



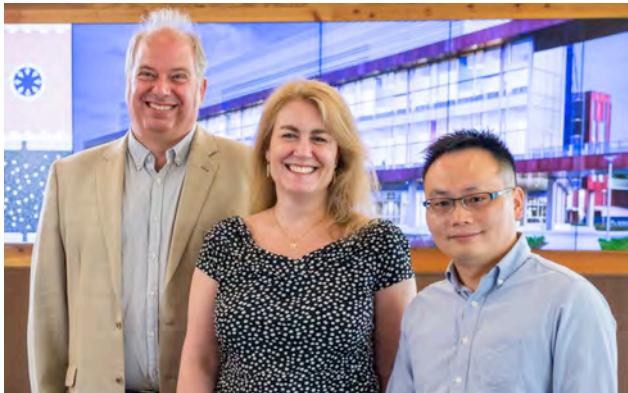
# STABILIZING AND CONTROLLING ROOM TEMPERATURE MAGNETIC SKYRMIONS

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Materials Science Division  
Argonne National Laboratory

Supported by the U.S. Department of Energy, Office of Science,  
Materials Sciences and Engineering Division.

# ACKNOWLEDGEMENTS



**Axel Hoffmann      Wanjun Jiang (now at Tsinghua U.)**

**Wei Zhang, M. Benjamin Jungfleisch, H. Somaily, Frank Y. Fradin, John E. Pearson, Olle Heinonen**

Argonne National Laboratory

**Xichao Zhang, Yan Zhou**

University of Hong Kong

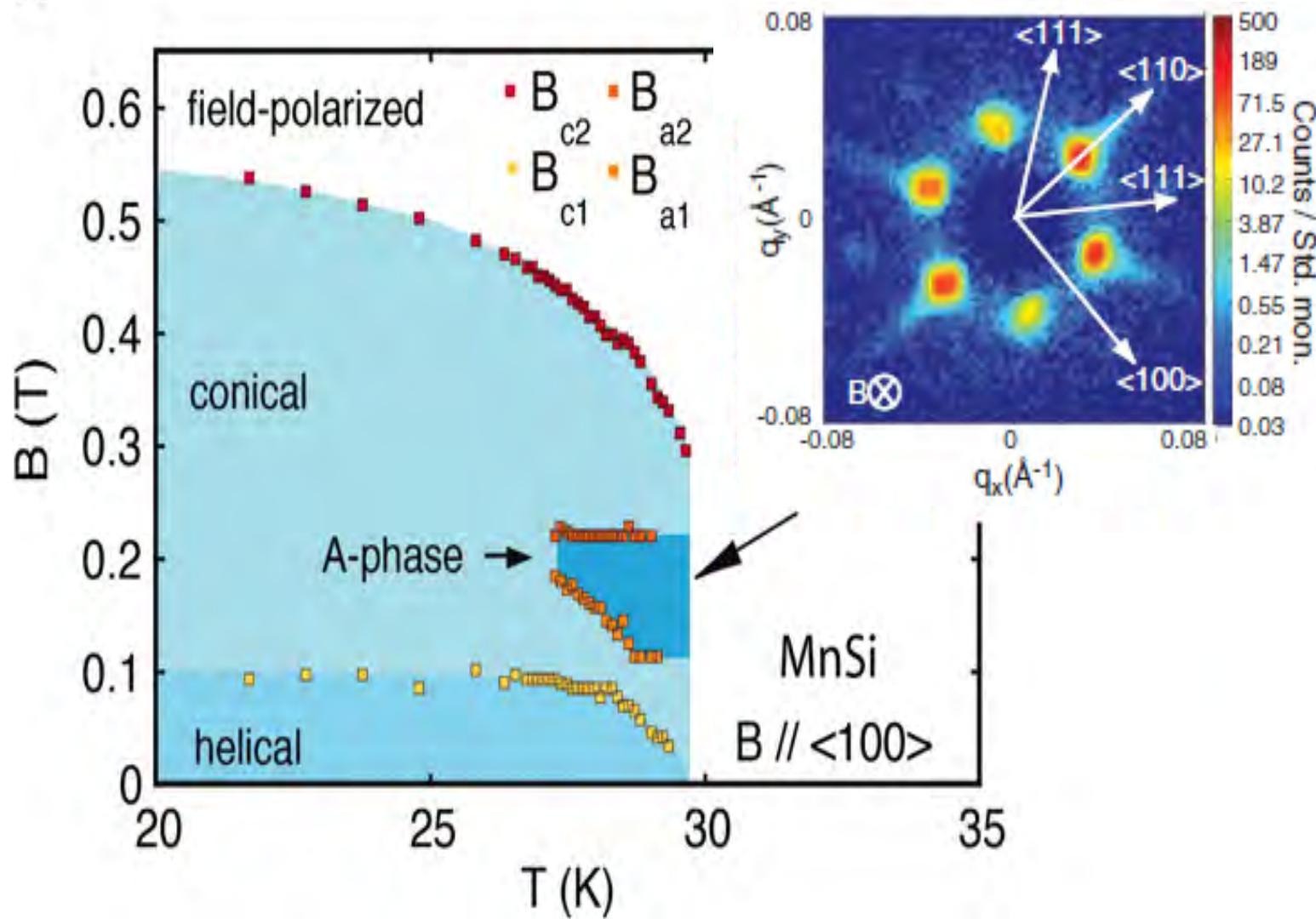
**Xiao Wang, Xuemei Cheng**

Bryn Mawr College

**Pramey Upadhyaya, Guoqiang Yu, Yaroslav Tserkovnyak, and Kang L. Wang**

University of California Los Angeles

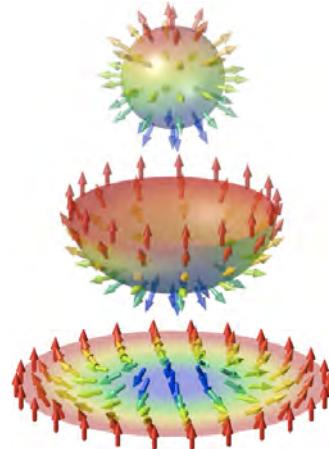
# DISCOVERY OF MAGNETIC SKYRMIONS



S. Mühlbauer *et al.*, Science 323, 915 (2009)

# MAGNETIC SKYRMIONS

Non-Trivial Topology

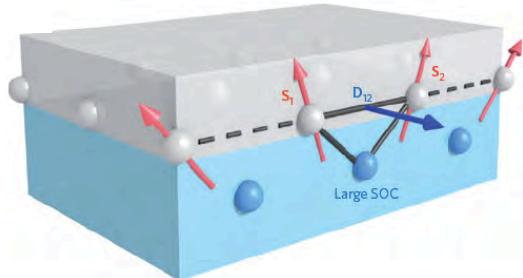


Topological Charge

$$Q = \frac{1}{4\pi} \int \mathbf{m} \cdot (\partial_x \mathbf{m} \times \partial_y \mathbf{m}) dx dy$$

$$Q = \pm 1$$

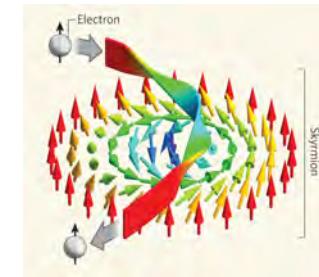
Stabilized by Chiral  
Dzyaloshinskii-Moriya Interaction



$$\mathbf{D}_{ij} \cdot (\mathbf{S}_i \times \mathbf{S}_j)$$

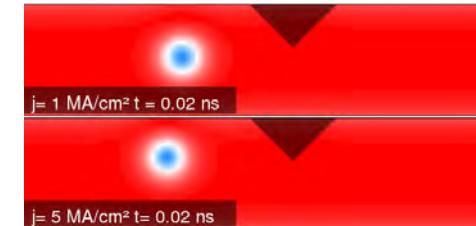
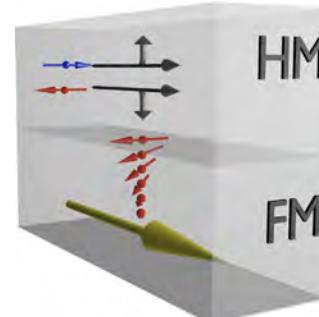
Efficient Manipulation  
by Electric Charge Currents

Emergent Magnetic Field



$$h \approx \frac{\Phi_0}{\pi R^2} \approx \frac{\Phi_0}{\pi a^2} \left(\frac{D}{J}\right)^2 \quad h \sim 100 \text{ T}$$

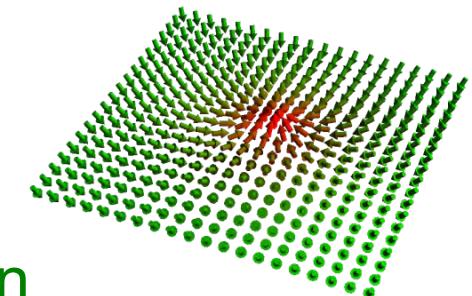
Motion due to Spin Hall Effect



# SKYRMIONS ARE STABILIZED BY CHIRAL INTERACTIONS

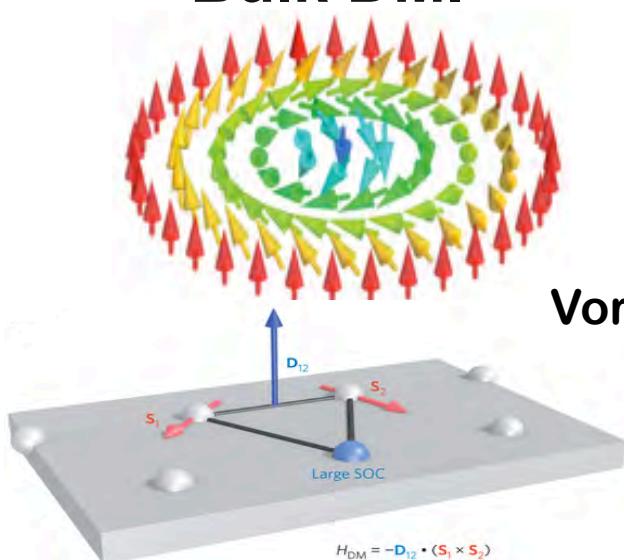
$$H = \sum_{\langle ij \rangle} -JS_i \cdot S_j + D_{ij} \cdot (S_i \times S_j) - \sum_i B \cdot S_i$$

Ferromagnetic    Helical Spiral    Skyrmion



Dzyaloshinskii-Moriya Interaction (DMI) requires Inversion Symmetry Breaking

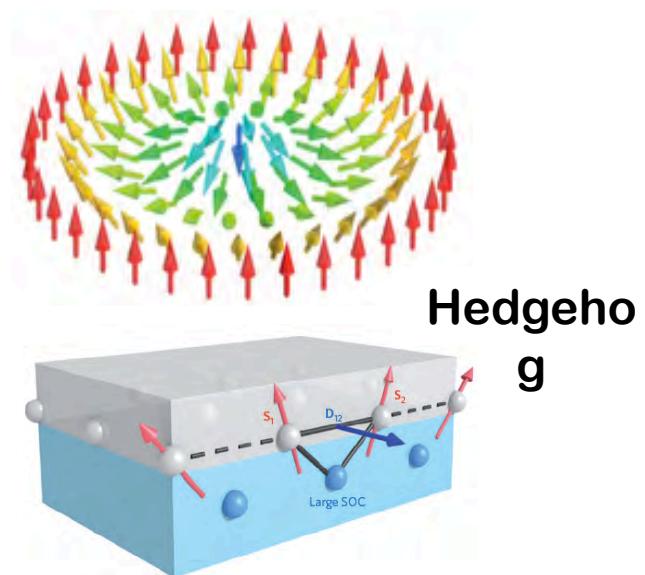
Bulk DMI



e.g., B20 compounds (MnSi, etc.)

Breaking

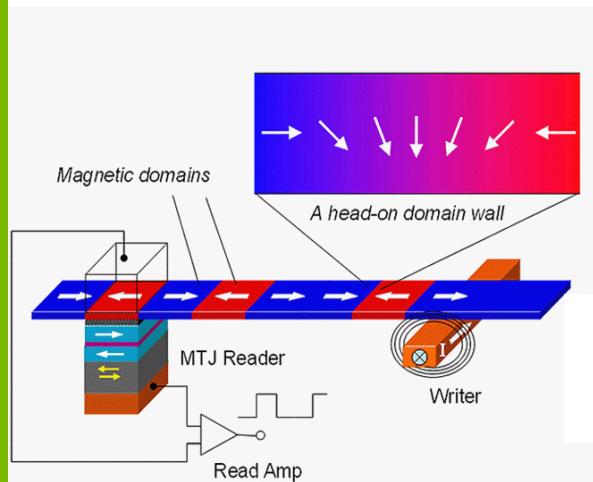
Interfacial DMI



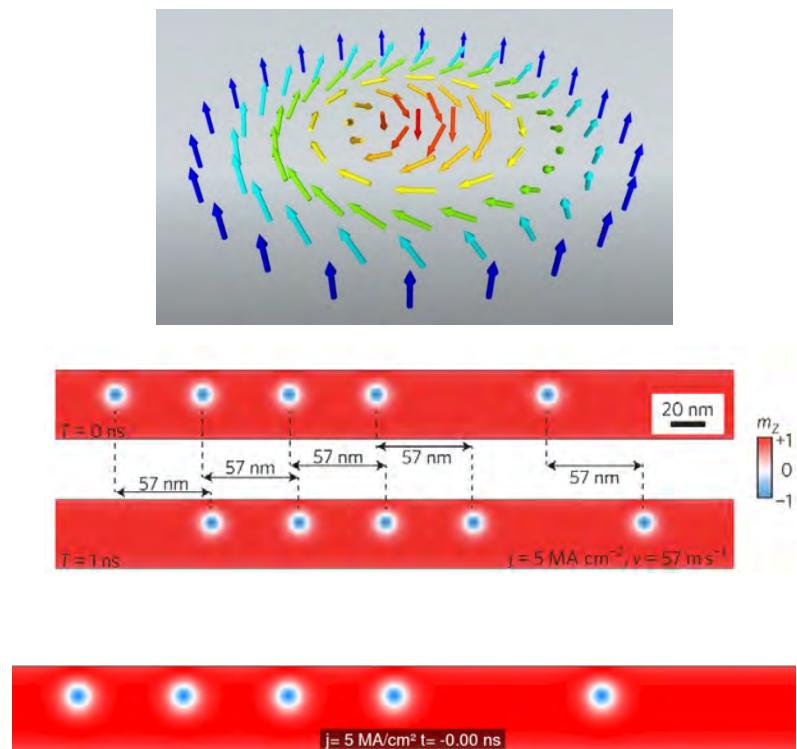
e.g., Co/Pt, Ni<sub>80</sub>Fe<sub>20</sub>/Ta, etc.

# ENCODING INFORMATION IN SPIN TEXTURES

## Racetrack Memory



## Skyrmions

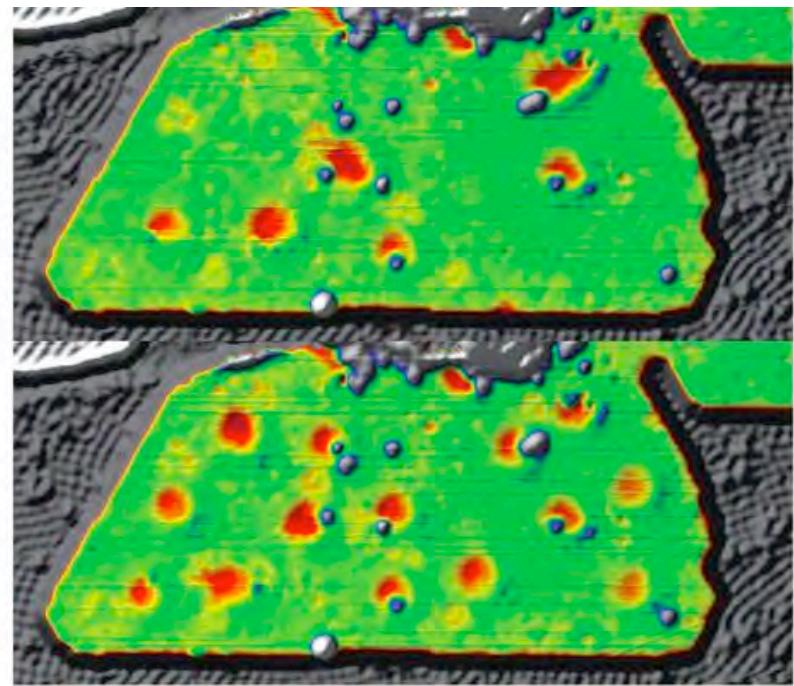
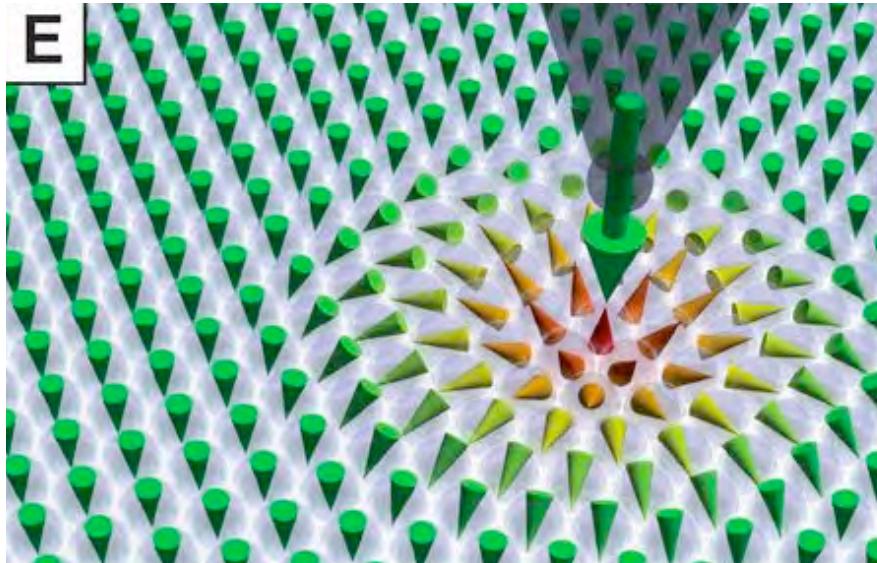


S. S. P. Parkin, M. Hayashi, and L. Thomas,  
Science 320, 190 (2008)

A. Fert *et al.*, Nature Nano. 8, 152 (2013)

# GENERATING INDIVIDUAL SKYRMIONS

Using spin-polarized scanning tunneling microscope



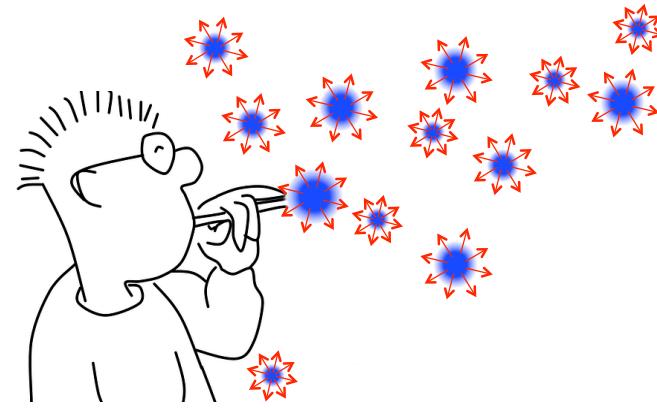
N. Romming, et al., Science 341, 636 (2013)

Spin-transfer torque switches skyrmion core reversibly

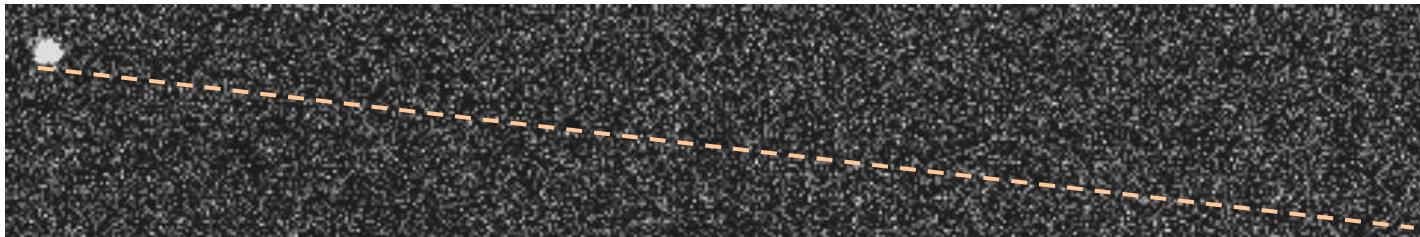
HOW TO  
EXPERIMENTALLY  
CREATE AND/OR MANIPULATE  
ROOM-TEMPERATURE  
MOBILE SKYRMIONS  
IN COMMON MATERIAL SYSTEMS.

