

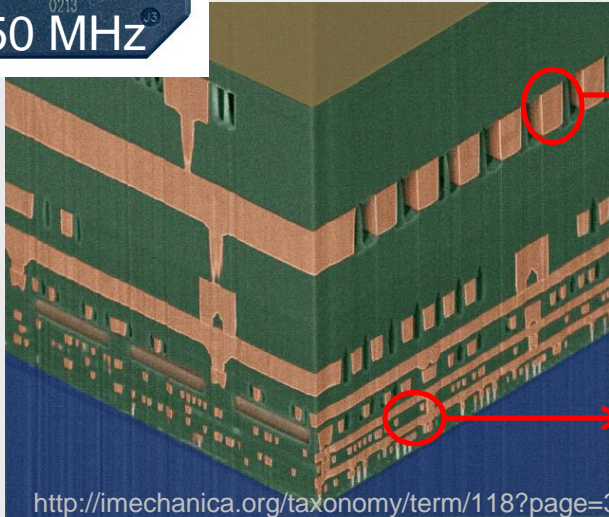


Exotic spin-orbit torques from topological materials and ferrimagnets

Hyunsoo Yang

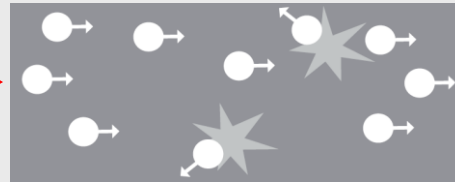
Department of Electrical and Computer Engineering
National University of Singapore

Charge electronics → Spin electronics

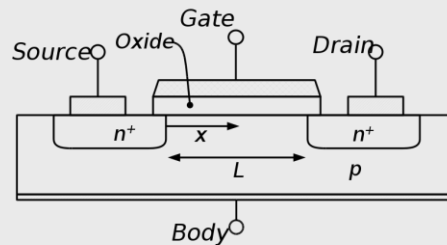


<http://imechanica.org/taxonomy/term/118?page=3>

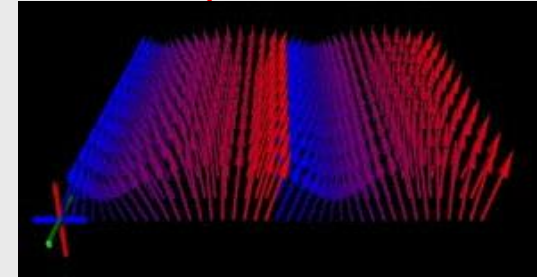
Information transfer
= electron transfer



Information processing
= processing electron flow



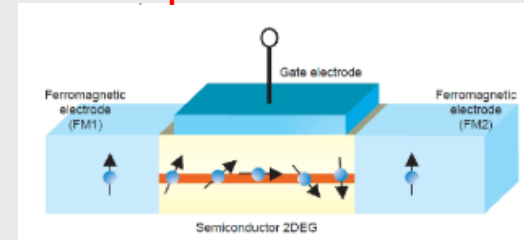
Spin wave



MTJ memory

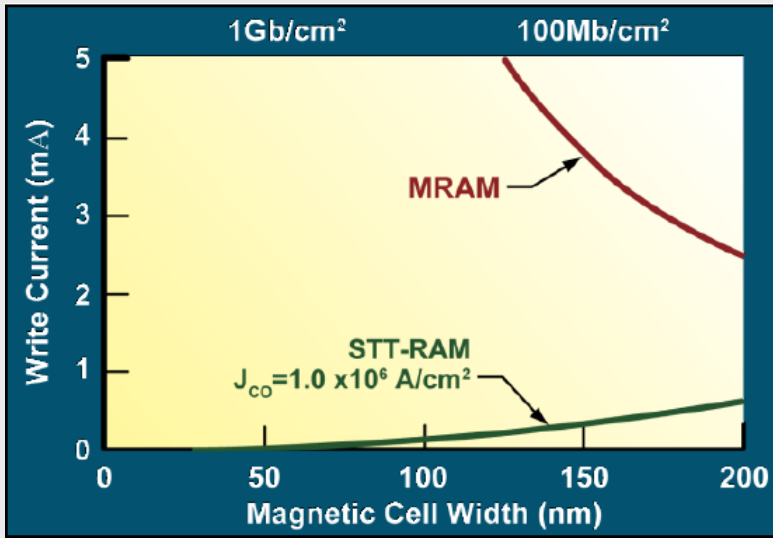


Spin transistor



Charge transfer and processing energy loss is huge
→ All spin electronics

Spin torque MRAM



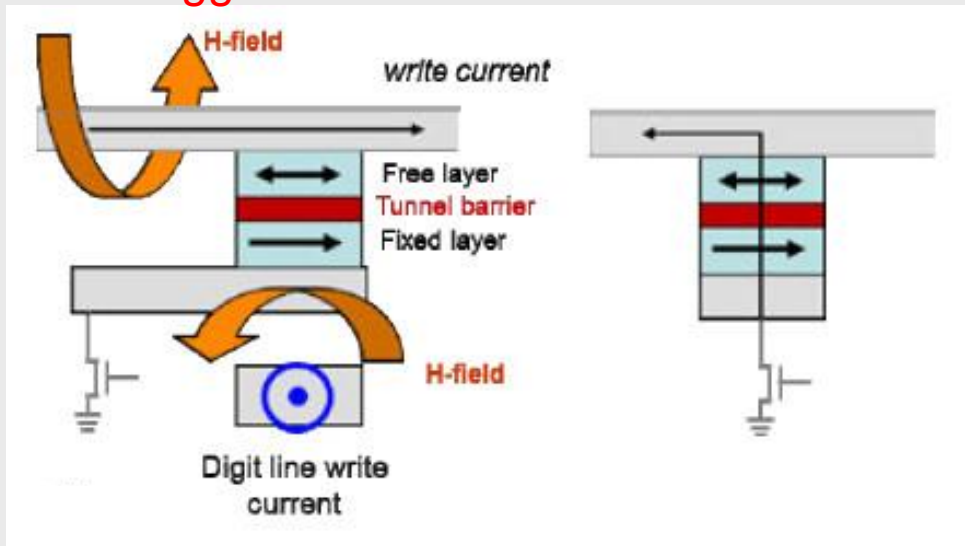
- Present STT-MRAM (Everspin) uses 90 nm node
- GF started to use 40, 28, 14 nm node
- TSMC, IBM, Samsung, TDK, Hynix, Sony, Avalanche, Toshiba, Intel, Qualcomm...

2019 March
Samsung
28 nm FD-SOI
embedded



Toggle-MRAM

STT-MRAM



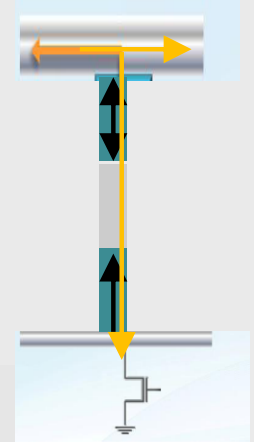
16 Mb

256 Mb



Everspin with
GlobalFoundries

SOT-MRAM

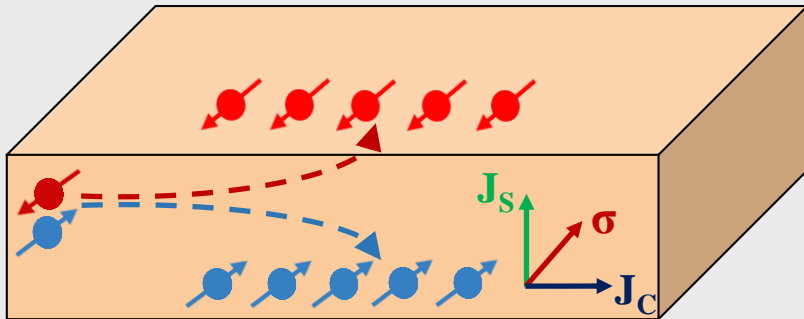


Spin Hall and Inverse Spin Galvanic effects

Spin Hall effect

- Spin orbit coupling in the bulk of non-magnet.
- Can be from intrinsic or extrinsic (impurities).
- Quantified by spin Hall angle (SHA or θ_{sh})

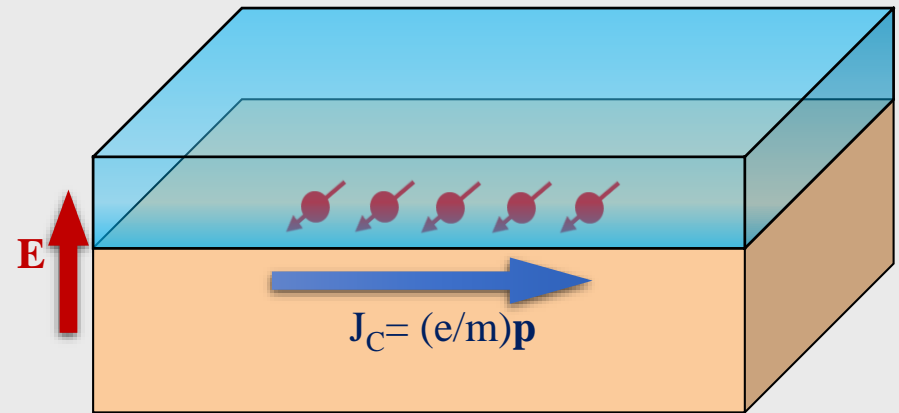
$$\vec{j}_s \propto \theta_{sh} (\vec{j}_c \times \vec{\sigma})$$



Rashba-Edelstein effect

- Structures with broken inversion symmetry
- Interfacial spin orbit coupling.
- Depends on the nature of the interface.
- Quantified by Rashba coefficient (\propto magnitude of $\mathbf{E}=-\nabla V$)

$$\hat{H}_{so} = \xi \hat{\sigma} \cdot (\nabla V \times \hat{p})$$

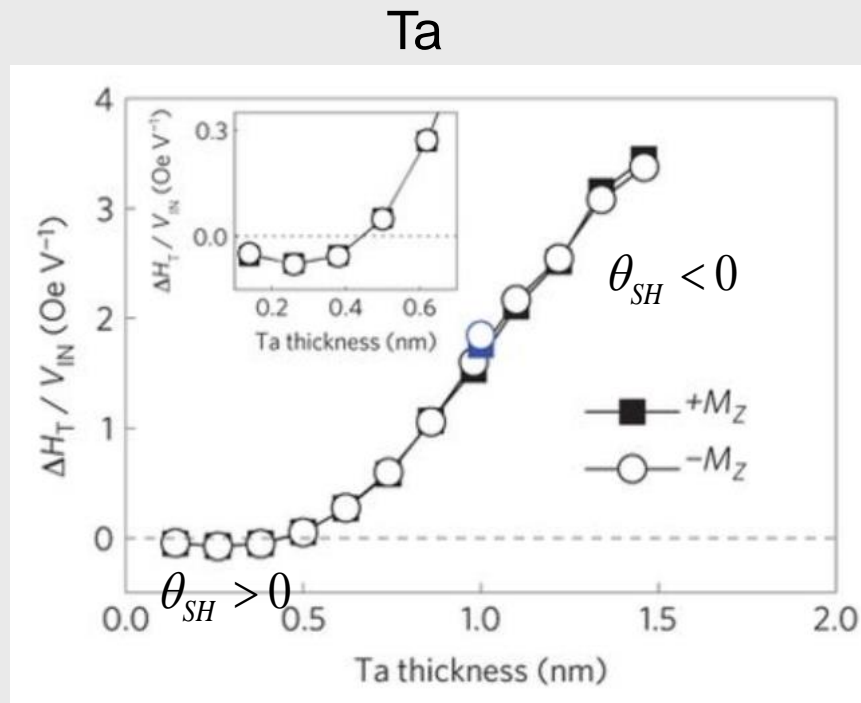


Is spin accumulation direction the same for both cases?

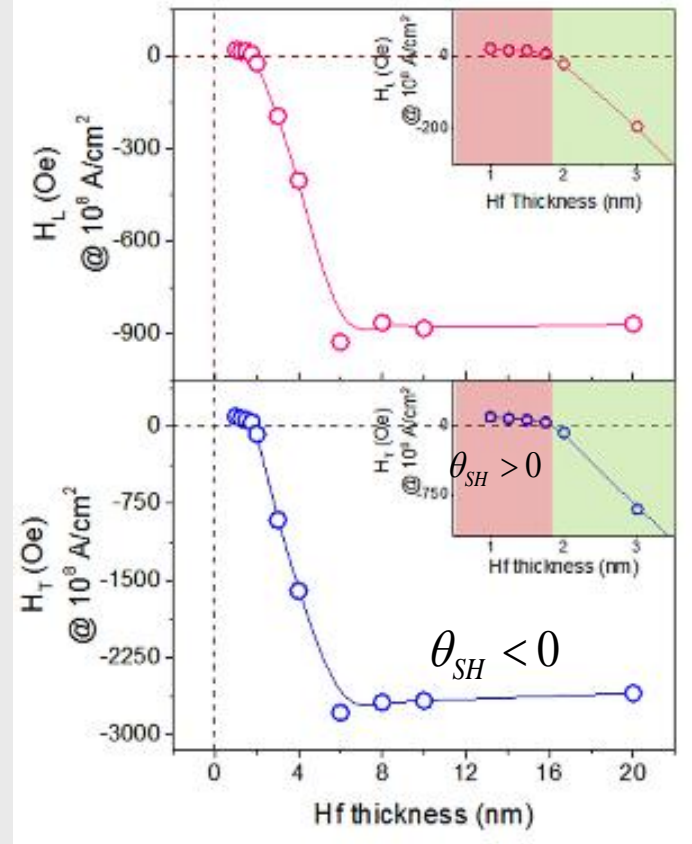
Theory: D'yakanov and Perel (1971), Hirsch (1999), Edelstein (2000), Zhang (2000), Murakami, Nagaosa, Zhang (2003), Sinova *et al.* (2004),...Manchon and Zhang (2008).
Experiments: (semiconductors) Ganichev *et al.* (2002), Kato *et al.* (2004), Wunderlich *et al.* (2005)
(metals) Valenzuela and Tinkham (2006), Saitoh *et al.* (2006), Kimura *et al.* (2007).

