

Glassy dynamics

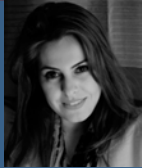
in geometrically frustrated Coulomb liquids

Louk Rademaker, KITP, Santa Barbara
Bad Metals Conference, 30 June 2015, Mainz

Thanks to...



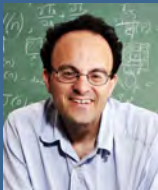
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Introduction into glass

Slow dynamics

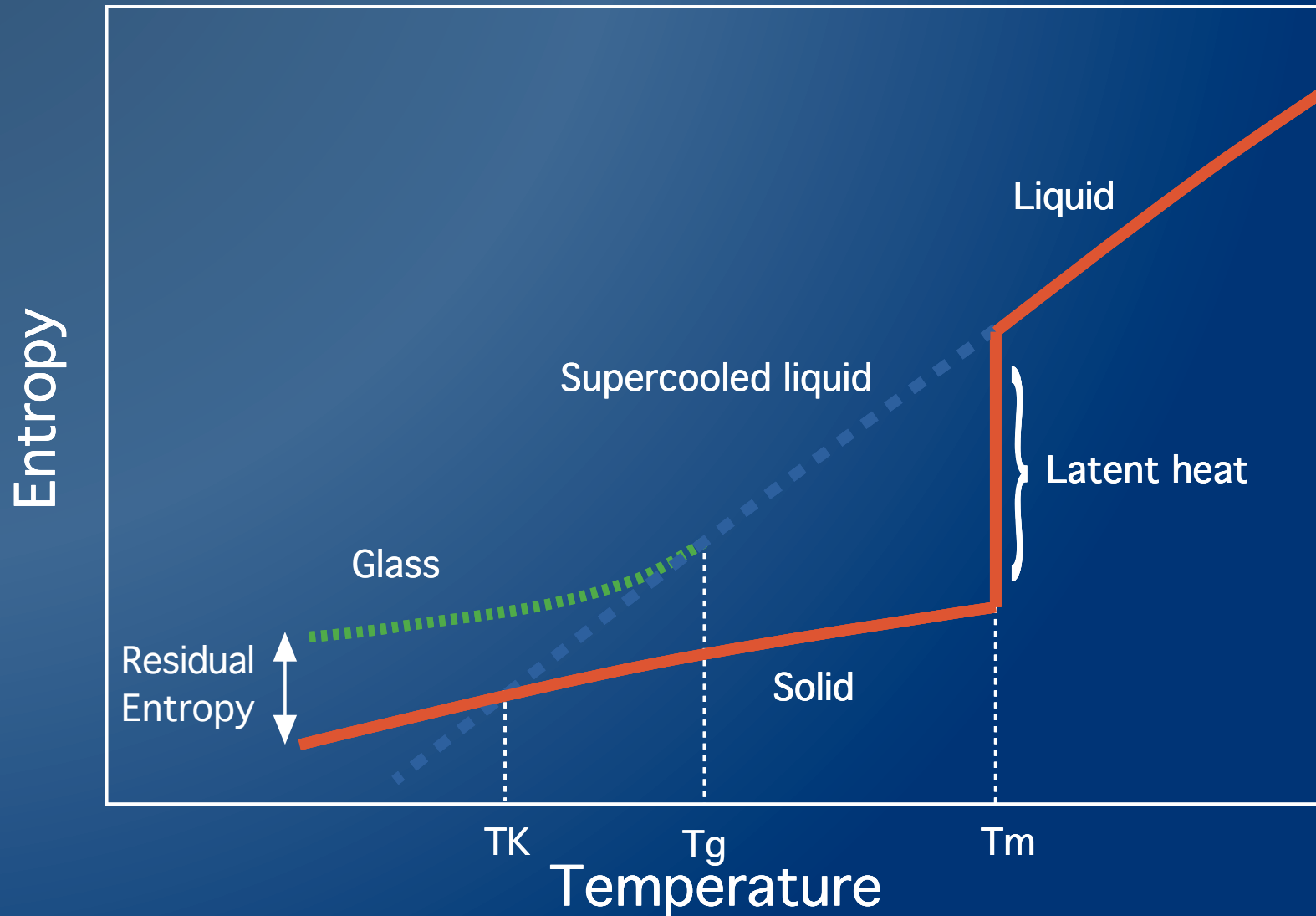


Wild energy landscape

Soft excitation gap

Short-range correlations

Supercooling a liquid



Electronic Glasses

Free Fermi gas, Landau Fermi liquid

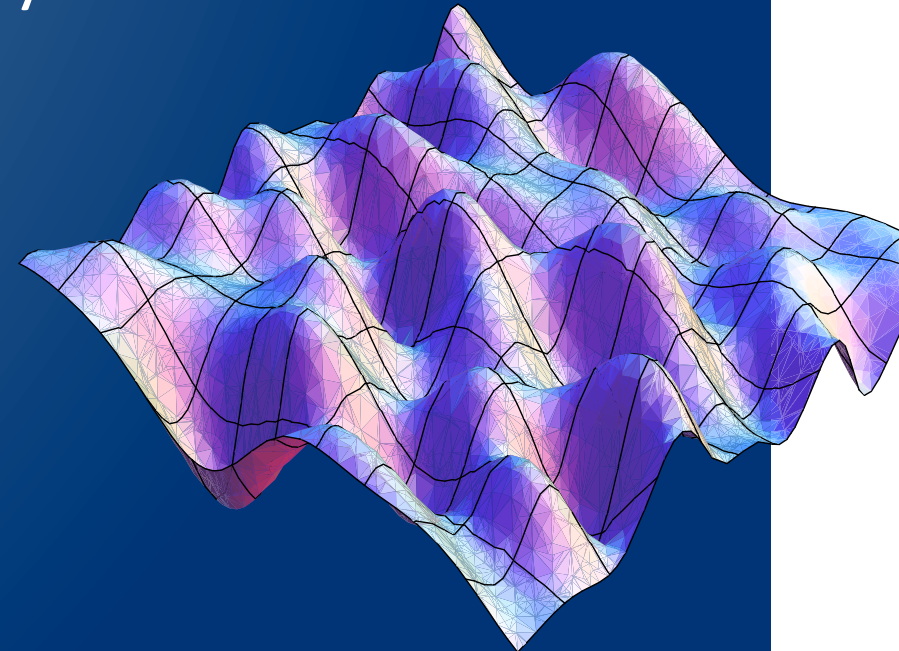
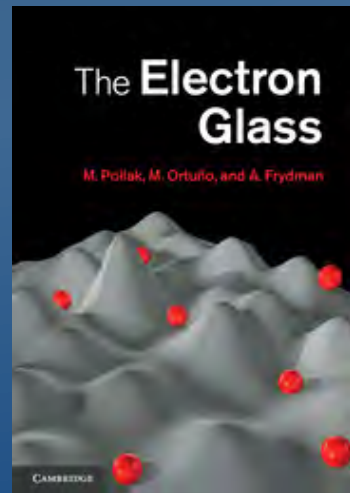
Solid: Wigner crystal, charge-density-wave