# TOPICS

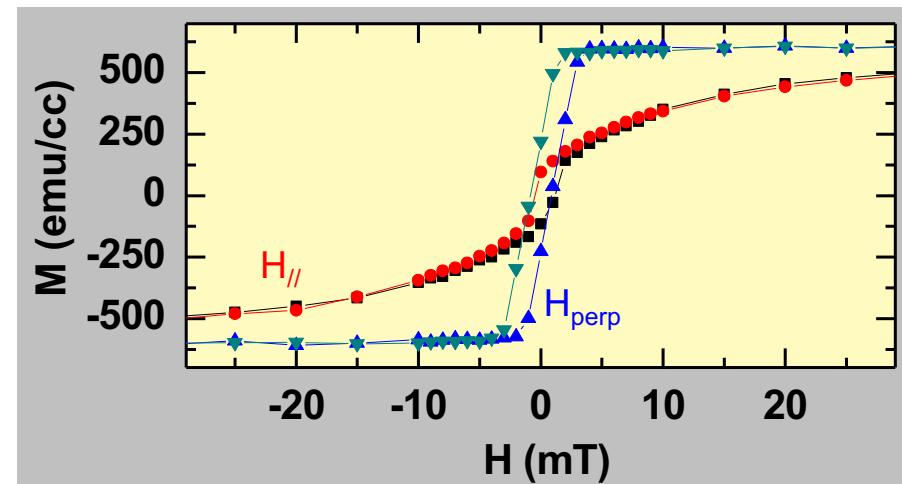
## Blowing Magnetic Skyrmion Bubbles



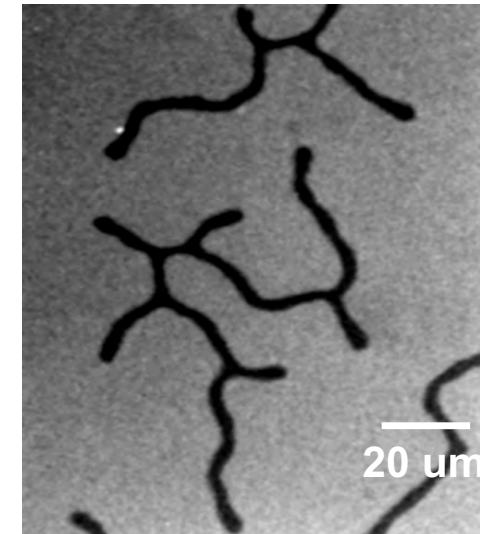
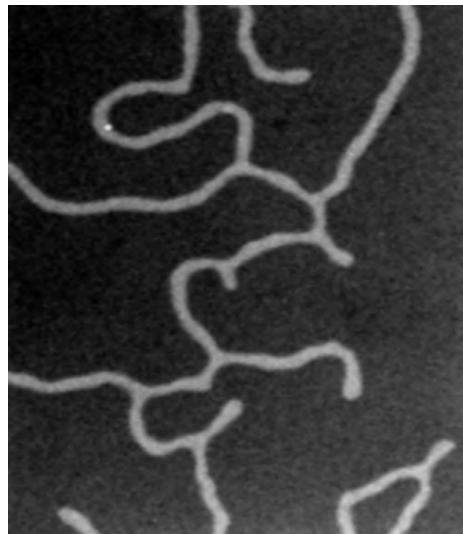
## Skyrmion Hall Effect



# FERROMAGNETIC FILMS WITH PERPENDICULAR ANISOTROPY

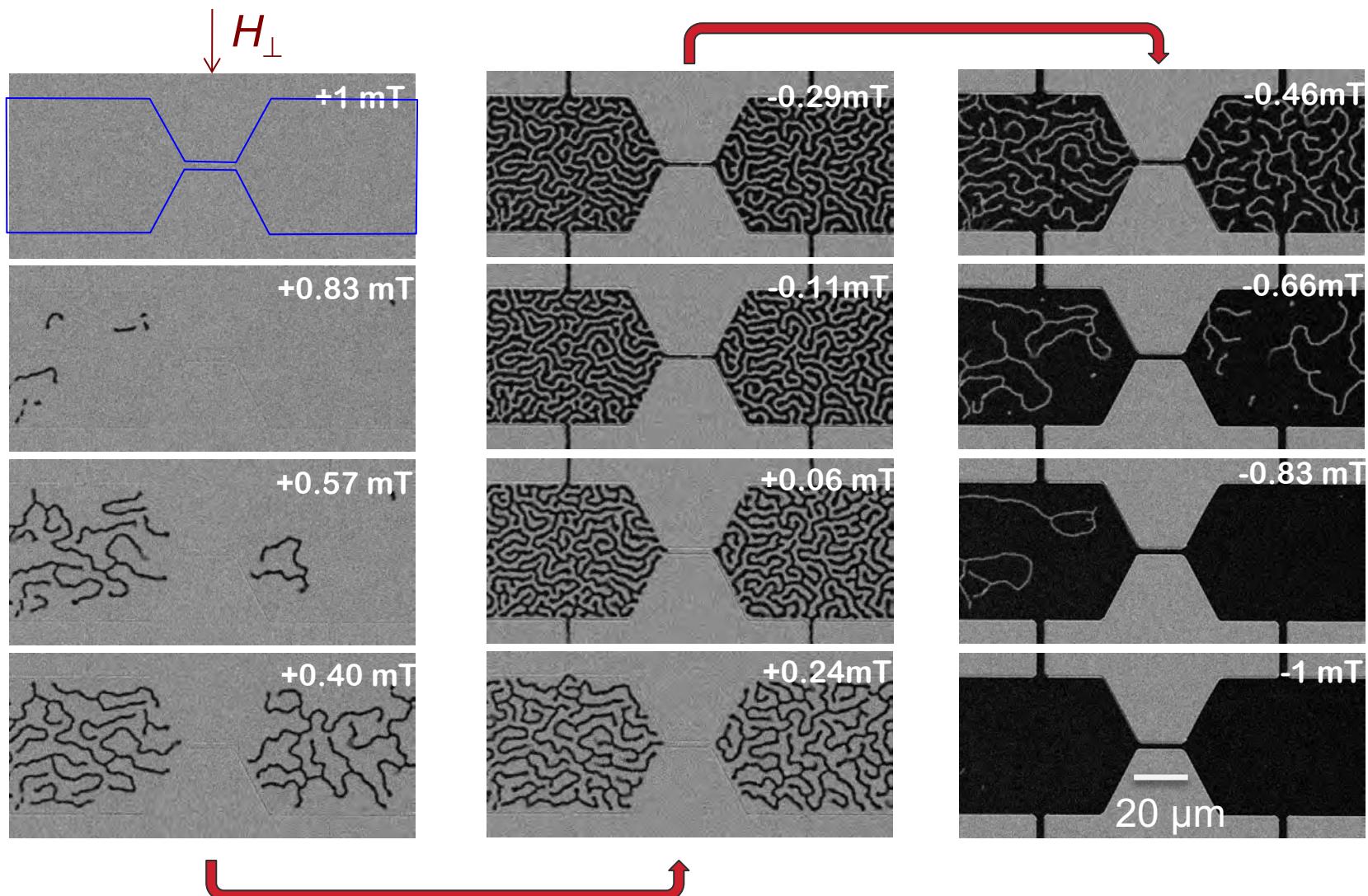


Polar  
Magneto  
Optic  
Kerr  
Effect  
(MOKE)  
Imaging



Increasing Magnetic Field

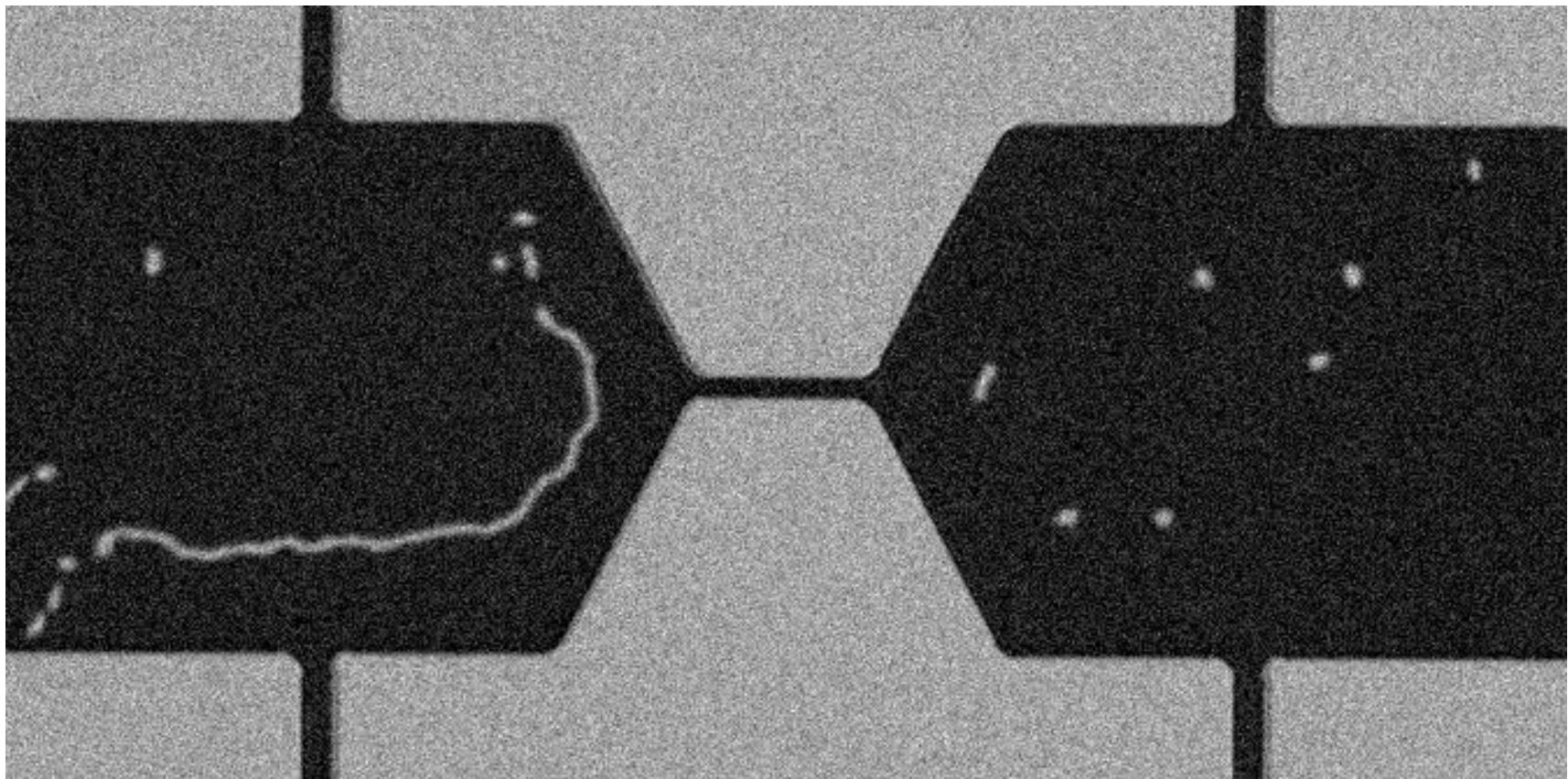
# DOMAIN EVOLUTION WITH MAGNETIC FIELD



# BLOWING MAGNETIC SKYRMION BUBBLES

$$H_{\perp} = -0.5 \text{ mT}$$

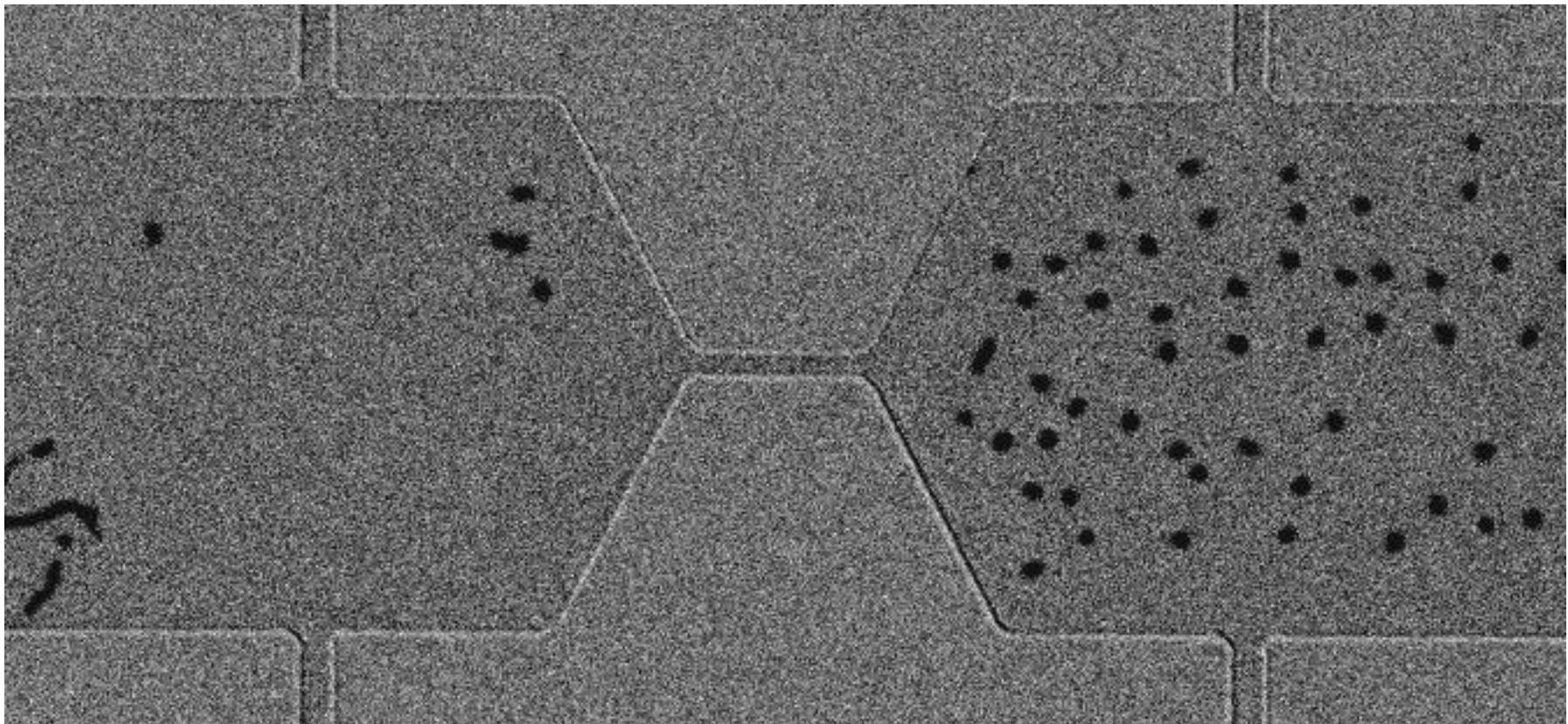
$$J_c = +0.5 \text{ MA/cm}^2 \quad \longrightarrow$$



