ANTIFERROMAGNETIC SPINTRONICS FROM TOPOLOGY TO NEUROMORPHIC COMPUTING



Hideo S Ohno F

Shunsuke Fukami

Workshop October 7th - 10th, 2019 Schloss Waldthausen, Mainz, Germany

ORGANIZERS:



Jairo Sinova



Elena [Hilp k

Denise Kornbrust

Robert-Andre Vettel Hjördis Karin Pusch Everschor-Sitte

Big data storage



Internet (PC & cloud IT)



Internet of Things (edge IT)







Sony/IBM tape (330TB)



Seagate HDD (16TB)



Sony optical disc (3.3TB)

Samsung Flash SSD (30TB)

Big data processing: 1. Artificial neural networks

Mass applications – Google Brain (2012 – image recognition, 2016 – language translation)



Synchronous: memory & processor under global clock

Hebbian learning $\Delta W_{ij} \sim X_i \cdot Y_j$

Big data processing: 1. Artificial neural networks

Mass applications – Google Brain (2012 – image recognition, 2016 – language translation)



Synchronous: memory & processor under global clock

Big data processing: 1. Artificial neural networks



Logic-in-memory

Prezioso et al. Nature 521, 61 (2015) Hu et al. Nature Elec. 1, 52 (2018) Ambrogio et al. Nature 558, 60 (2018)

Non-CMOS resistive memories



Non-CMOS multi-level resistive memories: electrical switching



Non-CMOS multi-level resistive memories: optical switching

down to ~100 fs, ~mJ/cm² pulses



Resistive



Wright et al. Adv. Mater. 23, 3408 (2011)



Chakravarty et al. Appl. Phys. Lett. 114, 192407 (2019), Nemec, Fiebig, Kampfrath, Kimel, Nature Phys. 14, (2018)

Kaspar et al. arXiv:1909.09071

Big data processing: 2. Spiking neural networks

Asynchronous spiking; order/delay between spikes; energy saving



Spiking time dependent plasticity of synapse



Leaky-sum-and-fire neuron

j=1



Gerstner & Kistler, Spiking Neuron Models, Cambridge University Press (2002)

 $u_i(t)$

Kurenkov et al. Adv. Mater. 31, 1900636 (2019)

Big data processing: 2. Spiking neural networks

Asynchronous spiking; order/delay between spikes; energy saving





Leaky-sum-and-fire neuron



Benjamin et al. Proceedings of the IEEE 102, 699 (2014)

Non-CMOS spiking analog memories



Kurenkov et al. Adv. Mater. 31, 1900636 (2019)

Magnetic vs. non-magnetic origin of resistive switching



Conductivity

$$\vec{j} = \overleftarrow{\sigma} \vec{E}$$

Linear response:
Invariant under inversion
$$P\vec{\sigma} = \vec{\sigma}$$

$$\begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} \sigma_{xx}^{S} & \sigma_{xy}^{S} & \sigma_{xz}^{S} \\ \sigma_{xy}^{S} & \sigma_{yy}^{S} & \sigma_{yz}^{S} \\ \sigma_{xz}^{S} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \end{bmatrix} + \begin{bmatrix} 0 & \sigma_{xy}^{a} & \sigma_{xz}^{a} \\ -\sigma_{xy}^{a} & 0 & \sigma_{yz}^{a} \\ -\sigma_{xz}^{a} & -\sigma_{yz}^{a} & 0 \end{bmatrix}$$

$$\begin{bmatrix} \vec{\sigma}_{xz} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \\ \vec{\sigma}_{xz}^{S} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \end{bmatrix} + \begin{bmatrix} 0 & \sigma_{xy}^{a} & \sigma_{xz}^{a} \\ -\sigma_{xz}^{a} & -\sigma_{yz}^{a} & 0 \end{bmatrix}$$

$$\begin{bmatrix} \vec{\sigma}_{xz} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \\ \vec{\sigma}_{xz}^{S} & \sigma_{zz}^{S} & \sigma_{zz}^{S} \end{bmatrix} + \begin{bmatrix} 0 & \sigma_{xy}^{a} & \sigma_{xz}^{a} \\ -\sigma_{xz}^{a} & -\sigma_{yz}^{a} & 0 \end{bmatrix}$$

