F. Santander-Syro,^{*,†} and P. Le Fèvre[‡]

To appear in Phys. Rev. B (R) 2015



Orbital symmetry reconstruction and strong mass renormalization in the two-dimensional electron gas at the surface of KTaO3

A. F. Santander-Syro, et al., Phys. Rev. B 86, 121107(R), 2012

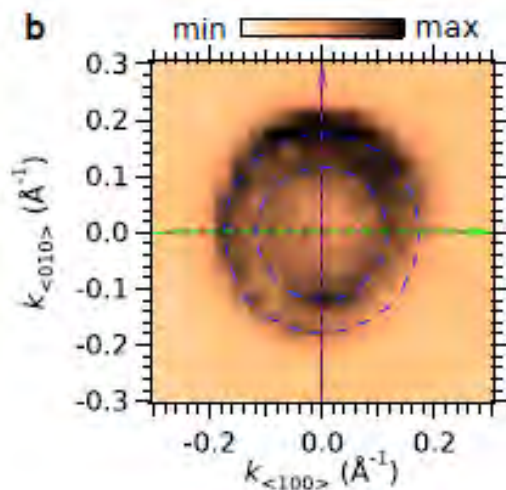
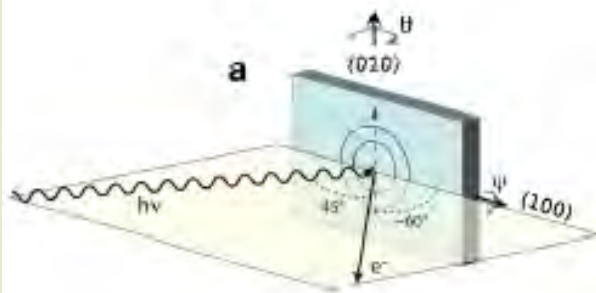


The complete valence band spectrum, measured on resonance ($h\nu = 460.20$ eV), is depicted in Fig. 1(c). The off-resonance spectrum ($h\nu = 445.95$ eV) shows no spectral weight at the chemical potential. Moving through the resonance two structures appear—indicating that both are of Ti 3d character—a broad structure at a binding energy of ≈ 1.3 eV (A) and a structure which is cut off by the Fermi-Dirac distribution function [marked by B in

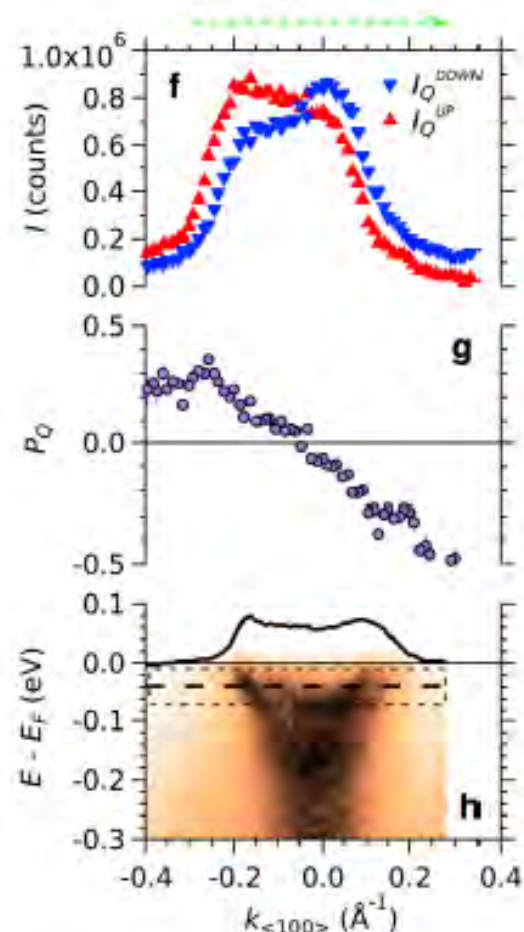
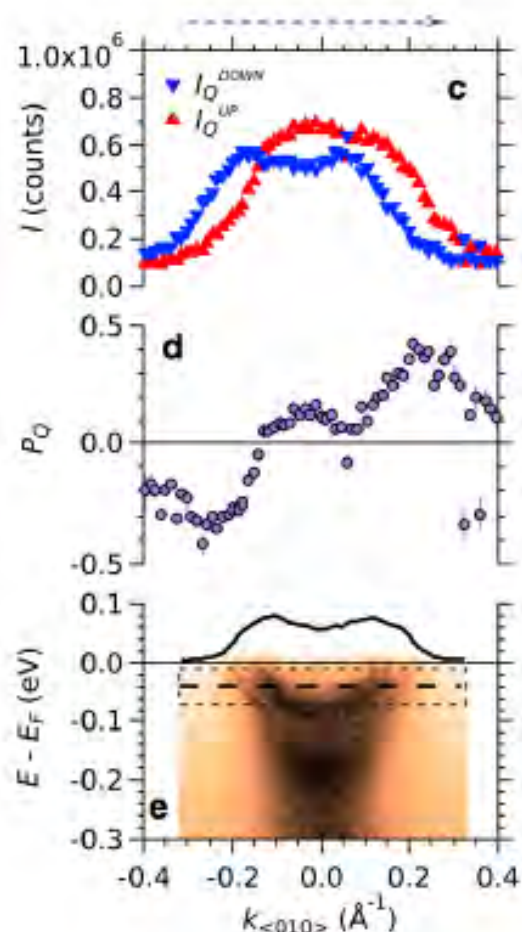
PRL 110, 247601 (2013)

Direct k -Space Mapping of the Electronic Structure in an Oxide-Oxide Interface

G. Berner,¹ M. Sing,¹ H. Fujiwara,² A. Yasui,³ Y. Saitoh,³ A. Yamasaki,⁴ Y. Nishitani,⁴ A. Sekiyama,² N. Pavlenko,^{5,6,7} T. Kopp,⁵ C. Richter,^{5,7} J. Mannhart,⁷ S. Suga,⁸ and R. Claessen¹

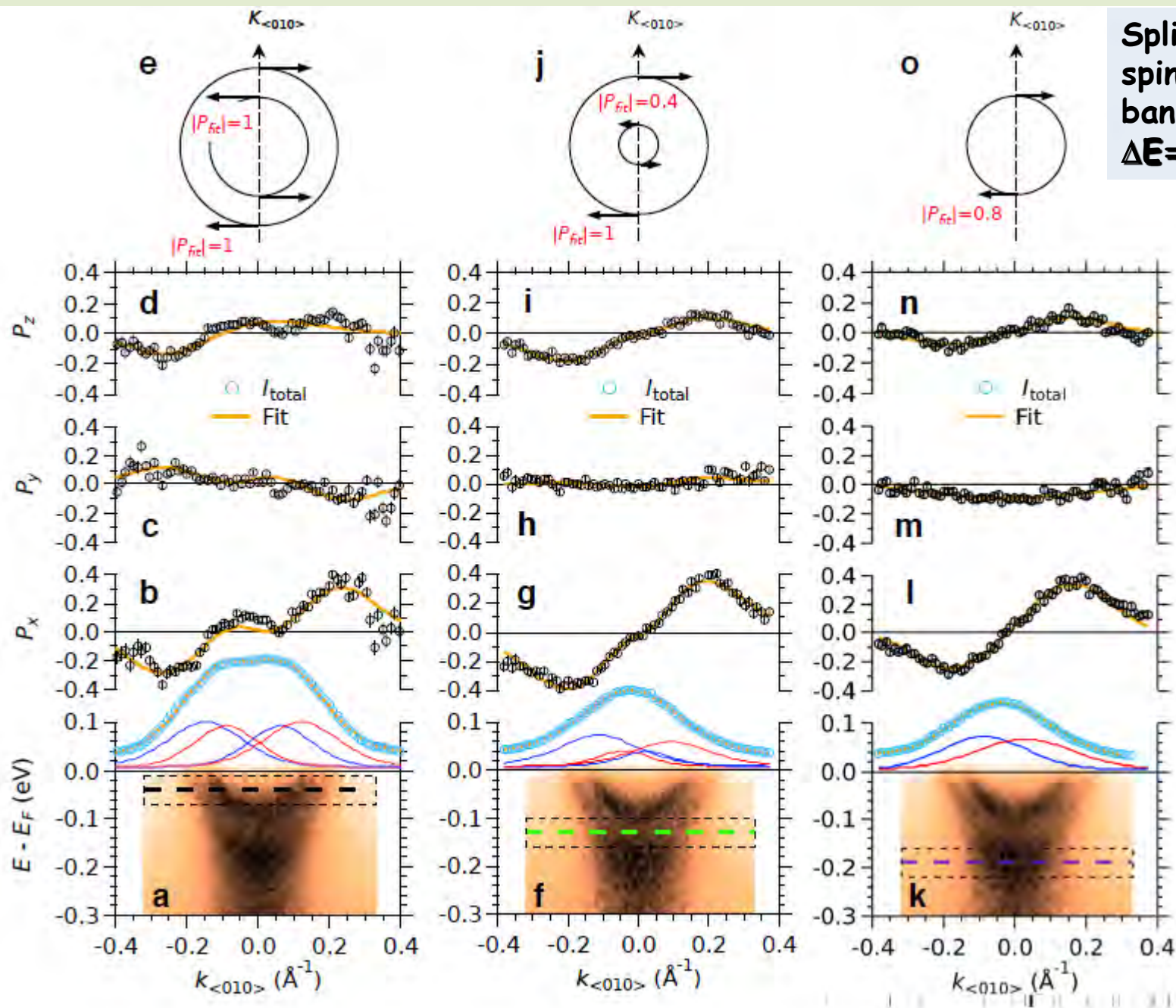


A. F. Santander-Syro, et al.
Nature Materials 13,
1085-1090 (2014)



$$\Delta k_F \sim 0.1 \text{ \AA}^{-1}$$

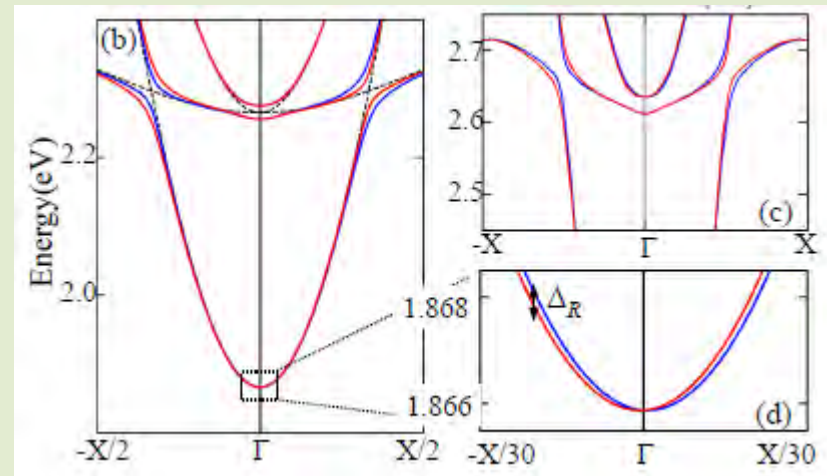
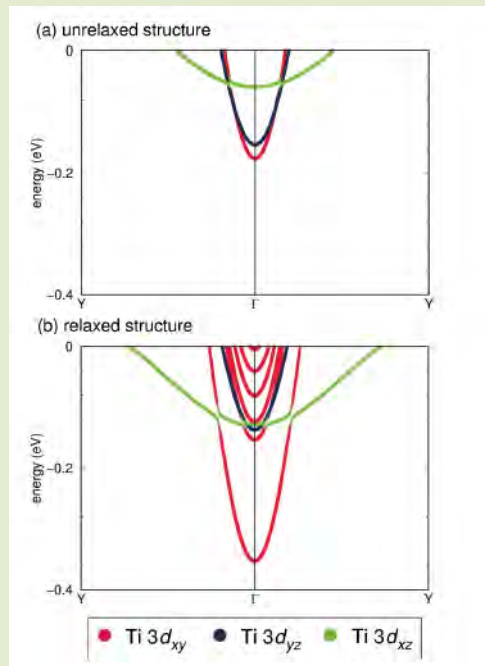
**Splitting of the
spin-polarized
bands:
 $\Delta E=100$ meV**



Oxygen vacancies at the boundary could be a unifying mechanism responsible for the presence of the 2DEG

What role, if any do O vac play in

- Spin-orbit effects (via relaxations or other processes)



Zhicheng Zhong, Anna Tóth, and Karsten Held

Phys. Rev. B 87, 161102(R) (2013)

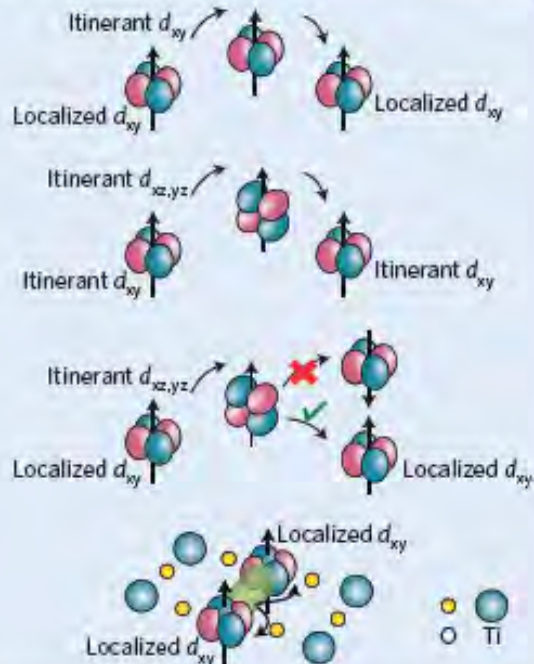
Juan Shen, Hunpyo Lee, Roser Valentí, and Harald O. Jeschke

PHYSICAL REVIEW B 86, 195119 (2012)

What role, if any do O vac play in

□ Magnetism 1

Table 1 | Selected explanations for the magnetism at the LAO/STO interface.