# BLOWING MAGNETIC SKYRMION BUBBLES

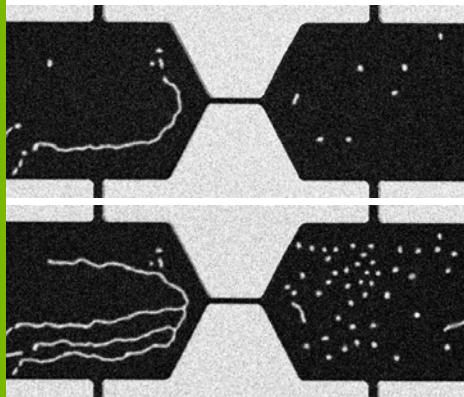
$H_{\perp} = +0.5 \text{ mT}$

$J_c = +0.5 \text{ MA/cm}^2$   $\longrightarrow$

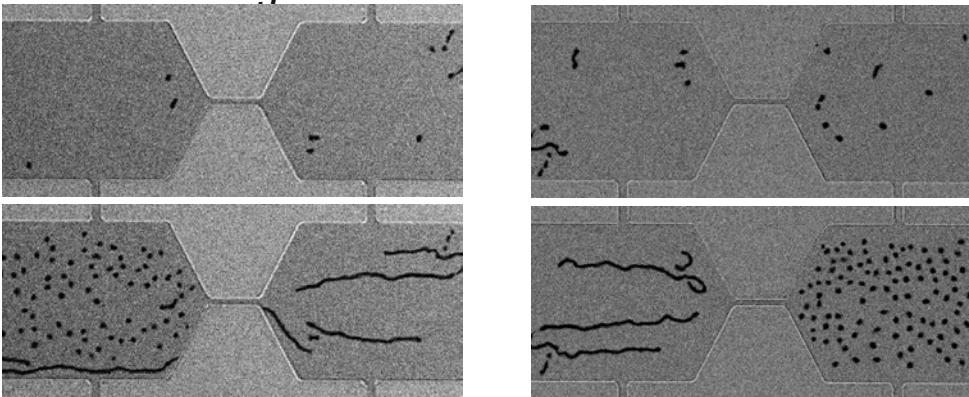


# SKYRMION GENERATION PHASE DIAGRAM

$H_{\perp} = -0.5 \text{ mT}$      $+J_c \longrightarrow$

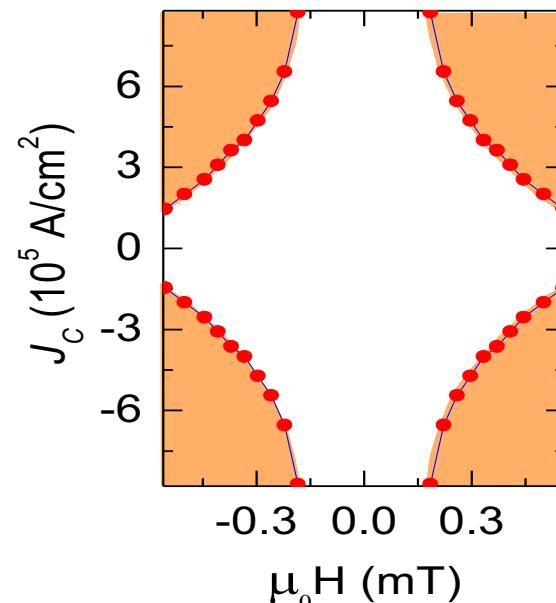


$H_{\perp} = +0.5 \text{ mT}$      $+J_c \longrightarrow$



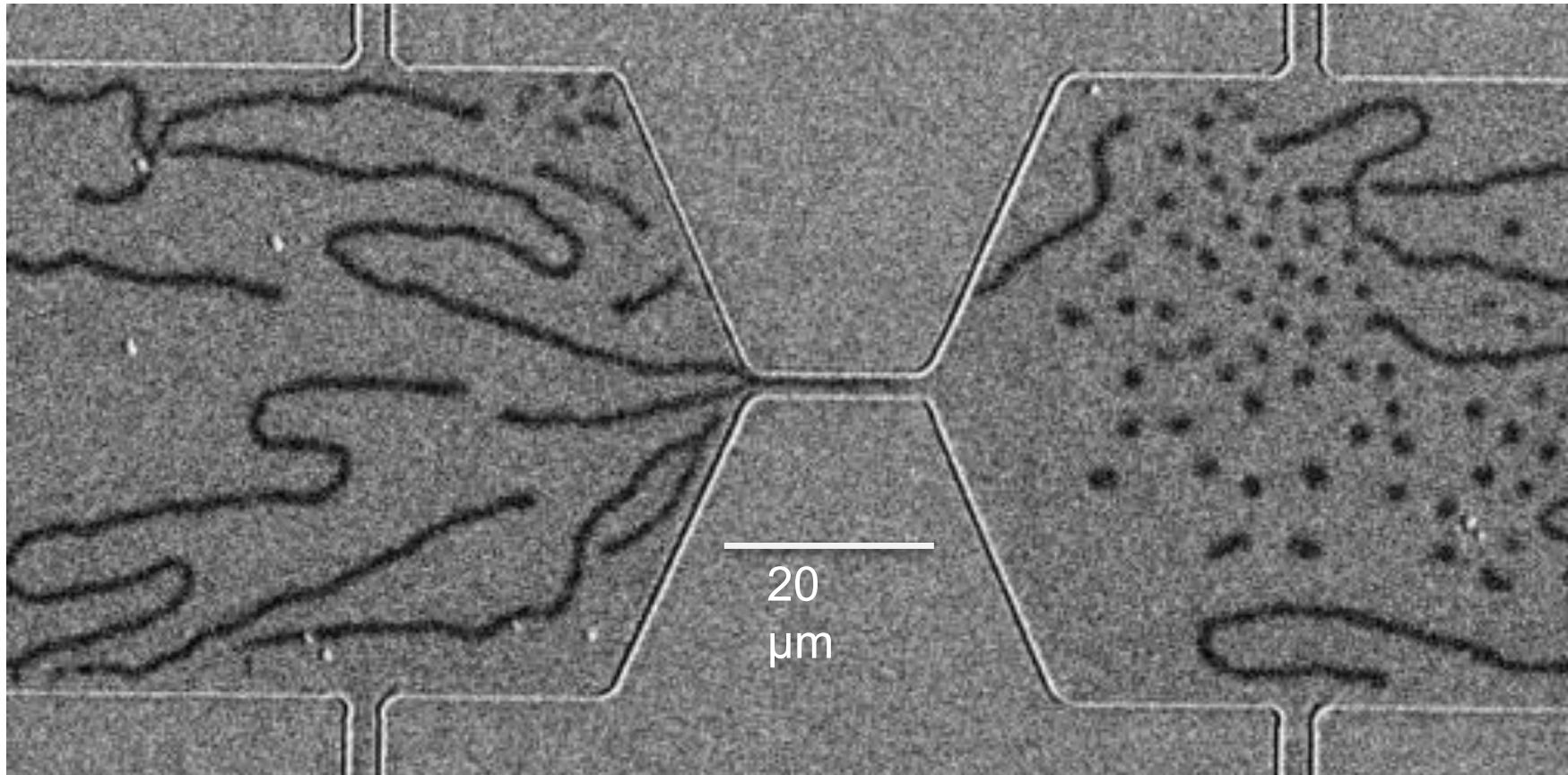
Before  
current pulse

After  
current pulse



# LOW CURRENT TRANSFORMATION DYNAMICS

$B_{\perp} = +0.46 \text{ mT}$ , DC current  $J_e = + 0.068 \text{ MA/cm}^2$

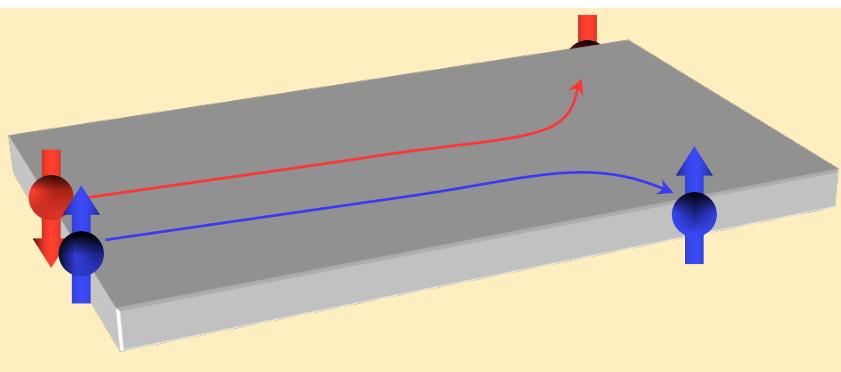


# SPIN HALL EFFECT

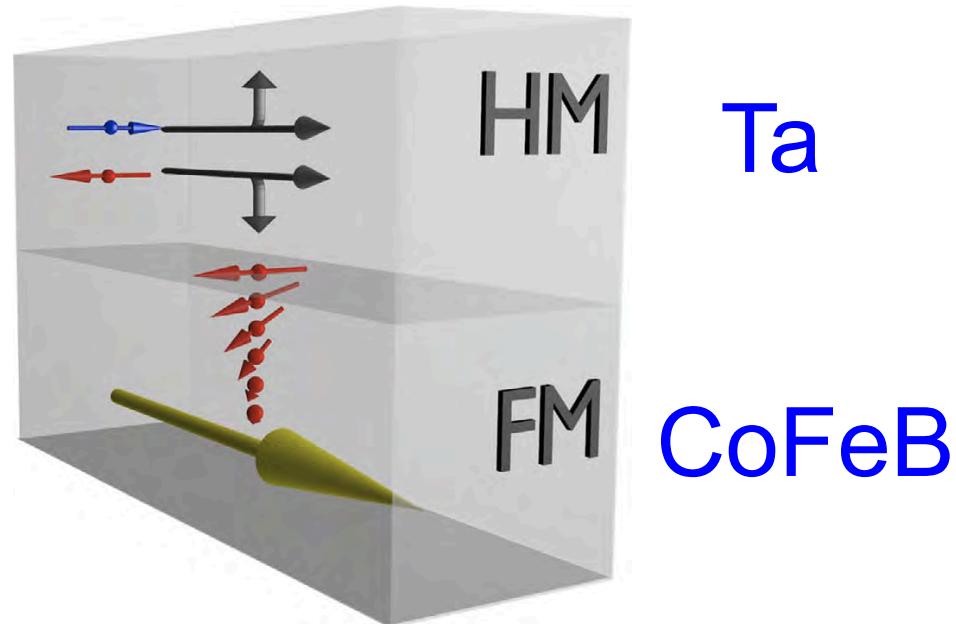
Spin Hall  
Charge Current



Transverse  
Spin Imbalance

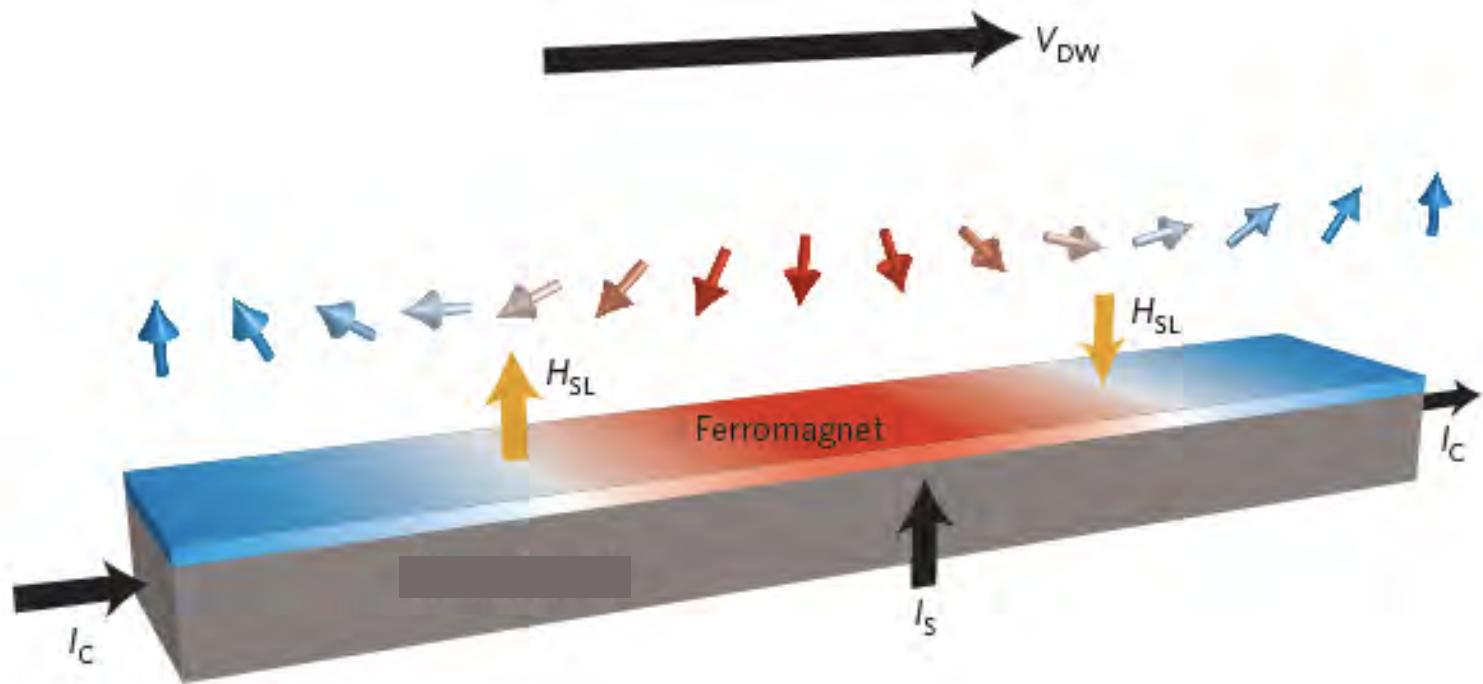


Spin Orbit Torque  
On Adjacent Ferromagnet



Recent Review: Axel Hoffmann,  
IEEE Trans. Magn. 49, 5172 (2013)

# CHIRAL DOMAIN WALL TORQUES

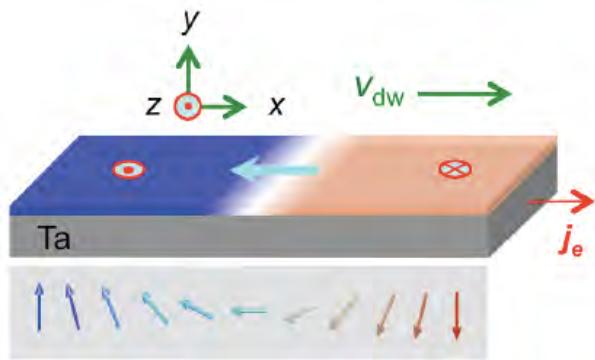


**DMI + Spin Hall  
= Efficient Torque**

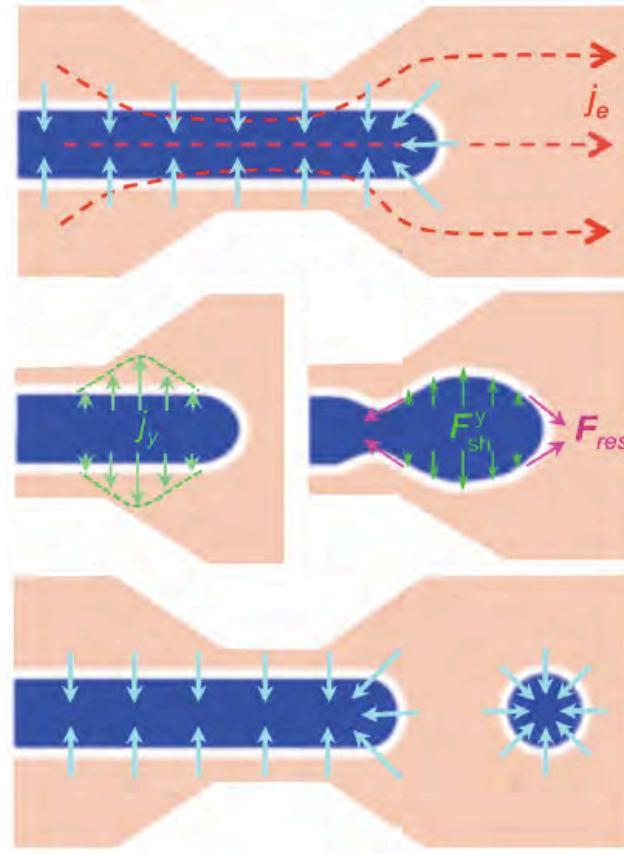
S. Emori, et al., Nat. Mater. **12**, 611 (2013)

K.-S. Ryu, et al., Nat. Nanotechn. **8**, 527 (2013)

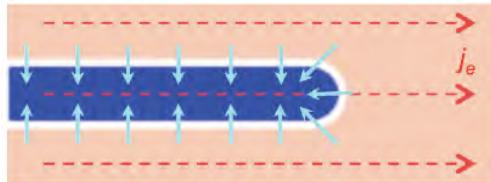
# INHOMOGENEOUS CHIRAL SPIN ORBIT TORQUES