Sign change of effective spin Hall angle

Hf



Nature Mater. 12, 240 (2013)

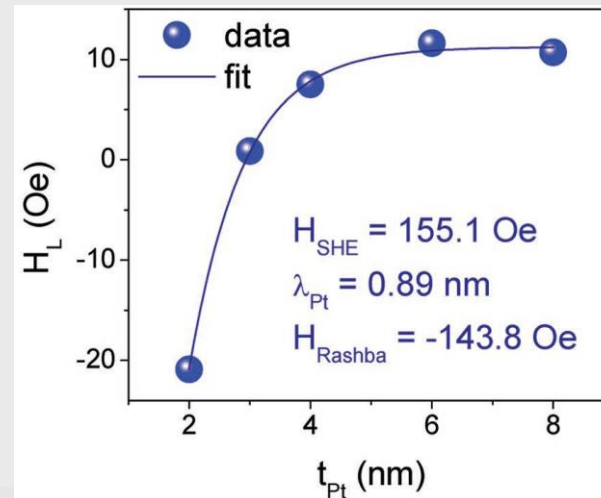
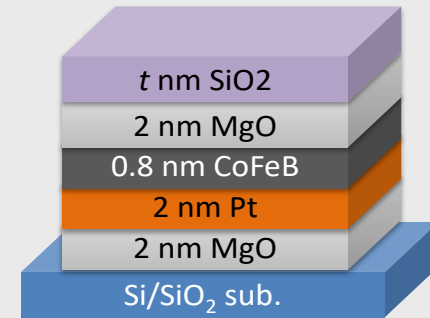
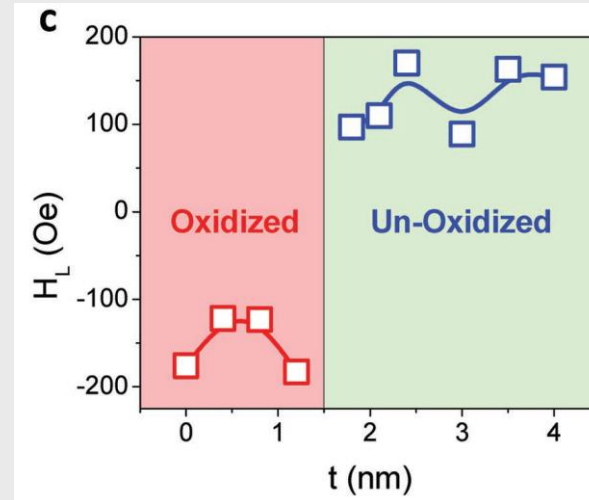
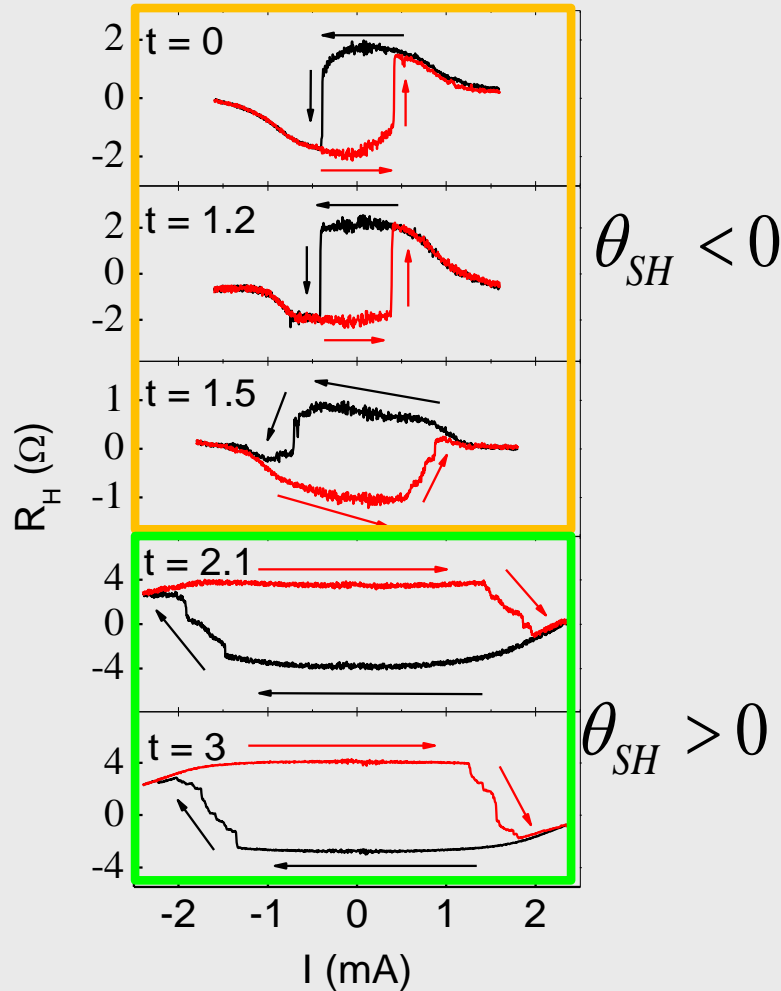


Appl. Phys. Lett. 108, 202406 (2016)

- The SOT polarity is a function of heavy metal thickness
- Thick regime: spin Hall dominant
- Thin regime: **interfacial** spin-orbit with opposite sign to spin Hall

Is the sign of θ_{SH} fixed for a given material?

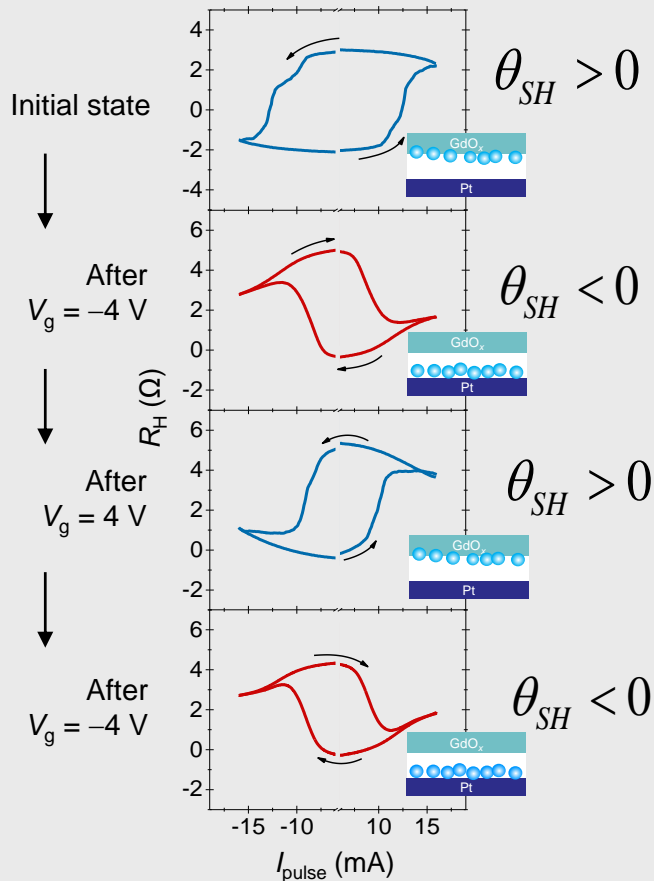
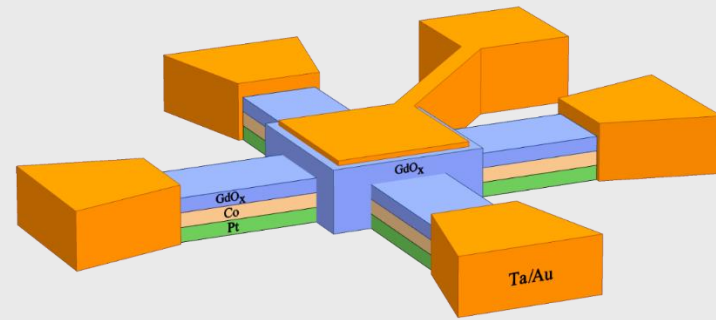
Reverse switching polarity by oxygen engineering



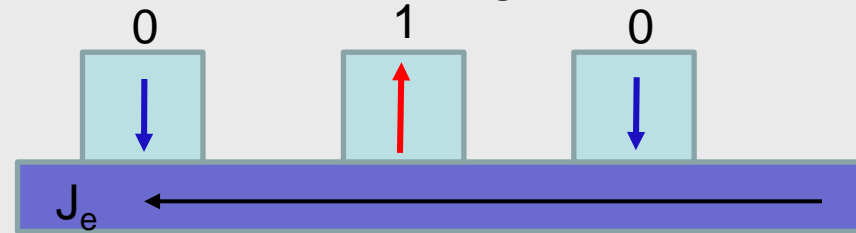
- Sign of spin Hall angle changes across a transition thickness of SiO₂ ($t = 1.5$ nm)
- Cannot be understood by spin Hall physics → suggest the role of **interface**

Electric-field control of effective spin Hall sign

GdO _x (3 nm)
Co (0.8 nm)
Pt (1.5 - 2 nm)

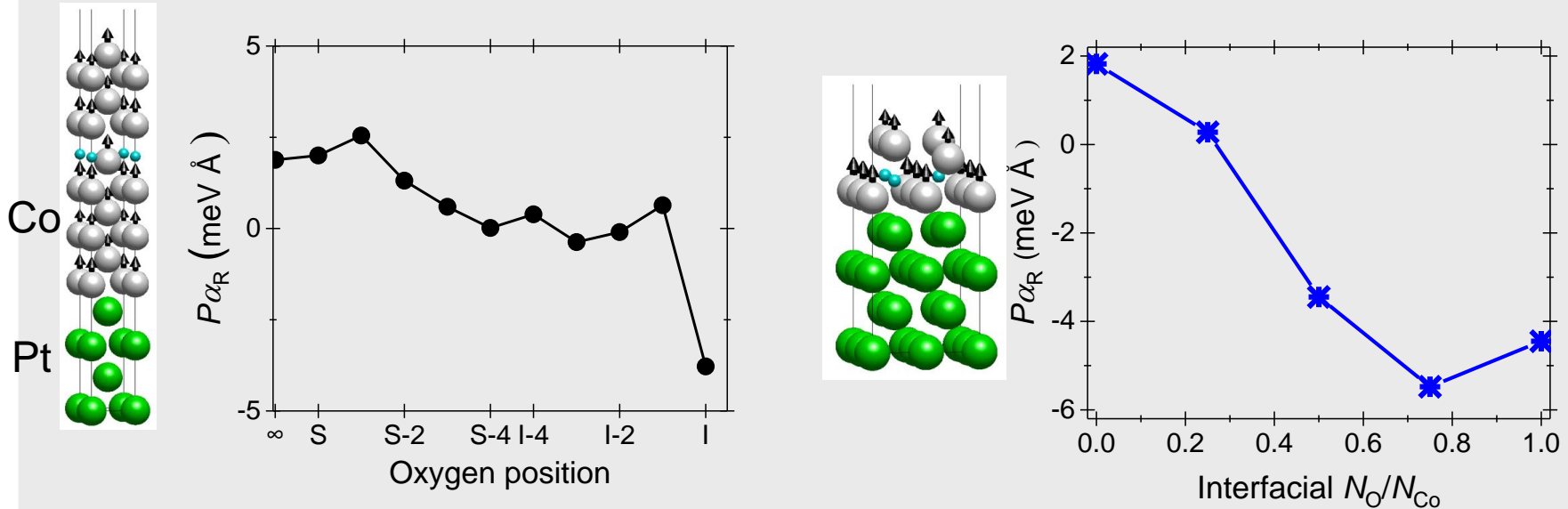


- Competition b/w bulk spin Hall vs. **interfacial SOC** with opposite sign
- Device can switch to either up or down depending on its programmed SOT state → reconfigurable **spin logic**



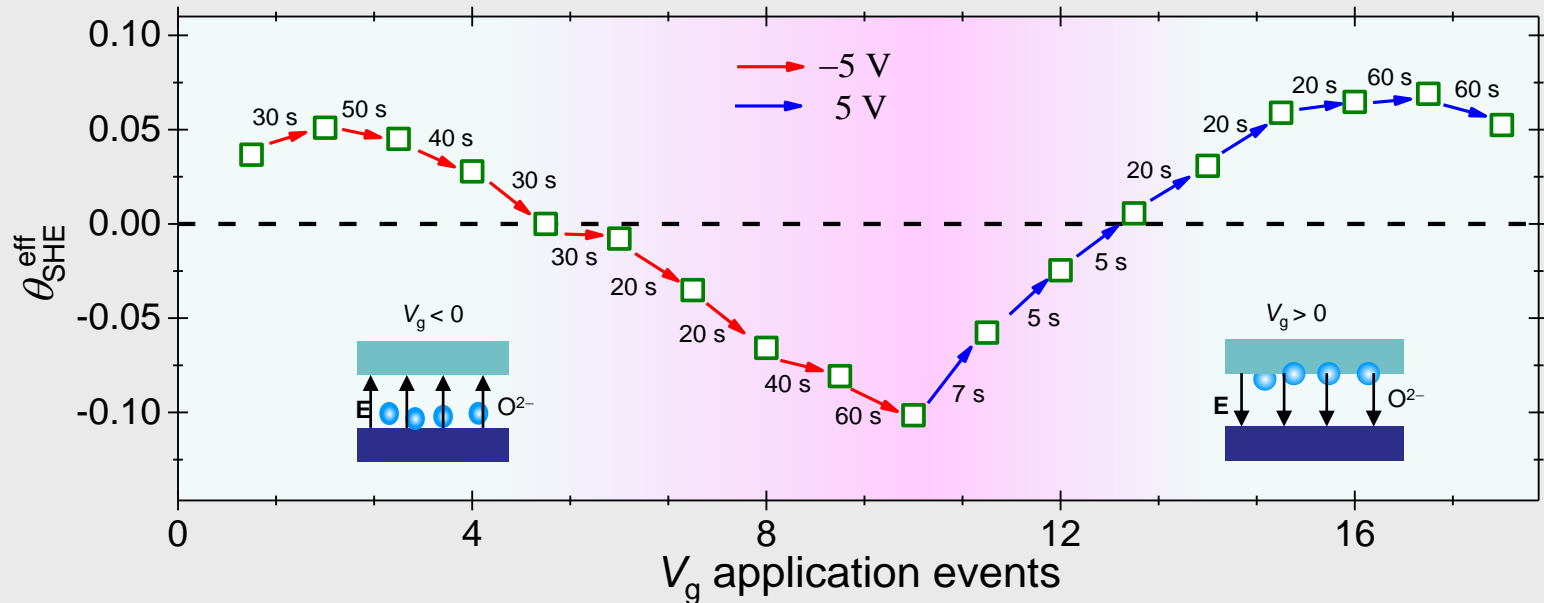
First principle calculations

First principle calculations performed in Prof. Nicholas Kioussis's group, *California State University*.



- The effective Rashba parameter ($P\alpha_R$) changes sign as oxygen is filled at the **Pt/Co interface**.
- The spin accumulation direction changes sign when the negative effective Rashba torque exceeds the spin Hall torque.
- Effective Rashba coefficient changes its sign for **30%** oxidization.

