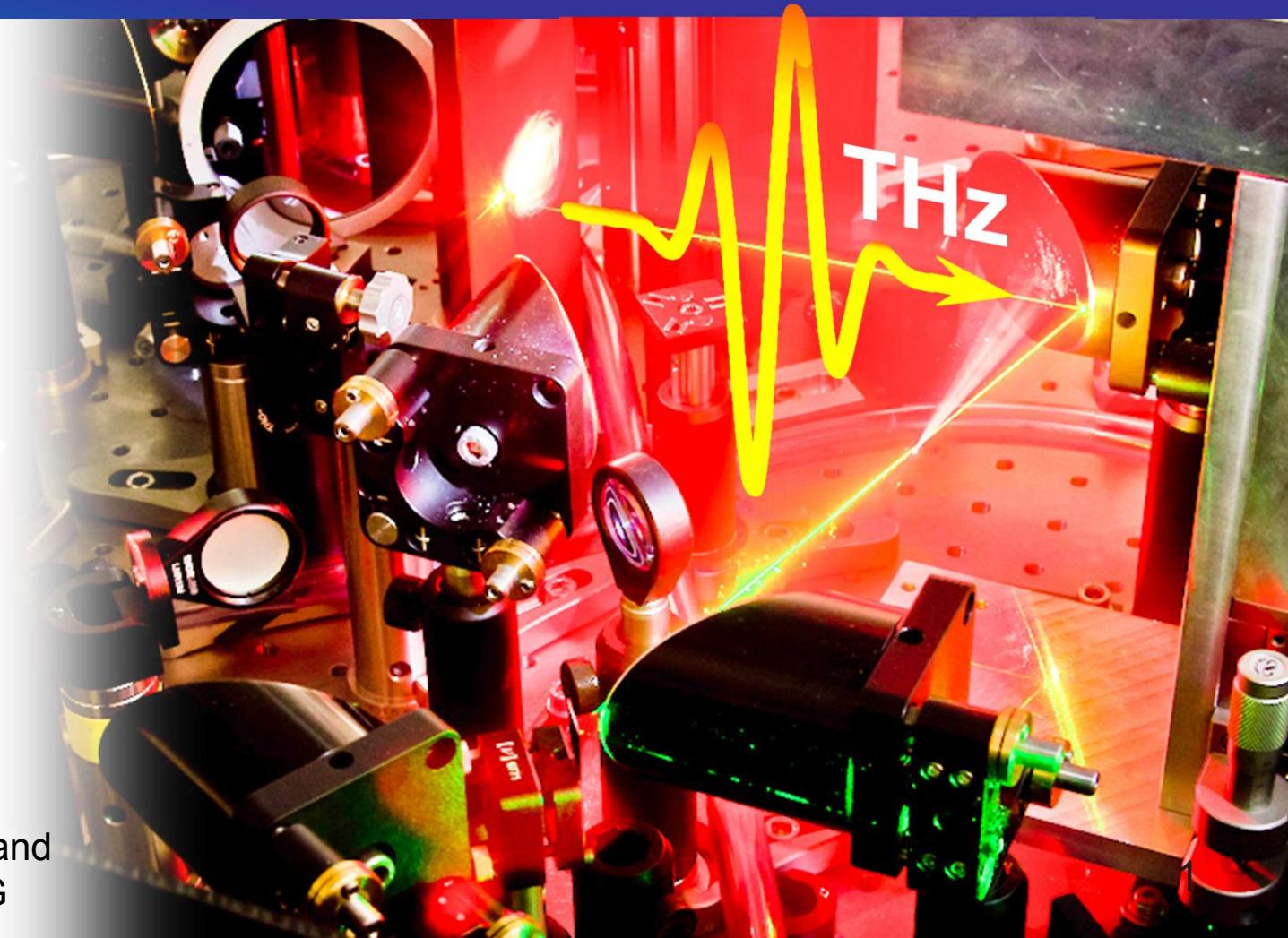
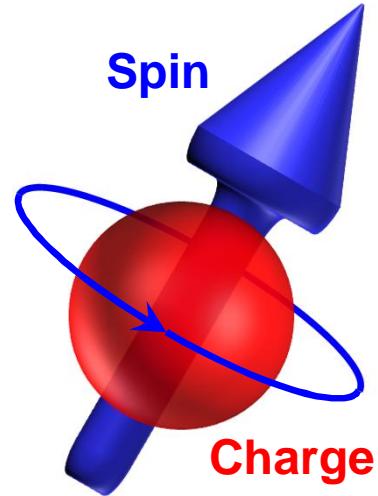


Terahertz spin transport in magnetic nanostructures



Tobias Kampfrath
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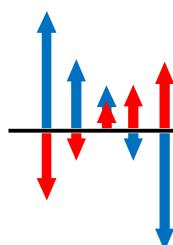


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UNIVERSITÄT
HALLE-WITTENBERG

Liane Brandt
Georg Woltersdorf



Richard Schlitz
S. Goennenwein



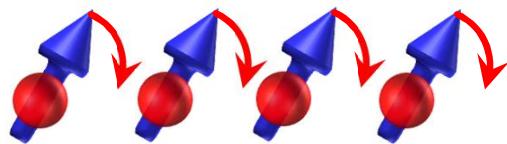
CRC / TRR 227

Ultrafast Spin Dynamics



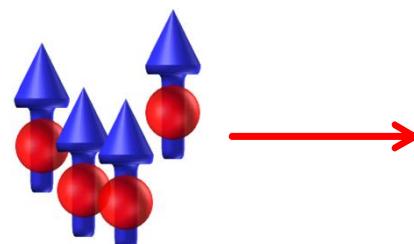
Spintronics: 3 basic operations

1. Turn spins around



→ Torque

2. Transport spins



→ Spin current

3. Detect spin dynamics



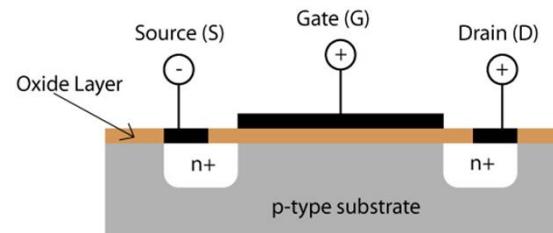
Goal: Reach speed of other information carriers: $1 \text{ THz} = 10^{12} \text{ Hz}$

Photons in fibers: $>10 \text{ Tbit/s}$



Hillerkuss *et al.*, Nature Phot. (2011)