Electron glass: Intrinsic **Disorder!**

Anderson Localization

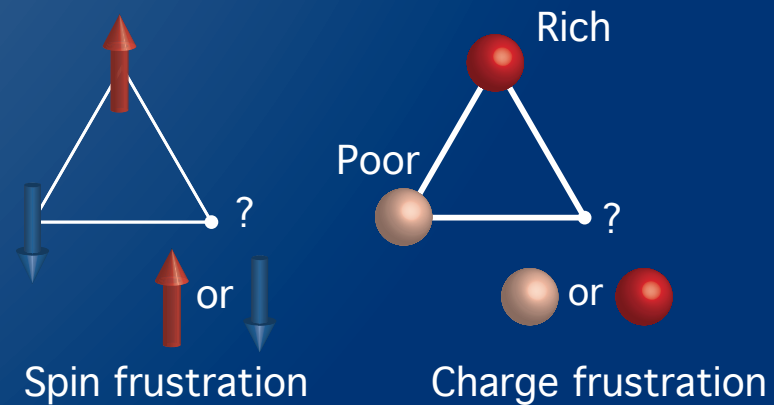
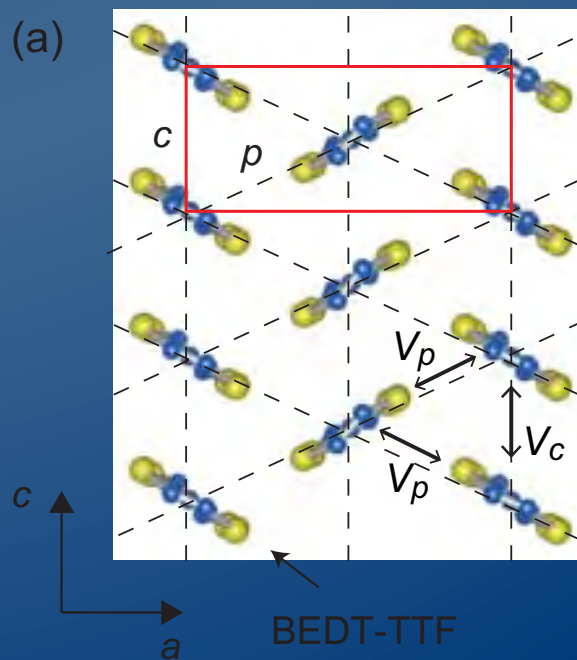
Many Body Localization

Coulomb glass



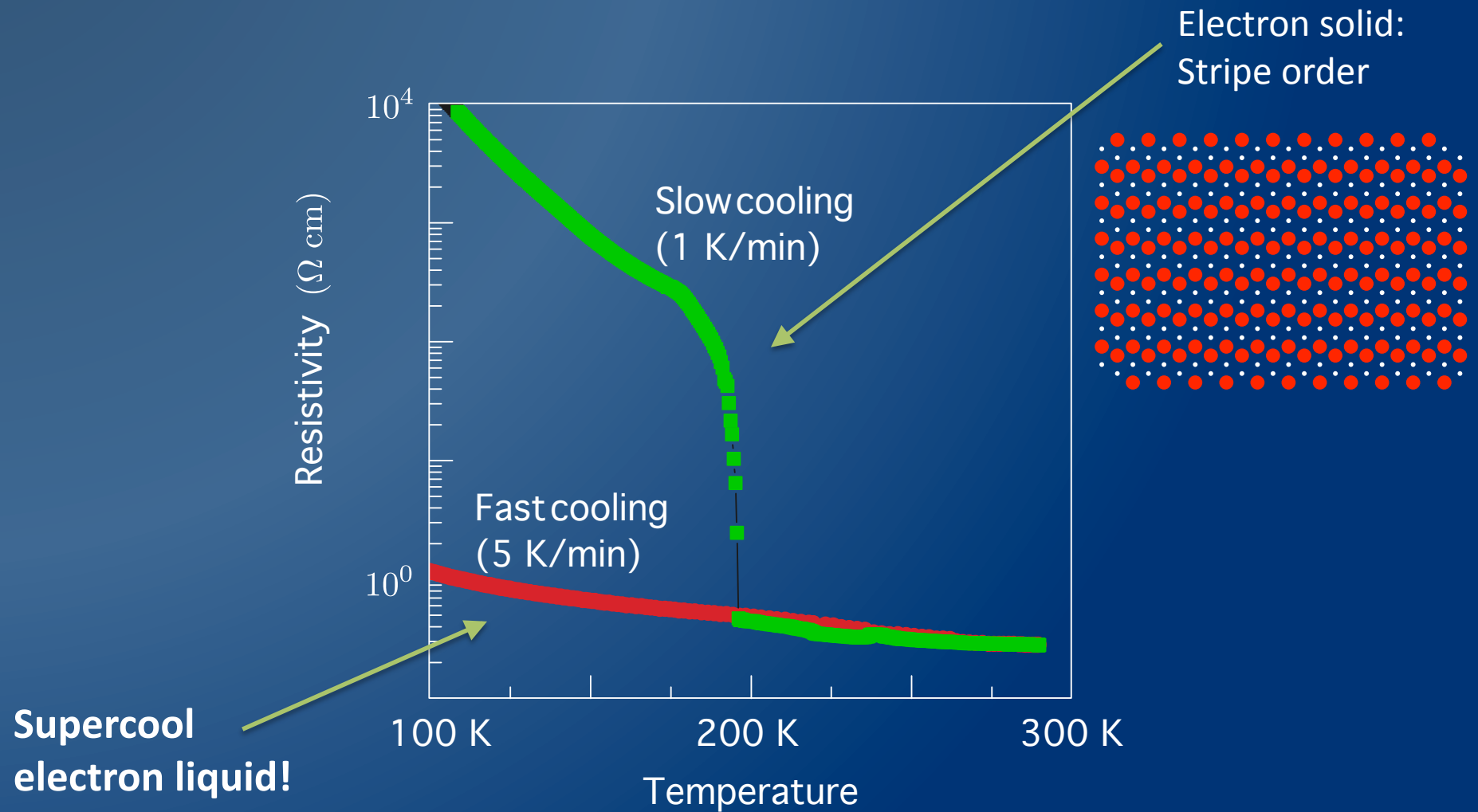
Supercool electron glass?

- Kagawa et al., Nat. Phys. 9, 413 (2013)
(and Sato et al, JPSJ 2014; Sato et al, PRB 2014;)
- Material: Clean θ -(BEDT-TTF)₂RbZn(SCN)₄

