



Axel Hoffmann

Department of Materials Science and Engineering University of Illinois at Urbana-Champaign

SPICE-SPIN+X, August 26, 2020







MAGNETIC MATCHMAKING: HYBRID MAGNON MODES

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SPICE-SPIN+X = XICE





OUTLINE

Introduction

- On-chip magnon-photon coupling with superconducting resonator
- Magnon-magnon coupling in magnetic bilayer thin film
- Manipulating phonon transport with

magnons

Conclusions



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Moore's Law



Moore's Law



Moore's Law



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Moore's Law



COMPUTING OF TODAY









COMPUTING OF TODAY











COMPUTING OF TOMORROW (MAYBE)

Quantum computing



COMPANIES > GOOGLE (ALPHABET)

Google Says Quantum Computer Beat 10,000-Year Task in Minutes





















Kurizki, et al. PNAS 112, 3866 (2015)





Kurizki, et al. PNAS 112, 3866 (2015)





X. Zhang et al. PRL 113, 156401 (2014)





X. Zhang et al. PRL 113, 156401 (2014)





X. Zhang et al. PRL 113, 156401 (2014)







Uniform mode (k=0)









Landau-Lifshitz-Gilbert equation:

$$\frac{d\mathbf{m}}{dt} = -\gamma \mathbf{m} \times \mathbf{H}_{\text{eff}} + \alpha \mathbf{m} \times \frac{d\mathbf{m}}{dt}$$



Uniform mode (k=0)



Landau-Lifshitz-Gilbert equation:

$$\frac{d\mathbf{m}}{dt} = \left[-\gamma\mathbf{m} \times \mathbf{H}_{eff}\right] + \alpha\mathbf{m} \times \frac{d\mathbf{m}}{dt}$$



Uniform mode (k=0)



Landau-Lifshitz-Gilbert equation:

$$\frac{d\mathbf{m}}{dt} = -\gamma \mathbf{m} \times \mathbf{H}_{eff} + \alpha \mathbf{m} \times \frac{d\mathbf{m}}{dt}$$

Gyromagnetic ratio: $\gamma/2\pi \sim 28$ GHz/Tesla Gilbert damping: $\alpha \sim 10^{-3}$

(Q factor~ $1/\alpha$)



Spin wave (k>0)







Spin wave (k>0)





Magnon dispersion:

 $\omega = \gamma H_{eff} + Dk^2$

D: Exchange constant (magnon: $Dk^2 \sim 3$ GHz for $\lambda \sim 100$ nm)



Spin wave (k>0)





Magnon dispersion:D: Exchange constant
(magnon:
$$Dk^2 \sim 3$$
 GHz for $\lambda \sim 100$ nm) $\omega = \gamma H_{eff} + Dk^2$ (Photon/phonon dispersion:)(Photon/phonon dispersion:)(photon: $\omega \sim 3$ GHz for $\lambda \sim 100$ mm)
(phonon: $\omega \sim 3$ GHz for $\lambda \sim 1 \mu$ m) $\omega = ck$



Spin wave (k>0)





Small + more flexible

Magnon dispersion:D: Exchange constant
(magnon: $Dk^2 \sim 3$ GHz for $\lambda \sim 100$ nm) $\omega = \gamma H_{eff} + Dk^2$ (photon: $\omega \sim 3$ GHz for $\lambda \sim 100$ nm)
(phonon: $\omega \sim 3$ GHz for $\lambda \sim 100$ nm)
(phonon: $\omega \sim 3$ GHz for $\lambda \sim 1 \mu$ m)



Advantage of Magnons for coherent information processing:

• Strong Coupling (sub-GHz)





Advantage of Magnons for coherent information processing:

- Strong Coupling (sub-GHz)
- Quantum Transduction





Advantage of Magnons for coherent information processing:

- Strong Coupling (sub-GHz)
- Quantum Transduction
- Wave Propagation



Magnonic logic



Advantage of Magnons for coherent information processing:

- Strong Coupling (sub-GHz)
- Quantum Transduction
- Wave Propagation
- On-chip Integration



Microwave quantum circuits





QUANTUM MAGNONICS



Y. Tabuchi et al. Science 349, 405 (2015) D. Lachance-Quirion, et al. Science 367, 425 (2020)
QUANTUM MAGNONICS

Goal of quantum magnonics: microwave quantum circuits





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Most previous demonstrations: <u>Bulk FM (YIG)</u> or <u>3D cavity</u>



Tabuchi et al. PRL 113, 083603 (2014) Bai et al. PRL 114, 227201 (2015) Morris et al. Sci. Rep. 7, 11511 (2017)



ON-CHIP MAGNON-PHOTON HYBRID SYSTEM

Coplanar superconducting resonator





ON-CHIP MAGNON-PHOTON HYBRID SYSTEM

Coplanar superconducting resonator



Photon system: NbN coplanar superconducting resonator

<u>Magnon system:</u> **Ni₈₀Fe₂₀** device (permalloy, Py)

• High T_c (14 K)

Large magnetization ($\mu_0 M_s = 1 T$)

