

Nanosystems from ions, spins and electrons

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Spin-Orbitronics for advanced magnetic memories: chiral domain walls and anti-skyrmions

Max Planck Institute for Microstructure Physics, Halle



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N[†]ISE

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Alexander von Humboldt Stiftung/Foundation

> Horizon 2020 European Union funding for Research & Innovation

Atomically engineered materials; Spintronics; Cognitive devices; 3D structures; Room temperature superconductivity; Topological matter; Combinatorial materials discovery....

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Spintronics – from materials and phenomena to applications



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IBM/DARPA Advanced Magnetic R andom Access Memory



Exchange-biased magnetic tunnel june to nonvolatile magnetic random acces

S. S. P. Parkin,^{a)} K. P. Roche, M. G. Samant, P. IBM Research Division, Almaden Research Center, San Jos

R. E. Scheuerlein IBM Storage Division, San Jose, California 95120

E. J. O'Sullivan, S. L. Brown, J. Bucchigano, D. V. P. L. Trouilloud, R. A. Wanner, and W. J. Gallagh IBM T. J. Watson Research Center, Yorktown Heights, New

Exchange biased magnetic tunnel junction (MTJ) struc forming magnetic memory storage elements in a novel been developed which exhibit very large magnetores temperature, with specific resistance values ranging do values enhanced by moderate thermal treatments. L





Figure 10

VO

elements with areas as small as 0.17 $(\mu m)^2$. The (a) Photograph of 16-Mb product demonstrator chip, with indicacharacteristics of an MTJ element integrated with a s tions of main functional subunits. From [58], with permission; properties of the MTJ element itself. © 1999 Ameri ©2005 IEEE. (b) SEM cross section showing three cells: M1, M2, [S0021-8979(99)77508-0] and M3 are segments of three global wiring layers; VA, V1, and CA are segments of three via layers; PC is a segment of the

CA are segments of three via layers; PC is a segment of the transistor gate layer; JA is a segment of the tunnel-junction definition layer; and MA is a segment of the local wiring layer.

Innately 2D!

2016 Samsung MICRON Global Foundries Intel Everspin Avalanche Crocus Spin Transfer Technologies

MRAM widely available in 2018!

. . . .

1995

Amplitude (arb. unit)

-2 0 2

MTJ resistive response to 5

Accomplishments over las

easy-axis current pulse strength

-Achieved MR>40%

-Resistance of MT.I

-1ns write time for

-MRAM1 design co cells; fabrication co

Time (ns)



19



Spin polarized current \rightarrow manipulate domain walls!

- electrical current in ferromagnetic metals is innately spin-polarized

Magnetization



- Current driven domain wall motion



- → Giant magnetoresistance
- \rightarrow Spin transfer switching
- \rightarrow Current induced magnetization precession

- → Domain wall moves in direction of flow of spin angular momentum
- → For Ni and Co based soft magnetic alloys domain wall moves in direction of *electron flow*
- → All domain walls move in same direction!

First demonstration!

Current-Controlled Magnetic Domain-Wall Nanowire Shift Register

Masamitsu Hayashi, Luc Thomas, Rai Moriya, Charles Rettner, Stuart S. P. Parkin*

The controlled motion of a series of domain walls along magnetic nanowires using spin-polarized current pulses is the essential ingredient of the proposed magnetic racetrack memory, a new class of potential non-volatile storage-class memories. Using permalloy nanowires, we achieved the successive creation, motion, and detection of domain walls by using sequences of properly timed, nanosecond-long, spin-polarized current pulses. The cycle time for the writing and shifting of the domain walls was a few tens of nanoseconds. Our results illustrate the basic concept of a magnetic shift register that relies on the phenomenon of spin-momentum transfer to move series of closely spaced domain walls.