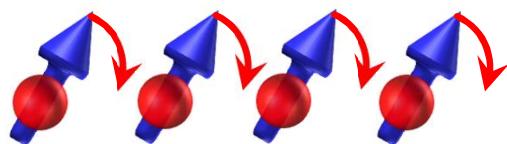
Electrons in a FET: $\sim 1 \text{ THz}$ cut-off



Del Alamo *et al.*, Nature (2011)

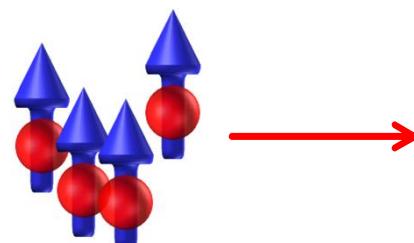
Spintronics: 3 basic operations

1. Turn spins around



→ Torque

2. Transport spins



→ Spin current

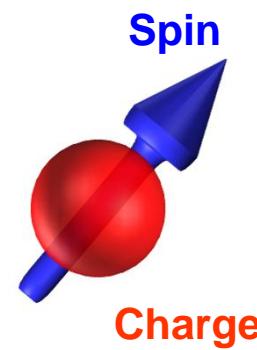
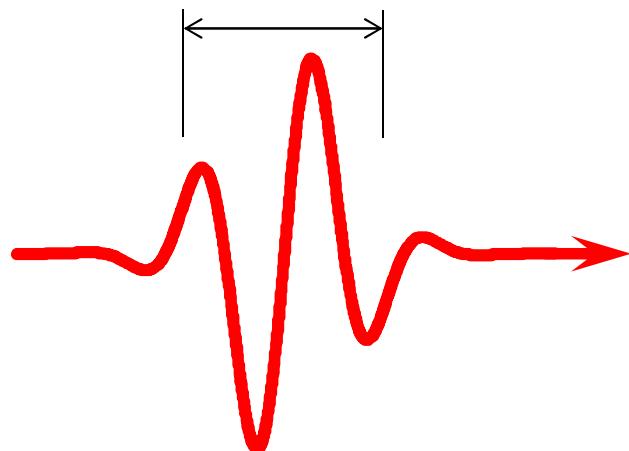
3. Detect spin dynamics



Goal: Reach speed of other information carriers: $1 \text{ THz} = 10^{12} \text{ Hz}$

Idea: Use THz electromagnetic pulses

$$1 \text{ ps} = (1 \text{ THz})^{-1} = 300 \mu\text{m}/c$$



THz + spin transport = useful?

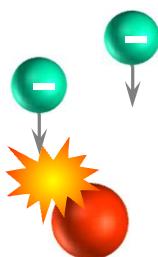
Why THz spin transport?

1) Test the speed of spintronic effects, reveal initial steps

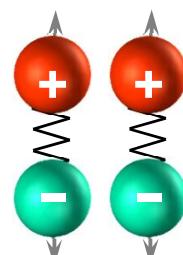
E.g. spin-Hall, spin-Seebeck and GMR

2) New physics, new methods

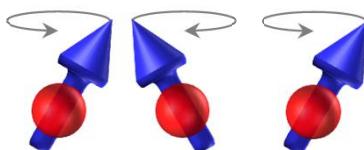
As THz coincides with many fundamental modes



Intraband transport



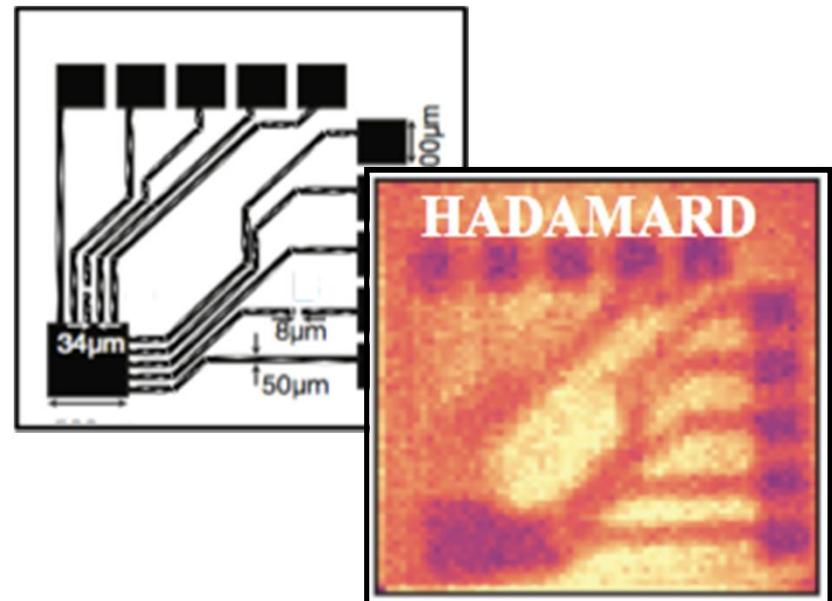
Phonons



Magnons

3) Reward for THz photonics

E.g. THz sources and modulators for spectroscopy and imaging



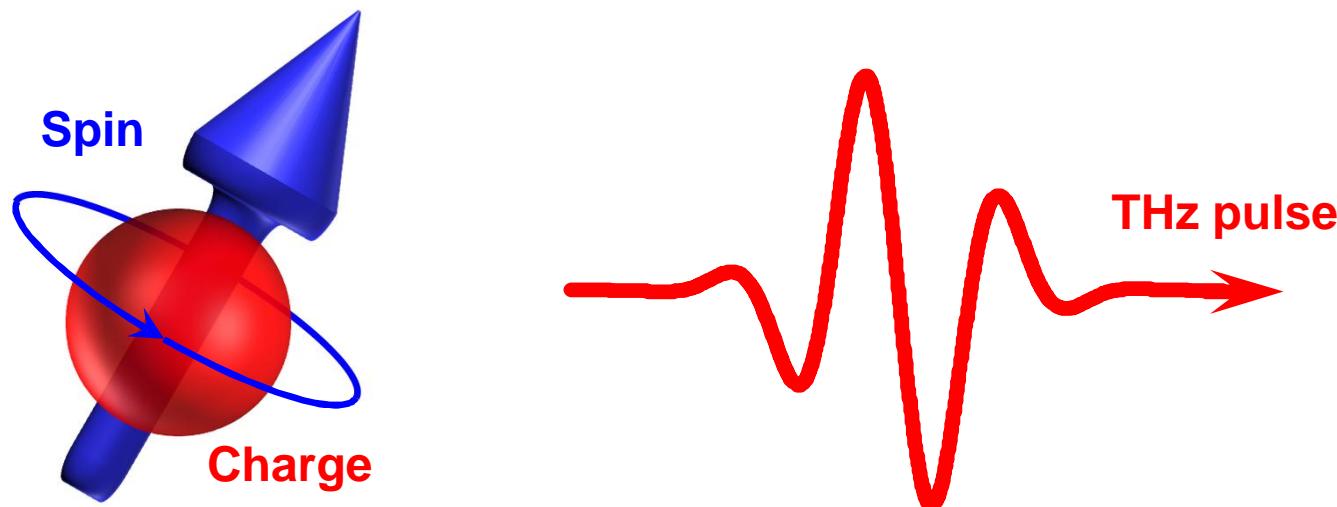
Stantchev *et al.*, Science Adv. (2016)

How to drive spin transport?