Yi Li, et al. PRL 123, 107701 (2019)

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SUPERCONDUCTING CIRCUIT





SUPERCONDUCTING CIRCUIT











PHOTON SYSTEM: UNLOADED NbN SUPERCONDUCTING RESONATOR





(0 dBm = 1 mWatt) (-55 dBm = 4 nWatt)

PHOTON SYSTEM: UNLOADED NbN SUPERCONDUCTING RESONATOR







] [

PHOTON SYSTEM: UNLOADED NbN SUPERCONDUCTING RESONATOR





MAGNONS SYSTEM: Ni₈₀Fe₂₀ (Py) STRIPE



NbN resonator



T=1.4 K Yi Li, et al. P_{in}=-55 dBm PRL 123, 107701 (2019)



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Photon damping rate: $\kappa_p/2\pi$ =2.0 MHz Q=2500 NbN resonator NbN resonator + Py device 4.96 Photon 4.94



P_{in}=-55 dBm



Photon damping rate: $\kappa_p/2\pi$ =2.0 MHz Q=2500 NbN resonator NbN resonator + Py device 5.08 4.96 Photon 4.94 5.06 (2H2 4.92 4.90 (CHZ) μ2/σ Magnon $S_{21}(dB)$ $S_{21}(dB)$ Hybrid -30 5.00 4.88 mode -40 -50 With Py $\theta = 0^{\circ}$ without Py 4,98 4.86 50 -50 50 -100-50 100 -100100 0 0 $\mu_0 H_B$ (mT) $\mu_0 H_B$ (mT) H_{B} magnon Coupling energy: $E = \mu_0 \vec{M} \cdot \vec{h}_{rf}$ HT=1.4 K Yi Li, et al. photon PRL 123, 107701 (2019) P_{in}=-55 dBm





















TUNE COUPLING EFFICIENCY



TUNE COUPLING EFFICIENCY



NUMBER OF SPINS



62



NONLINEARITY FROM SUPERCONDUCTING VORTICES



NONLINEARITY FROM SUPERCONDUCTING VORTICES





NONLINEARITY FROM SUPERCONDUCTING VORTICES







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y Microwave quantum circuits



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Magnon hybrid system





MAGNETIC BILAYER FOR COHERENT INFORMATION PROCESSING



Klingler et al. PRL 120, 127201 (2018) Chen et al. PRL 120, 217202 (2018) Qin et al. Sci. Rep. 8, 5755 (2018)



MAGNETIC BILAYER FOR COHERENT INFORMATION PROCESSING



Klingler et al. PRL 120, 127201 (2018) Chen et al. PRL 120, 217202 (2018) Qin et al. Sci. Rep. 8, 5755 (2018)



In-plane Kittel equation:

$$\frac{\omega}{\gamma} = \sqrt{H_B(H_B + M_s)}$$




Klingler et al. PRL 120, 127201 (2018) Chen et al. PRL 120, 217202 (2018) Qin et al. Sci. Rep. 8, 5755 (2018)



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Klingler et al. PRL 120, 127201 (2018) Chen et al. PRL 120, 217202 (2018) Qin et al. Sci. Rep. 8, 5755 (2018)



In-plane Kittel equation:

$$\frac{\omega}{\gamma} = \sqrt{(H_B + H_{ex})(H_B + M_s + H_{ex})}$$

$$\mu_0 H_{ex} = D_{ex} k^2$$





Perpendicular standing spin wave mode

$$k = \frac{n\pi}{t}$$

Klingler et al. PRL 120, 127201 (2018) Chen et al. PRL 120, 217202 (2018) Qin et al. Sci. Rep. 8, 5755 (2018)



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Structure: YIG(100 nm)/Py(5-60 nm)



Y₃Fe₅O₁₂ (YIG):

- Sputtering on Gd₃Ga₅O₁₂ (GGG)
- Post-annealing in air at 850°C

Ni₈₀Fe₂₀ (Py):

- Ion milling of YIG surface
- Sputtering of Py on YIG



Structure: YIG(100 nm)/Py(5-60 nm) H_B YIG (n=0)YIG (n=2)Py(n=0)

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Py

• ω



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- Ion milling of YIG surface
- Sputtering of Py on YIG









Yi Li, et al. PRL 124, 117202 (2020)





Yi Li, et al. PRL 124, 117202 (2020)







YIG(100 nm)/Py(7.5 nm)





YIG(100 nm)/Py(7.5 nm)





YIG(100 nm)/Py(7.5 nm)



Yi Li, et al. PRL 124, 117202 (2020)







MACROSPIN MODEL: COHERENT COUPLING

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Macrospin model:

 $\frac{d\mathbf{m}_i}{dt} = -\mu_0 \gamma_i \mathbf{m}_i \times \mathbf{H}_{eff} + \alpha_i \mathbf{m}_i \times \frac{d\mathbf{m}_i}{dt} \\ -\gamma_i \mathbf{m}_i \times \frac{J}{M_i t_i} \mathbf{m}_j + \Delta \alpha_i (\mathbf{m}_i \times \frac{d\mathbf{m}_i}{dt} - \mathbf{m}_j \times \frac{d\mathbf{m}_j}{dt})$

