Towards topological superconductivity in 2D TIs Vlad Pribiag

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SPICE Workshop, Mainz – August 3, 2015













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Overview

1. Topological Superconductivity and Majoranas

$$\gamma^{\dagger} = \gamma$$

2. 2D TI phase in InAs/GaSb quantum wells



3. Towards topological superconductivity in InAs/GaSb





Iopological protection of Majoranas

A pair of **spatially-separated** Majoranas forms a **non-local** fermionic state

 $c = \frac{1}{2}(\gamma_1 + i\gamma_2)$

→ state is **protected against local perturbations**



Bound Majorana modes are predicted to be **non-abelian** quasiparticles

- \rightarrow Majorana exchange follows braiding group
- \rightarrow topologically-protected braiding-based gates



Topological superconductivity in helical nanowires



J. Alicea. Rep. Prog. Phys. 75, 076501 (2012).

Ingredients for topological superconductivity:

- 1. material with large SOC
- 2. s-wave SC
- 3. B-field

$$E_z > \sqrt{\Delta^2 + \mu^2}$$

J.D. Sau, et al. *PRL* 104,040502 (2010).
R.M. Lutchyn, et al., *PRL* 105,077001 (2010).
Y. Oreg, et al. *PRL* 105,177002 (2010).
V. Mourik et al. Science 336, 1003 (2012).

Experimental signatures of Majorana bound states in nanowires

Kouwenhoven group, Science 336, 1003 (2012).



InSb nanowire

Zero-bias peak is a signature of a localized Majorana mode.

Similar results using InAs nws:

- Das et al., Nature Phys. 8, 887 (2012)
- H. O. H. Churchill *et al., PRB* 87, 241401 (2013).



Also, STM DOS spectroscopy of Fe nanowires on Pb substrate:

Yazdani group, Science, (2014).

Promising complementary approach: 2D topological insulator + superconductor

Topological insulators (2D)

Conventional insulator: n = 0

Insulating state	Energy	Conduction band
	B	Valence band
		Momentum

Quantum Hall (topological) insulator:



- Insulating bulk
- Broken time-reversal symmetry
- Chiral edge modes

2D topological insulator (Quantum Spin Hall insulator)



- Insulating bulk
- Time-reversal symmetric
- <u>Helical</u> edge modes (counter-propagating, spin-polarized)

C. L. Kane and E. J. Mele, *Science* **314** 1692 (2006). Thouless, Kohmoto, Nightingale and de Nijs, *PRL* **49**, 405 (1982).

2D topological insulators

Inverted band-structure in the bulk:



Need opposite band symmetries (e.g. Γ_6 and Γ_8 bands)

Subband alignment vs. thickness for HgTe/HgCdTe QWs



BHZ model for 2D TIs (a.k.a. Quantum Spin Hall Insulators):

- 4-band Dirac-like effective Hamiltonian, H_{BHZ} , derived from the 6-band Kane model
- E1 and H1 are effective s-like electron and p-like hole subbands in the QW
- At topological phase transition, mass term in H_{BHZ} changes sign

C.L. Kane and E.J. Mele, *PRL* **95** 226801 (2005). C.L. Kane and E.J. Mele, *PRL* **95** 146802 (2005). B.A. Barnevig, T.L. Hughes and S.C. Zhang, *Science* **314**, 1757 (2006).

Topological superconductivity in helical 1D systems

Common idea: superconducting coupling between helical modes



$$E_z > \sqrt{\Delta^2 + \mu^2}$$

Ingredients:

- material with large SOC
- s-wave SC •
- **B**-field \bullet

J.D. Sau, et al. PRL 104,040502 (2010). R.M. Lutchyn, et al., PRL 105,077001 (2010). Y. Oreg, et al. PRL 105,177002 (2010). V. Mourik et al. Science 336, 1003 (2012).



Ingredients:

- topological insulator
- s-wave SC

Helical modes are intrinsic Induced SC is protected from disorder by TRS

J. Alicea. Rep. Prog. Phys. 75, 076501 (2012).

J. Nilsson et al., PRB 101, 120403 (2008).

L. Fu and C.L. Kane, PRL 100, 096407 (2008).

L. Fu and C.L. Kane, PRB 79, 161408(R) (2009).

Majorana modes in a 2D TI

$$E_z < \sqrt{\Delta^2 + \mu^2}$$

- Superconductivity dominates the gap \rightarrow 1D topological superconductor at the edge
- Zeeman energy dominates gap → magnetically-gaped edge mode

Majorana bound states are localized at the domain wall:



Origin of subband inversion in InAs/GaSb:



Broken-gap band alignment:

E1-H1 band offset of ~150 meV

C. Liu et al. PRL, 100, 236601 (2008).



C. Liu et al. PRL, 100, 236601 (2008).

- 2D TI phase arises from a modified BHZ Hamiltonian
- Includes strong Bulk Inversion and Structure Inversion asymmetries:

 $H = H_{BHZ} + H_{BIA} + H_{SIA}$

- 4-band effective Hamiltonian, H, derived from 8-band Kane model, including the split-off bulk band, Γ_7 .
- **Topological phase transition** as function of thicknesses of GaSb (d_1) and InAs (d_2) :







Characteristics of InAs/GaSb:

- Based on III-V semiconductors

 → advantageous for growth and nanofabrication
- Gate-tunability between topological and trivial phases
- Non-trivial exciton physics
- Possible electron correlation effects





BHZ model \rightarrow topologicallyprotected helical edge modes

(C. Liu et al. PRL 100, 236601 (2008))

Imaging edge modes of InAs/GaSb

Scanning SQUID imaging:











E.M.Spanton et al., PRL 113 026804 (2014).

Quantized Conductance in InAs/GaSb



$$G_{14,32} = 4e^2/h$$





L. Du et al., *PRL* **114**, 096802 (2015). M. Büttiker, *PRB* **38**, 9375 (1988).

Phase diagram of InAs/GaSb



Experiment: F. M. Qu et al. PRL 115, 036803 (2015)

Theory: C. X. Liu et al. PRL 100, 236601 (2008)

Gate-tunability between topological and trivial phases in InAs/GaSb

Josephson Junctions based on InAs/GaSb

First step towards Majorana physics:

Supercurrent carried by the edge-modes of InAs/GaSb V.S. Pribiag et *al.*, *Nature Nano*. **10**, 593 (2015).



Previous work on hybrid 2D TIs / superconductor devices:

- Andreev reflection in InAs/GaSb Knez et al. PRL, **109**, 186603 (2012).
- Induced superconductivity in the edge modes of HgTe/CdTe S. Hart, H. Ren *et al.*, *Nat.Phys.* **10**, 638 (2014).

S-InAs/GaSb-S junctios







Gate-tunable Josephson junctions based on the 2D TI InAs/GaSb

Normal state: $B_z = 0.1T$



DC Josephson effect



AC Josephson effect







Shapiro steps:
$$V = n \cdot \frac{hf_{RF}}{2e}$$

 $n = 0, \pm 1, \pm 2, \dots$

Superconducting quantum interference





Extracting the spatial distribution

measurement extract spatial distribution $I_c(B_Z) \equiv \max[I_s(B_Z)] = \left| \int_{-\infty}^{\infty} J_c(x) e^{ikx} dx \right|$ $k = 2\pi L B_Z / \Phi_0$ L is the effective junction length

<u>Dynes and Fulton (PRB 1971)</u> \rightarrow estimate the complex phase:

- $J_c(x) = J_{even}(x) + J_{odd}(x)$
- If $J_c(x) \approx J_{even}(x)$, can interpolate near the nodes to estimate $J_{odd}(x)$

A more general treatment gives good agreement with the Dynes&Fulton method:

Proximity-induced superconductivity and Josephson critical current in quantum spin Hall systems Phys. Rev. B **90**, 224517 – Published 18 December 2014

Hoi-Yin Hui, Alejandro M. Lobos, Jay D. Sau, and S. Das Sarma





InAs control device



InAs-only \rightarrow bulk supercurrent

Edge-mode width

Fourier limited: $\Delta x \sim W \Phi_0 / \Delta \Phi$



IS and Majorana detection in a 2D TI

Conductance measurements:

- Experimental Signature: ZBCP in the gate-averaged conductance
- insensitive to quasiparticle poisoning



Proposal by S. Mi et al., *PRB* **87** 241405(R) (2013).

IS and Majorana detection in a 2D TI

Conductance measurements:

- Experimental Signature: ZBCP in the gate-averaged conductance
- insensitive to quasiparticle poisoning

Current-phase relation:

- Majoranas γ_1 and $\gamma_2 \rightarrow$ single-electron can be transported across Josephson junction
- Doubling of current-phase (flux) period: $\frac{h}{2e} \rightarrow \frac{h}{e}$





Proposal: L. Fu and C.L. Kane, PRB 79, 161408(R) (2009).

Summary

- Demonstrated edge-mode supercurrent in InAs/GaSb
- <u>Next</u>: Detection of Majoranas and TS in InAs/GaSb

Thank you!



Q.2 Q.1 266 nm + 257 nm 266 nm + 257 nm -8 -6 -4 -2 0 2 4 6 8 x (μm)

V.S.Pribiag et al., Nature Nano. 10, 593 (2015)