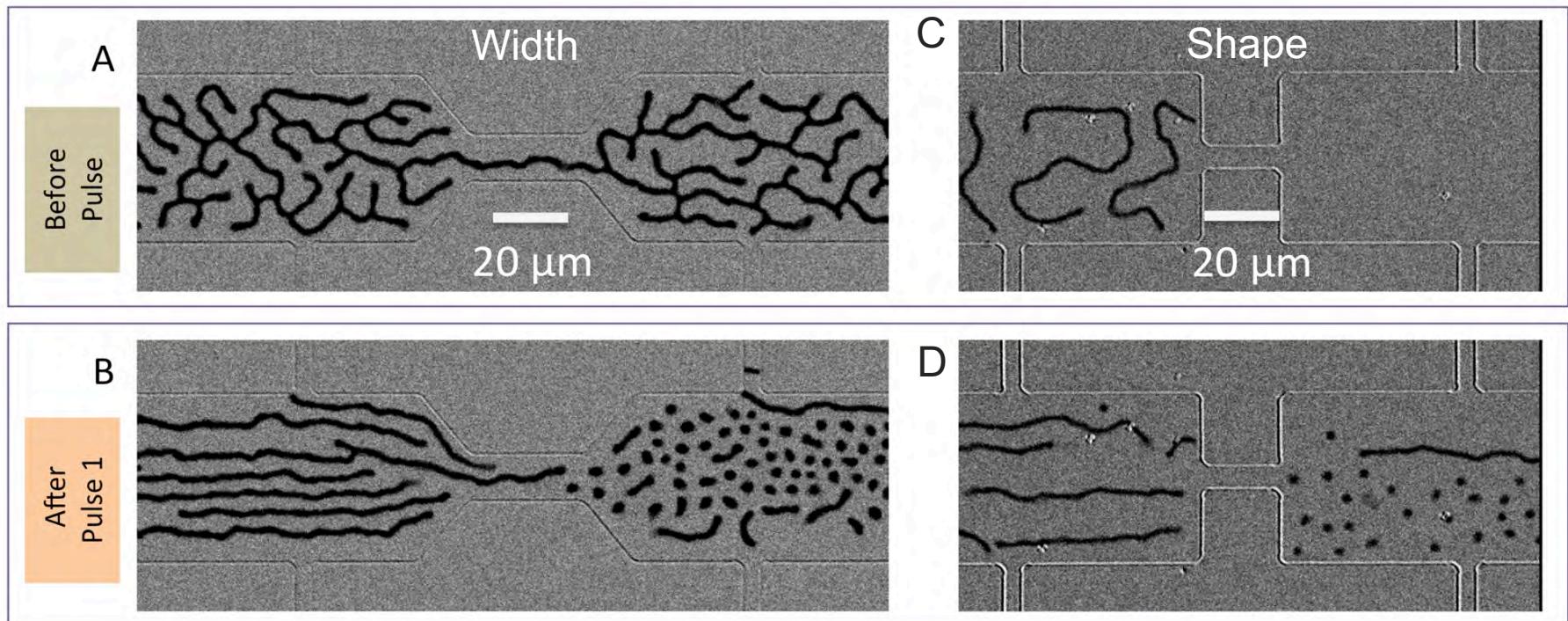
Stripe Domain with  
Inhomogeneous  
Current



Stripe Domain with  
Homogeneous Current



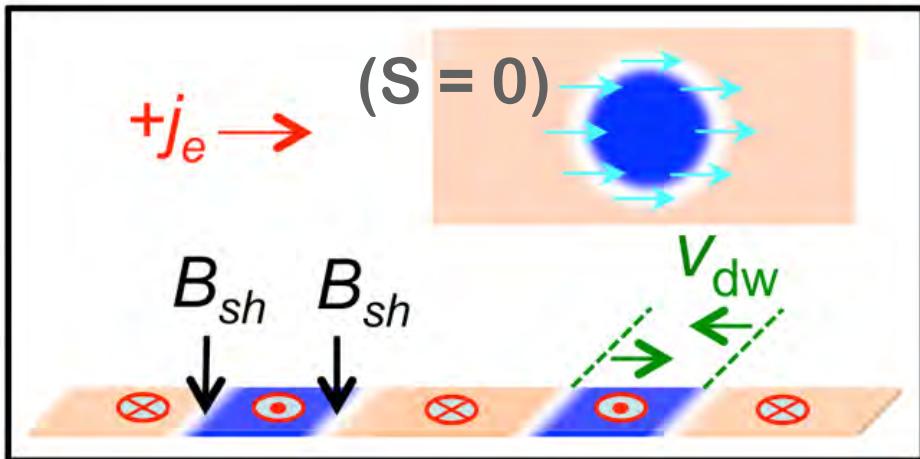
# DIFFERENT GEOMETRIES



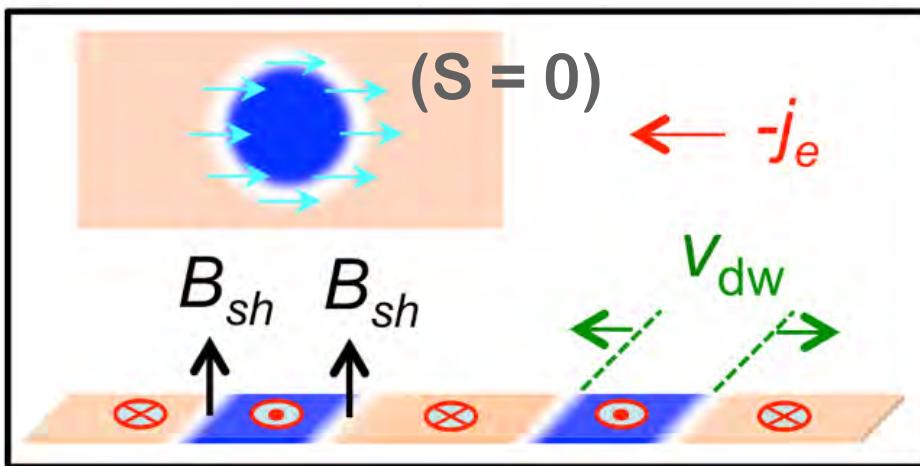
Skymion generation is robust  
Spatially divergent currents are the key

# TOPOLOGICALLY TRIVIAL BUBBLES

(S = 0) bubble stabilized by in-plane magnetic fields



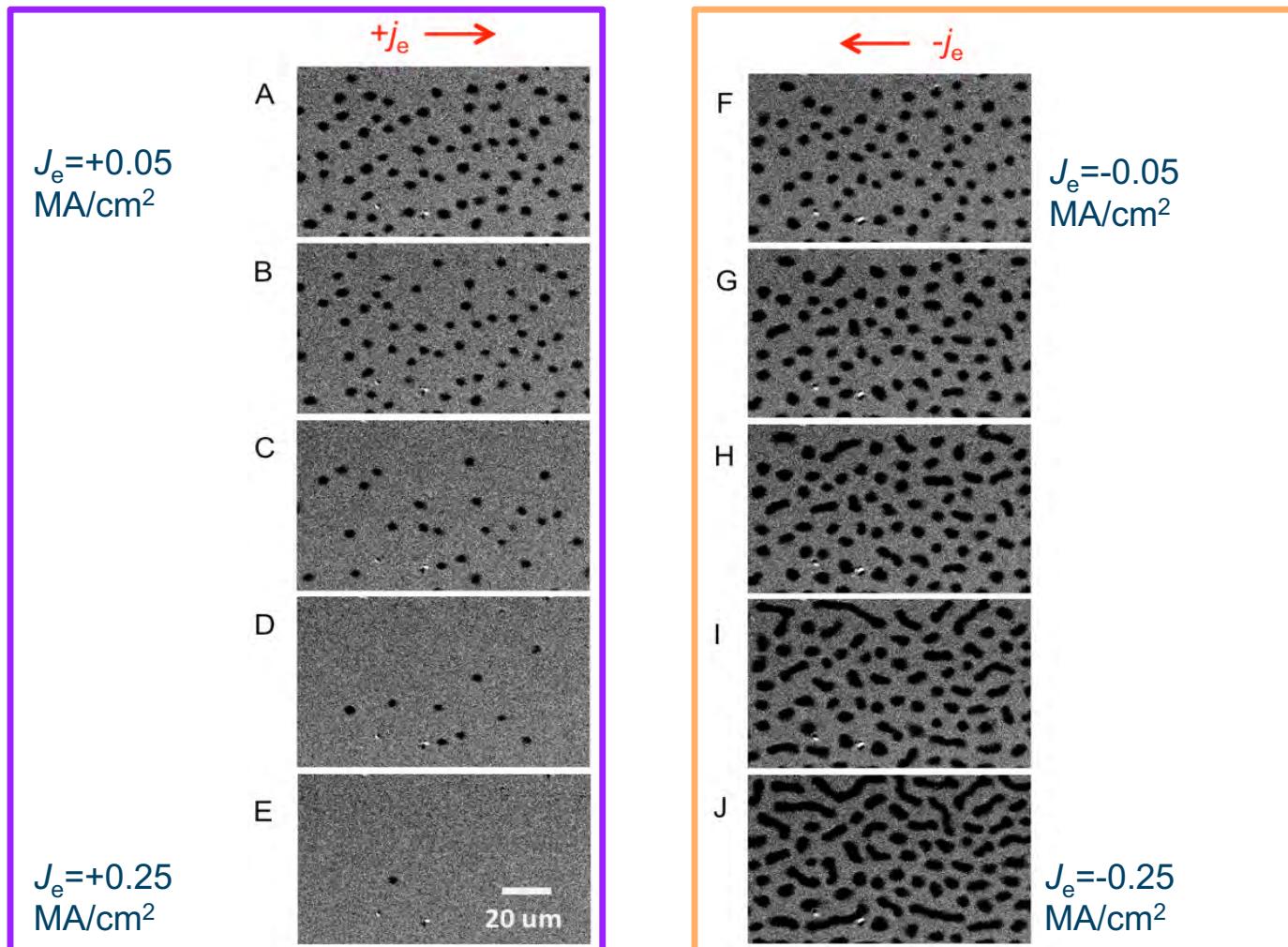
$$\vec{B}_{sh} = B_{sh}^0 (\hat{m} \times (\hat{z} \times \hat{j}_e))$$



- J. Choi, et al., PRL (2007)  
O. Lee et al., PRB (2014)  
S. Emori, et al., Nat. Mater. (2013)  
K.-S. Ryu, et al., Nat. Nano. (2013)  
S. Emori, et al., PRB. (2014)

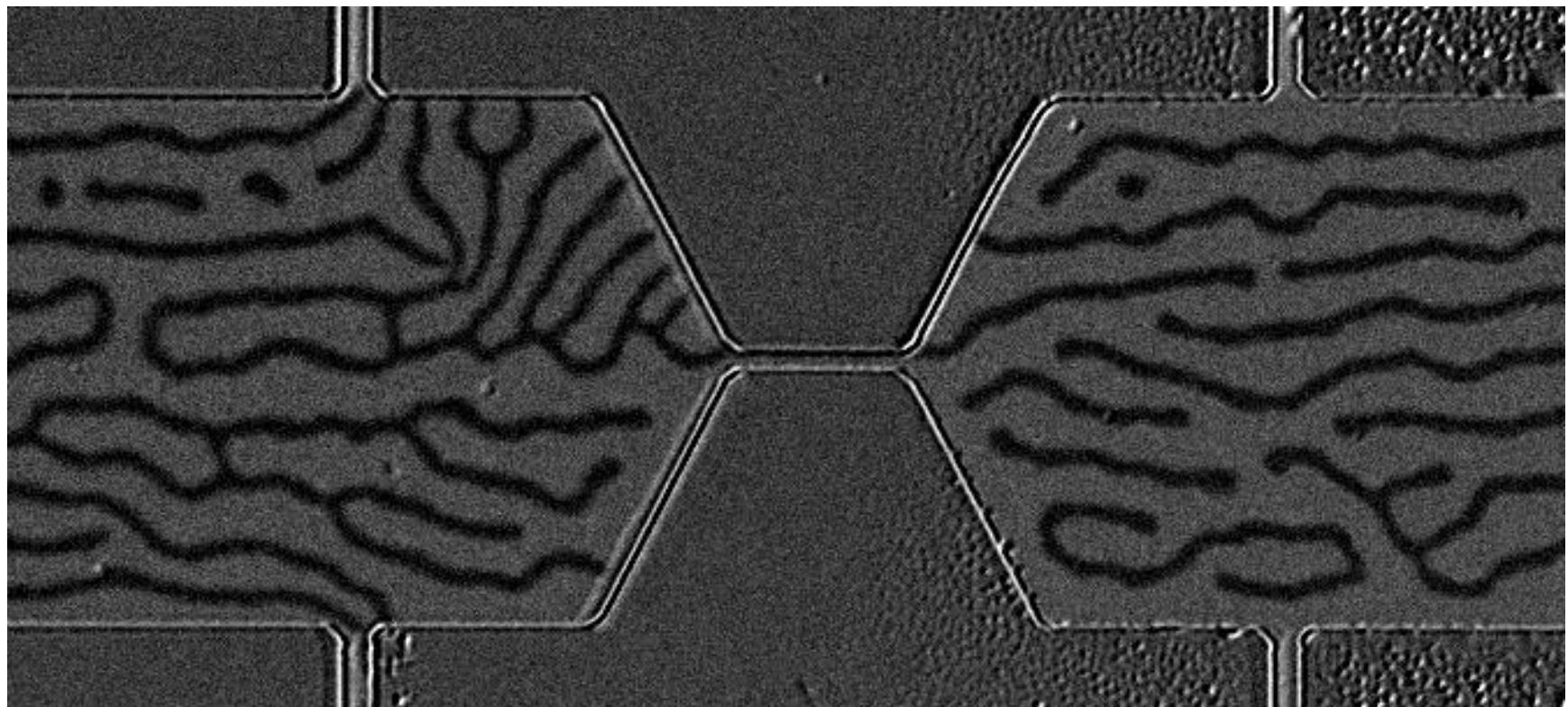
# RESPONSE TO CURRENTS WITH IN-PLANE MAGNETIC FIELDS

$B_{\parallel} = 10 \text{ mT}$

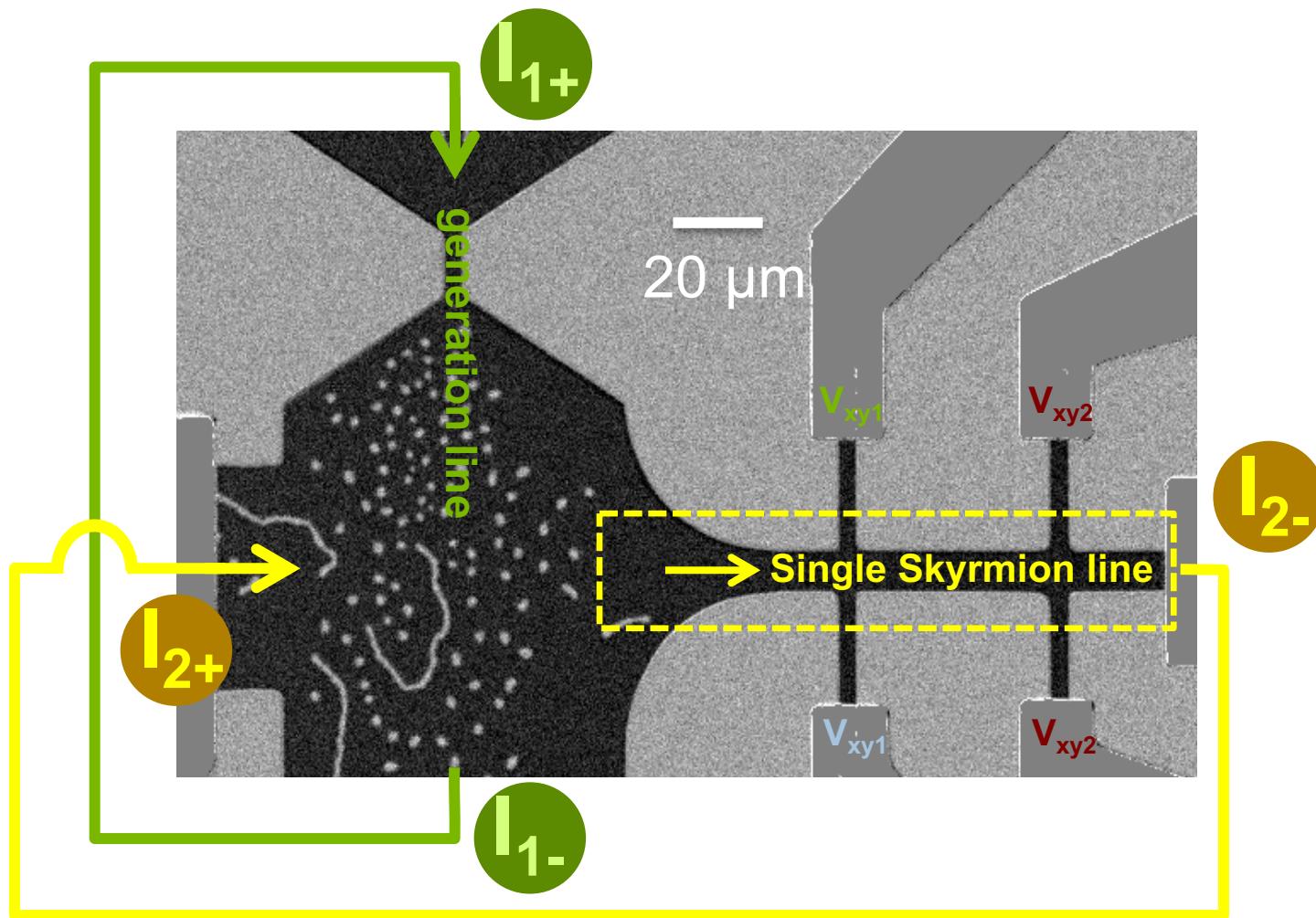


# NO BUBBLE FORMATION WITHOUT UNIFORM CHIRALITY

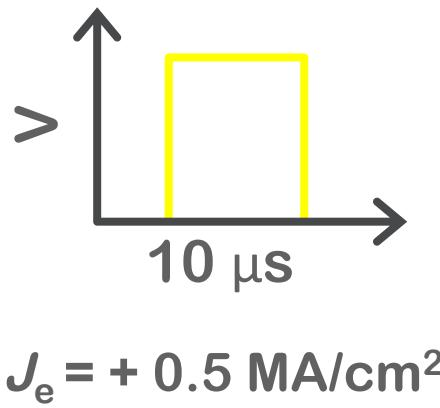
$$B_{\parallel} = 5 \text{ mT} \quad J_e = +0.5 \text{ MA/cm}^2$$