SCIENCE VOL 320 11 APRIL 2008

Racetrack Memory 1.0



Fig. 4. Demonstration of a three-bit unidirectional magnetic DW shift register. (**A**) Data are encoded by the magnetization direction of three domains in the nanowire. (**B**) Nanowire resistance variation when a pulse sequence is used to write and shift along the register the sequence 010111 two times in succession. The light and dark regions indicate writing and shifting operations, respectively. The table shows the corresponding evolution of the states of the three bits during these operations. The highlighted digits show how the input bit sequence is transferred to the output after two write/shift operations. (**C**) Schematic illustration of the shift-register operation. Black squares and white squares represent HH and TT DWs, respectively. Black arrows represent the magnetization direction within each domain. Blue and red arrows represent the electron flow direction.



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Racetrack Memory 3.0



- \rightarrow 20 domain walls moved in lock step with current pulses
- \rightarrow High velocity at low current densities
- → Narrow domain walls (~6 nm)
- \rightarrow Very thin racetracks (~1 nm)

Ryu et al. Nature Nanotechnology (2013) Yang et al. Nature Nanotechnology (February 2015) Parkin et al. Nature Nanotechnology (March 2015)



Chiral domain walls from Dzyaloshinskii-Moriya Interaction

(#1) Perpendicular Magnetic anisotropy (PMA)

- materials: Co/Ni, Co/Pt, Co/Pd, RE/TM multilayers, low symmetry magnetic materials
- racetracks magnetized perpendicular to the plane of the racetrack



(#4) Spin Hall effect: conversion of charge to spin current





Measurement of Spin orbit torque \rightarrow Spin Hall angle

ST-FMR: Spin Torque Ferromagnetic Resonance



Spin Hall torque and Oersted Field torque

Liu et al., Phys. Rev. Lett. (2011)

 \blacktriangleright Ratio \rightarrow charge to spin conversion efficiency

ST-FMR: Spin Torque Ferromagnetic Resonance



Interface is important!



→ Transparency of NiFe/Pt and Co/Pt interfaces not 100 % → NiFe/Pt: 25%; Co/Pt: 50%
→ Correct for transparency then Spin Hall Angle: SHA of Pt = 0.20

Zhang et al., Nature Phys. (2015)

IrMn₃ – fcc structure - a triangular antiferromagnet





Lattice without magnetism Four equivalent axes

- [111] [1-11]
- [11-1]
- [1-1-1]

Lattice with noncollinear AF $[111] \rightarrow$ one special plane

[1-11] [11-1] [1-1-1]



Zhang et al. Sci. Adv. (2016)

Strong effect on SHA of field annealing in *perpendicular* field \rightarrow change in AF domain structure in IrMn₃



- \rightarrow Tunes the AFM domain configuration at the IrMn₃ interface
- → Only effective for thicknesses of IrMn3 for which the AF domains can be reset i.e. the blocking temperature $T_B < T_{AN}$

N.B. in-plane annealing \rightarrow changes in-plane exchange bias field \rightarrow no effect on SHA

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IrMn₃: Calculation of AHC and SHC



Calculations by Binghai Yan et al.

Zhang et al. Sci. Adv. (2016)

AHE in *hexagonal* Mn₃Ge: a triangular AF



- \rightarrow Results on single crystals
- \rightarrow Facet dependent AHE conductivity in Mn₃Ge: a non-collinear AFM
- → Anticipate strong SHE/ SHC in this and related materials (*h*-Mn₃Sb, *h*-Mn₃Sn)



Nayak et al. Sci. Adv. (2016)

Charge to spin conversion vias the Spin Hall Effect

- Spin Hall effect: dramatic evolution
- Novel materials show large spin Hall effect
 - Triangular antiferromagnets
 - Novel crystalline phases
- Important technological applications
 - Racetrack DW velocities > 1,000 m/sec
 - 3T Single DW Racetrack replacement for SRAM



Zhang et al. Nat. Phys. (2015) Demasius et al., Nat. Comm. (2016) Nayak et al. Sci. Adv. (2016) Zhang et a. Sci. Adv. (2016)



Chiral Spin Torque \rightarrow drives domain walls very fast!

