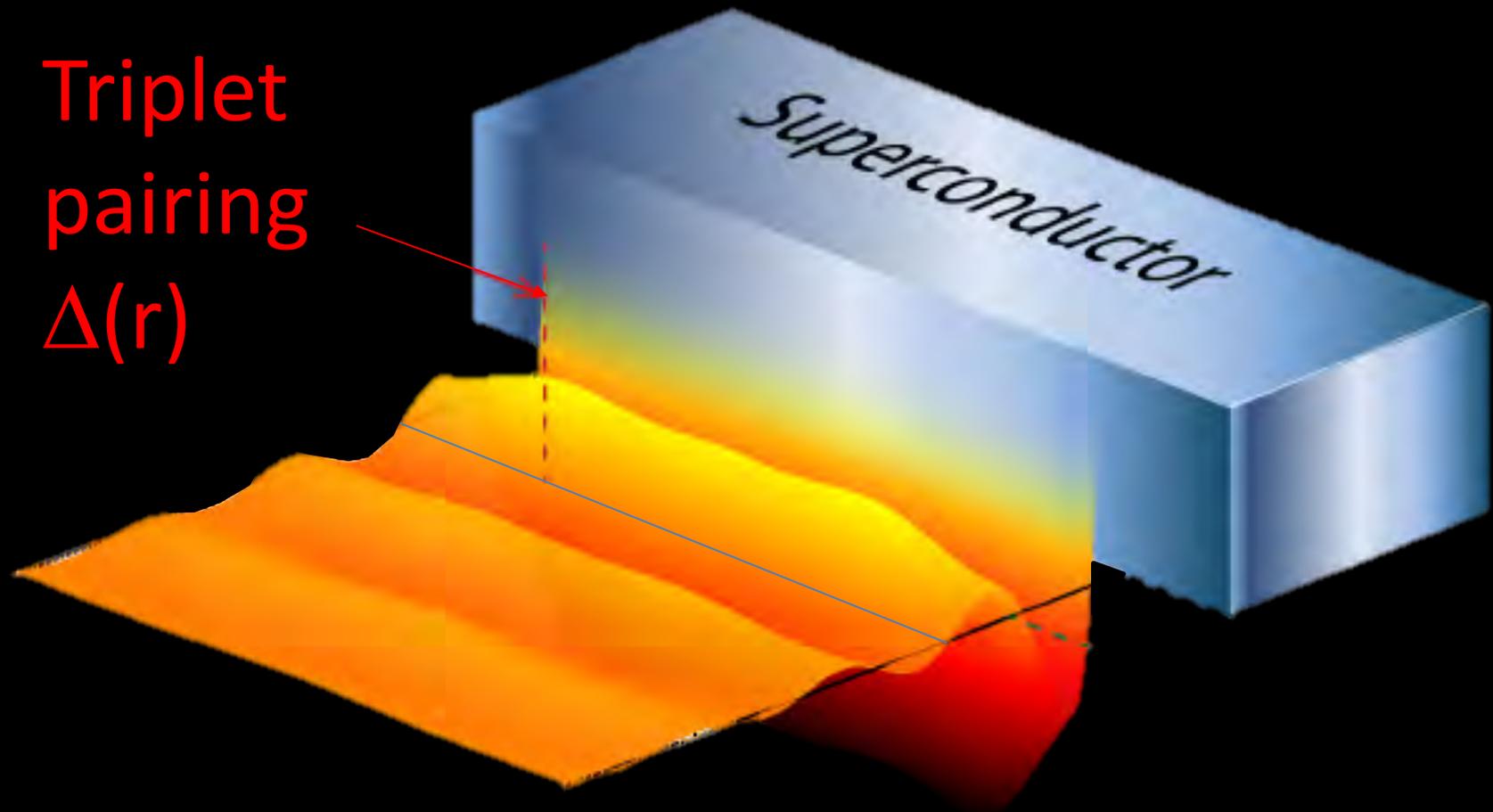


# A New Spin on Superconductivity

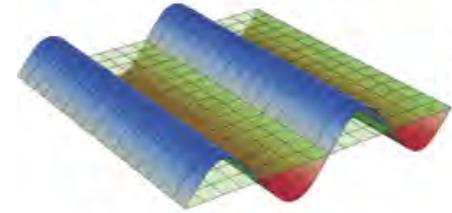
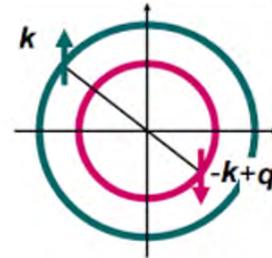
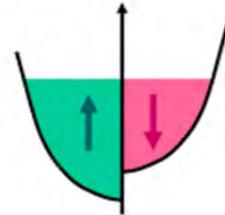
Amir Yacoby, Harvard University



# Unconventional Superconductivity: Finite momentum pairing -Triplet pairing

Breaking of an additional symmetry: translational symmetry; spin symmetry...

1964 – Fulde-Ferrell; Larkin-Ovchinnikov



Spatially varying order parameter

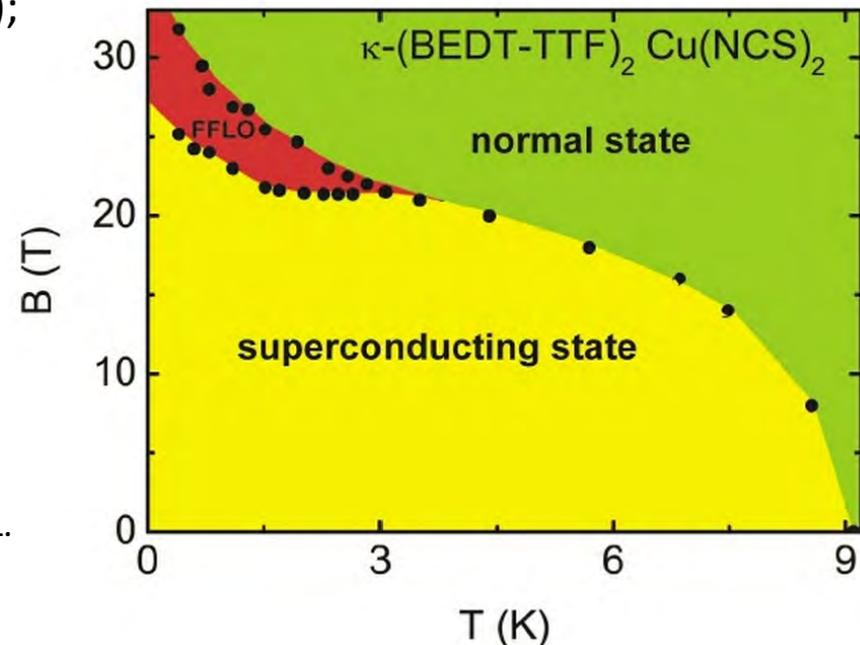
P. Fulde, R. A. Ferrell, Phys. Rev. 135, A550 (1964);  
A. I. Larkin, Y. N. Ovchinnikov, Sov. Phys. JETP 20,  
762 (1965).

Organic SC:

H. Shimahara, in A.G. Lebed (ed.): The Physics of Organic  
Superconductors and Conductors, Springer, Berlin (2008)

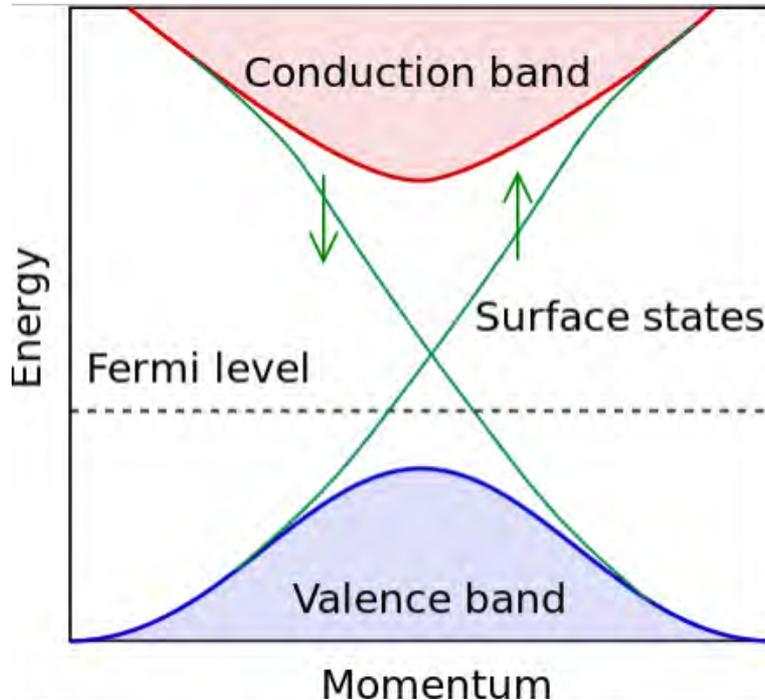
heavy-fermion material  $\text{CeCoIn}_5$

Bianchi, A.; Movshovich, R.; Capan, C.; Pagliuso, P.G.; Sarrao, J.L.  
Phys. Rev. Lett 91, 03'



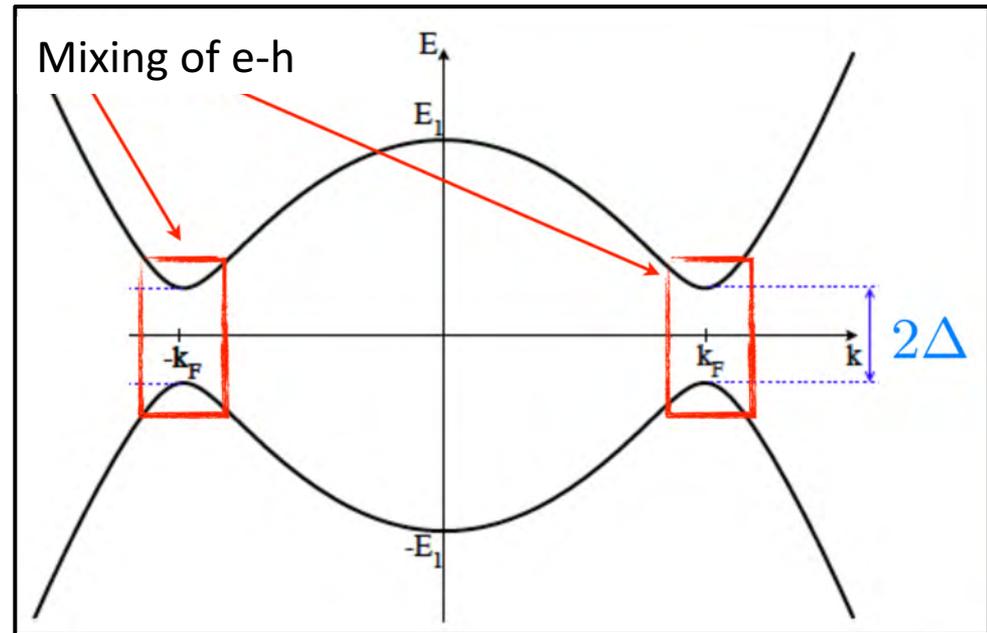
# Topological Insulators and Topological Superconductors

Insulators



Topological insulators have **electronic** surface states

Superconductors



Topological superconductors have **Majorana** surface states

# Bulk Topological Superconductivity

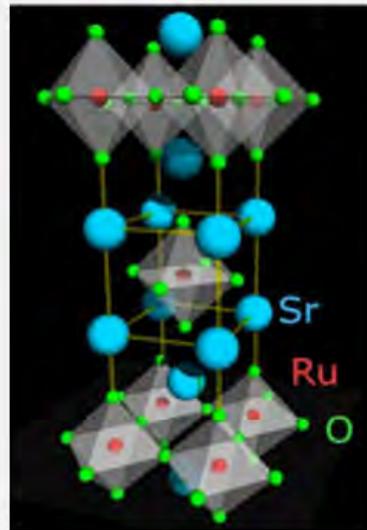
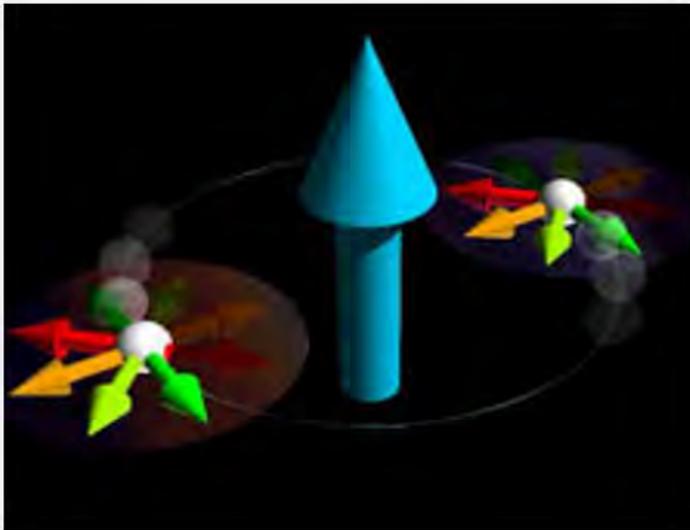
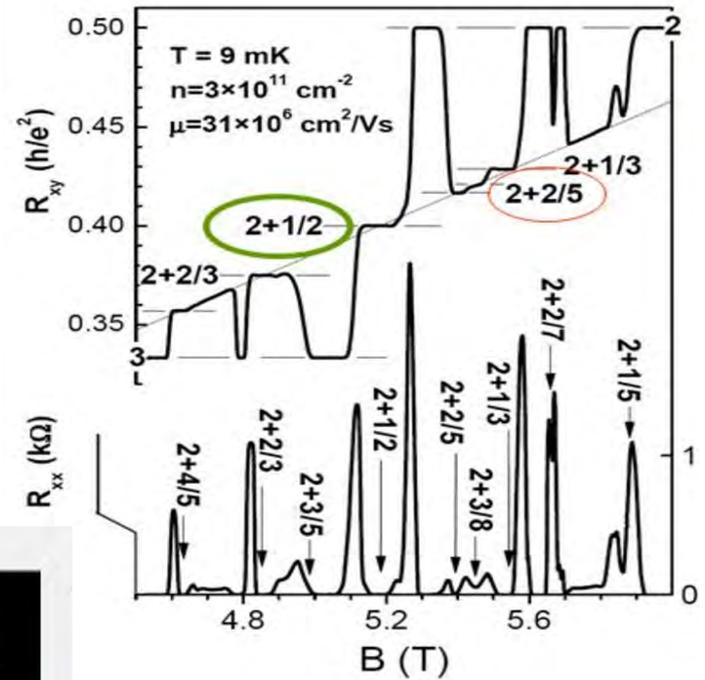
$\nu=5/2$  FQHE – Read and Green

Non Abelian particles – when exchanged wave function changes completely

$$e^* = e/4.$$

M. Heiblum et al, 2010  
V. Venkatachalam, AY, Nature 2010

J.S. Xia et al., PRL (2004).