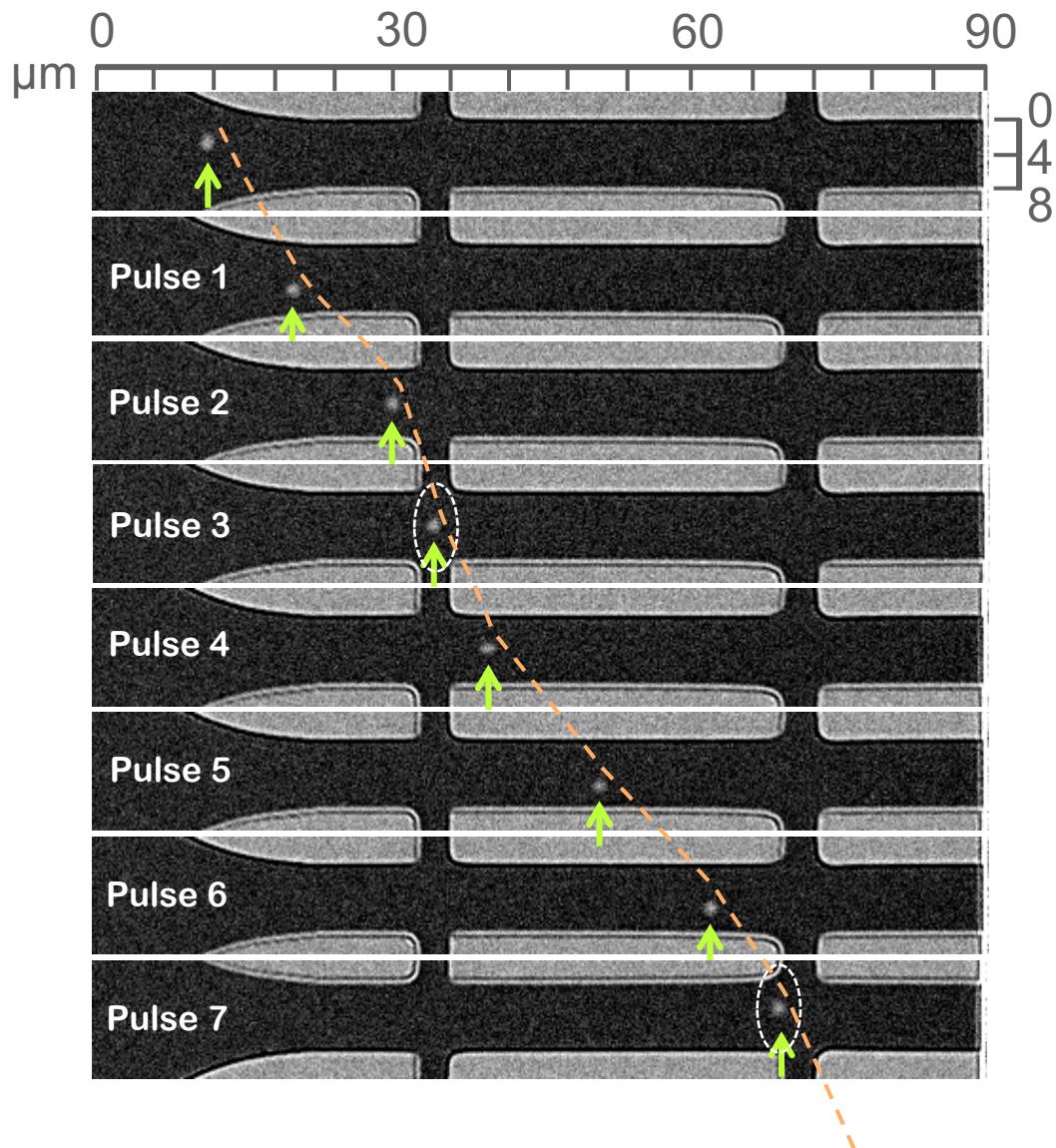
# INVESTIGATING INDIVIDUAL SKYRMIONS



# MOTION OF INDIVIDUAL SKYRMIONS

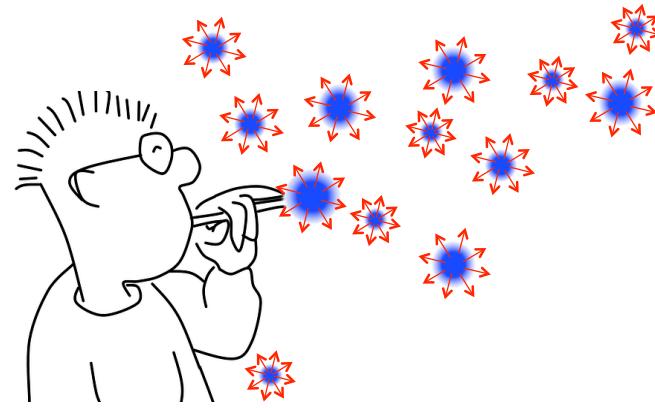


$$J_e = + 0.5 \text{ MA/cm}^2$$

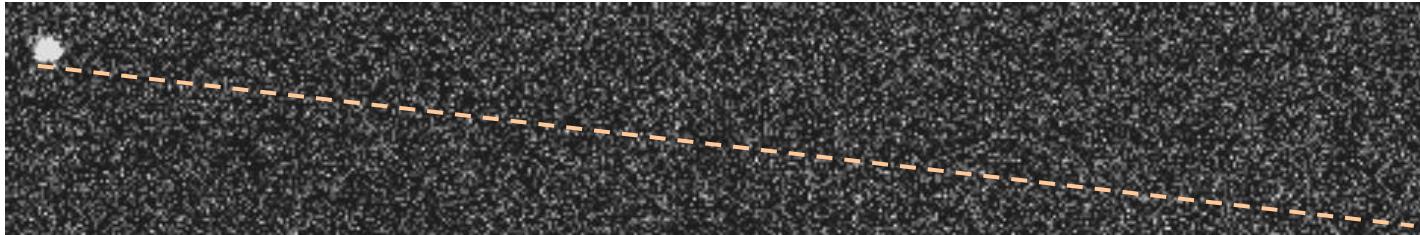


# TOPICS

## Blowing Magnetic Skyrmion Bubbles



## Skyrmion Hall Effect

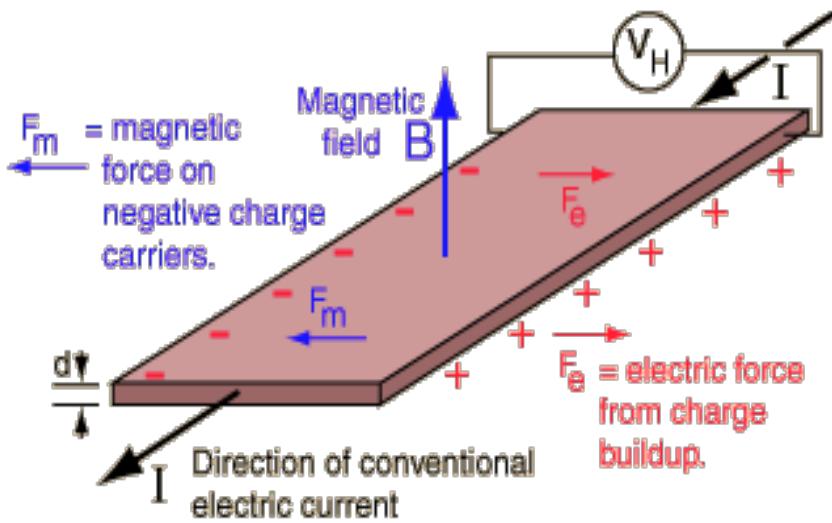


# SKYRMION HALL EFFECT

Classic Hall effect

*Electric charge  $q_e$*

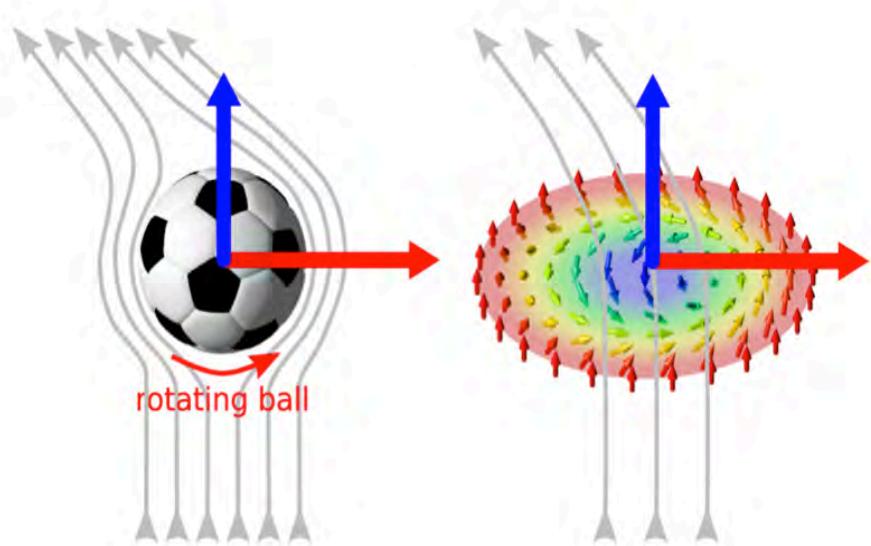
*Lorentz force  $q_e(v \times B)$*



Skymion Hall effect

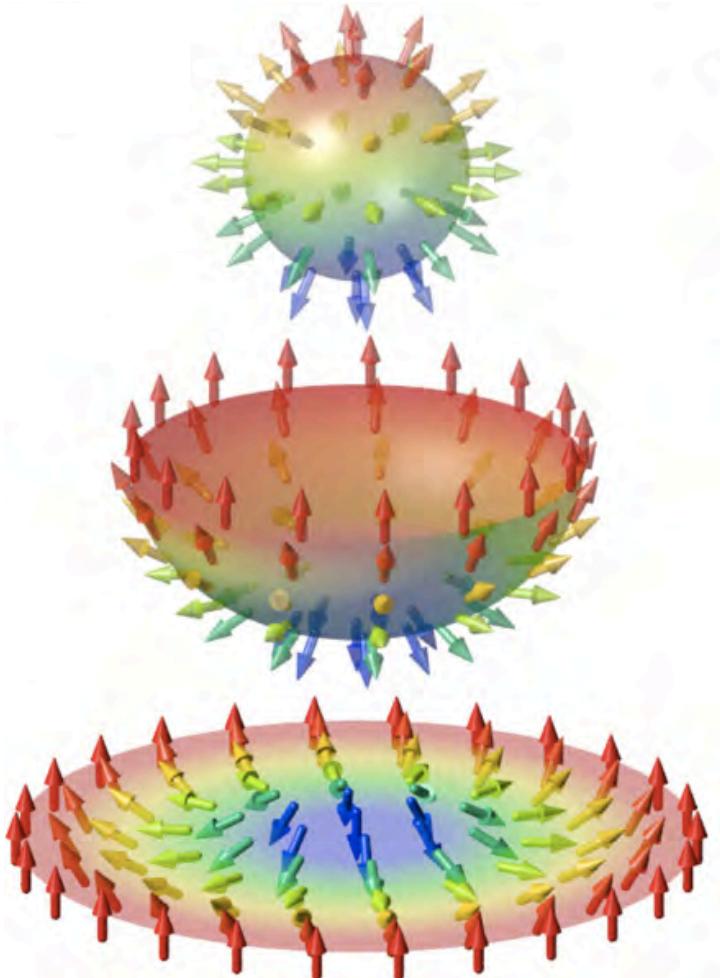
*Topological charge  $q_t$*

Magnus force  $4\pi q_t(v \times e_z)$



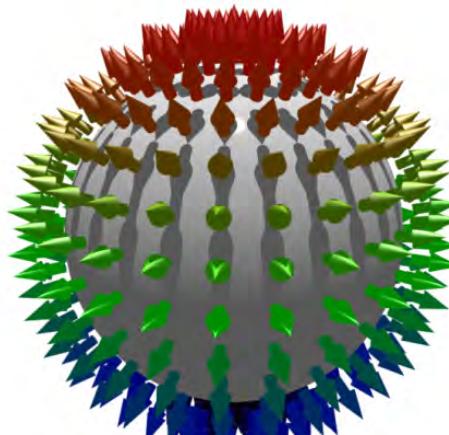
K. Everschor-Sitte and M. Sitte,  
J. Appl. Phys. **115**, 172602 (2014)

# TOPOLOGICAL CHARGE

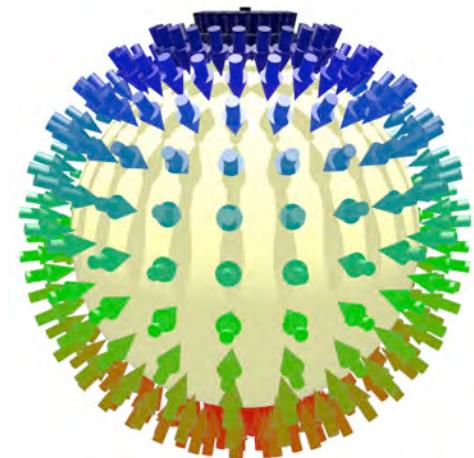


K. Everschor-Sitte and M. Sitte

Is there a Hall effect due  
to  
topological charge?



$$Q = -1$$

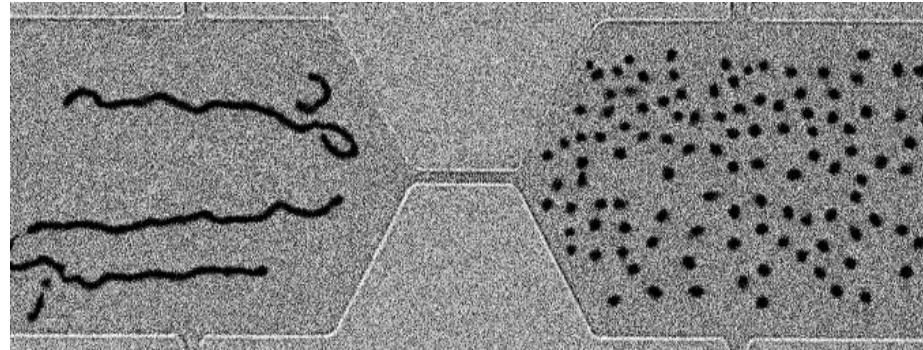
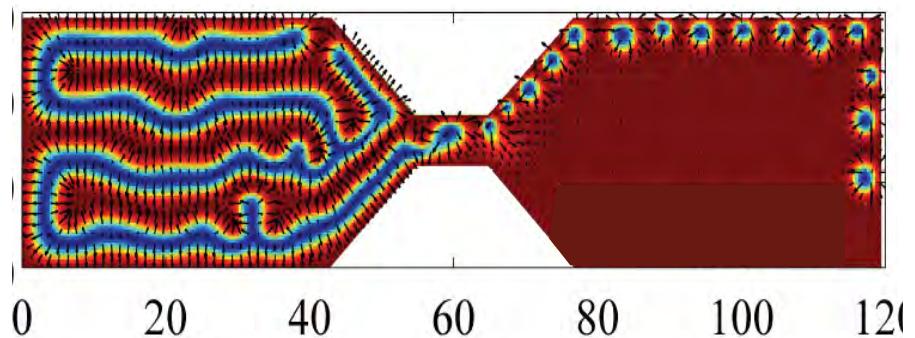


$$Q = +1$$

Electrons (-1) / holes (+1)

Shizeng Lin

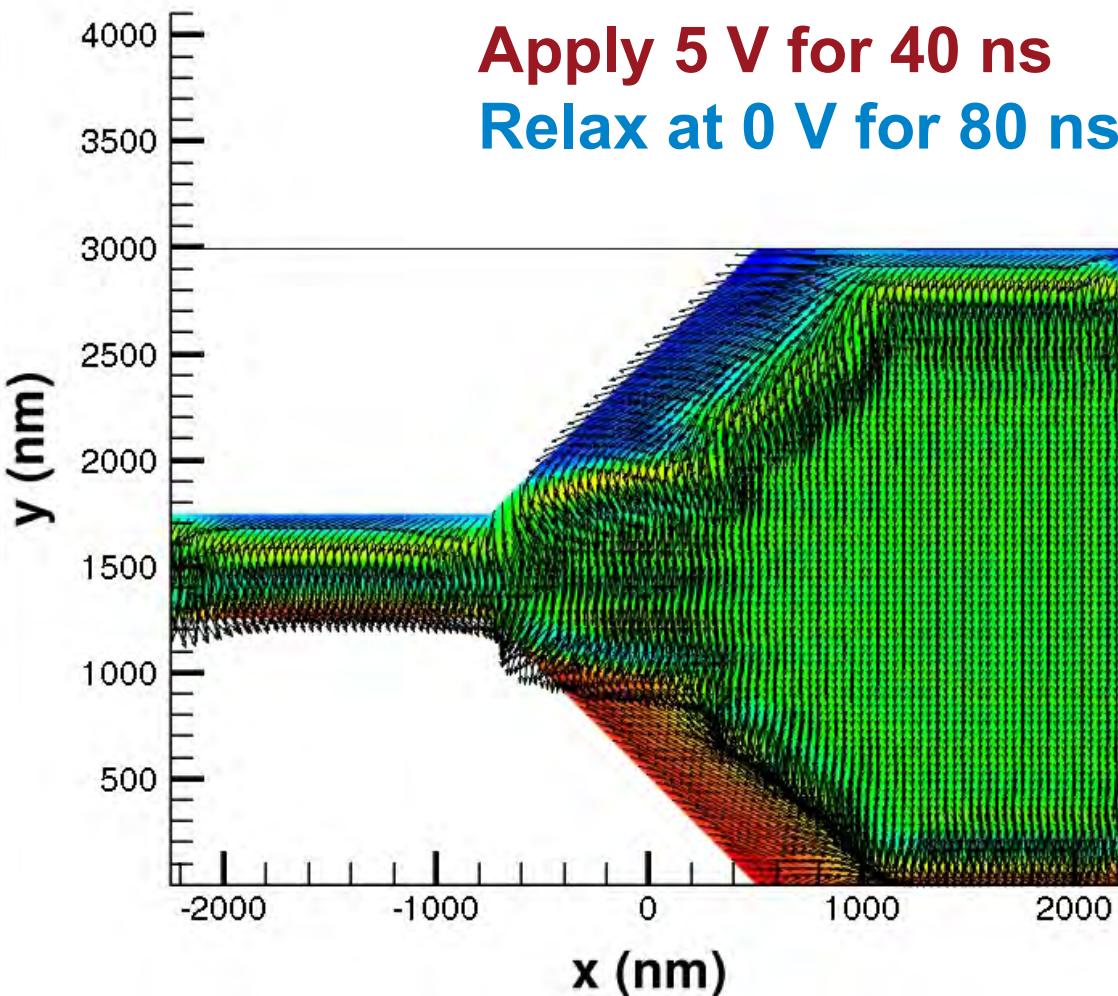
# Skyrmion Hall Effect



Shi-Zeng Lin, Phys. Rev. B **94**, 020402 (2016)

