

Magnon Transport Both in Ferromagnetic and Antiferromagnetic Insulating Magnets

Kouki Nakata

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KN, S. K. Kim (UCLA), J. Klinovaja, D. Loss (2017)

arXiv:1707.07427

See also review article [KN, Simon (Paris) & Loss, J. Phys. D (2017)] with [KN, JK & DL, PRB (2017)] on magnonic “quantum” Hall effect and Wiedemann-Franz law in a topological ferromagnet

Magnonic Topological Insulator in Antiferromagnet

Kouki Nakata

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KN, S. K. Kim (UCLA), J. Klinovaja, D. Loss (2017)

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See also review article [KN, Simon (Paris) & Loss, J. Phys. D (2017)] with [KN, JK & DL, PRB (2017)] on magnonic “quantum” Hall effect and Wiedemann-Franz law in a topological ferromagnet

GOAL

Topological insulator

Kane & Mele, PRL (2005, 2005).

Bernevig & Zhang, PRL (2006).

Magnonic Topological insulator: A bosonic analog

KN, Kim, Klinovaja & Loss (2017).

arXiv:1707.07427

BASIC IDEA

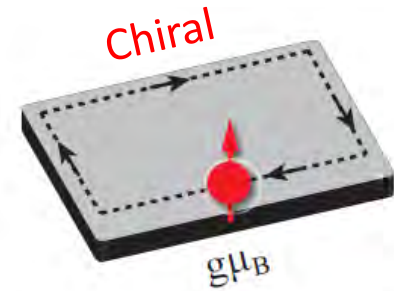
Antiferromagnet (AF): Néel Order

AF = Independent copies of FM

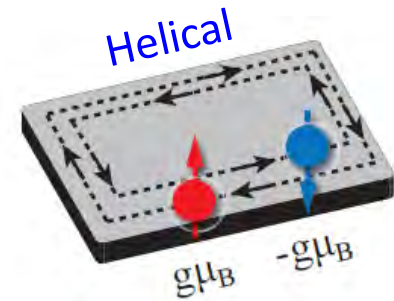
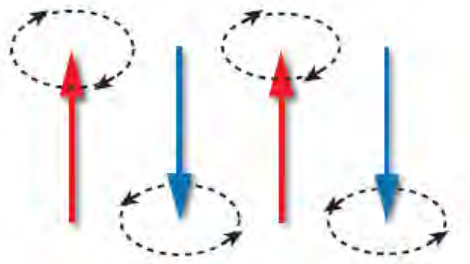
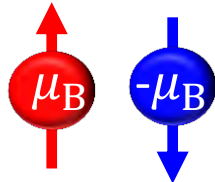
Anderson, Phys. Rev. (1952)
Kubo, Phys. Rev. (1952)

$\sigma g\mu_B$: Up- ($\sigma = 1; \uparrow$) & down- ($\sigma = -1; \downarrow$) magnons \rightarrow Direction of cyclotron motion; **opposite**

FM:



AF:



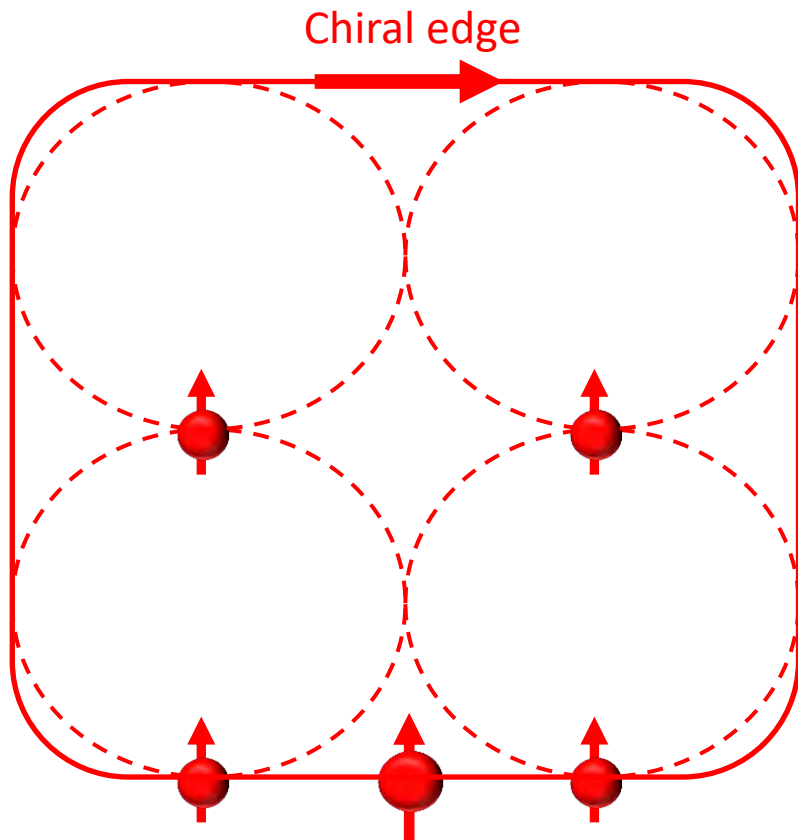
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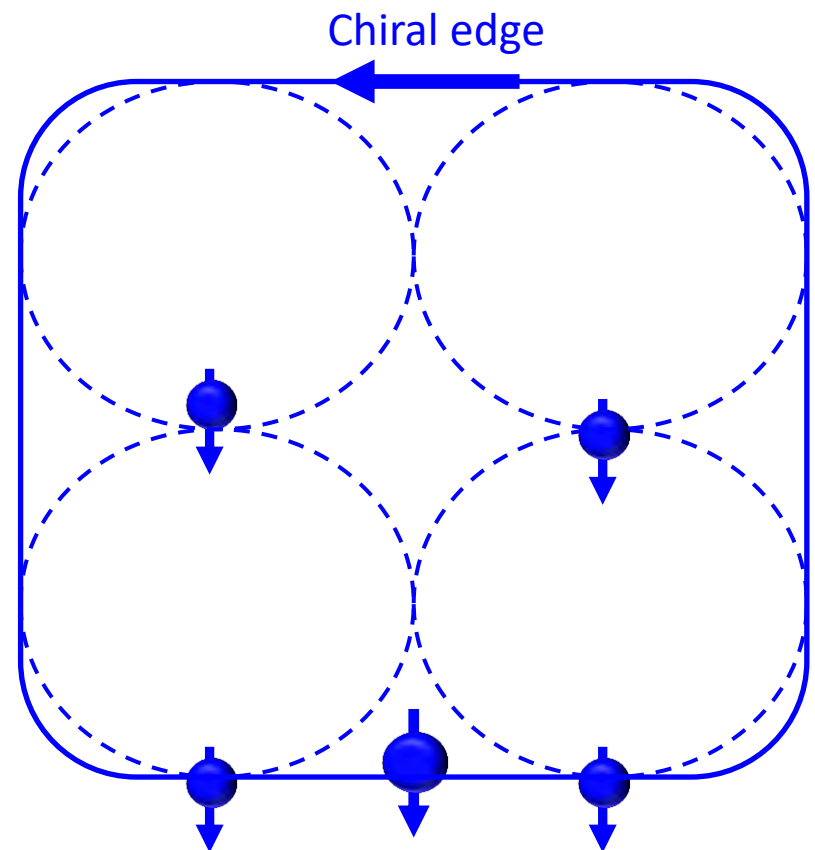
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Up-magnon: $\sigma = 1; \uparrow$



Down-magnon: $\sigma = -1; \downarrow$



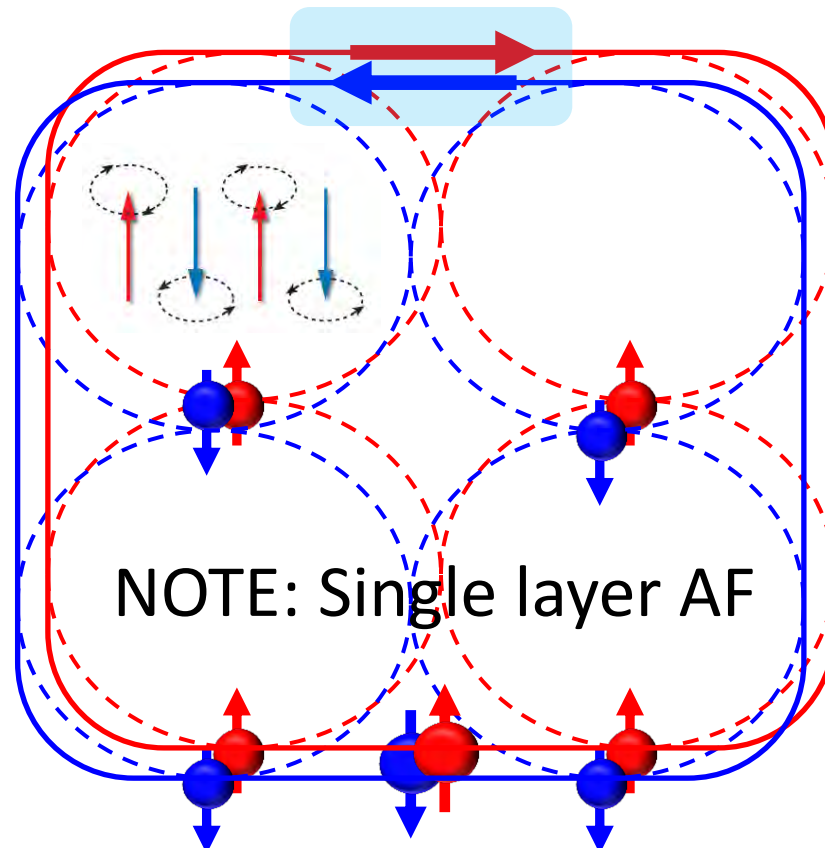
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(Single layer) AF: Helical edge magnon