Fieldlike: exchange Dampinglike: spin pumping



MACROSPIN MODEL: COHERENT COUPLING

Macrospin model:

 $\frac{d\mathbf{m}_i}{dt} = -\mu_0 \gamma_i \mathbf{m}_i \times \mathbf{H}_{eff} + \alpha_i \mathbf{m}_i \times \frac{d\mathbf{m}_i}{dt} \\ -\gamma_i \mathbf{m}_i \times \frac{J}{M_i t_i} \mathbf{m}_j + \Delta \alpha_i (\mathbf{m}_i \times \frac{d\mathbf{m}_i}{dt} - \mathbf{m}_j \times \frac{d\mathbf{m}_j}{dt})$

Fieldlike: exchange Dampinglike: spin pumping

Solution:

Fieldlike: Dampinglike:

$$g_c = \sqrt{\frac{\mathbf{J}}{M_1 t_1} \cdot \frac{\mathbf{J}}{M_2 t_2}}$$
$$\kappa_c = \sqrt{\frac{\mathbf{J'}}{M_1 t_1} \cdot \frac{\mathbf{J'}}{M_2 t_2}}$$

 $J'(\omega) = \frac{g^{\uparrow\downarrow}}{4\pi} \hbar \omega$ g^{↑↓}: spin mixing conductance, ~ 42 nm⁻² for YIG/Py ₉₂ Yi Li, et al. PRL 124, 117202 (2020)



MACROSPIN MODEL: COHERENT COUPLING



$$\kappa_c(\omega) = \beta \mu_0 \Delta H_{sp}^{\rm Py}(\omega)$$













Antiferromagnetic coupling: $g_c < 0$







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Yi Li, et al. PRL 124, 117202 (2020)





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ELASTICALLY DRIVEN MAGNETIZATION DYNAMICS



Free Energy for magnetization with Surface Acoustic Wave (SAW):

$$F^{0}(\mathbf{m}) = -\mu_{0}\mathbf{H}\cdot\mathbf{m} + B_{d}m_{z}^{2} + B_{u}(\mathbf{u}\cdot\mathbf{m})^{2} + B_{1}\varepsilon(\mathbf{x},t)m_{x}^{2}$$

Zeeman Demagentizing Anisotropy Magnetoelastic

M. Weiler, et al. PRL 106, 117601 (2011)



ELASTICALLY DRIVEN MAGNETIZATION DYNAMICS



M. Weiler, et al. PRL 106, 117601 (2011)

FMR given by broad maxima

Why is the linewidth so large?



CAN MAGNONS MODIFY PHONON TRANSPORT?



127.86° Y-X-cut LiNbO₃ substrate 50-nm thick Ni film

Two types of devices

- As grown
- Annealed in vacuum at 400°C for 30 min.



C. Zhao, et al. PRAppl. 13, 054032 (2020)

DETECT FERROMAGNETIC RESONANCE



Annealed



C. Zhao, et al. PRAppl. 13, 054032 (2020)



PHONON TRANSMISSION





PHONON TRANSMISSION



C. Zhao, et al. PRAppl. 13, 054032 (2020)

PHONON TRANSMISSION

I. A. Privorotskii, IEEE Trans. Mag. **16**, 666 (1980)
$$E = E_0 \frac{v_f}{c} \frac{B_2^2}{C_{44}} \frac{\left|h_{eff}^{\perp}\right|}{\Delta H^2} (1+p^2)^{-\frac{1}{2}} [1+(p+\beta)^2]^{-\frac{1}{2}} \exp\left[\frac{-(\eta/2)}{1+(p+\beta)^2}\right] \left[1+\exp\left(\frac{-\eta_0 H}{2}\right)\right]$$

Includes both phonon attenuation and generation

 η Is proportional to magnetoelastic coupling



C. Zhao, et al. PRAppl. 13, 054032 (2020)

LINEWIDTH



FMR linewidth from fit agrees well with inductively measured FMR of Ni on LiNbO₃

Increased apparent linewidth in phonon transmission is due to magnon-phonon interaction and not due to extrinsic inhomogeneities

MAGNETIC FILMS GROUP AT ARGONNE




MAGNETIC FILMS GROUP AT ARGONNE





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CONCLUSION

Magnon-photon

Magnon-magnon

Magnon-phonon



Yi Li, et al. PRL 123, 107701 (2019) Yi Li, et al. PRL 124, 117202 (2020) Chenbo Zhao, et al. PRAppl 13, 054032 (2020)

CONCLUSION

Magnon-photon



Magnon-magnon



H_B YIG(n=2) Py(n=0) h_{rf}

On-chip applications

Yi Li, et al. PRL 123, 107701 (2019) Yi Li, et al. PRL 124, 117202 (2020)

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Chenbo Zhao, et al. PRAppl 13, 054032 (2020)_

KSAM

100 µm



IDT

Magnon-phonon



LiNbO₃

Ni film

IDT