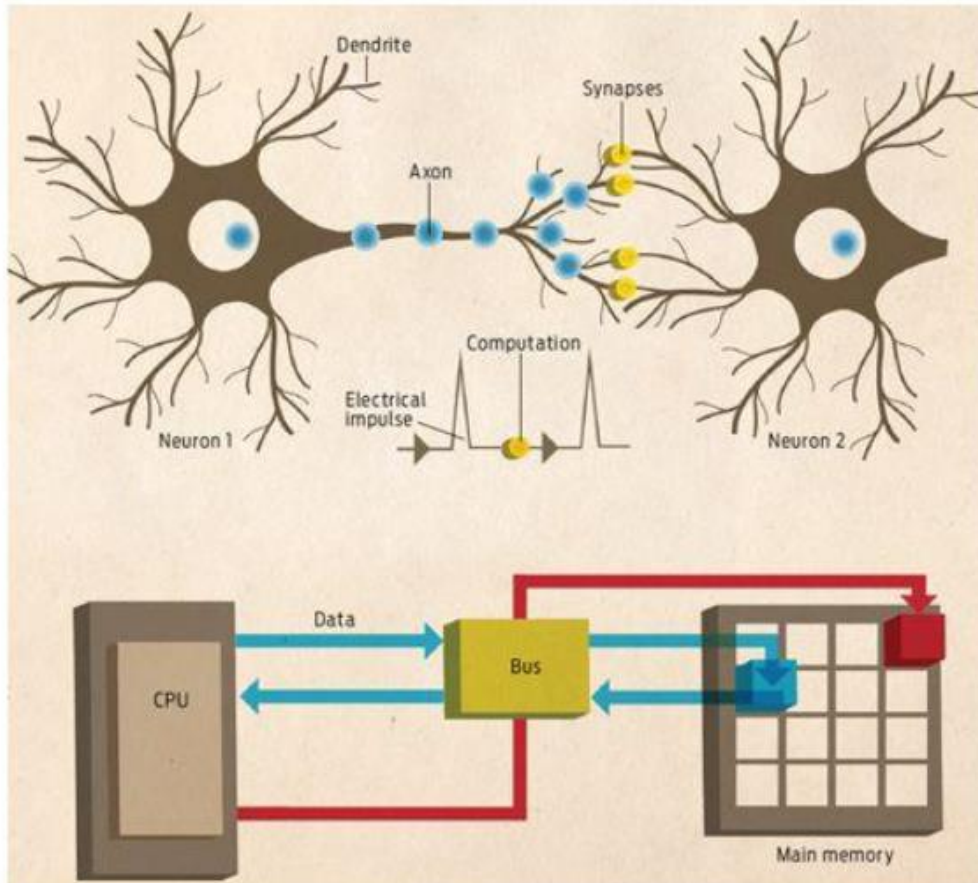
Gradual modulation of SOTs



- Application of V_g pulses on a Pt (1.5 nm)/Co (0.8 nm) device.
- Thin gate oxide results in modulation at room temperature.
- The sign of effective spin Hall angle of the material can be changed using electric-field.
- Gradual modulation of SOTs can be applied to **neuromorphic devices**.

Neuromorphic Computation

Even the simplest brain is superior to a super computer,
the secret: ARCHITECTURE!



Human brain:

- 10^6 neurons / cm^2
- 10^{10} synapses / cm^2
- 2 mW / cm^2

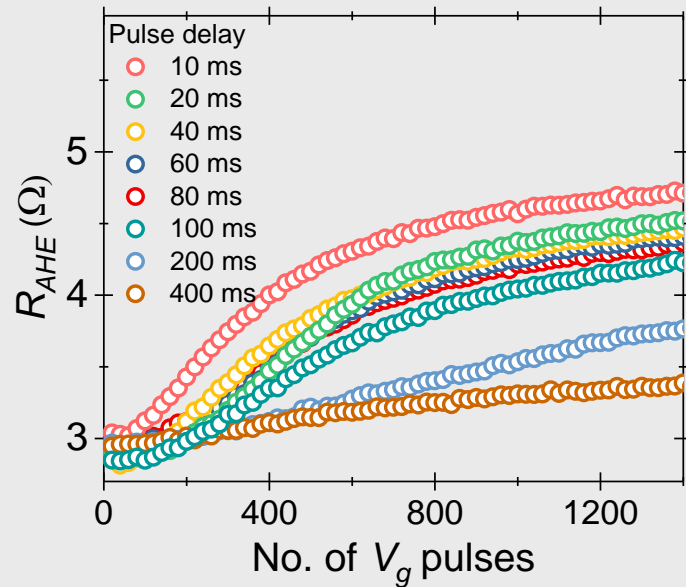
Total power consumption: 20 Watts

Memristors:

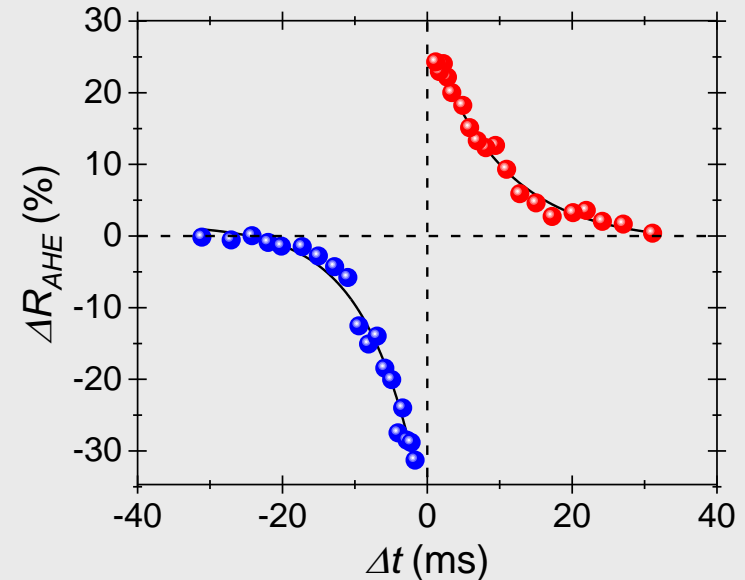
- Cheap
- Power efficient
- Small

Spike temporal dependence

Spike-rate dependence



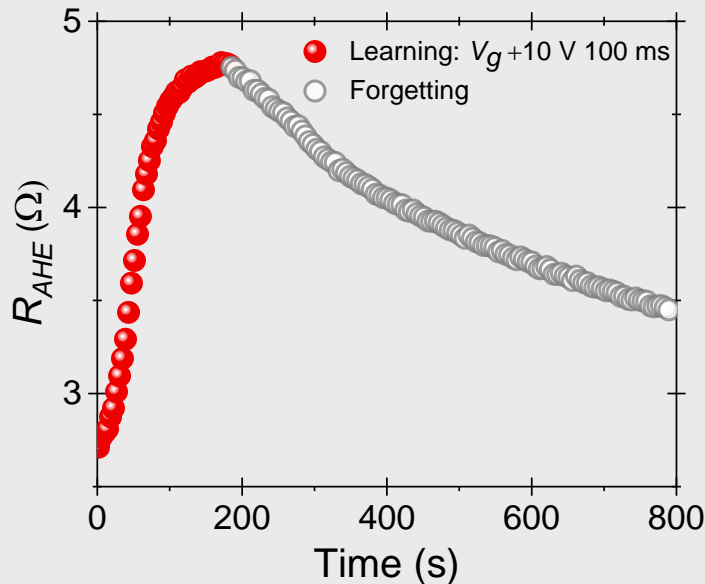
Spike-timing dependence plasticity



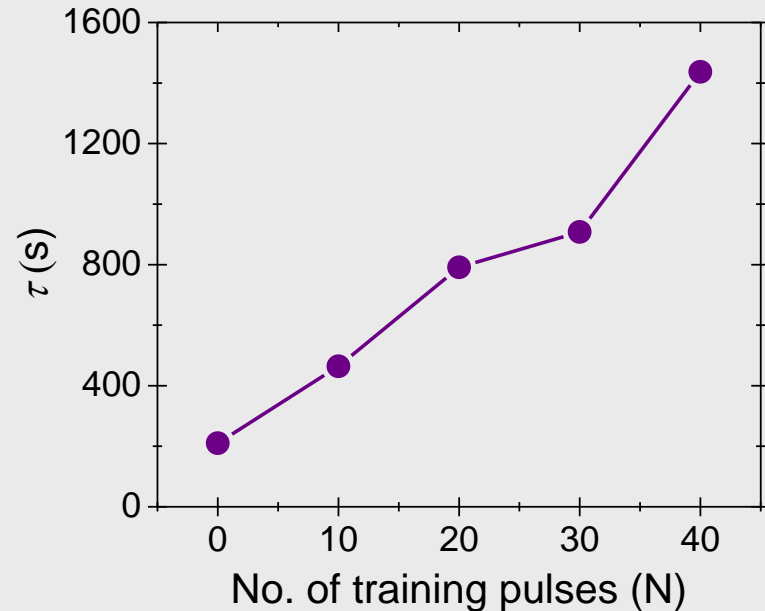
- Frequent stimulation results in larger modulation of weight.
- The magnetic synapse follows the STDP curve.
- Closely spaced pre- and post-synaptic spikes results in larger weight modulation.

Learning, forgetting & training

Learning & forgetting

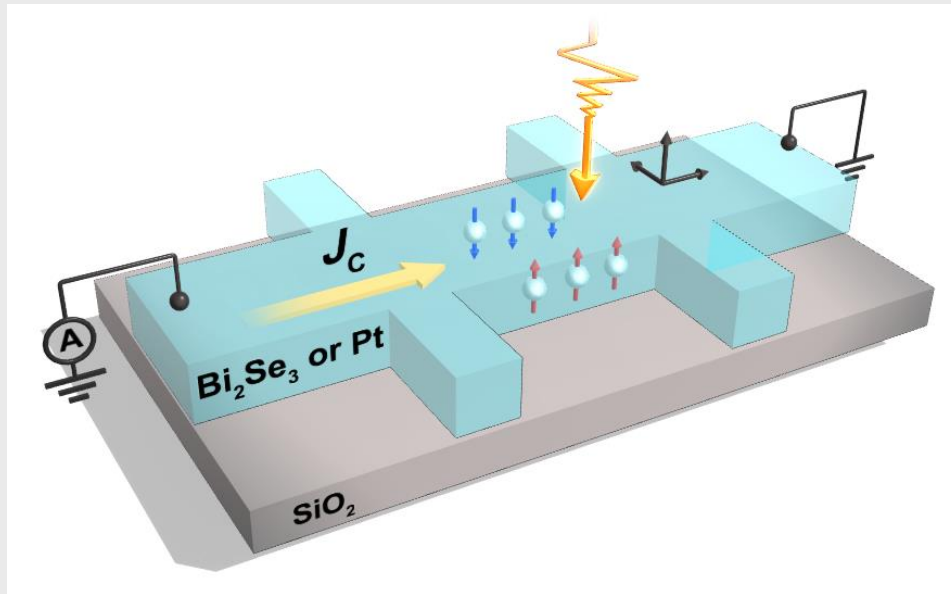


Training



- After removing the stimulation the weight starts decreasing again similar to forgetting behavior of human brain.
- With increasing number of training pulses the forgetting rate decreases.

Scanning photovoltage microscope with currents



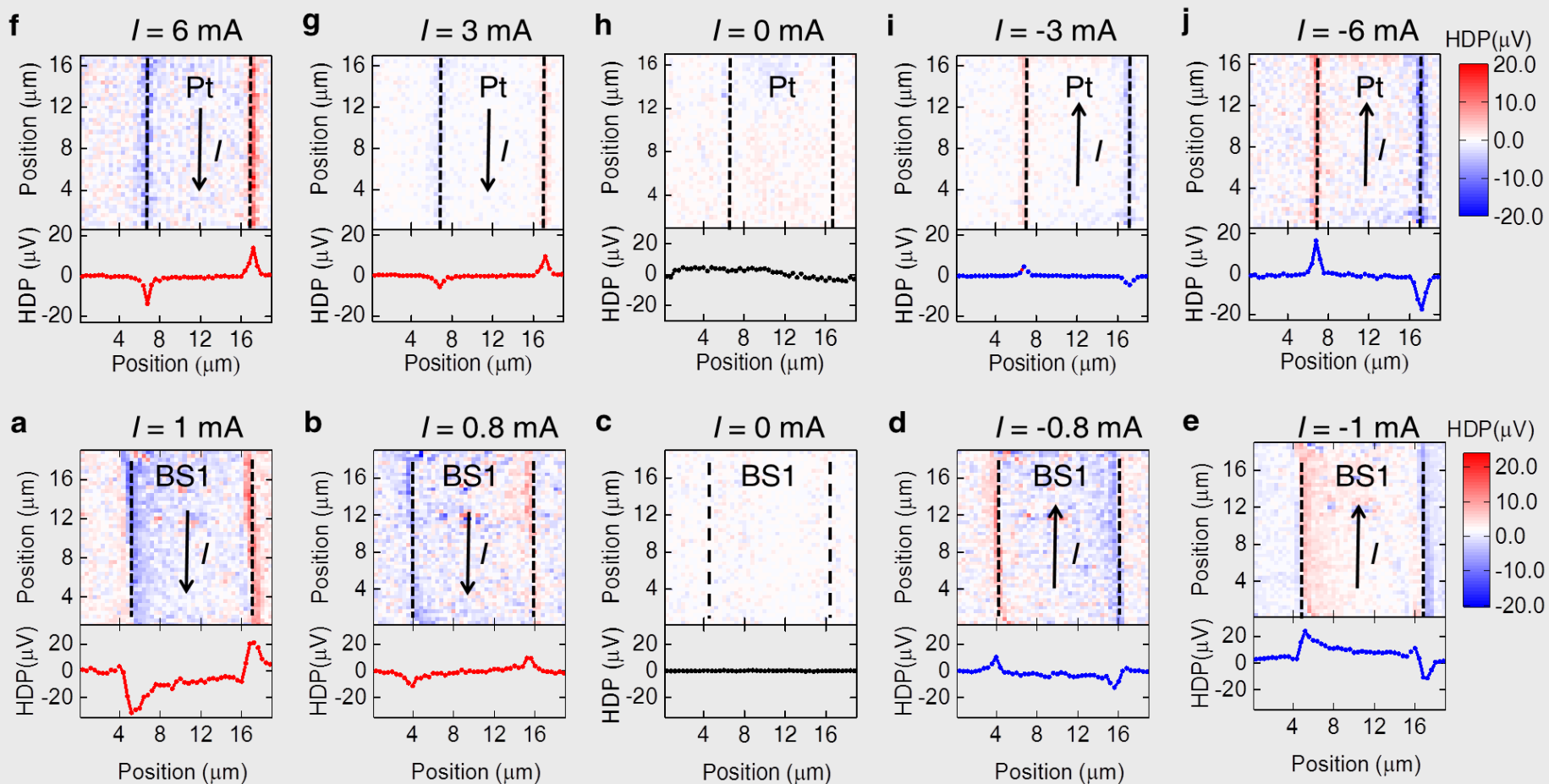
- ❖ DC current is applied to induce spin accumulation.
- ❖ Circularly polarized light normally incidents on the sample.
- ❖ Magnetic circular dichroism.
- ❖ Photovoltages are detected by lock-in amp.
- ❖ Piezo sample stage enables mapping.

$$\text{❖ } V_{\text{photovoltage}} = V_{RCP} - V_{LCP}$$

RCP light excites **spin up electron**, while **LCP light** excites **spin down electron**.