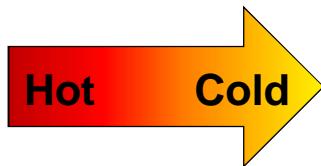
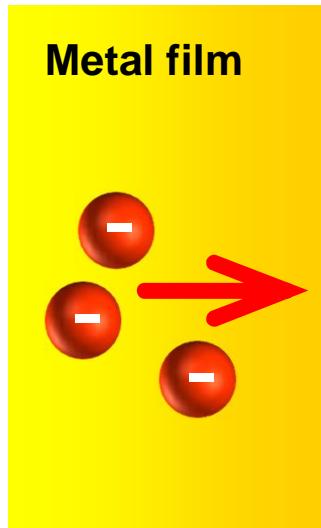
Important drivers of spin transport

- Electric fields: E.g. spin Hall effect
- Heat: Spin-caloric effects

How to transfer to the THz range?



Heat-driven: The Seebeck effect



Temperature
gradient

Thomas Seebeck (1821):

A temperature gradient drives an electron current

Uchida and Saitoh (2008):

In ferromagnets, the Seebeck current
is spin-dependent

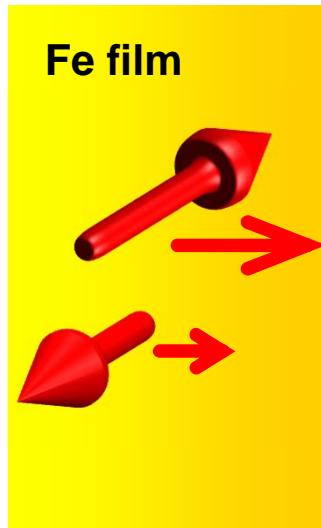
Spin-dependent Seebeck effect (SDSE)



\uparrow and \downarrow electrons have very different transport properties

Uchida, Saitoh *et al.*, Nature (2008)
Bauer, Saitoh, Wees, Nature Mat. (2013)

Spin-dependent Seebeck effect (SDSE)



↑ and ↓ electrons have very different transport properties

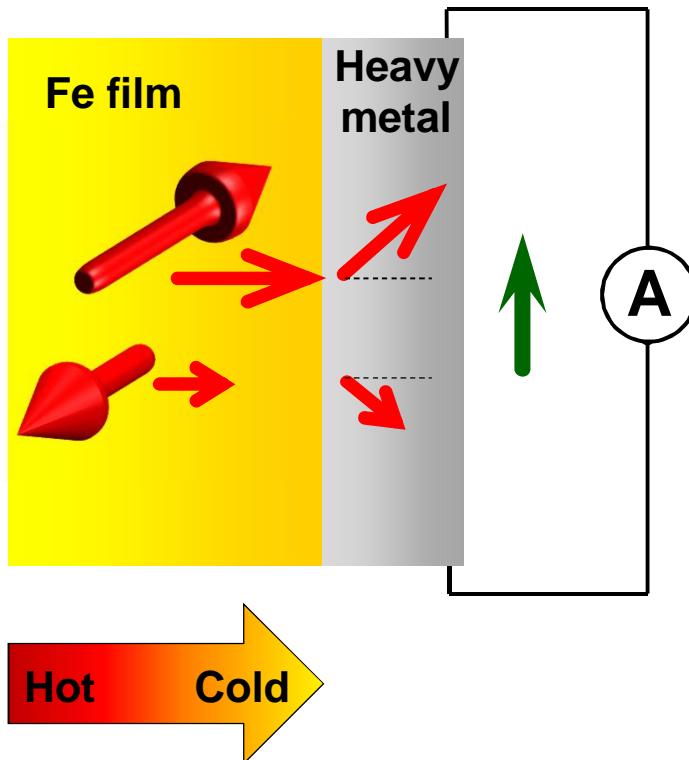
Uchida, Saitoh *et al.*, Nature (2008)
Bauer, Saitoh, Wees, Nature Mat. (2013)

⇒ Spin-polarized current

Detection with the
inverse spin Hall effect

Temperature
gradient

Inverse spin Hall effect (ISHE)

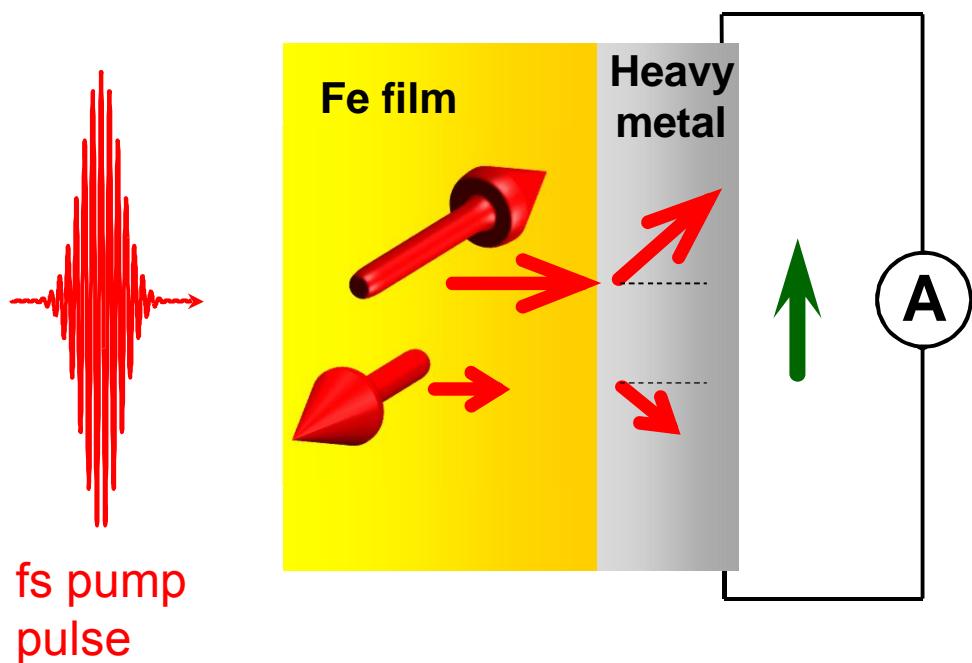


Temperature
gradient

Spin-orbit coupling
deflects electrons
⇒ Transverse charge current
⇒ Spin-to-charge (S2C)
conversion
Saitoh *et al.*, APL (2006)