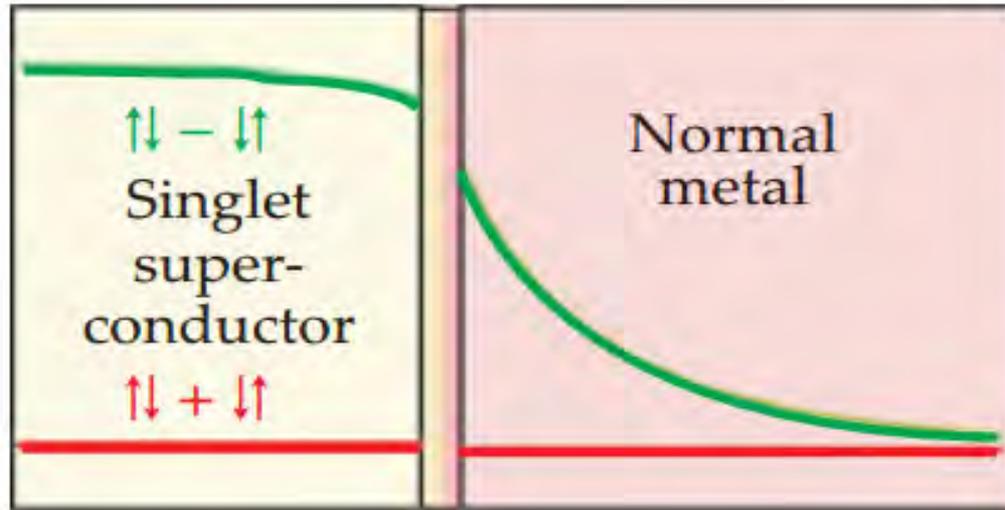


$\text{SrRu}_2\text{O}_2$   
Y. Maeno

# Unconventional Superconductivity: Finite momentum pairing -Triplet pairing

- Superconductor – Ferromagnetic junctions

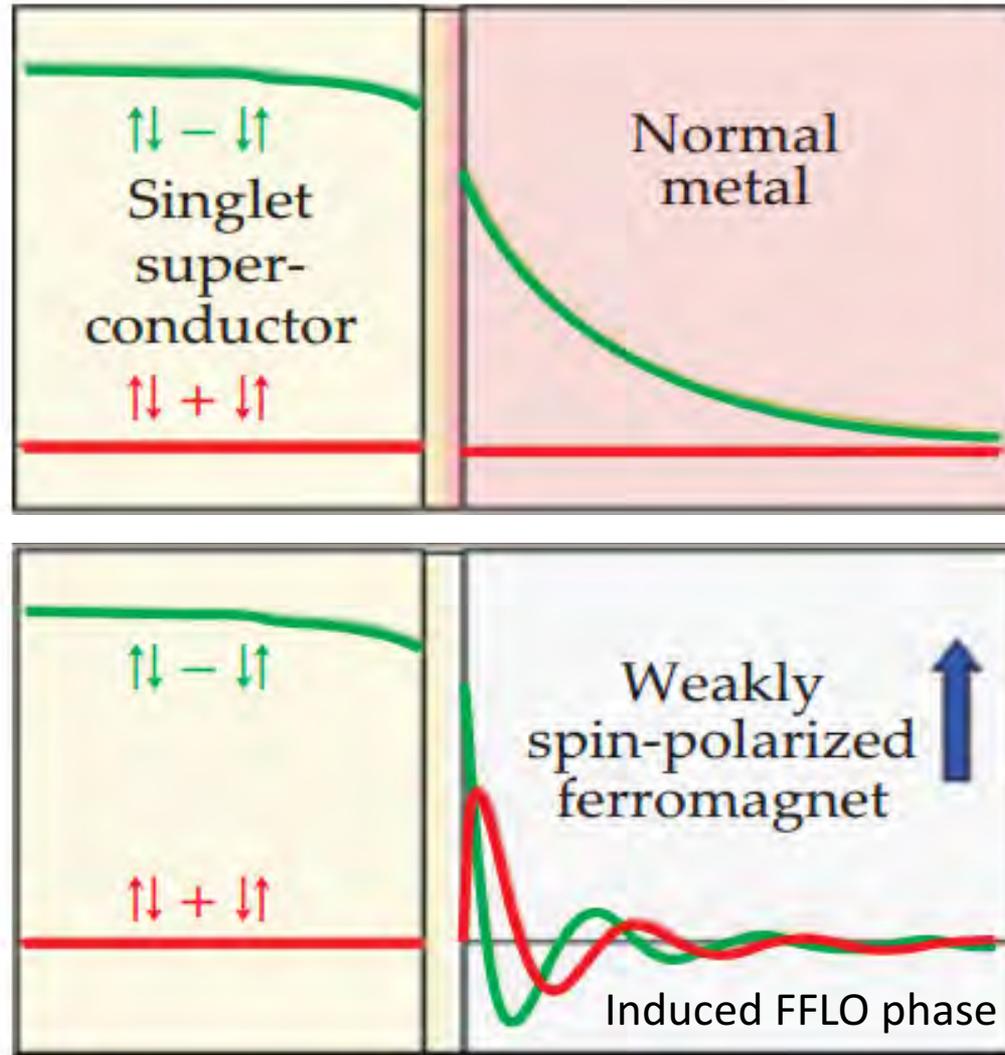
See review: Matthias Eschrig, Phys. Today 64(1), 43 (2011)



# Unconventional Superconductivity: Finite momentum pairing -Triplet pairing

- Superconductor – Ferromagnetic junctions

See review: Matthias Eschrig, Phys. Today 64(1), 43 (2011)



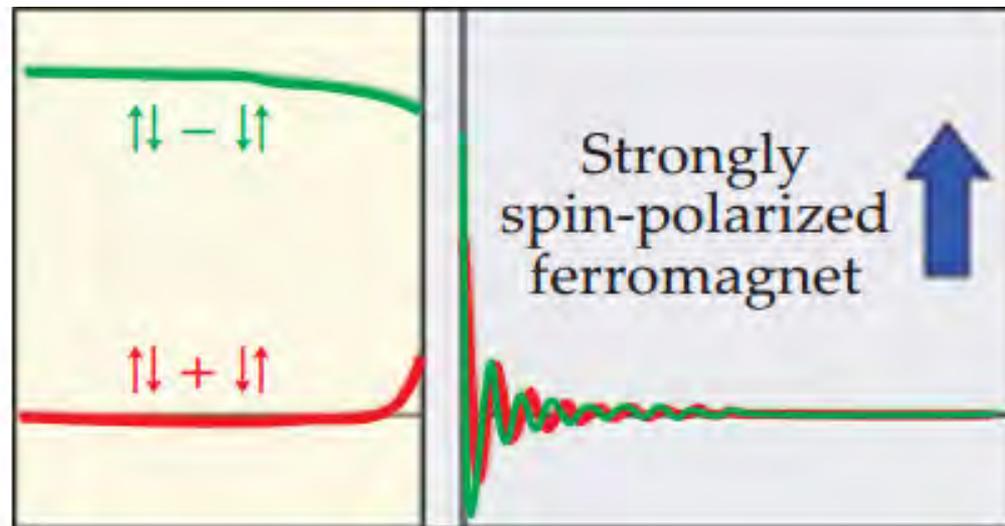
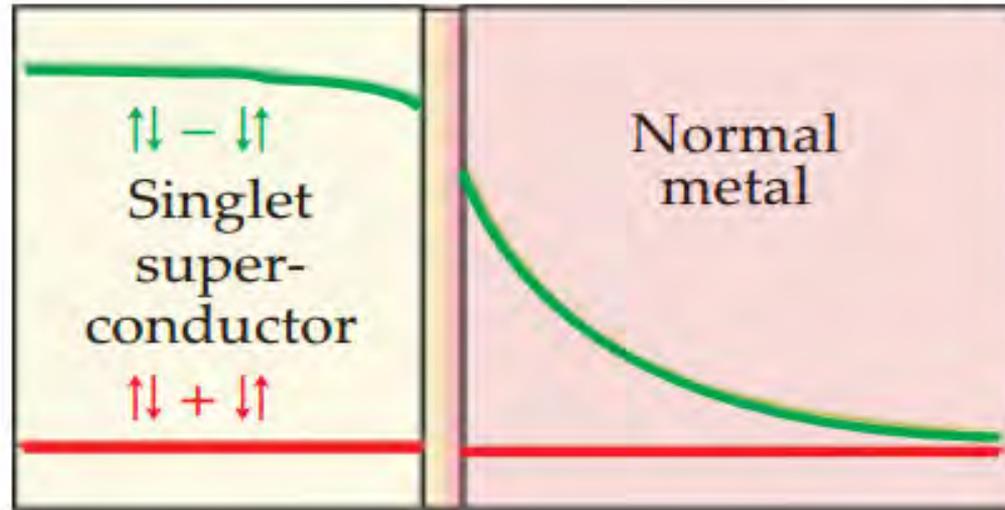
Theory:  
Buzdin 1982

Experiments:2001-2  
Ryazanov et al;  
Golubov et al; Kontos  
et al and Palevski

# Unconventional Superconductivity: Finite momentum pairing - Triplet pairing

- Superconductor – Ferromagnetic junctions

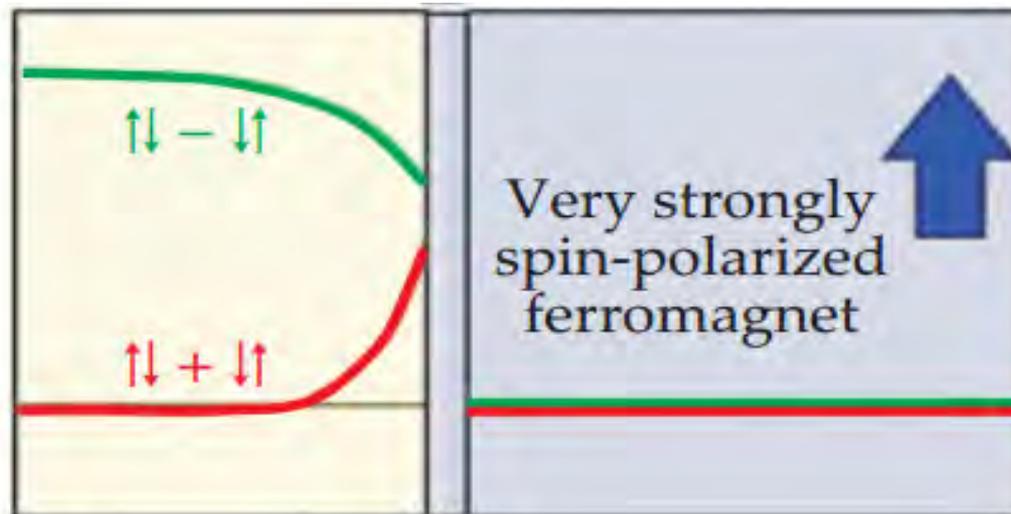
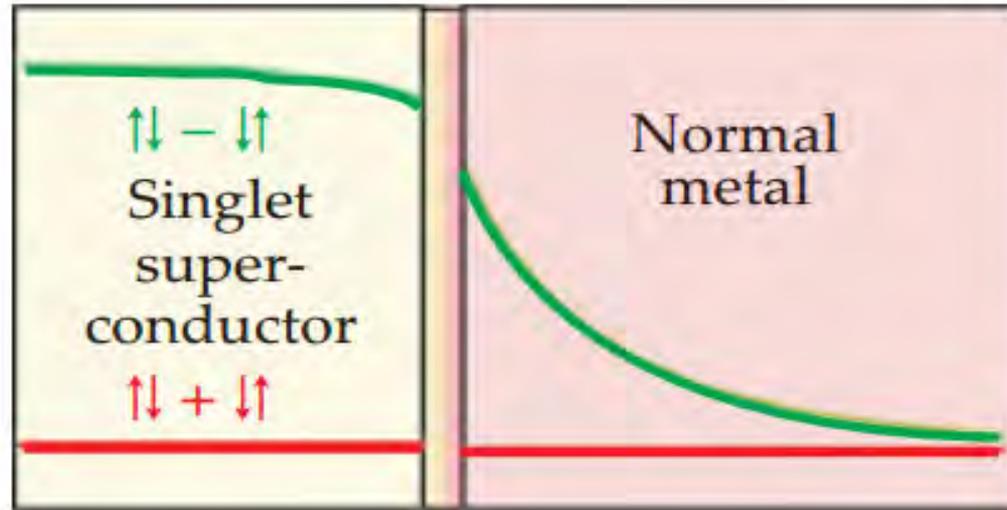
See review: Matthias Eschrig, Phys. Today 64(1), 43 (2011)



# Unconventional Superconductivity: Finite momentum pairing - Triplet pairing

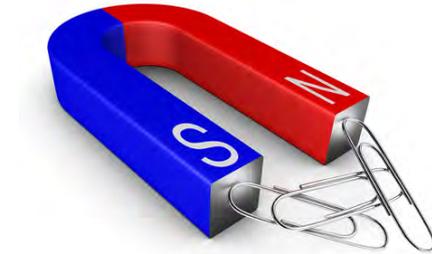
- Superconductor – Ferromagnetic junctions

See review: Matthias Eschrig, Phys. Today 64(1), 43 (2011)

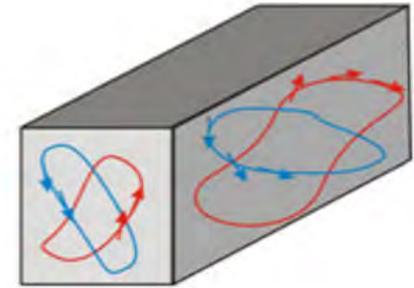
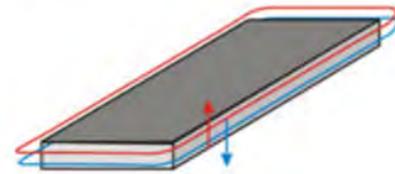
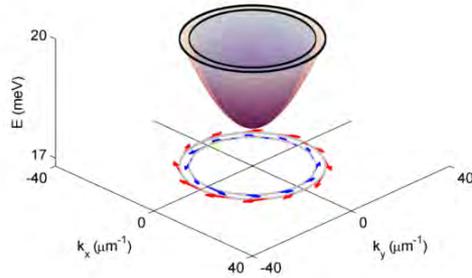


# Unconventional Superconductivity from Hybrid Devices

'Conventional Phases'



Materials with strong spin momentum locking



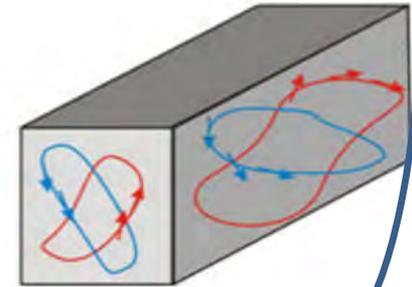
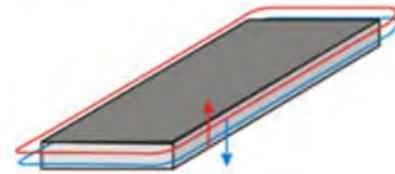
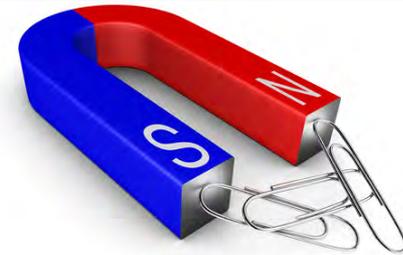
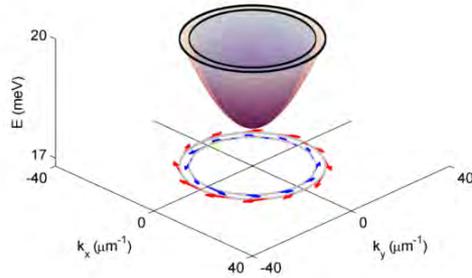
Topological Insulators

# Unconventional Superconductivity from Hybrid Devices

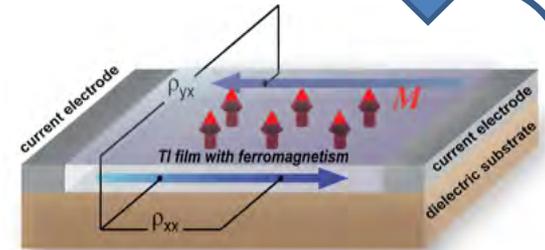
'Conventional Phases'



Materials with strong spin momentum locking



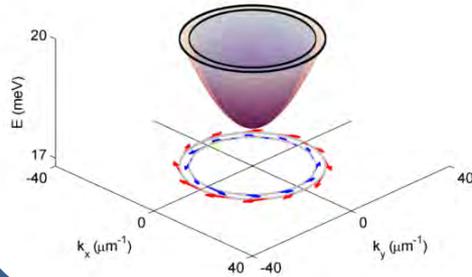
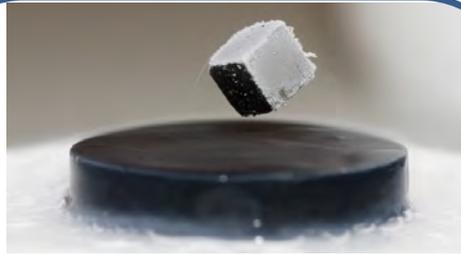
Topological Insulators



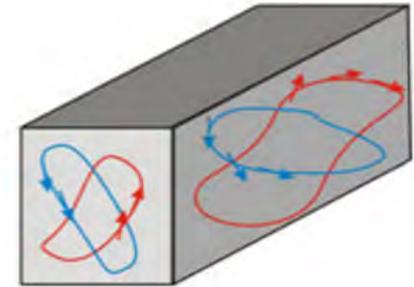
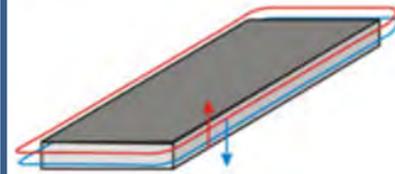
Quantum Anomalous Hall Effects:  
Dissipationless transport and **chiral Majorana modes**

# Unconventional Superconductivity from Hybrid Devices

'Conventional Phases'