At low concentrations itinerant d_{xy} electrons mediate a ferromagnetic interaction between the d_{xy} spins localized at the interface¹².

The quasi one-dimensional nature of the d_{xz} (or d_{yz}) electrons allows a large magnetic polarization to develop on these orbitals. Hund's rule transfers this polarization to the d_{xy} electron¹³.

For low concentrations of electrons, Hund's rule mediates a ferromagnetic interaction between the d_{xy} spins localized at the interface. With Rashba spin-orbit coupling a spiral state develops⁵.

At the interface, an oxygen vacancy (identified by the green region) releases an electron in each of the neighbouring Ti d_{xy} orbitals. Their spins couple ferromagnetically¹⁴.

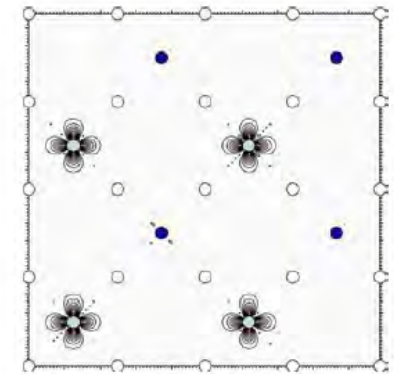


FIG. 11. (Color online) Charge density plot of the charge ordered state localized in the TiO_2 layer of the n -type interface. Ti^{3+} with an occupied d_{xy} orbital and Ti^{4+} marked by light and dark blue (gray) circles, respectively, are ordered in a (45°) checkerboard manner. Oxygens are marked by white circles. The plotted area corresponds to four $c(2 \times 2)$ -simulation cells.

R. Pentcheva,
PRB74, 035112 (2006)

14. Pavlenko, N. *et al.* *Phys. Rev. B* 85, 020407(R) (2012)

Marc Gabay and Jean-Marc Triscone

Nature physics 10.1038/nphys2737 (2013)

Oxygen vacancies at the boundary could be the unifying mechanism responsible for the presence of the 2DEG

What role, if any do O vac play in

□ Magnetism 2

M. Salluzzo et al.

Phys. Rev. Lett. 111, 087204 (2013)

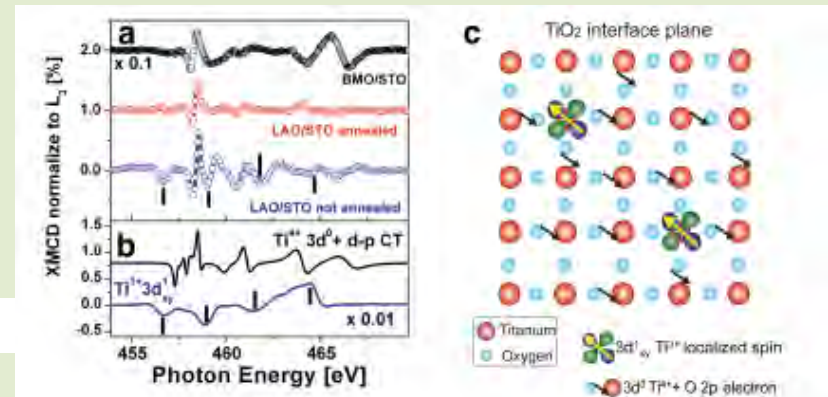


FIG. 4: (color online) Comparison between experimental and theoretical XMCD spectra. (a) Normalized Ti XMCD of BMO/STO (black circles, scaled by a factor 0.1), of standard LAO/STO (red circles) and of oxygen-vacancy rich LAO/STO (blue open circles). (b) Single ion multiplet scattering calculations of the XMCD spectra for Ti^{4+} (black line) and Ti^{3+} (blue line), using an exchange field of -10 meV and +10 meV respectively. (c) A cartoon illustration of the proposed mechanism of magnetism at the titanate heterostructures. In LAO/STO oxygen vacancies lead to localized Ti^{3+} sites, holding sizable spin and orbital magnetic moments thanks to their $3d_{xy}$ electron. The q2DEG gets polarized by these magnetic impurities and mediates a long range magnetic interaction leading to ferromagnetism at low temperature.

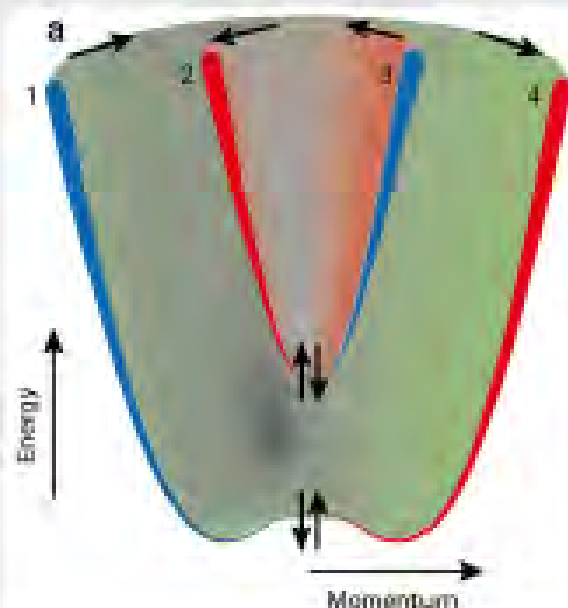
Rashba splitting and magnetism

Energy gap

Hamiltonian and eigenvalues:

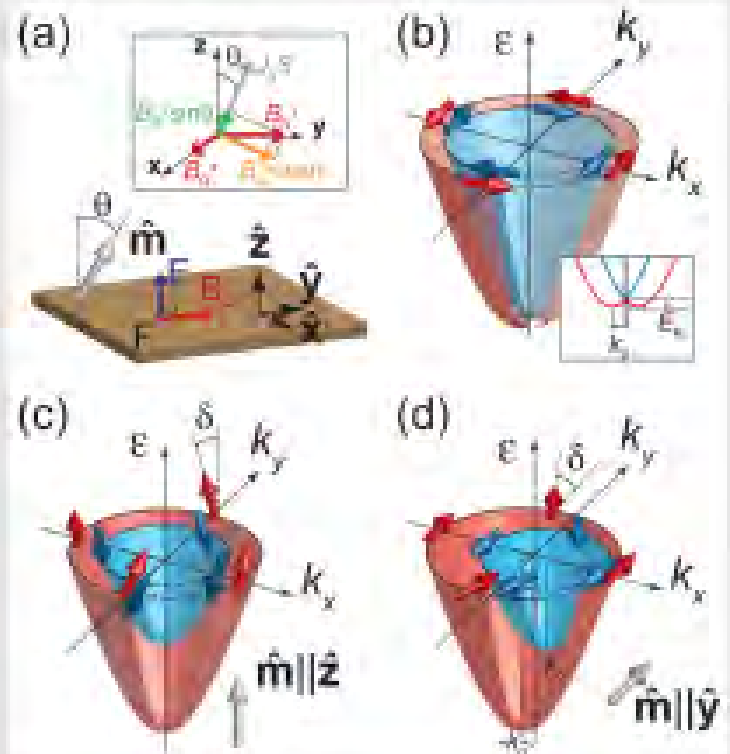
$$\mathcal{H} = \frac{\vec{p}^2}{2m} + \frac{\alpha}{\hbar} (\vec{\sigma} \times \vec{p}) \cdot \hat{z} + \Delta \sigma_z$$

$$\epsilon_{\pm}(\vec{k}) = \frac{\hbar^2 k^2}{2m} \pm \sqrt{\alpha^2 k^2 + \Delta^2}$$



Spin canting

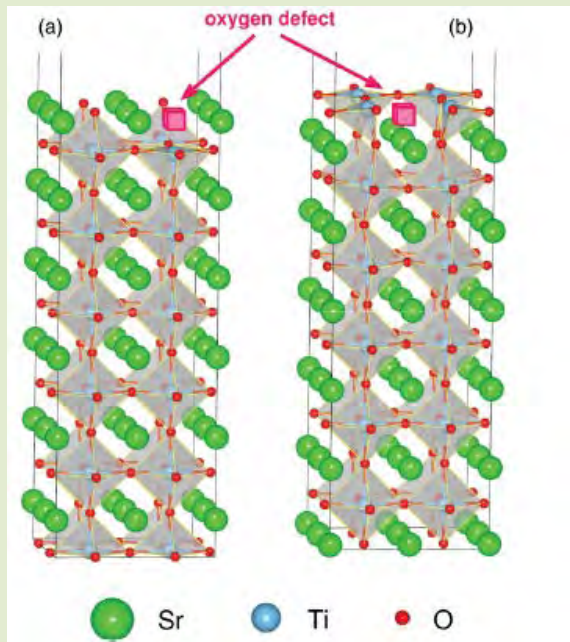
$$\delta = \tan^{-1} \left(\frac{\alpha_R (k_x^2 \cos^2 \theta + k_y^2)^{1/2}}{J_0 S} \right)$$



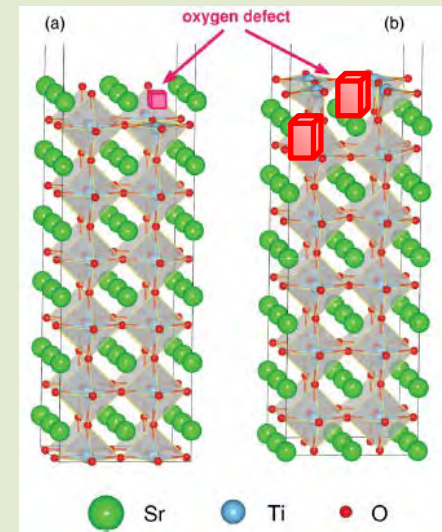
Oxygen vacancies at the boundary could be the unifying mechanism responsible for the presence of the 2DEG

Modeling by

- ❖ DFT for slabs in the presence of O vac (relaxed structures)
- ❖ Finite cluster analytical calculations (assess most relevant orbitals)



Single O vac

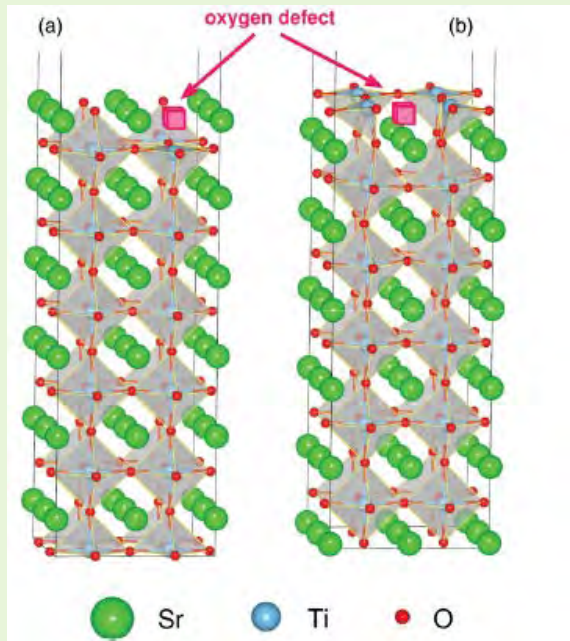


Double O vac

Oxygen vacancies at the boundary could be the unifying mechanism responsible for the presence of the 2DEG

Modeling by

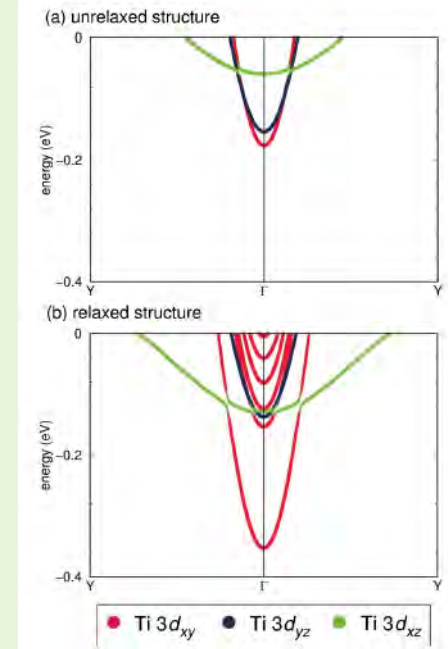
- ❖ DFT for slabs in the presence of O vac (relaxed structures)
- ❖ Finite cluster analytical calculations (assess most relevant orbitals)