$$\begin{bmatrix} \vec{\sigma}_{xz} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \\ \vec{\sigma}_{xz}^{S} & \sigma_{zz}^{S} & \sigma_{zz}^{S} \end{bmatrix} + \begin{bmatrix} \vec{\sigma}_{xz}^{S} & \sigma_{xz}^{S} \\ \vec{\sigma}_{xz}^{S} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \end{bmatrix}$$

$$\begin{bmatrix} \vec{\sigma}_{xz}^{S} & \sigma_{zz}^{S} & \sigma_{zz}^{S} \\ \vec{\sigma}_{xz}^{S} & \sigma_{zz}^{S} & \sigma_{zz}^{S} \end{bmatrix} + \begin{bmatrix} \vec{\sigma}_{xz}^{S} & \sigma_{xz}^{S} \\ \vec{\sigma}_{xz}^{S} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \end{bmatrix}$$

vector:
$$P\vec{j} = -\vec{j}$$

 $P\vec{E} = -\vec{E}$
pseudo-vector: $P\vec{h} = \vec{h}$

$$\vec{j} = \overleftrightarrow{\sigma} \vec{E}$$



Ordinary conductivity

$$\vec{j} = \overleftrightarrow{\sigma} \vec{E}$$

Linear response:
Invariant under inversion
$$P\vec{\sigma} = \vec{\sigma}$$

$$\begin{pmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{pmatrix} = \begin{pmatrix} \sigma_{xx}^{s} & \sigma_{xy}^{s} & \sigma_{xz}^{s} \\ \sigma_{xy}^{s} & \sigma_{yy}^{s} & \sigma_{yz}^{s} \\ \sigma_{xz}^{s} & \sigma_{yz}^{s} & \sigma_{zz}^{s} \end{pmatrix} + \begin{pmatrix} 0 & \sigma_{xy}^{a} & \sigma_{xz}^{a} \\ -\sigma_{xy}^{a} & 0 & \sigma_{yz}^{a} \\ -\sigma_{xz}^{a} & -\sigma_{yz}^{a} & 0 \end{pmatrix}$$
Effective time-reversal symmetry PT
Wadley et al., Science '16, Bodnar et al., Nature Commun. '18, Meinert et al. Phys. Rev. Appl. '18, Zhou et al. Phys. Rev. Appl. '18}



Anistropic magnetoresistance

Weak spin-orbit (Dirac) interaction crystal spin

Shick et al. PRB '10, Park et al. Nature Mater '11, Wang et al. Phys. Rev. Lett. '12, Marti, TJ et al. Nature Mater. '14, Moriyama et al. Appl. Phys. Lett. '15

Ordinary conductivity



Nano-texture magnetoresistance

Strong exchange (Coulomb) interaction

Maca et al. PRB 96, 094406 (2017), Kaspar et al. arXiv:1909.09071



Disorder resistance







Wadley et al., Science '16, Bodnar et al., Nature Commun. '18, Meinert et al. Phys. Rev. Appl. '18, Zhou et al. Phys. Rev. Appl. '18



$$\vec{j} = \overleftrightarrow{\sigma} \vec{E}$$

Linear response:
Invariant under inversion
$$P\vec{\sigma} = \vec{\sigma}$$

$$\begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} \sigma_{xx}^{S} & \sigma_{xy}^{S} & \sigma_{xz}^{S} \\ \sigma_{xy}^{S} & \sigma_{yy}^{S} & \sigma_{yz}^{S} \\ \sigma_{xz}^{S} & \sigma_{yz}^{S} & \sigma_{zz}^{S} \end{bmatrix} + \begin{bmatrix} 0 & \sigma_{xy}^{a} & \sigma_{xz}^{a} \\ -\sigma_{xy}^{a} & 0 & \sigma_{yz}^{a} \\ -\sigma_{xz}^{a} & -\sigma_{yz}^{a} & 0 \end{bmatrix}$$

$$\begin{bmatrix} Hall \\ \vec{j}_{H} = \vec{h} \times \vec{E} \\ \vec{h} = (\sigma_{zy}^{a}, \sigma_{xz}^{a}, \sigma_{yx}^{a}) \\ Hall (pseudo)-vector \\ T\vec{h}(\vec{s}) = \vec{h}(-\vec{s}) = -\vec{h}(\vec{s}) \end{bmatrix}$$