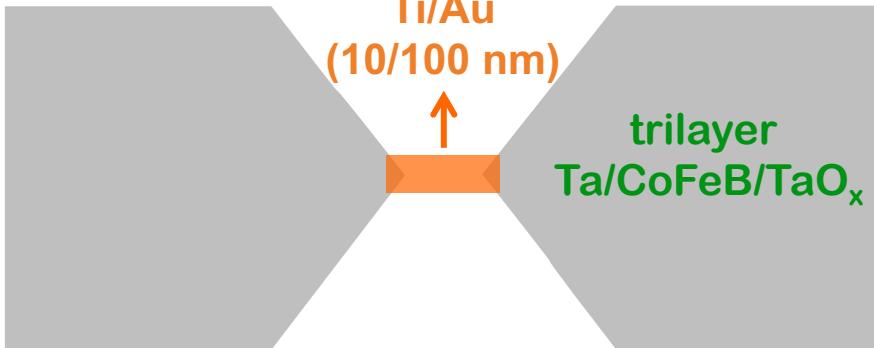
- Small  $J \rightarrow$  creep motion
- Bigger  $J \rightarrow$  steady motion (inaccessible)
- Complication from inhomogeneous currents

# MICROMAGNETIC SIMULATION OF TRANSFORMATION



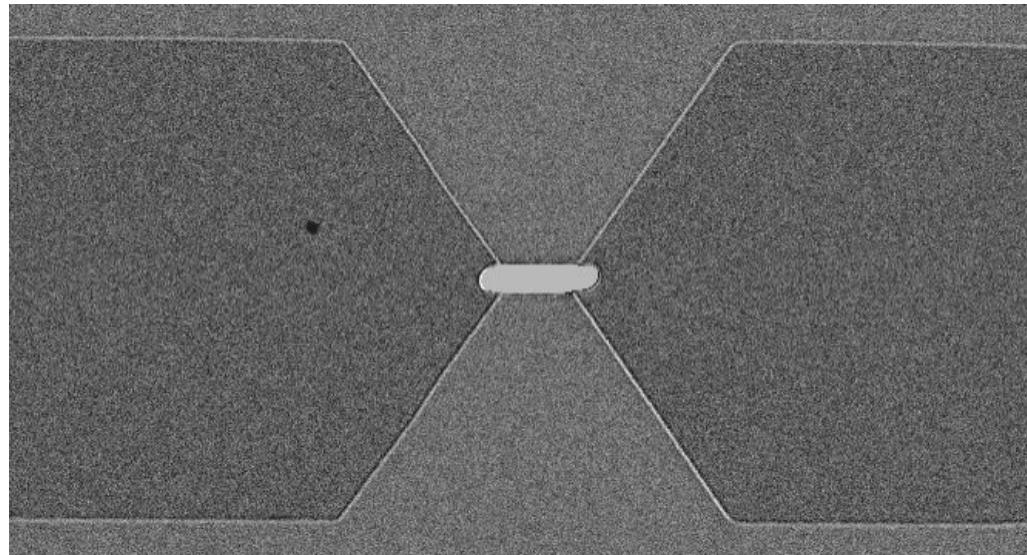
$H = 5 \text{ Oe}$   
 $M_s = 650 \text{ emu/cm}^3$   
 $H_a = 8868 \text{ Oe}$   
 $A = 3 \mu\text{erg/cm}$   
 $\text{DMI} = 0.5 \text{ erg/cm}^2$   
 $\alpha = 0.02$   
 $\sigma_{Ta} = 0.83 \text{ MS}$   
 $\theta_{sh} = 10\%$

# CREATION OF SKYRMIONS VIA NONMAGNETIC CONTACTS

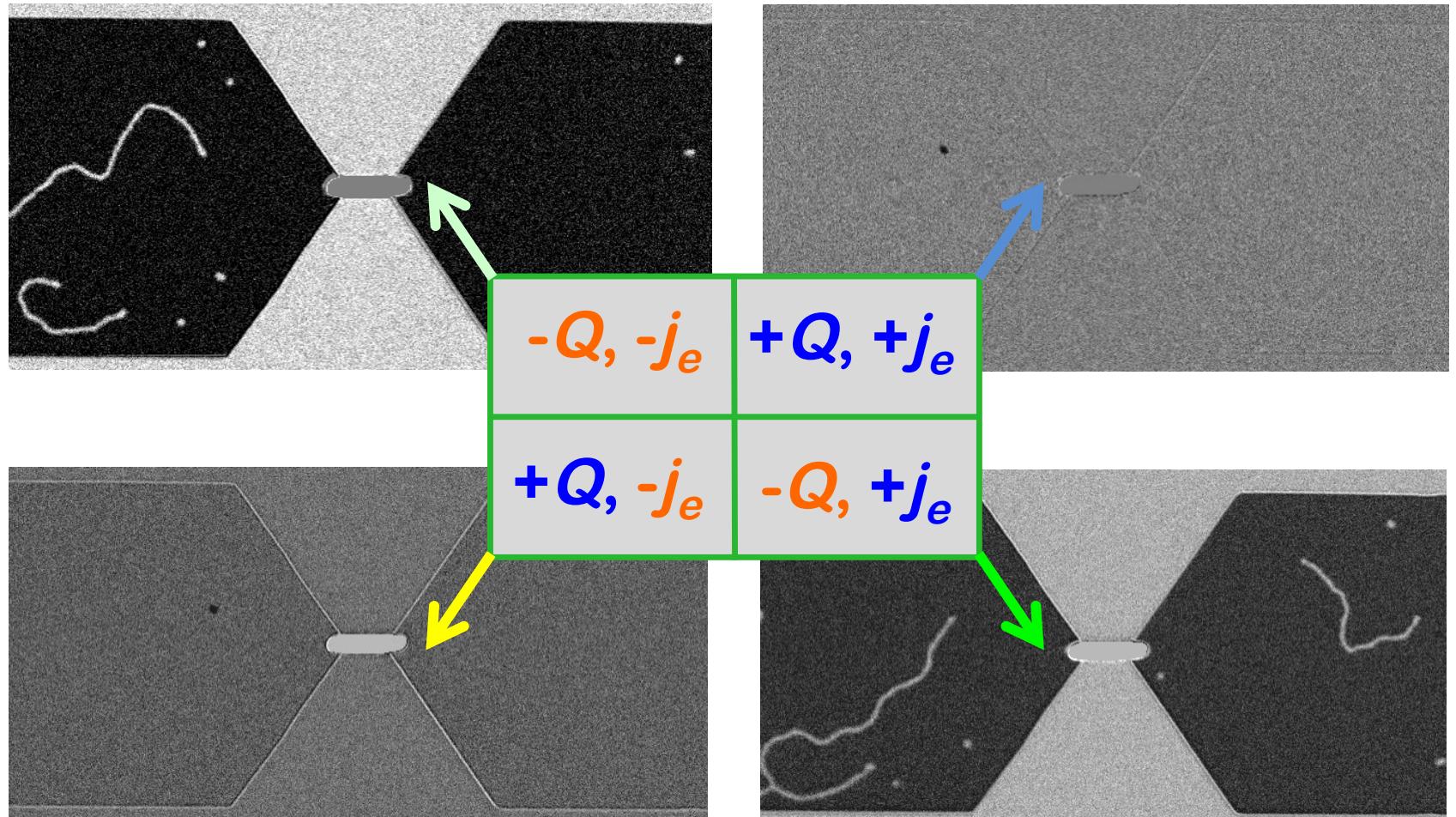


No stripe domains through  
No heat involved  
**Only divergence of current**  
Requires Larger Currents

Pulse current: 15 V of duration 1 ms at 1 Hz



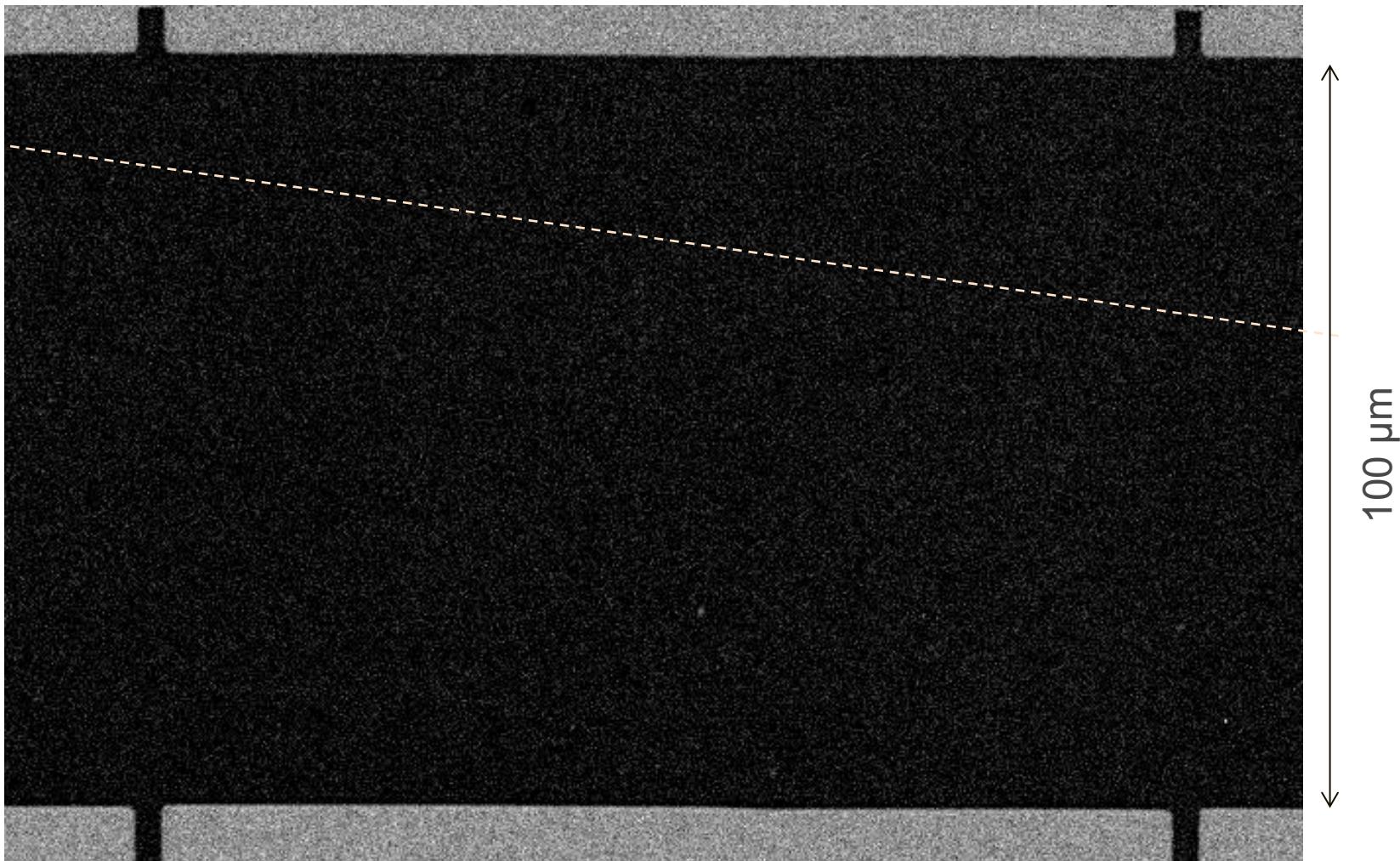
# Skyrmion Hall Effect



Topological charge behaves as electronic charge  
Skyrmion Hall effect in real space

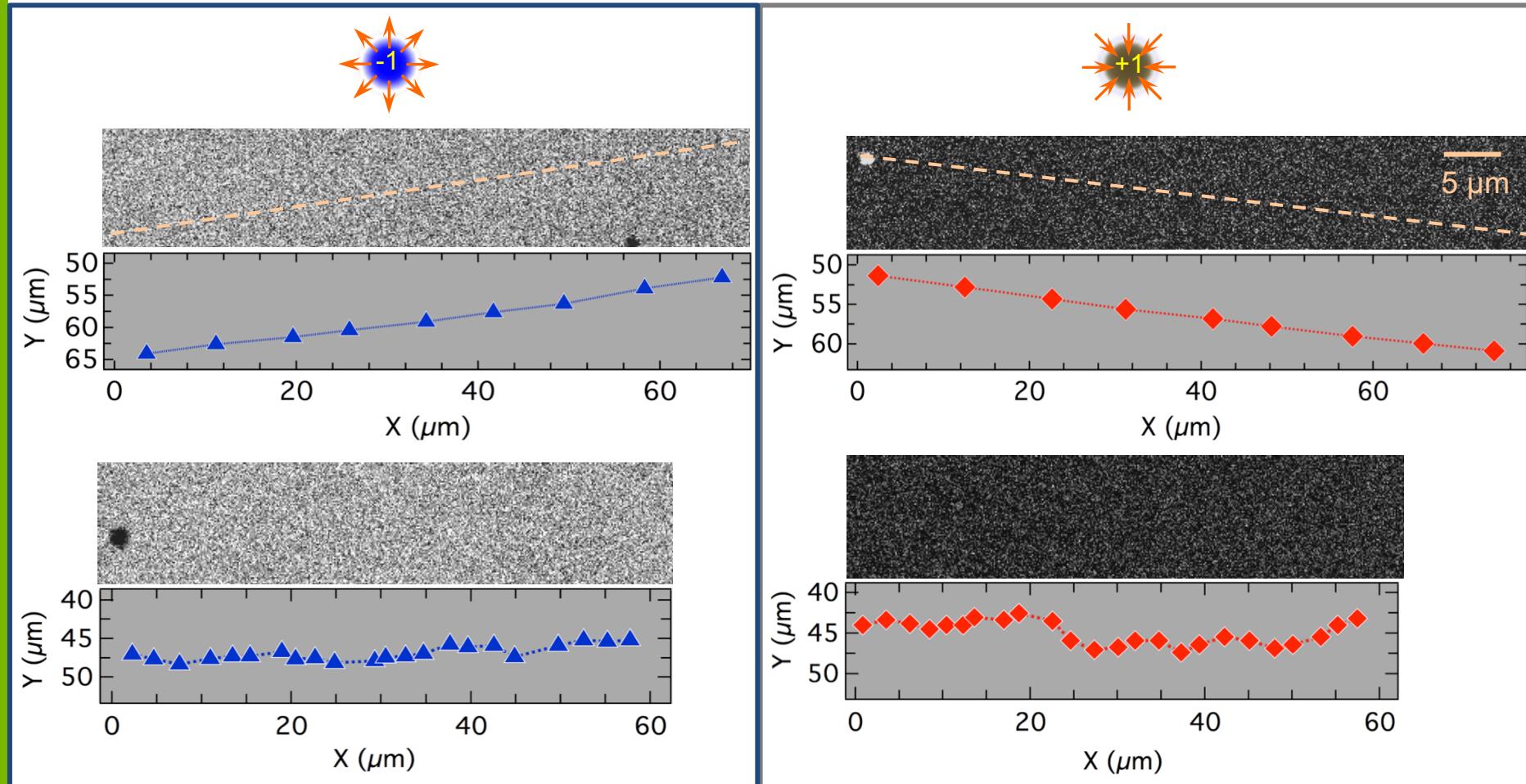
# SKYRMION MOTION WITH HOMOGENEOUS CURRENT