Single O vac



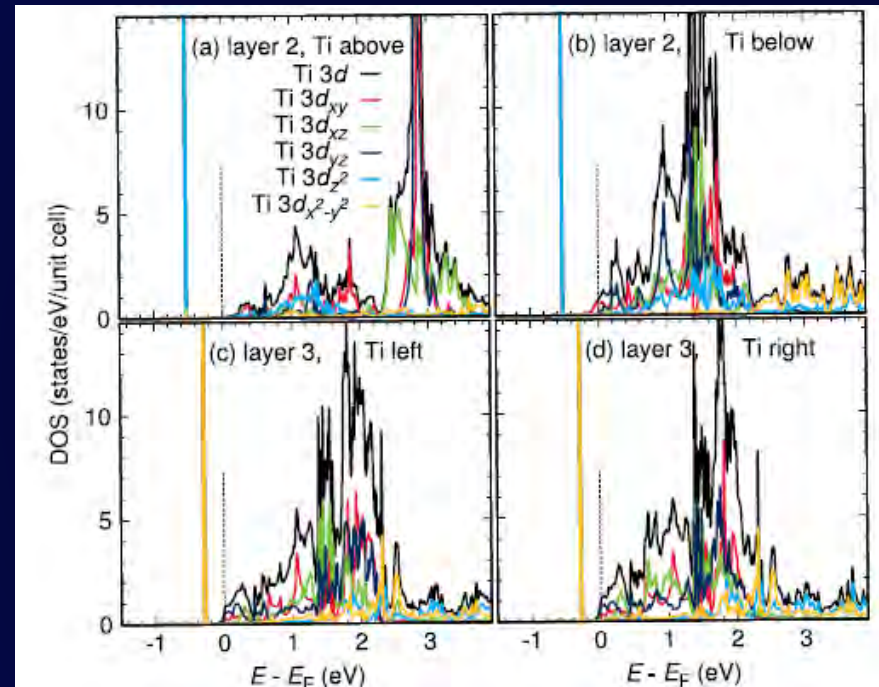
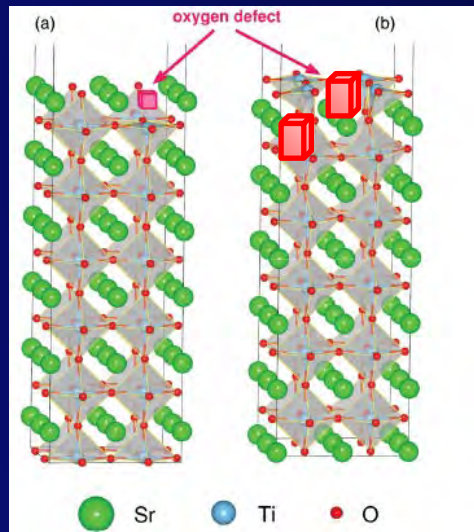
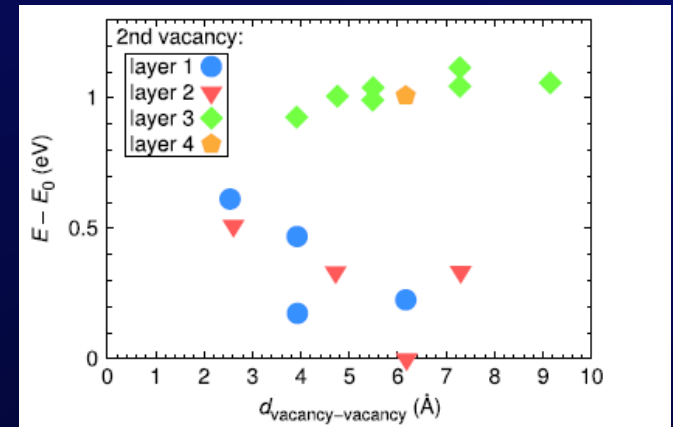
2DEG but no in-gap state



Oxygen vacancies at the boundary could be the unifying mechanism responsible for the presence of the 2DEG

divacancy

- TiO₂ termination
- 1 vacancy in top layer
- + 1 nearby

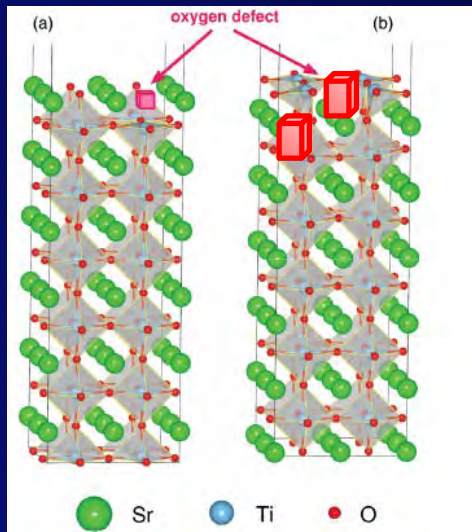


Oxygen vacancies at the boundary could be the unifying mechanism responsible for the presence of the 2DEG

divacancy

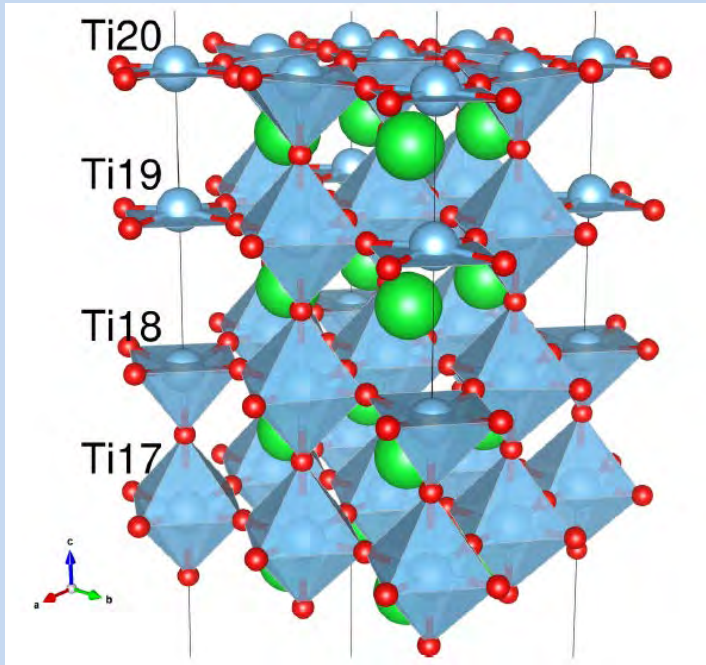
- TiO₂ termination
- 1 vacancy in top layer
- +1 nearby

Qualitative agreement with non spin-resolved ARPES spectrum

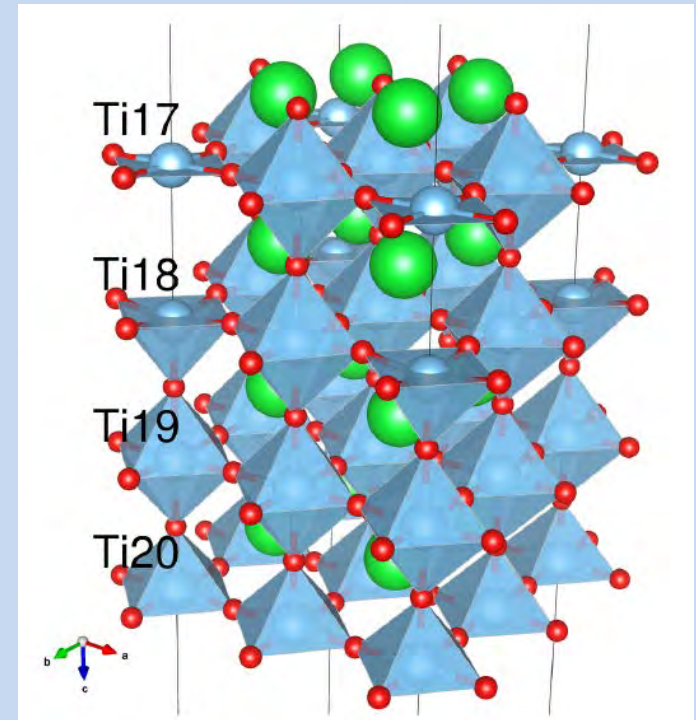


- ☐ Can we recover 1.3 eV peak?
- ☐ Where do the extra states come from?
- ☐ How about the spin sector?

Vertical di-vacancy



TiO₂ terminated



SrO terminated

Experimentally, termination is a mixture of both



Sr



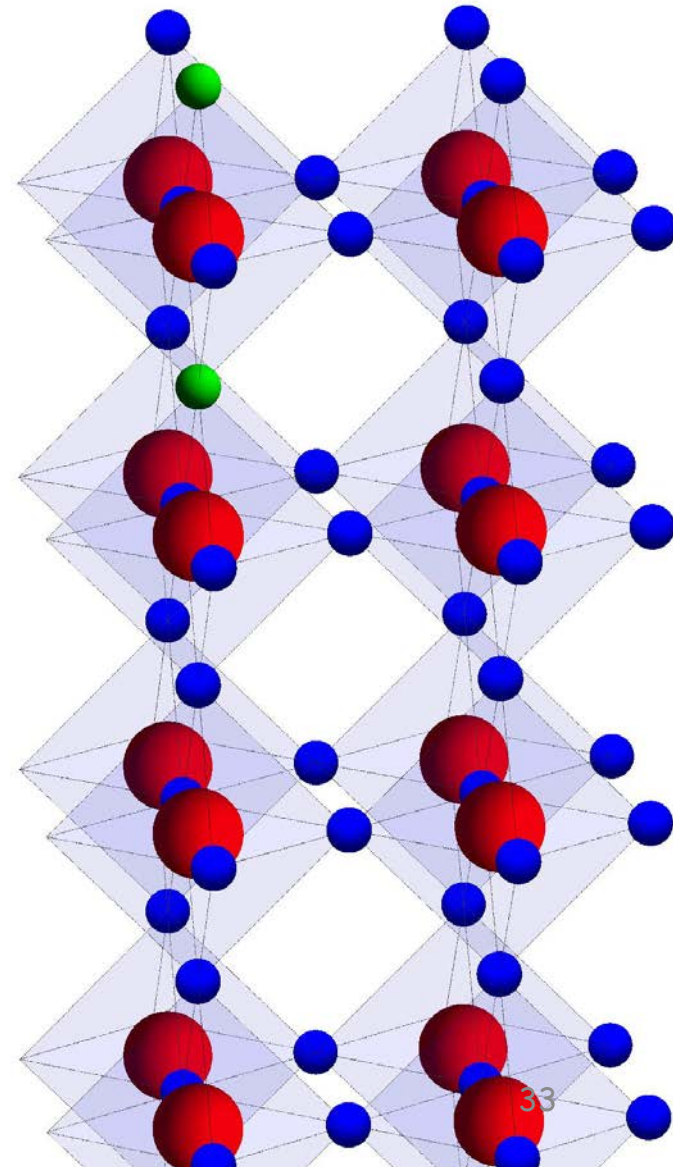
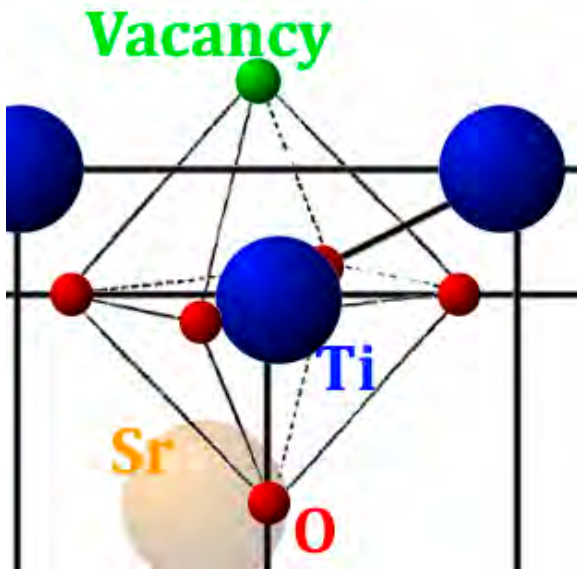
Ti