Odd under time-reversal



Suzuki et al. Phys. Rev. B 95, 094406 (2017) Zelezny et al. PRL 119, 187204 (2017) Net ferromagnetic (pseudo)-vector

$$\vec{j}_{H} = \vec{h} \times \vec{E}$$
$$\vec{h} = (\sigma_{zy}^{a}, \sigma_{xz}^{a}, \sigma_{yx}^{a})$$
Hall (pseudo)-vector

 $T\vec{h}(\vec{s}) = \vec{h}(-\vec{s}) = -\vec{h}(\vec{s})$ Odd under time-reversal



No spin-orbit coupling

 \dot{h} invariant under pure spin rotation R^s_{arphi}



Net ferromagnetic (pseudo)-vector



Effective time-reversal symmetry $R_{\pi}^{s}T$

Suzuki et al. Phys. Rev. B 95, 094406 (2017) Zelezny et al. PRL 119, 187204 (2017)

 $\vec{j}_{H} = \vec{h} \times \vec{E}$ $\vec{h} = (\sigma_{zy}^{a}, \sigma_{xz}^{a}, \sigma_{yx}^{a})$ Hall (pseudo)-vector $T\vec{h}(\vec{s}) = \vec{h}(-\vec{s}) =$ $-\bar{h}(\vec{s})$ Odd under time-reversal



Spin-orbit coupling

 $ec{h}$ not invariant under pure spin rotation R^s_arphi

Net ferromagnetic (pseudo)-vector

$$\vec{j}_{H} = \vec{h} \times \vec{E}$$

$$\vec{h} = (\sigma_{zy}^{a}, \sigma_{xz}^{a}, \sigma_{yx}^{a})$$

Hall (pseudo)-vector

$$T\vec{h}(\vec{s}) = \vec{h}(-\vec{s}) = -\vec{h}(\vec{s})$$

Odd under time-reversal



No spin-orbit coupling

 \dot{h} invariant under pure spin rotation R^s_{arphi}



No (or negligible) net ferromagnetic moment



Suzuki et al. Phys. Rev. B 95, 094406 (2017) Zelezny et al. PRL 119, 187204 (2017)



Spin-orbit coupling

 \dot{h} not invariant under pure spin rotation R^s_{arphi}

No (or negligible) net ferromagnetic moment

Chen, Niu, MacDonald, PRL '14 Nakatsuji, Kiyohara, Higo, Nature '15 Nayak et al. Science Adv. '16



 $\vec{j}_{H} = \vec{h} \times \vec{E}$ $\vec{h} = (\sigma_{zy}^{a}, \sigma_{xz}^{a}, \sigma_{yx}^{a})$ Hall (pseudo)-vector

 $T\vec{h}(\vec{s}) = \vec{h}(-\vec{s}) = -\vec{h}(\vec{s})$ Odd under time-reversal



Šmejkal et al. arXiv:1901.00445

Spin-orbit coupling

 \dot{h} not invariant under pure spin rotation R^s_arphi

No (or negligible) net ferromagnetic moment

 $\vec{j}_H = \vec{h} \times \vec{E}$ $\vec{h} = (\sigma_{zy}^a, \sigma_{xz}^a, \sigma_{yx}^a)$ Hall (pseudo)-vector $T\vec{h}(\vec{s}) = \vec{h}(-\vec{s}) = -\vec{h}(\vec{s})$

Odd under time-reversal



Strong exchange (Coulomb) coupling



Even under time-reversal



No magnetism in graphene No *PT* symmetry in ferromagnets

Smejkal, Mokrousov, Yan, MacDonald, Nature Phys. 14, 242 (2018)

Tang et al. Nature Phys. 12, 1100 (2016) Smejkal et al. PRL 118, 106402 (2017)

Even under time-reversal

Topological Dirac semimetal



No magnetism in graphene No *PT* in ferromagnets

Smejkal, Mokrousov, Yan, MacDonald, Nature Phys. 14, 242 (2018)

Tang et al. Nature Phys. 12, 1100 (2016) Smejkal et al. PRL 118, 106402 (2017)

Odd under time-reversal





Smejkal, Mokrousov, Yan, MacDonald, Nature Phys. 14, 242 (2018)

Data – the oil of the 21st century

Data mining

Neuromorphic

amazon

Spiking analog

Logic-in-memory

UBER

Current & light

Imaging

Symmetry & topology

Non-CMOS storage & memory

