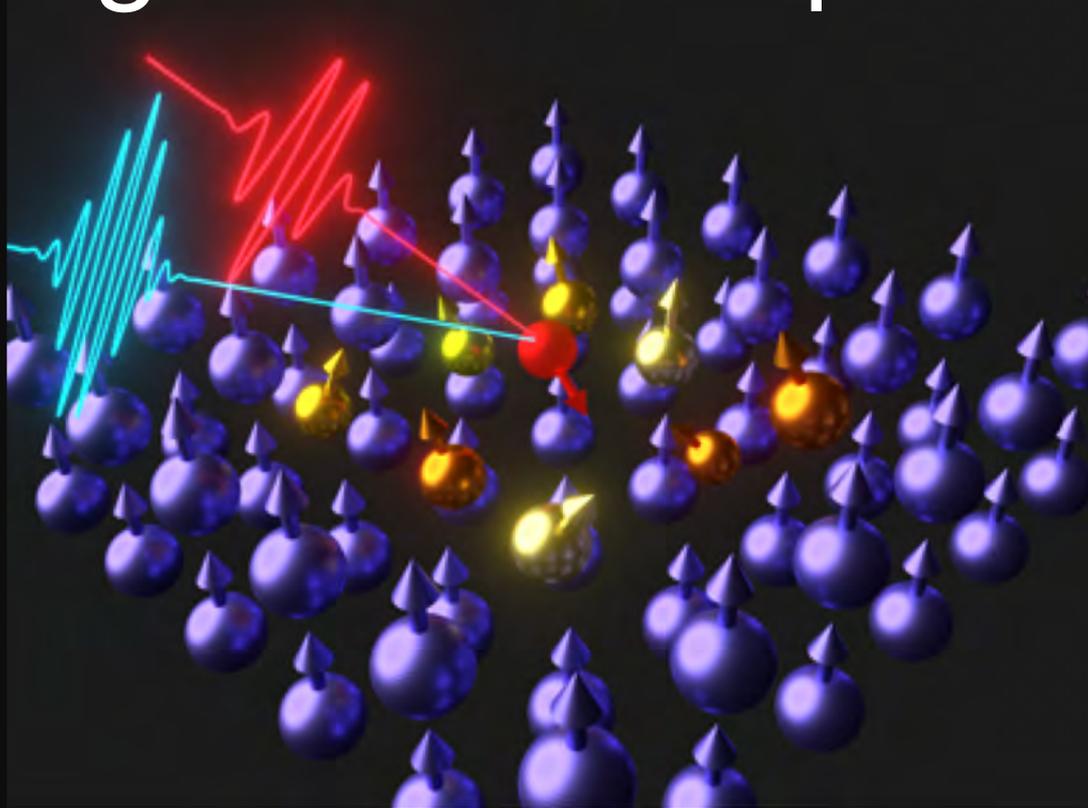


Femtosecond measurement of direct light-induced spin transfer



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Ultrafast Ferromagnetism

- Previous understanding:
 - Ultrafast demagnetization takes place in 100-200 fs

- Introduction to extreme ultraviolet Magneto-optical Kerr Spectroscopy
 - Element Specific Information
 - 10fs bursts, very high time resolution

- New Understanding: Spin excitations within few femtoseconds!
 - Ultrafast magnetic phase transition in Ferromagnetic Nickel
 - Direct optical spin transfer between different elements in Co₂MnGe

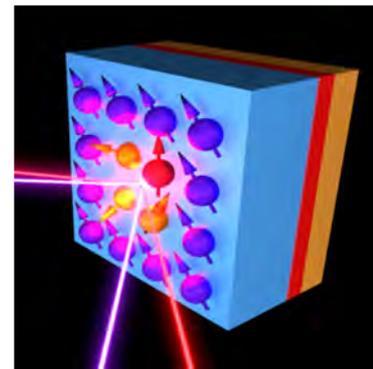
Zusin, D., Tengdin, P., et. al. PRB 97, 024433 (2018)

Tengdin, P. et. al. Sci. Adv. 4 : eaap9744 (2018)

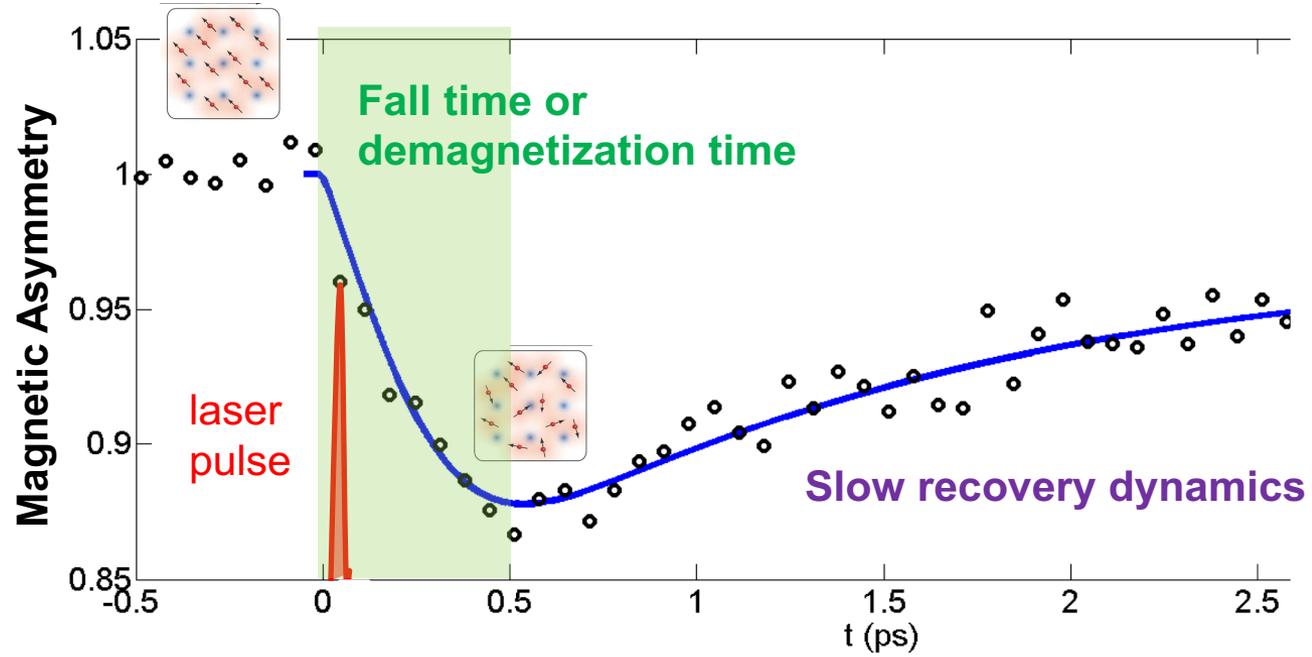
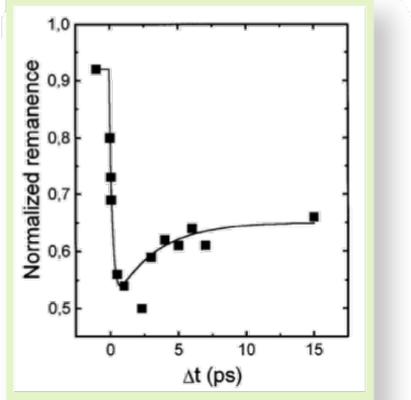
You, W., Tengdin, P., et. al. PRL 121, 077204 (2018)

Tengdin, P. et al., Sci. Adv. in press (2019)

Hofherr, M., et al., Sci. Adv. in press (2019)