Antiferromagnet (AF): Néel Order

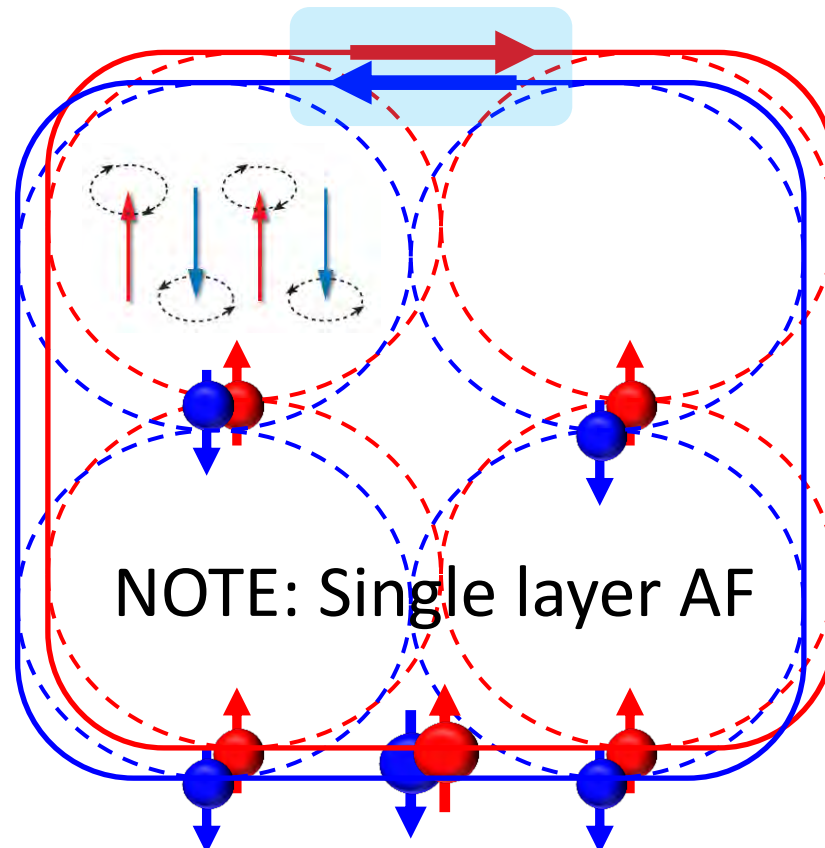
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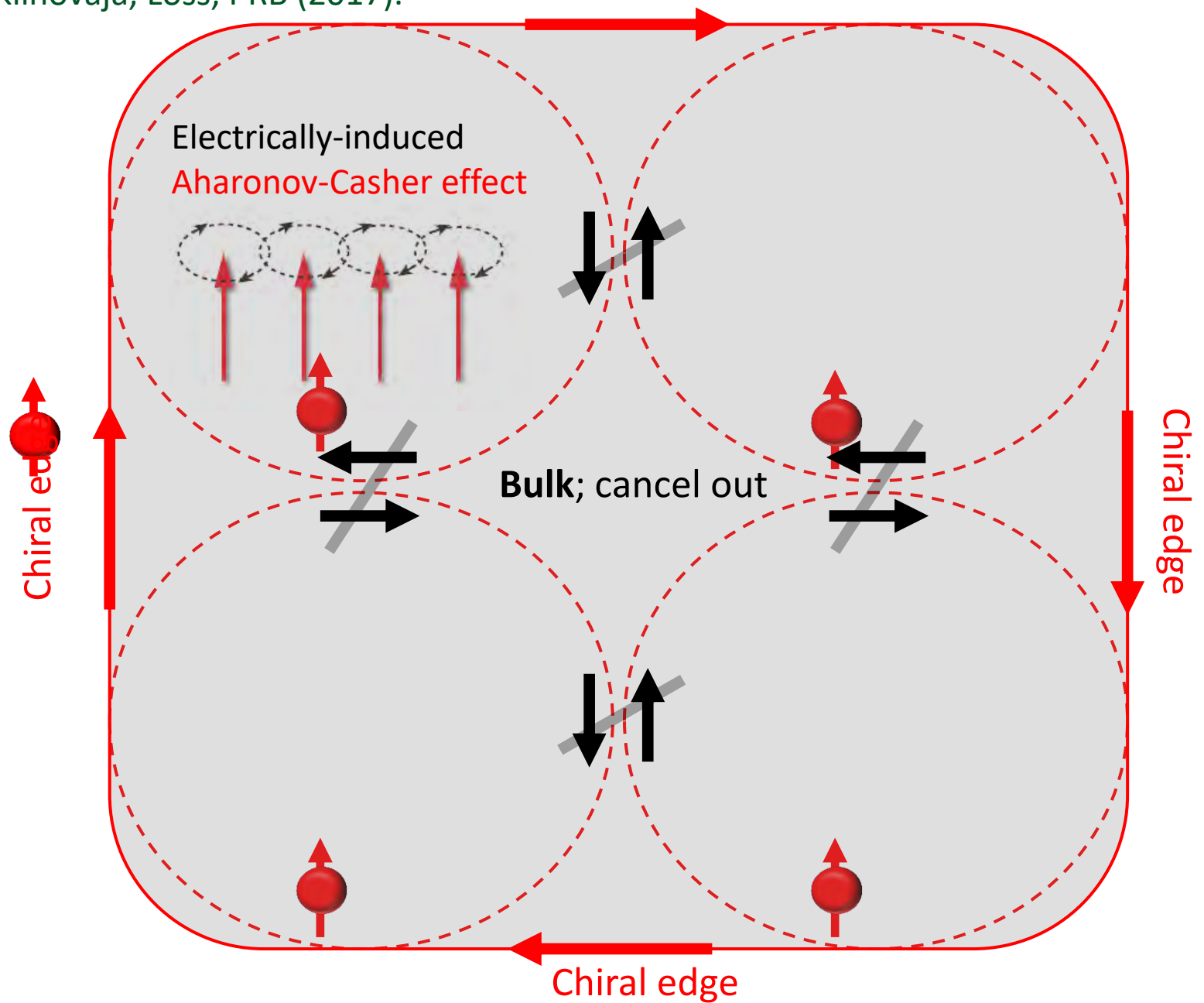
- Q. How to realize such a cyclotron motion of each magnon ?
A. Aharonov-Casher (AC) effect on magnons in electric field

(Single layer) AF: Helical edge magnon

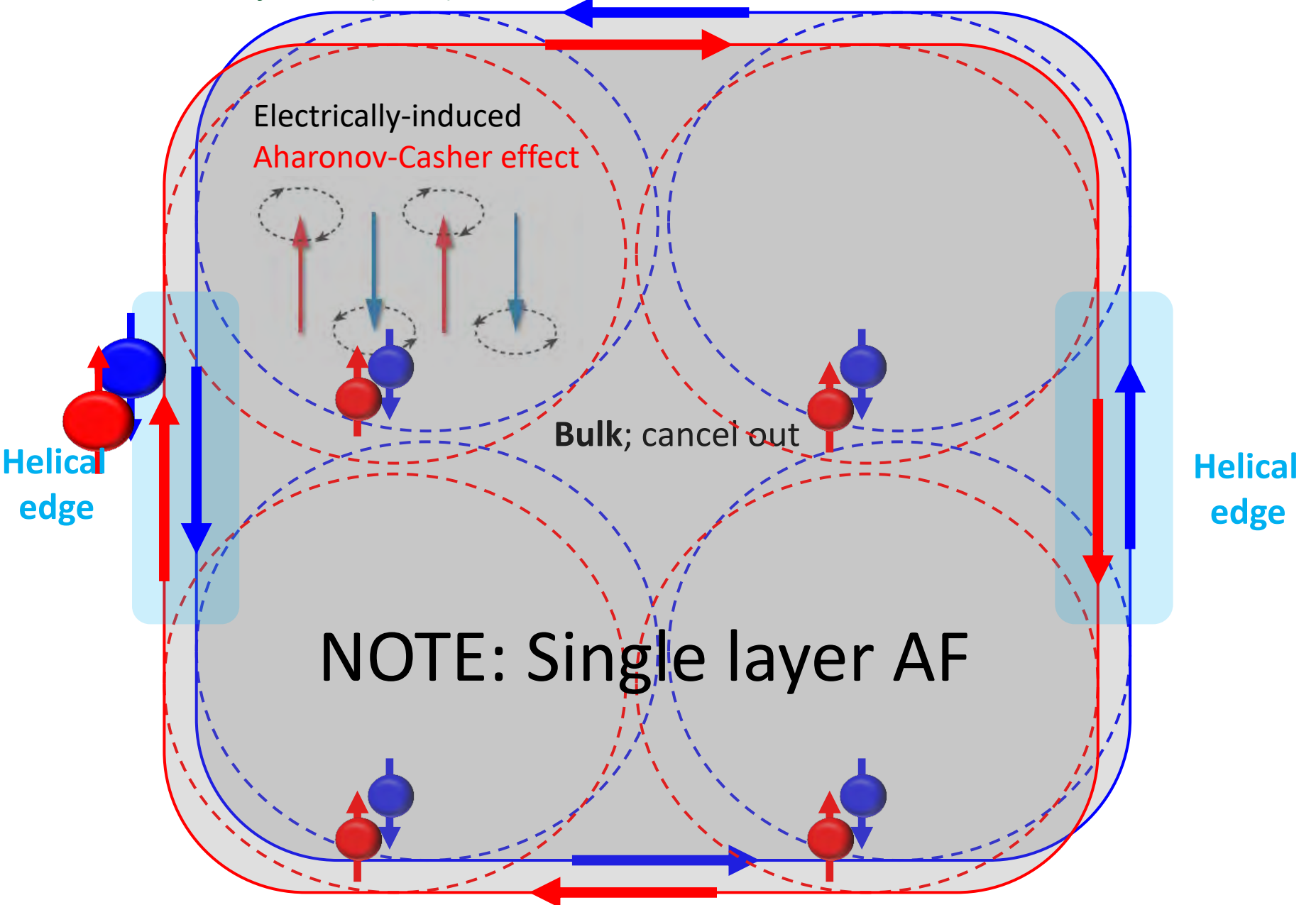


“Quantum” Hall effect of magnon in ferromagnet (FM): Chiral edge by AC effect

KN, Klinovaja, Loss, PRB (2017).



“Quantum” spin Hall effect of magnons in AF (single layer): **Helical edge** by AC effect
KN, Kim, Klinovaja, Loss (2017). = **Magnonic topological insulator (TI)**



PURPOSE OF TALK

To share the “qualitative” understanding of such mechanism in topological AFs:

See [KN, Kim, Klinovaja, Loss, arXiv:1707.07427] for details

- Q. How to realize such a 2-dim (single layer) **topological AF** ???
- A. **Aharonov-Casher effect** on magnons in electric field

Magnonic topological insulator (TI) in AF

1) Helical edge magnon state

2) Topological invariant:

Topological Hall effect of bulk magnons
(*e.g.*, spin, thermal, Nernst & Ettinghausen effects)

3) Bosonic Wiedemann-Franz (WF) law:

Universal thermomagnetic properties

OUTLINE OF TALK

I. Topological FM

KN, Klinovaja & Loss, PRB (2017)

Chiral edge

II. Topological AF

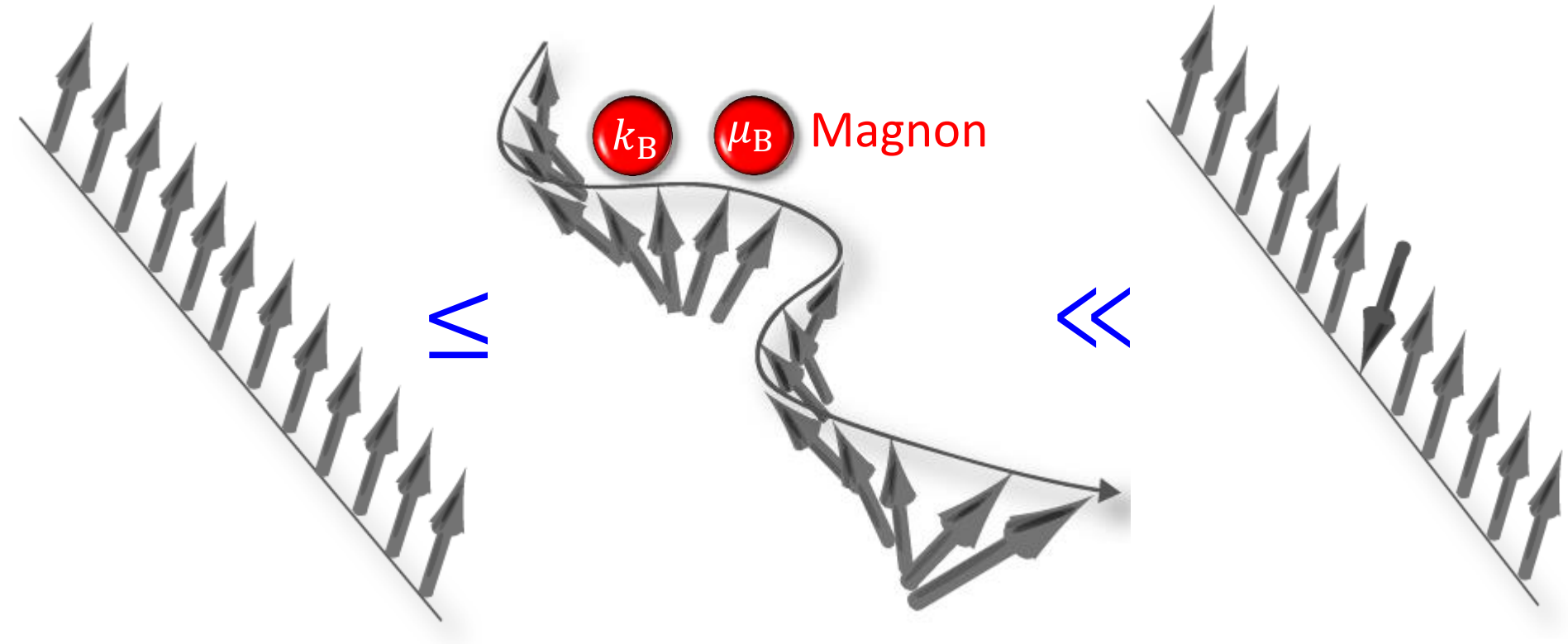
KN, Kim, Klinovaja & Loss (2017)

Helical edge

Let's start !!

Magnon Carries μ_B & k_B

Low-energy collective excitation in insulating FM & AF



QUESTION

Q. Can magnon μ_B (boson) transport be similar to electron e (fermion) transport ?

Yes !! KN, Simon & Loss, J. Phys. D (2017): Review article

Electron e
Fermion

Magnon μ_B
Boson

Wiedemann-Franz (WF) law

Franz & Wiedemann, Annalen der Physik (1853)

Magnonic Wiedemann-Franz law

KN, Simon & Loss, PRB (2015)

Josephson effect

Josephson, Phys. Lett. (1962)

Magnonic Josephson effect

KN, Hoogdalem, Simon & Loss, PRB (2014)

KN, Simon & Loss, PRB (2015)

Integer quantum Hall effect (IQHE)

Klitzing *et al.*, PRL (1980)

TKNN, PRL (1982) / Kohmoto, Ann. Phys. (1985)

Magnonic ``quantum'' Hall effect: Chiral edge

KN, Klinovaja & Loss, PRB (2017)

Topological ferromagnet (FM)

Topological insulator (TI)

Kane & Mele, PRL (2005, 2005).