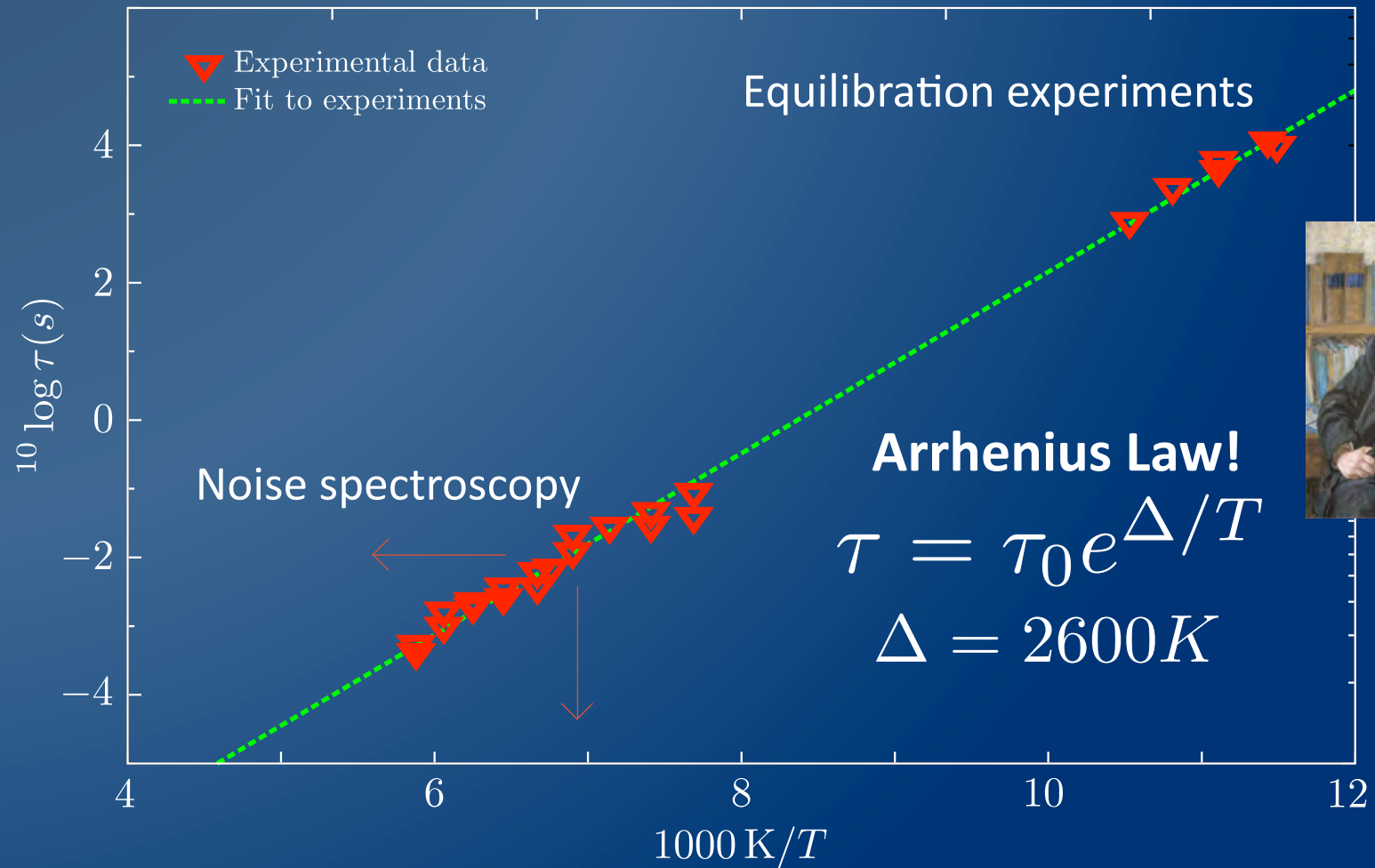


See talk by Simone Fratini, yesterday

Bad metal

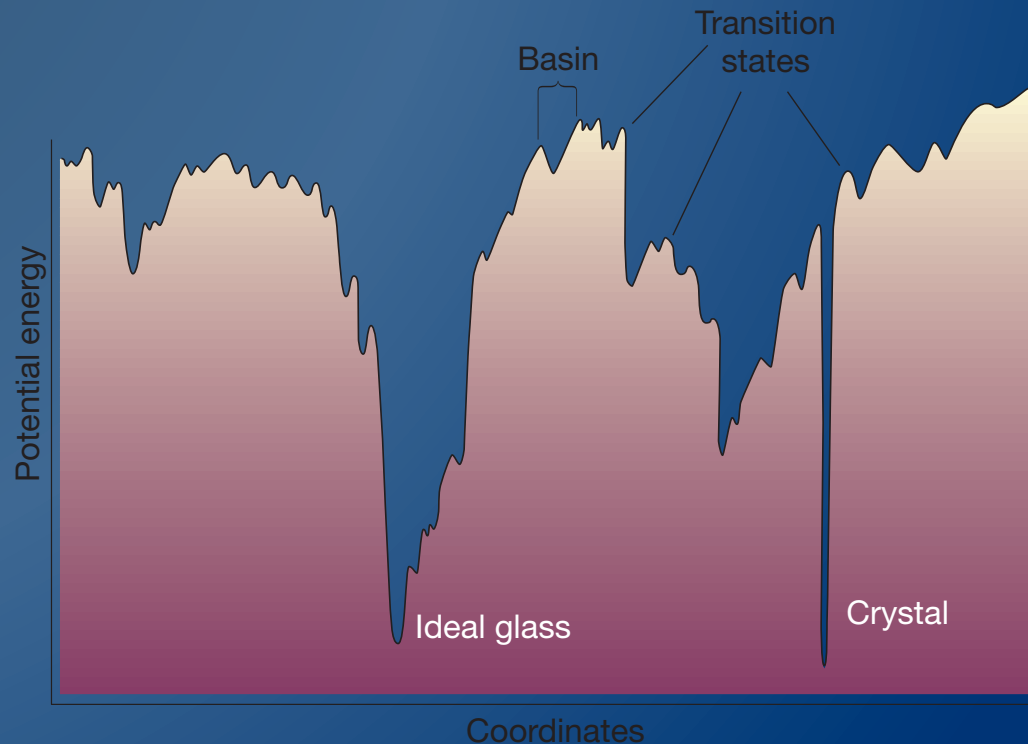


Slow dynamics



1903

Complex energy landscape



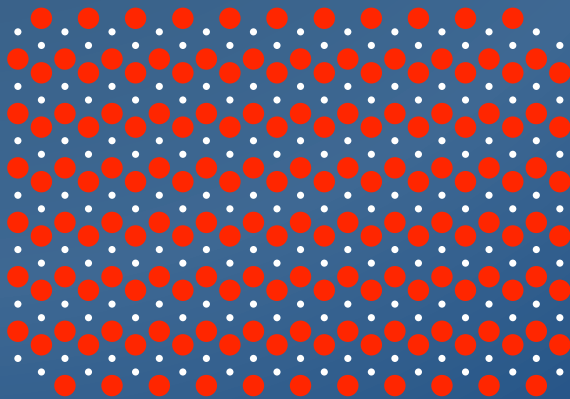
Debenedetti & Stillinger, Nature 2001

But: Electrons can **tunnel!**

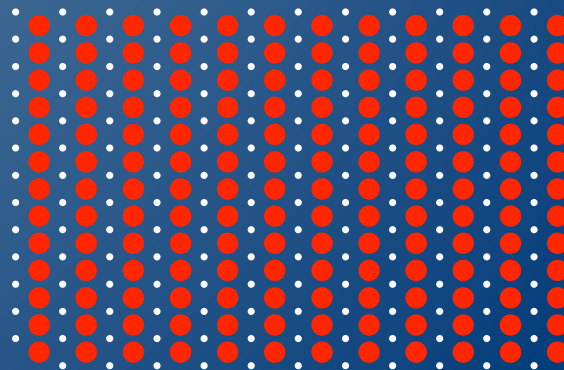
Need **frustration**

Geometric Frustration

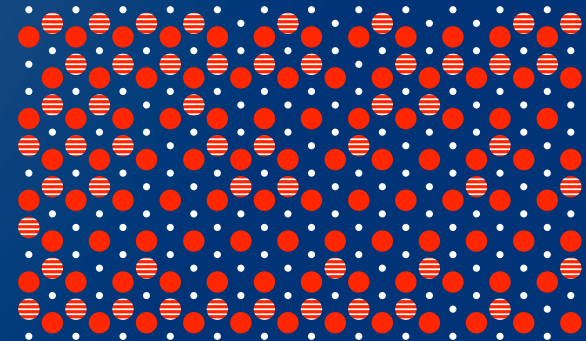
Triangular **Ising** Model $H = \frac{1}{2} \sum_{ij} V_{ij} \left(n_i - \frac{1}{2} \right) \left(n_j - \frac{1}{2} \right)$
Exponentially many Ground States $E = -V/4$



Zigzag Stripes



Straight Stripes



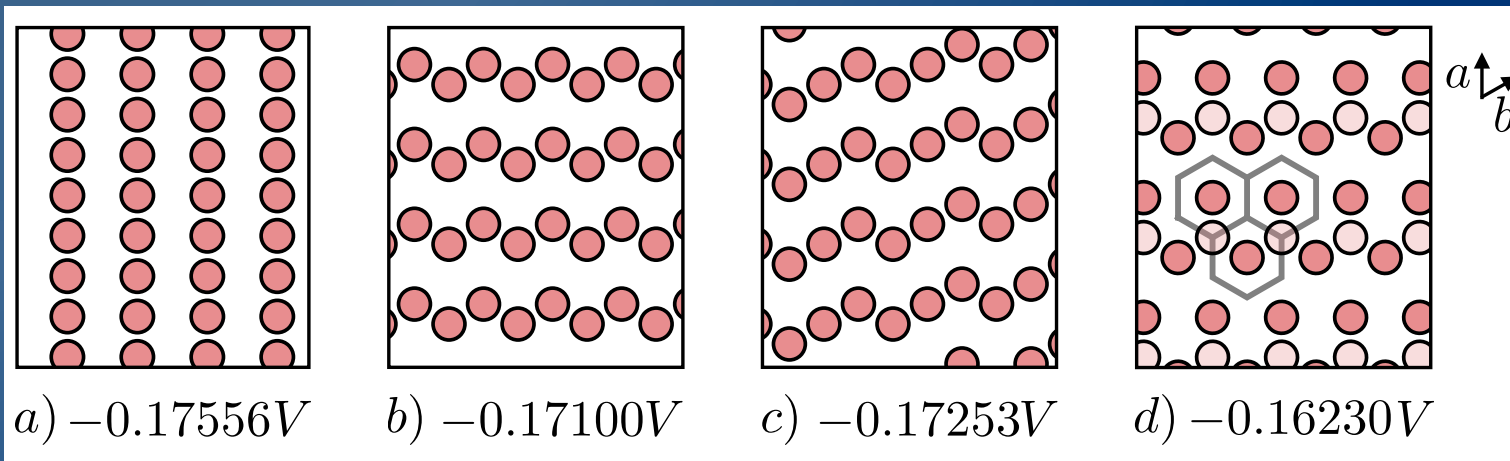
Three sublattice structure

... but not **metastable!**

Long-range interactions

Including **Long-Range Coulomb** interactions lifts degeneracy

$$V_{ij}^{\text{LR}} = \frac{V}{|r_i - r_j|}$$



Macroscopically many states metastable!