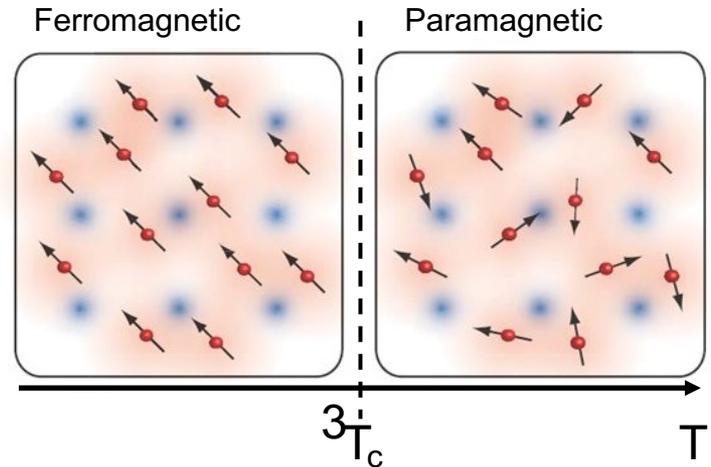
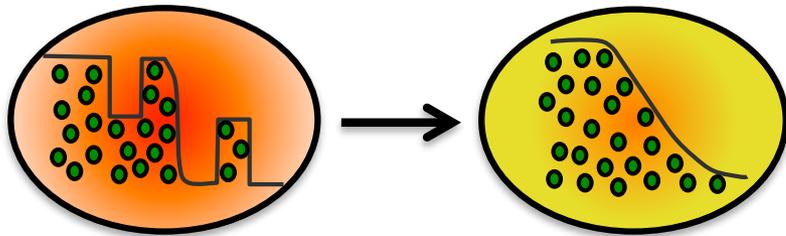


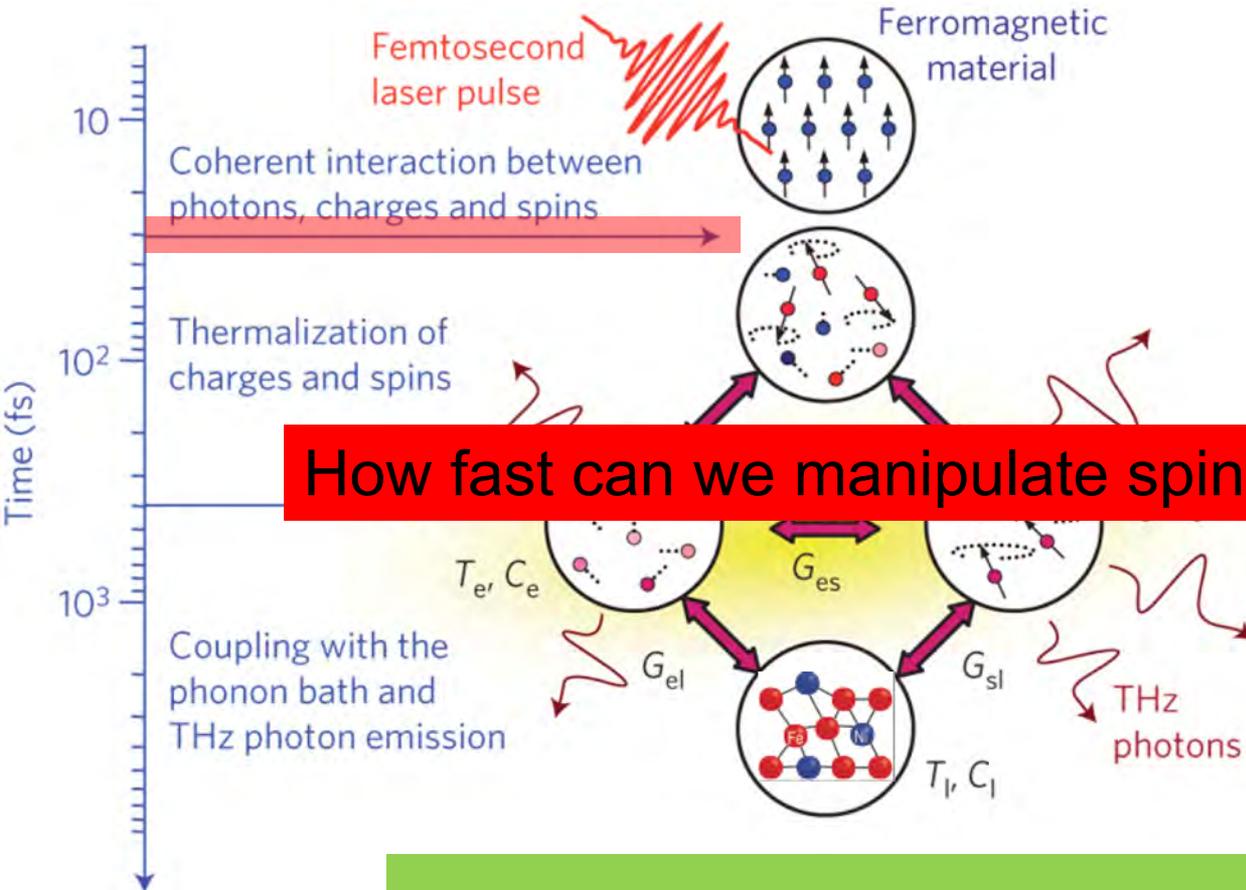
Demagnetization by fs Laser
E Beaurepaire et. al, PRL, 76, 4250 (1996)



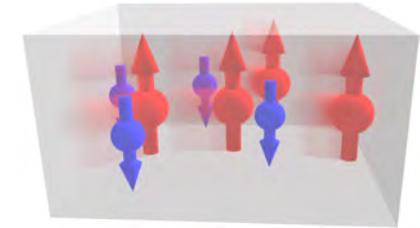
optically excited electrons

thermalized hot electron gas





How fast can we manipulate spins with light?

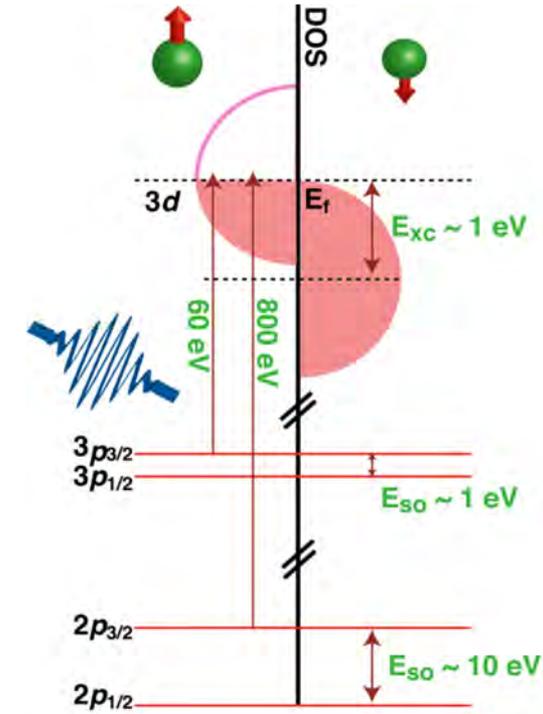
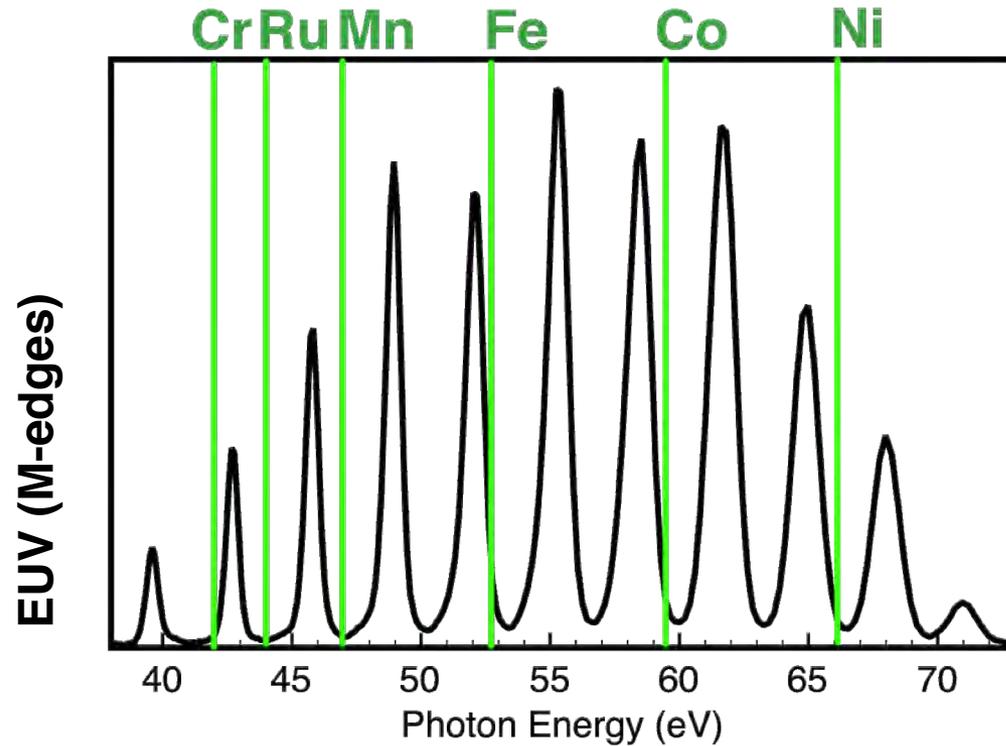