Bernevig & Zhang, PRL (2006).

Quantum spin Hall effect (QSHE)

Magnonic topological insulator: Helical edge

KN, Kim, Klinovaja & Loss (2017) arXiv:1707.07427

Topological antiferromagnet (AF):

Magnonic quantum spin Hall effect

QUESTION

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Magnon μ_B
Boson

Magnonic Wiedemann-Franz law

KN, Simon & Loss, PRB (2015)

Magnonic Josephson effect

KN, Hoogdalem, Simon & Loss, PRB (2014)

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Magnonic ``quantum'' Hall effect: Chiral edge

KN, Klinovaja & Loss, PRB (2017)

Topological ferromagnet (FM)

Magnonic topological insulator: Helical edge

KN, Kim, Klinovaja & Loss (2017) arXiv:1707.07427

Topological antiferromagnet (AF):

Magnonic quantum spin Hall effect

Aharonov-Casher (AC) phase:

Aharonov & Casher, PRL (1984)

Observation for magnons:

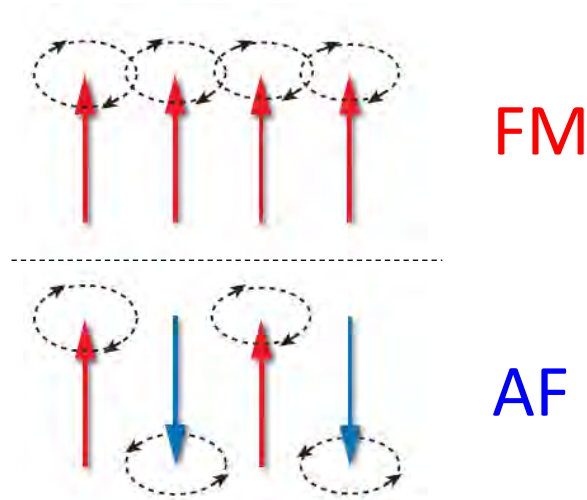
Zhang *et al.* (Yale), PRL (2014)

Electric-field coupling to magnons

QUESTION

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Magnon μ_B
Boson

Magnonic Wiedemann-Franz law

KN, Simon & Loss, PRB (2015)

Magnonic Josephson effect

KN, Hoogdalem, Simon & Loss, PRB (2014)

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Magnonic "quantum" Hall effect: Chiral edge

KN, Klinovaja & Loss, PRB (2017)

Topological ferromagnet (FM)

Magnonic topological insulator: Helical edge

KN, Kim, Klinovaja & Loss (2017) arXiv:1707.07427

Topological antiferromagnet (AF):

Magnonic quantum spin Hall effect

Topological FM

KN, Klinovaja & Loss, PRB (2017)

STRATEGY

Magnonic classical Hall effect in Aharonov-Casher phase: No topological edge

Meier & Loss, PRL (2003)



Topological invariant: Chern (TKNN) integer
TKNN, PRL (1982)
Kohmoto, Ann. Phys. (1985)

Magnonic “quantum” Hall effect in topological FM: Chiral edge

KN, Klinovaja & Loss, PRB (2017)



\mathbb{Z}_2 topological number
Hasan & Kane, RMP (2010) *etc.*

Magnonic “quantum” spin Hall effect in topological AF: Helical edge

KN, Kim, Klinovaja & Loss (2017)

Geometric Phases

<h2>Aharonov-Bohm (AB) phase</h2> <p>Aharonov & Bohm, Phys. Rev. (1959)</p>	<h2>Aharonov-Casher (AC) phase</h2> <p>Aharonov & Casher, PRL (1984) Mignani, J. Phys. A (1991). Meier & Loss, PRL (2003).</p>
<p>(Electrically) charged particle: e</p>	<p>Magnon = Magnetic dipole: $\vec{\mu} = g\mu_B \mathbf{e}_z$</p>
<p>Magnetic field coupling with electrons:</p> $\theta_{AB} = \frac{e}{\hbar c} \oint d\mathbf{l} \cdot \mathbf{A}$	<p>Electric field coupling with magnons:</p> $\theta_{AC} = \frac{g\mu_B}{\hbar c^2} \oint d\mathbf{l} \cdot (\mathbf{E} \times \mathbf{e}_z) = \frac{g\mu_B}{\hbar c} \oint d\mathbf{l} \cdot \mathbf{A}_m$
<p>Magnetic vector potential \mathbf{A}</p> $\nabla \times \mathbf{A} = \mathbf{B}$	<p>Electric vector potential: $\mathbf{A}_m \equiv \mathbf{E} \times \mathbf{e}_z / c$</p> $\nabla \times \mathbf{A}_m = \varepsilon \mathbf{e}_z / c \quad \text{for } \mathbf{E} = \varepsilon(-x, 0, 0)$

Observation of AC effect on magnons:
 Zhang *et al.* (Yale), PRL (2014)
 Electric-field coupling to magnons

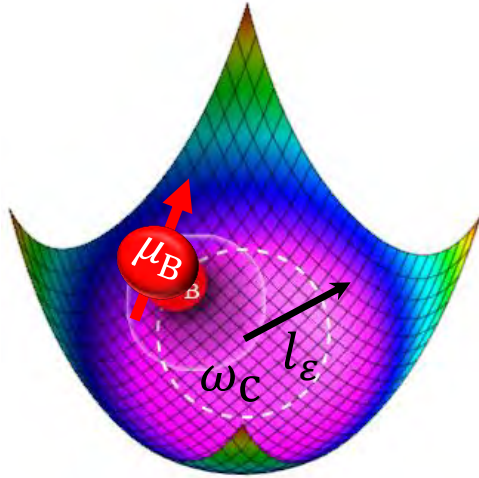
Hoogdalem *et al.*, PRB (2013) [Katsura *et al.*, PRL (2005)]:
 DM int. \rightarrow An analogue of artificial gauge field
 \rightarrow AC effect

FM: Landau Level in AC Effect

KN, Klinovaja & Loss, PRB (2017)

External electric field gradient \mathcal{E}

KN, Klinovaja & Loss, PRB (2017)



AC effect-induced topological FM:

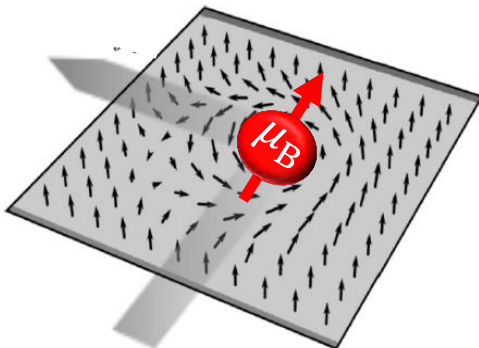
$$\nabla \times \mathbf{A}_m = \frac{\mathcal{E}}{c} \mathbf{e}_z \quad \left\{ \begin{array}{l} \mathbf{E}(\mathbf{r}) = \mathcal{E}(-x, 0, 0) \\ \mathbf{A}_m(\mathbf{r}) = \frac{1}{c} \mathbf{E} \times \mathbf{e}_z \\ \mathbf{A}_m(\mathbf{r}) = (\mathcal{E}/c)(0, x, 0) \end{array} \right.$$

Landau energy level: $E_n = \hbar\omega_c \left(n + \frac{1}{2} \right) \quad \omega_c = \frac{g\mu_B \mathcal{E}}{mc^2}$

Cyclotron motion: $l_\epsilon \equiv \sqrt{\hbar c^2 / g\mu_B \mathcal{E}}$

Cf. Skyrmion lattice induced by DM int.

Hoogdalem, Tserkovnyak & Loss, PRB (2013)



DM int. \rightarrow Vector potential analogous to \mathbf{A}_m

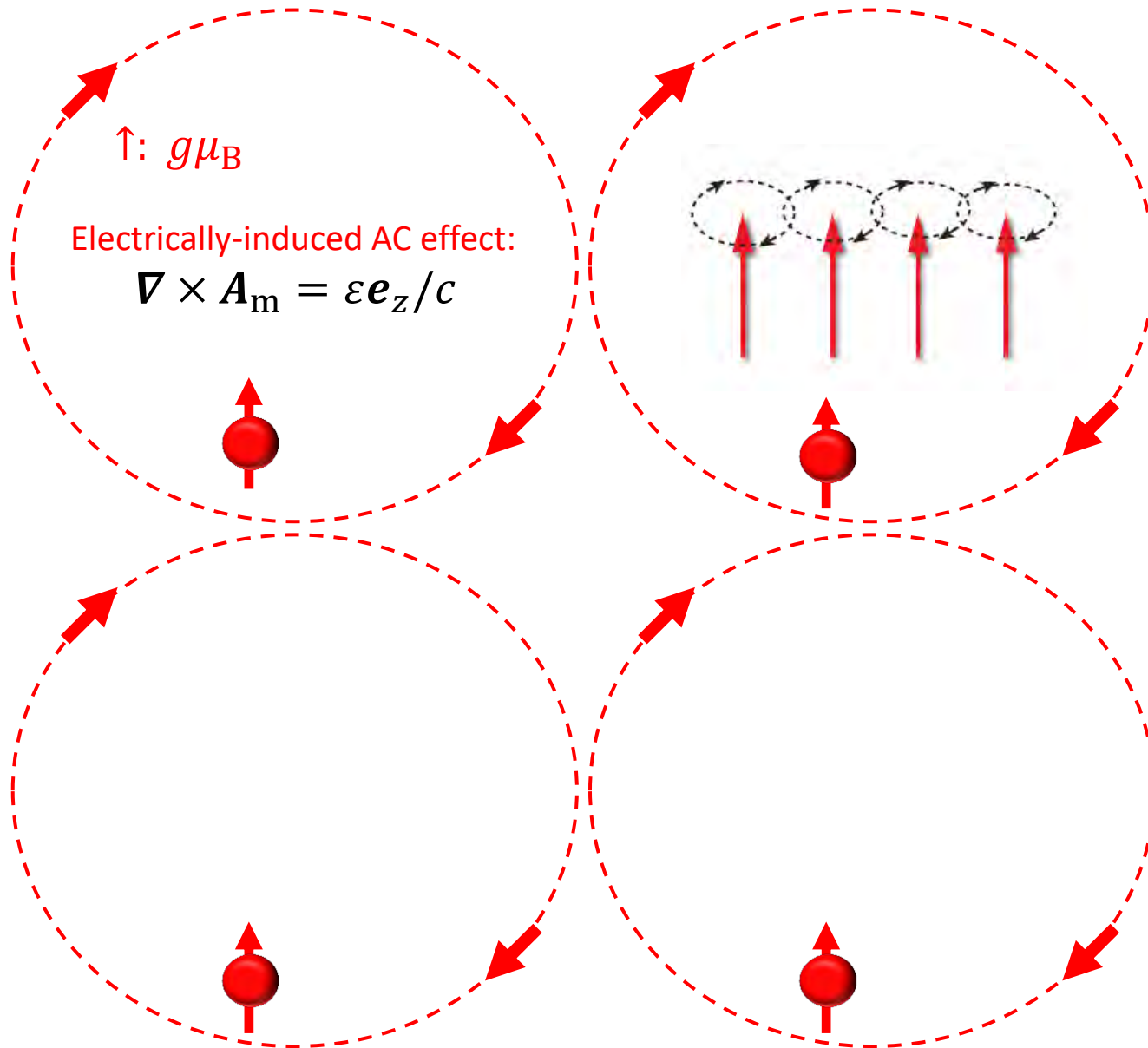
Average fictitious field (textured magnetization)

Landau gap: $\Delta E_n = 2.5 \text{ meV} = 18 \text{ K}$

Within experimental reach: Nagaosa & Tokura, Nat. Nano. (2013)