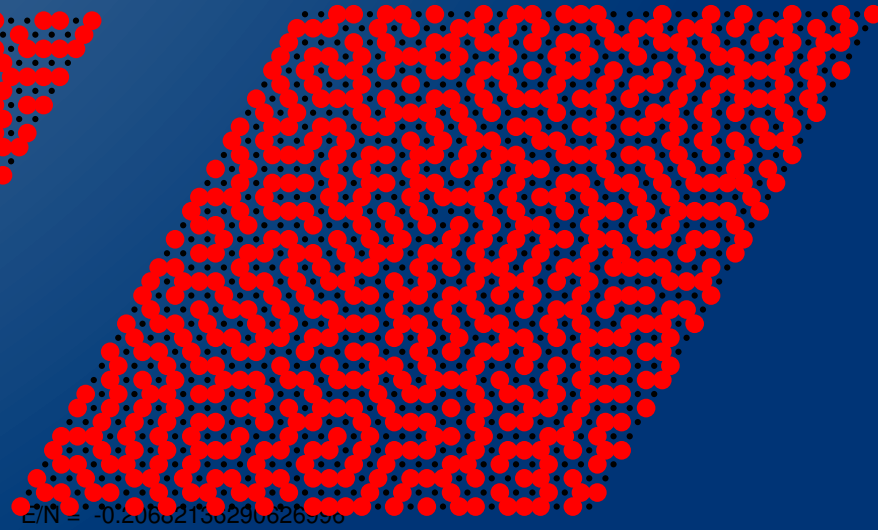
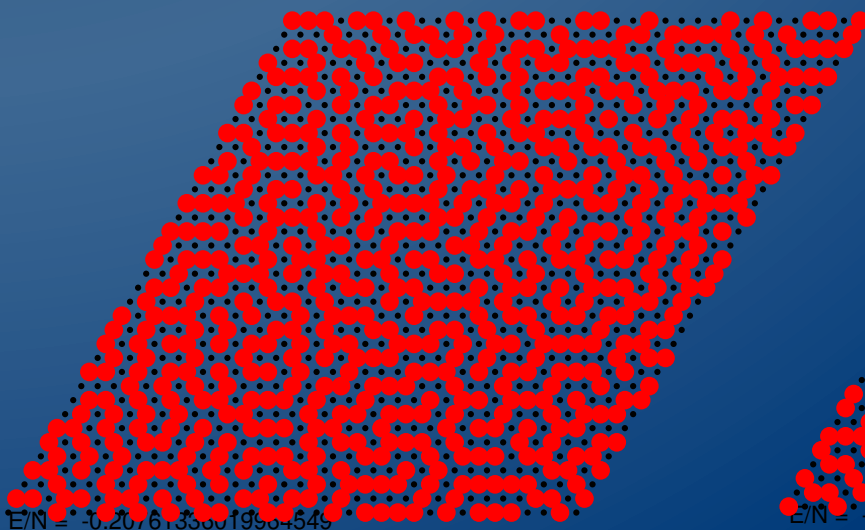
Glassiness due to long-range interactions: Schmalian&Wolynes, PRL 85 (2000)

Counting of Metastable States

Stability criterion (move particle from i to j)

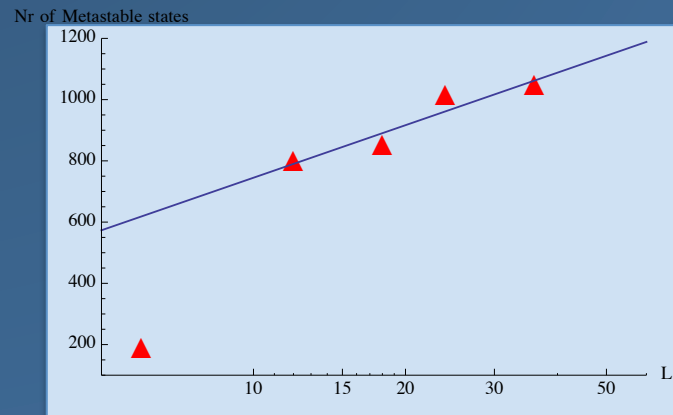
$$\Delta E = \epsilon_j - \epsilon_i - \frac{V}{|r_{ij}|} > 0$$

Exponentially many metastable states

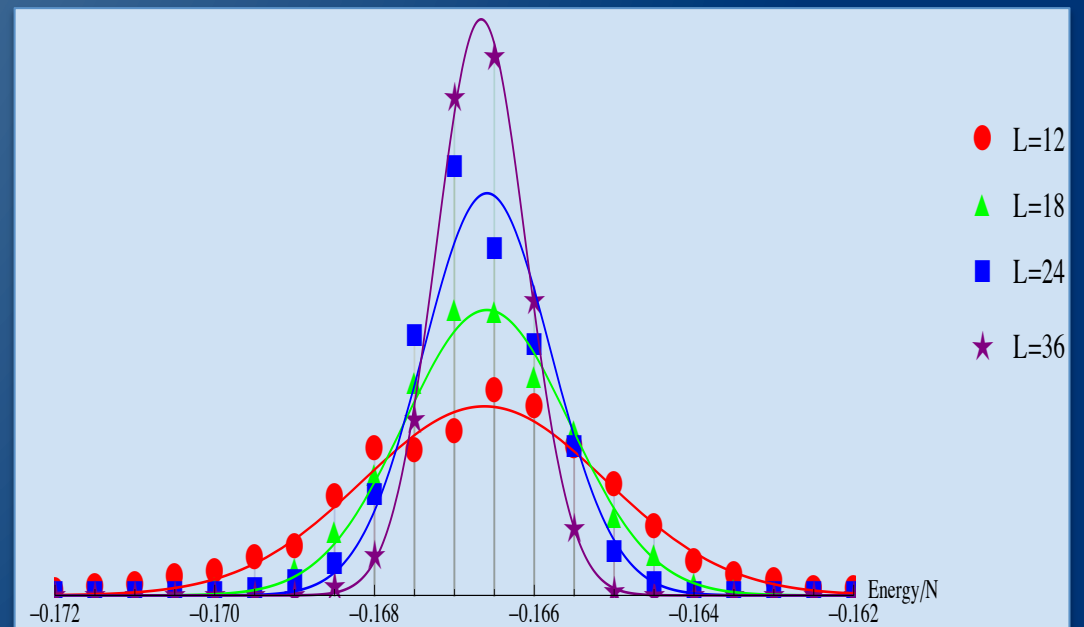


'Inherent' states

Only a small set of MS is relevant



L	N_{ms} (1440 sweeps)	expected total N_{ms}
6	183	185
12	695	798
18	708	849
24	775	1012
36	778	1046



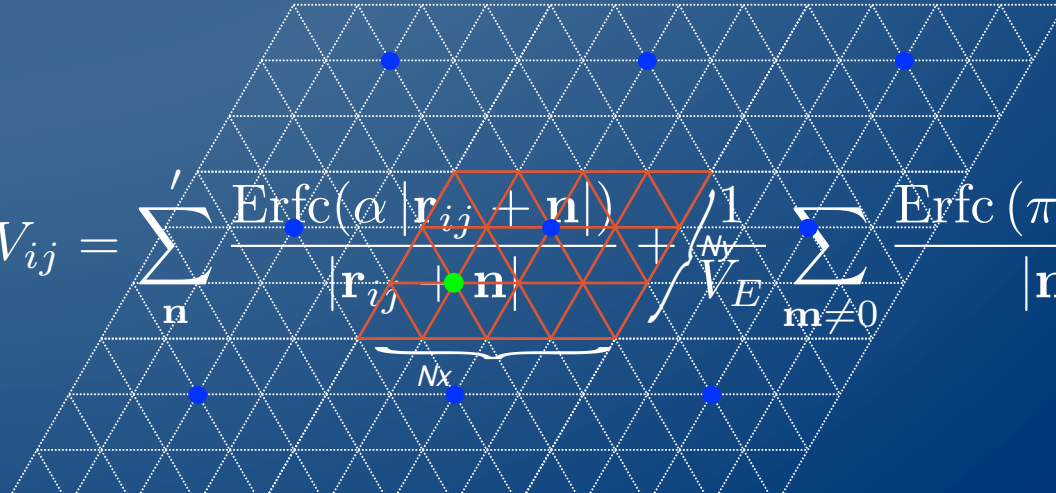
Higher energy than the ground state

Metastable states are indistinguishable in thermodynamic limit

Finite Temperature Monte Carlo

- **Metropolis:** accept local move with probability $P = \min(e^{-\Delta E/T}, 1)$
- Infinite range interactions on finite lattice?