Superdiffusive spin transport : Battiato *et al.*, *PRL* **105**, 27203 (2010)

Bigot et al, *Nature Physics* **5**, 515 (2009)
Koopmans et al., *Nature Mater.* **9**, 259 (2010)

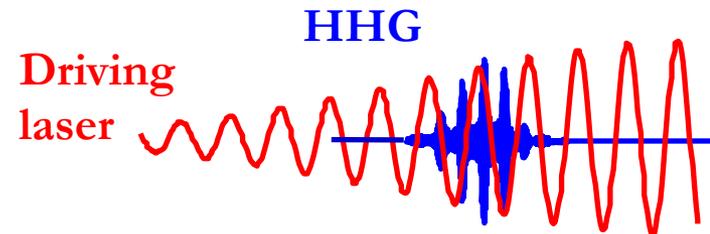
- Spin-flip scattering: *Koopmans et al.*, *Nature Mater.* **9**, 259 (2010)
- Spin-orbit coupling: *Töws et al.*, *PRL* **115**, 217204 (2015)
- Superdiffusive spin transport : *Battiato et al.*, *PRL* **105**, 27203 (2010)
- *Rev. Mod. Phys.* **82**, 2731 (2010)

Photon Energy



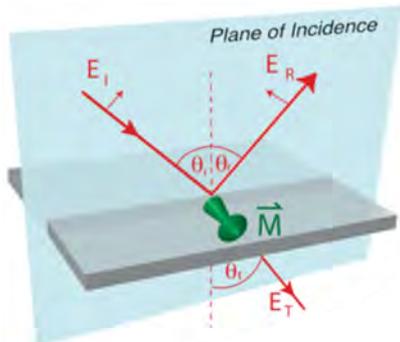
UNIQUE ADVANTAGES

- Element selective
- Multiple elements simultaneously
- $\ll 10$ fs time resolution, no jitter
- Span entire M edges, L edges



attosecond pulses

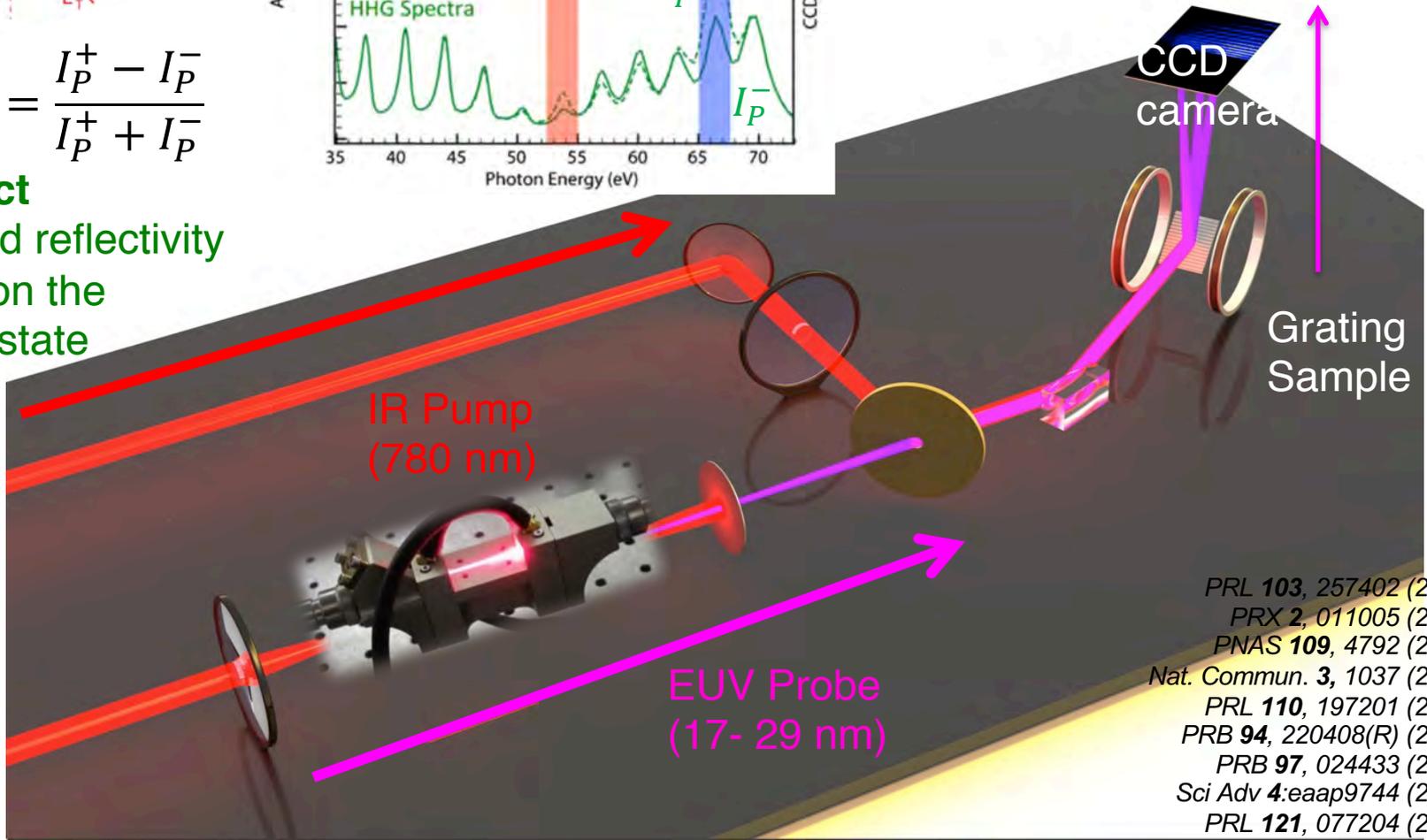
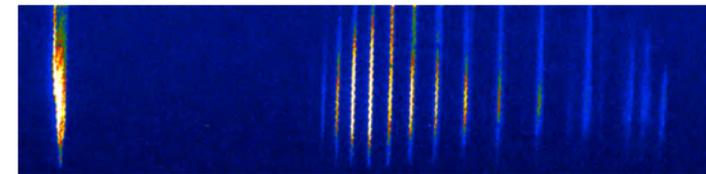
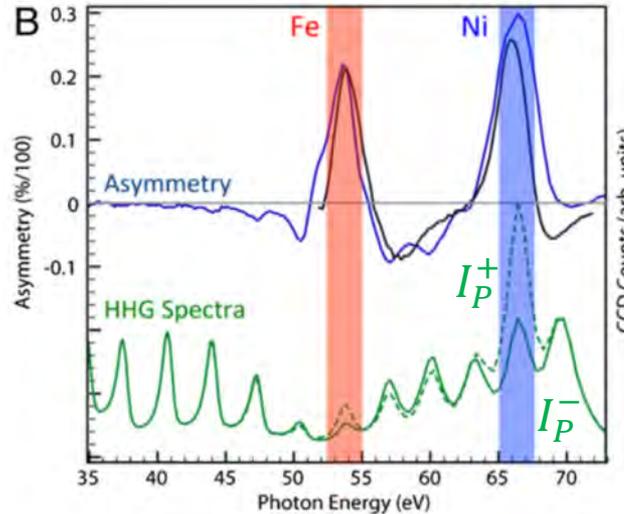
Rundquist et al, *Science* **280**, 1412 (1998)
 Chen et al., *PRL* **105**, 173901 (2010)
 Popmintchev et al, *Science* **336**, 1287 (2012)