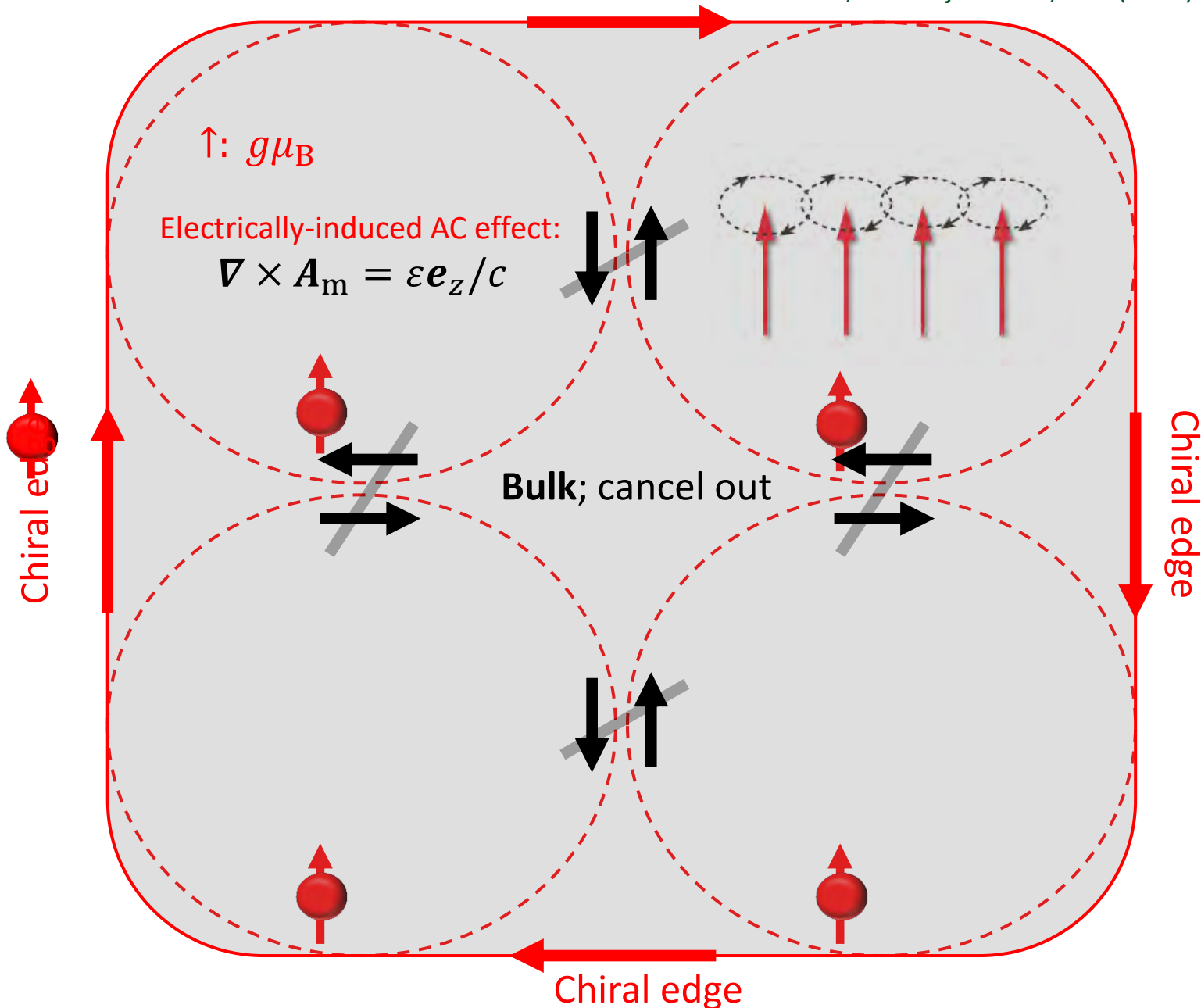
“Quantum” Hall effect of up-magnon in FM: Chiral edge by cyclotron motion

KN, Klinovaja & Loss, PRB (2017)



“Quantum” Hall effect of up-magnon in FM: Chiral edge by cyclotron motion

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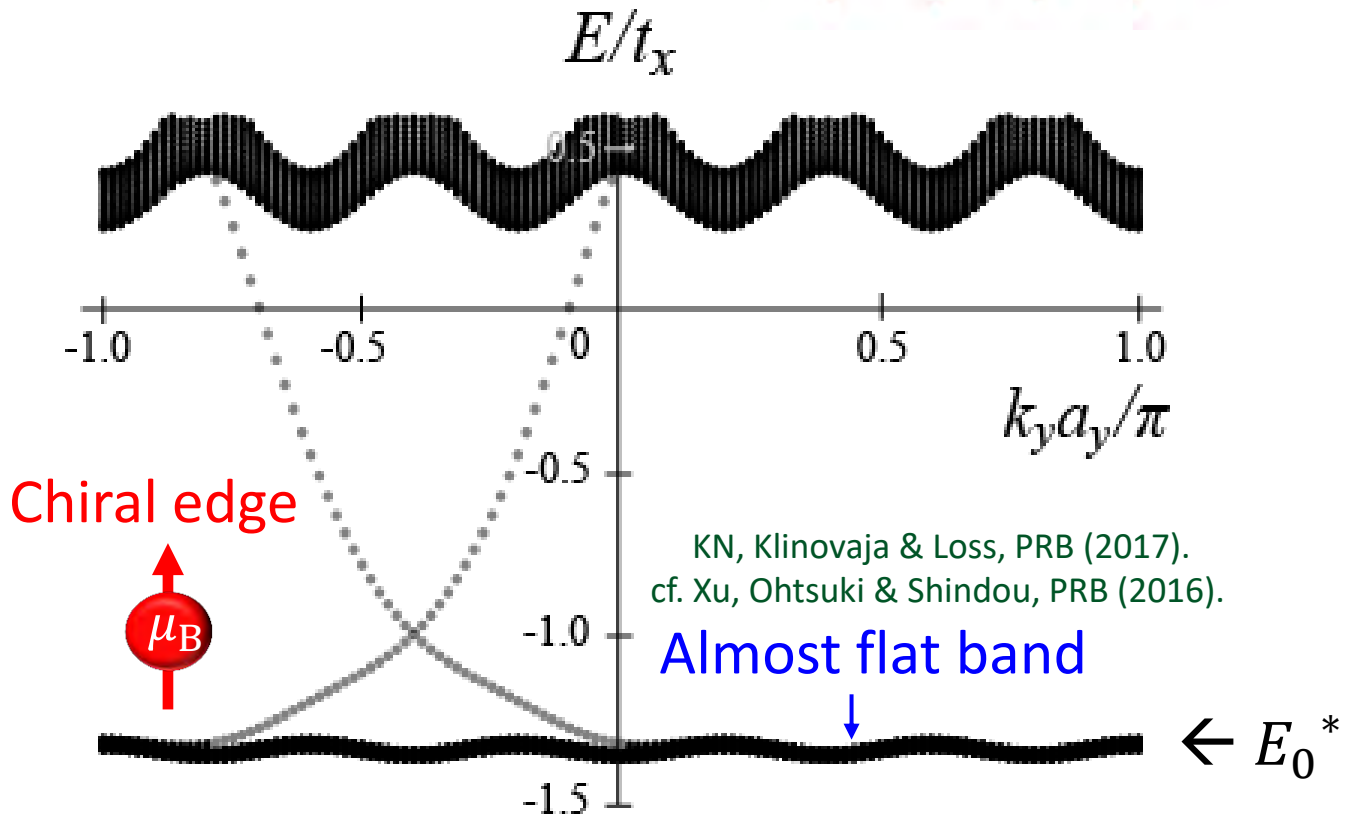
Chiral Edge Magnon States: Topological FM

KN, Klinovaja & Loss, PRB (2017)

$$\mathcal{H}_m = - \sum_{\langle ij \rangle} t_{ij} e^{i\theta_{ij}} (a_i a_j^\dagger + \text{H.c.}) \quad \theta_{ij} = (g\mu_B / \hbar c^2) \int_{\mathbf{x}_i}^{\mathbf{x}_j} d\mathbf{r} \cdot (\mathbf{E} \times \mathbf{e}_z)$$

Electrically-induced AC effect: $\mathbf{A}_m \equiv \mathbf{E} \times \mathbf{e}_z / c$

$$\nabla \times \mathbf{A}_m = \varepsilon \mathbf{e}_z / c$$



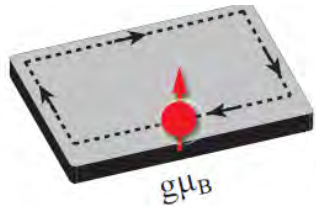
Magnonic WF Law in Topological FM: Bulk

KN, Klinovaja & Loss, PRB (2017)

Chiral edge:

Chern integer:

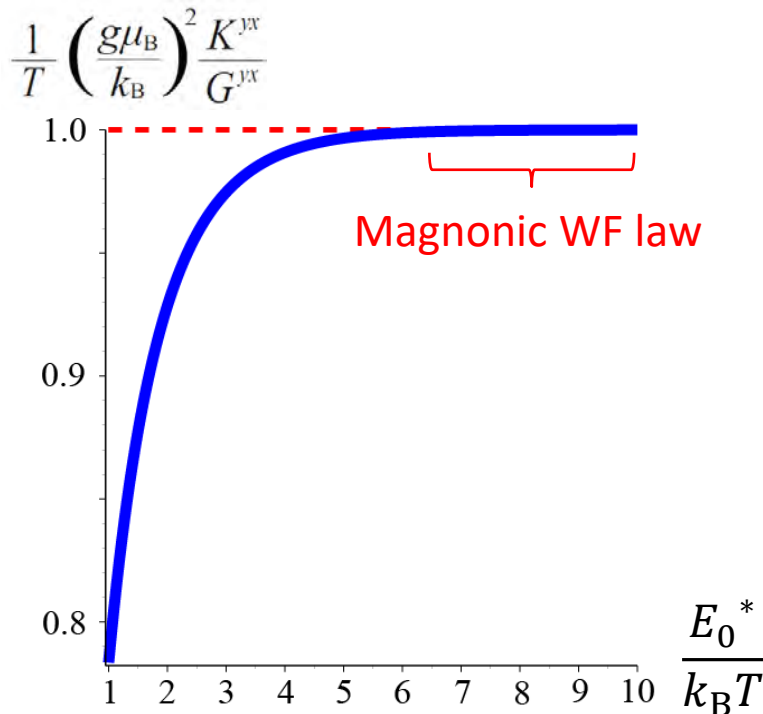
Hall coefficients: $L_{ij}^{yx} \propto N_{0\uparrow}$ in almost flat band



$$\mathcal{N}_{0\uparrow} = +1$$

$$\begin{pmatrix} \langle j_y \rangle \\ \langle j_y^Q \rangle \end{pmatrix} = \begin{pmatrix} L_{11}^{yx} & L_{12}^{yx} \\ L_{21}^{yx} & L_{22}^{yx} \end{pmatrix} \begin{pmatrix} -\partial_y B \\ -\partial_y T/T \end{pmatrix}$$

cf. Matsumoto & Murakami, PRL (2011)



Magnonic WF law in 'quantum' Hall system:

Universal at low temperature ($k_B T \ll E_0^*$)

$$\frac{K^{yx}}{G^{yx}} = \frac{1}{T} \frac{L_{22}^{yx} - \frac{(L_{12}^{yx})^2}{L_{11}^{yx}}}{L_{11}^{yx}} \cong \left(\frac{k_B}{g\mu_B} \right)^2 T \propto T$$

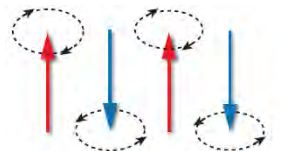
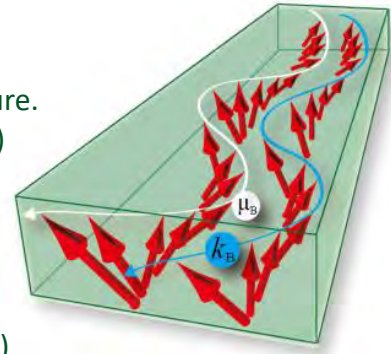
Counter-current by magnetization gradient