Materials with strong spin momentum locking

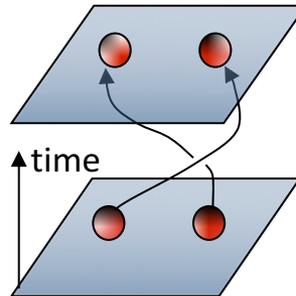


Topological Insulators



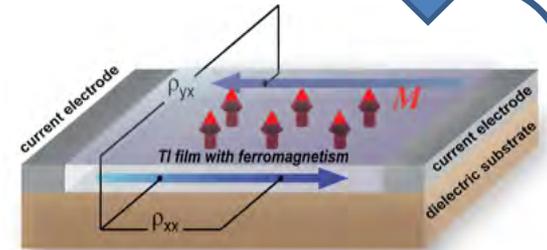
Triplet Cooper pairs

$$\psi(r_1, r_2) \neq e^{i\varphi} \psi(r_2, r_1)$$



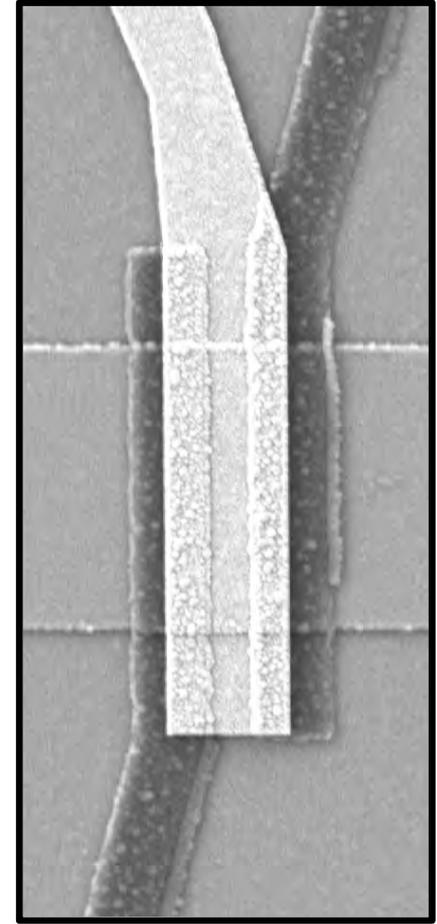
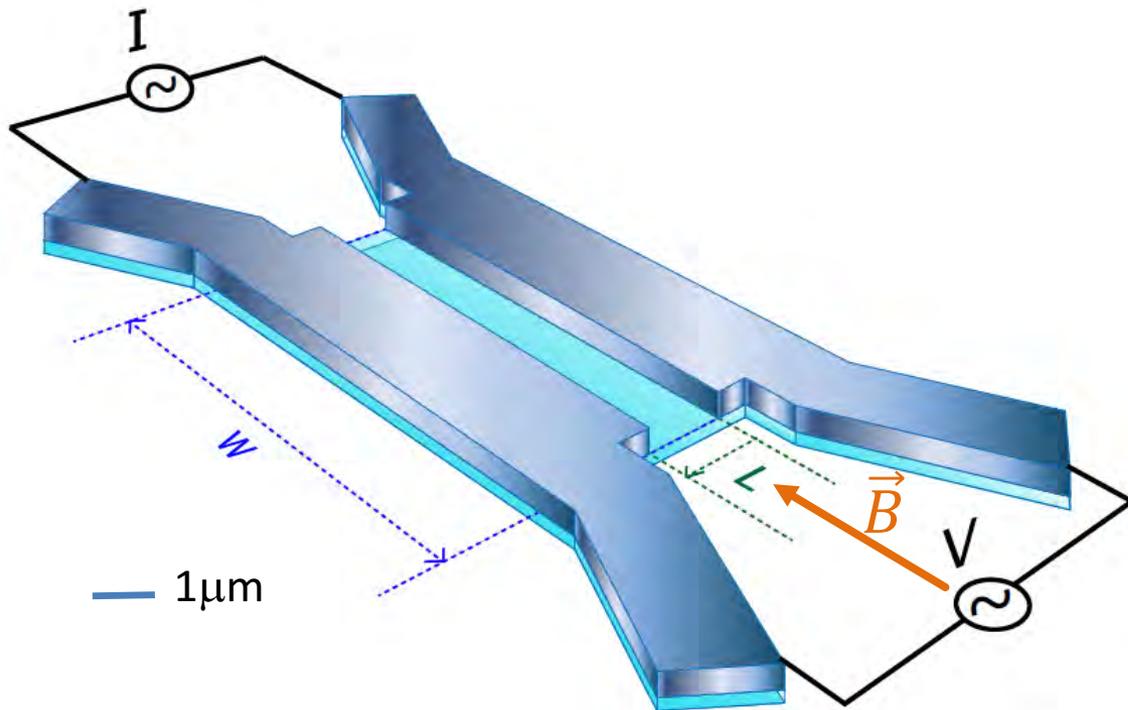
Topological Superconductivity:  
Non Abelian Statistics -  
Topological quantum computing

Triplet superconductivity:  
Unconventional superconductivity

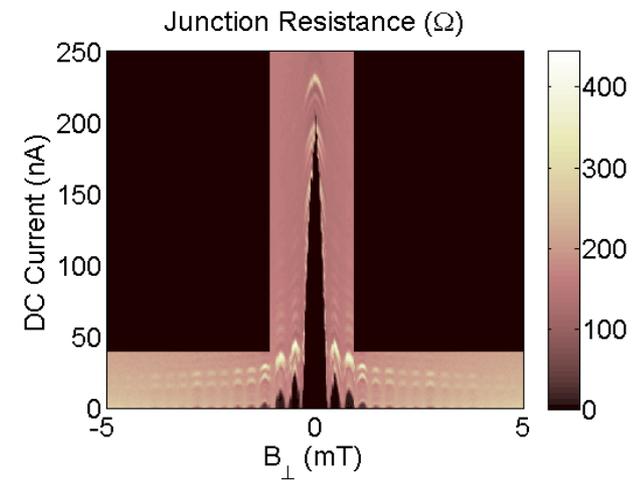
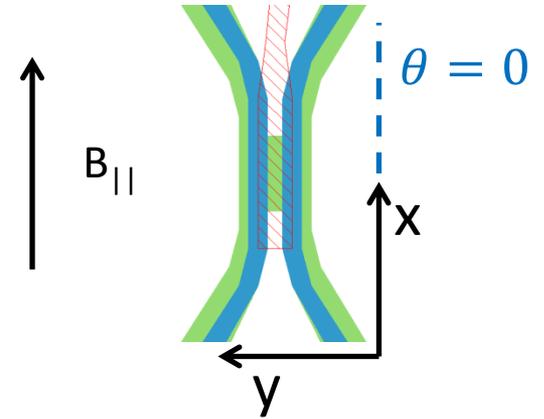


Quantum Anomalous Hall Effects:  
Dissipationless transport and chiral  
Majorana modes

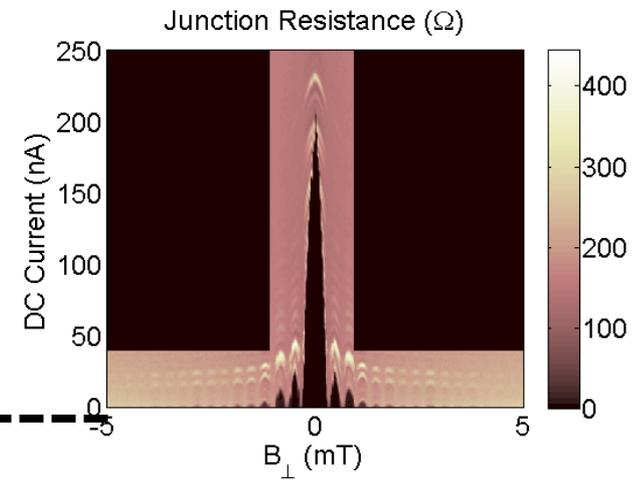
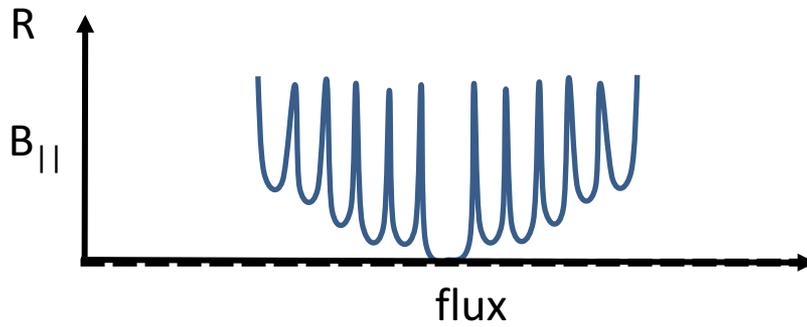
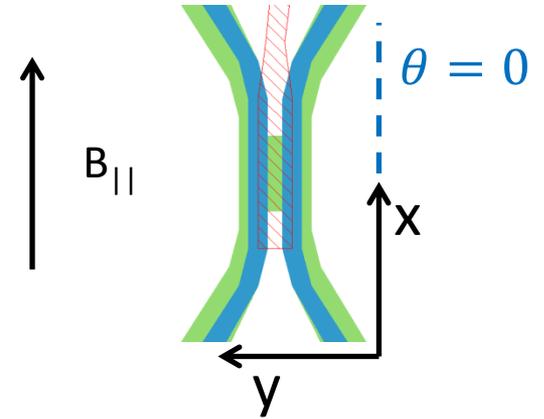
# Experimental Geometry



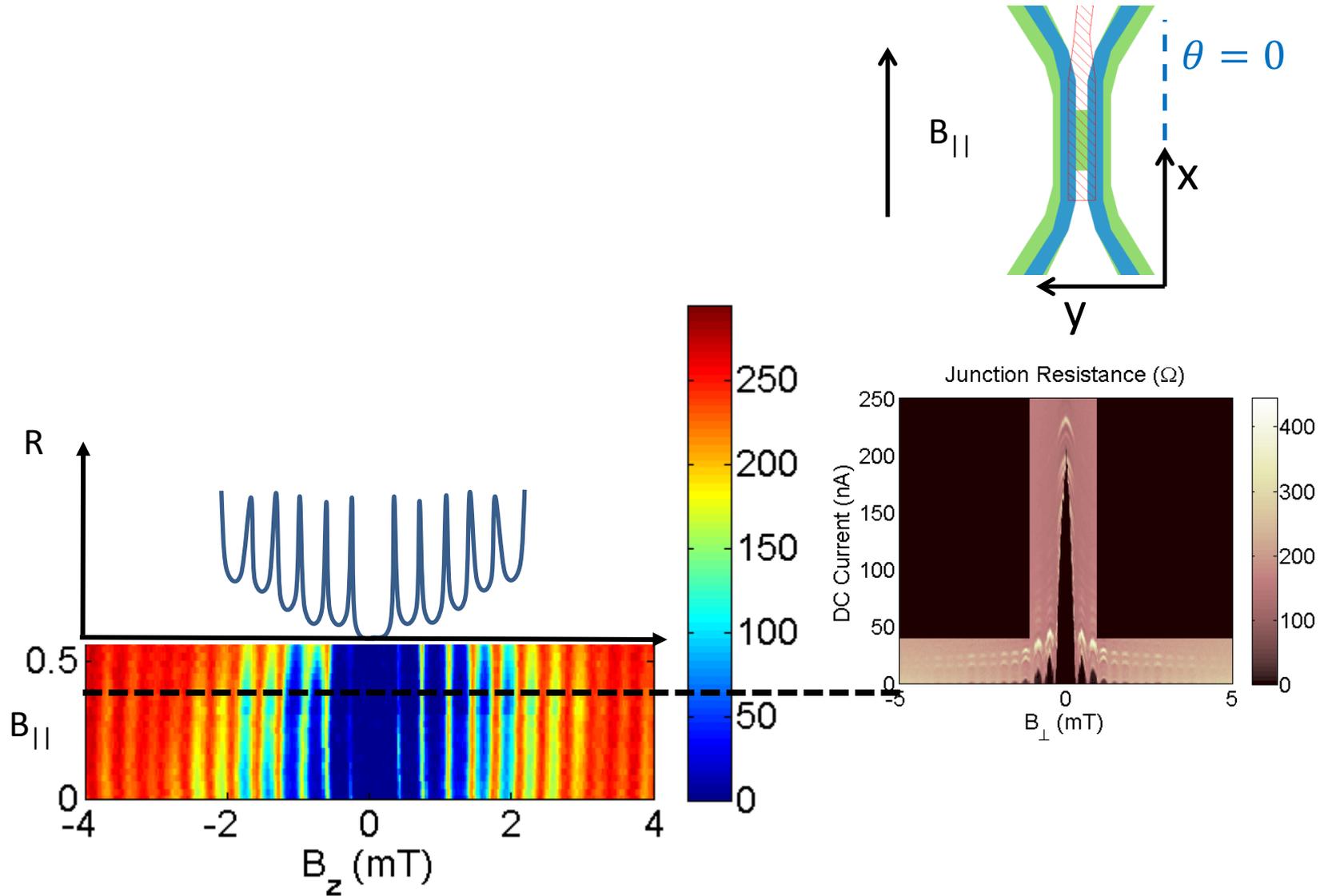
# Experimental Geometry



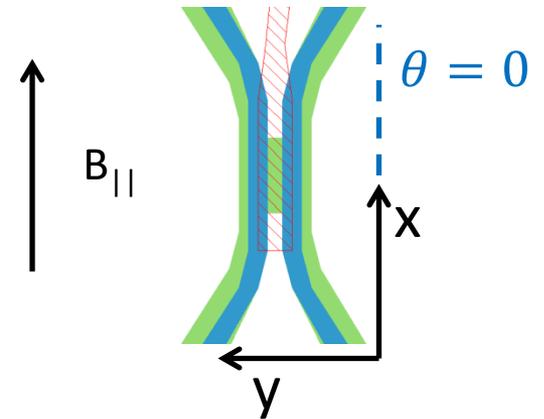
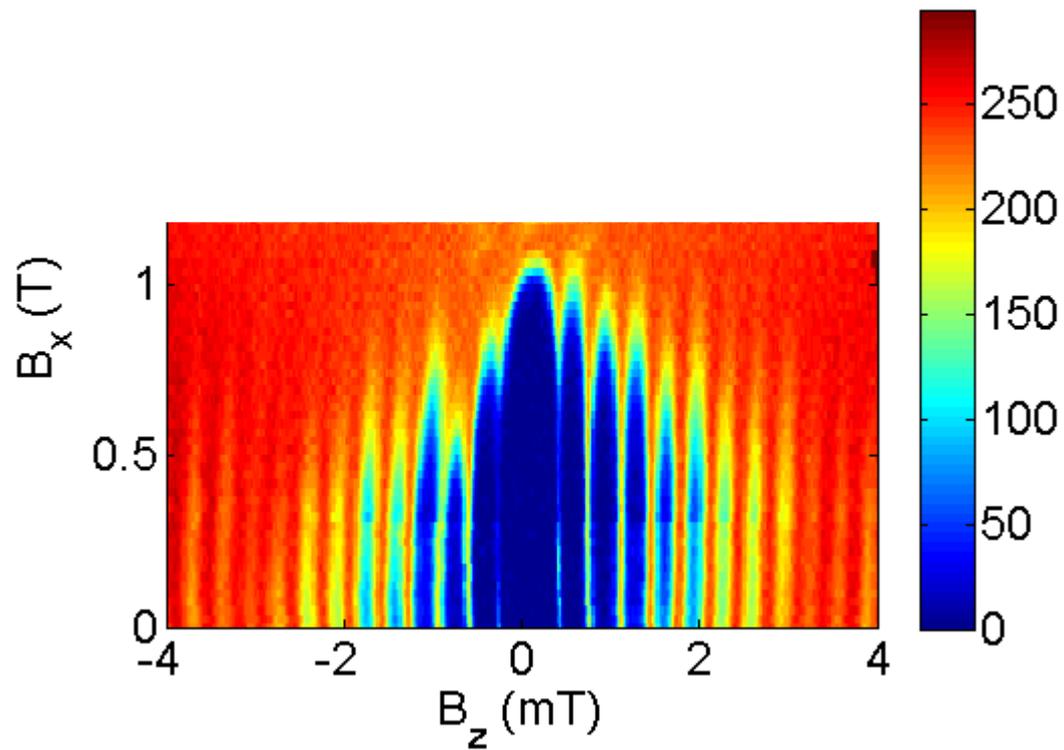
# Experimental Geometry – In Plane Magnetic Field



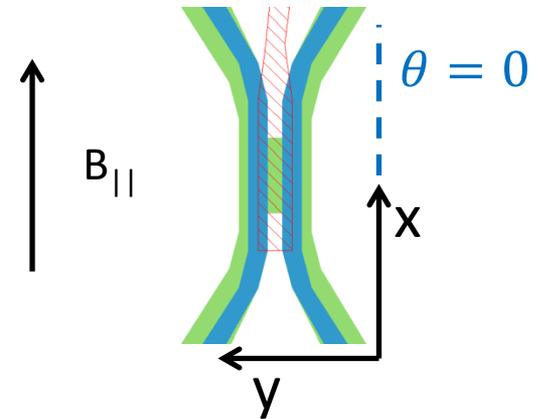
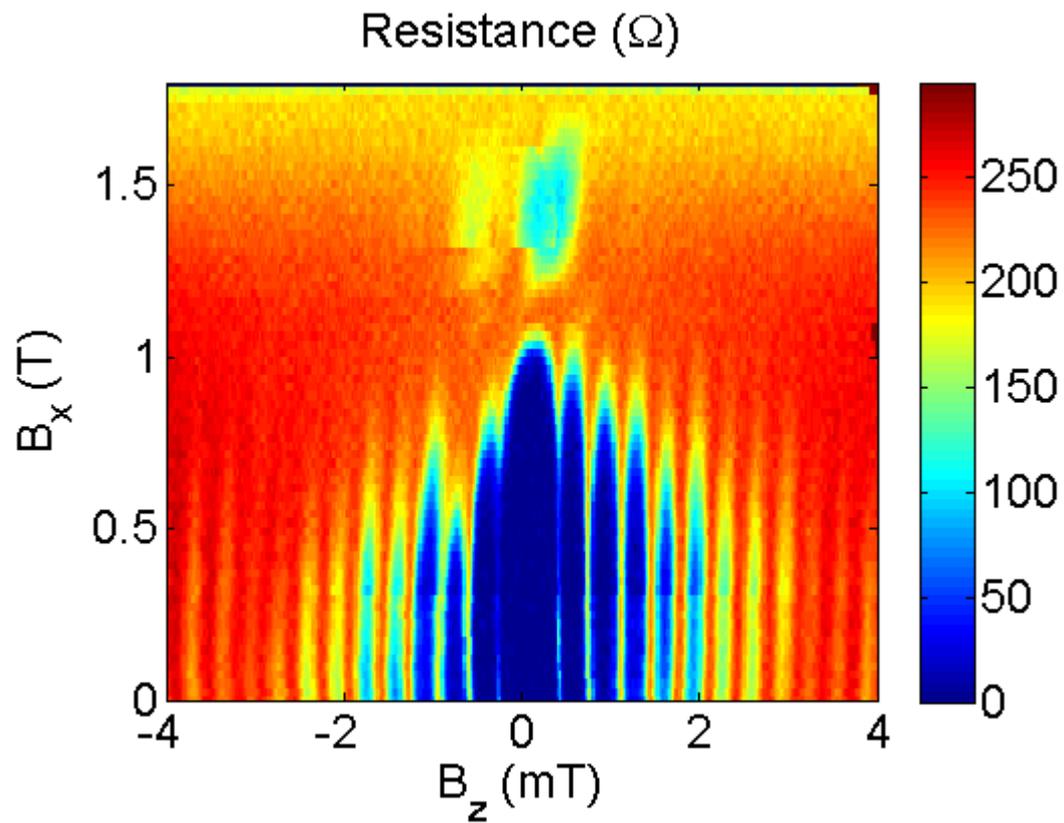
# Experimental Geometry – In Plane Magnetic Field