O

(001) – 2x2x4 cluster

- Include: Rotations, **Oxygen-vacancies**
- Tight-binding model
- XY - periodic



Remove Vacancy

Remove d - p and
 p - p

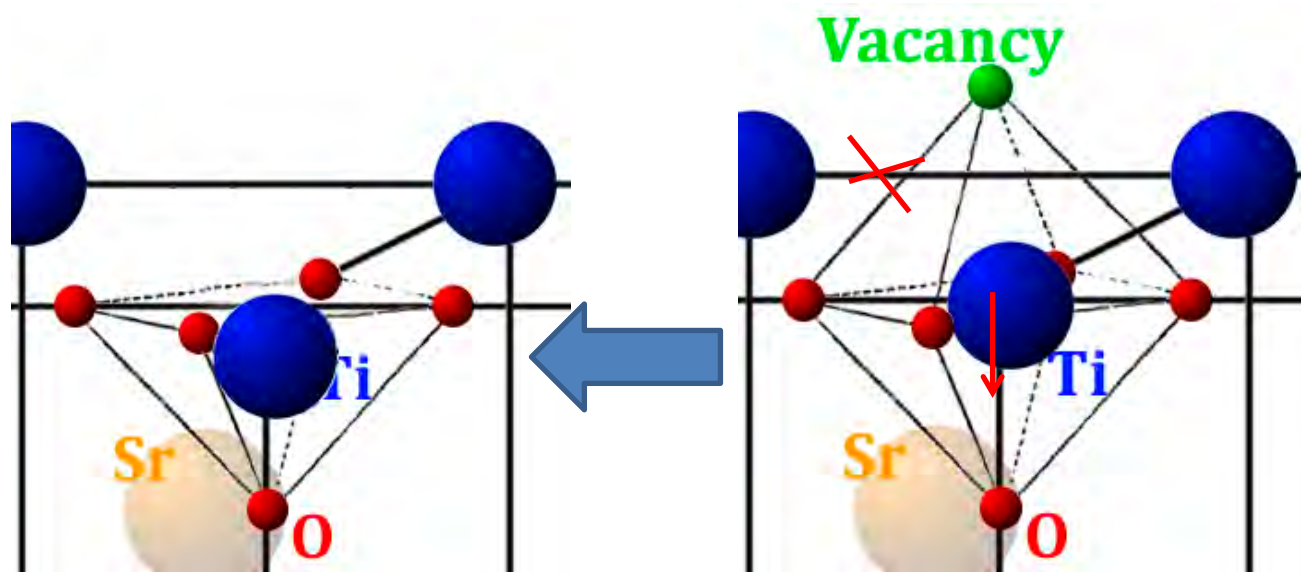
Displace Ti

Increase E_p

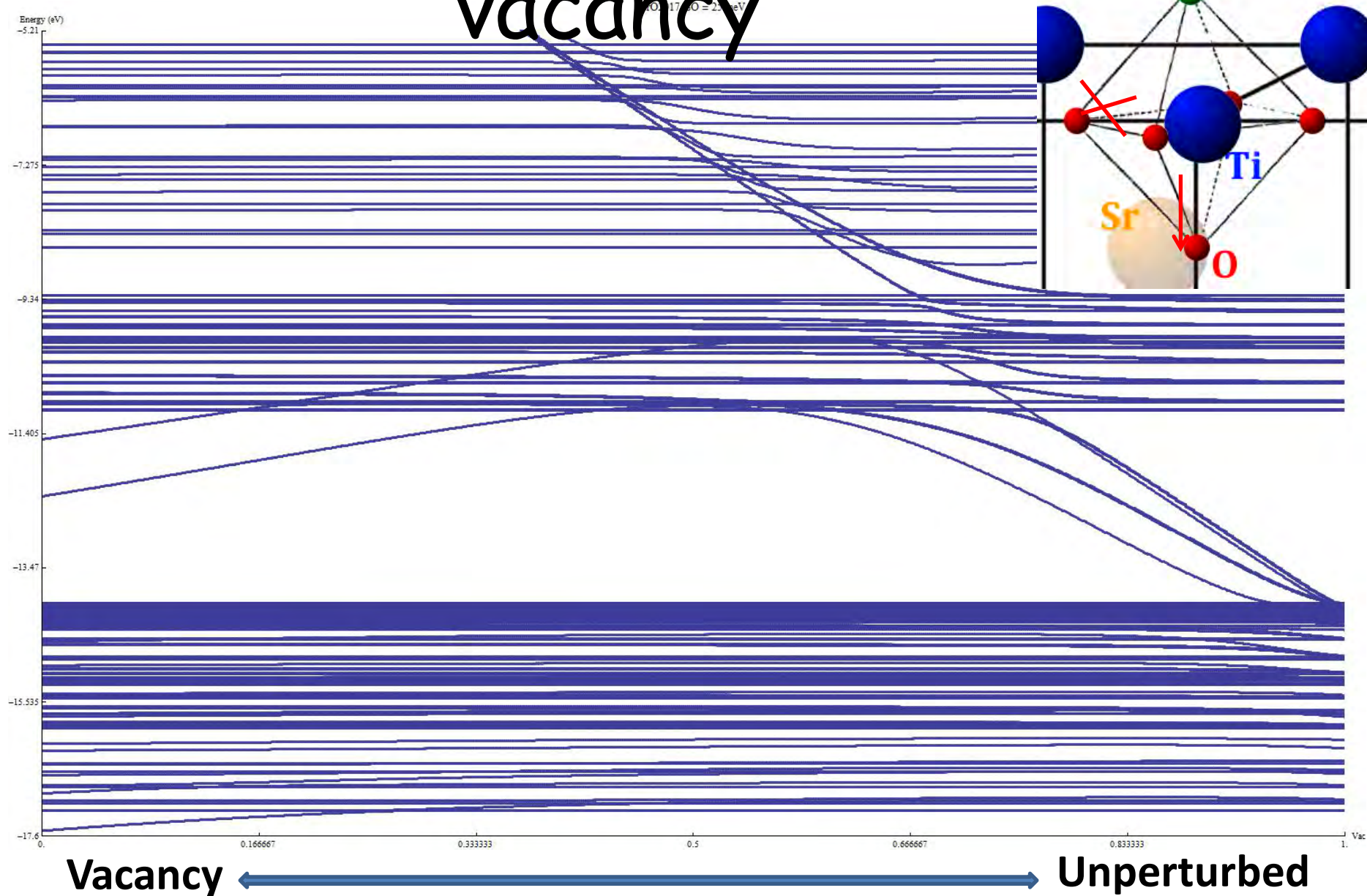
Avoid having e
sitting on site of
vacancy

Change d_{z^2} - d_{xz} -
 d_{yz} energies

Some d orbitals
'see' the effect of
the new Ti local
neighbours



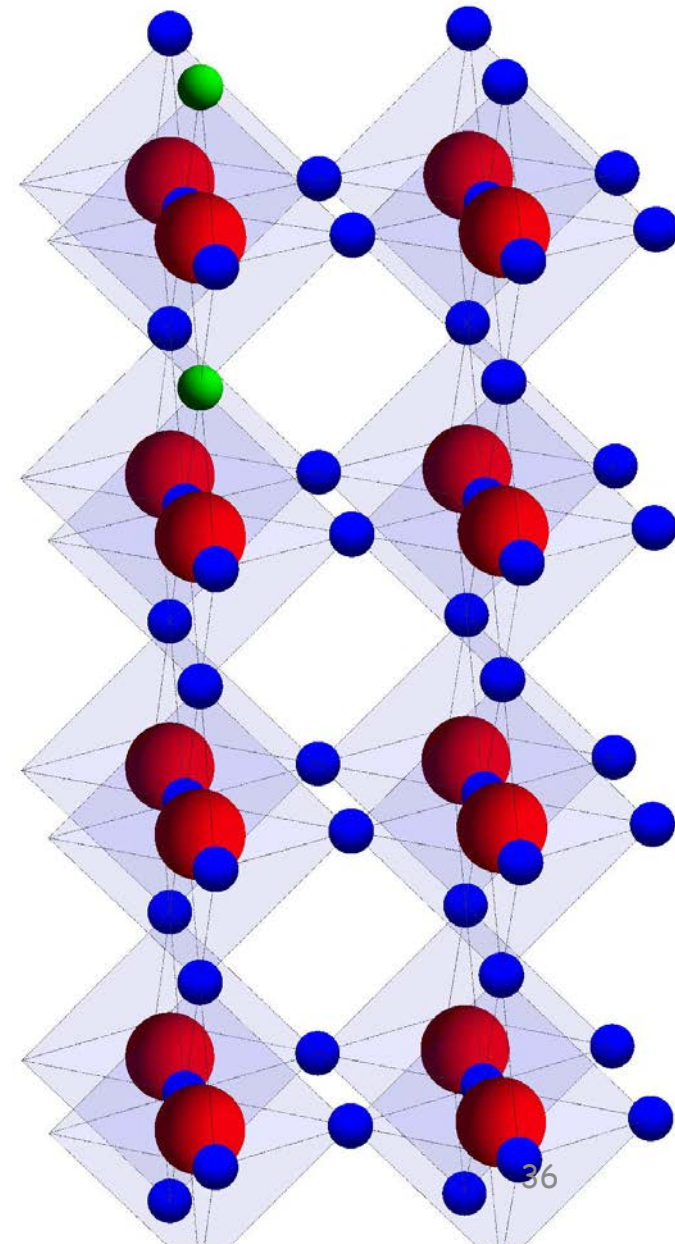
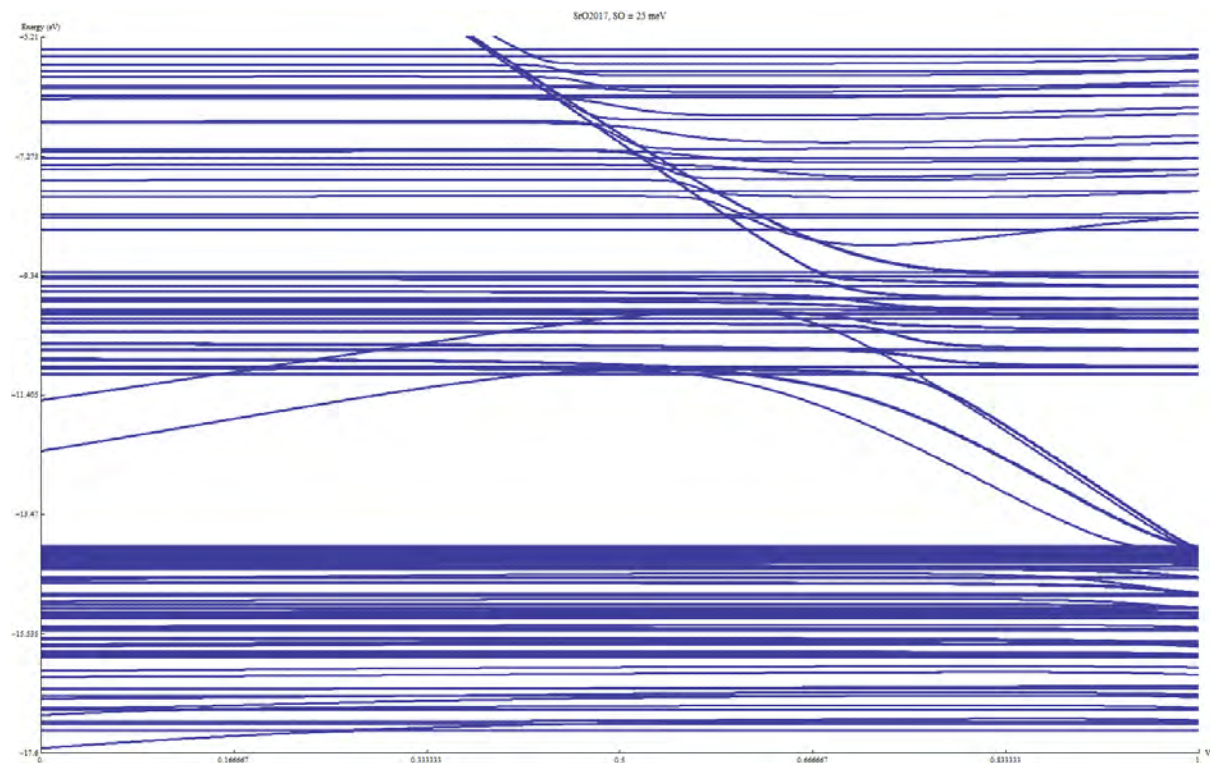
"Smoothly" removing the vacancy



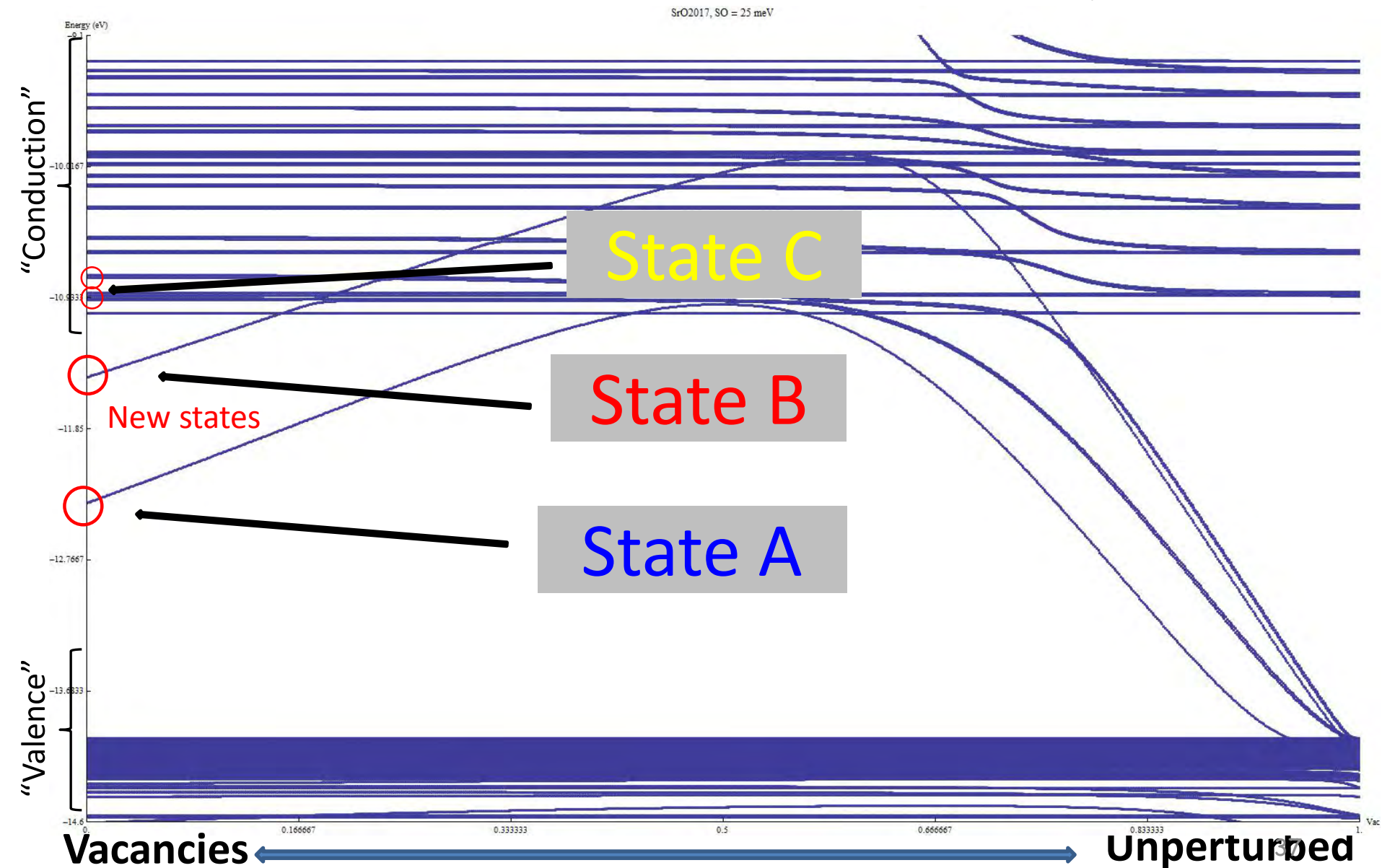
Results for 2 vacancies vertically aligned