NOTE: $K^{yx} \neq L_{22}^{yx}/T$ for boson

Observation of Magnonic Edge State & WF Law

Within experimental reach

2006	Quasi-equilibrium magnon-BEC	Demokritov <i>et al.</i> , Nature
	Inverse spin-Hall effect	Saitoh <i>et al.</i> , APL
2008	Spin-Seebeck effect	Uchida <i>et al.</i> ('08, '10, '11), Nature. Theory: Adachi <i>et al.</i> , PRB (2011)
2010	Spin-wave spin current	Kajiwara <i>et al.</i> , Nature
	Magnon thermal Hall effect: Magnonic thermal conductivity	Onose <i>et al.</i> , Science Theory: Katsura <i>et al.</i> , PRL (2010) Matsumoto & Murakami, PRL (2011) Shindou <i>et al.</i> , PRB ('13): Chiral edge magnon mode
2014	Aharonov-Casher effect on magnon	Yale-group, PRL: Observation. Aharonov & Casher, PRL (1984)
2016	Snell's law for spin-waves	Tanabe <i>et al.</i> APE (2014). Stigloher <i>et al.</i> , PRL (2016).
	Magnonic spin conductivity	Cornelissen <i>et al.</i> , PRB (2016).
	Magnon planar Hall effect	Liu <i>et al.</i> , PRB (2017).
	Spin Seebeck effect in AF	Seki <i>et al.</i> , (2015). Theory: Ohnuma <i>et al.</i> , PRB (2013)
	Magnonic spin Nernst effect in AF	Shiomi <i>et al.</i> , arXiv:1706.03978. cf. Theory [Cheng <i>et al.</i> , PRL (2016)] & [Zyuzin <i>et al.</i> , PRL (2016)]



Remark:

Thermal (Hall) Conductance for Boson

KN, Simon & Loss, PRB (2015)

KN, Klinovaja & Loss, PRB (2017)

KN, Simon & Loss, J. Phys. D (2017): Review article

Thermal Conductivity $K \approx L^{22}$ for Fermions

Textbook by Ashcroft & Mermin

$$\begin{array}{l} \text{Charge} \\ \text{Heat} \end{array} \begin{pmatrix} \mathbf{J}_e \\ \mathbf{J}_Q \end{pmatrix} = \begin{pmatrix} L^{11} & L^{12} \\ L^{21} & \underbrace{L^{22}}_K \end{pmatrix} \begin{pmatrix} \mathbf{E} \\ -\nabla T \end{pmatrix}$$

WF law: Electron = Fermion

*Lifshitz & Pitaevskii (Vol. 10)

$$\frac{K}{\sigma} \approx \frac{L^{22} + \mathcal{O}((k_B T / \epsilon_F)^2)^*}{L^{11}} \Rightarrow \frac{\pi^2}{3} \left(\frac{k_B}{e} \right)^2 T$$

Thermal Conductivity $K \neq L^{22}$ for Magnons

KN, Simon & Loss, PRB (2015)

$$\begin{array}{l} \text{Magnon} \\ \text{Heat} \end{array} \begin{pmatrix} \mathbf{I}_m \\ \mathbf{I}_Q \end{pmatrix} = \begin{pmatrix} L^{11} & L^{12} \\ L^{21} & \underbrace{L^{22}}_K \end{pmatrix} \begin{pmatrix} \nabla B \\ -\nabla T \end{pmatrix}$$

WF law: Magnon = Boson

$$\frac{K}{G} \equiv \frac{L^{22} - L^{21}L^{12}/L^{11}}{L^{11}} \Rightarrow \left(\frac{k_B}{g\mu_B} \right)^2 T$$

Textbook by Ashcroft & Mermin Eq. (13.56): K is measured under conditions of no quasi-particle current

$$\text{Thermal conductivity } K: \mathbf{I}_Q \equiv -K \cdot \nabla T \quad \text{with} \quad \mathbf{I}_m \stackrel{!}{=} 0$$

$$\mathbf{I}_m = L^{11} \nabla B - L^{12} \nabla T \stackrel{!}{=} 0 \quad \rightarrow \quad \nabla B^* = \frac{L^{12}}{L^{11}} \nabla T$$

$$\mathbf{I}_Q = L^{21} \nabla B^* - L^{22} \nabla T = - \underbrace{\left(L^{22} - L^{21}L^{12}/L^{11} \right)}_K \nabla T$$

Magnetization gradient:

Johnson & Silsbee (1987)

Basso *et al.* (2016)

Topological AF

KN, Kim, Klinovaja, Loss (2017)

STRATEGY

Magnonic classical Hall effect in Aharonov-Casher phase: No topological edge

Meier & Loss, PRL (2003)



Topological invariant: Chern (TKNN) integer

TKNN, PRL (1982)

Kohmoto, Ann. Phys. (1985)

Magnonic “quantum” Hall effect in topological FM: Chiral edge

KN, Klinovaja & Loss, PRB (2017)



\mathbb{Z}_2 topological number

Hasan & Kane, RMP (2010) *etc.*

Magnonic “quantum” spin Hall effect in topological AF: Helical edge

KN, Kim, Klinovaja & Loss (2017)

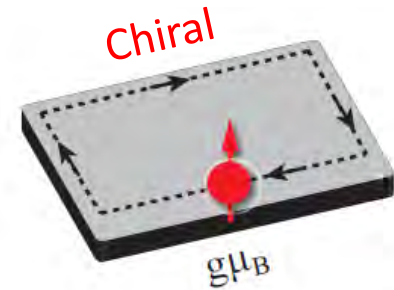
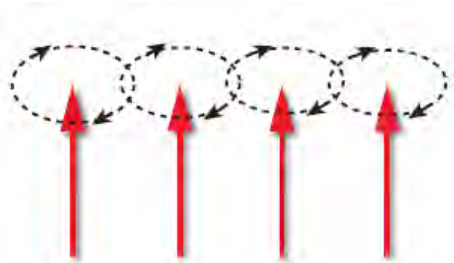
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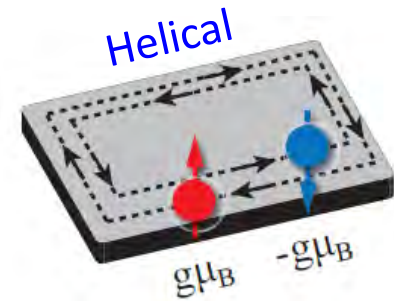
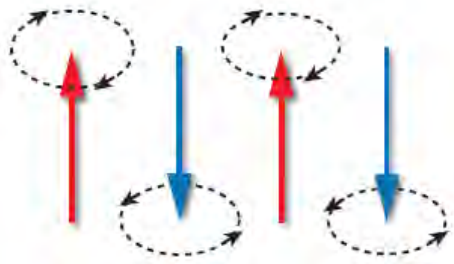
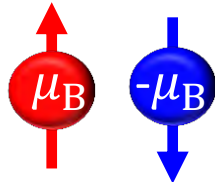
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FM:



AF:



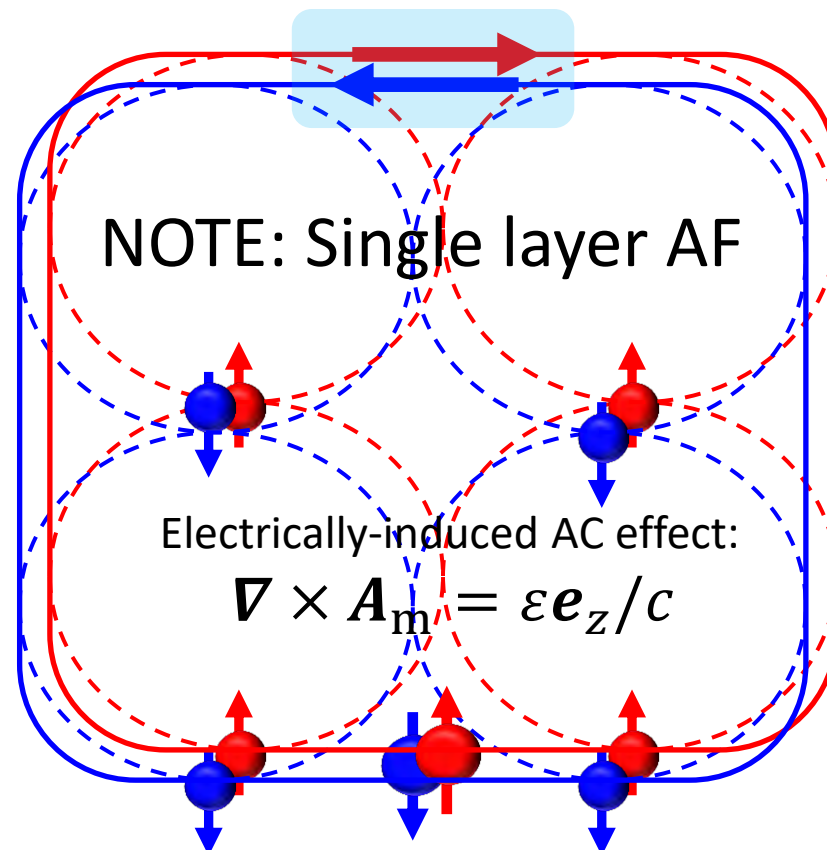
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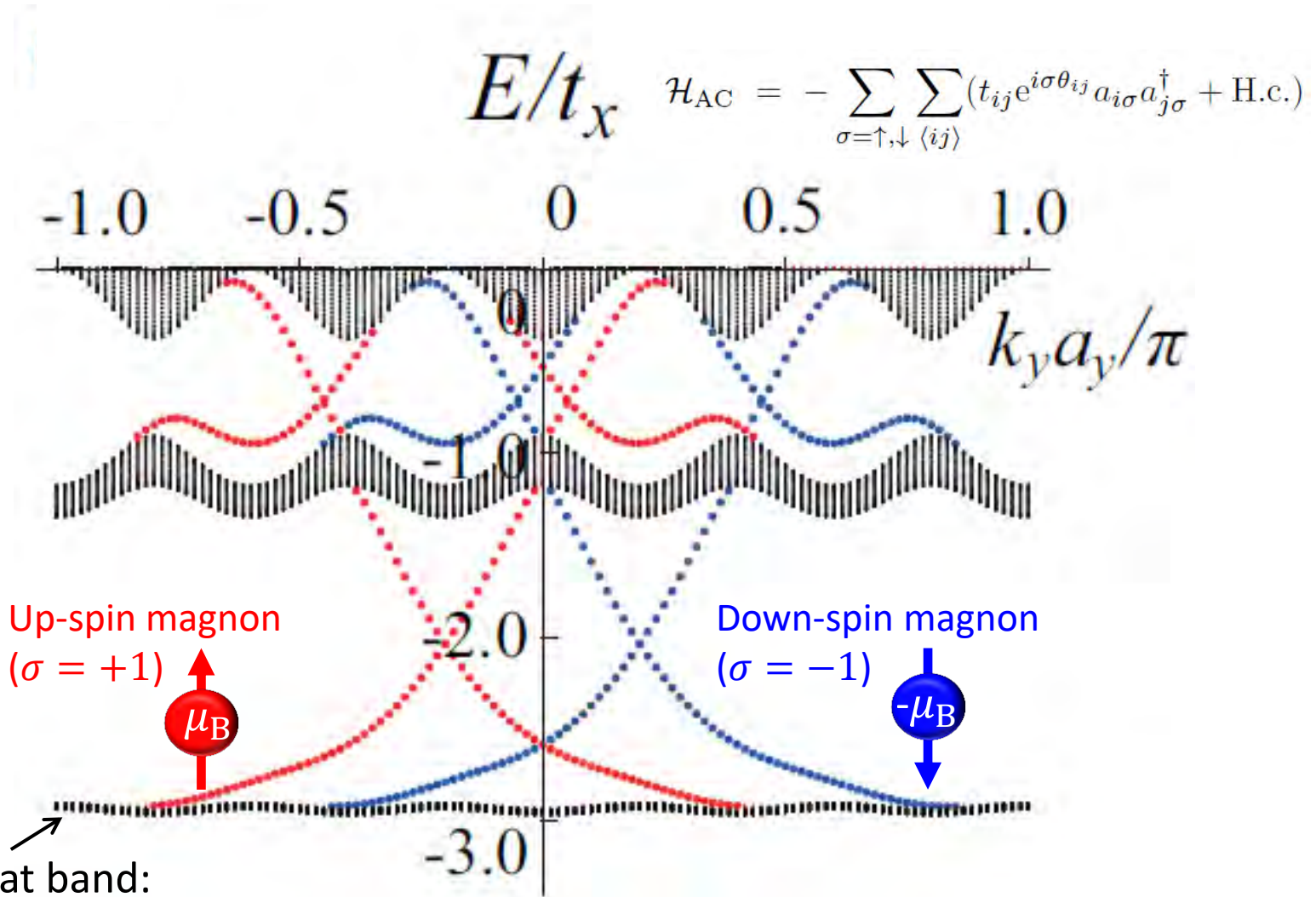
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(Single layer) AF: Helical edge magnon