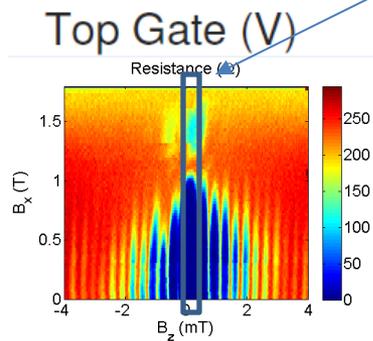
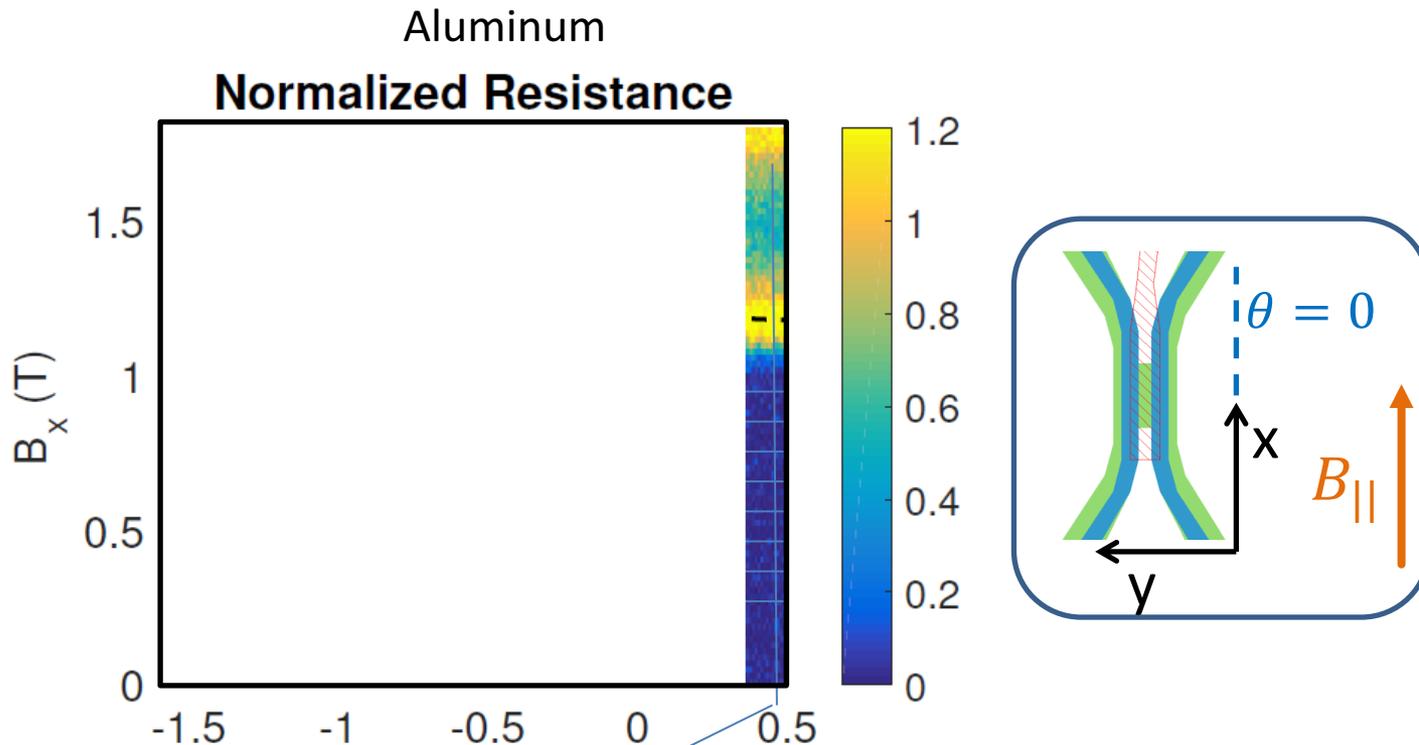
# Experimental Geometry – In Plane Magnetic Field



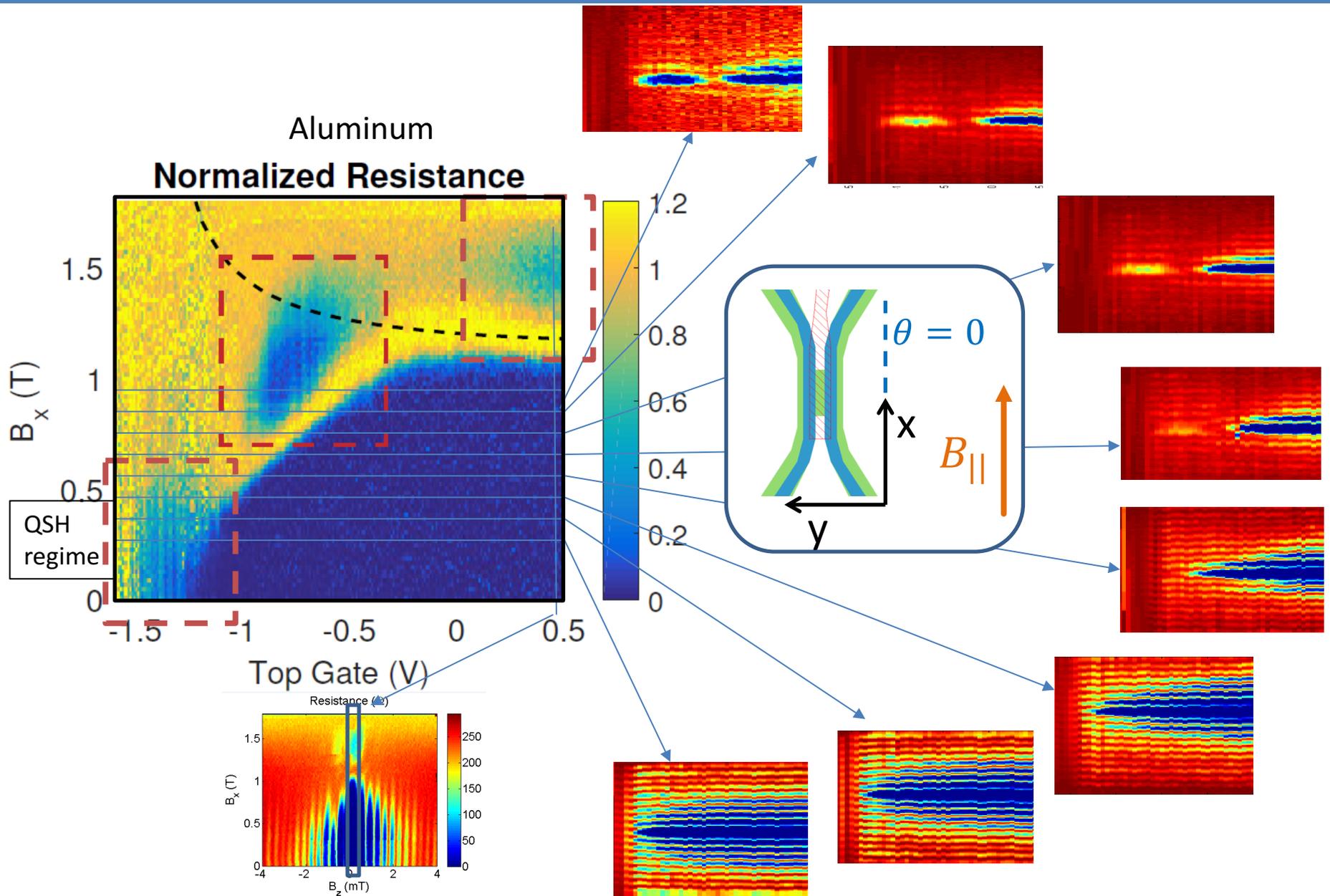
# Experimental Geometry – In Plane Magnetic Field



# Phase Diagram of $B_{||}$ perpendicular to current flow



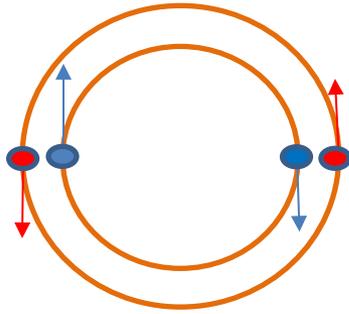
# Phase Diagram of $B_{||}$ perpendicular to current flow



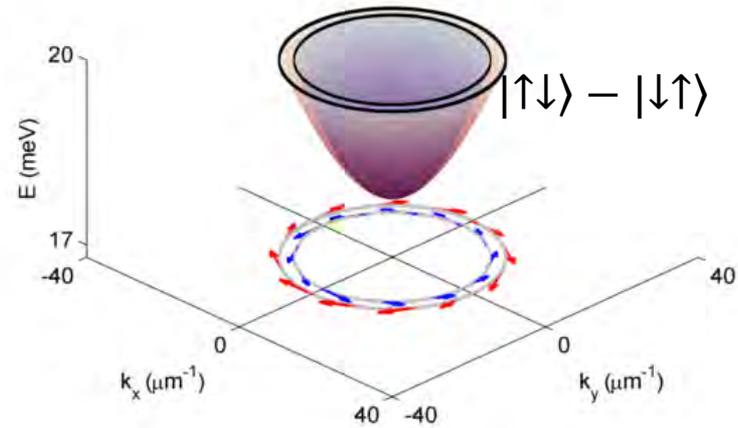
# Expectations with Rashba Spin-Orbit Interaction

B=0

Pairing with each Fermi surface



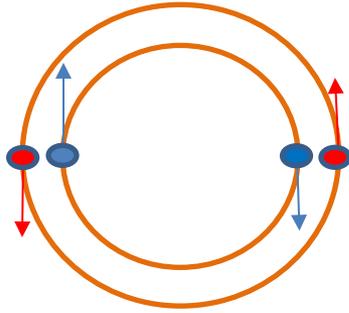
$$\vec{B} = (0 \text{ T}, 0 \text{ T}, 0 \text{ T})$$



# Expectations with Rashba Spin-Orbit Interaction

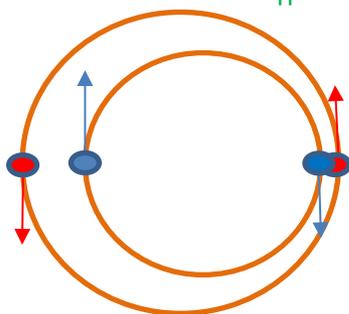
B=0

Pairing with each Fermi surface



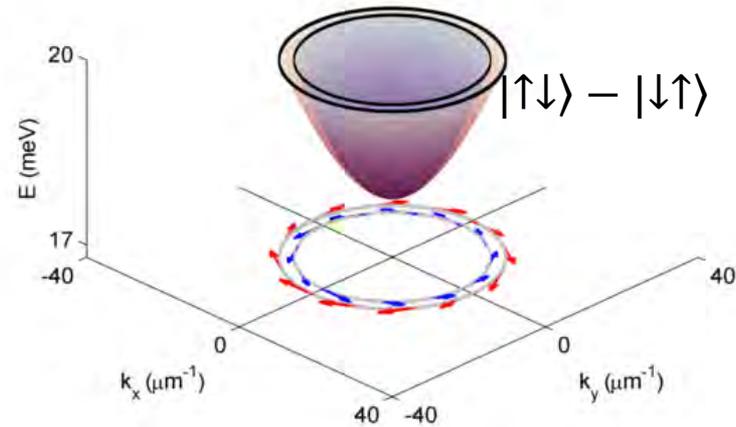
B>0

Some pairs acquire momentum  $\Delta K \perp B_{||}$



$$e^{i\vec{\Delta K} \cdot \vec{r}} |\uparrow\downarrow\rangle - e^{-i\vec{\Delta K} \cdot \vec{r}} |\downarrow\uparrow\rangle$$

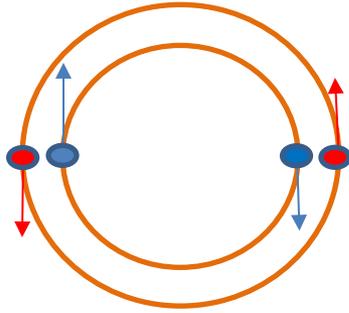
$$\vec{B} = (0 T, 0 T, 0 T)$$



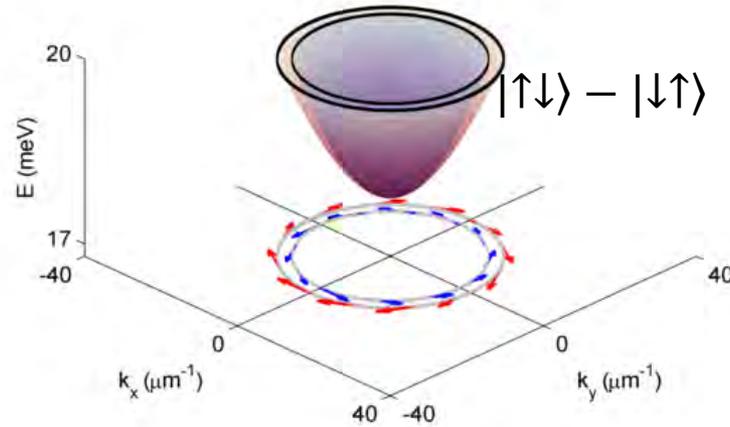
# Expectations with Rashba Spin-Orbit Interaction

B=0

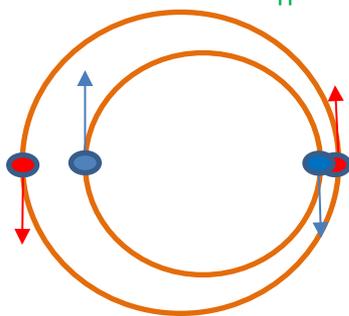
Pairing with each Fermi surface



$$\vec{B} = (0 T, 0 T, 0 T)$$



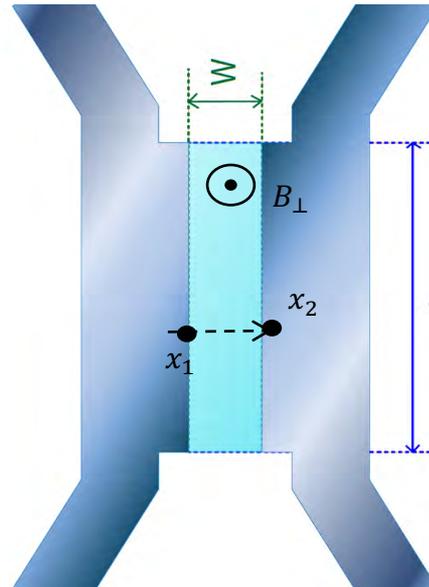
Some pairs acquire momentum  $\Delta K \perp B_{||}$



B>0

$$\Delta K = \frac{E_Z}{\hbar V_F}$$

$$e^{i\vec{\Delta K} \cdot \vec{r}} |\uparrow\downarrow\rangle - e^{-i\vec{\Delta K} \cdot \vec{r}} |\downarrow\uparrow\rangle$$

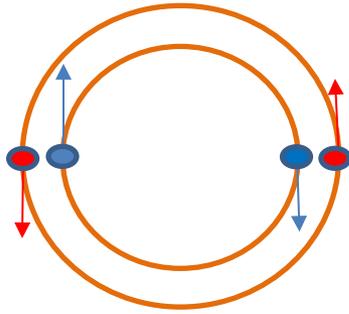


$$\Delta_{\text{ind}} = \text{Cos}(\vec{\Delta K} \cdot W) |S\rangle + i \text{Sin}(\vec{\Delta K} \cdot W) |T0\rangle$$