$$j_e = +2.8 \times 10^6 \text{ A/cm}^2$$



# CURRENT DEPENDENCE OF MOTION

$$j_e = +2.8 \times 10^6 \text{ A/cm}^2 \rightarrow$$



$$j_e = +1.3 \times 10^6 \text{ A/cm}^2 \rightarrow$$

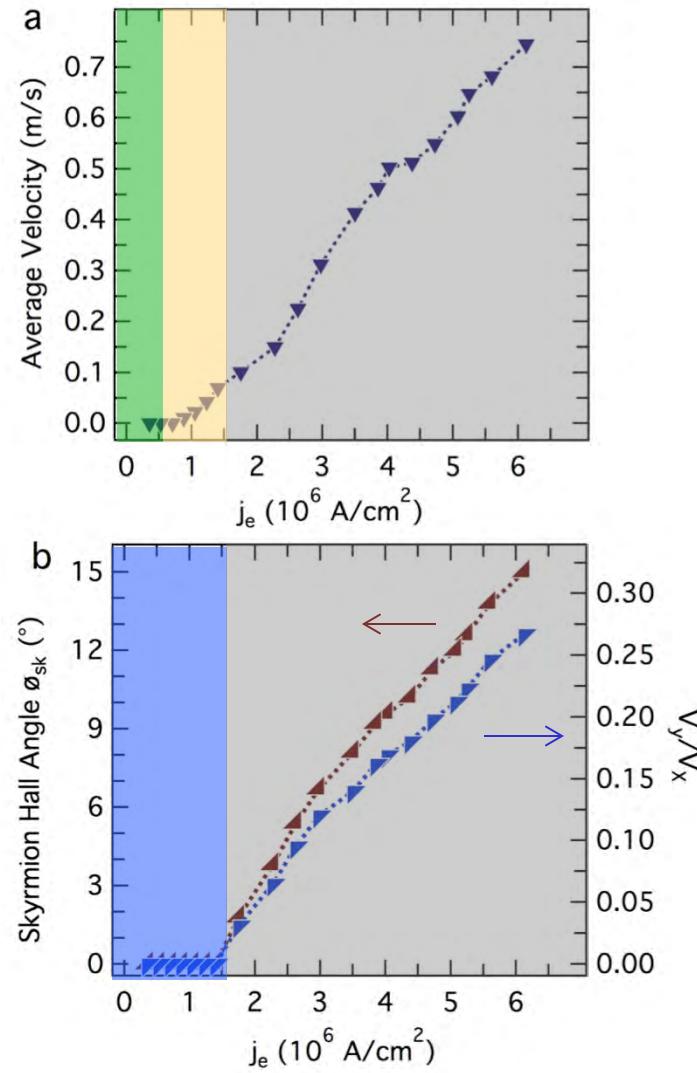
# EVOLUTION OF VELOCITY/HALL ANGLE ON CURRENTS

Skymion is pinned  
(stochastic) hopping

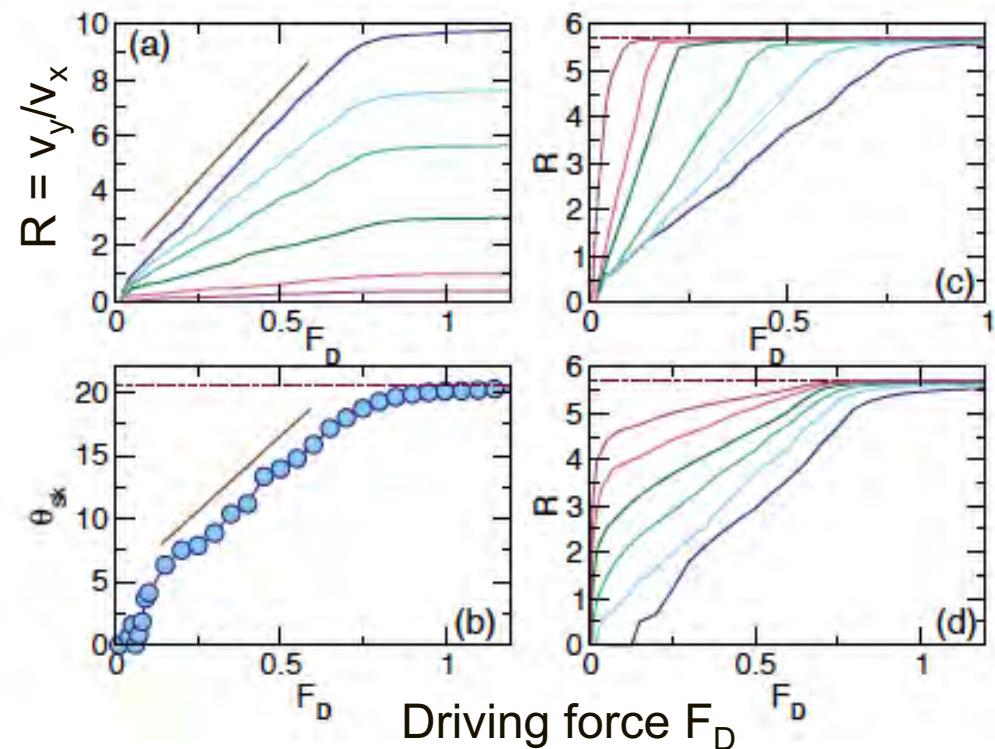
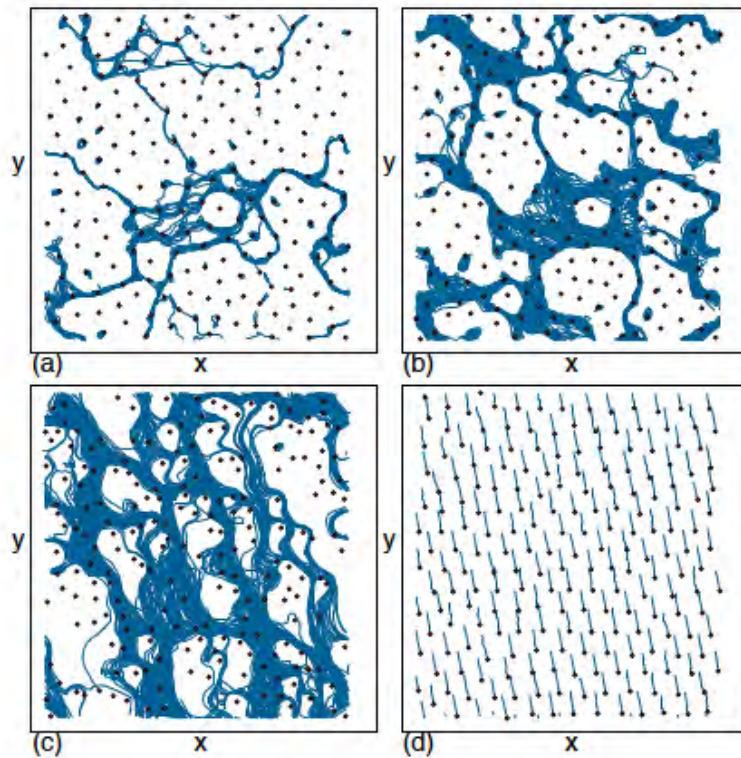
Drive-dependent  
Skymion Hall angle

rather than predicted  
constant value

$$\frac{v_y}{v_x} = \tan(\Phi_{sk}) = \frac{1}{\alpha D}$$

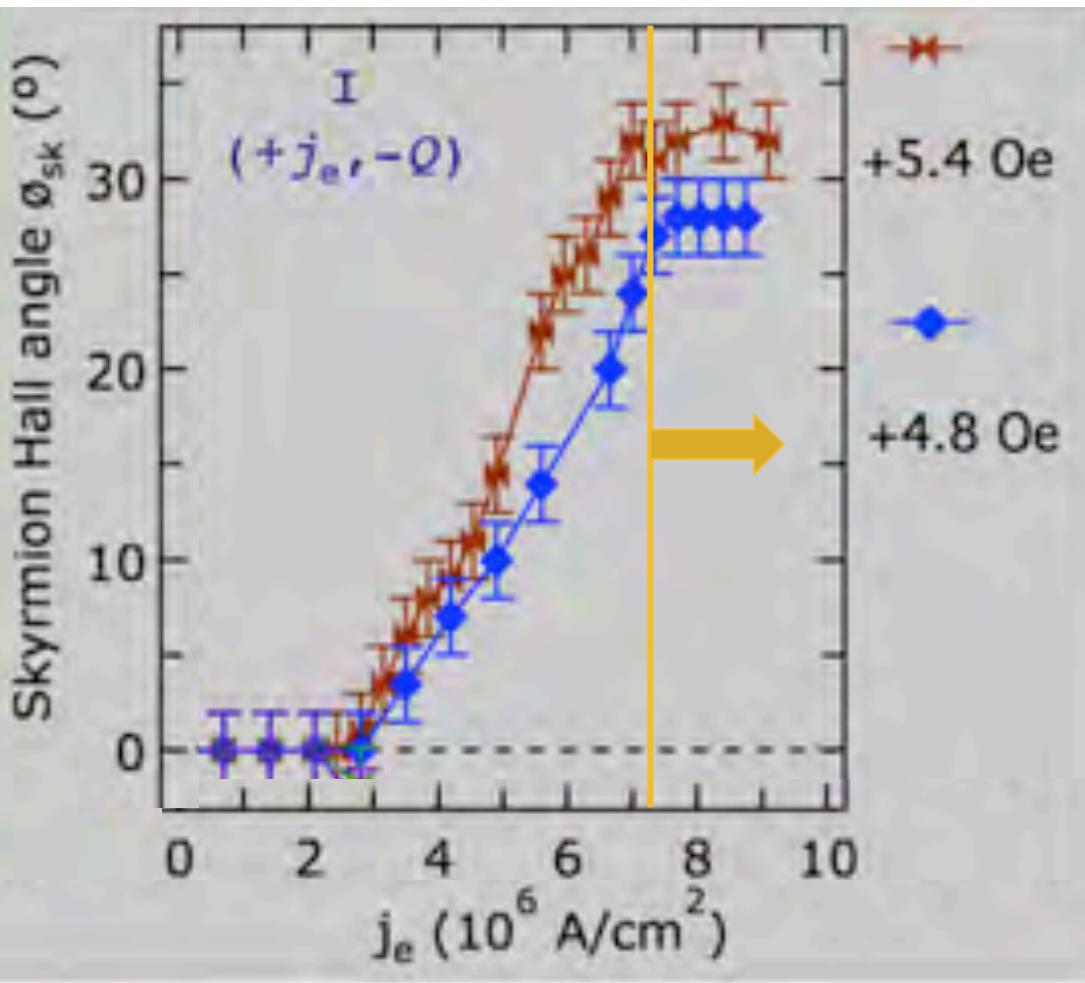


# PARTICLE-BASED SIMULATION MODEL



- For small intrinsic Hall angles, the Hall angle increases linearly with external drive.
- At sufficiently high drives, the skyrmions enter a free flow regime,  
**=>  $R$  saturates to the disorder-free limit.**

# DRIVE-DEPENDENT SKYRMION HALL ANGLE



Skyrmion Hall angle eventually increases

Flow regime is reached

$H = 5.4 \text{ Oe}, d = 800 \pm 300 \text{ nm}$

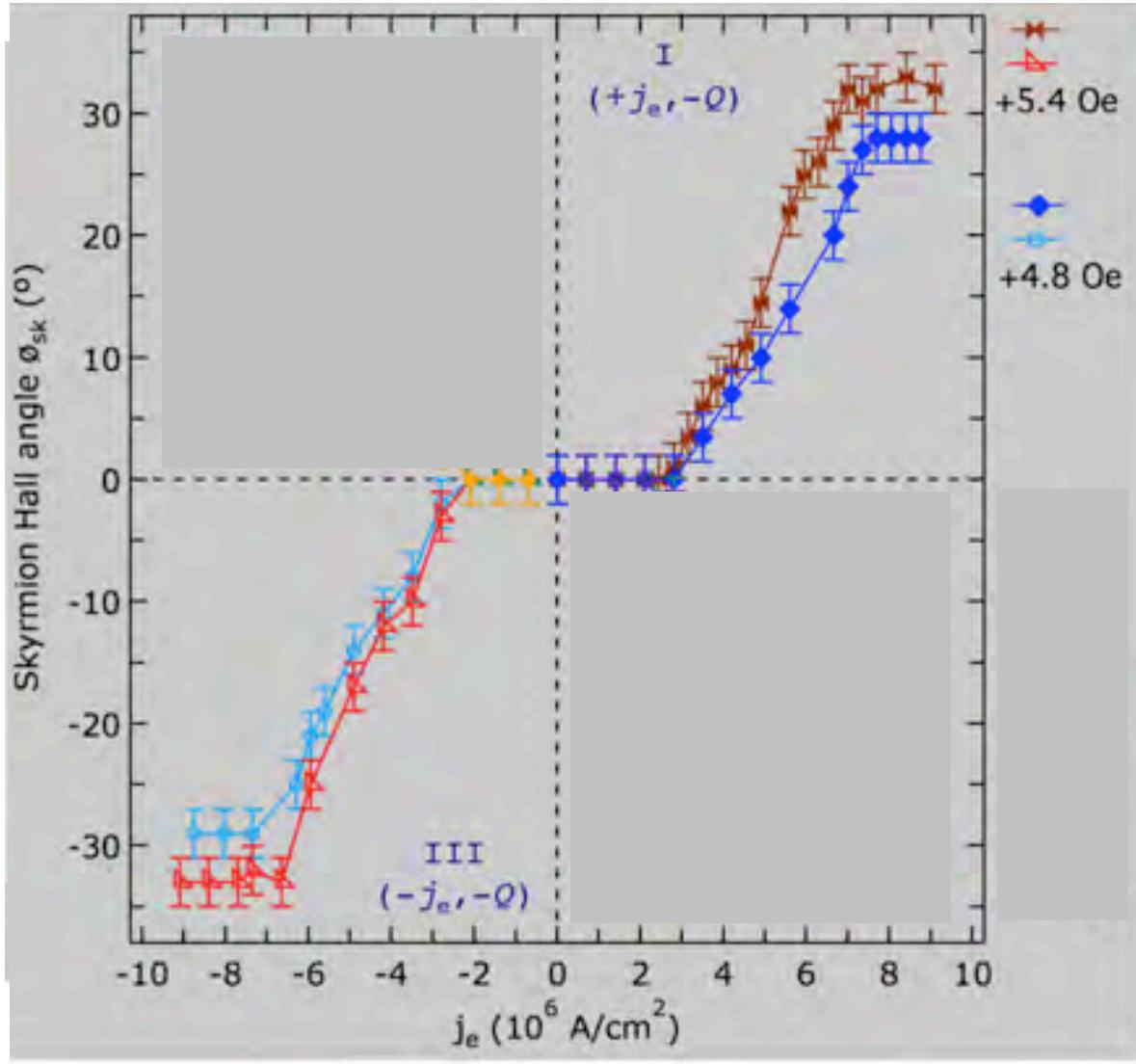
$$\Phi_{sk} = 32 \pm 2^\circ$$

$H = 4.8 \text{ Oe}, d = 1100 \pm 300 \text{ nm}$

$$\Phi_{sk} = 28 \pm 2^\circ$$

Final Skyrmion Hall angle depends on the skyrmion size

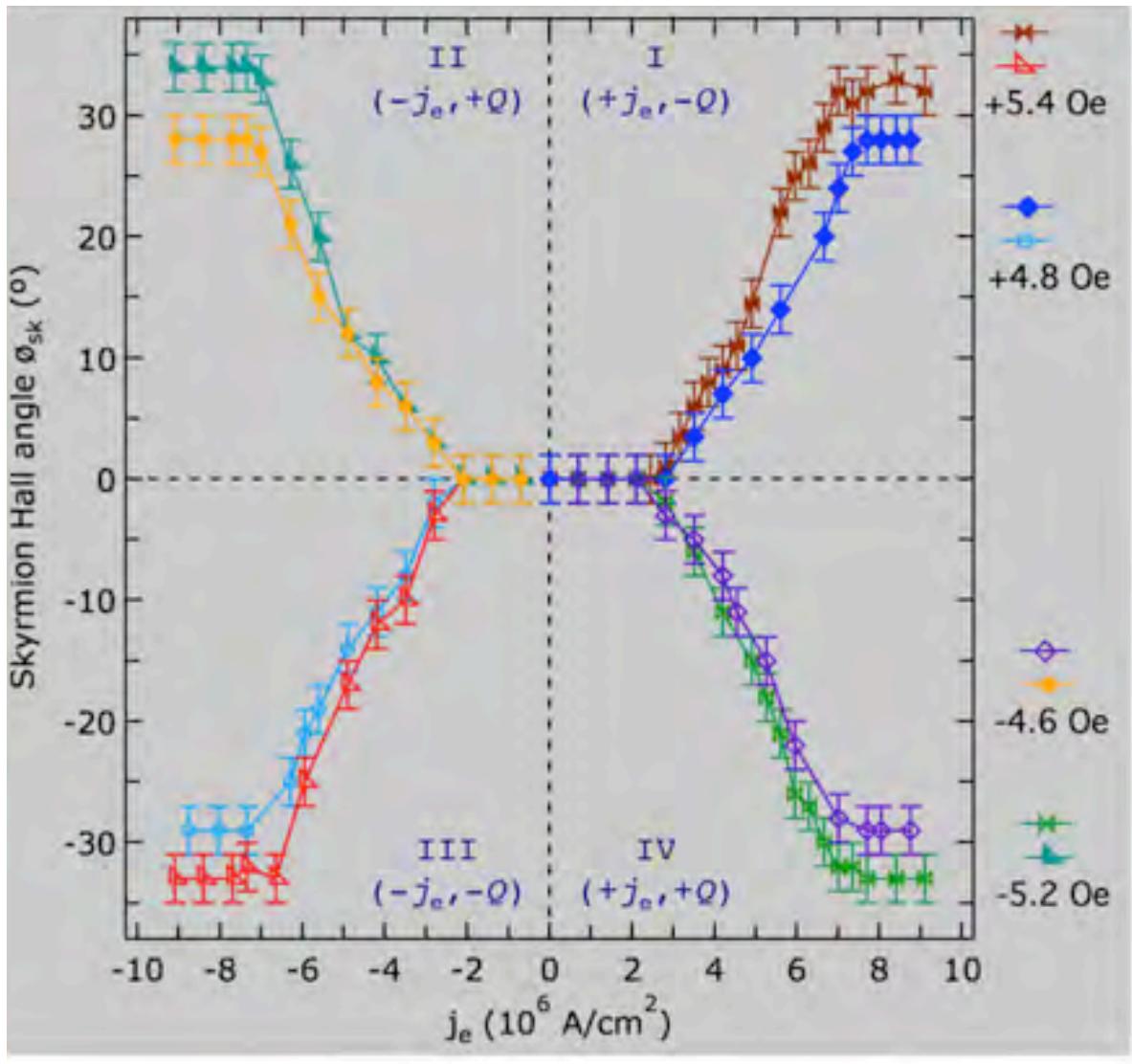
# DRIVE-DEPENDENT SKYRMION HALL ANGLE



The sign of the Skyrmion Hall angle changes with

- Current direction

# DRIVE-DEPENDENT SKYRMION HALL ANGLE



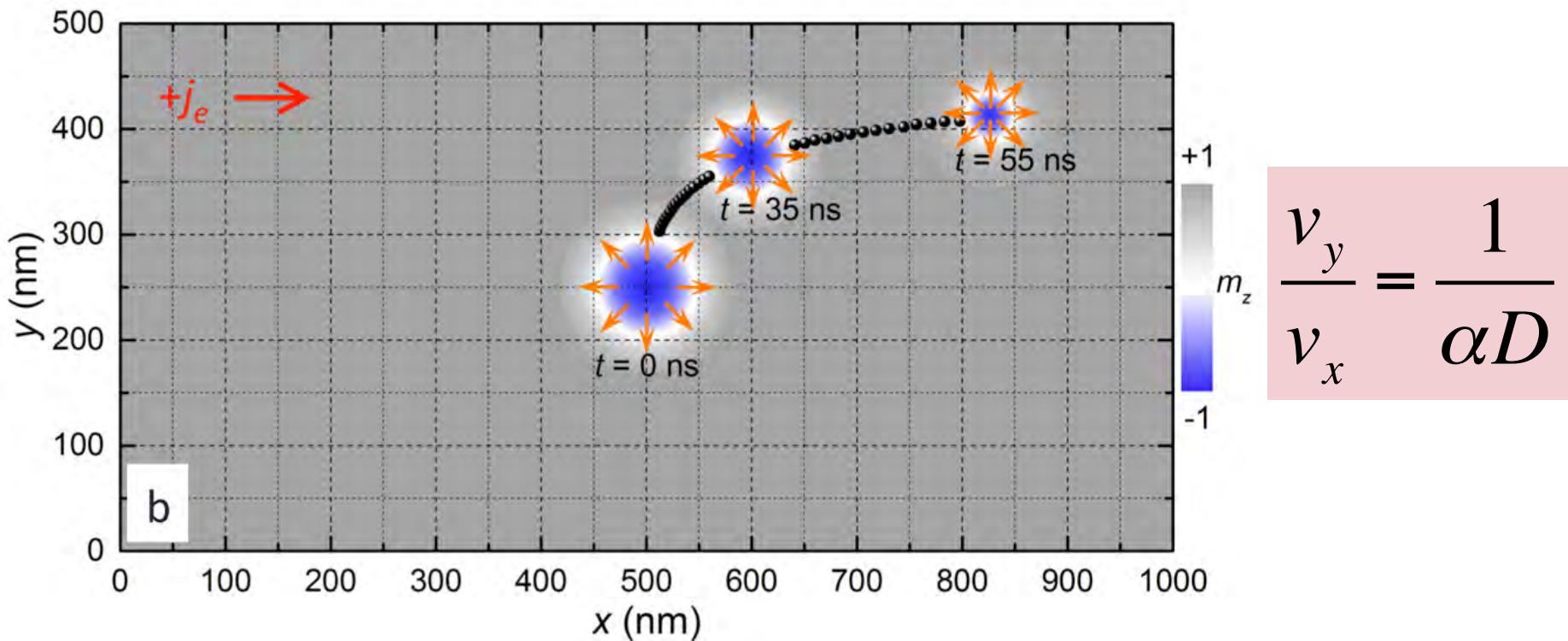
The sign of the Skyrmion Hall angle changes with