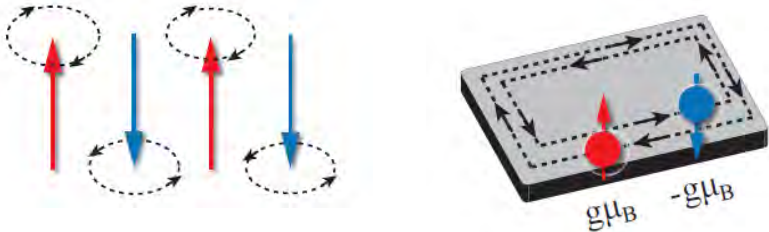
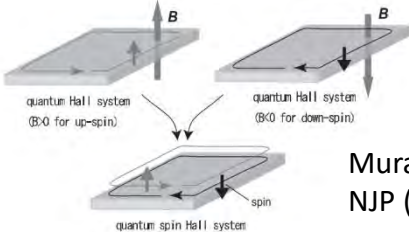
Helical Edge Magnon States: Topological AF

KN, Kim, Klinovaja & Loss (2017) arXiv:1707.07427



cf. KN, Klinovaja & Loss, PRB (2017)

Magnonic TI: AF in AC effect

$\nabla \times \mathbf{A}_m$ $= \epsilon \mathbf{e}_z / c$	Up- & down-magnons $\sigma g \mu_B$ along the opposite direction $\rightarrow N_{0\sigma} = \pm 1 = \sigma$	
FM	<p>TKNN integer: Chiral edge</p> <p>KN, Klinovaja & Loss, PRB (2017)</p> $\mathcal{N}_{0\uparrow} = +1$	 <p>Chiral edge:</p>
AF	<p>Magnonic TI: Helical edge</p> <p>KN, Kim, Klinovaja & Loss (2017) $\sigma = \pm 1 = \uparrow, \downarrow$</p> $\mathcal{N}_{0\sigma} = \sigma$	 <p>Helical edge:</p>
<p>Total Chern number: $\sum_{\sigma} \mathcal{N}_{0\sigma} = \mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow} = 0$</p> <p>$\mathbb{Z}_2$ topological invariant: $\mathcal{Z}_0 \equiv \frac{(\mathcal{N}_{0\uparrow} - \mathcal{N}_{0\downarrow})}{2} = 1$</p> <p>Hasan & Kane, RMP (2010) etc.</p>		<p>cf. An electronic TI by AB phase</p>  <p>Murakami, NJP (2007)</p>

Magnonic TI: AF in AC effect

cf. KN, Klinovaja & Loss, PRB (2017)

Spin

Nernst*: Ensured by \mathbb{Z}_2 invariant

$$\begin{matrix} \text{Spin:} \\ \text{Heat:} \end{matrix} \begin{pmatrix} \langle \mathcal{J}_y \rangle \\ \langle \mathcal{J}_y^Q \rangle \end{pmatrix} = \begin{pmatrix} L'_{11}(\mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow}) & L'_{12}(\mathcal{N}_{0\uparrow} - \mathcal{N}_{0\downarrow}) \\ L'_{21}(\mathcal{N}_{0\uparrow} - \mathcal{N}_{0\downarrow}) & L'_{22}(\mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow}) \end{pmatrix} \begin{pmatrix} \partial_x B \\ -\frac{\partial_x T}{T} \end{pmatrix}$$

Ettinghausen

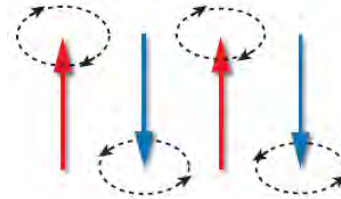
Heat

AF

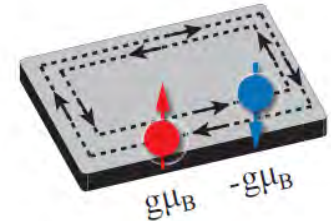
Magnonic TI: Helical edge

KN, Kim, Klinovaja & Loss (2017) $\sigma = \pm 1 = \uparrow, \downarrow$

$$\mathcal{N}_{0\sigma} = \sigma$$



Helical edge:

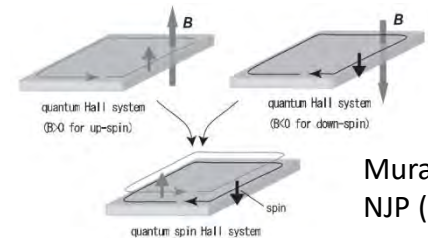


Total Chern number: $\sum_{\sigma} \mathcal{N}_{0\sigma} = \mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow} = 0$

\mathbb{Z}_2 topological invariant: $\mathcal{Z}_0 \equiv \frac{(\mathcal{N}_{0\uparrow} - \mathcal{N}_{0\downarrow})}{2} = 1$

Hasan & Kane, RMP (2010) etc.

cf. An electronic TI by AB phase



Murakami, NJP (2007)

NOTE. AF magnon Nernst*: Report of observation in a AF [Shiomi, Takashima & Saitoh, arXiv:1706.03978].
cf. Theory [Cheng *et al.*, PRL (2016)] & [Zyuzin *et al.*, PRL (2016)]

LAST QUESTION

Is magnonic Wiedemann-Franz law satisfied in insulating AFs ?

FM: [KN, Simon & Loss, PRB (2015)] & [KN, Klinovaja & Loss, PRB (2017)]

“No” in topological AF (magnonic TI):

$$G_{\text{AF}}^{yx} = 0, \quad K_{\text{AF}}^{yx} = 0 \quad : \text{Topological invariant} \quad \sum_{\sigma} \mathcal{N}_{0\sigma} = \mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow} = 0$$
$$\therefore G_{\text{AF}}^{yx} \propto \mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow} = 0 \quad K_{\text{AF}}^{yx} \propto \mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow} = 0$$

“Yes” in non-topological AF:

$$\text{WF law for bulk magnons} \quad \text{AF: } \frac{K_{\text{AF}}}{G_{\text{AF}}} \cong \frac{5}{2} \left(\frac{k_{\text{B}}}{g\mu_{\text{B}}} \right)^2 T \quad \text{FM: } \frac{K_{\sigma}}{G_{\sigma}} \cong \frac{5}{2} \left(\frac{k_{\text{B}}}{g\mu_{\text{B}}} \right)^2 T$$

∴ AF = Independent copies of FM

GOAL

Topological insulator

Kane & Mele, PRL (2005, 2005).

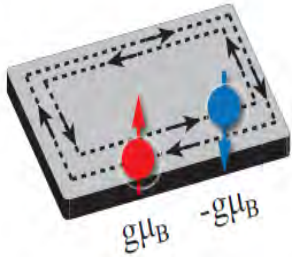
Bernevig & Zhang, PRL (2006).

Magnonic Topological insulator: A bosonic analog

KN, Kim, Klinovaja & Loss (2017).

arXiv:1707.07427

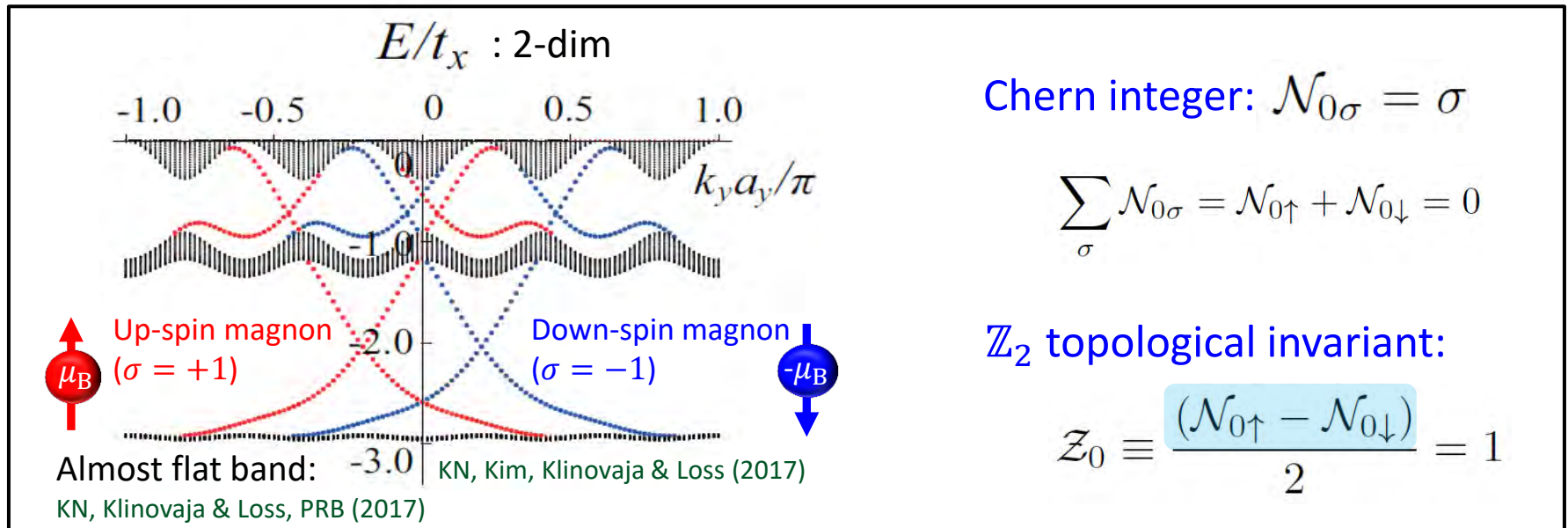
SUMMARY



Magnonic Topological Insulator in Antiferromagnet:

Helical edge magnon state in Aharonov-Casher effect

KN (Basel), S. K. Kim (UCLA), J. Klinovaja, D. Loss. arXiv:1707.07427



Spin

Nernst*: Ensured by \mathbb{Z}_2 invariant

Magnonic WF law:

$$\text{Spin: } \begin{pmatrix} \langle \mathcal{J}_y \rangle \\ \langle \mathcal{J}_y^Q \rangle \end{pmatrix} = \begin{pmatrix} L'_{11}(\mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow}) & L'_{12}(\mathcal{N}_{0\uparrow} - \mathcal{N}_{0\downarrow}) \\ L'_{21}(\mathcal{N}_{0\uparrow} - \mathcal{N}_{0\downarrow}) & L'_{22}(\mathcal{N}_{0\uparrow} + \mathcal{N}_{0\downarrow}) \end{pmatrix} \begin{pmatrix} \partial_x B \\ -\frac{\partial_x T}{T} \end{pmatrix}$$

KN, Simon & Loss, PRB (2015)

No, in topological AF.

Yes, in trivial AF.

Ettinghausen

Heat

*Report of observation in a AF [Shiomi, Takashima & Saitoh, arXiv:1706.03978]

APPENDIX

NOTE: *pptx-file for animation is available in the link:*

https://www.dropbox.com/s/9be58ryqrdhsyu5/AFmagnonicZ2TI_Animation_KoukiNakata.pptx?dl=0

All the responsibilities of this slide rest with Kouki Nakata (July 2017)

Onsager Coefficients: Bulk Magnon Transport

KN, Kim, Klinovaja, Loss (2017)

Topologically trivial AF

Hamiltonian: $\mathcal{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{\mathbf{k}} \hbar\omega_{\mathbf{k}\sigma} a_{\mathbf{k}\sigma}^\dagger a_{\mathbf{k}\sigma}$

$$\begin{pmatrix} \langle j_{x\sigma} \rangle \\ \langle j_{x\sigma}^Q \rangle \end{pmatrix} = G_\sigma \begin{pmatrix} L_{11\sigma} & L_{12\sigma} \\ L_{21\sigma} & L_{22\sigma} \end{pmatrix} \begin{pmatrix} \partial_x B \\ -\partial_x T/T \end{pmatrix}$$

Particle current: $j_{x\sigma}^P = j_{x\sigma}/(\sigma g\mu_B)$

Thermal conductivity: $K_\sigma = \frac{1}{T} \left(L_{22\sigma} - \frac{L_{12\sigma}L_{21\sigma}}{L_{11\sigma}} \right)$

WF law for bulk magnons:

AF: $\frac{K_{AF}}{G_{AF}} \Rightarrow \frac{5}{2} \left(\frac{k_B}{g\mu_B} \right)^2 T$ FM: $\frac{K_\sigma}{G_\sigma} \Rightarrow \frac{5}{2} \left(\frac{k_B}{g\mu_B} \right)^2 T$