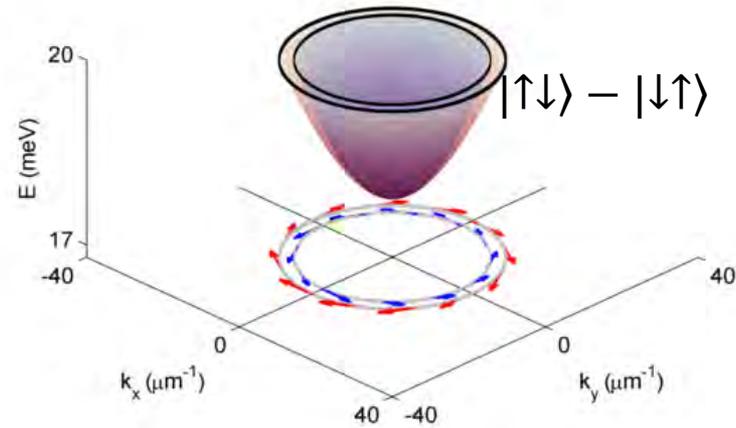
# Expectations with Rashba Spin-Orbit Interaction

B=0

Pairing with each Fermi surface

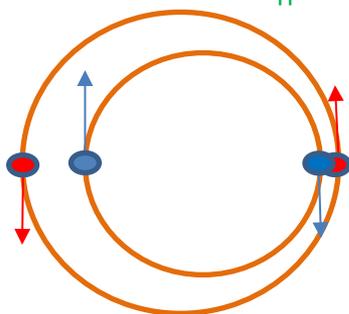


$$\vec{B} = (0 T, 0 T, 0 T)$$



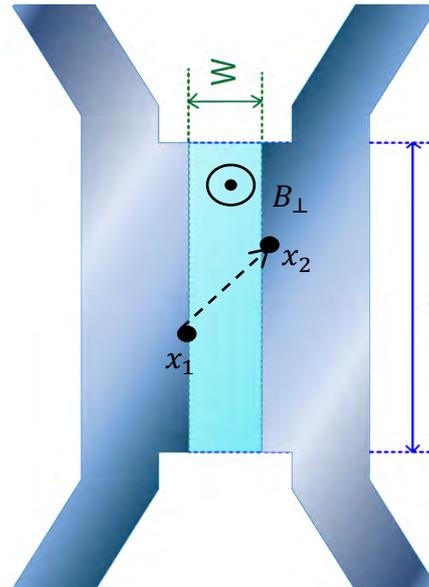
Some pairs acquire momentum  $\Delta K \perp B_{||}$

B>0



$$\Delta K = \frac{E_Z}{\hbar V_F}$$

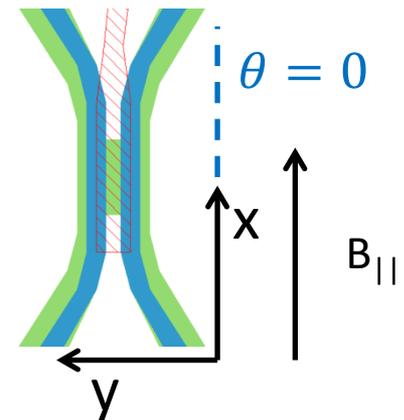
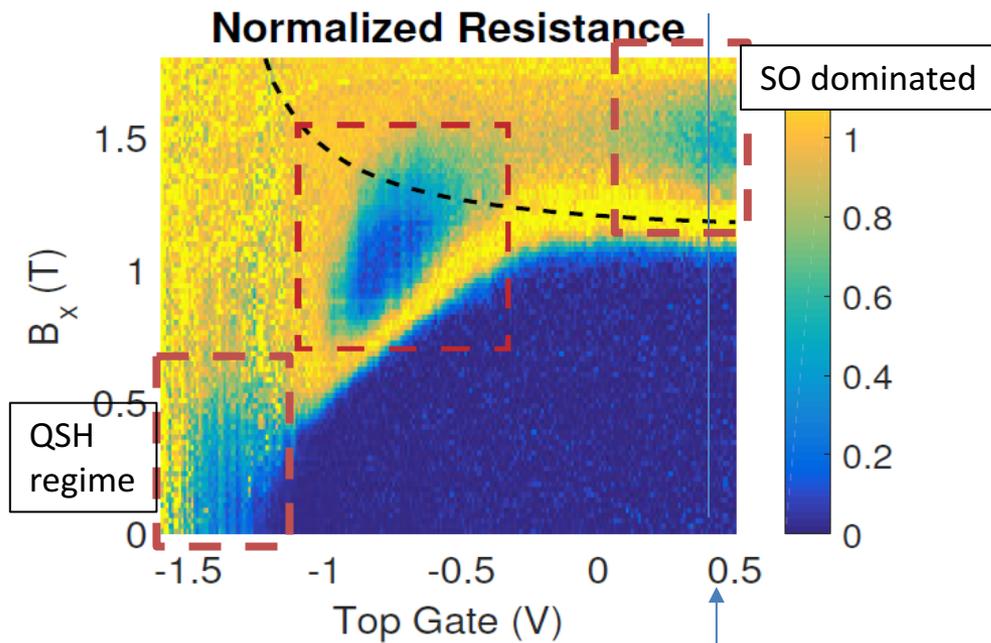
$$e^{i\vec{\Delta K} \cdot \vec{r}} |\uparrow\downarrow\rangle - e^{-i\vec{\Delta K} \cdot \vec{r}} |\downarrow\uparrow\rangle$$



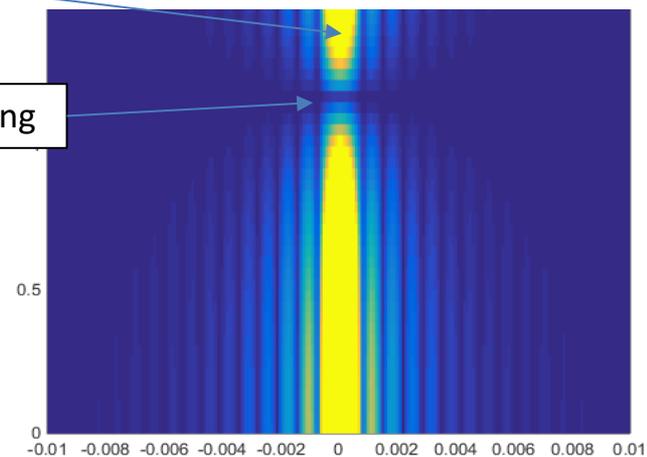
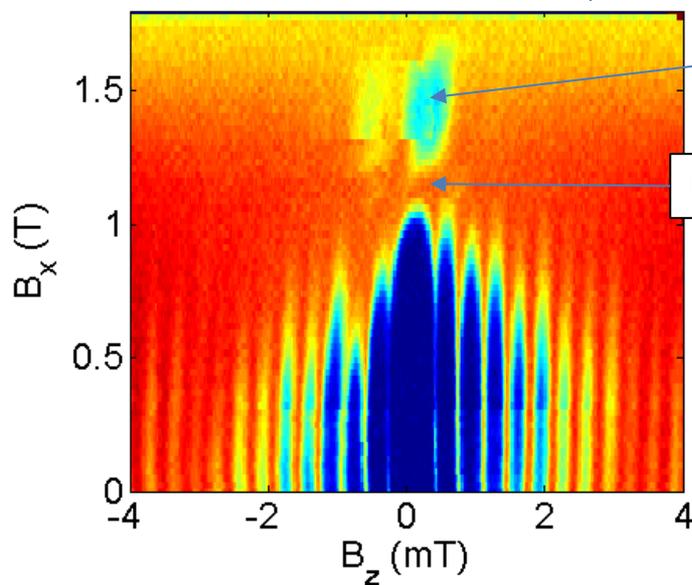
Summing over all trajectories

$$\Delta_{\text{ind}} = \text{Cos}(\vec{\Delta K} \cdot W) |S\rangle + i \text{Sin}(\vec{\Delta K} \cdot W) |T0\rangle$$

# Spin-Orbit Dominated Regime

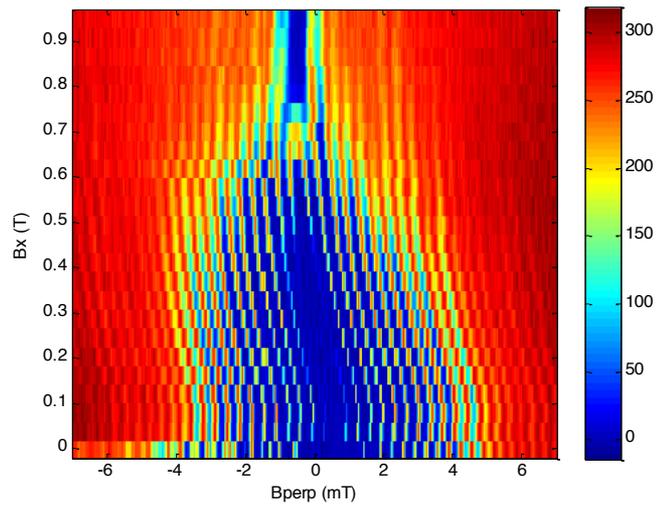
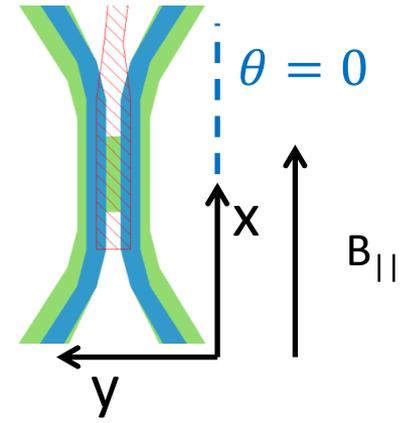
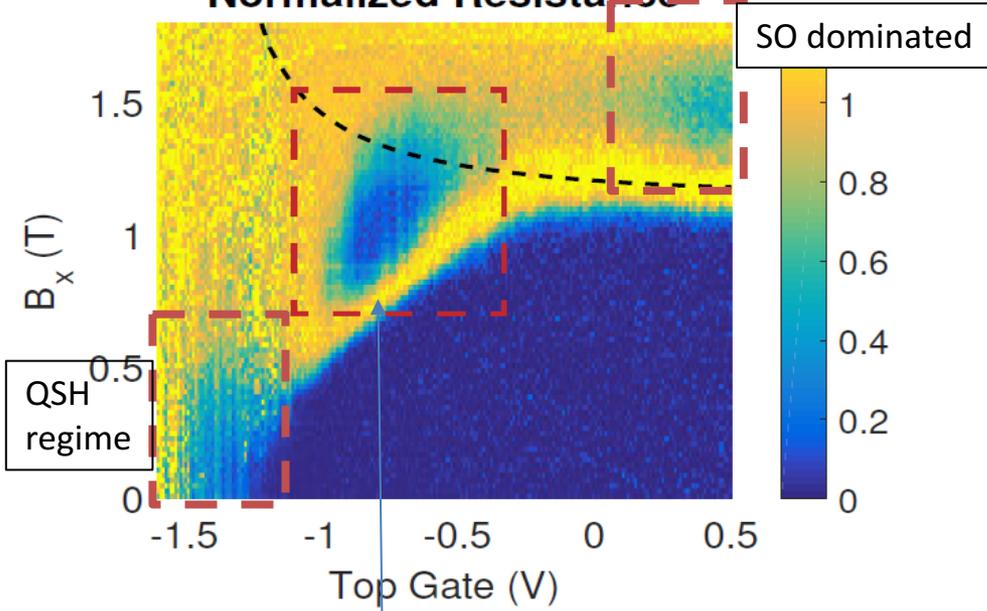


$$\Delta_{\text{ind}} = \text{Cos}(\vec{\Delta K} \cdot \vec{W}) |S\rangle + i \text{Sin}(\vec{\Delta K} \cdot \vec{W}) |T0\rangle$$



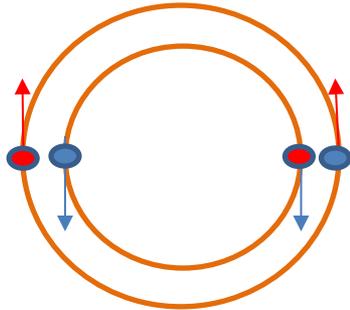
# Lower Density Regime

## Normalized Resistance



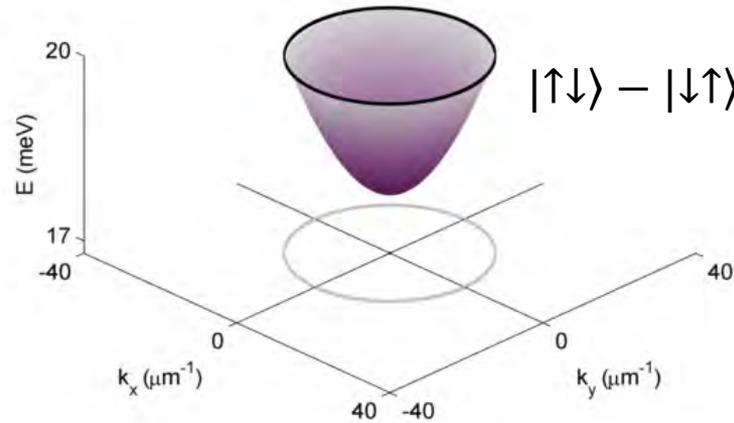
# Zeeman Only

All pairs acquire momentum



$$e^{-i\vec{\Delta K} \cdot \vec{r}} |\uparrow\downarrow\rangle - e^{i\vec{\Delta K} \cdot \vec{r}} |\downarrow\uparrow\rangle$$

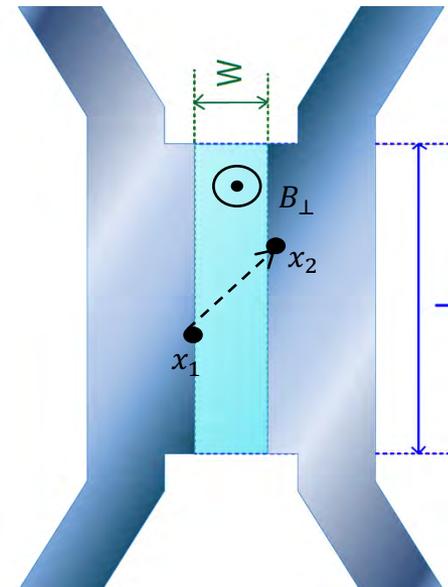
$$\vec{B} = (0 T, 0 T, 0 T)$$



Resembles S-F-S  
Junctions – tunable

‘Induced’ - FFLO

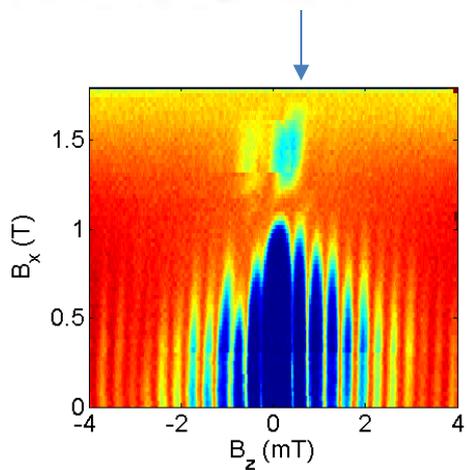
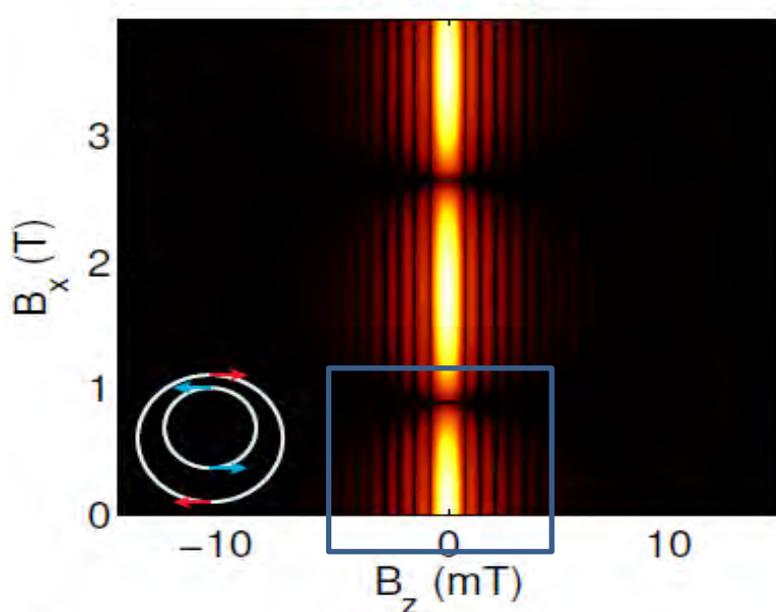
$$\Delta_{\text{ind}} = \text{Cos}(\vec{\Delta K} \cdot (\vec{X}_2 - \vec{X}_1)) |S\rangle + i \text{Sin}(\vec{\Delta K} \cdot (\vec{X}_2 - \vec{X}_1)) |T0\rangle$$