**How can we induce an
imbalance as fast as possible?**

Inverse spin Hall effect (ISHE)



Ultrafast spin transport

Malinowski *et al.*, Nature Phys. (2008)
Melnikov, Bovensiepen *et al.*, PRL (2011)
Rudolf *et al.*, Nat. Comm. (2012)
Eschenlohr *et al.*, Nature Mat. (2013)
Turgut *et al.*, PRL (2013)
Schellekens *et al.*, Nat. Comm. (2014)
Ganichev *et al.*, Nature (2002)

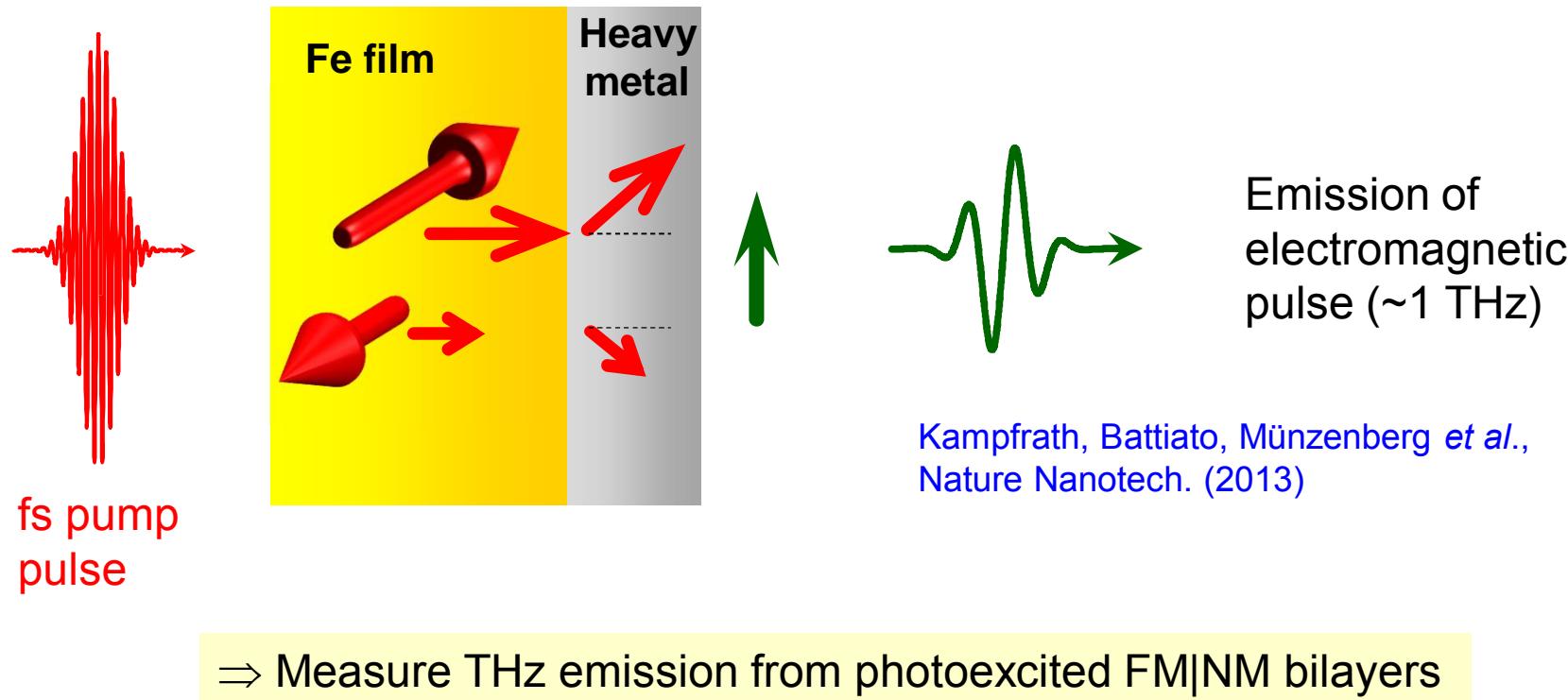
Goal:

Measure quasi-electrically

Technical challenge:

- Electric detection has cutoff at < 50 GHz
- But expect bandwidth > 10 THz

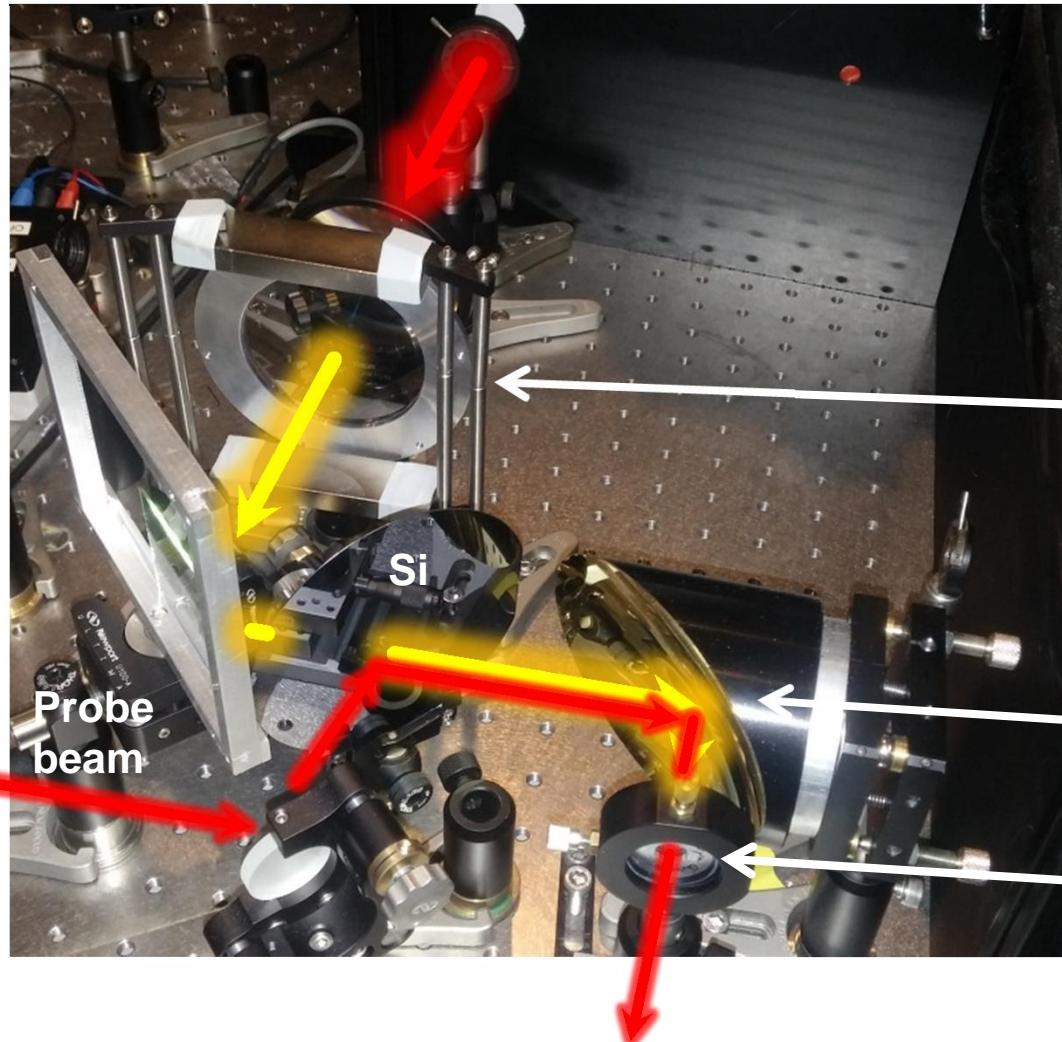
Inverse spin Hall effect (ISHE)



Samples: Polycrystalline films (from Kläui, Münzenberg, Woltersdorf labs)

Pump pulses: From Ti:sapphire oscillator (10 fs, 800 nm, 2.5 nJ, 80 MHz)

Simple THz emission setup in lab



Optical pump beam

Spintronic sample

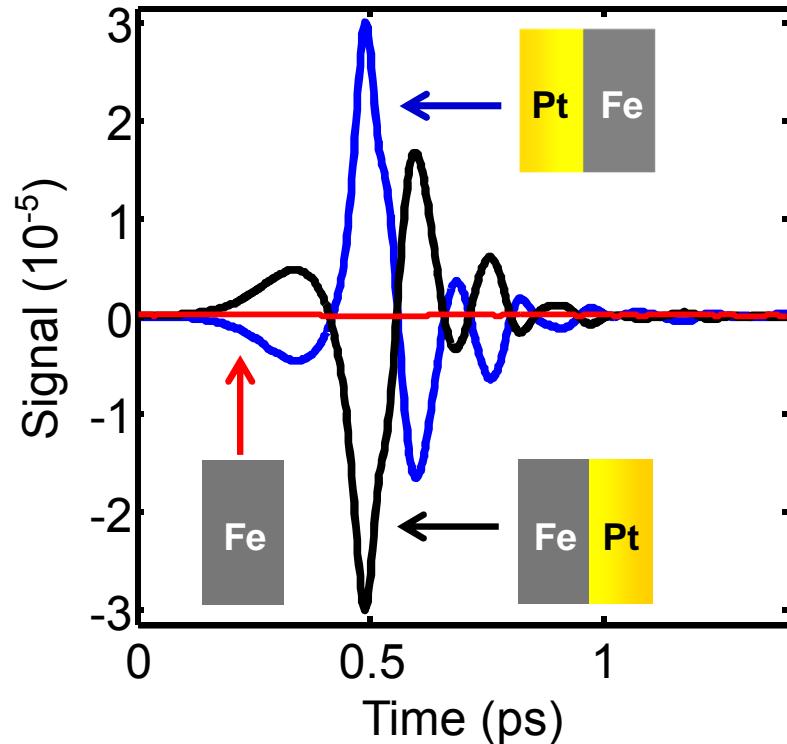
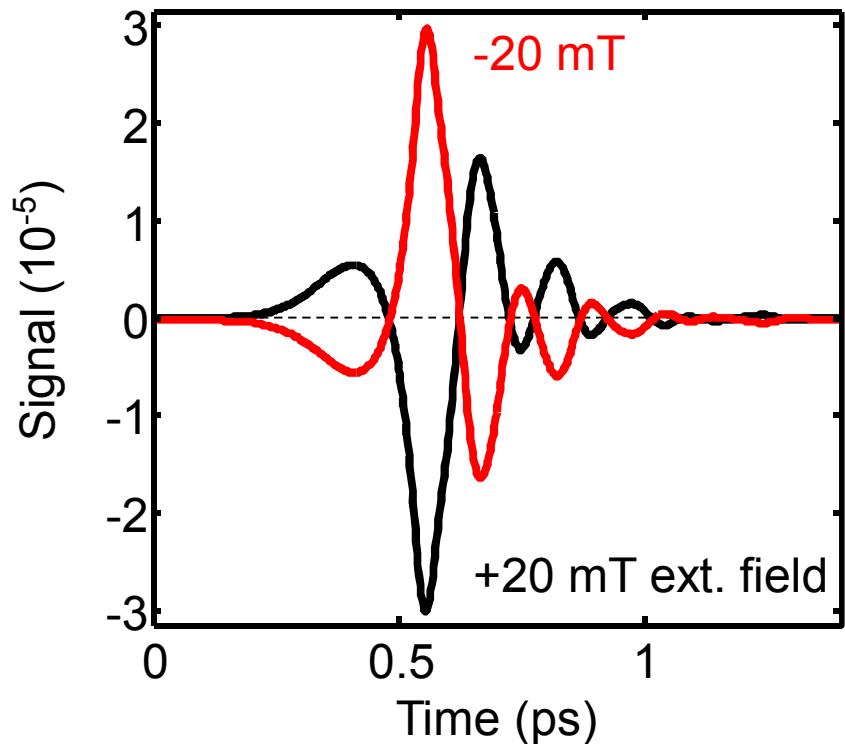
Parabolic mirror