Topological AF

Hamiltonian: $\mathcal{H}_{AC} = - \sum_{\sigma=\uparrow,\downarrow} \sum_{\langle ij \rangle} (t_{ij} e^{i\sigma\theta_{ij}} a_{i\sigma} a_{j\sigma}^\dagger + \text{H.c.})$

i.e., $\mathcal{H}_{m\sigma} = \frac{1}{2m} \left(\mathbf{p} + \sigma \frac{g\mu_B}{c} \mathbf{A}_m \right)^2$

$$\begin{pmatrix} \langle j_{y\sigma} \rangle \\ \langle j_{y\sigma}^Q \rangle \end{pmatrix} = G_\sigma^{yx} \begin{pmatrix} L_{11\sigma}^{yx} & L_{12\sigma}^{yx} \\ L_{21\sigma}^{yx} & L_{22\sigma}^{yx} \end{pmatrix} \begin{pmatrix} \partial_x B \\ -\partial_x T/T \end{pmatrix}$$

Particle Hall current: $j_{y\sigma}^P = j_{y\sigma}/(\sigma g\mu_B)$

$$\langle j_{y\sigma}^P \rangle = \sigma \mathcal{N}_{0\sigma} \frac{L'_{11}}{g\mu_B} \partial_x B - \mathcal{N}_{0\sigma} \frac{L'_{12}}{g\mu_B} \frac{\partial_x T}{T}$$

Spin & thermal Hall conductances:

$$G_{AF}^{yx} = 0, \quad K_{AF}^{yx} = 0 \quad K_\sigma^{yx} = \left(L_{22\sigma}^{yx} - \frac{L_{21\sigma}^{yx} L_{12\sigma}^{yx}}{L_{11\sigma}^{yx}} \right) / T$$

Antiferromagnet	Topologically trivial bulk: Sec. II	Topological bulk: Sec. IIIB
Magnetic field gradient:	Helical magnon transport. Spin: $\sum_{\sigma} L_{11\sigma} \neq 0$. Heat: $\sum_{\sigma} L_{21\sigma} = 0$.	Magnon Hall transport. Spin: $\sum_{\sigma} L_{11\sigma}^{yx} = 0$. Heat: $\sum_{\sigma} L_{21\sigma}^{yx} \neq 0$.
Thermal gradient:	Magnon transport. Spin: $\sum_{\sigma} L_{12\sigma} = 0$. Heat: $\sum_{\sigma} L_{22\sigma} \neq 0$.	Helical magnon Hall transport. Spin: $\sum_{\sigma} L_{12\sigma}^{yx} \neq 0$. Heat: $\sum_{\sigma} L_{22\sigma}^{yx} = 0$.

Cyclotron Motion Along the Opposite Direction

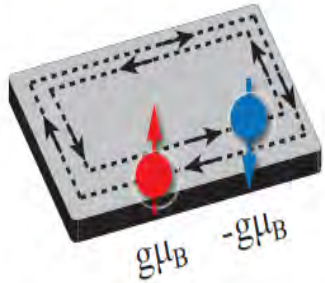


Fig. 1 (b)

Indeed, they form the same Landau levels⁴³ with the principal quantum number $n_\sigma \in \mathbb{N}_0$,

$$E_{n_\sigma} = \hbar\omega_c \left(n_\sigma + \frac{1}{2} \right) + \Delta \quad \text{for } n_\sigma \in \mathbb{N}_0, \quad (15)$$

and the two magnons $\sigma g\mu_B \mathbf{e}_z$ perform cyclotron motions with the same frequency⁴³

$$\omega_c = \frac{g\mu_B \mathcal{E}}{mc^2} \quad (16)$$

and same electric length⁴³ $l_\mathcal{E}$, defined by

$$l_\mathcal{E} \equiv \sqrt{\hbar c^2 / g\mu_B \mathcal{E}}, \quad (17)$$

but along opposite direction, cf. Fig. 1 (b),

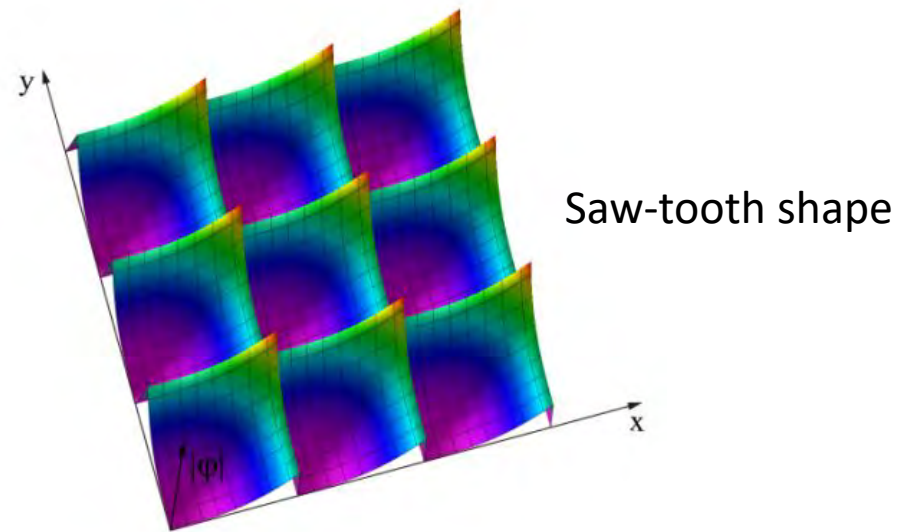
$$\frac{d}{dt}(\mathcal{R}_{x\sigma} + i\mathcal{R}_{y\sigma}) = i\sigma\omega_c(\mathcal{R}_{x\sigma} + i\mathcal{R}_{y\sigma}), \quad (18)$$

where $\mathbf{R}_{\mathcal{E}\sigma} = (\mathcal{R}_{x\sigma}, \mathcal{R}_{y\sigma})$ is the relative coordinate⁸². The factor σ in Eq. (18) is rooted in the magnetic dipole moment $\sigma g\mu_B \mathbf{e}_z$ of a magnon. The source of cyclotron motion is the electric field gradient \mathcal{E} [Eqs. (14) and (16)], which is common to the both modes.

where $\mathbf{R}_{\mathcal{E}\sigma} = (\mathcal{R}_{x\sigma}, \mathcal{R}_{y\sigma}) \equiv (-l_\mathcal{E}^2 \Pi_{y\sigma} / \hbar, l_\mathcal{E}^2 \Pi_{x\sigma} / \hbar)$ and $\mathbf{\Pi}_\sigma \equiv \mathbf{p} + \sigma g\mu_B \mathbf{A}_m / c$

Q. Helical edge magnon state:

Still exist in 'periodic' electric vector potential A_m ?

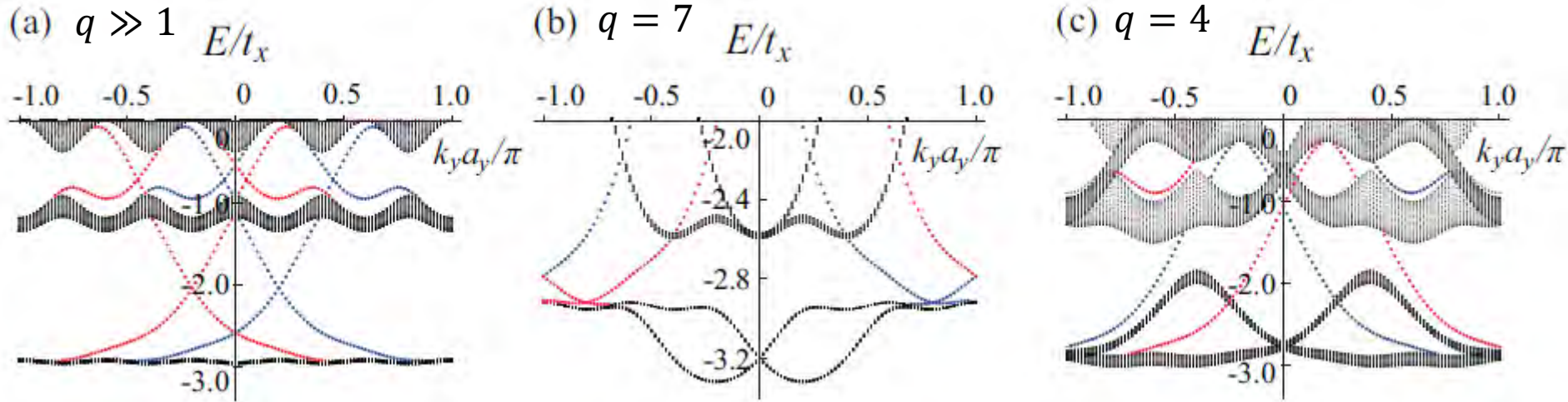


ANSWER: YES.

KN, S. K. Kim, J. Klinovaja, D. Loss (2017)

Helical Edge Magnon States in Periodic potential A_m

KN, Kim, Klinovaja, Loss (2017)



AC phase: $\theta_n = n\theta_0$ $\theta_0 \equiv (g\mu_B/\hbar c^2)\mathcal{E}a_x a_y$

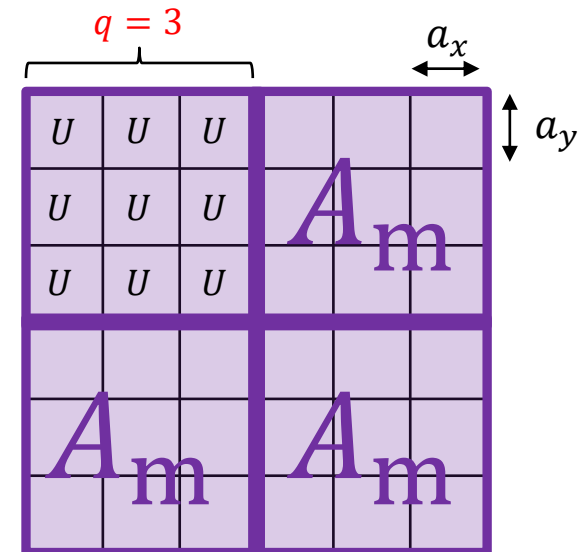
Periodicity: $R_q = qa_x$ $\mathbf{A}'_m \rightarrow \mathbf{A}'_{mq} = (\mathcal{E}R_q/c)(0, \{x/R_q\}, 0)$

Landau gauge: $\mathbf{A}'_m = (\mathcal{E}/c)(0, x, 0) \rightarrow$ Spectrum: $E = E(k_y)$

$$\mathcal{H}_{AC} = - \sum_{\sigma=\uparrow,\downarrow} \sum_{\langle ij \rangle} (t_{ij} e^{i\sigma\theta_{ij}} a_{i\sigma} a_{j\sigma}^\dagger + \text{H.c.})$$

$$\mathcal{H}_{AC} = \sum_{k_y} H_{k_y}$$

$$H_{k_y} = -t_x \sum_{n,\sigma} (a_{k_y, n+1, \sigma}^\dagger a_{k_y, n, \sigma} + \text{H.c.}) - 2t_y \sum_{n,\sigma} [\cos(k_y a_y + \sigma\theta_n)] a_{k_y, n, \sigma}^\dagger a_{k_y, n, \sigma}$$



An Intuitive Criterion for Topological Edges

Key: Magnons experience geometric (AC) phase **globally & macroscopically** (*)

→ Otherwise no topological edge modes (#)

35

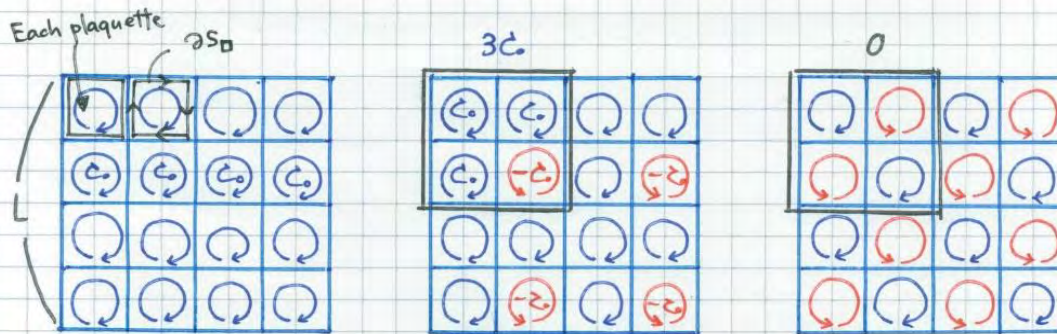
► Conclusion: theoretical proposal for a criterion to generate topological edge

$$\nabla \times \mathbf{A}_{lm} = \frac{\epsilon}{c} \mathbf{e}_z$$

↳ globally or macroscopically constant so that (in the source)

$$\lim_{L \rightarrow \infty} \frac{\sum_{\square} \int dS_{\square} \cdot \nabla \times \mathbf{A}_{lm}}{L} = \lim_{L \rightarrow \infty} \frac{\sum_{\square} \oint_{\partial S_{\square}} d\mathbf{l} \cdot \mathbf{A}_{lm}}{L} \neq 0$$

... (*)



→ Topological edge → Topological edge → No topological edge

* Frequency of cyclotron motion $\omega_c \propto \epsilon \propto |\nabla \times \mathbf{A}_{lm}|$
 i.e., $\nabla \times \mathbf{A}_{lm} = \frac{\epsilon}{c} \mathbf{e}_z$ in Landau or symmetric gauge.