SrO-terminated

$$\lambda = 25 \text{ meV}$$

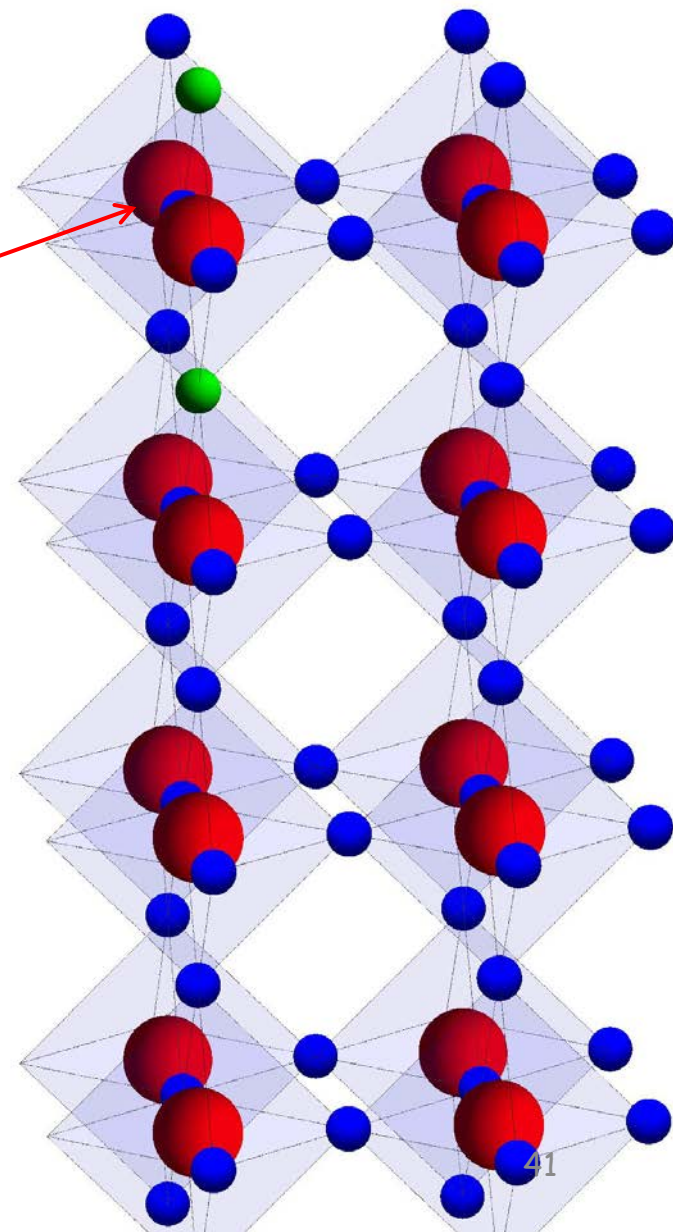
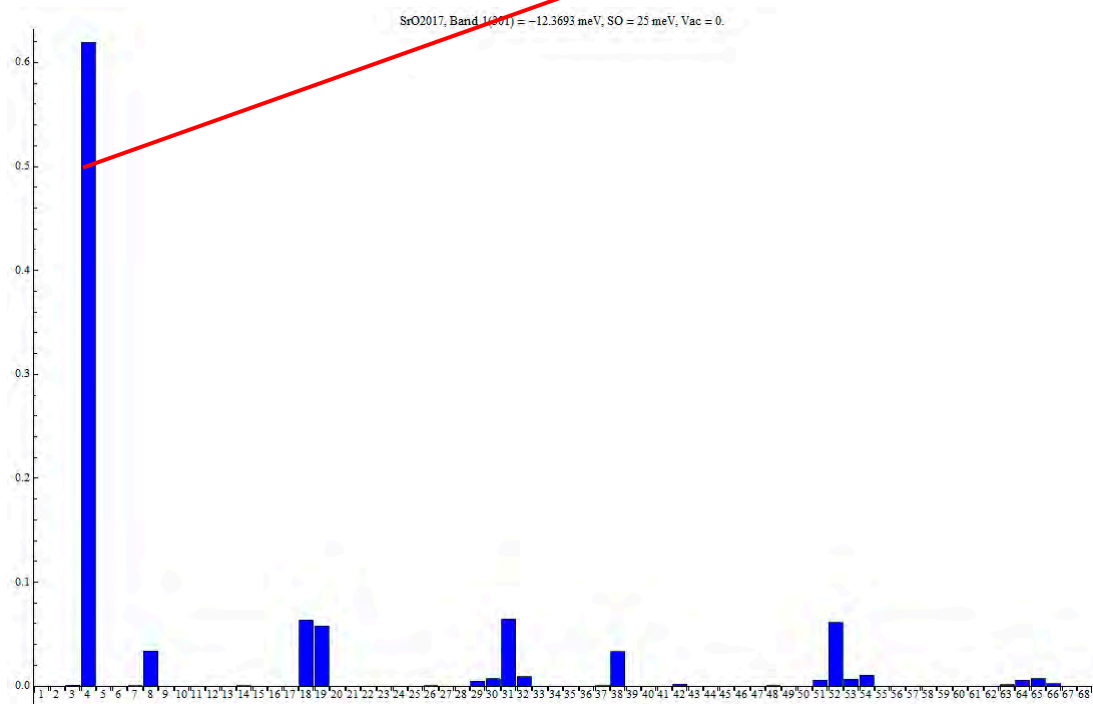
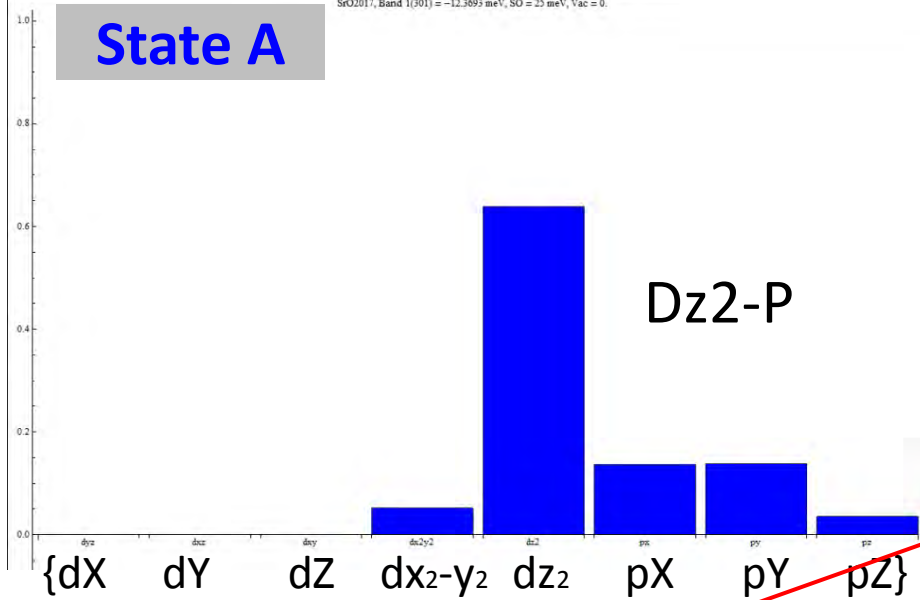


"Smoothly" removing the vacancies, significant change in some energy states



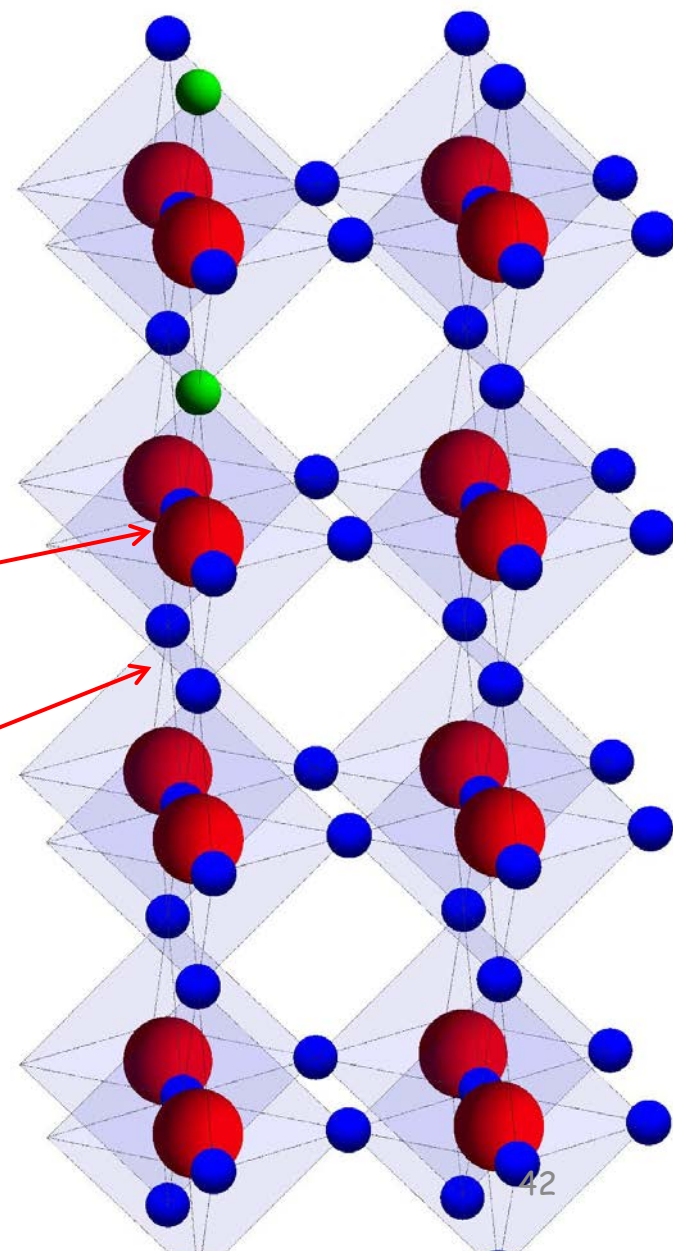
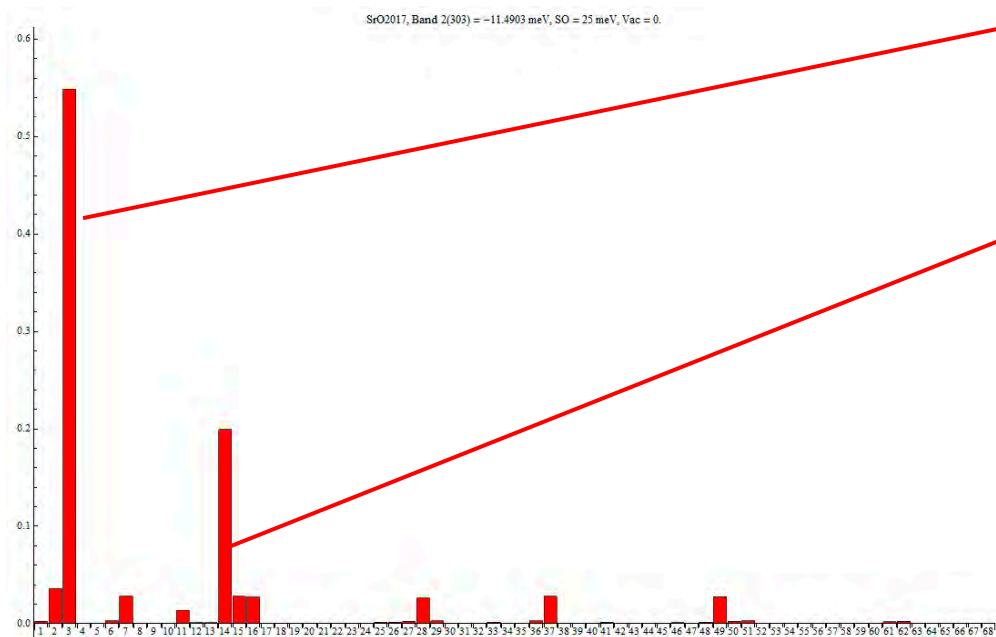
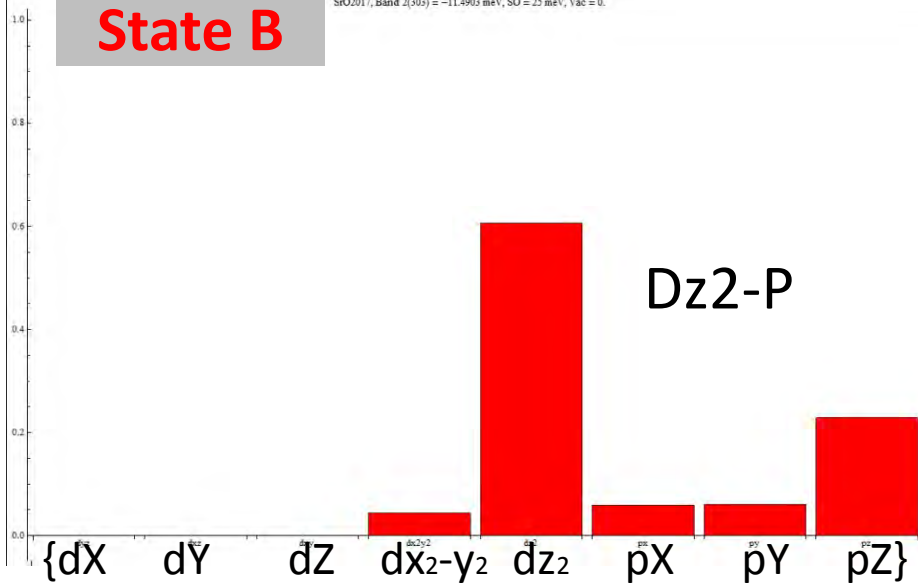
State A

SrO2017, Band 1(301) = -12.3693 meV, SO = 25 meV, Vac = 0.