Electro-optic crystal for sampling of
the THz electric field

To detection of
probe ellipticity

Typical signals

Typical THz waveforms from Fe|Pt bilayers



Further findings

- Signal \propto pump power
- THz electric field \perp sample magnetization

Consistent with scenario
spin transfer + ISHE

Need more evidence for the spin Hall scenario

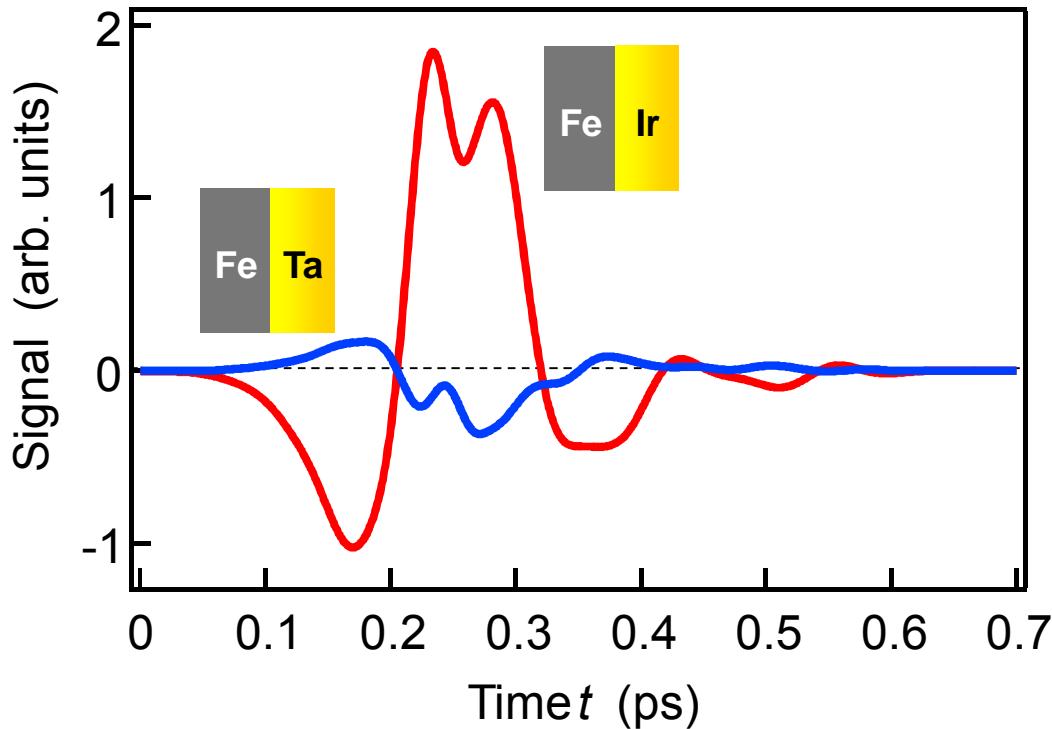
Ultrafast inverse spin Hall effect

Idea:

Vary nonmagnetic cap layer

Ta vs Ir:

Opposite spin Hall angles, Ir larger



The inverse spin Hall effect is still operative at THz frequencies

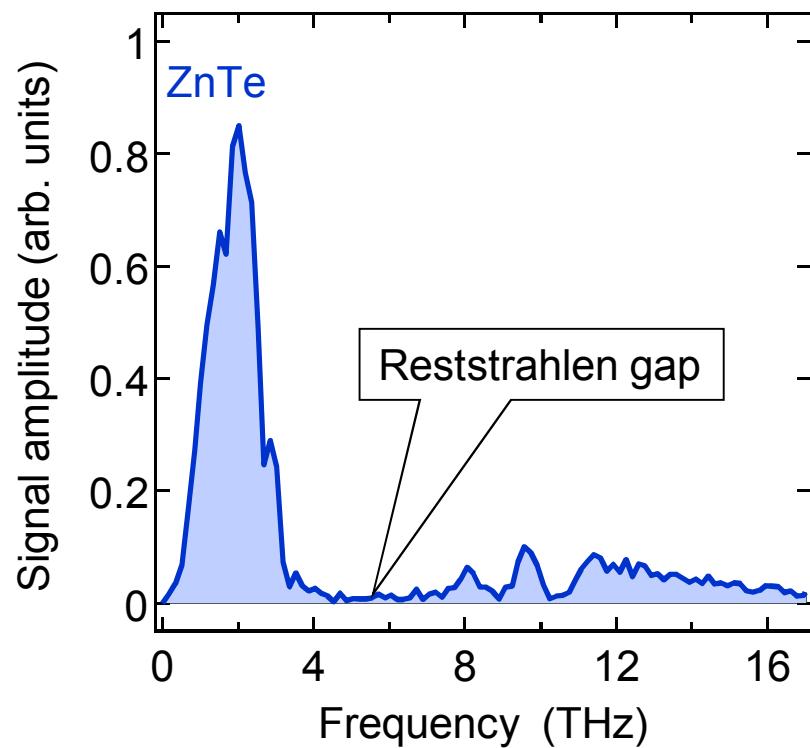
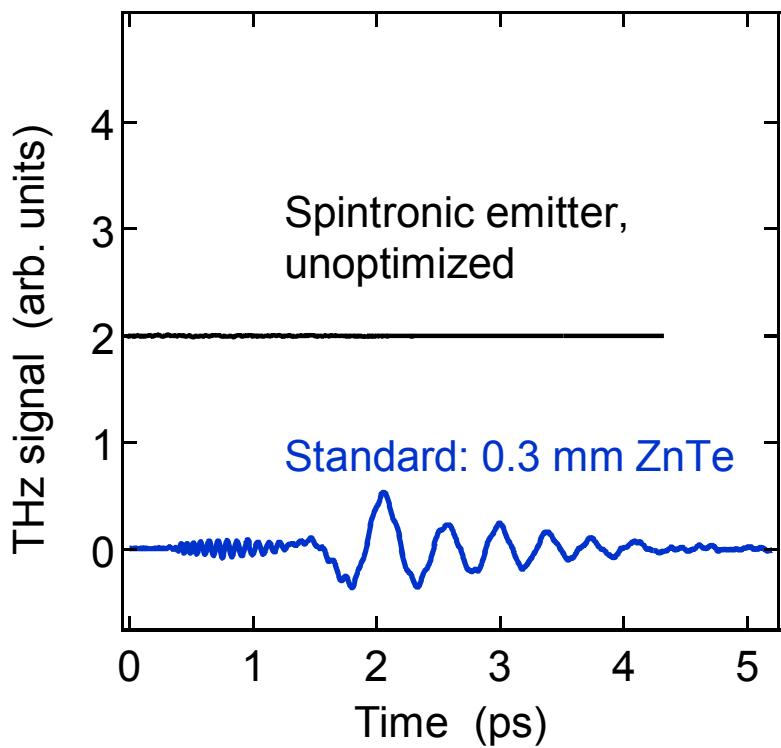
Kampfrath, Battiato, Oppeneer, Freimuth, Mokrousov, Radu, Wolf, Münzenberg *et al.*, Nature Nanotech. (2013)