Remember AB effect: $\nabla \times \mathbf{A}_0 = \mathbf{B} \mathbf{e}_z$ ← Source of cyclotron motion.

An Intuitive Criterion for Topological Edges

Key: Magnons experience geometric (AC) phase **globally & macroscopically** (*)

→ Otherwise no topological edge modes (#)

* To produce topological edge modes, the form of A_m is important

①, ②
↓

An origin of topological edge mode
= cyclotron motion

$$A_m \neq 0$$

No edge modes :
Not satisfying (*)

→ No cyclotron motion.

Topological edge

$$\nabla \times A_m = \underline{c}_0 \otimes \underline{e}_z = \text{satisfying } (*)$$

e.g.,

$$A_m = \left\{ \begin{array}{l} \text{Landau gauge} \\ \text{Symmetric gauge} \end{array} \right\} \nabla \times A_m = \underline{c}_0 \otimes \underline{e}_z$$

→ Cyclotron motion since \underline{c}_0 is the source.

Magnonic ``Quantum'' Hall Effect & WF Law

KN, Klinovaja & Loss, Phys. Rev. B 95, 125429 (2017)

See also [KN, Simon & Loss, Phys. Rev. B 92, 134425 (2015)]
for magnonic WF law in insulating FM

Chiral Edge Magnon State: Isotropy

Tight-binding model: $\mathcal{H}_m = \sum_{k_y} H_{k_y} \quad t_{ij} = J_{ij}S$

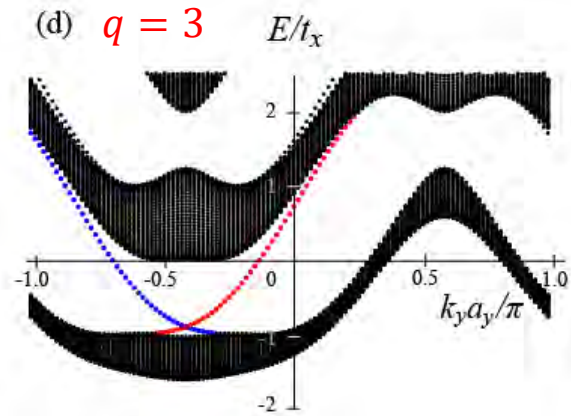
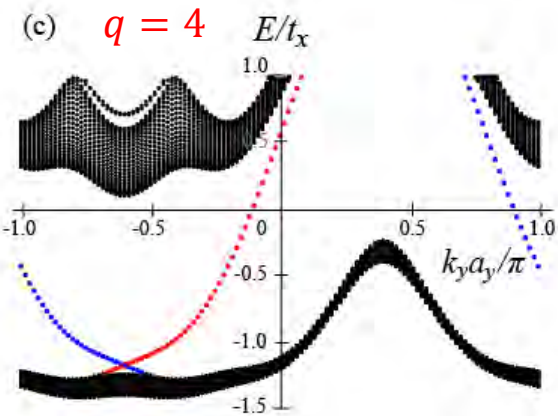
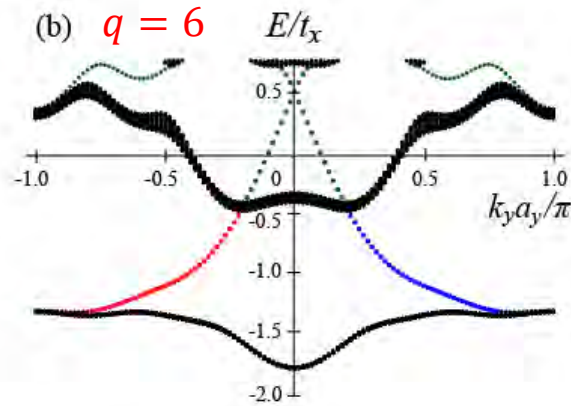
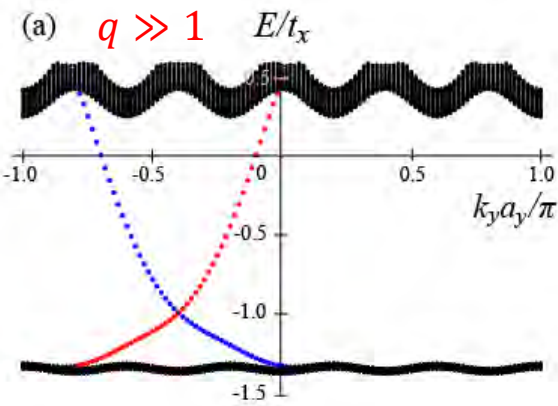
$$H_{k_y} = -t_x \sum_n (a_{k_y, n+1}^\dagger a_{k_y, n} + \text{H.c.}) - 2t_y \sum_n [\cos(k_y a_y + \theta_n)] a_{k_y, n}^\dagger a_{k_y, n}$$

AC phase: $\theta_n = n\theta_0 \quad \theta_0 \equiv (g\mu_B/\hbar c^2)\mathcal{E}a_x a_y$

Landau gauge: $\mathbf{A}'_m = (\mathcal{E}/c)(0, x, 0) \rightarrow$ Spectrum: $E = E(k_y)$

Periodicity: $R_q = qa_x \quad \mathbf{A}'_m \rightarrow \mathbf{A}'_{mq} = (\mathcal{E}R_q/c)(0, \{x/R_q\}, 0)$

cf., Spin Hamiltonian: $t_{ij} = J_{ij}S \quad \mathcal{H}_m = - \sum_{\langle ij \rangle} t_{ij} e^{i\theta_{ij}} (a_i a_j^\dagger + \text{H.c.})$



Isotropic case: $J_x = J_y$

$$\theta_0 = 2\pi/5 < \pi$$

(a)-(d): Chiral edge states

(a)-(b): NOT flat bulk gap

(c)-(d): Bulk gap "closed"
 \rightarrow Gapless

NOTE) Weak disorder:
 Edge mode will not couple to bulk

\sim Weyl systems

cf., Weyl magnon in AF [Li et al., Nat. Comm.(2016)]

Magnonic Wiedemann-Franz Law

KN, Simon, and Loss: Phys. Rev. B **92**, 134425 (2015)

See also [KN, Simon & Loss J. Phys. D. 50, 114004 (2017)] for review article

Universal Thermomagnetic Properties

Magnonic Wiedemann-Franz Law:

KN, Simon, and Loss, Phys. Rev. B **92**, 134425 (2015)

Ratios of L^{ij} , WF law, Seebeck, and Peltier coefficients are material independent

e VS μ_B	Electron (metal)	Magnon (FI)
	Franz and Wiedemann, Annalen der Physik 165 , 497 (1853)	KN, Simon, and Loss, Phys. Rev. B 92 , 134425 (2015)
Statistics	Fermion	Boson
WF law (Low temp.)	$\frac{K}{\sigma} \approx \frac{L^{22} + \mathcal{O}((k_B T / \epsilon_F)^2)}{L^{11}} \Rightarrow \frac{\pi^2}{3} \left(\frac{k_B}{e}\right)^2 T$ (Free electron at low temp.)	$\frac{K}{G} \equiv \frac{L^{22} - L^{21}L^{12}/L^{11}}{L^{11}} \Rightarrow \left(\frac{k_B}{g\mu_B}\right)^2 T$ Low temp.: $\hbar/(2\tau) \ll k_B T \ll g\mu_B B$
Lorenz number	$\mathcal{L} = \frac{\pi^2}{3} \left(\frac{k_B}{e}\right)^2$	$\mathcal{L}_m = \left(\frac{k_B}{g\mu_B}\right)^2$
Seebeck & Peltier	$S \equiv L^{12}/L^{11} \quad \Pi \equiv L^{21}/L^{11}$	$S \overset{\rightarrow}{=} \frac{B}{T} \quad \Pi \overset{\rightarrow}{=} B$
Onsager-Thomson relation	$L^{21} = T L^{12} \quad \Pi = T S$	$L^{21} = T L^{12} \quad \Pi = T S$

OTHERS

Electron e : Fermion

Magnon μ_B : Boson

Review article: KN, Simon & Loss.

J. Phys. D: 50, 114004 (2017). arXiv:1610.08901

Wiedemann-Franz (WF) law

Franz and Wiedemann, Annalen der Physik (1853)

Magnonic Wiedemann-Franz law

KN, Simon & Loss, Phys. Rev. B 92, 134425 (2015).

arXiv:1507.03807

Slide: https://www.dropbox.com/s/5n40xudfu51ibj3/MagnonicWFlaw_KoukiNakata.pdf?dl=0

Josephson effect

Josephson, Phys. Lett. (1962)

Magnonic Josephson effect

KN, Hoogdalem, Simon & Loss, Phys. Rev. B 90, 144419 (2014).

arXiv:1406.7004

KN, Simon & Loss, Phys. Rev. B 92, 014422 (2015).

arXiv:1502.03865

Slide: https://www.dropbox.com/s/704l39hto8pzz79/MagnonBECtransport_KoukiNakata.pdf?dl=0

Integer quantum Hall effect

Klitzing *et al.*, PRL (1980)

TKNN, PRL (1982) / Kohmoto, Ann. Phys. (1985)

Magnonic 'quantum' Hall effect & the WF law: Chiral edge

KN, Klinovaja & Loss, Phys. Rev. B 95, 125429 (2017).

arXiv:1611.09752

Topological ferromagnet (FM)

Slide: https://www.dropbox.com/s/tmryrxua10v5vgm/MagnonicQHE_KoukiNakata.pdf?dl=0

Topological insulator (TI)

Kane and Mele, PRL (2005, 2005).

Bernevig and Zhang, PRL (2006).

Quantum spin Hall effect (QSHE)

Magnonic topological insulator: Helical edge

KN, Kim, Klinovaja, Loss (2017). arXiv:1707.07427

Topological antiferromagnet (AF):

Magnonic quantum spin Hall effects

Slide: https://www.dropbox.com/s/delxgvyumuiuh99/AFmagnonicZ2TI_KoukiNakata.pdf?dl=0

Topological Phenomena

Electron: Quantum Hall Effects & Topological Insulator

- 1982: Thouless, Kohmoto, Nightingale, and Nijs (TKNN), PRL
- 1982: Halperin, PRB (1982). Hatsugai, PRL (1993)
- 1985: Kohmoto, Ann. Phys.
- 1985: Niu, Thouless, and Wu, PRB [Xiao, Chang, and Niu, RMP (2010)]

- 2005: Kane and Mele, PRL (2005, 2005): [Hasan and Kane, RMP (2010)]
- 2006: Bernevig and Zhang, PRL (2006): [Qi and Zhang, RMP (2011)]



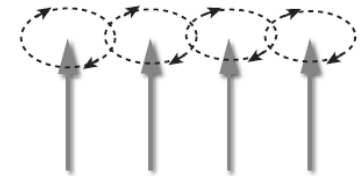
Picture from Google search

FM insulators: Magnonic Hall Effects + ...

- 1995: Haldane & Arovas, PRB
- 2009: Fujimoto, PRL

- 2010: Onose *et al.*, Science
- 2010: Katsura *et al.*, PRL
- 2011: Matsumoto & Murakami, PRL & PRB

- 2013: Shindou *et al.*, PRB *etc.* (2013, 2013, 2014, 2016)
- 2013: Zhang *et al.*, PRB
- 2014: Mook *et al.*, PRB (2014, 2014, 2015)



Phase twist & Berry curvature
in magnonic system

Observation of the magnon Hall effect
& the theories

Magnonic Chern insulators: Chiral edge

NOTE: See [Haldane and Arovas, PRB (1995)] & [Xu, Ohtsuki, and Shindou, PRB (2016)] for disordered quantum Hall systems, and [Matsumoto & Murakami, PRL & PRB (2011)], [Shindou *et al.*, PRB (2013, 2014)], & their review [Murakami & Okamoto, JPSJ (2017)] for chiral edge states in dipolar int. and the bulk-edge correspondence.