# Rashba vs Zeeman

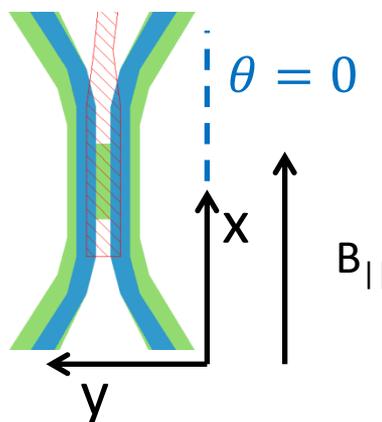
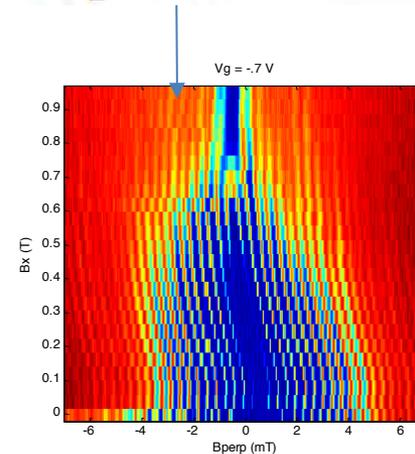
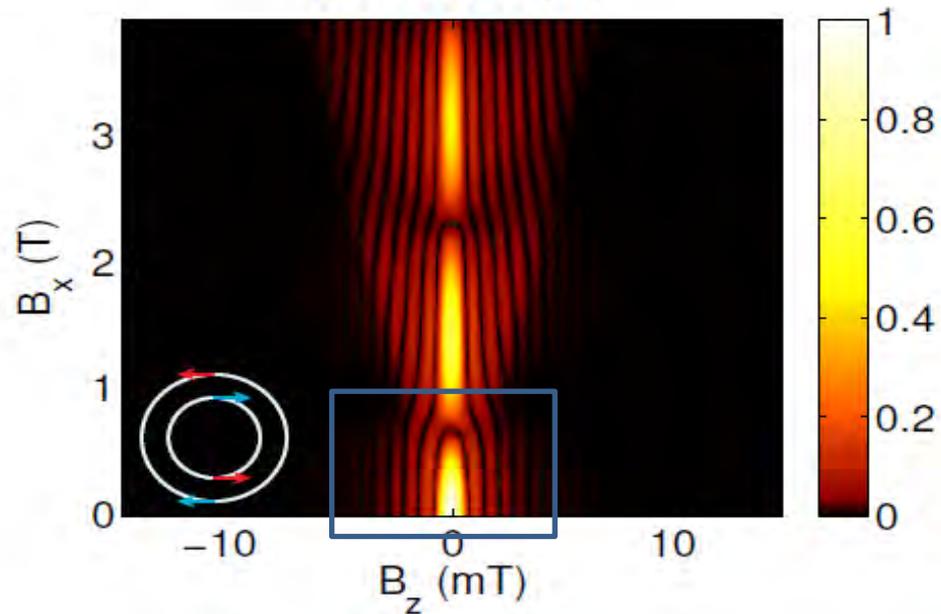
Rashba dominated

Critical Current



Zeeman only

Critical Current



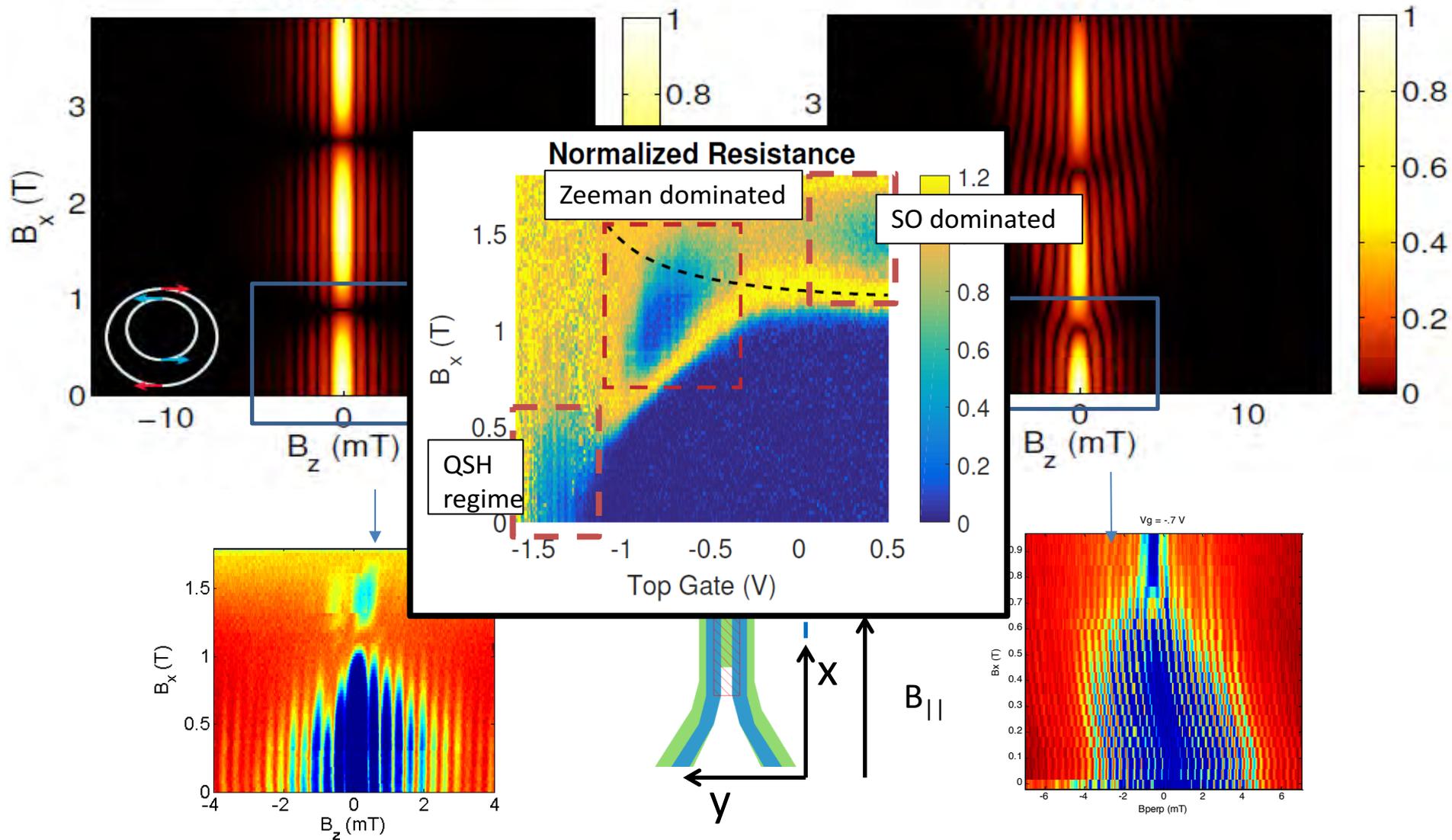
# Rashba vs Zeeman

Rashba dominated

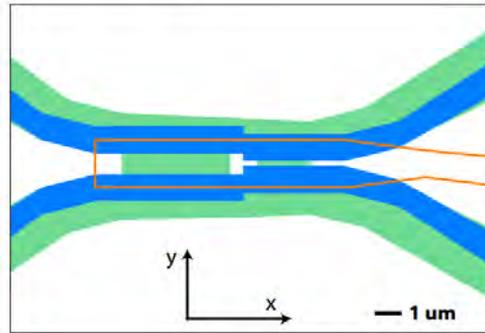
Zeeman only

Critical Current

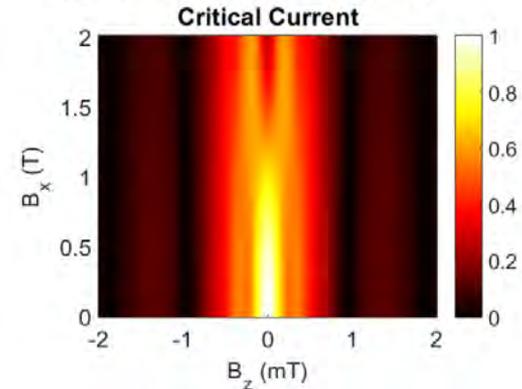
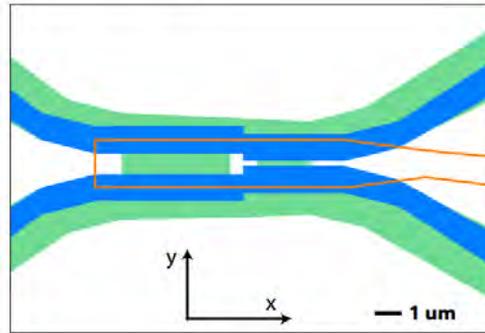
Critical Current



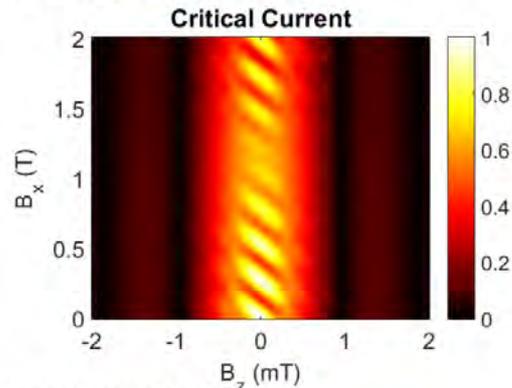
# $\pi$ phase shift and $\pi$ junction



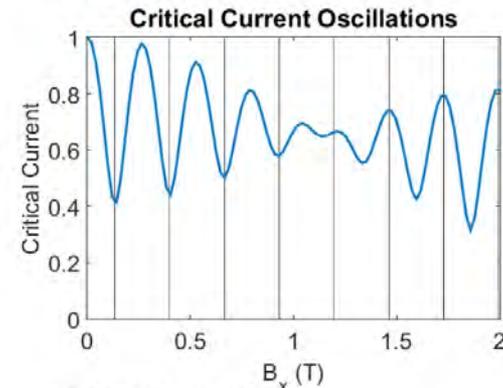
# $\pi$ phase shift and $\pi$ junction



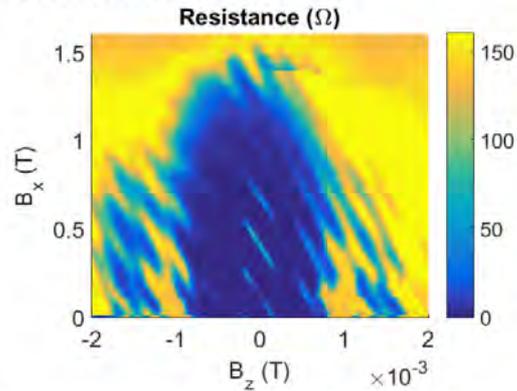
c. Simulation with Screening



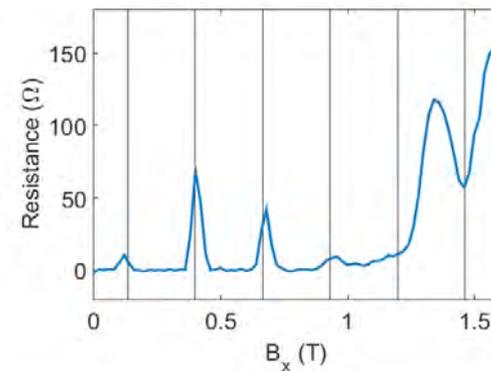
d. Simulation of Linecut



e. Resistance Data



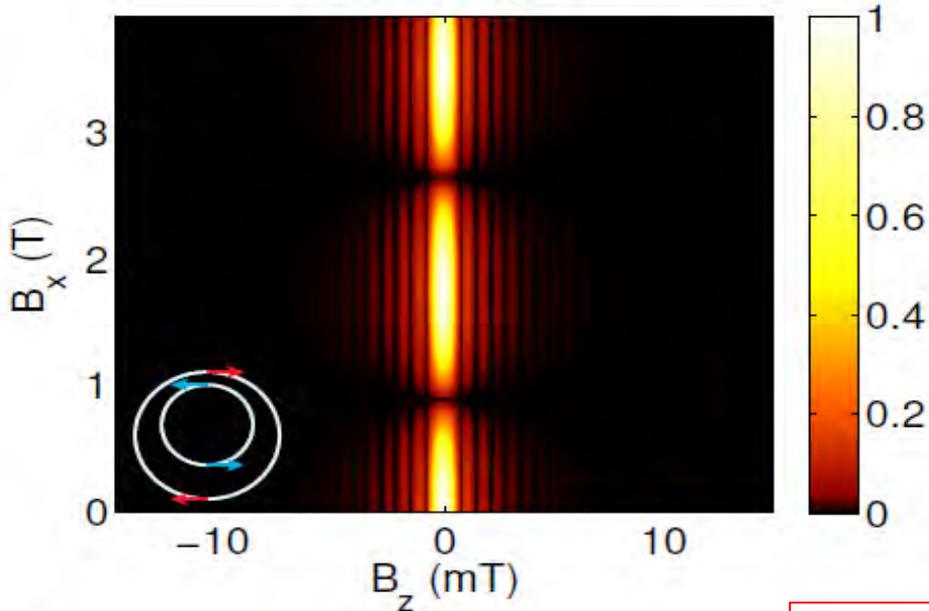
f. Linecut of Data



# Rashba vs Zeeman

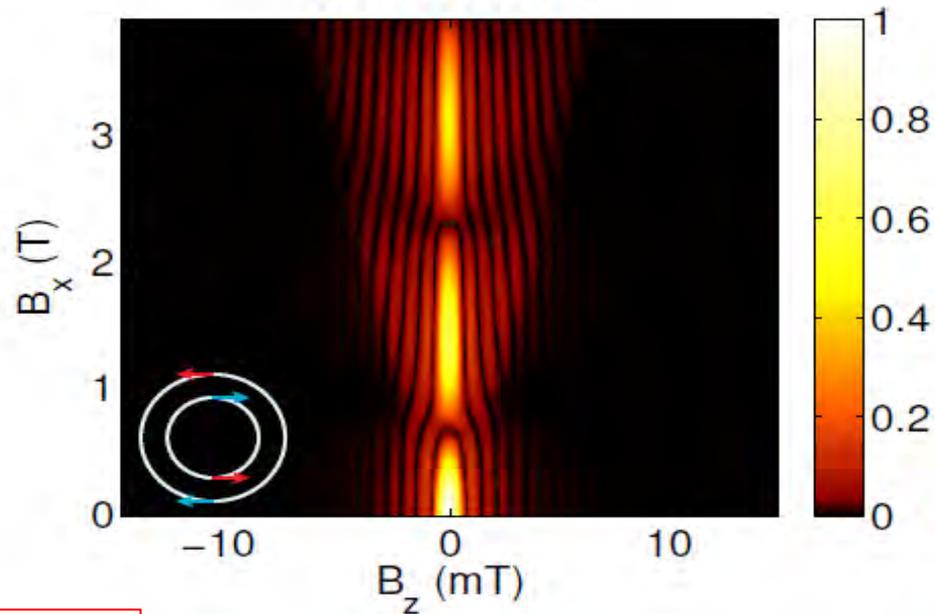
Rashba dominated

Critical Current



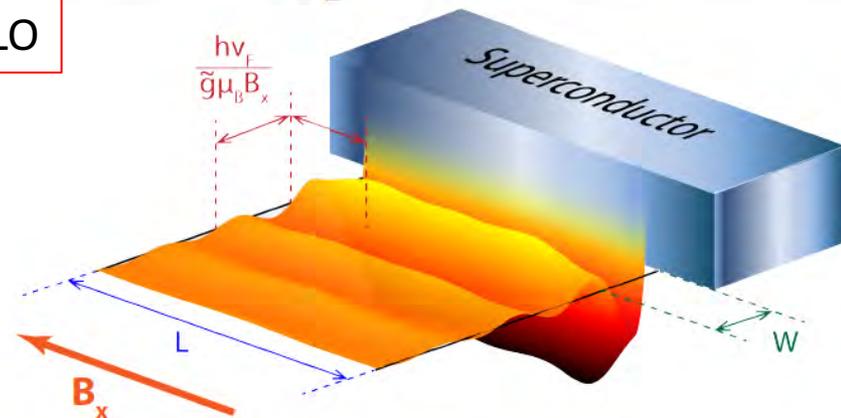
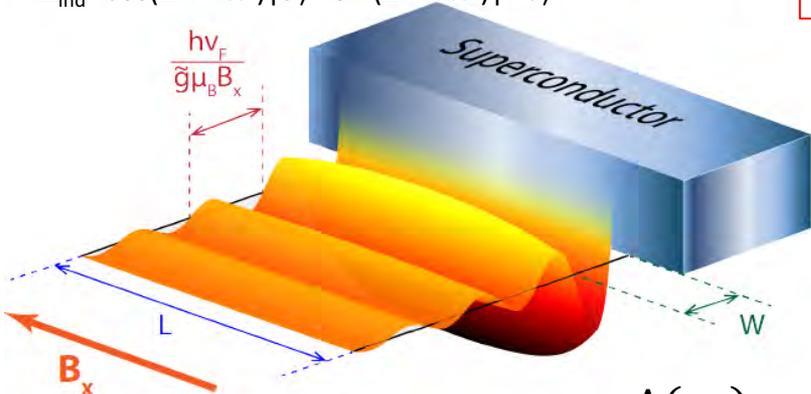
Zeeman Dominated