Ewald summation:



The diagram shows a 2D lattice of sites. A central site is highlighted in green, and its position vector is \mathbf{r}_{ij} . A surrounding region of sites is outlined in red, with a width of N_x sites. The lattice is overlaid with a grid of dashed lines. The Ewald summation formula is shown to the right of the lattice.

$$V_{ij} = \sum_{\mathbf{n}}' \frac{\text{Erfc}(\alpha |\mathbf{r}_{ij} + \mathbf{n}|)}{|\mathbf{r}_{ij} + \mathbf{n}|} + \frac{1}{\sqrt{V_E}} \sum_{\mathbf{m} \neq 0} \frac{\text{Erfc}(\pi |\mathbf{m}| / \alpha)}{|\mathbf{m}|} \frac{1}{E} \cos(2\pi \mathbf{m} \cdot \frac{\mathbf{r}_{ij}}{|\mathbf{r}_{ij} + \mathbf{n}| \sqrt{\pi}}) \frac{q_i q_j}{|\mathbf{r}_{ij} + \mathbf{n}| \sqrt{\pi}} \delta_{ij}$$

Relaxation Time

Autocorrelation function

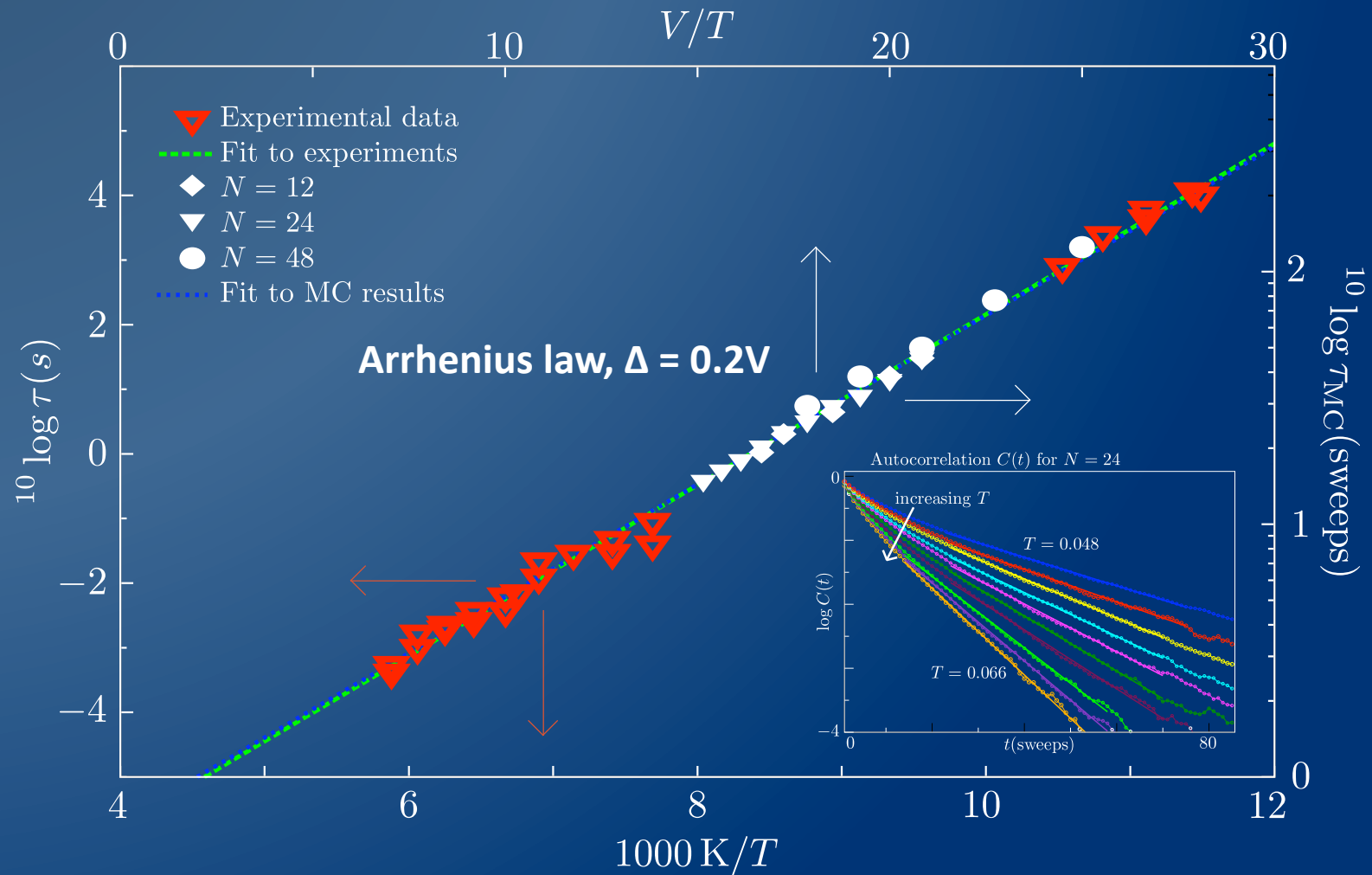
$$C(t + t_w, t_w) = \frac{2}{N} \sum_i \langle \delta n_i(t + t_w) \delta n_i(t_w) \rangle$$

If system equilibrates $C(t + t_w, t_w) \rightarrow C(t)$

Exponential tail of $C(t)$ gives relaxation time

Relaxation Time

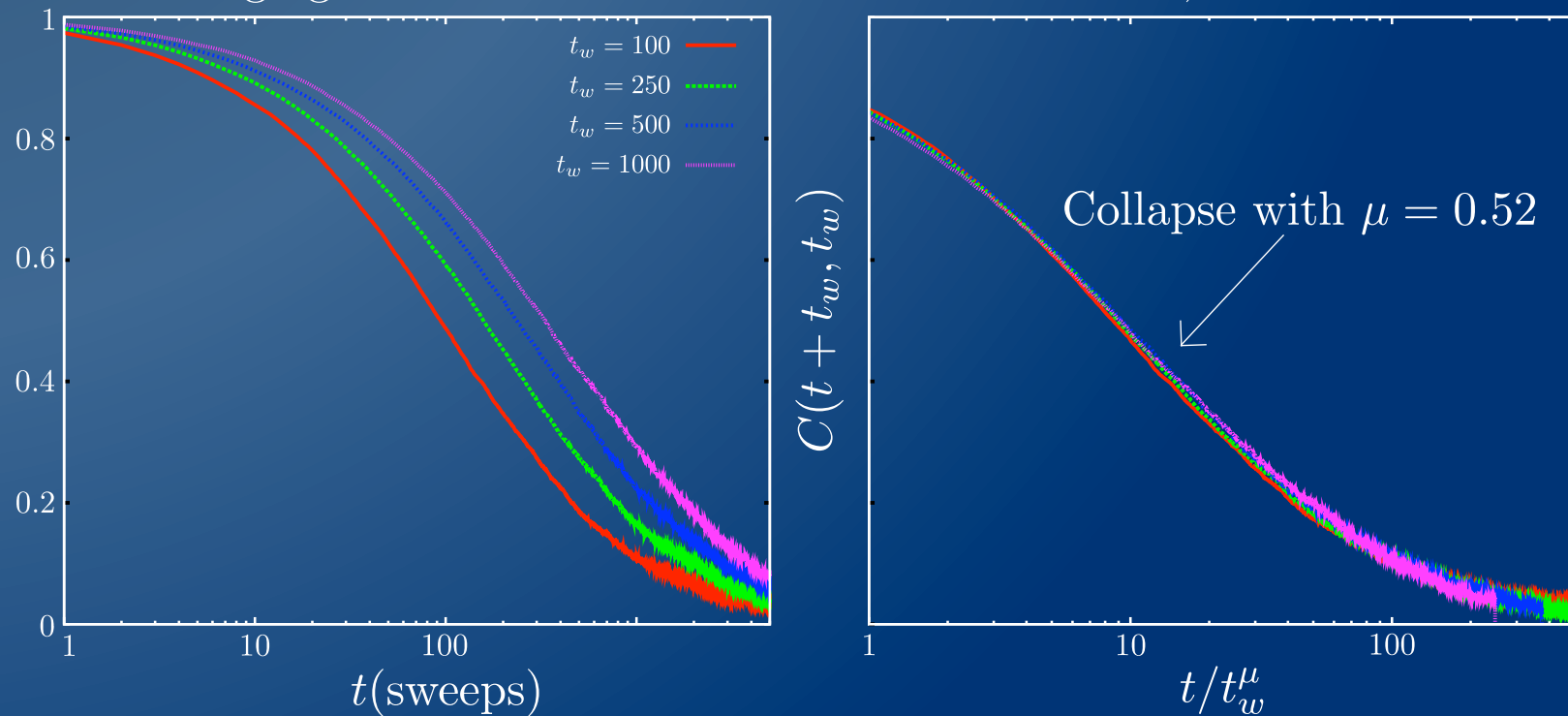
Experimental and Monte Carlo relaxation time