Top view: Counter-clockwise chiralty



- DMI local field \rightarrow Neel domain walls
- DMI field **H**_{DMI} sets a specific chirality the Neel domain Walls
- spin Hall torque causes **M** to rotate towards spin accumulation **S**
- DMI field torque moves all the DWs along the current flow direction.

Ryu et al., Nat. Nanotech. (2013); Yang et al., Nat. Nanotech. (2015)

Dipolar fringing fields emitted from Domains



RM 4.0: Very High DW Speeds in SAF PMA racetracks



- → SAF: Synthetic antiferromagnet: upper racetrack = exact mirror image of lower
- → DW velocity increases as degree of compensation of moments in upper and lower racetracks increased
- \rightarrow 4x highest DW velocity yet reported \rightarrow Speeds exceed 1.5- 5 km/sec
- → AF exchange field > DMI field

Yang et al. Nature Nanotechnology (2015)

Racetrack Memory 4.0



- \rightarrow 20 domain walls moved in lock step with current pulses
- \rightarrow High velocity at low current densities
- → Narrow domain walls (~6 nm)
- \rightarrow Very thin racetracks (~1 nm)
- \rightarrow Giant domain wall velocities in Synthetic Antiferromagnet racetracks



Ryu et al. Nature Nanotechnology (2013) Yang et al. Nature Nanotechnology (February 2015) Parkin et al. Nature Nanotechnology (March 2015)

SORBET - complex interplay of 4 spin-orbit derived phenomena

- 1. Perpendicular magnetic anisotropy from broken symmetry from interfaces
- 2. Proximity induced magnetization
- 3. Chiral domain walls DMI
- 4. Spin currents from spin Hall effect (SHE)







Dramatic dependence of DW velocity on Curvature





 $R = 7 \ \mu m$ and $w = 2 \ \mu m$



DW motion after 2 pulses (0.6 x 10⁸ A/cm² for 100ns)

 \rightarrow Domain size decreases or increase by 10 times depending on curvature

- DW velocities are different for all J.
- CW (CCW) DW motion : $\kappa > 0$ ($\kappa < 0$)
- $\bigcirc | \otimes \kappa > 0 = \otimes | \bigcirc \kappa < 0 \rightarrow \text{slower}$
- $\odot | \otimes \kappa < 0 = \otimes | \odot \kappa > 0 \rightarrow \text{faster}$
- DW velocity in straight wire ($\kappa = 0$) is in-between $\kappa > 0$ and $\kappa < 0$.

Garg et al. Science Advances (May 5, 2017)

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Quasi 2D Model for curved wires with radius dependent current





- DW is not straight but curved due to the cylindrical symmetry of experimental parameters such as electrical current, SHE and DMI.
- Key point → transformation between cylindrical moving frame and cartesian static coordinates → DW magnetization and DW tilting angle in static coordinates are a function of DW position.

Quasi Two Dimensional Model Results for curved wires: time-resolved



- Terminal velocity for $\bigcirc | \otimes$ with $\kappa > 0$ is larger than that for $\otimes | \bigcirc$ with $\kappa > 0$.
- It takes shorter for $\odot | \otimes$ with $\kappa > 0$ (~ 10 ns) to reach the terminal velocity than for $\otimes | \odot$ with $\kappa > 0$ (~20 ns).
- DW magnetization reaches steady state much faster (< 1ns) than DW tilting angle (> 20 ns).

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- DW velocity (torque) at outer rim is faster than at inner rim.
- Angular velocity is the same along the DW in steady state.

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DW motion in synthetic antiferromagnetic curved wires \rightarrow no change !



- Curvature does not affect DW motion in SAF nanowires.
- Exchange coupling torque dominant \rightarrow insensitive to DW tilting
- Curvature induced lock-step motion breakdown problem resolved !!

Garg *et al.* Science Advances (May 5, 2017) Stuart Parkin – Spin on Electronics!