Critical Current



$$\Delta_{ind} = \text{Cos}(\Delta \vec{K} \cdot \vec{W}) |S\rangle + i \text{Sin}(\Delta \vec{K} \cdot \vec{W}) |T0\rangle$$

Induced FFLO

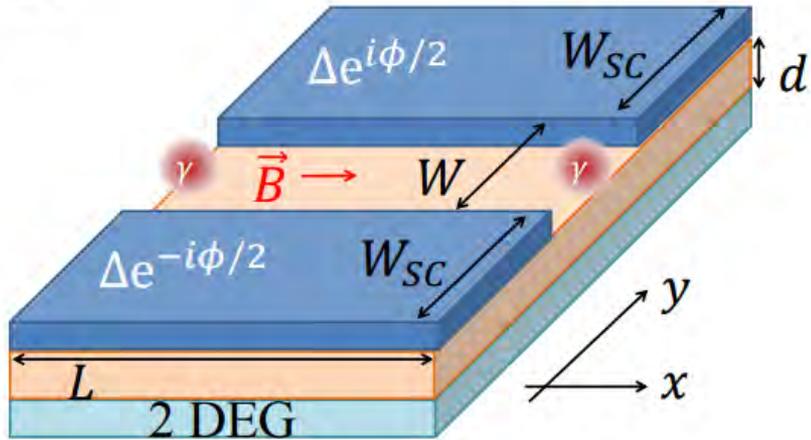


Decay independent of B

$$\Delta(x_2)_{ind} = \left| \int_{-L/2}^{L/2} dx_1 [\lambda_1(x_1) F_0(x_1, x_2)] \right|$$

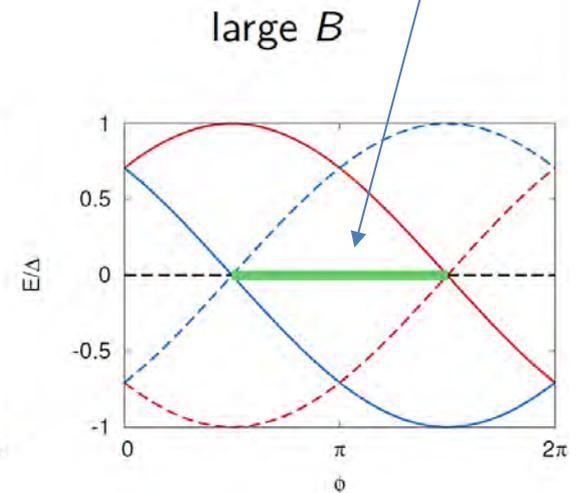
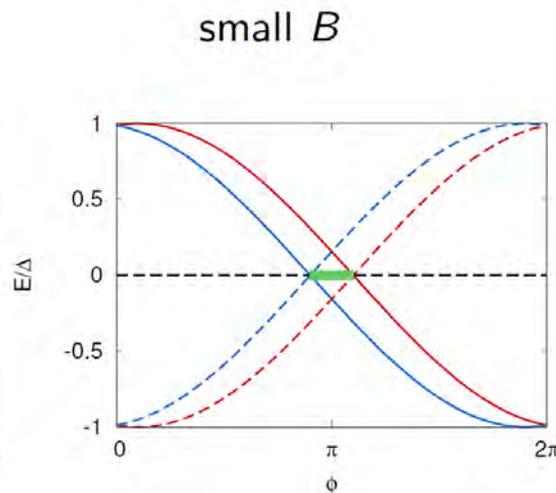
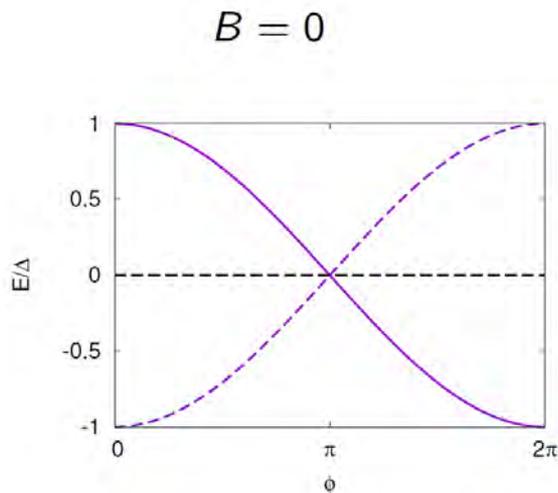
# New Theoretical Insight

Topological transition only governed by  $Kx=0$



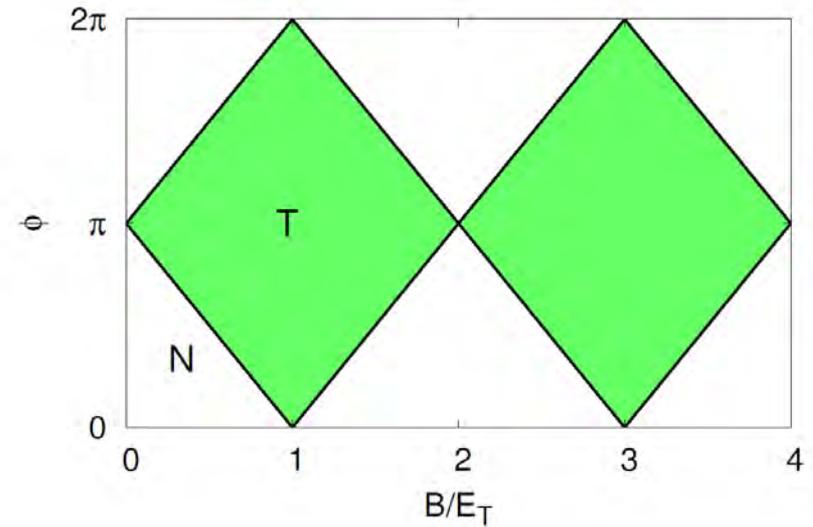
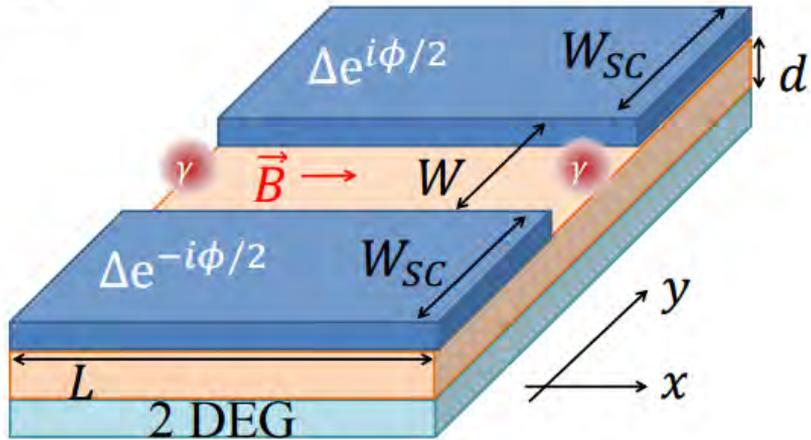
- Narrow junction
- Neglect normal reflections

- Topological superconductor



# Topological Switch

Topological transition only governed by  $Kx=0$



- $\phi = 0$  system *always* trivial
- $\phi = \pi$  system *always* topological

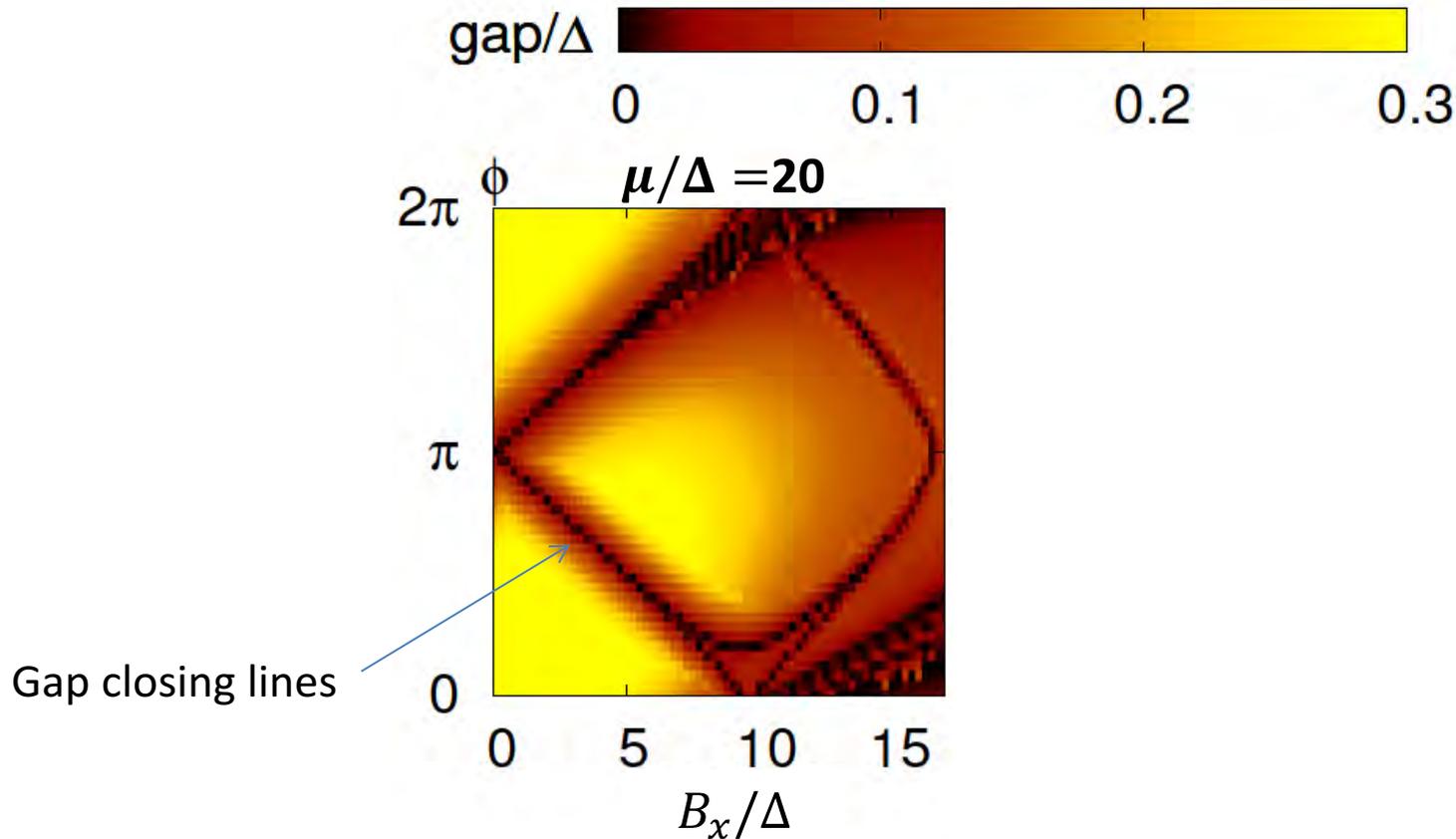
topology at  $\phi = 0, \pi$   
independent of  $\mu, B, \Delta$  and  $W$



topology switch

# Magnitude of the Gap

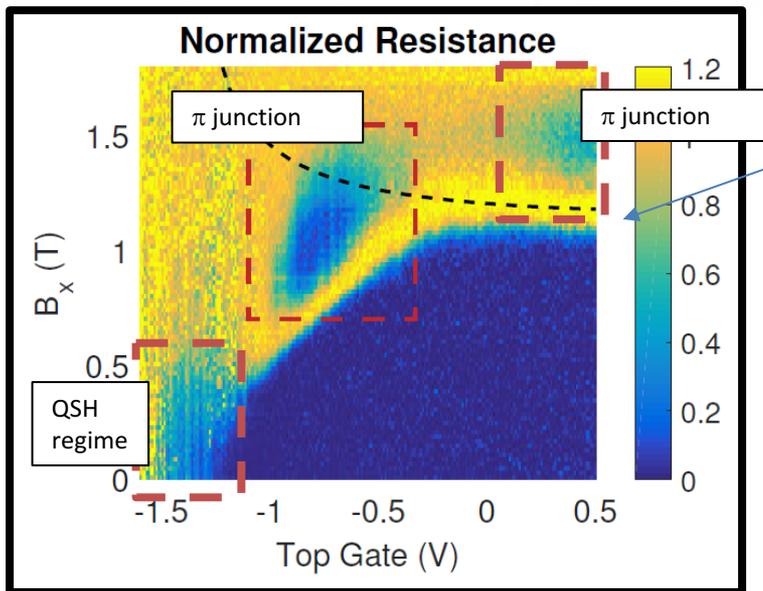
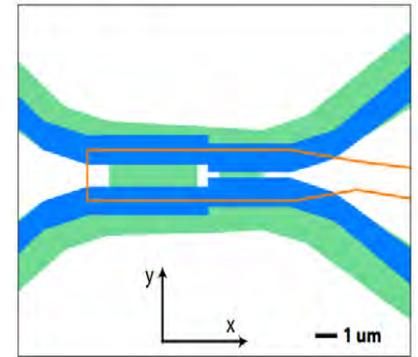
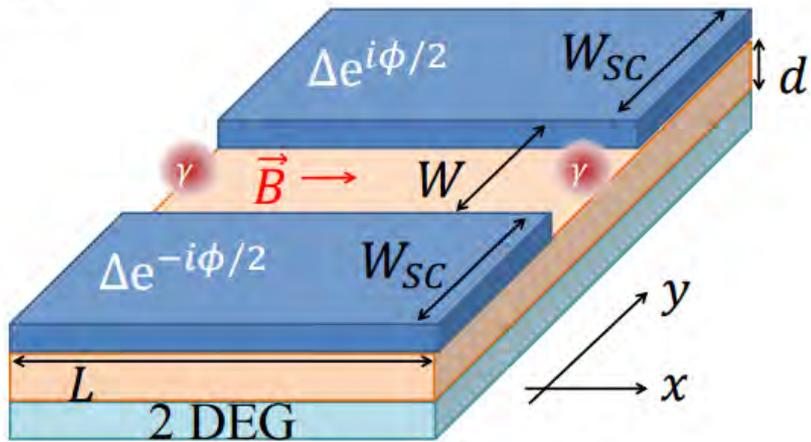
Gap obtained numerically using scattering matrix approach,  $\Delta = 1/mW^2$



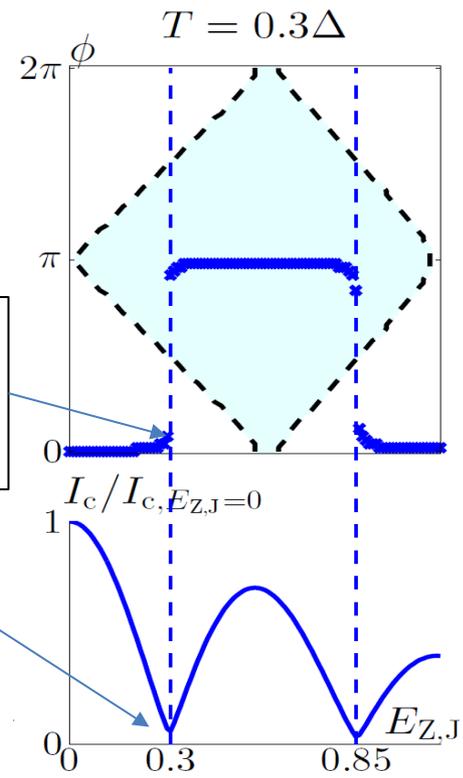
Gap of order  $\Delta$  in the topological region!

Size of the gap is insensitive to  $\mu$ !