State B

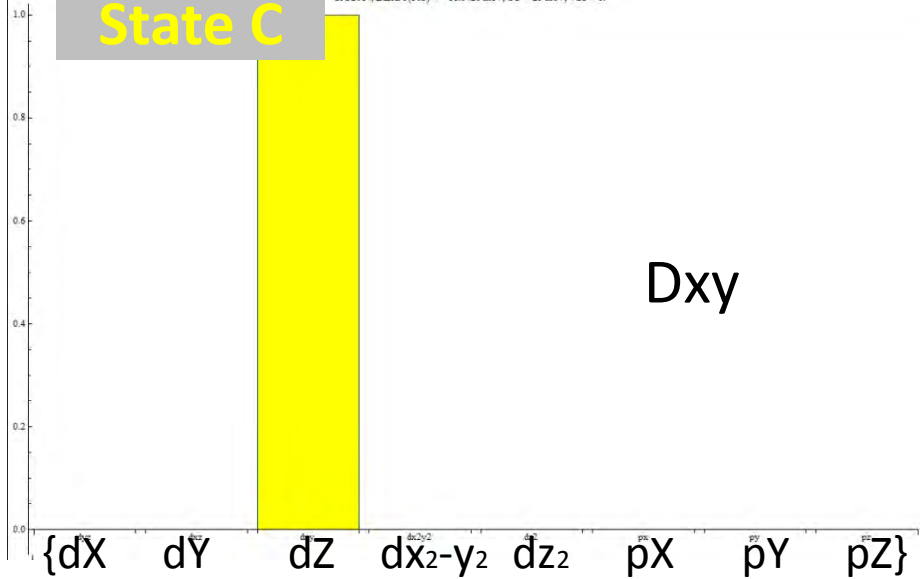
SrO2017, Band 2(303) = -11.4903 meV, SO = 25 meV, Vac = 0.



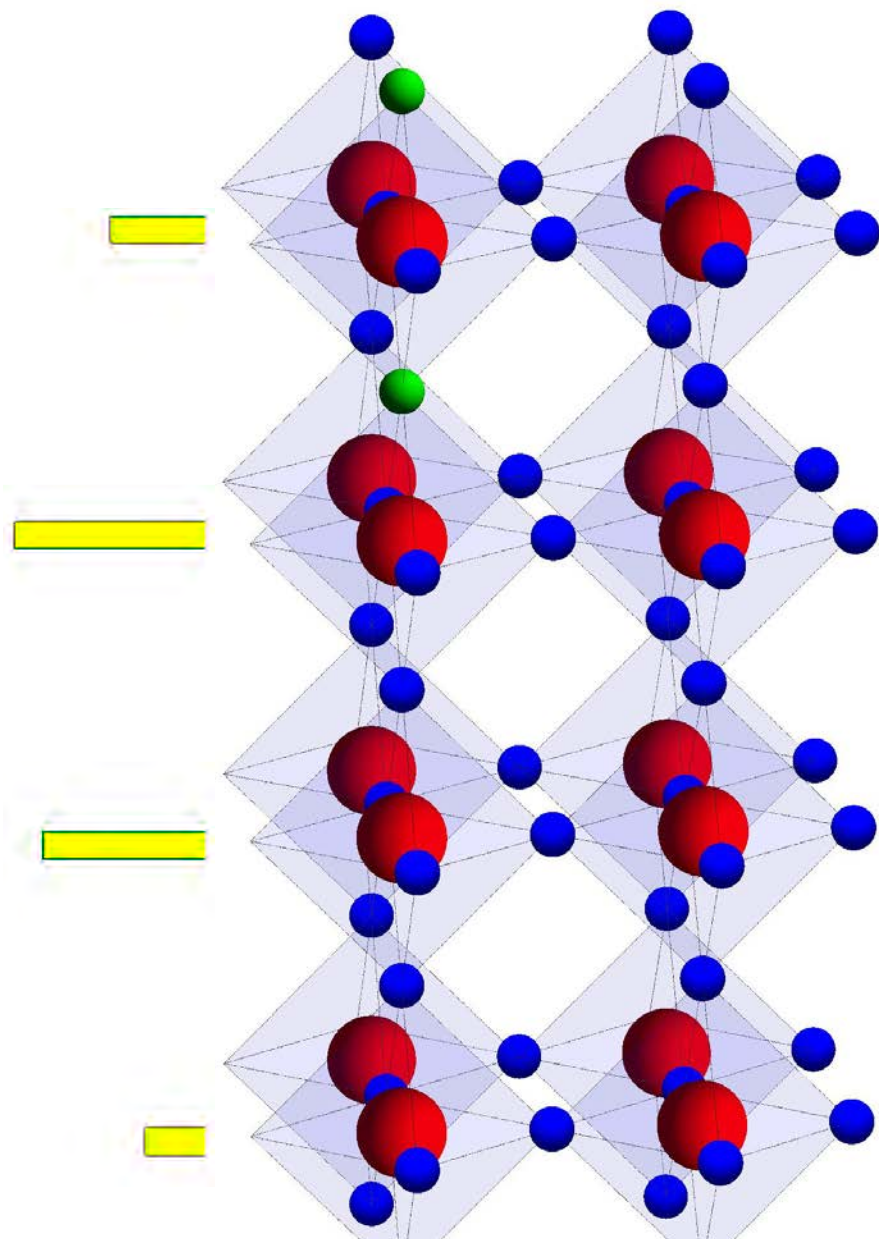
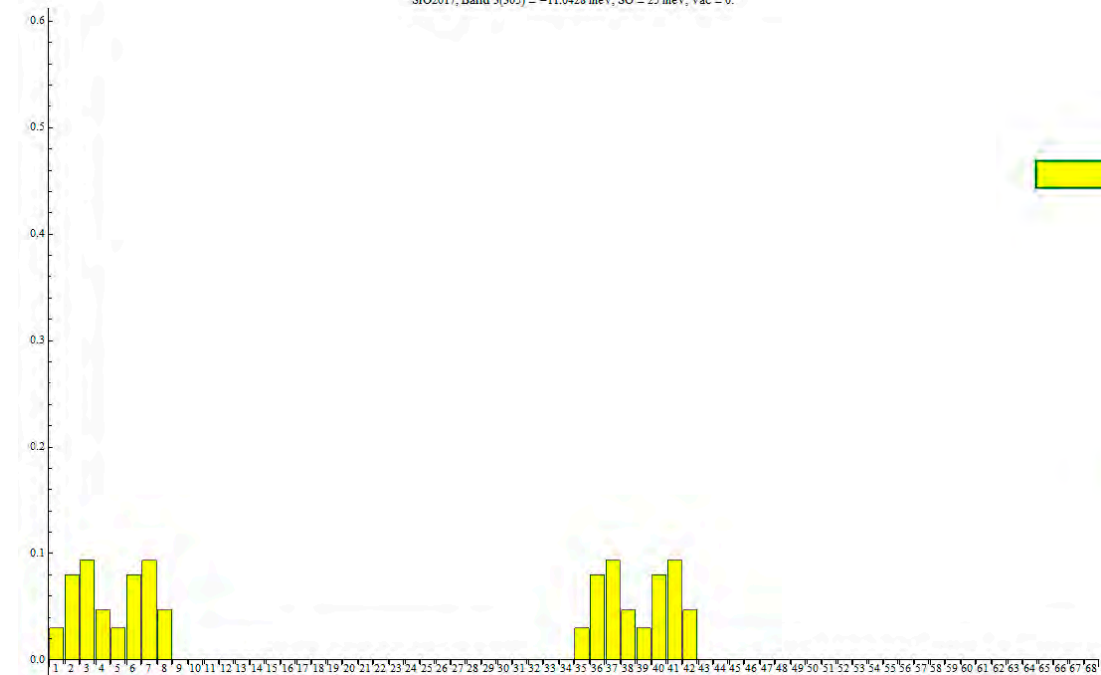
State C

SrO2017, Band 3(305) = -11.0428 meV, SO = 25 meV, Vac = 0.

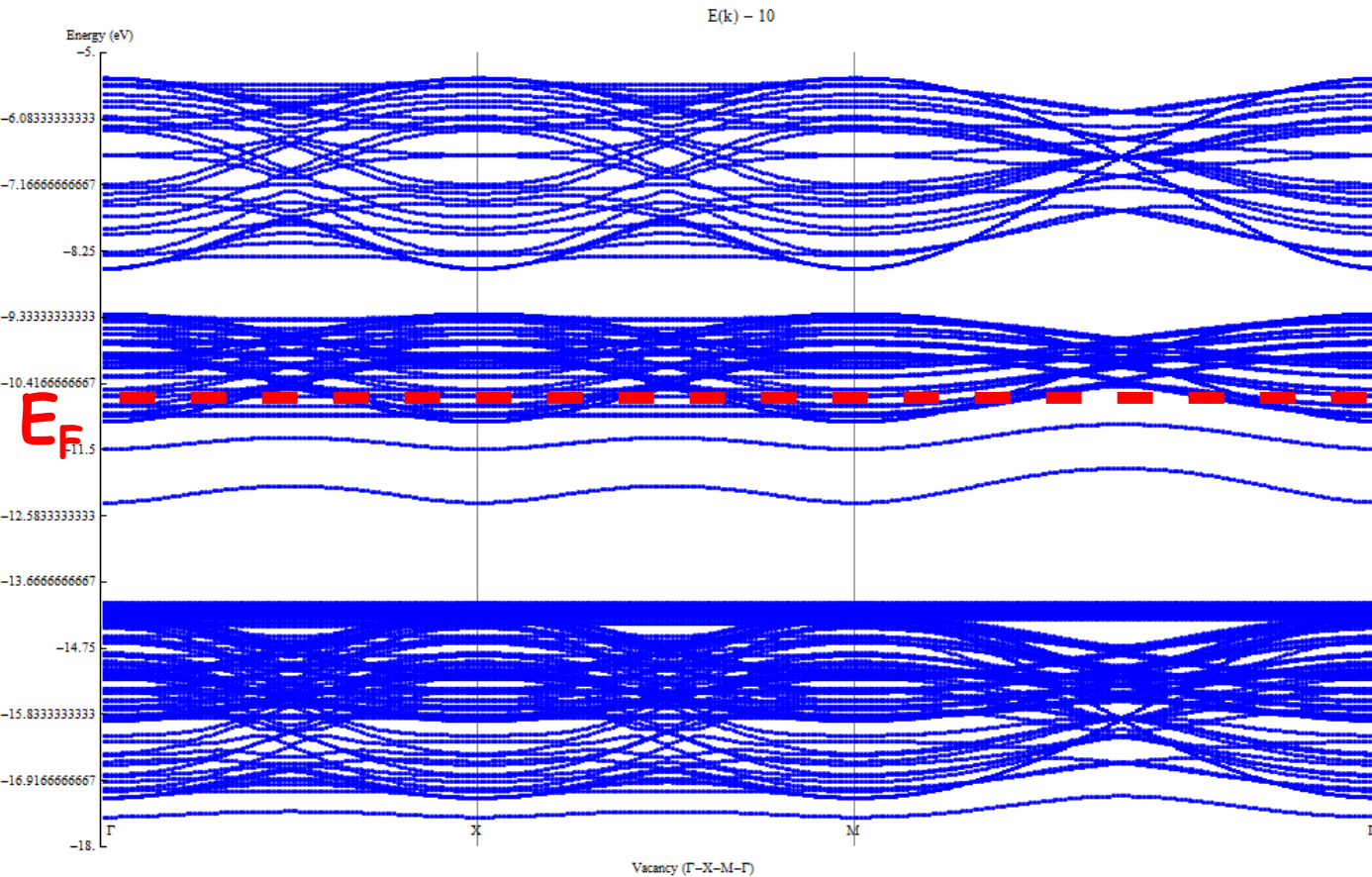
Dxy



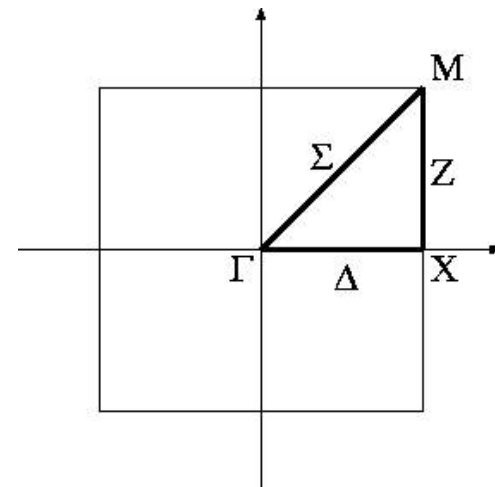
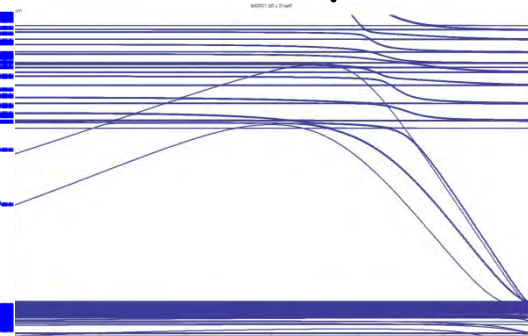
SrO2017, Band 3(305) = -11.0428 meV, SO = 25 meV, Vac = 0.