$$\vec{M} \propto A = \frac{I_P^+ - I_P^-}{I_P^+ + I_P^-}$$

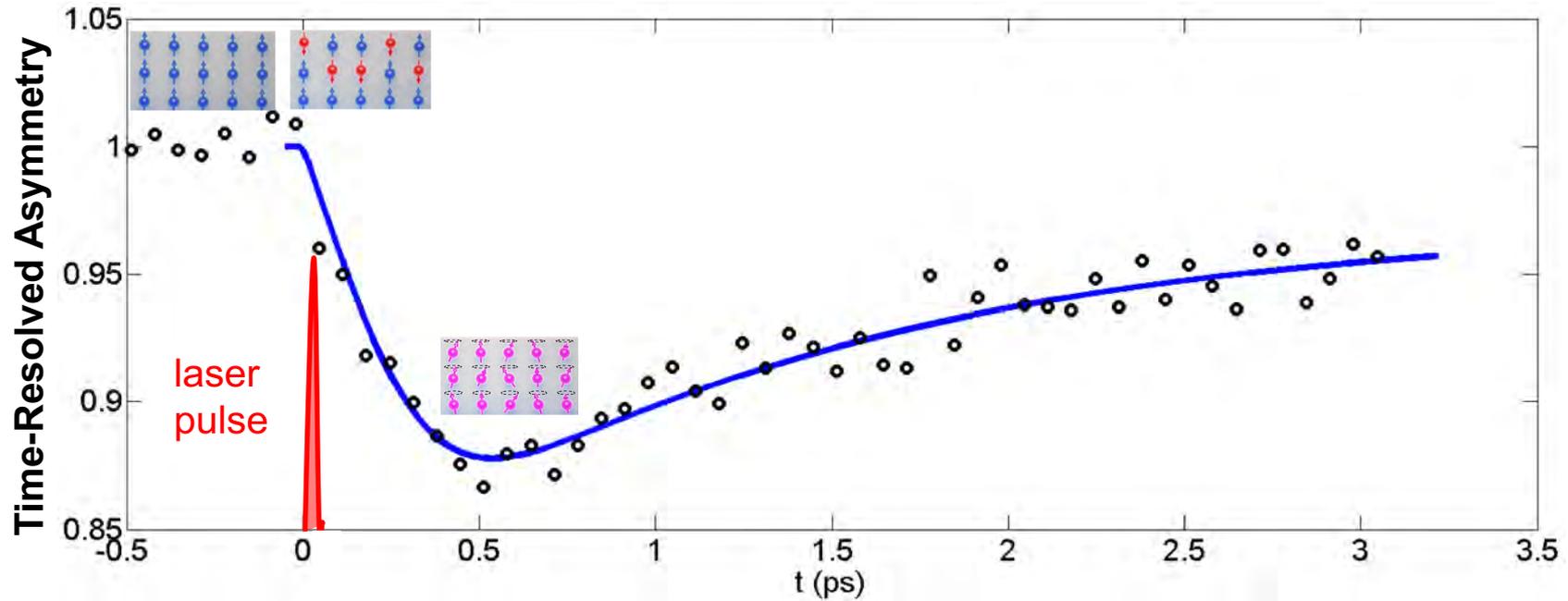
Kerr Effect

p-polarized reflectivity depends on the magnetic state



PRL **103**, 257402 (2009)
 PRX **2**, 011005 (2012)
 PNAS **109**, 4792 (2012)
 Nat. Commun. **3**, 1037 (2012)
 PRL **110**, 197201 (2013)
 PRB **94**, 220408(R) (2016)
 PRB **97**, 024433 (2018)
 Sci Adv **4**:eaap9744 (2018)
 PRL **121**, 077204 (2018)

Direct manipulation of spins: ultrafast spin excitation in two different materials



20 fs

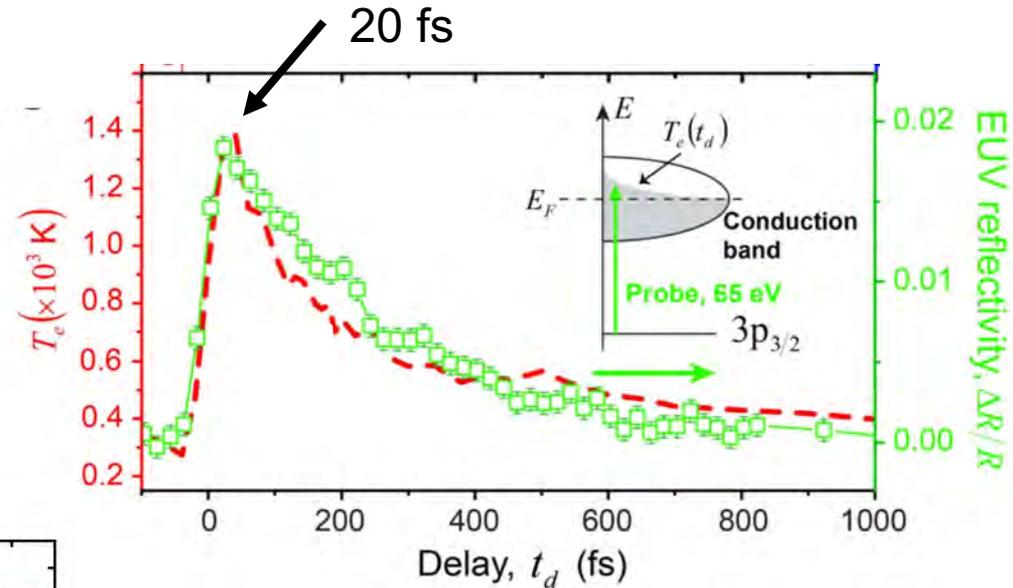
200 fs

Ferromagnetic Nickel

Half-metal Heusler
Alloy: Co_2MnGe^7

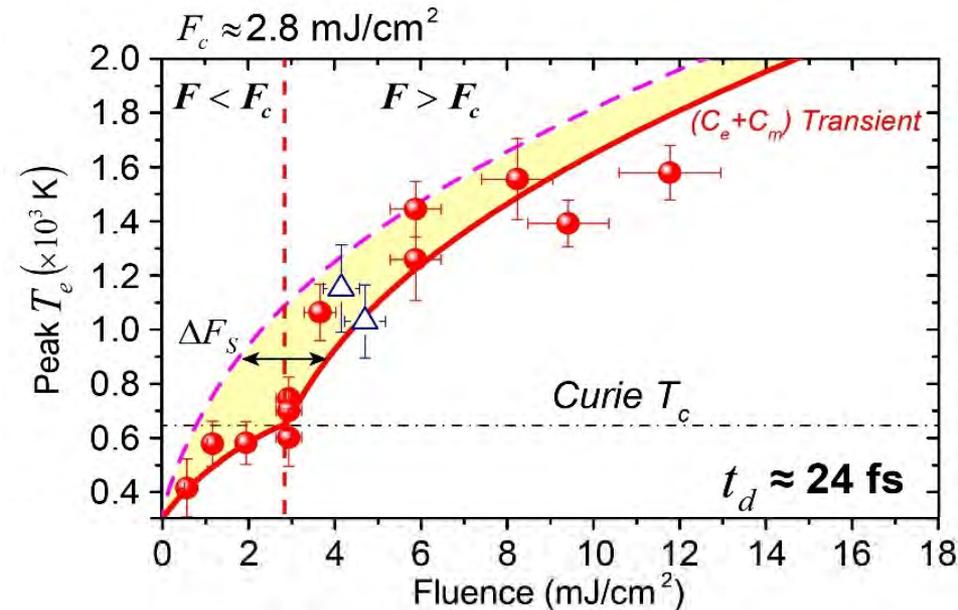
Probed spin dynamics in Ni using tabletop EUV high harmonics