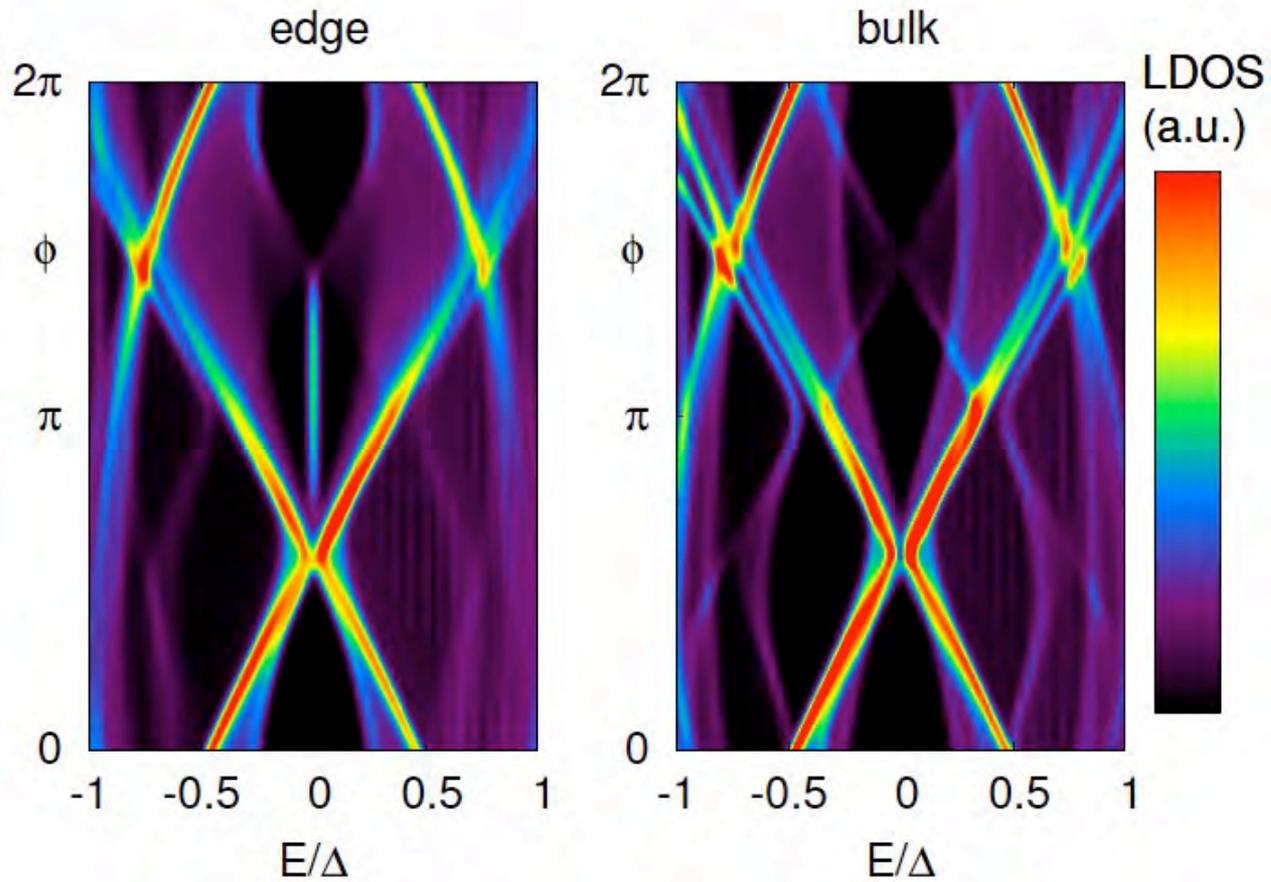
# Self Tuning into Topological Phase



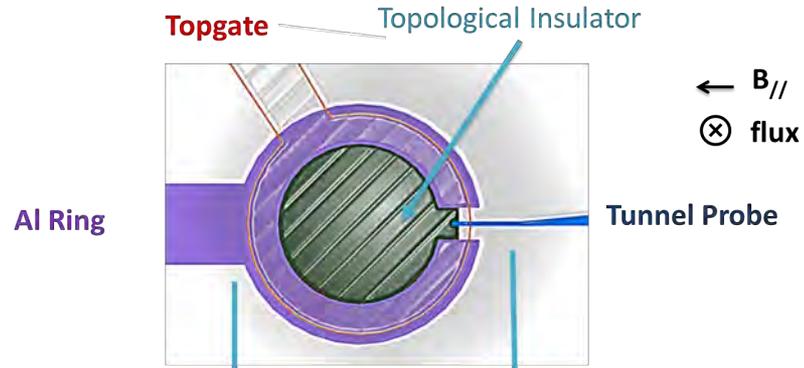
First order topological phase transition



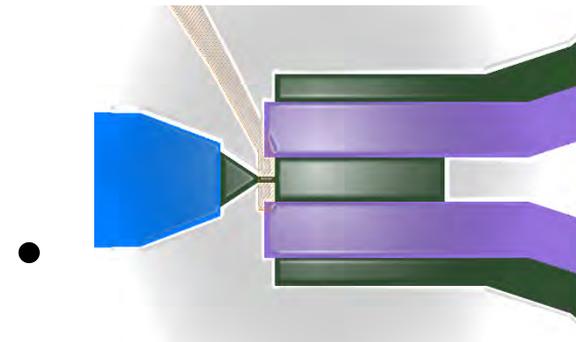
# Search for Majorana End States



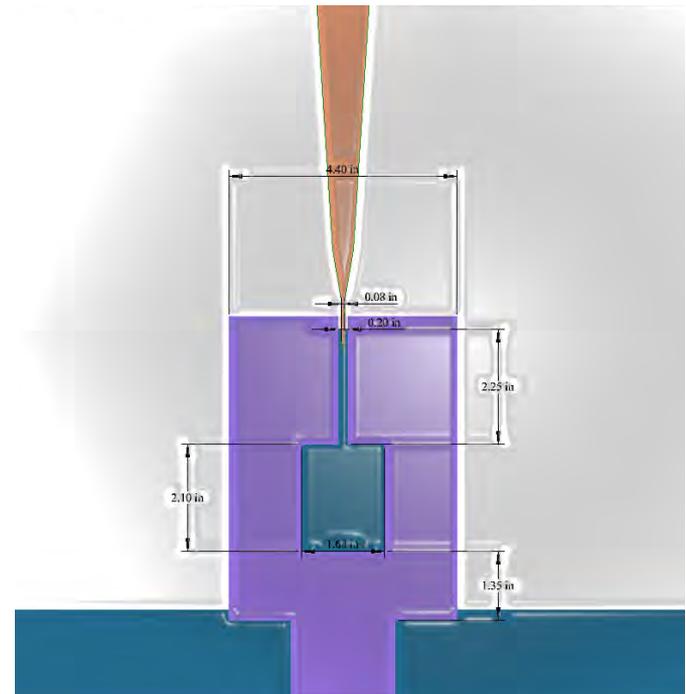
# Tunneling Spectroscopy



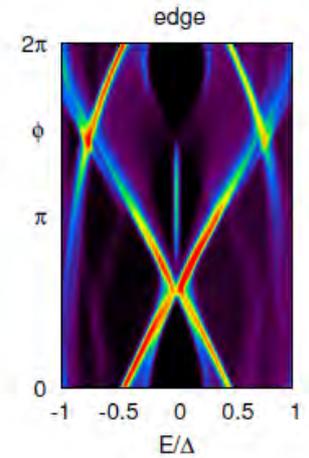
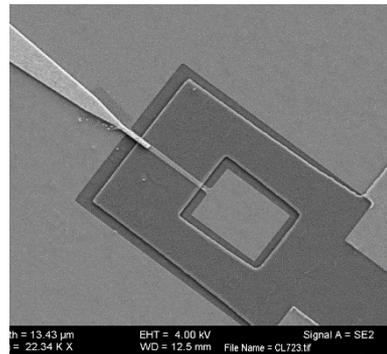
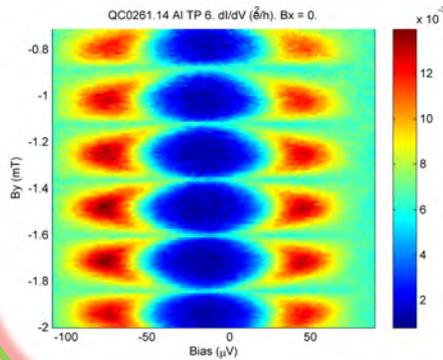
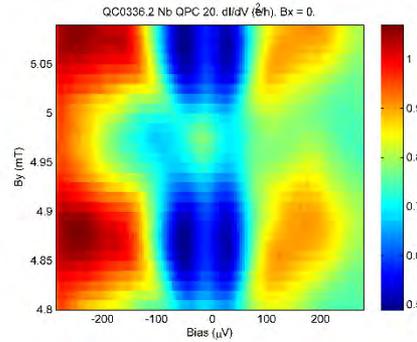
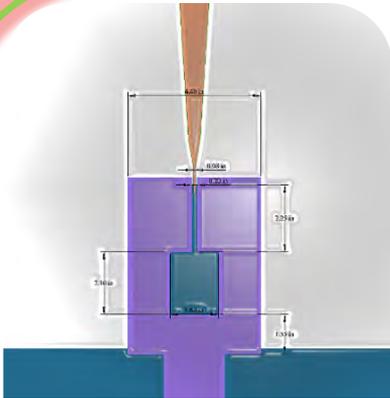
Measure differential conductance between tunnel probe and superconductor



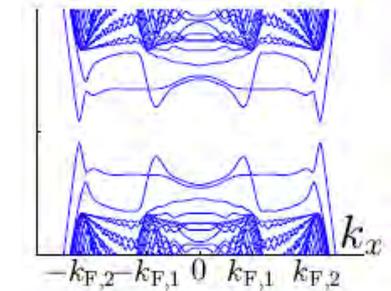
Flux control – ‘a Topology Switch’



# Local Density of States Measurement and Calculation

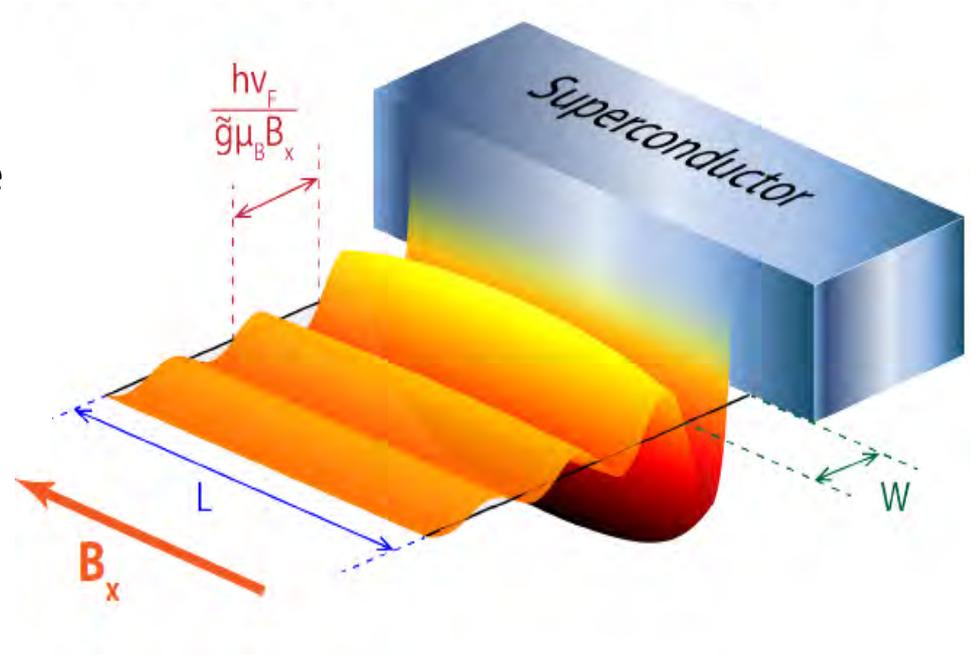


c)  $\phi = \pi, E_{Z,J} = 0.5$



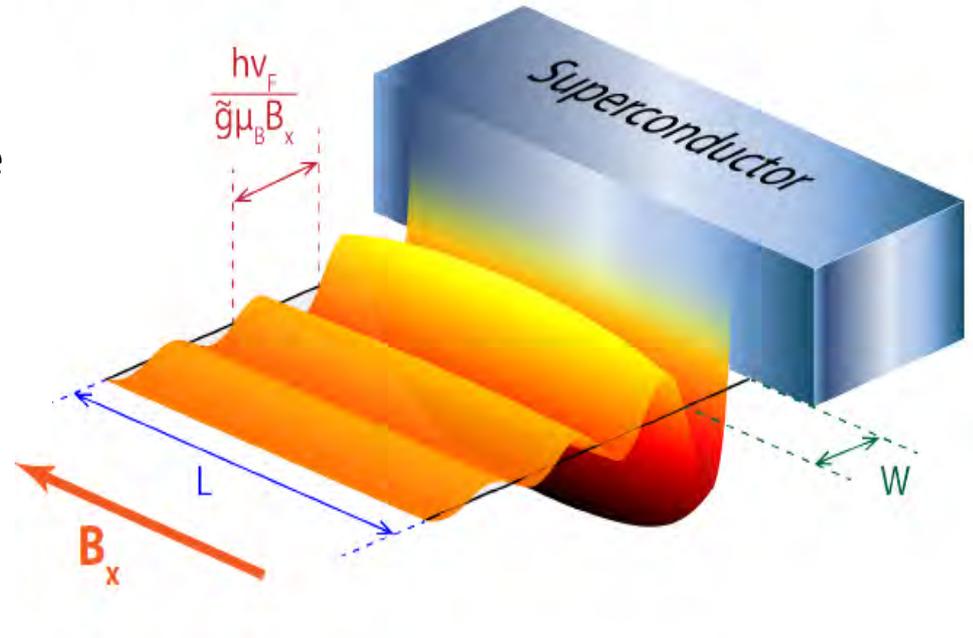
## Outlook:

- Can we use this measurement approach to determine the presence of p-type order?
- Can we determine what is the underlying p-type order (T+, T-, To)?

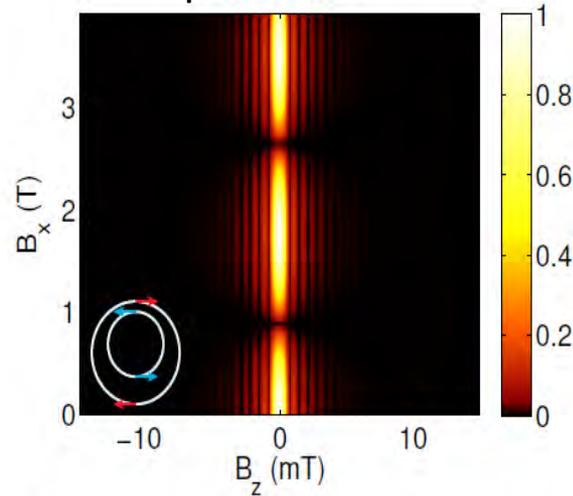


## Outlook:

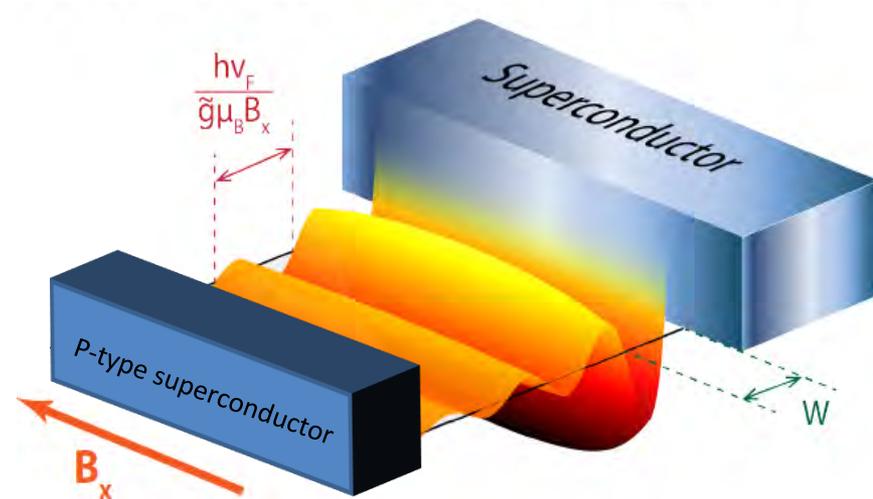
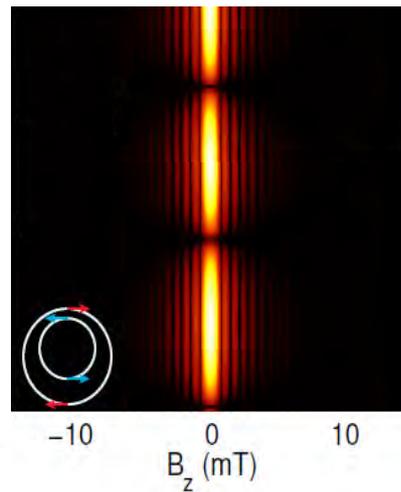
- Can we use this measurement approach to determine the presence of p-type order?
- Can we determine what is the underlying p-type order ( $T+$ ,  $T-$ ,  $T_0$ )?



### S-superconductor



### P-superconductor

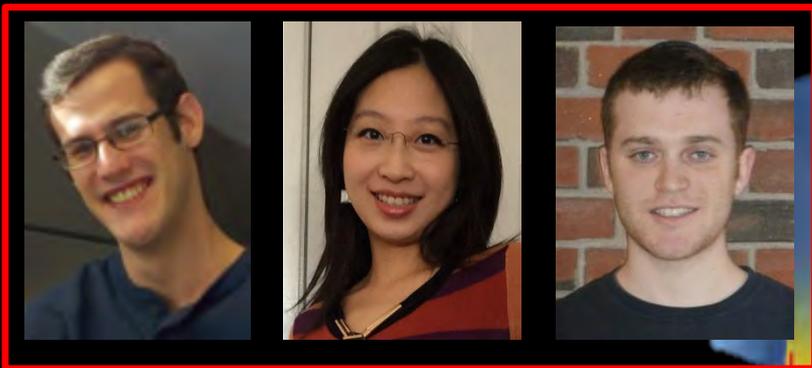


# A New Spin on Superconductivity

Sean Hart, Hechen Ren, Michael Kosowsky, Bert Halperin, Harvard University

L. Molenkamp's group, University of Wurzburg

Triplet pairing



In collaboration with:  
F. Pientka, A. Keselman, A. Stern,  
E. Berg