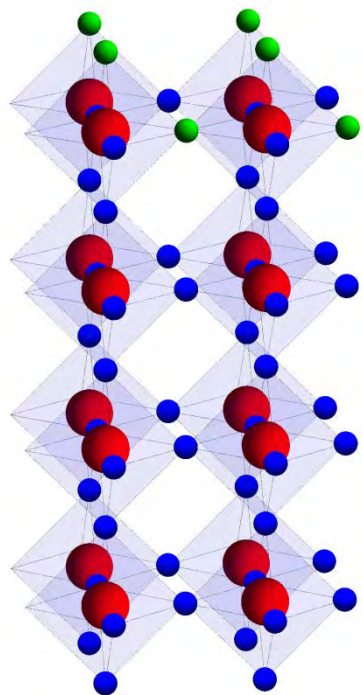
- K-dependence?



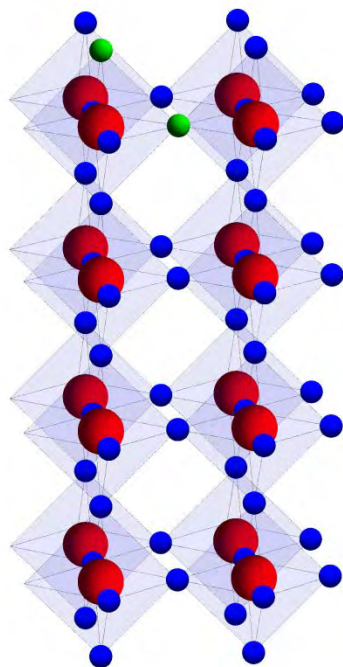
Calculated at
the Γ point



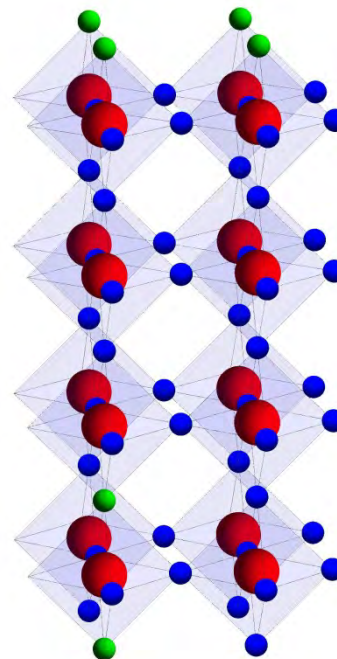
Other configurations



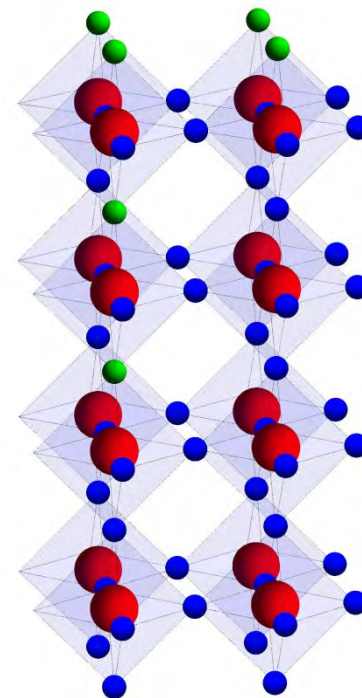
TiO_2 terminated;
horizontal di-vacancy



SrO terminated;



SrO terminated;
Vertical di-vacancy



TiO_2 terminated;
vertical di-vacancy

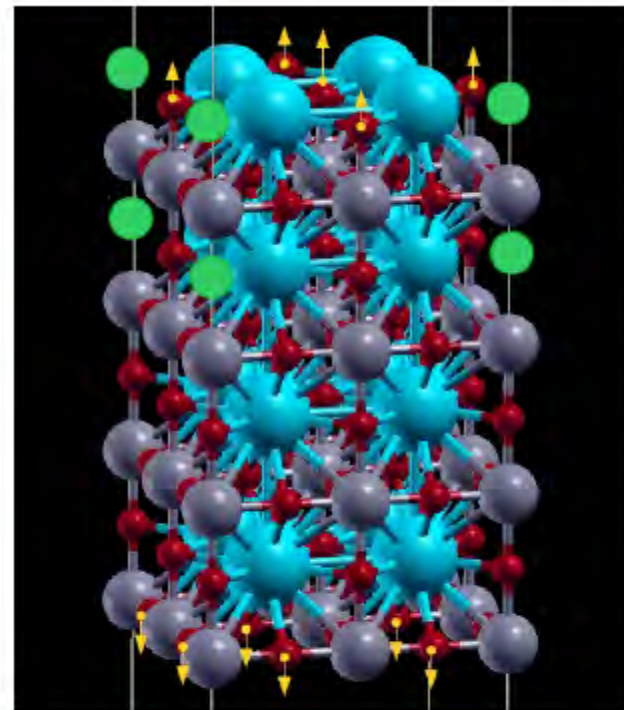
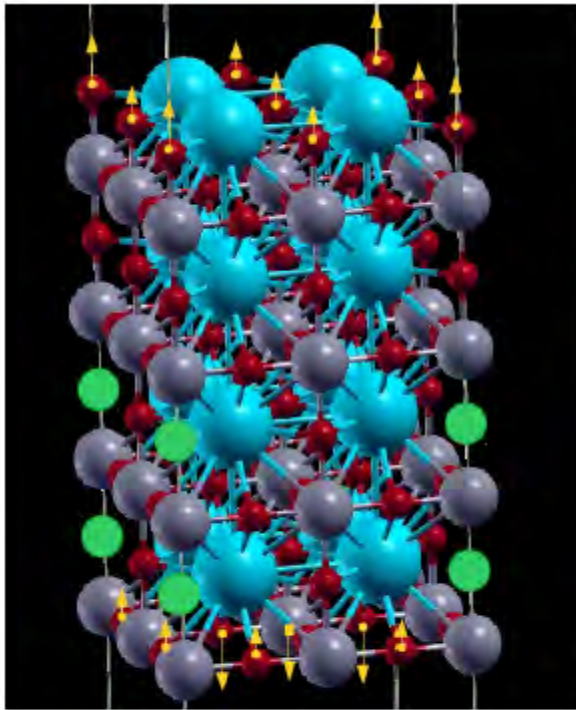


Figure 2 : Schematic view of the relaxed crystal structure. The yellow arrows represent the relative displacement of the ions w.r.t. to the Sr ones (top layer) and the corner Ti ones (bottom layer)

TiO₂ terminated

SrO terminated



Ti



O

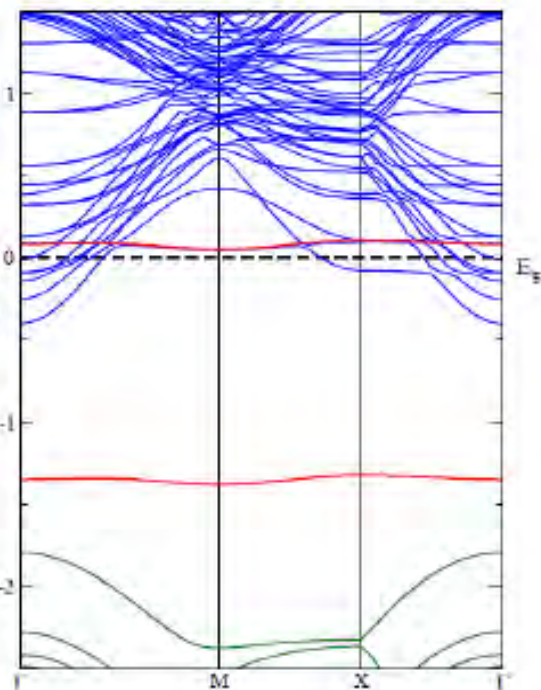


Sr

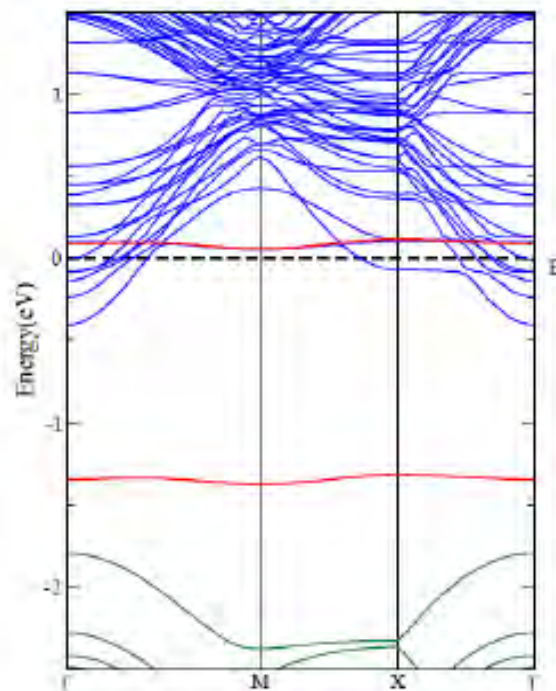


O-Vac

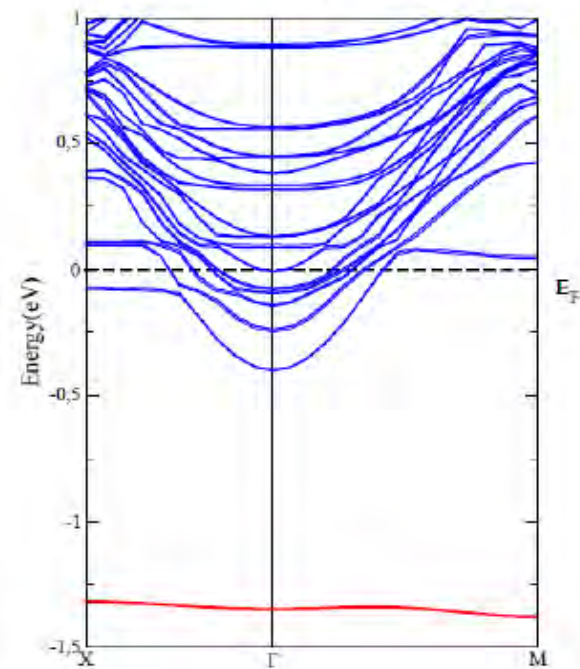
Case "TiO top" relaxed (down spin)



Case "TiO top" relaxed (up spin)

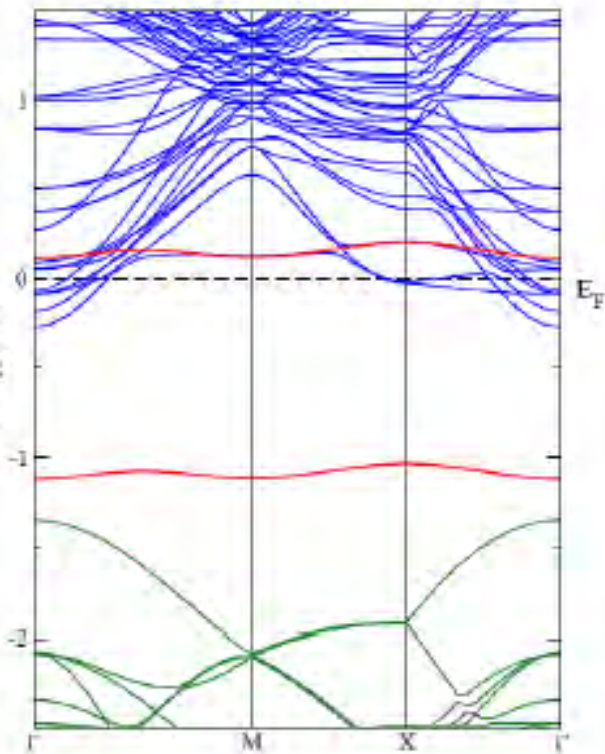


Case "TiO top" with SO

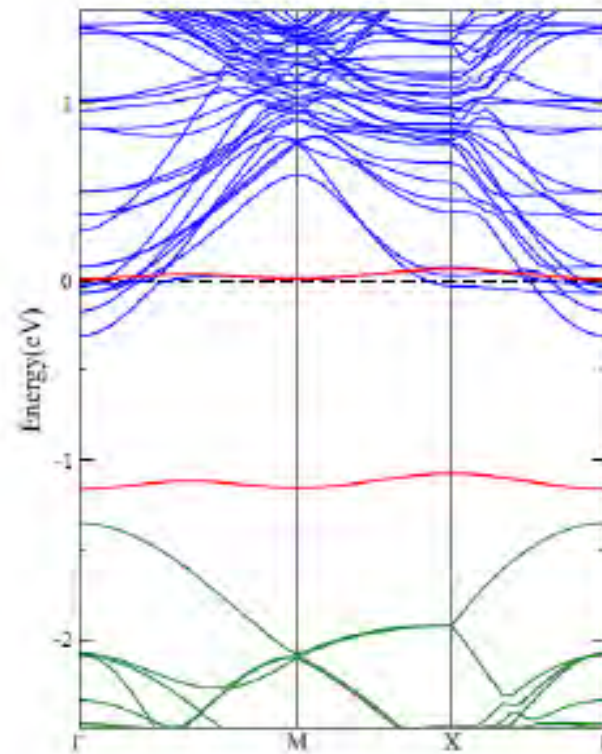


Magnetic Kohn-Sham band structure for a relaxed 2x2x4 SrTiO₃ slab with an oxygen divacancy case « TiO-top ».