Aging & Memory

- Below T_s aging scaling $C(t + t_w, t_w) \sim F(t/t_w^\mu)$

Aging behavior in the autocorrelation function, $T = 0.03V$

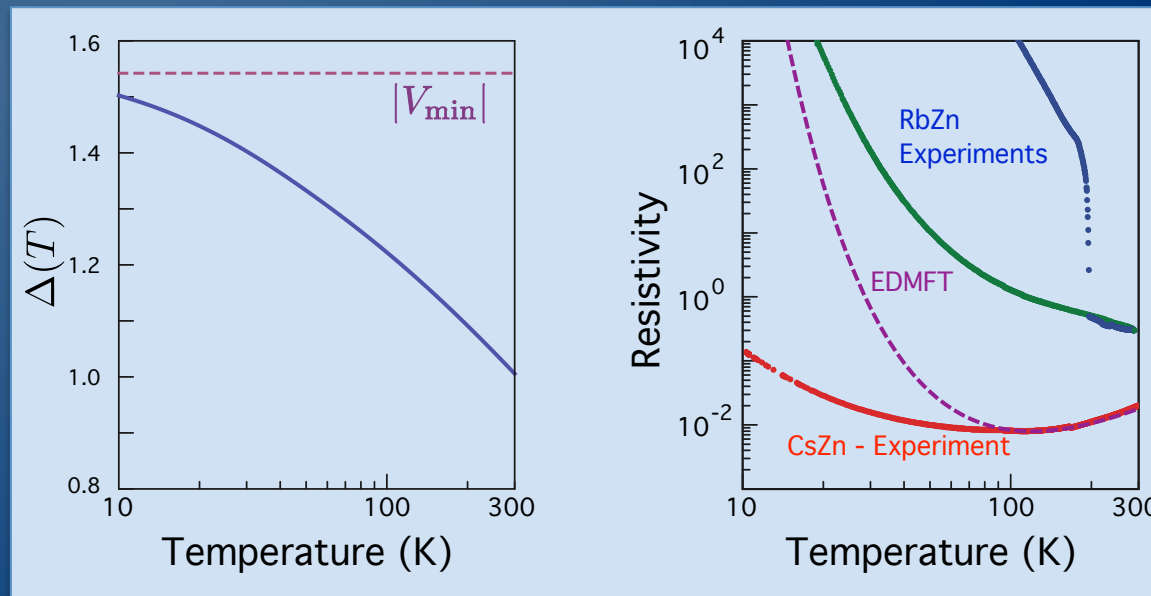


Conductivity

Extended Dynamical Mean Field Theory:

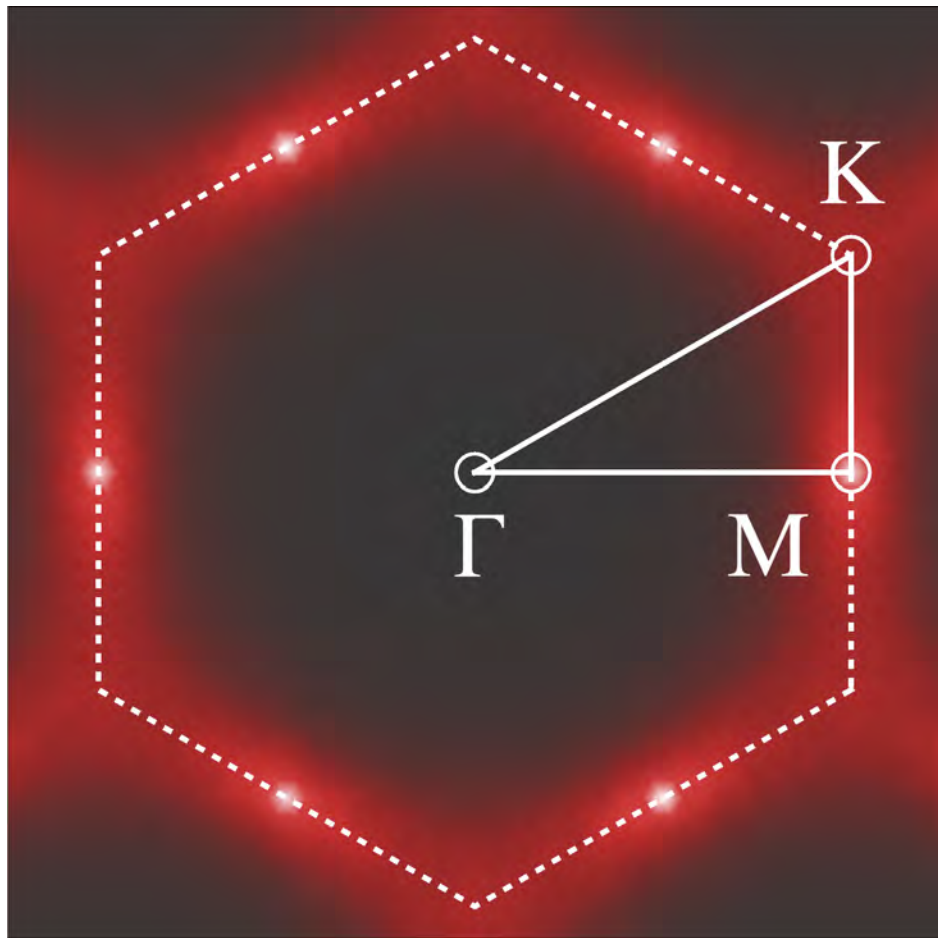
$$\Pi_{\mathbf{k}} = \sum_{\mathbf{r}_{ij}} \langle (n_i - \bar{n})(n_j - \bar{n}) \rangle e^{i\mathbf{k} \cdot \mathbf{r}_{ij}} = \frac{1}{(\bar{n} - \bar{n}^2)^{-1} + \beta(\Delta + V_{\mathbf{k}})}$$

$$\sigma_{DC} = \frac{\pi e^2}{\hbar a d} \int_{-\infty}^{\infty} d\omega \frac{\rho^2(\omega)}{4T \cosh^2 \frac{\omega}{2T}} = \frac{\beta e^2}{8\hbar a d} \sqrt{\frac{\beta\pi}{\Delta}} e^{-\frac{1}{4}\beta\Delta}$$