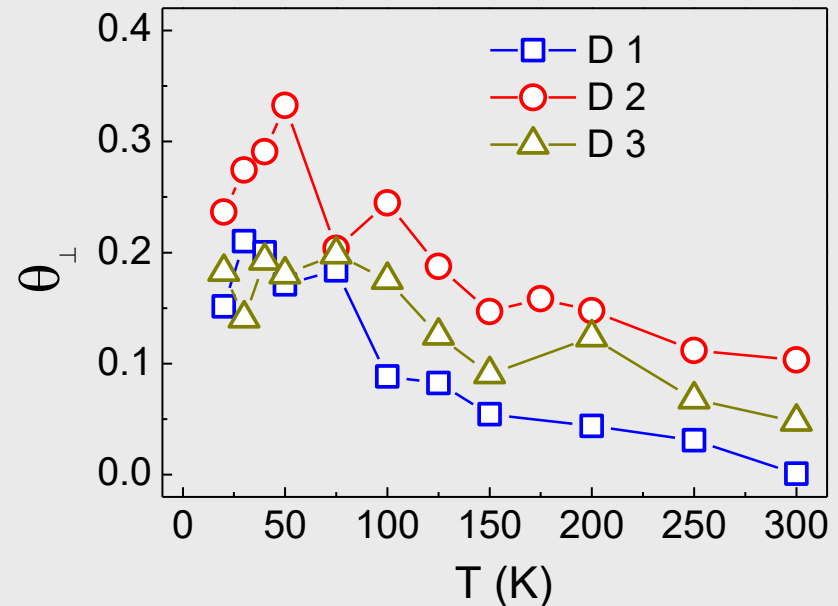
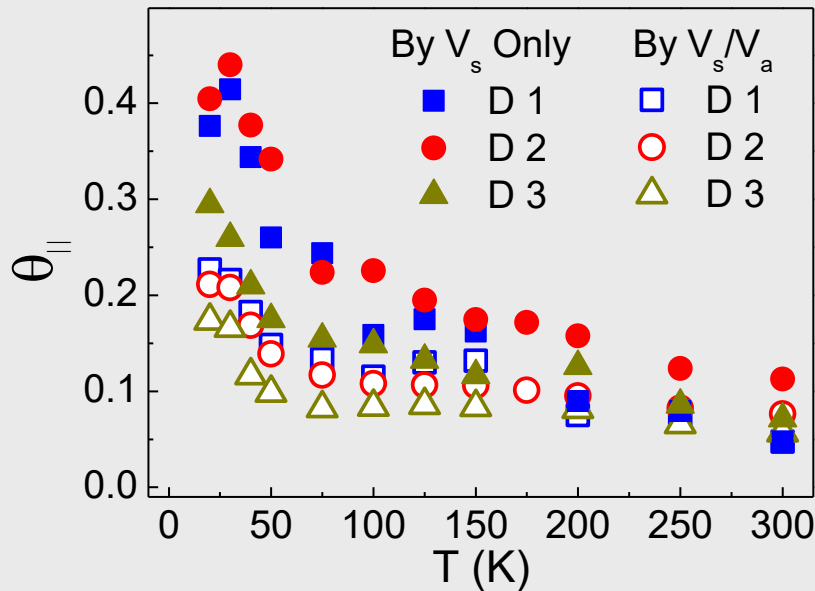
- $V_{\text{photovoltage}} > 0 \rightarrow$ local spin direction is **spin down**.
- $V_{\text{photovoltage}} < 0 \rightarrow$ local spin direction is **spin up**.

Accumulated spin imaging in Pt and Bi_2Se_3



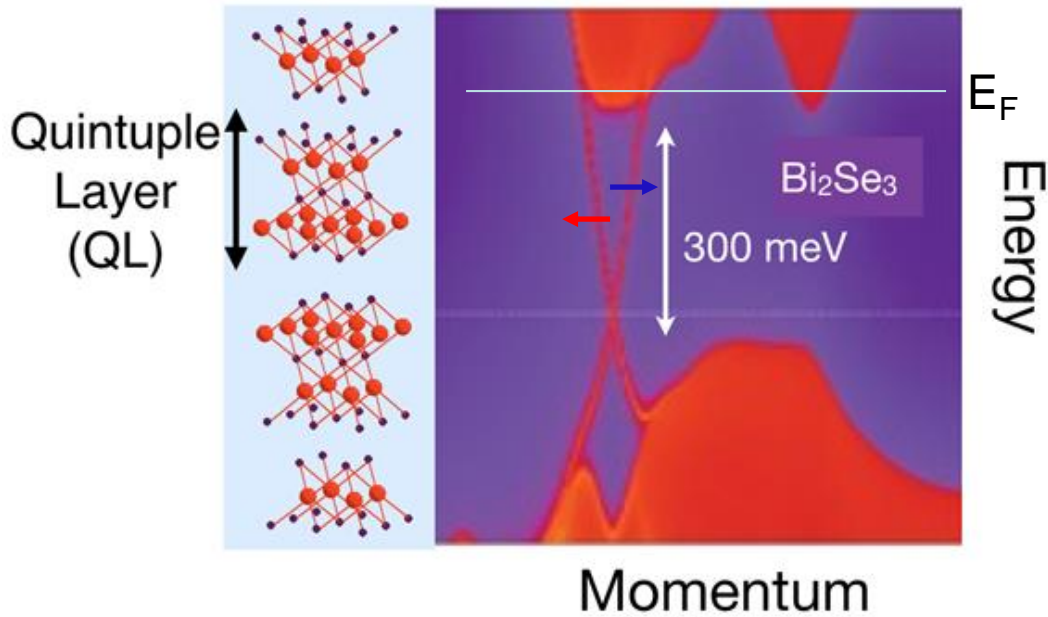
- ❖ Sign switches in opposite edges and with reversing currents.
- ❖ Both semiconductors and metals work.
- ❖ Can extract spin Hall angle and spin lifetime without a ferromagnet.

In-plane and out-of-plane SOT efficiency in Bi_2Se_3



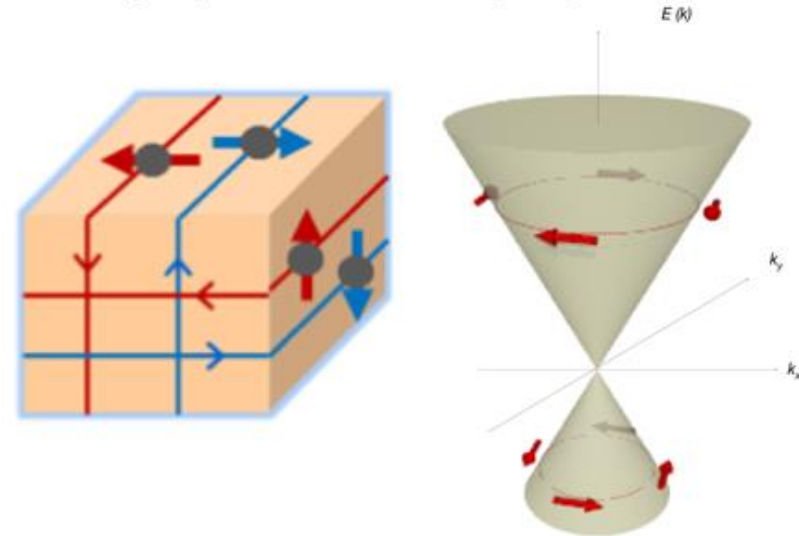
- θ_{\parallel} increase by 10 times at low temperature (upto $\sim \mathbf{0.42}$).
- θ_{\perp} has the same order of magnitude compared to θ_{\parallel} .
- Origin of the observed spin-orbit torques is topological surface states in Bi_2Se_3 .
- Hexagonal warping in the TSS of Bi_2Se_3 can account for θ_{\perp} .

3D topological insulators (TIs)



H. Zhang, C.-X. Liu, *et al.*, *Nature Physics* **5**, 438 (2009)

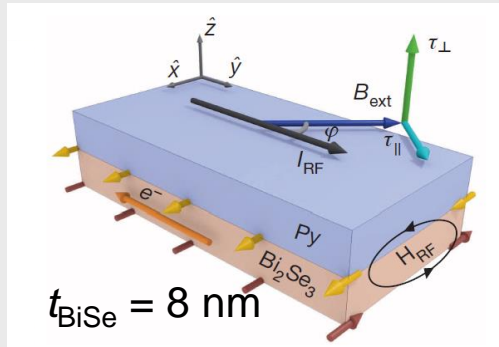
Fu & Kane, *Phys. Rev. B* **76**, 045302 (2007)
Moore & Balents, *Phys. Rev. B* **75**, 121306(R) (2007)
Roy, *Phys. Rev. B* **79**, 195321 (2009)



- ❑ Spin polarized surface currents
- ❑ Spin-momentum locking \rightarrow giant spin Hall angle?

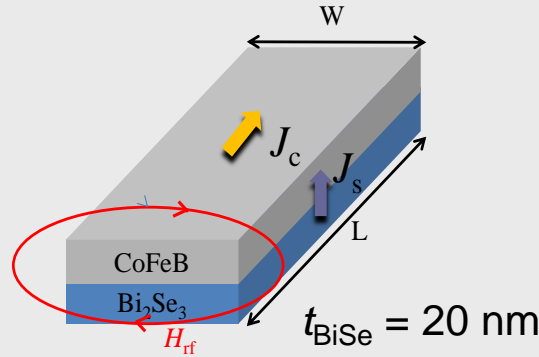
Exotic spin Hall angles from topological insulators

spin Hall angle (θ_{SH}) = 2~3.5
ST-FMR (Cornell)



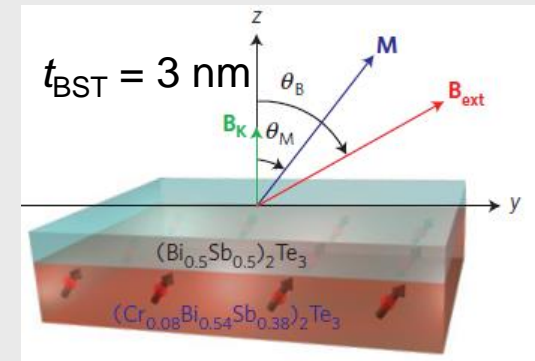
Nature **511**, 449 (2014)

$\theta_{SH} = 0.42$ (low temp)
ST-FMR (NUS)



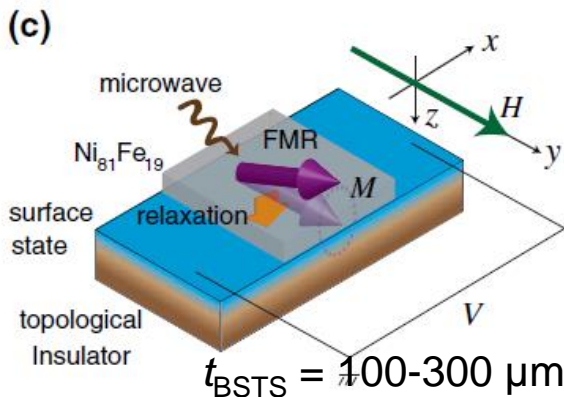
PRL **114**, 257202 (2015)

$\theta_{SH} = 140-425$ (low temp)
spin-orbit switching (UCLA)



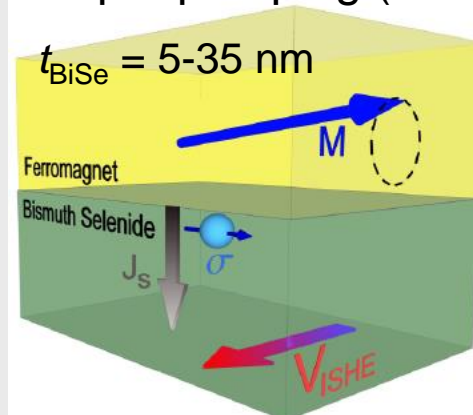
Nat. Mater. **13**, 699 (2014)

$\theta_{SH} \sim 1 \times 10^{-4}$
Spin-pumping (Tohoku)



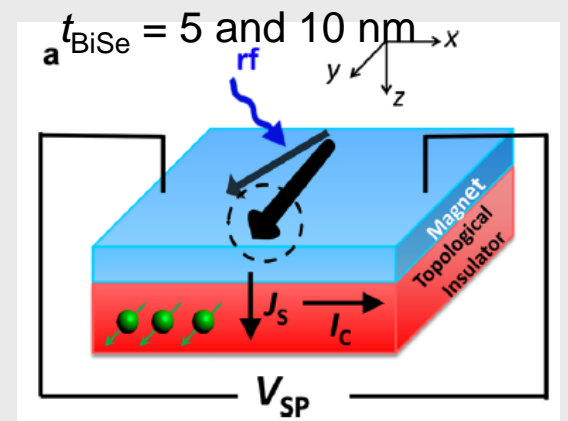
PRL **113**, 196601 (2014)

$\theta_{SH} = 0.01$
Spin-pumping (NUS)