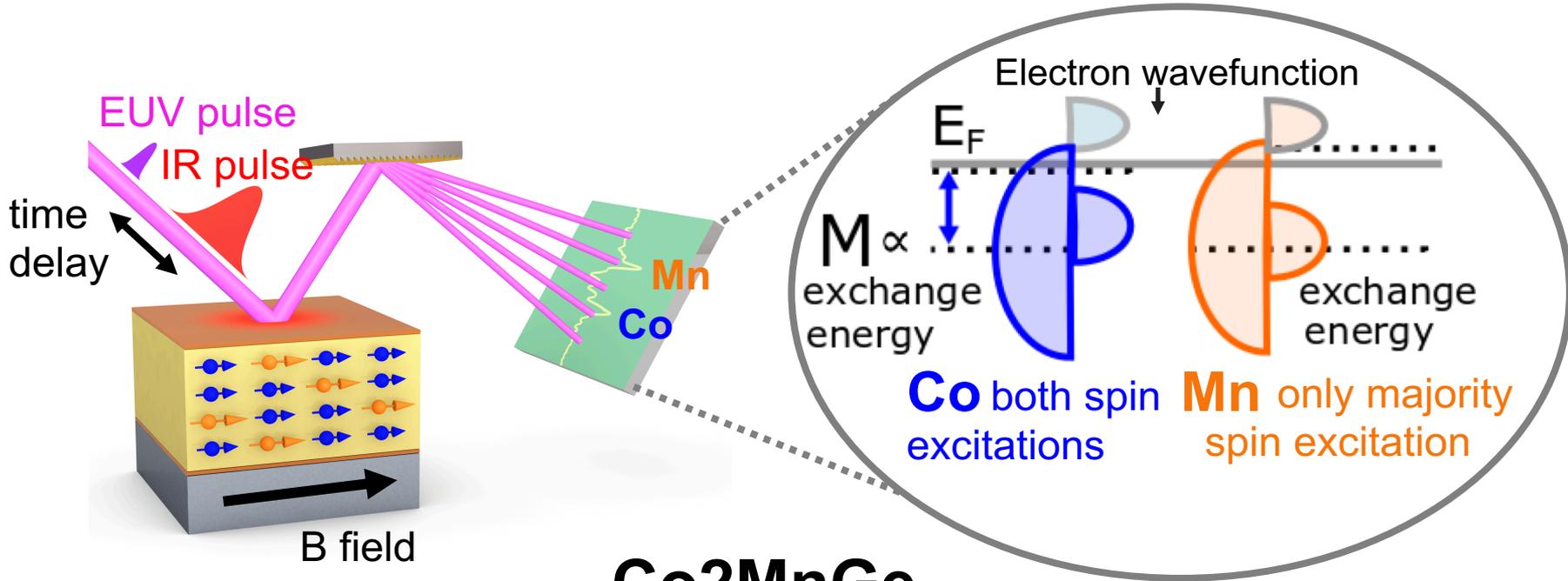
- (ARPES) to probe exchange splitting, electron temperature
- (MOKE) to probe magnetic order, reflectivity



Observed very surprising spin dynamics

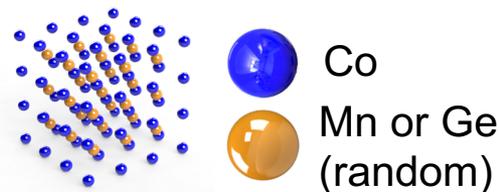
- Light excitation couples to spin system within 20fs after the pump pulse (new timescale for field) (ARPES)
- Initial super-excited state dictates future dynamics of the material for the full range of pump fluences (MOKE)





Co₂MnGe

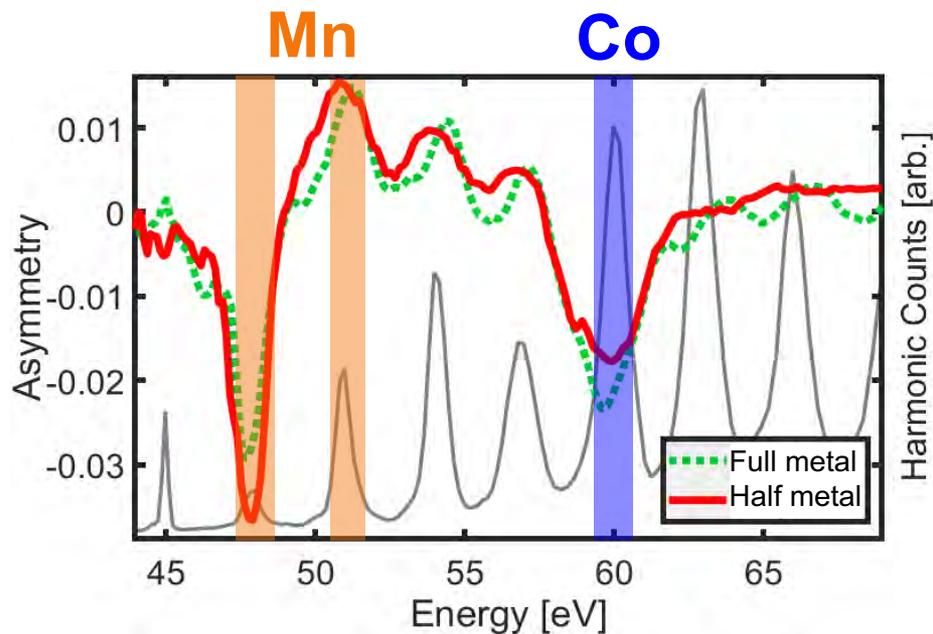
- Half metallic: one spin state conducting, the other insulating
- Half-metal gap in Mn is bigger than in Co
- Expect to see slower dynamics due to reduced scattering processes