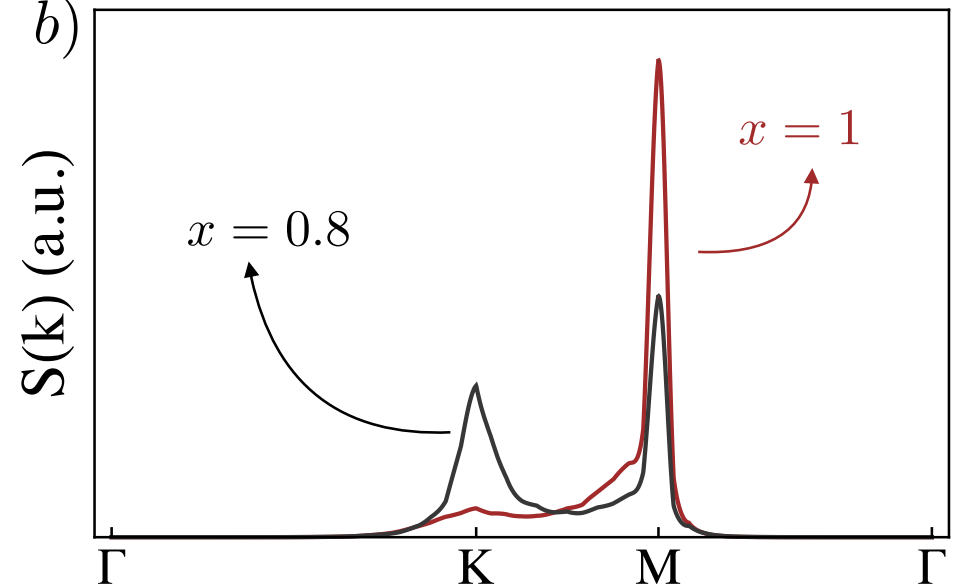


Density Correlations

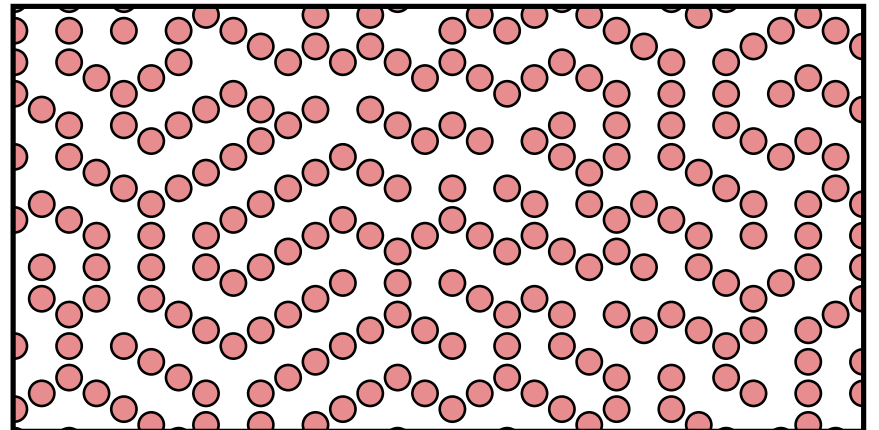
a)



b)



c)



Efros-Shklovskii gap (1)

Energy to remove or add a particle:

$$\epsilon_i = \mu_i + \frac{1}{2} \sum_{j \neq i} \frac{V}{|r_{ij}|} (n_j - \bar{n})$$

What is the **distribution** $g(\epsilon)$ of energies ϵ_i ?
(in the ground state of disordered system)

Stability when electron moves from i to j :

$$\Delta E = \epsilon_j - \epsilon_i - \frac{V}{|r_{ij}|} > 0$$

Efros-Shklovskii gap (2)

Distance between states close to Fermi level

$$|r_{ij}| > \frac{V}{\epsilon_j - \epsilon_i}$$

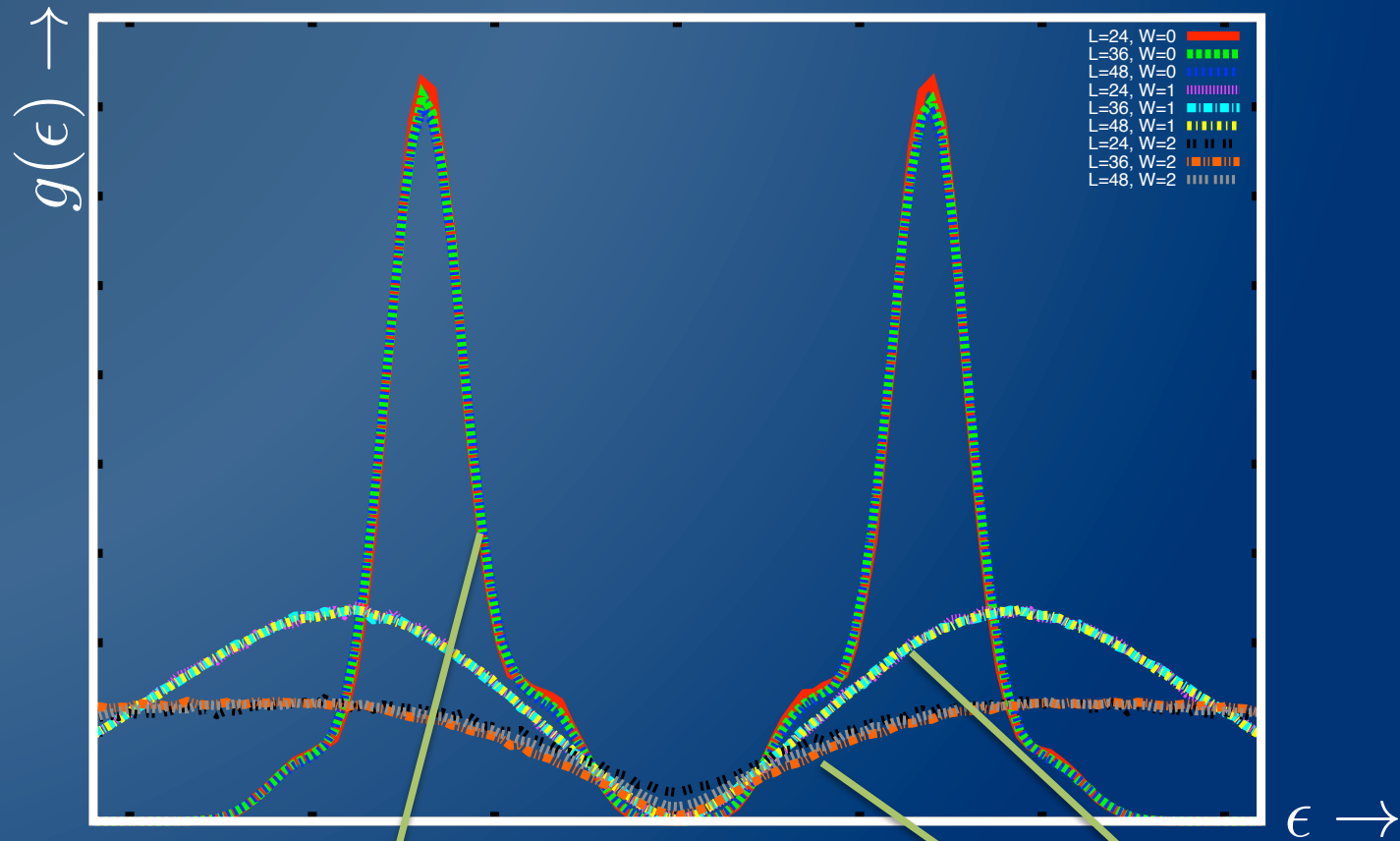
Stability sets upper **bound** for density of states

$$g(\epsilon) \leq \frac{d}{2\pi} |\epsilon|^{d-1}$$



Efros Shklovskii 1975

Coulomb gap without disorder



Stronger Coulomb gap
in absence of disorder!

$$g(\epsilon) \sim e^{-\Delta/|\epsilon|}$$

Linear ES gap
with disorder

Gap & Correlations (1)

Density correlations: $\Pi_{ij} = \langle (n_i - \bar{n})(n_j - \bar{n}) \rangle$

Chance to find electron at r: $P_e(\vec{r}) = \frac{1}{2} - 2\Pi(\vec{r})$

Relate local density of states: $P_e(\vec{r}) = \int_{-\infty}^0 d\epsilon' g_{\vec{r}}(\epsilon')$

Assumptions:

Density correlation decay like $\tilde{\Pi}(r) = \frac{3A}{\sqrt{2\pi r}} e^{-r/\xi}$