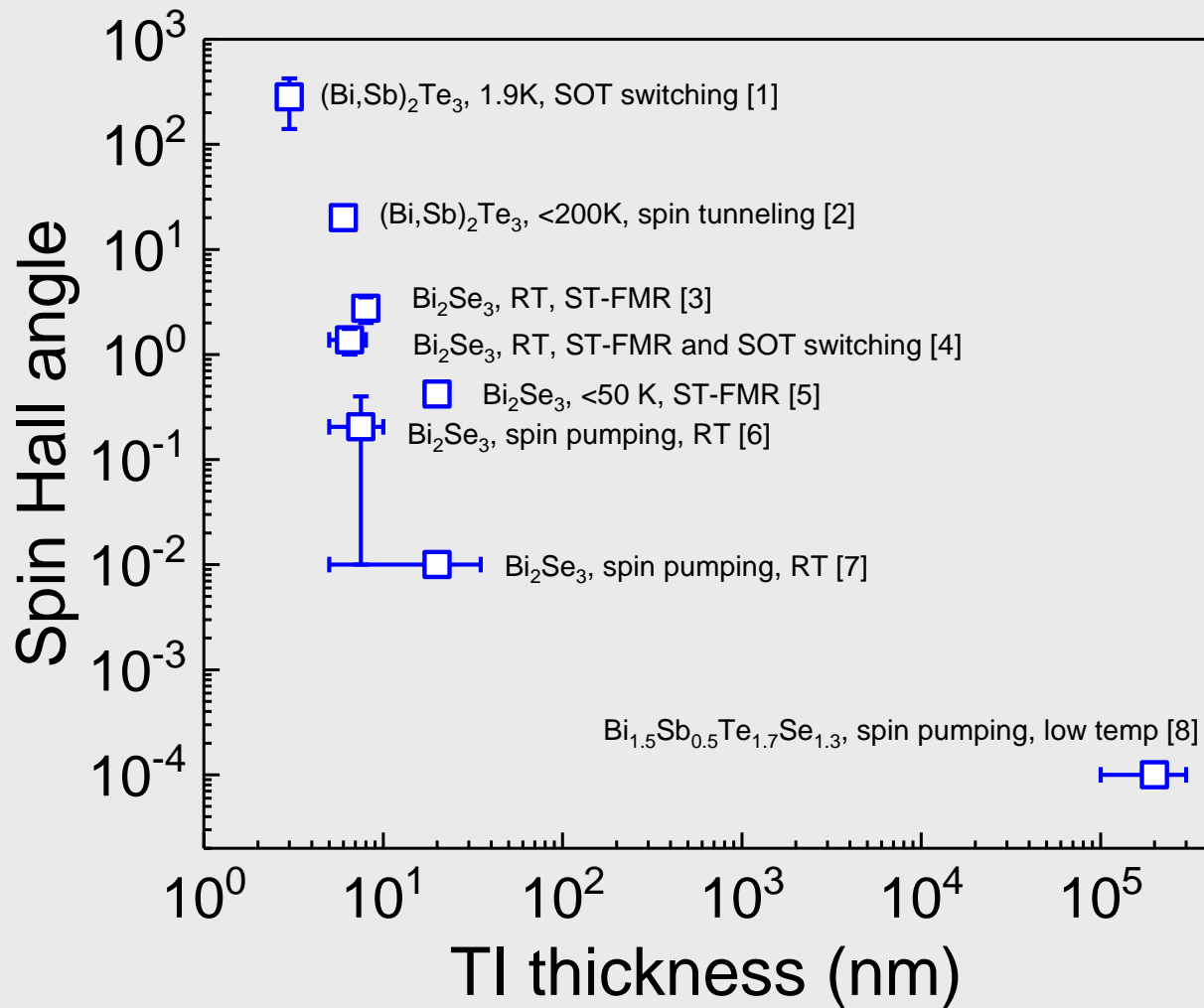


PRB **90**, 094403 (2014)

$\theta_{SH} = 0.01-0.4$
Spin-pumping (Minnesota)



Nano Lett **15**, 7126 (2015) 29



- TI is not a good insulator.
- E_F is pinned in conduction band.
- Current shunting thru bulk reduces the efficiency!

[1] Nat. Mater. **13**, 699 (2014)

[2] PRB **91** 235437 (2015)

[3] Nature **511**, 449 (2014)

[4] Nat. Commun. **8** 1364 (2017)

[5] PRL **114**, 257202 (2015)

[6] PRB **90**, 094403 (2014)

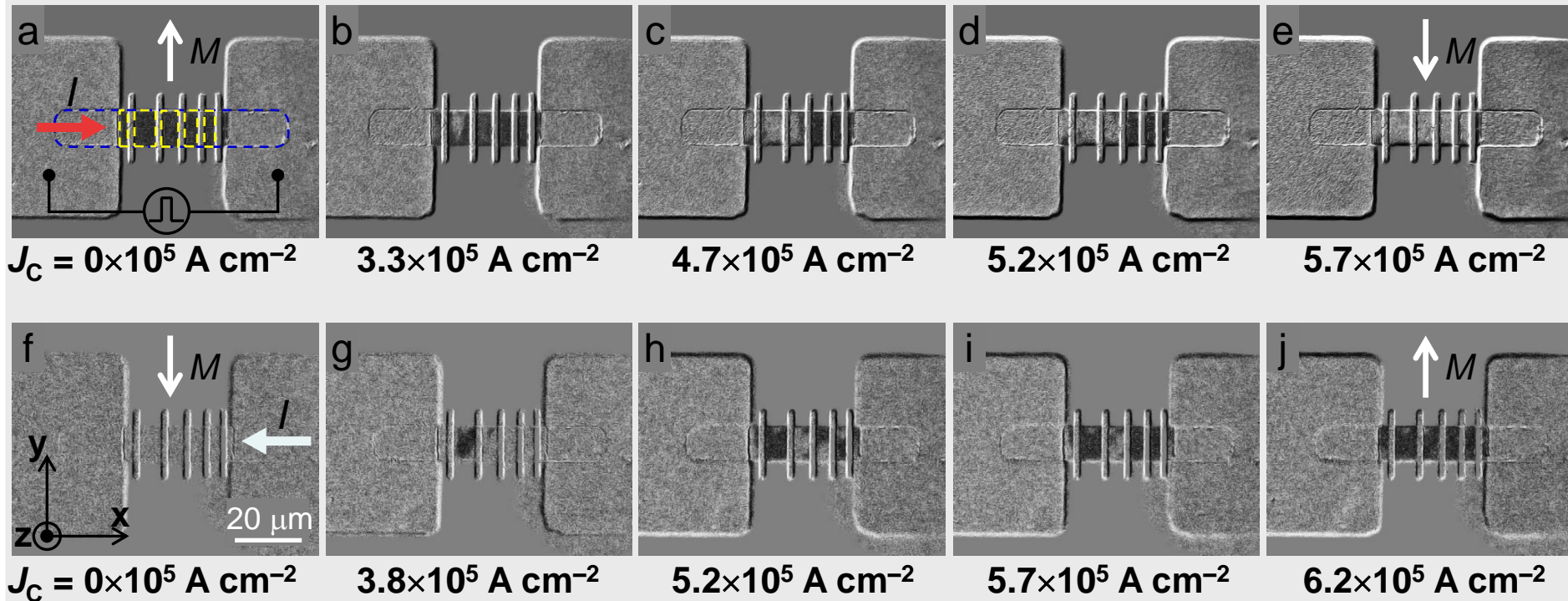
[7] Nano Lett **15**, 7126 (2015)

[8] PRL **113**, 196601 (2014)



Current induced magnetization switching in Bi₂Se₃/Py

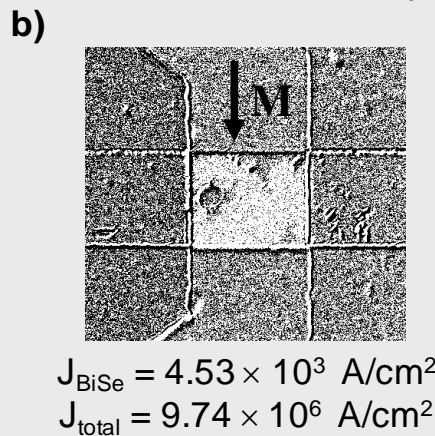
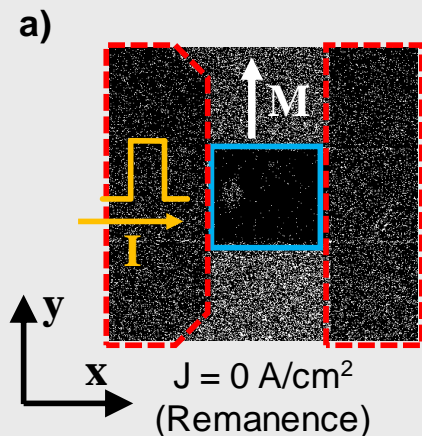
Bi₂Se₃ (8 QL)/NiFe (6 nm)/MgO (1 nm)/SiO₂ (4 nm)



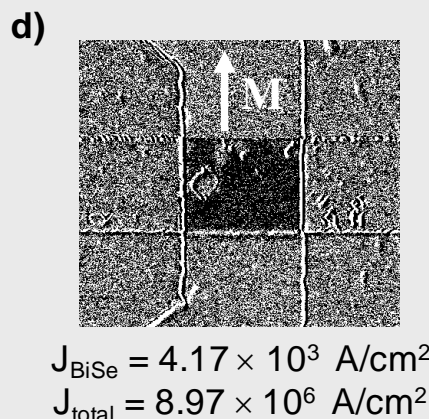
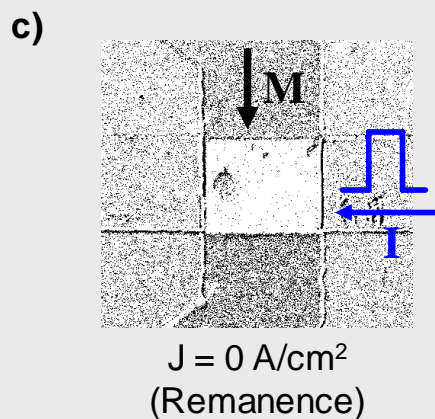
- TI magnetization switching reported in a Cr doped TI at 1.9 K with an external magnetic field [Nat. Mater. 13, 699 (2014)].
- Demonstrated magnetization switching of Bi₂Se₃/NiFe at room temp with a low critical current density ($J_C \sim 6 \times 10^5$ A/cm²) and without a magnetic field.
- A giant $\theta_{TI} = 1.75$.

Sputtered Bi_2Se_3

Si/SiO₂ sub/Bi₂Se₃ (t)/Py (6 nm)/SiO₂ (5 nm)



$$\theta_{\text{TI}} = 45 \text{ for } 10 \text{ nm}$$

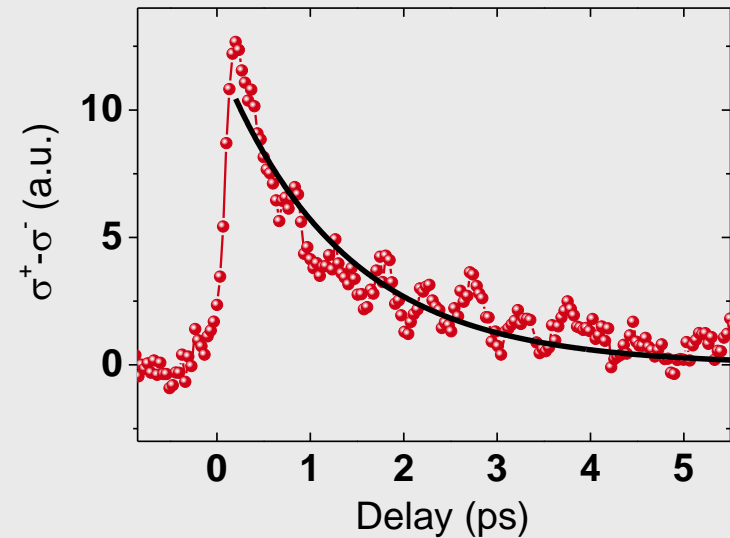
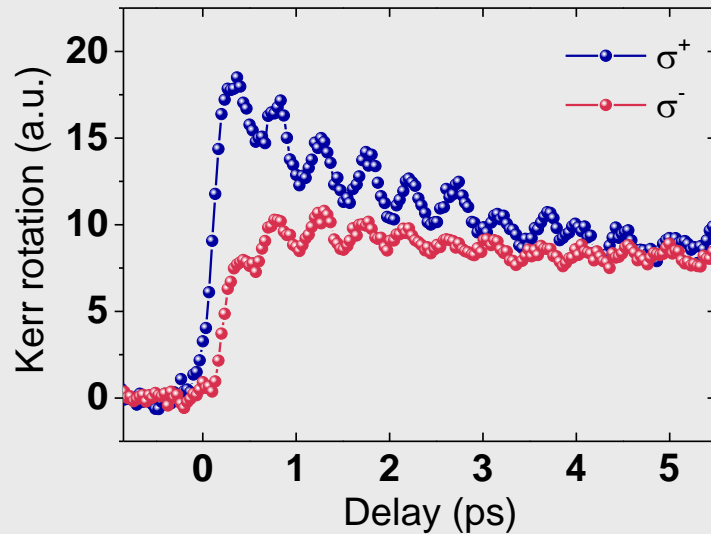


- Extremely large spin Hall angle, but large power consumption due to large resistivity
- J.P. Wang (Minnesota) also reported sputtered Bi_2Se_3 (Nat. Mater. 17, 800 (2018))
 $\theta = 18.6$, $J_c = 4.3 \times 10^5 \text{ A cm}^{-2}$

J. Phys. D: Appl. Phys. **52**, 224001 (2019)

Spin relaxation time in Bi_2Se_3

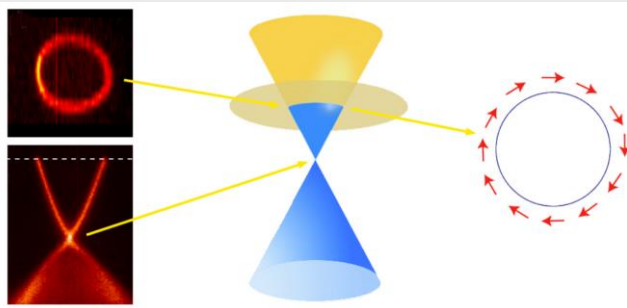
Time-resolved magneto optical Kerr effect



- Signal sensitive to bulk due to large penetration depth of light
- Oscillation frequency is 2.13 THz from coherent vibrations of the A_{1g} longitudinal optical phonons of Bi_2Se_3
- Exponentially decay with a characteristic time of **1.3 ps**

Weyl semimetal background

Topological insulator (TI):

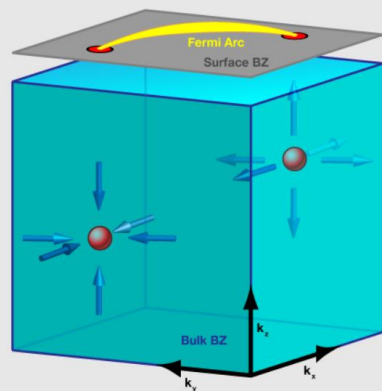


Conductive surface: spin-momentum locked surface states ($P \sim 20\text{-}50\%$)

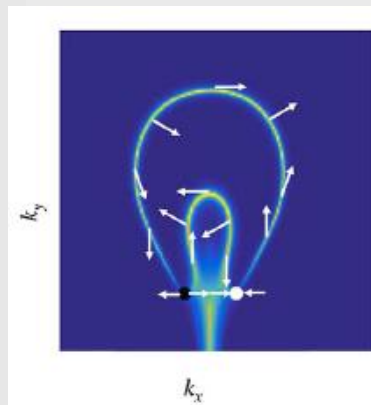
S. Xu, *et al.* Nat. Commun. **6**, 6870 (2015)

Td-WTe₂ : Weyl semimetal, strong Edelstein effect, good conductivity, 2D layered TMD material, less roughness than MBE grown TI such as Bi₂Se₃, etc.