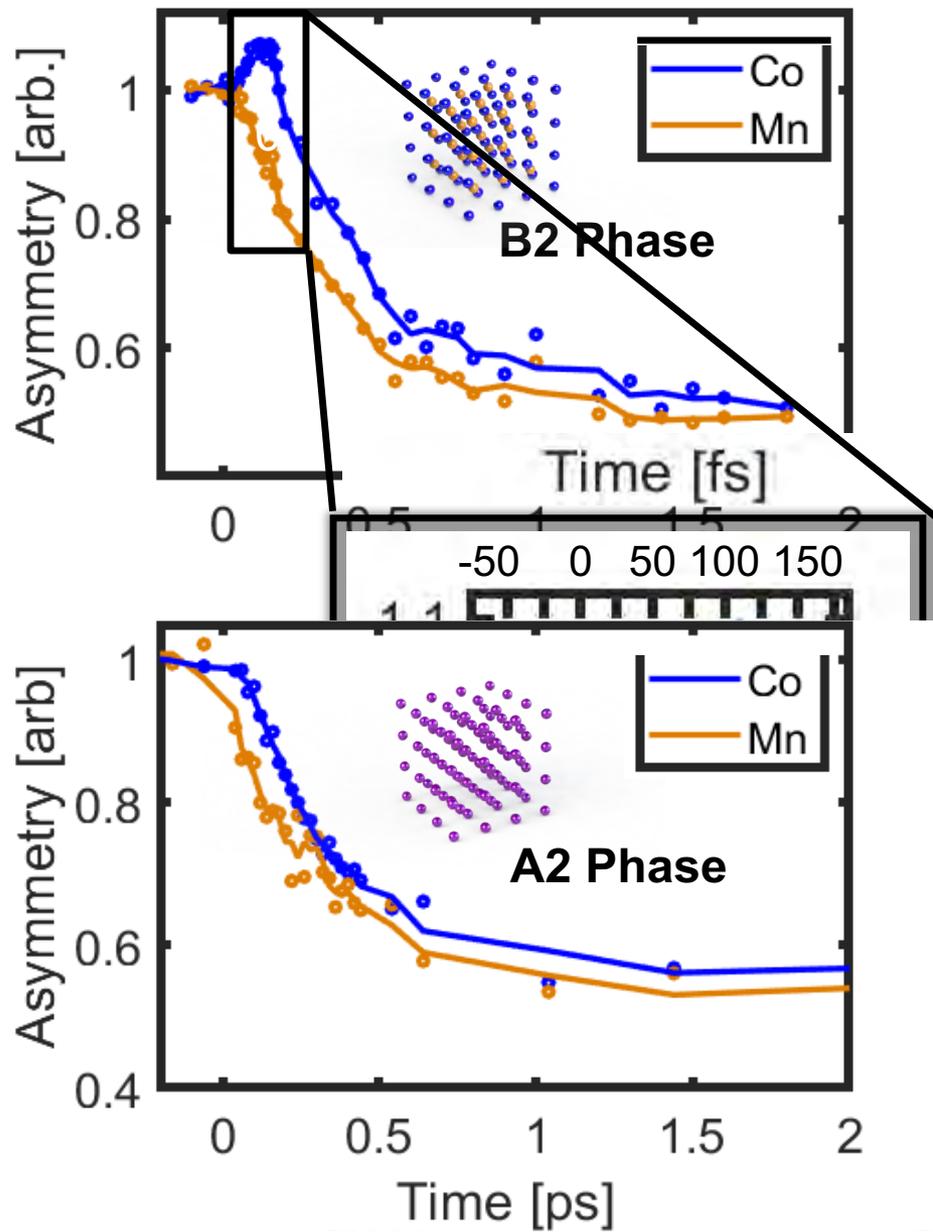
Half-metal Phase

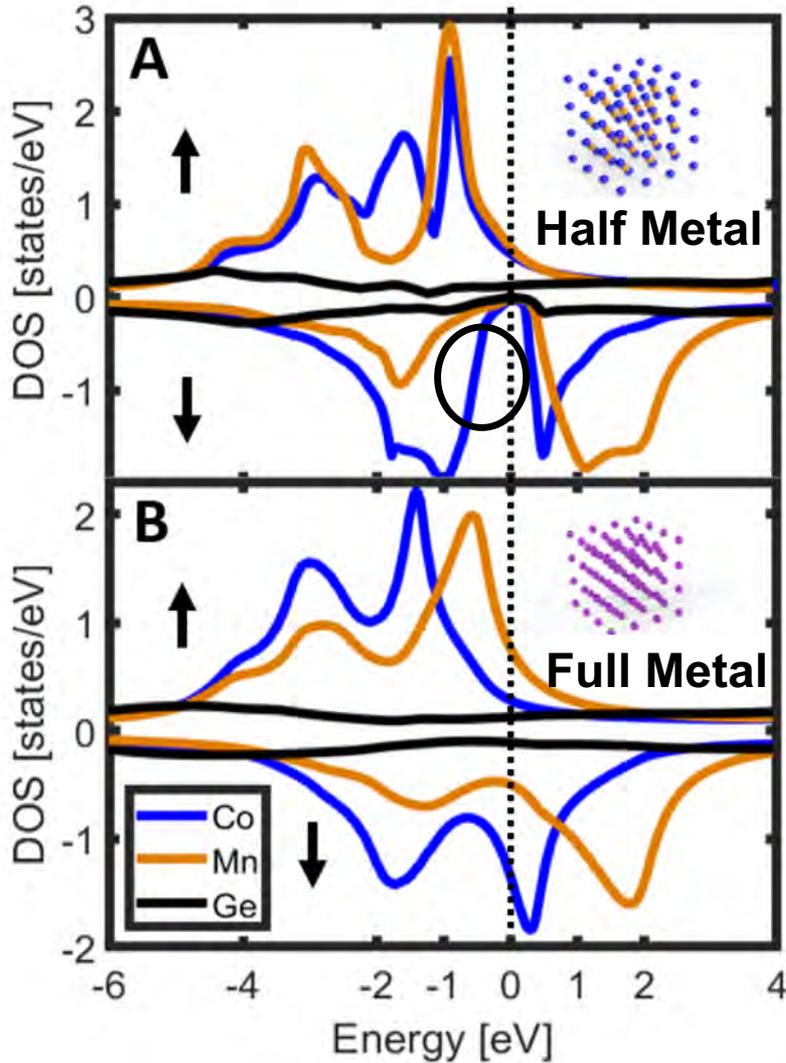


Full Metal Phase

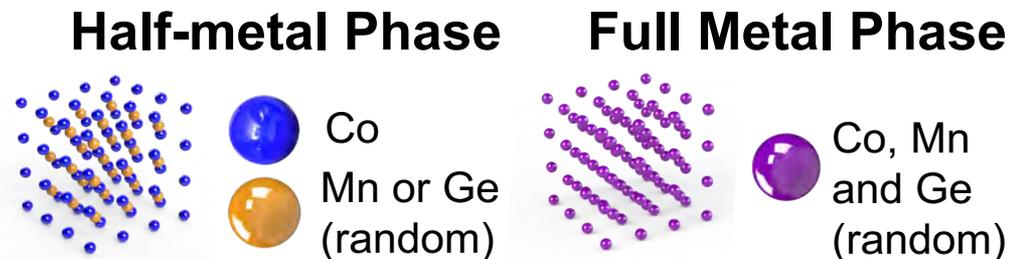
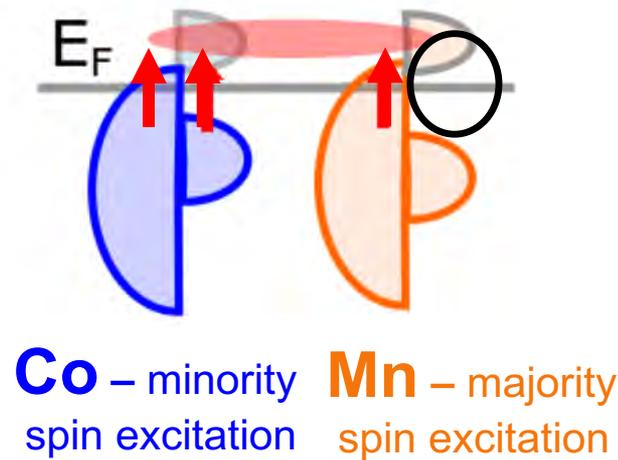