- Current direction
- Topological charge Q

# SKYRMION HALL EFFECT

Thiele Equation:  $G \times v - \alpha D \cdot v + 4\pi \vec{B} J_c = 0$

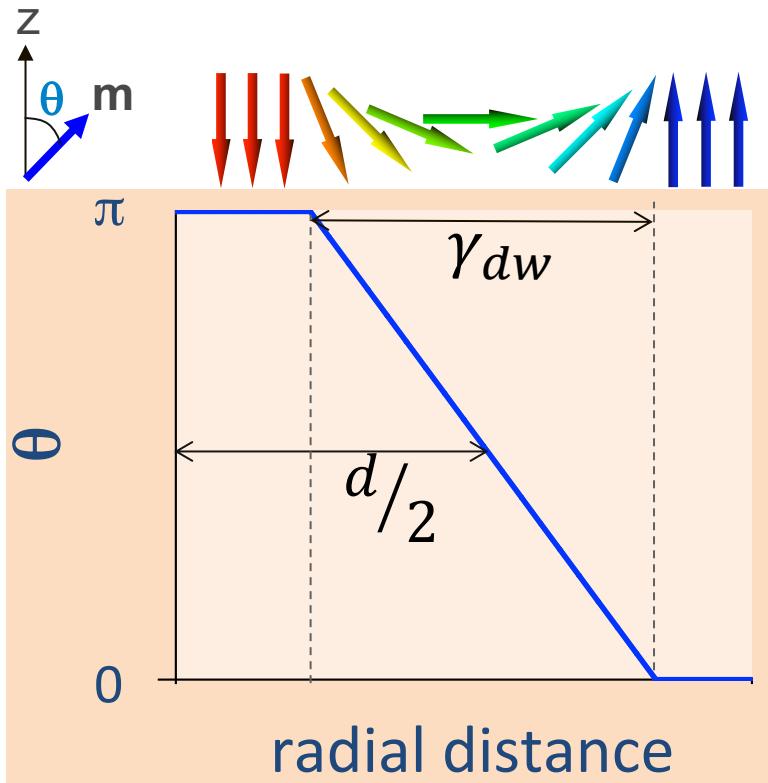
$$G = (0, 0, -4\pi Q) \quad Q = \frac{1}{4\pi} \int \mathbf{m} \cdot (\partial_x \mathbf{m} \times \partial_y \mathbf{m}) dx dy$$



$$\frac{v_y}{v_x} = \frac{1}{\alpha D}$$

# DISSIPATIVE FORCE TERM – ISOLATED SKYRMIONS

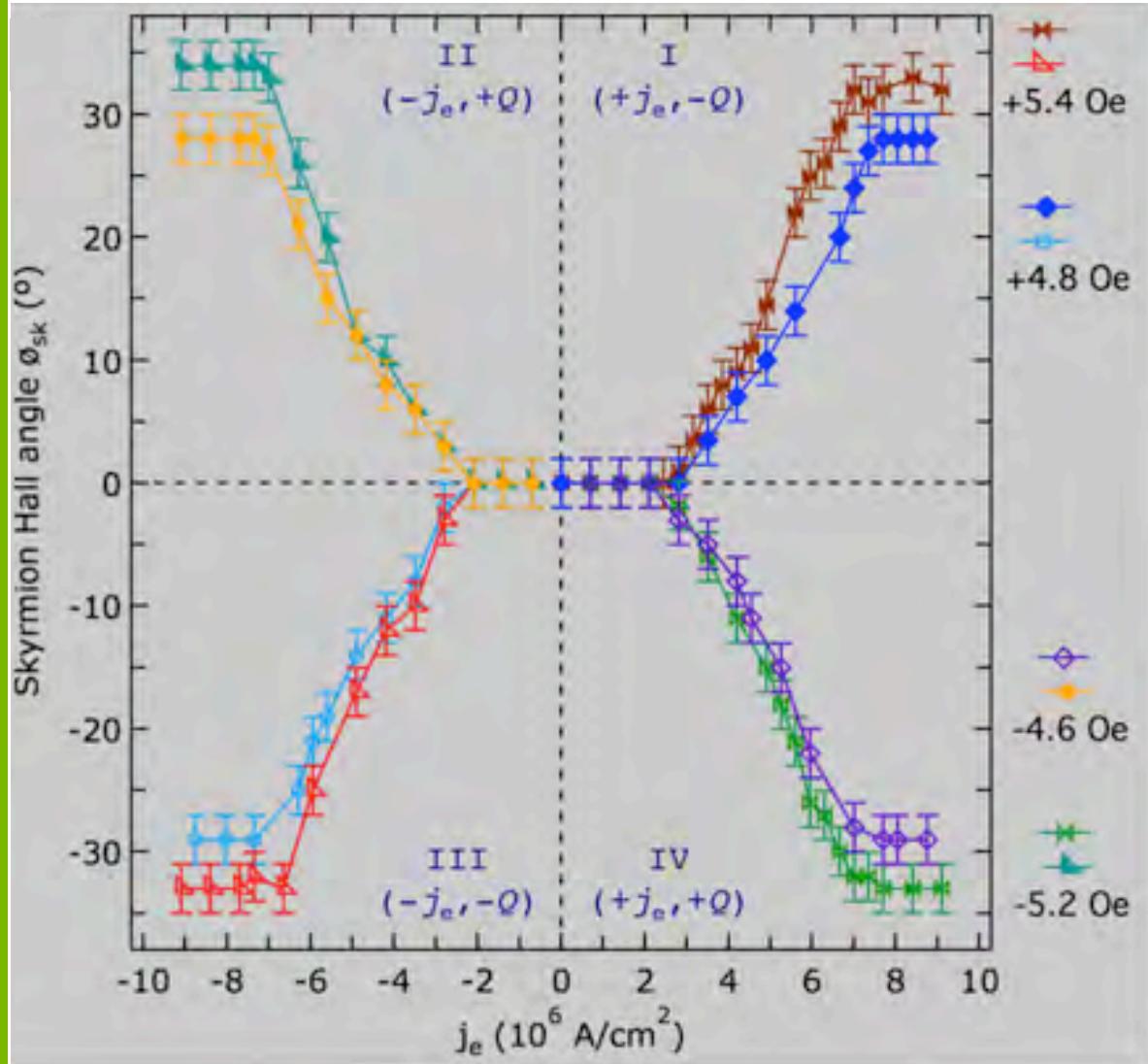
$$\vec{\mathcal{D}} = 4\pi \begin{pmatrix} \mathcal{D}_{ii} & 0 \\ \mathcal{D}_{ji} & \mathcal{D}_{jj} \end{pmatrix} \quad \mathcal{D} = \mathcal{D}_{ij} = \frac{1}{4\pi} \int_{\text{unit cell}} \frac{\partial \mathbf{m}}{\partial i} \cdot \frac{\partial \mathbf{m}}{\partial j} dx dy$$



$$\mathcal{D} = \frac{\pi^2 d}{8\gamma_{dw}}$$

$d$	$\gamma_{dw}$	$D$	$v_y/v_x$	$\Phi_{sk}$
10 nm	10 nm	1.2	40	89°
100 nm	10 nm	12	4	76°
1000 nm	10 nm	12 3	0.4	22°

# DRIVE-DEPENDENT SKYRMION HALL ANGLE



$$\gamma_{dw} \approx \sqrt{A/K_{eff}} \approx 21 \text{ nm}$$
$$\alpha \approx 0.02$$

$H = 5.4 \text{ Oe}, d = 800 \pm 300 \text{ nm}$

$$\Phi_{sk} = 32 \pm 2^\circ$$

$$\Phi_{sk} \approx 47 \pm 11^\circ$$

$H = 4.8 \text{ Oe}, d = 1100 \pm 300 \text{ nm}$

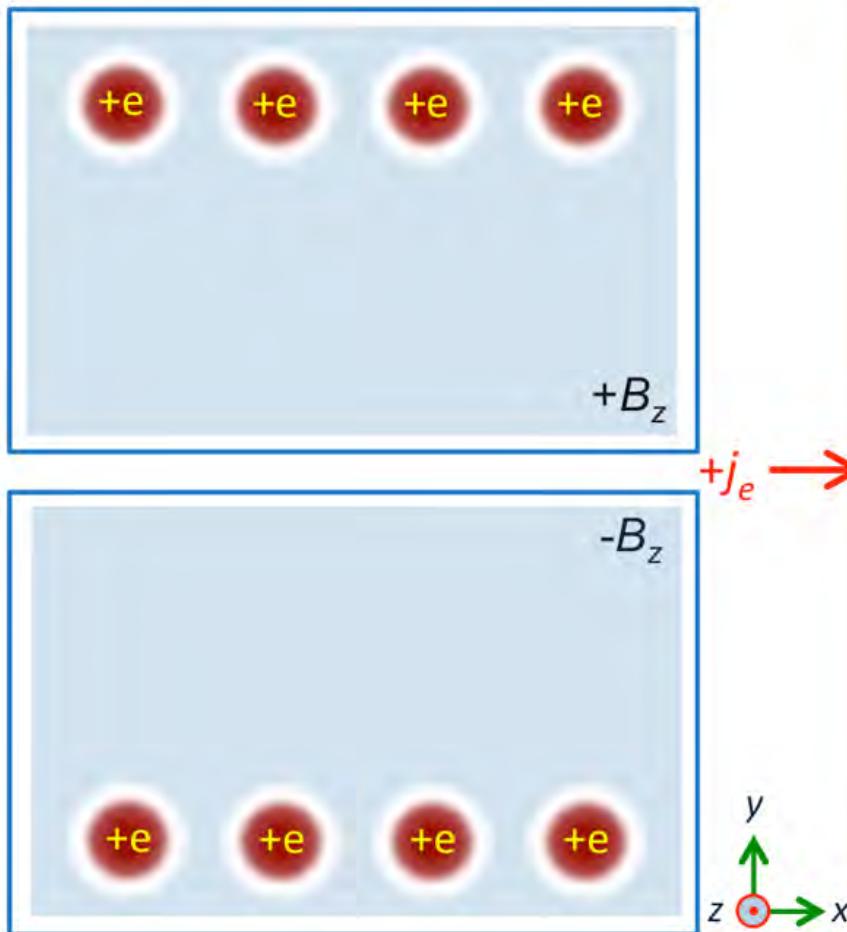
$$\Phi_{sk} = 28 \pm 2^\circ$$

$$\Phi_{sk} \approx 38 \pm 8^\circ$$

# Classic Hall effect

Electric charge  $q_e$

Lorentz force



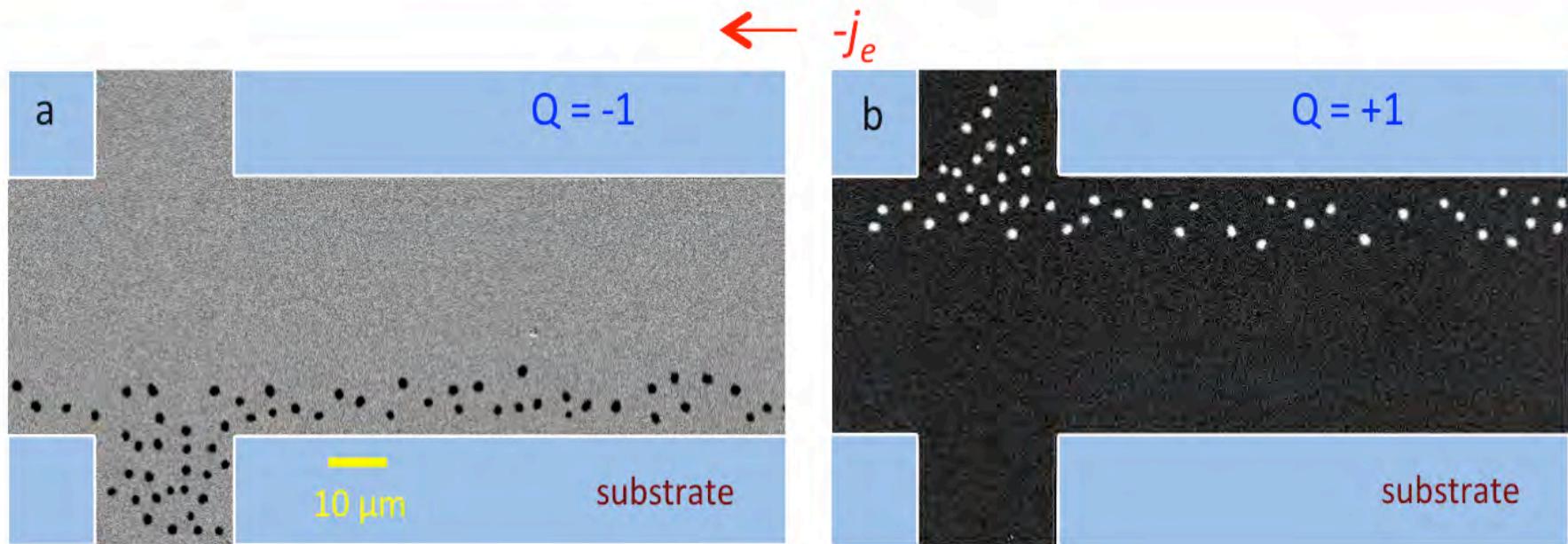
# Skymion Hall effect

Topological charge  $Q$

Magnus force

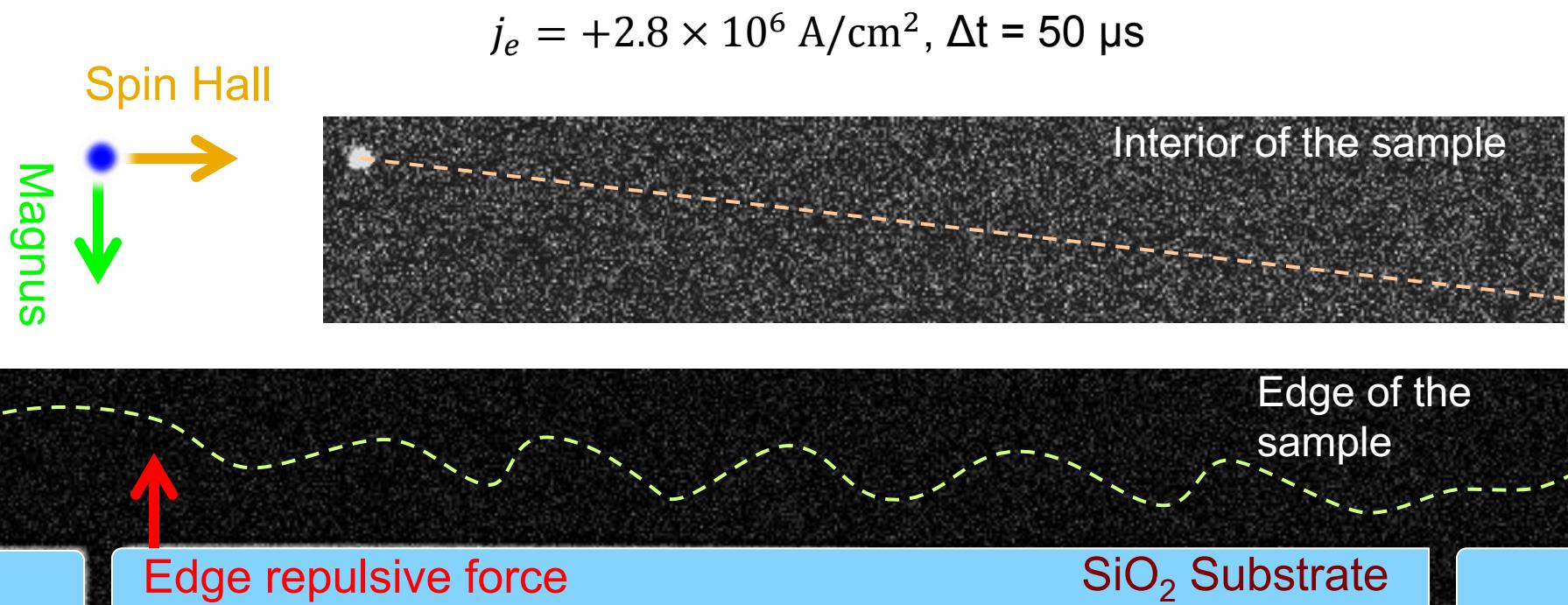
$$Q = \frac{1}{4\pi} \int \mathbf{m} \cdot (\partial_x \mathbf{m} \times \partial_y \mathbf{m}) dx dy$$

# SKYRMION ACCUMULATION



Accumulation of skyrmions along one wire edge is observed.

# SKYRMION MOTION: EDGE VS. INTERIOR



Oscillatory skyrmion edge transport:  
Competition between edge repulsive force and  
magnus force

# CONCLUSIONS

- Magnetic Skyrmions at RT
  - Use inhomogeneous currents for generating skyrmions
- Skyrmion Hall Effect
  - **Low current:** Linear dependence of skyrmion Hall angle (creep regime).
  - **High current:** Saturation of skyrmion Hall angle (flow regime).
  - Magnitude of skyrmion Hall angle  $\Phi_{sk}$  consistent with size of skyrmion.
  - Sign of  $\Phi_{sk}$  changes with sign of  $Q$ .