Weyl semimetal:

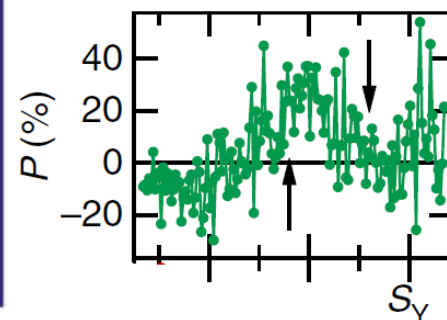


Conductive surface: SML Fermi arc states ($P \sim 80\%$)



Conductive bulk: strong spin orbit coupling ($P \sim 40\%$)

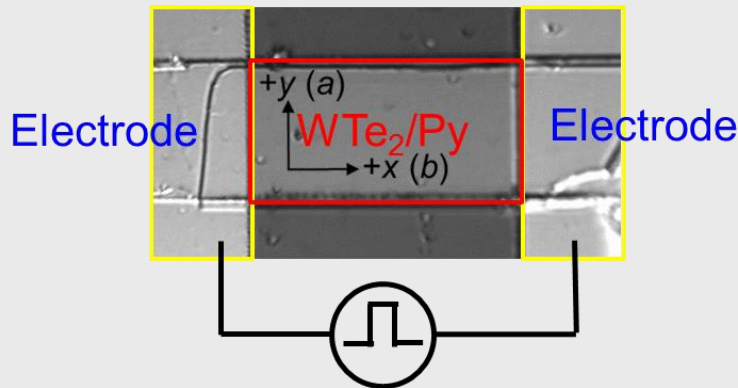
S. Xu, *et al.* Phys. Rev. Lett. **116**, 096801 (2016)
P. K. Das, *et al.* Nat. Commun. **7**, 10847 (2016)



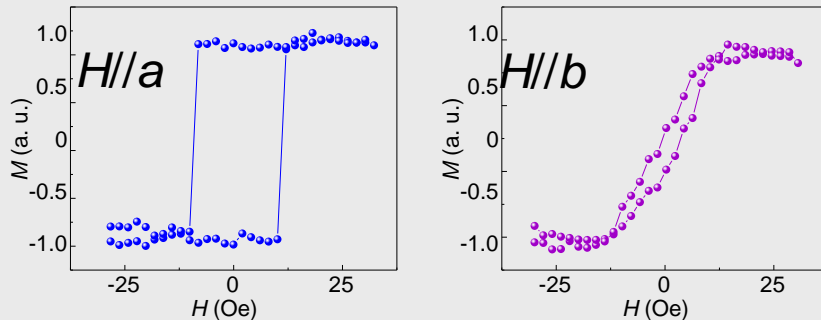
Ideal material to achieve the spin orbit torque driven magnetization switching.

SOT driven magnetization switching-*I* along *b*-axis

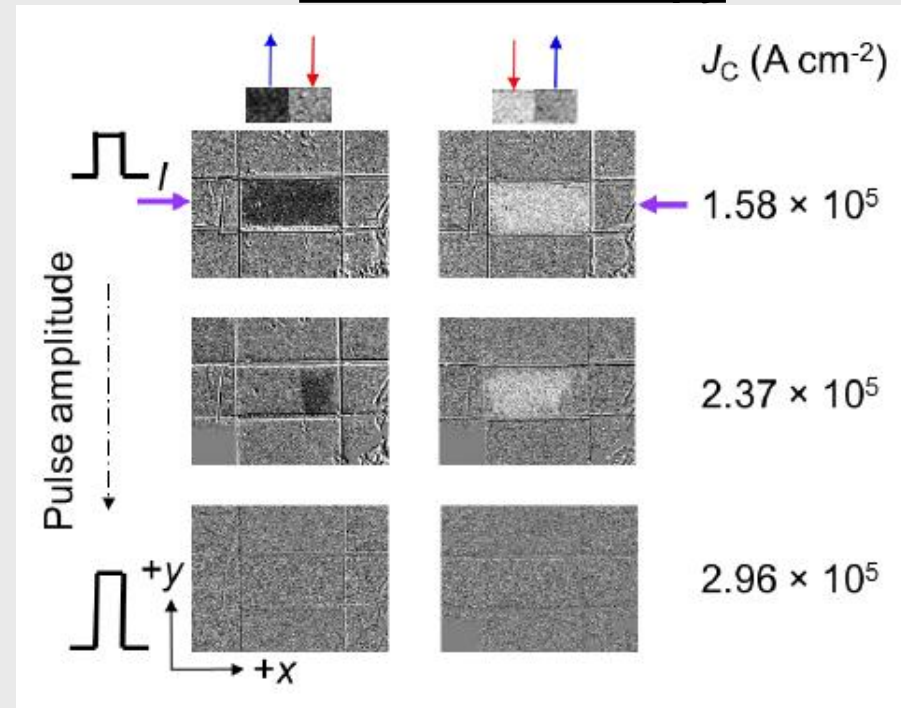
- Charge current is applied in the *b*-axis.
- Spin current is generated along the *a*-axis.



Hysteresis loops of the device

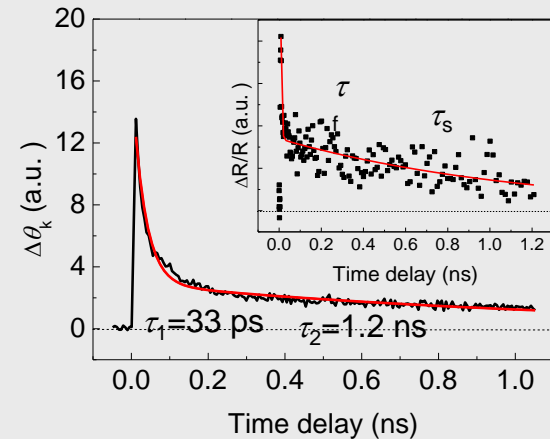
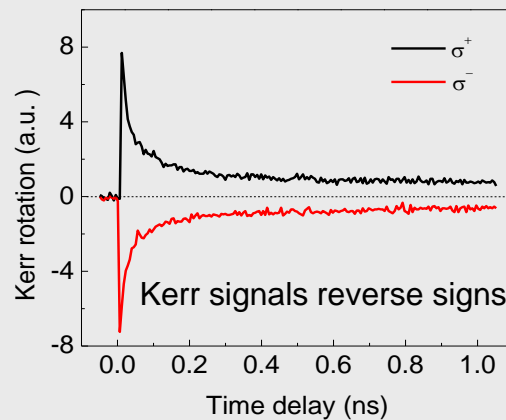
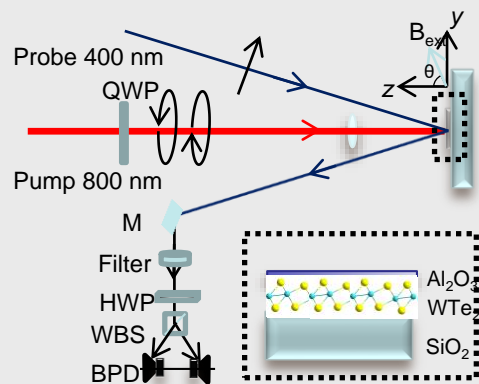
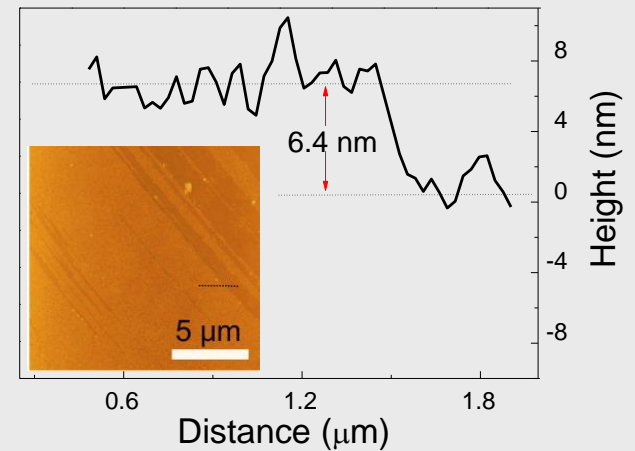
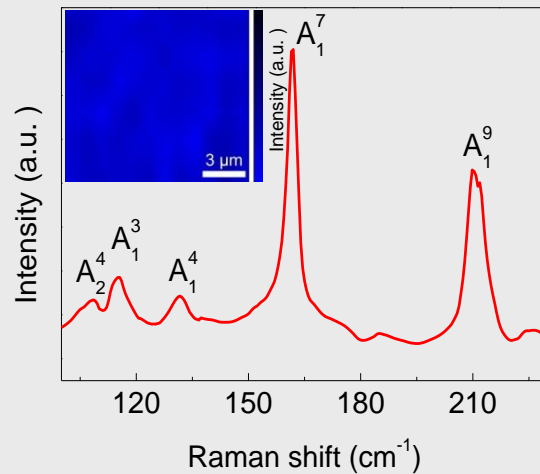
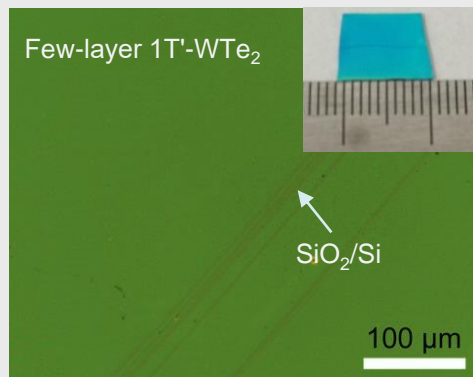


MOKE microscopy



- Switching starts to happen at 1.58×10^5 A/cm²
- Domain wall moves in the direction of J_C

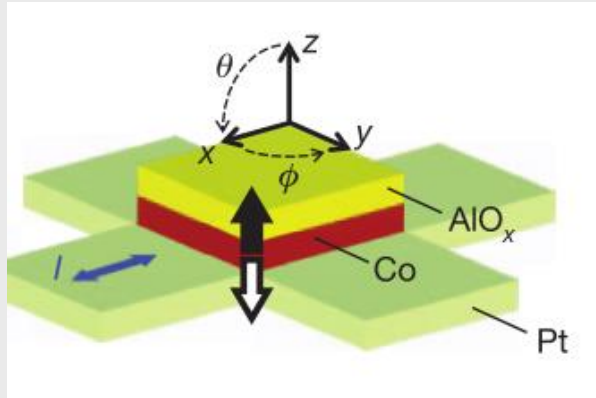
Room-temperature nanosecond spin relaxation



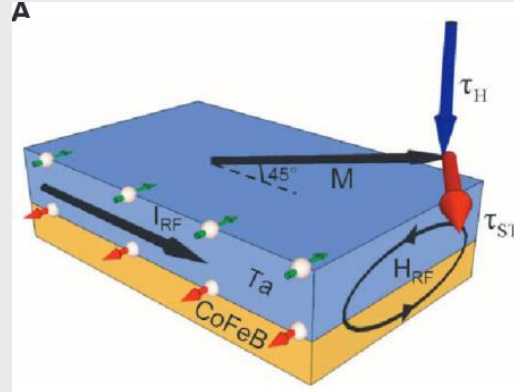
- Centimeter-scale, chemical vapor deposition (CVD)-grown few-layer 1T'-W(Mo)Te₂
- Room-temperature **1.2 ns** spin relaxation

Spin orbit torque in various magnets

Ferromagnet

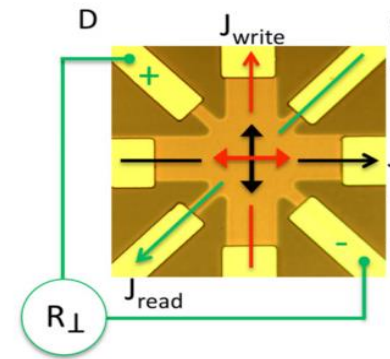
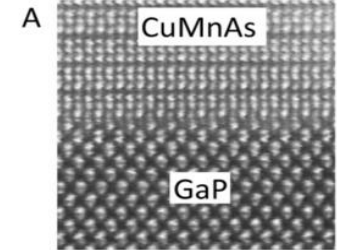


Nature **476**, 189 (2011)



Science **336**, 555 (2012)

Antiferromagnet



Science **351**, 587 (2016)

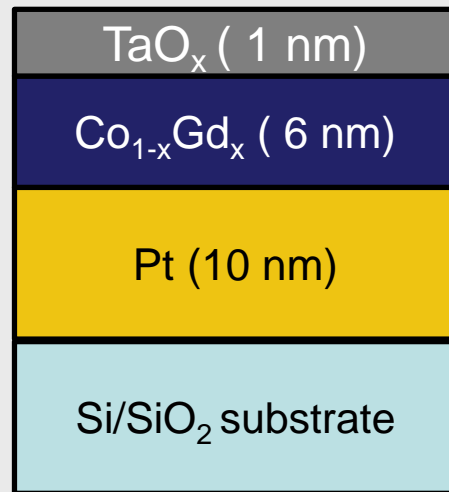
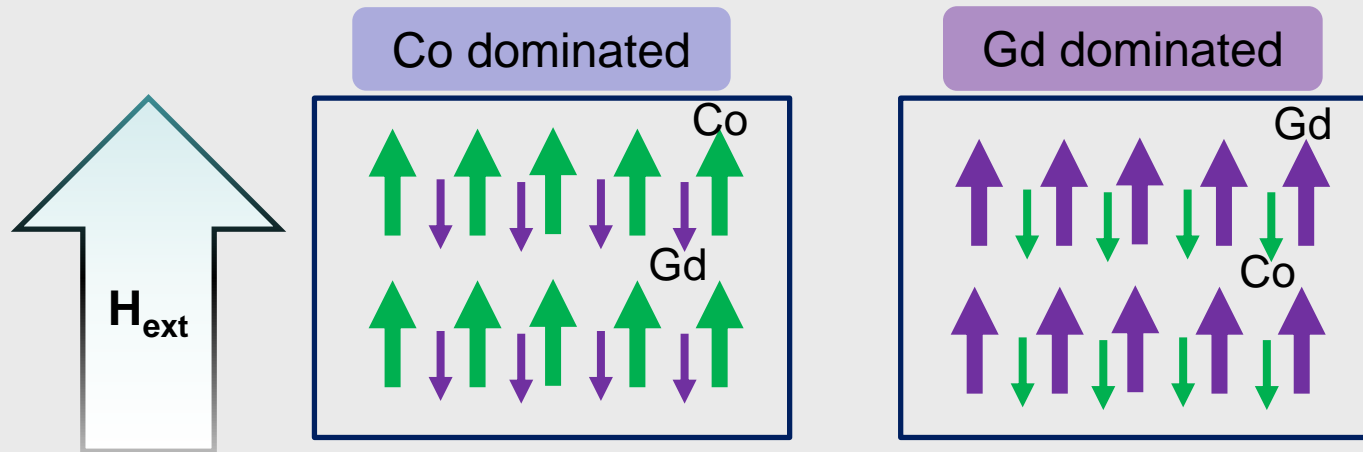
- $1/t_{\text{FM}}$ dependence of spin torque (surface torque)
- Limitation on FM thickness
- Thermal stability issue

- 90 degree rotation
- Not compatible with MTJs
- Small MR → slow reading

How about a **multilayer** or **ferrimagnet (FIM)**?