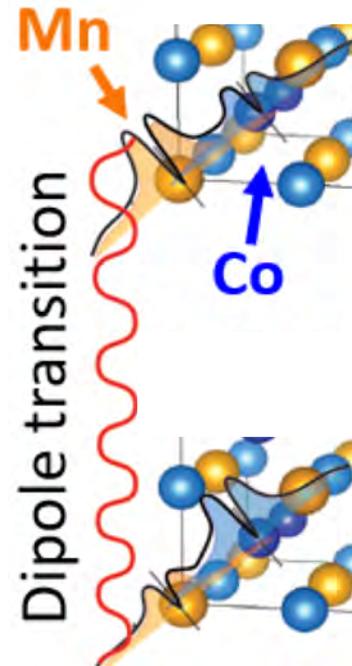
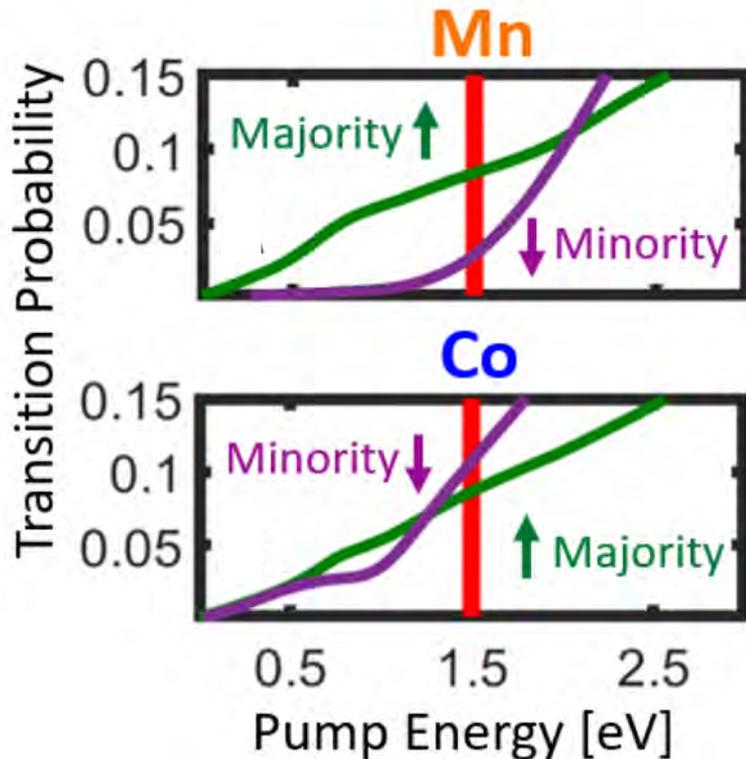
- Element specific measurements reveal transfer of magnetization between elemental sublattices ***within duration of pump pulse***
- Enhancement only present in half metal
- **Purely optically induced**





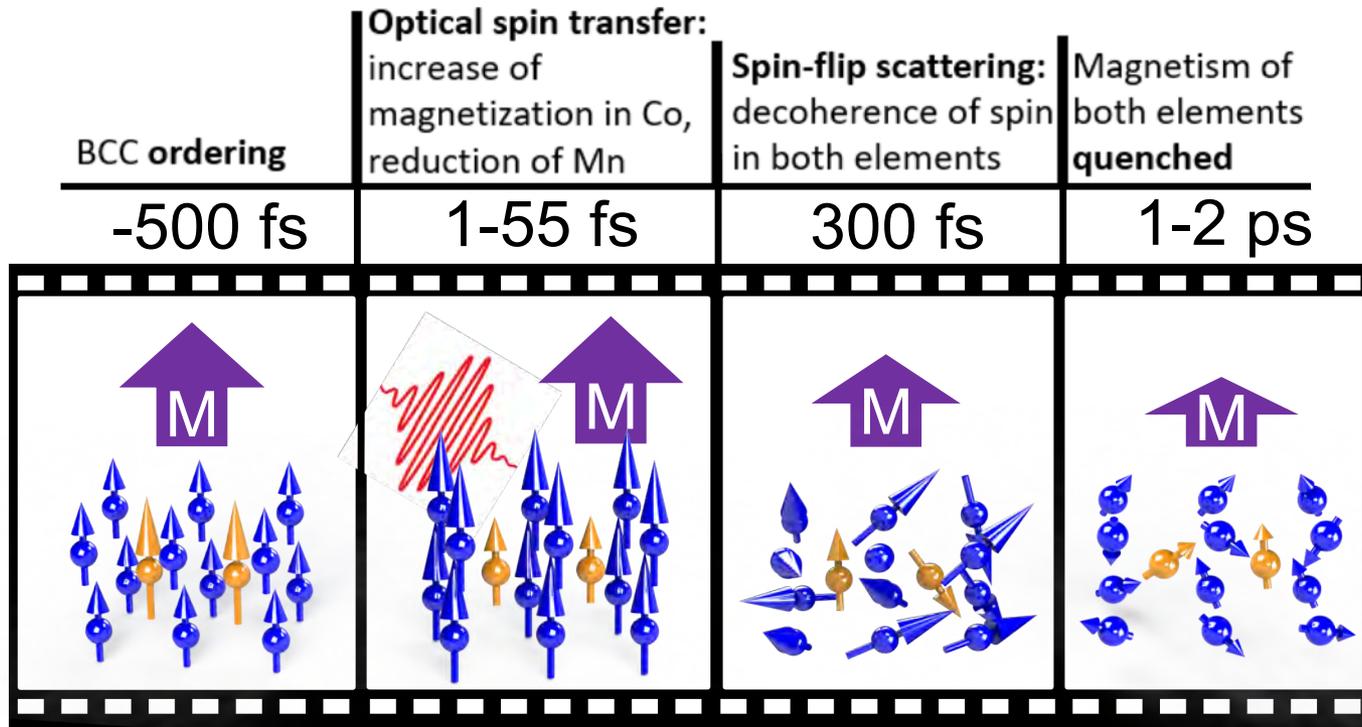
- **Full metal phase:** both spin states of both elements excited
- **Half metal phase:** Only minority carriers excited in Mn while both carriers are excited in Co





Microscopic Process

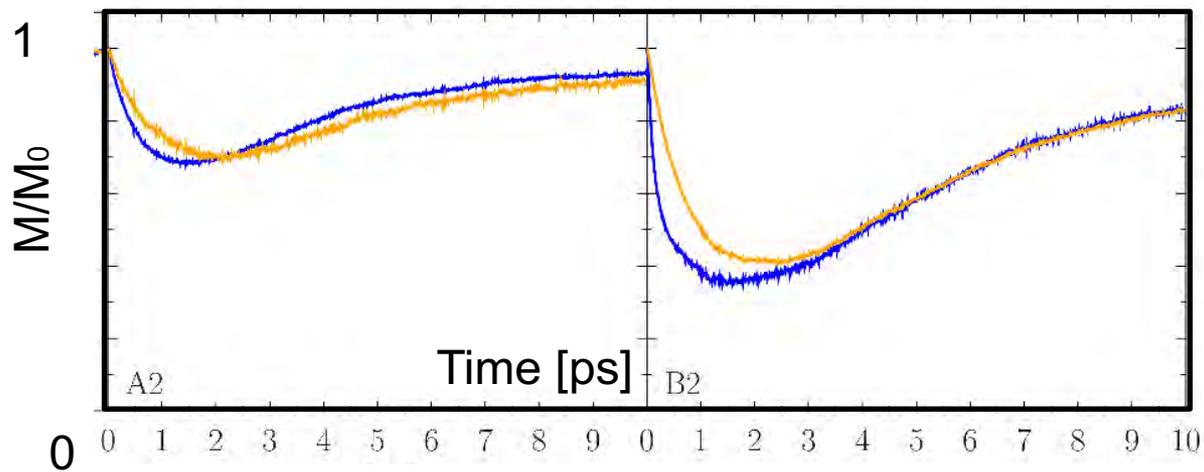
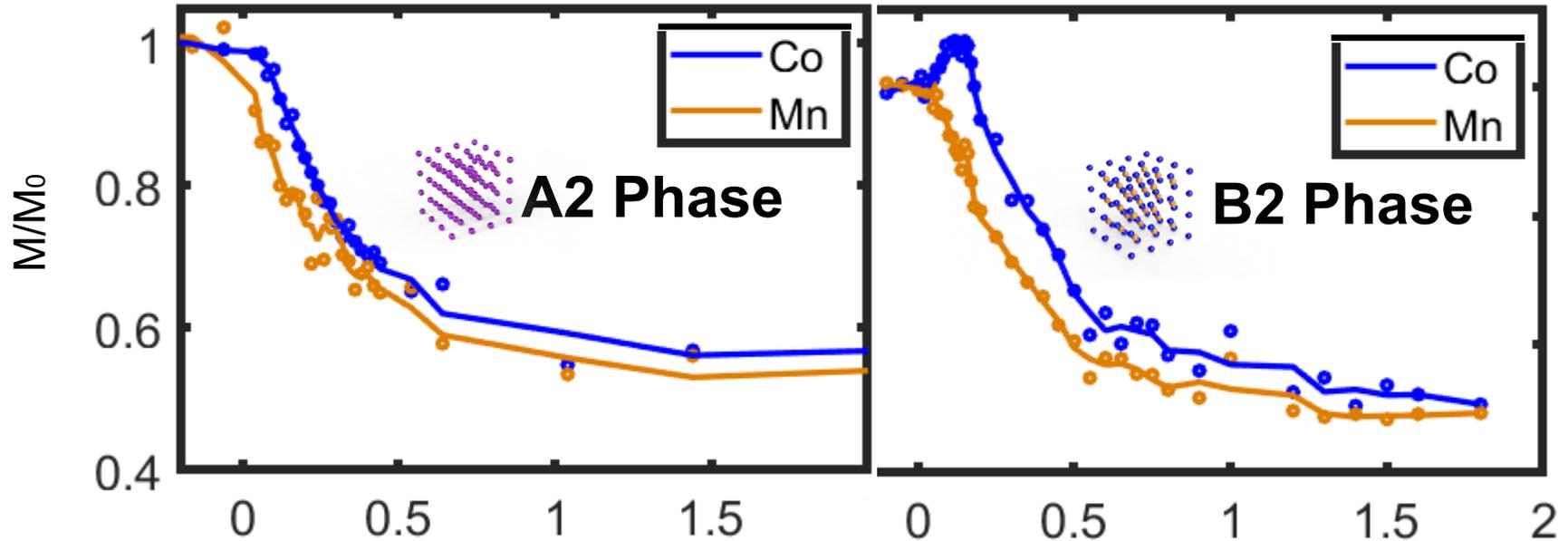
1. Preferential absorption of a single spin state in the valence band of Mn (wavefunction has primarily Mn character, electrons more localized)
2. Direct optical transfer via the conduction band (wavefunction has primarily Co character, electrons more itinerant)