3T single domain wall Racetrack \rightarrow SRAM replacement



- 100 picosecond writing speed
- Non-volatile & very low energy
- 40% size of SRAM

	ultra dense	1DWRM 6	3T-SOT- MRAM 6 Q	1DWRM 8	3T-SOT- MRAM 8 Q
	SRAM	Q layout	layout	Q layout	layout
	1	1	1	2	2
Feature size (F)	14	14	14	14	14
Fin pitch (FP) (nm)	42.67	42.67	42.67	42.67	42.67
contact poly pitch (CPP) (nm)	80	80	80	80	80
Bit line pitch (BLP) (FP)	8	4	4	4	4
Bit line pitch (BLP) (nm)	341.3	170.7	170.7	170.7	170.7
Word line pitch (WLP) (CPP)	2	1.5	1.5	2	2
Word line pitch (WLP) (nm)	160	120	120	160	160
Cell (Q)	16	6	6	8	8
Cell (nm ²)	54613	20480	20480	27307	27307
Cell (F ²)	278.6	104.5	104.5	139.3	139.3
Cell/SRAM	1	0.375	0.375	0.5	0.5
Energy (fJ)	96	9.6	9.6	19.2	19.2
Switching speed (ps)	667	< 200	< 200	< 200	< 200
Max transistor	30	30	30	60	60

Skyrmions at room temperature in Mn₂(Pd₄₀Pt₆₀)₁Sn





Nayak et al. (unpublished)



Non collinear spin structures in $Mn_{1.4}(Pt_{0.9}Pd_{0.1})$ Sn

Electronics!

Skyrmions and anti-skyrmions

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Anti-Skyrmions!

1st observation - Stable above 300 K

Nayak et al. (2017)

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Phase diagram: Mn_{1.4}(Pt_{0.9}Pd_{0.1}) Sn

Nayak et al. (submitted)

Anti-Skyrmions!

1st observation - Stable above 300 K

Nayak et al. (submitted)

Micromagnetic simulations - phase diagram

OOMMF simulation of phase change as a function of perpendicular field strength, with: (a) H_z =0.09T, helix+antiskyrmion phase, (b) H_z =0.15T, helix+antiskyrmion phase, (c) H_z =0.21T, antskyrmion phase, (d) H_z =0.39T, antikyrmion phase, (e) H_z =0.47T, antiskyrmion + spin polarized phase and (f) H_z =0.50T, spin polarized phase.

LTEM – anti-skyrmion lattice evolution

Field dependence of the antiskyrmion size at various temperatures. The error bar is the standard deviation of the size distribution. Field dependence of the antiskyrmion lattice mean angle at several different temperatures. The inset show the corresponding standard deviation of the lattice angles.

Analysis of a LTEM image of the antiskyrmion lattice at 200K under a perpendicular field of H=0.23T.

Magnetic bubbles / skyrmions in W | CoFeB films

Hz

 \otimes

2 Oe : Bubble

4 Oe : Single domain

Skyrmions on the "Racetrack"

"hedgehog" (Néel) type Examples: thin Fe/Ir → SOC + interface sets chirality

"vortex type" (Bloch) chiral phase examples: MnSi; films with perpendicular anisotropy plus magnetic dipolar coupling (achiral)

Skyrmions in CoFeB \rightarrow speed 0.000001 m/sec at j= 10⁵ A/cm²

Jiang et al. Science (2015)

Skyrmions in CoFeB thin films

Hz=0 @ different temperature

293.9K stripe 287.5K stripe 282.6K uniform 277.7K uniform 272.9K uniform

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Spin Orbitronics

- Chiral spin torque and giant exchange torque
 - Ultrafast current induced DW velocities in racetracks with no magnetization!
 - Curvature strongly affects DW velocity but eliminated in SAF racetrack
- Novel materials show large spin Hall effect
 - Triangular antiferromagnets⁺
 - Novel crystalline phases + +
- Important technological applications
 - Racetrack DW velocities > 1,000 m/sec
 - 3T Single DW Racetrack replacement for SRAM
 - Promises 3T device with SRAM performance but increased density, much lower energy consumption & non-volatility (fast start-up)
- Non collinear spin textures!
 - Anti-skyrmions...

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Cr, Bi2Se3: ~100

(2015)

Bi₂Se₃: 1-2

(2013)

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