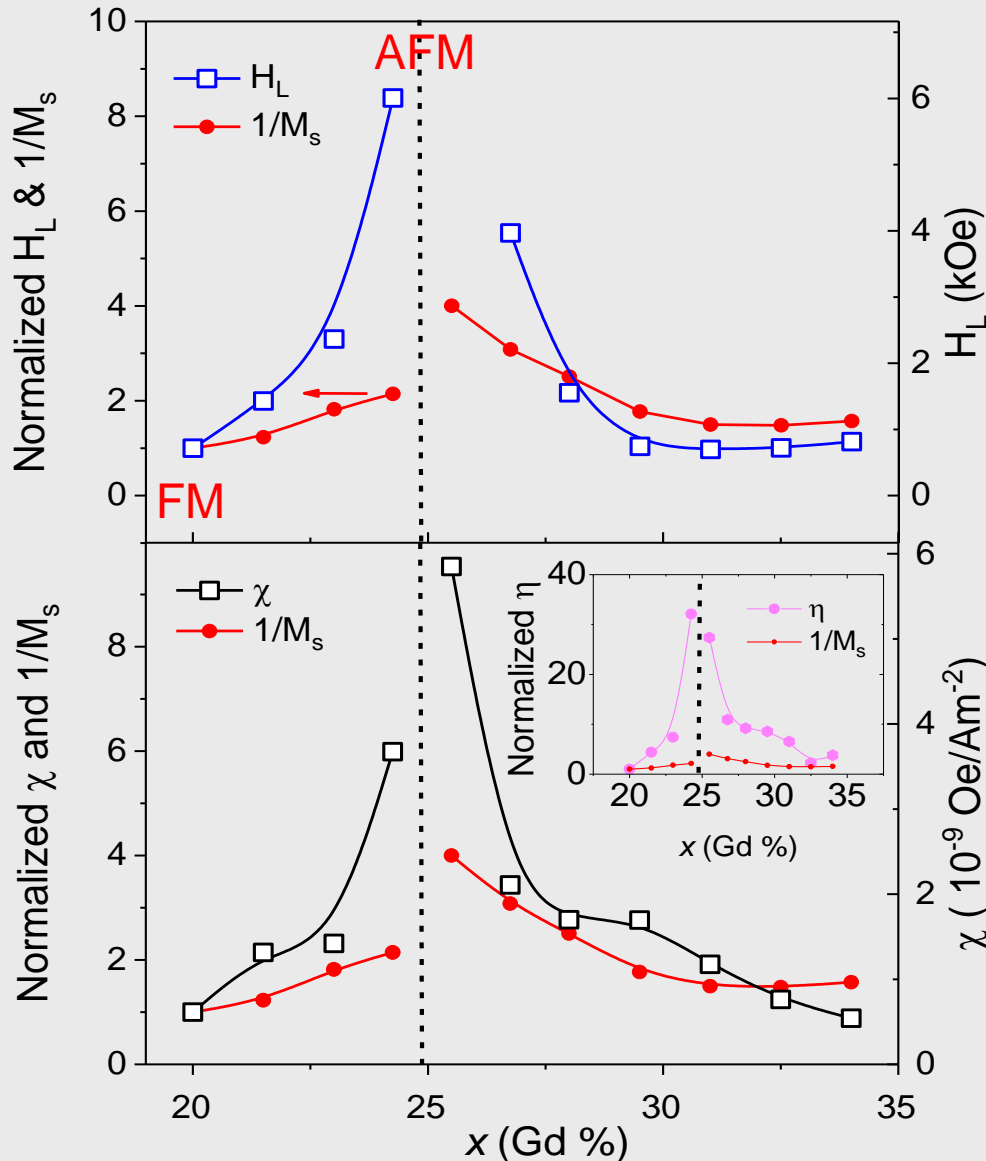
- Easy perpendicular anisotropy, enough magnetic volume

SOT in ferrimagnet : CoGd

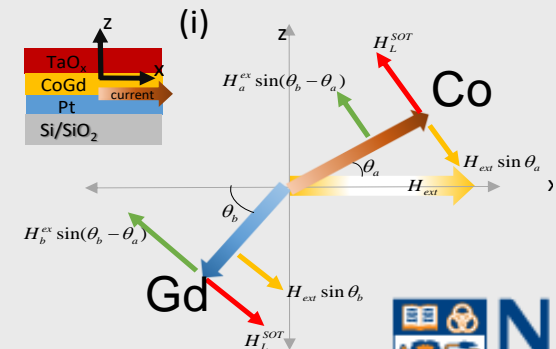


- ❖ Films deposited with varying Co and Gd compositions.
- ❖ CoGd has bulk PMA.
- ❖ Thermally stable thick magnetic layer is grown.

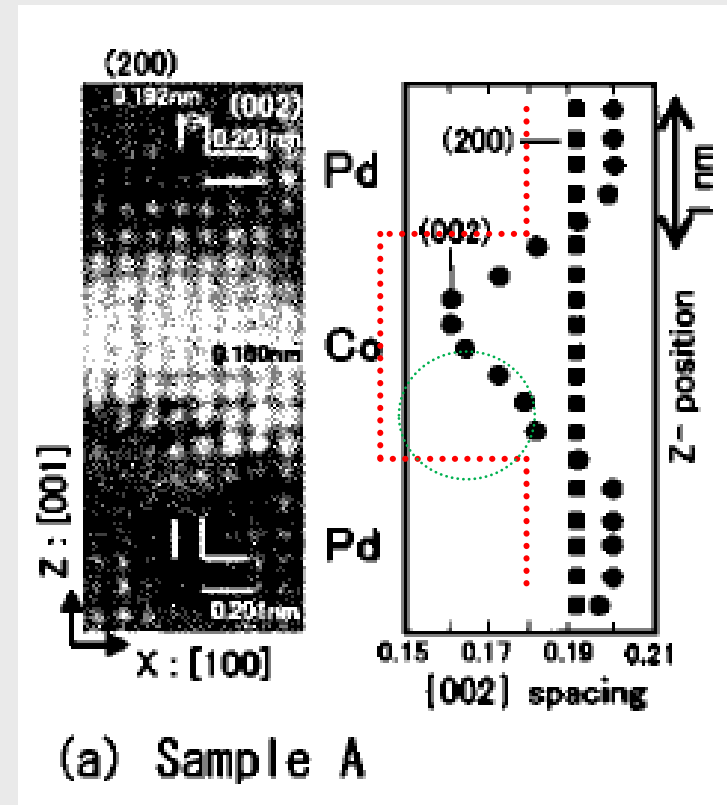
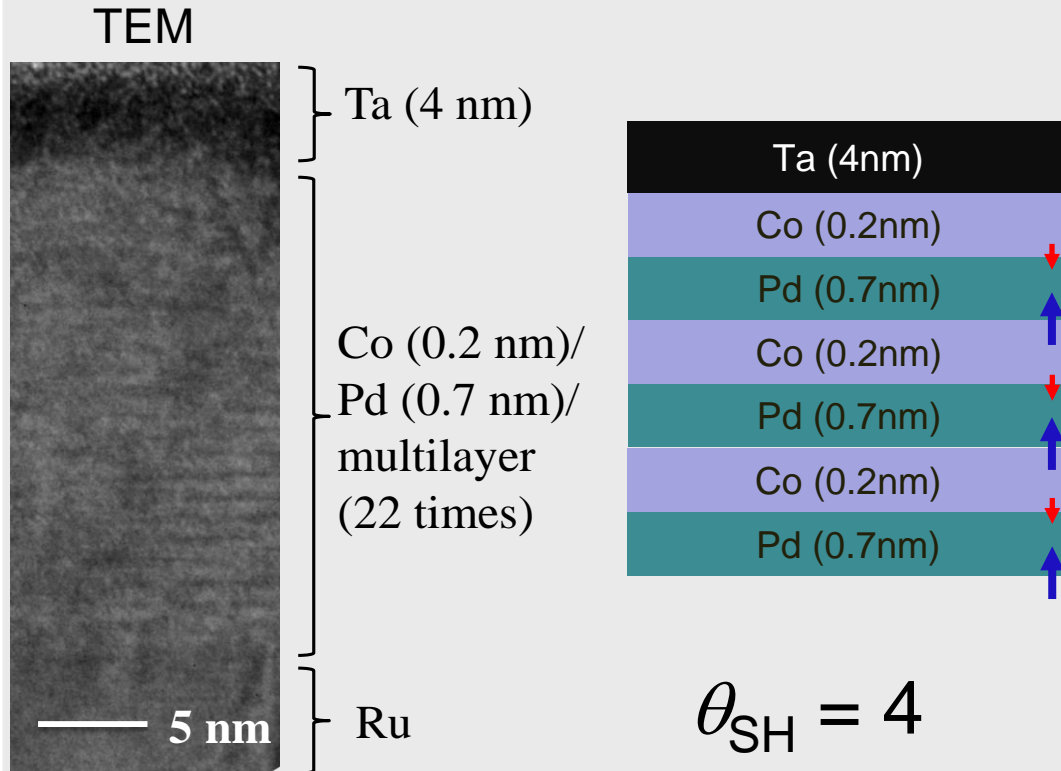
Anomalous scaling of SOT in ferrimagnets



- ❖ Switching efficiency ($\chi = H_p / J_s$), H_L and $1/M_s$ are normalized with respect to their respective values for $\text{Co}_{80}\text{Gd}_{20}$.
- ❖ Exceptional and disproportionate (to $1/M_s$) change of η and H_L cannot be explained by existing SOT understanding.
- ❖ **10 times** increase in θ_{SH} due to **negative** exchange coupling



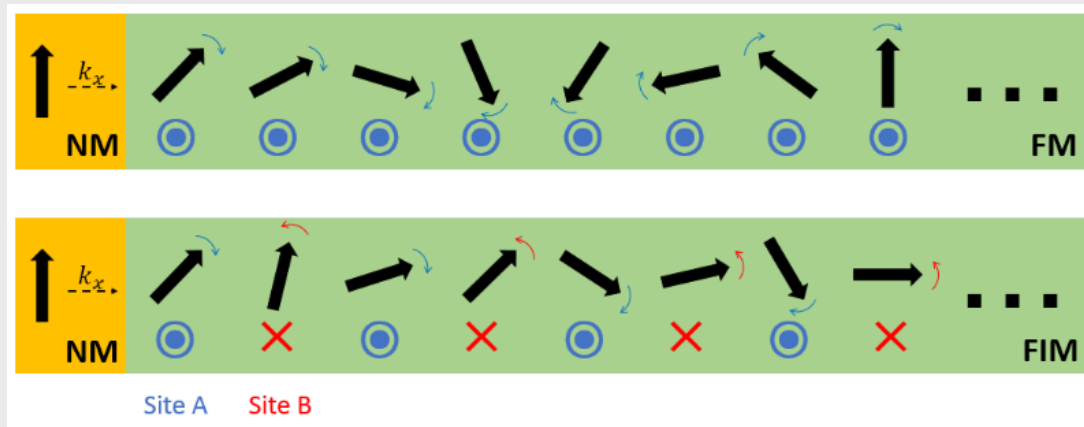
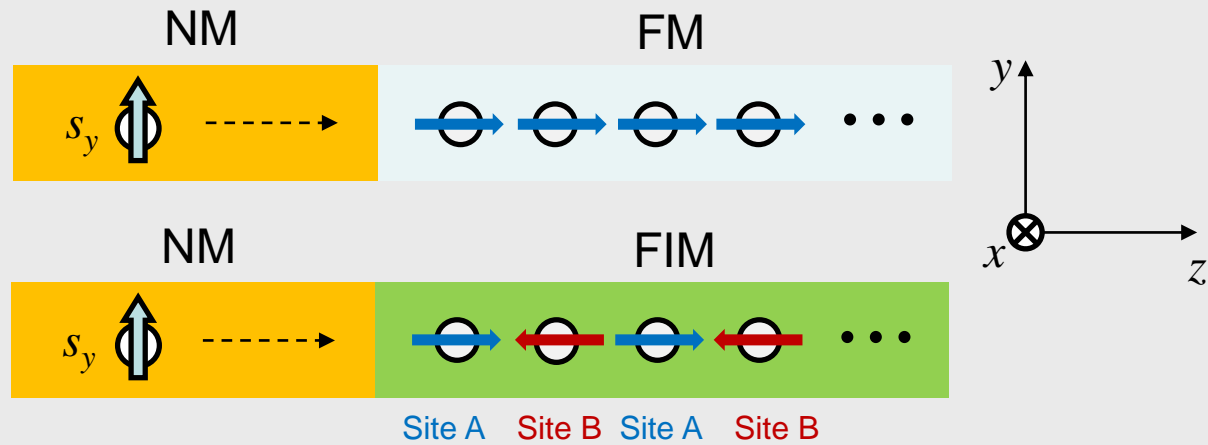
Multilayer: structural asymmetry can be added up



Maesaka, IEEE Trans. Magn. 38, 2676 (2002)
Kim, PRB 53, 11114 (1996)

- Two successive Co/Pd and Pd/Co interfaces are structurally dissimilar.
- Lattice mismatch (9%) between Pd and Co → **Strain engineering**
- This distortion is 30% stronger for Co/Pd than for Pd/Co interfaces