- **Transient enhancement of Co magnetization *within the duration of the pump pulse***
 - Multiple measurements confirm existence of the enhancement occurs only during pump pulse duration
 - **Enhancement possible *only for the half-metallic phase***
 - Direct measurement of nearly-instantaneous spin manipulation

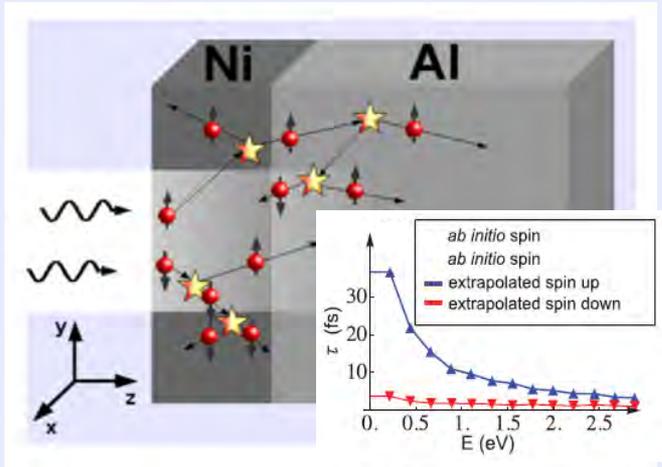
Unlike previous ultrafast measurements: behavior is not captured by LLG modelling!

Experiments



**Atomic
LLG**

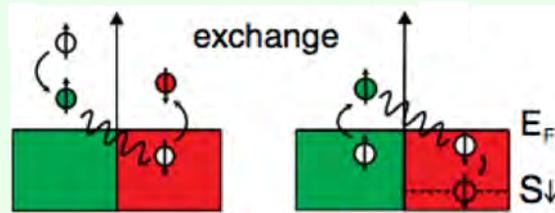
≈5-40 fs



Superdiffusive spin transport

Theory: PRL **105**, 027203 (2010)
Our work: PRL **110**, 197201 (2013)

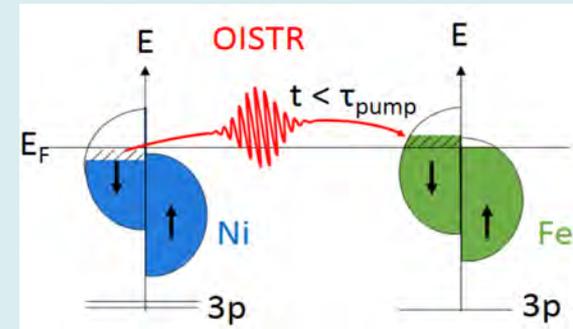
<20 fs



Spin flip scattering

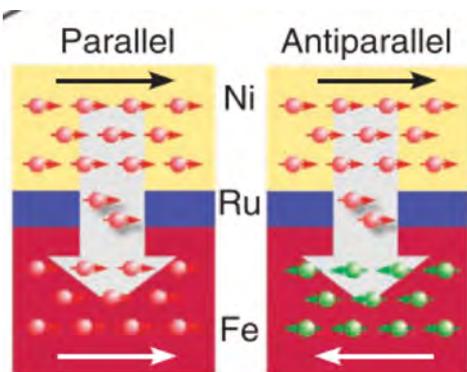
Theory: Nature Materials **9**, 259 (2010)
Our work: PRL **121**, 077204 (2018)
 Tengdin, P. et al. Sci. Adv. **4**, 9744 (2018)

<few-fs

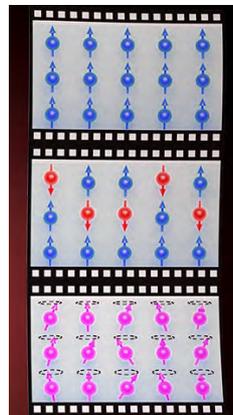


Ultrafast light-induced spin transfer

Theory: Nano Lett. **18**, 1842 (2018)
Our work: Tengdin et al., in press (2019)
 Hofherr et al., in press (2019)



FM/NM layers



ARPES/MOKE
 Measures:
 FM Nickel



Co₂MnGe Alloy



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 Strefan Mathias
 Lukas Hellbrück



GRPF &
 JILA Physics Frontier
 Center



Moore Foundation



DOE X-ray Scattering
 Program

We consider classical atomic magnetic moments $\{m_i\} = m_i\{e_i\}$ at site i . The dynamics of $\{m_i\}$ is governed by the atomistic Landau-Lifshitz-Gilbert (aLLG) equation (Antropov)

$$\frac{\partial m_i}{\partial t} = m \times \left(-\gamma(B_i + b_i \frac{\alpha}{m} \frac{\partial m_i}{\partial t}) \right)$$

Where γ is the gyromagnetic ratio and $B_i = \frac{\partial H}{\partial m_i}$ is the effective precession field related to the spin-Hamiltonian, employing a Heisenberg model where -

$$H = - \sum_{i,j} J_{i,j} m_i \cdot m_j - \mu_B B \sum_i m_i$$

Here, the magnetic moments at site i and j are coupled by the exchange parameter J_{ij} . The values of the Heisenberg exchange parameters were obtained from first principles electronic structure theory. In the expression above, B is the external magnetic field. From the fluctuation-dissipation theorem, thermal fluctuations enter by a stochastic field, b_i , that fulfills white noise properties, such as $\langle b_i \rangle = 0$ and $\langle b_i^\mu(t) b_k^\nu(t') \rangle = D \delta_{ij} \delta_{\mu\nu} \delta(t - t')$, with the fluctuation amplitude $D = 2\alpha k_b T / \gamma m$. The dissipation part enters the equation of motion via a viscous damping part scaled by the Gilbert damping constant α .