Interesting applications:

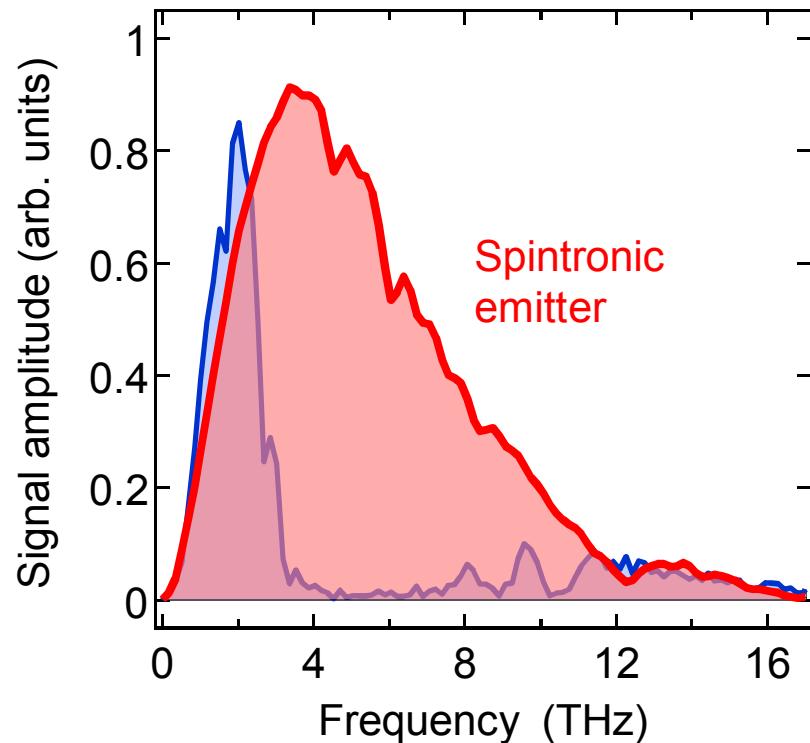
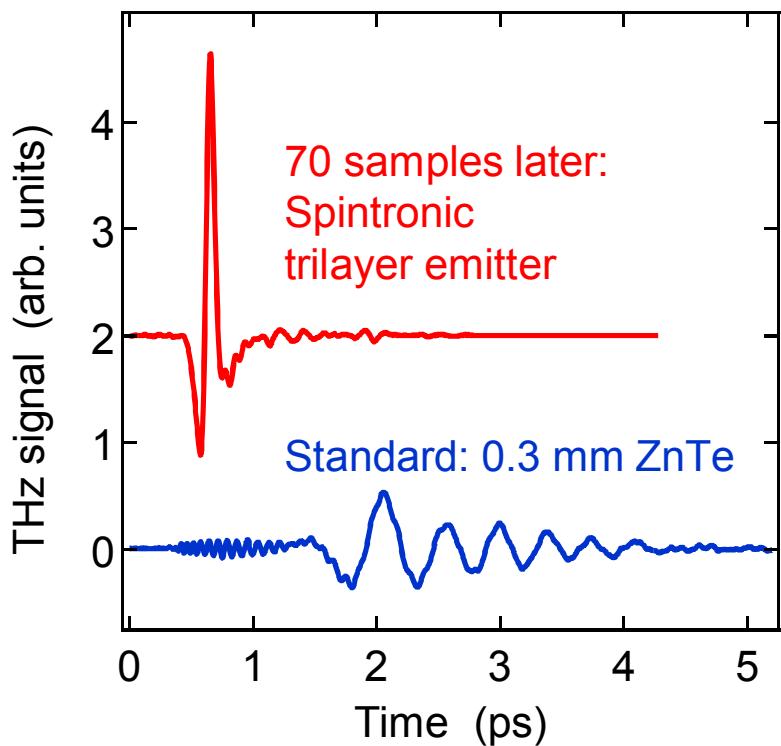
- 1) Rapid material characterization regarding spin-to-charge conversion
- 2) Measurement of spin-current dynamics
- 3) Generation of THz electromagnetic pulses

Compare to state-of-the-art emitter

Spintronic THz emitter



Spintronic THz emitter



More broadband, efficient and cheaper than standard emitters like ZnTe

Seifert, Radu, Kläui, TK *et al.*, Nature Photon. (2016)

Yang, Wu, Qi *et al.*, Adv. Opt. Mat. (2016)

Wu, Yang *et al.*, Adv. Mat. (2017)

Torosyan, Beigang, Papaioannou *et al.*, Sci. Rep. (2018)

Spintronic THz emitter in the lab



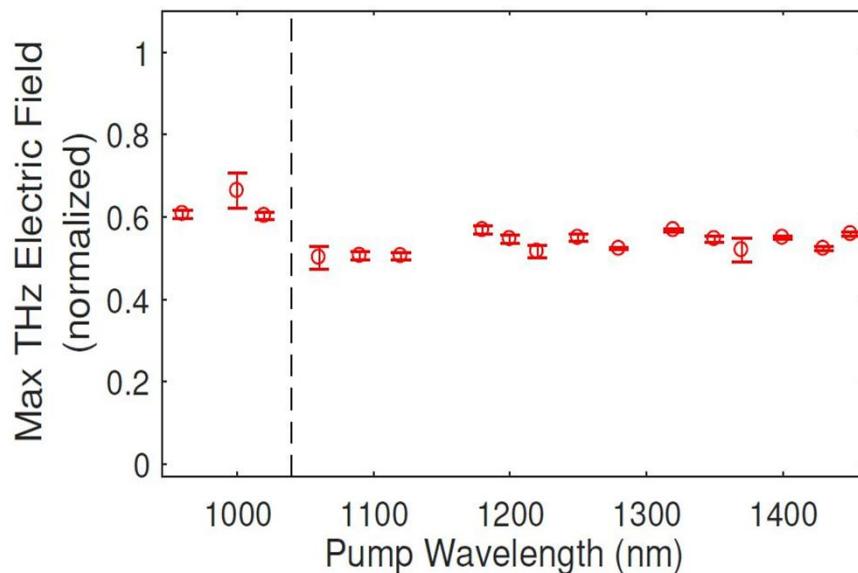
Movie by Oliver Gückstock
and Lukáš Nádvorník

Magnetism and thin metal
enable interesting features...