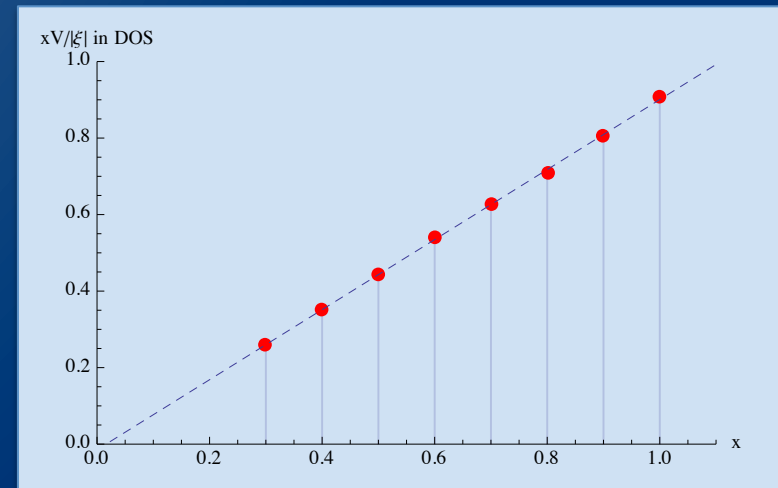
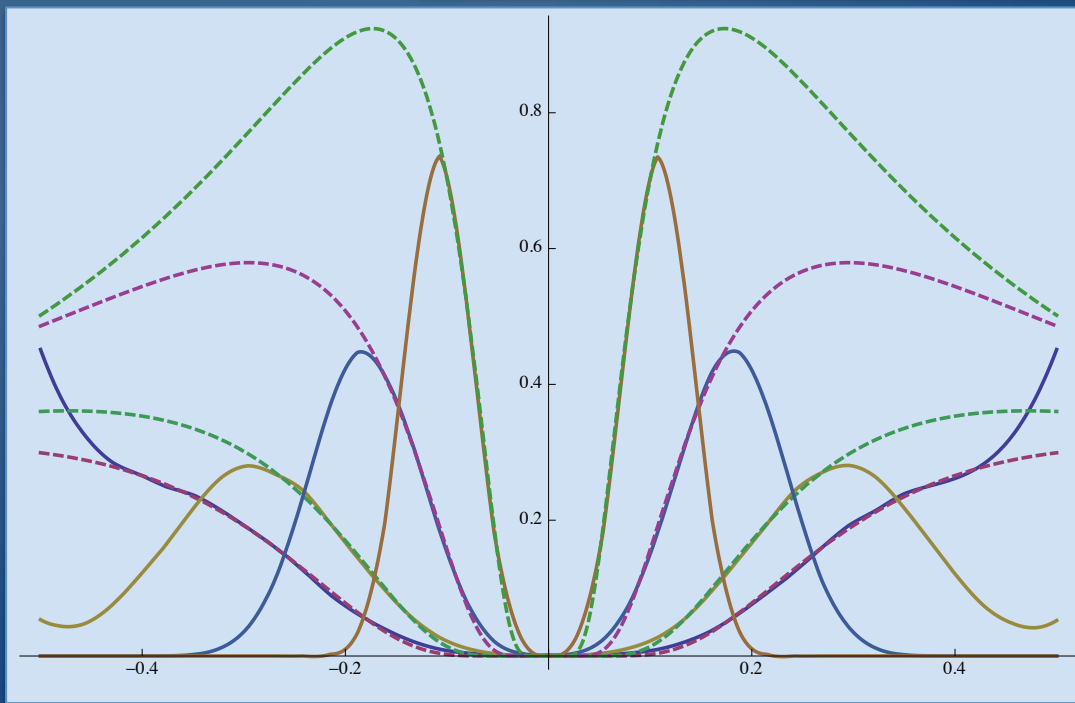
At long distances, LDOS becomes DOS:

$$\int_{-\infty}^{-\frac{V}{|r|}} d\epsilon' \tilde{g}_{\vec{r}}(\epsilon') \rightarrow \int_{-\infty}^{-\frac{V}{|r|}} d\epsilon' g(\epsilon')$$

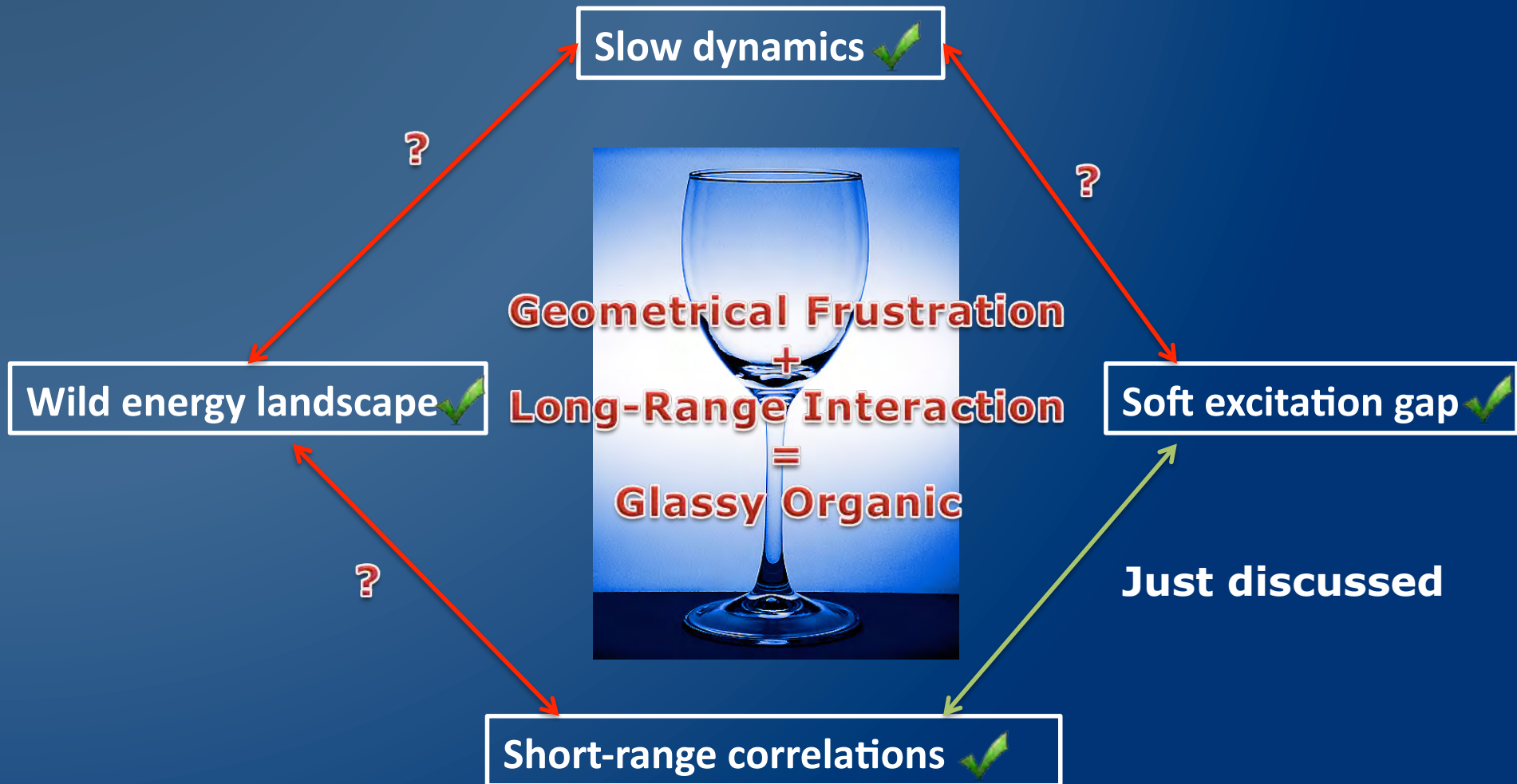
Gap & Correlations (2)

Hence DOS related to density correlations

$$g(\epsilon) = \frac{V}{r^2} g\left(\frac{\epsilon}{V}\right) \frac{V}{|\epsilon|}$$



Conclusions & Outlook



Reference: Mahmoudian, LR, Ralko, Fratini, Dobrosavljevic, arXiv1412.4441, PRL (2015)