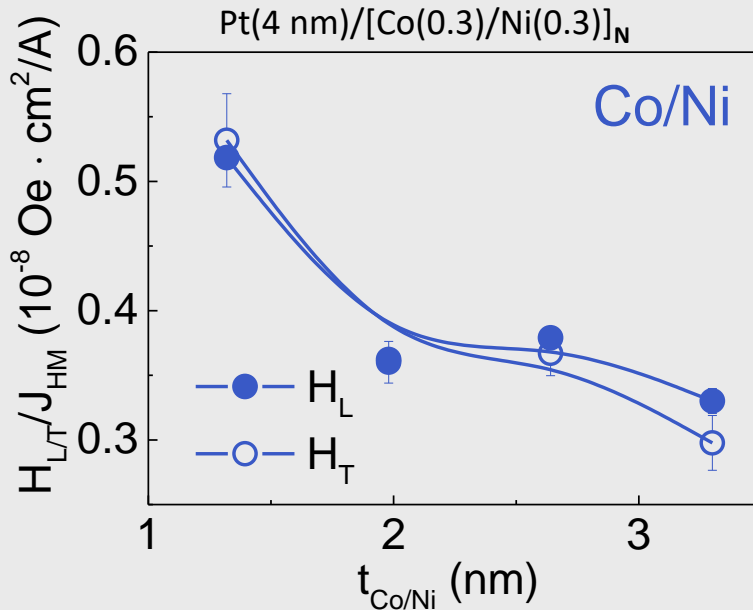
Ferromagnet vs. ferrimagnet



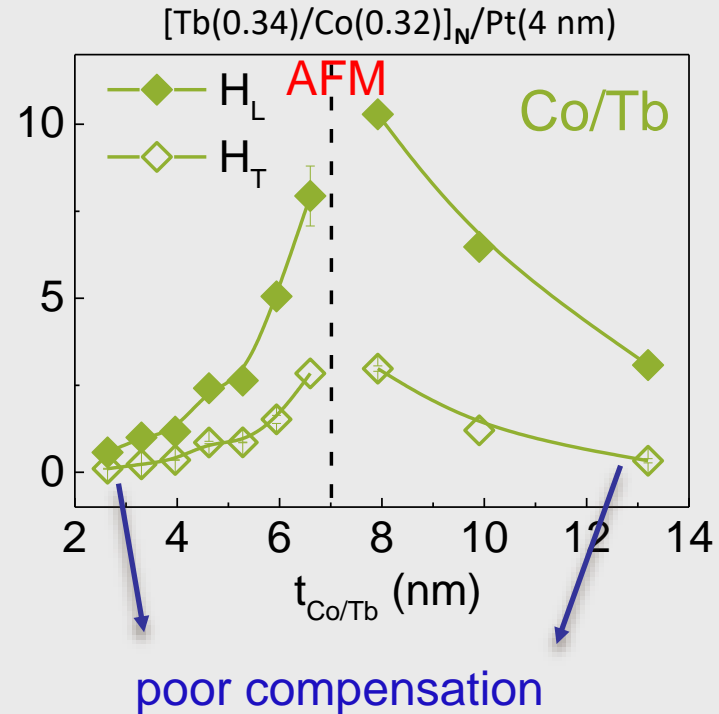
Alternating exchange fields in ferrimagnet (FIM) on an atomic scale
 → much less spin dephasing → long spin coherence length
 → Bulk-like torque in FIM (**20 times** enhanced SOT compared to FM)

SOT effective fields

Ferromagnet

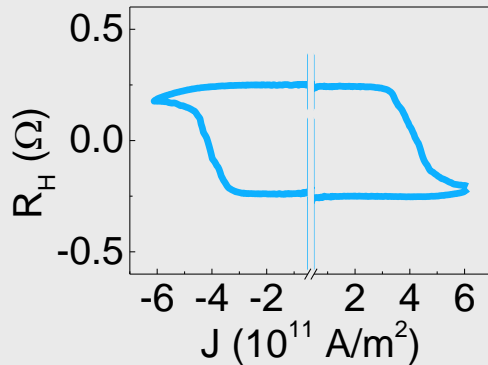


Ferrimagnet



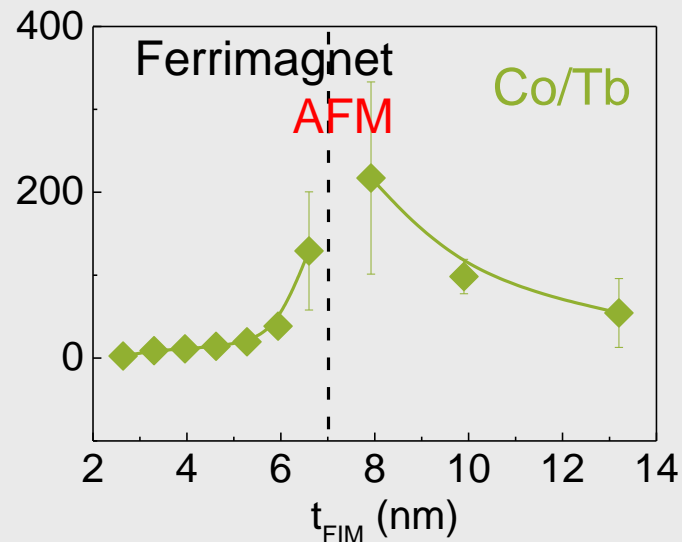
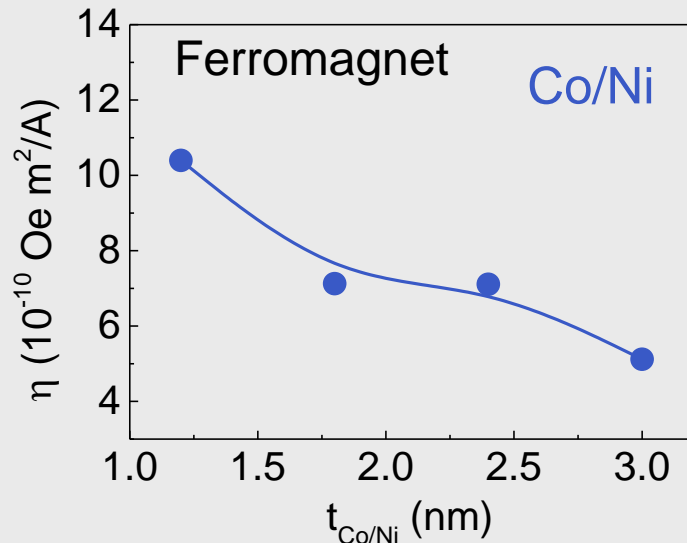
- FM Co/Ni: $t_{\text{Co/Ni}} \uparrow \rightarrow \text{SOT} \downarrow$.
 - FIM Co/Tb: SOT **diverges** at **compensation**.
 - SOT in FIM is **~20 times larger** than that in FM.
 - SOT in FIM shows **bulk-like-torque** characteristic.
- $$H_L = \frac{h}{2e} \frac{J_{HM}}{M_S t_{FM}} \theta_{HM} \left[1 - \text{sech} \left(\frac{t_{HM}}{\lambda_{HM}} \right) \right]$$

SOT current induced switching efficiency



Most switching through DW nucleation and propagation

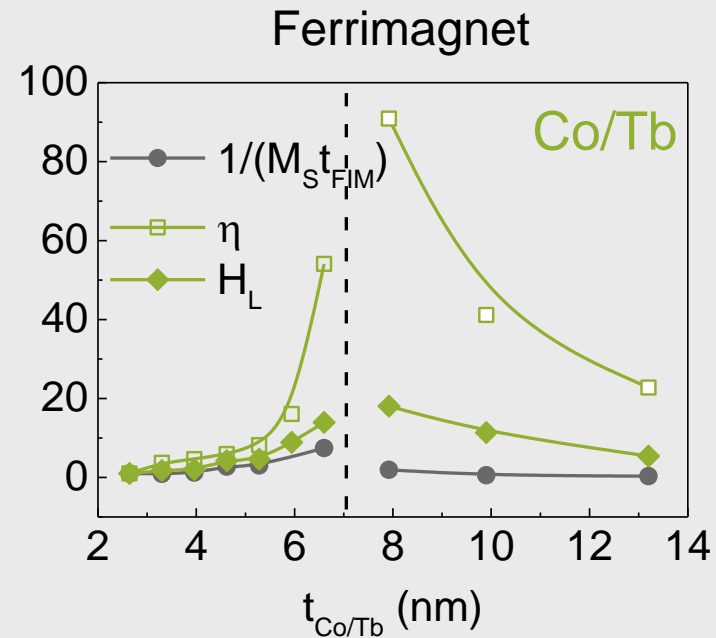
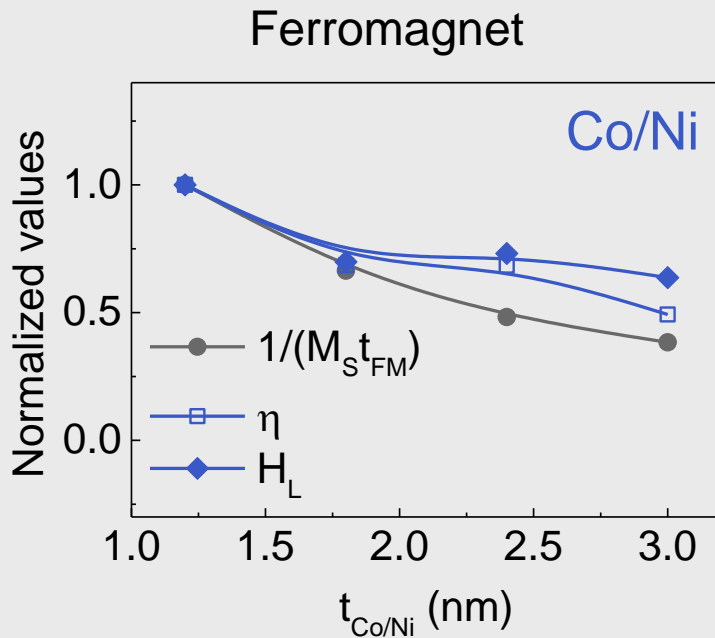
SOT switching efficiency: $\eta = H_P / J_W$



- η is in line with that of SOT effective fields, decreasing for FM and **diverging** at **compensation**.
- In FIM system, η is **~20 times higher** than FM system.

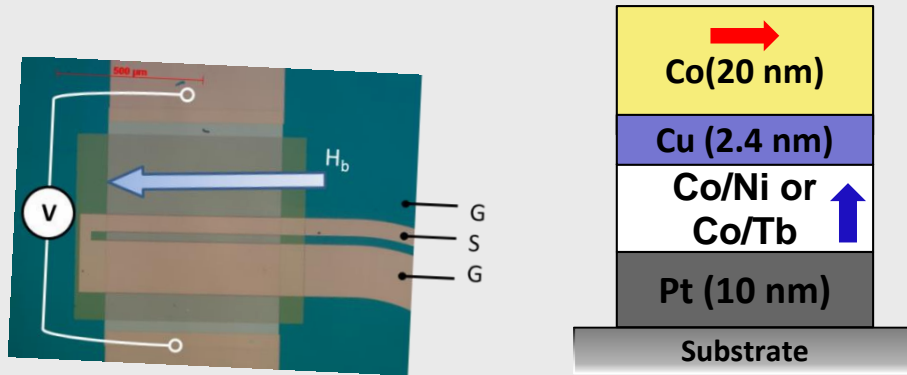
M_S & $t_{\text{FM}/\text{FIM}}$ effect?

$$H_L = \frac{\hbar}{2e} \frac{J_{\text{HM}}}{M_S t_{\text{FM}}} \theta_{\text{HM}} \left[1 - \text{sech} \left(\frac{t_{\text{HM}}}{\lambda_{\text{HM}}} \right) \right] \rightarrow \eta, H_L \propto 1 / (M_S t)$$

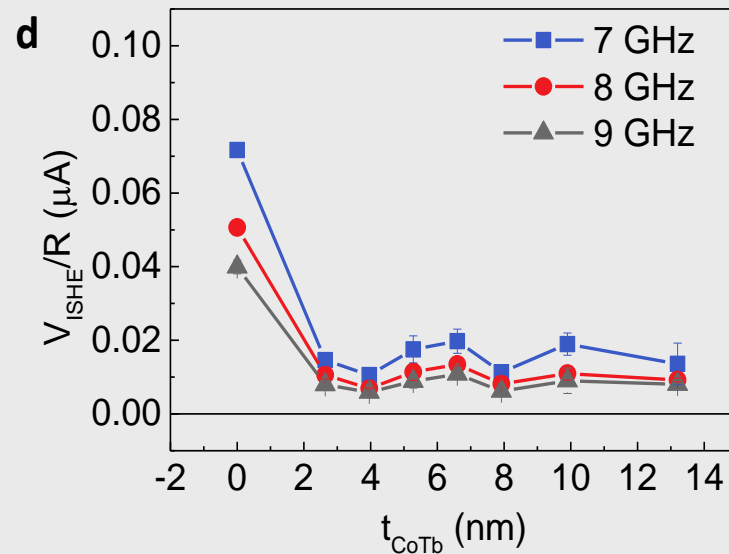
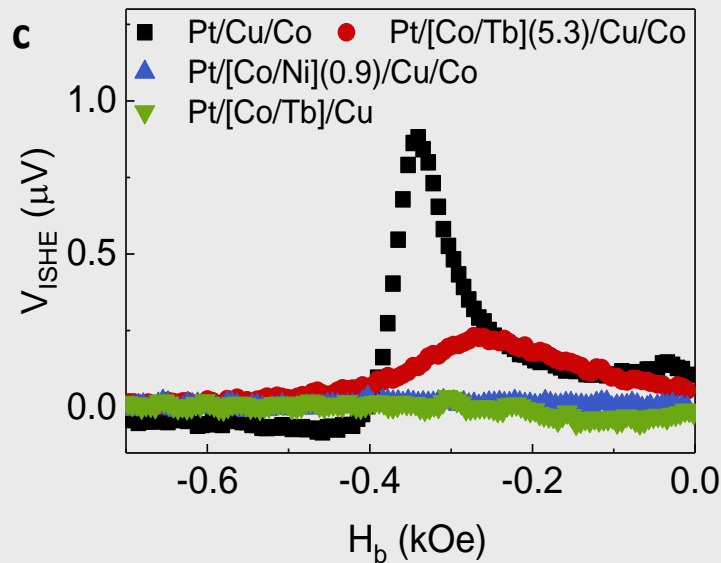


Surprisingly in FIM, the scaling trend of η and H_L **contradicts** SOT governing equation.

Spin pumping in a ferrimagnet



-If spin coherence length is long, a transverse spin current passes through the Co/Ni or Co/Tb layer.
 - V_{ISHE} is generated.



Spin dephasing length in Co/Tb > 13 nm
 Spin dephasing in FM is 1.2 nm (PRL 117, 217206 (2016))

Summary

- Spin orbit technologies
 - **Oxygen** modulation in HM/FM interface can switch spin Hall sign → neuromorphic devices
 - Spin accumulation **imaging** (www.tuotuo.com)
 - **Weyl** semimetal spin sources
 - **Ferri**magnetic spintronics
- Bulk-like