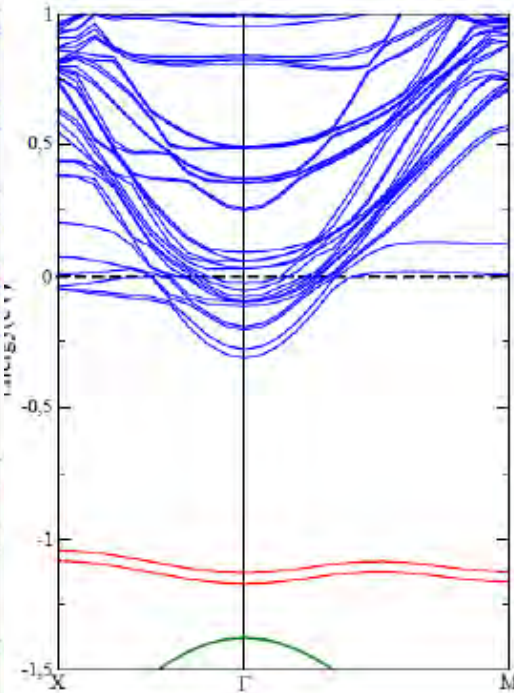
Case "SrO top" relaxed (down spin)



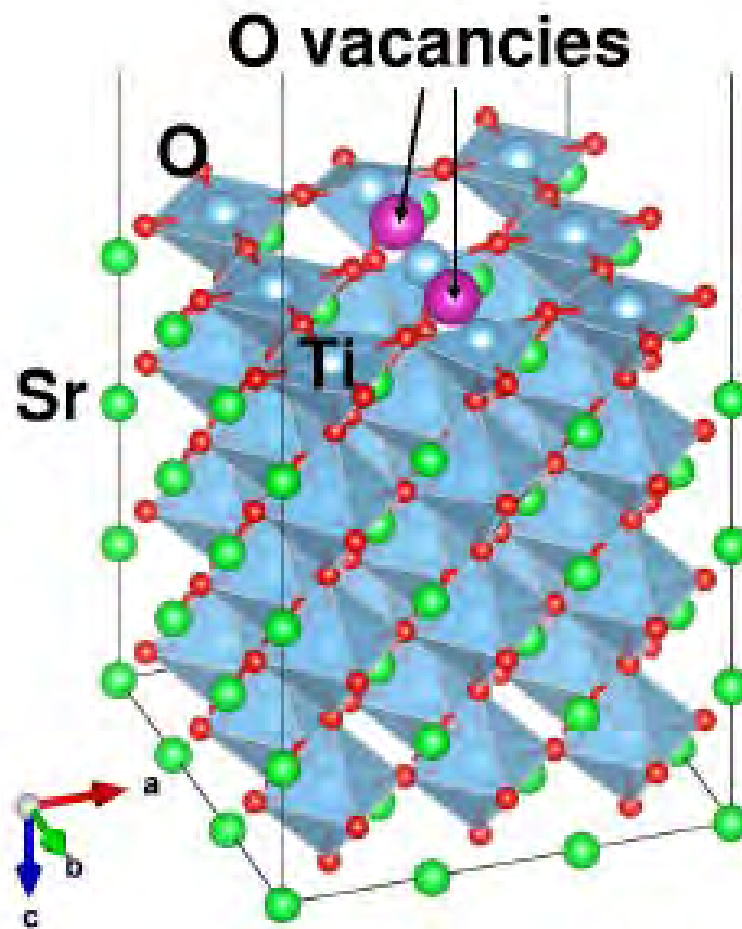
Case "SrO top" relaxed (up spin)



Case "SrO top" with SO



Magnetic Kohn-Sham band structure for a relaxed 2x2x4 SrTiO₃ slab with an oxygen divacancy case « SrO-top ».



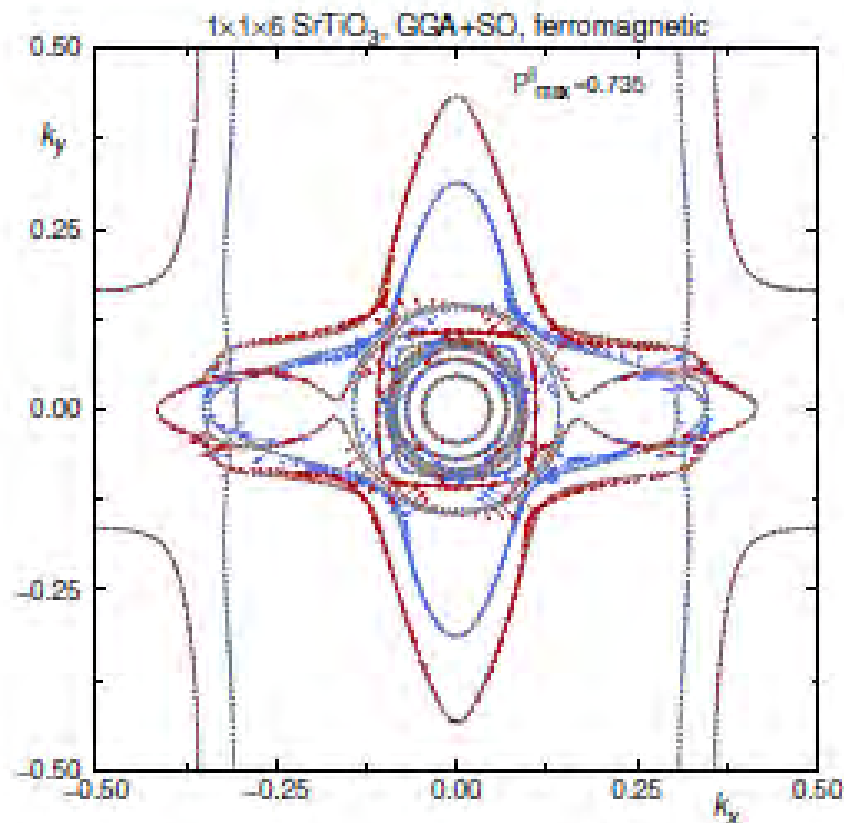
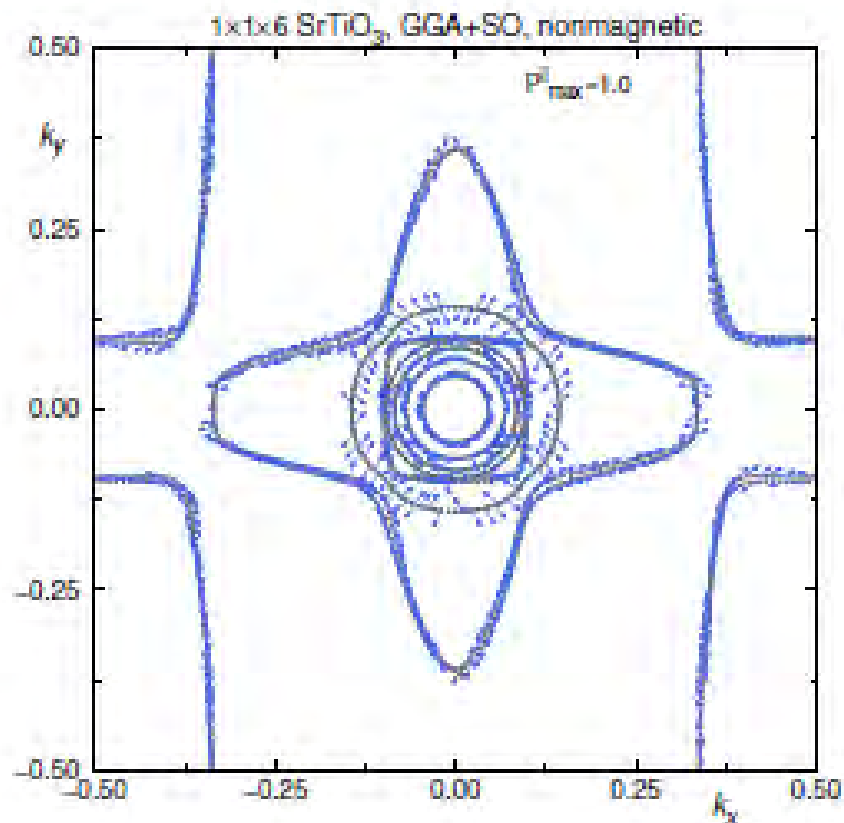
GGA+SOC nonmagnetic

GGA+SOC ferromagnetic with quantization in z/x direction

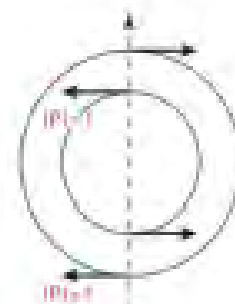
GGA+SOC+U on Ti nonmagnetic

GGA+SOC+U on Ti ferromagnetic with quantization in z/x direction

Relaxed 1x1x6 slab: nonmagnetic $\Leftrightarrow$ ferromagnetic

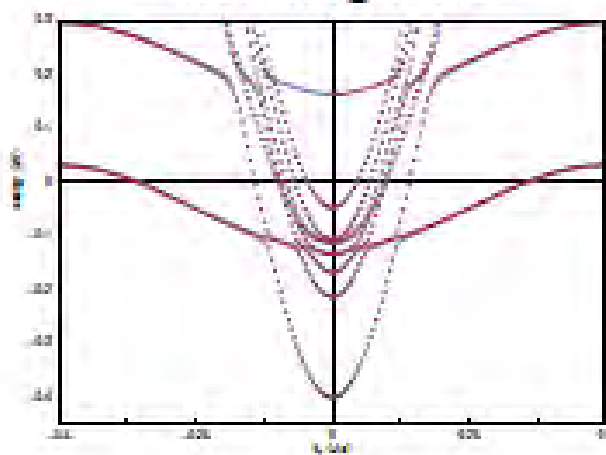


- $\Rightarrow$ Magnetism enhances spin splitting
- $\Rightarrow$ Ferromagnetism is energetically favorable

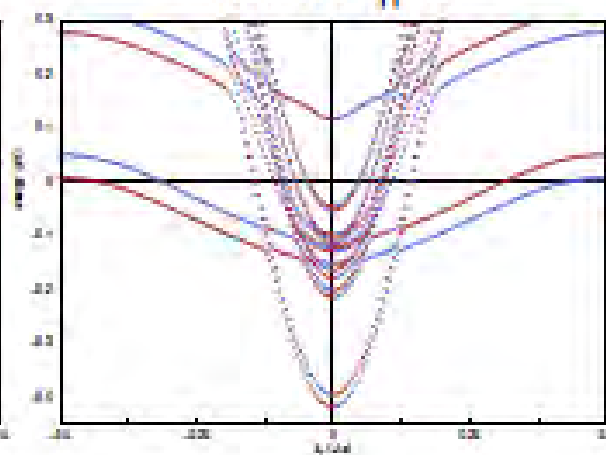


Summary GGA+SO+Magnetism: Experiment $\Leftrightarrow$ Theory

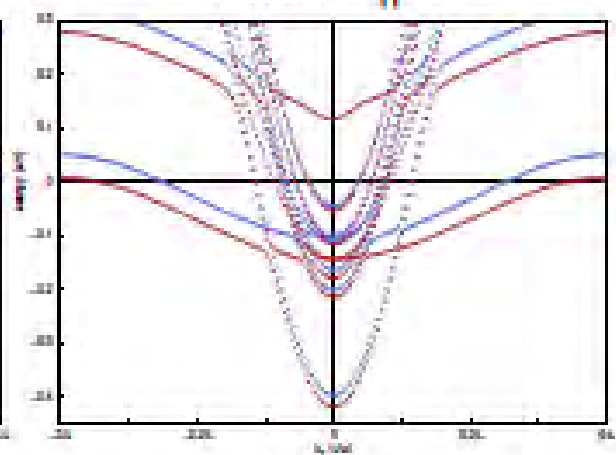
nonmagnetic



fm $\vec{m} \parallel \hat{z}$



fm $\vec{m} \parallel \hat{x}$



Experiment	In-plane polarization	Spin flip	Energy splitting
GGA+SO nm	✓	✓	✗
GGA+SO $\vec{m} \parallel \hat{z}$	✗	✓	✓
GGA+SO $\vec{m} \parallel \hat{x}$	✓	✗	✓

- ❑ Surface of STO harbors puddles of either magnetic or non magnetic regions
- ❑ Surface of STO is inhomogeneous : mixture of TiO_2 and SrO terminations.
- ❑ Orbital selectiveness (t_{2g} states lead to spin textures; e_g states to magnetism)
- ❑ Fermi surface with mixed orbital character