More features of spintronic THz emitters

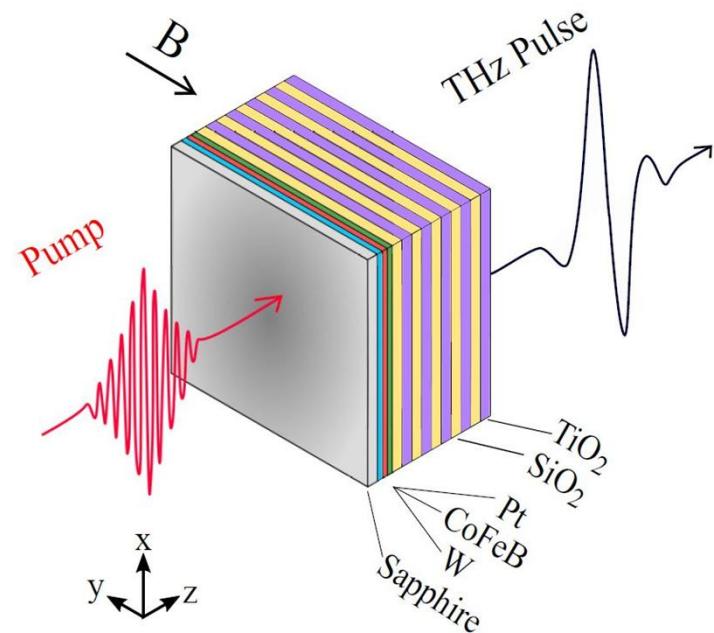
Emitter is insensitive to pump wavelength

Papaioannou, Beigang *et al.*, IEEE (2018)
Herapath, Seifert, TK, Hendry *et al.*, APL (2019)



Cavity pushes pump absorption from $\approx 50\%$ to $\approx 100\%$

Feng *et al.*, Adv. Opt. Mat (2018)
Herapath, Seifert, TK, Hendry *et al.*, APL (2019)

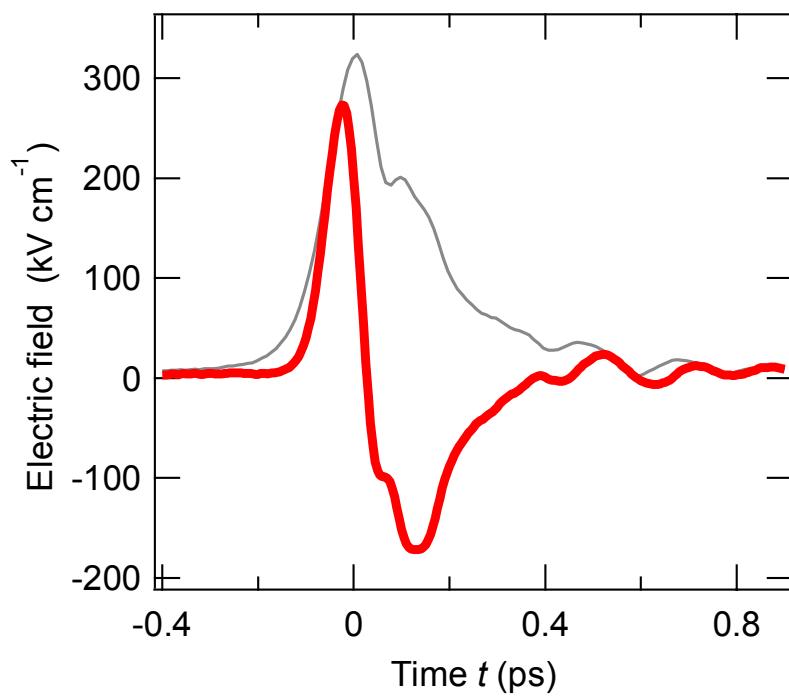


Larger emitters are easy and cheap to fabricate

More features of spintronic THz emitters

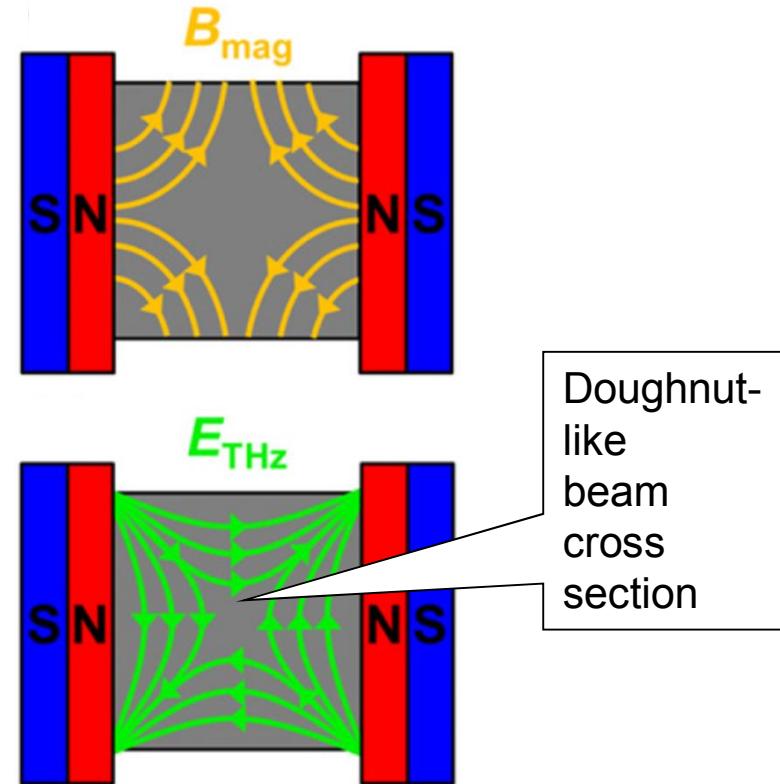
Upscaling to ~5 cm diameter
yields **0.3 MV/cm** peak field

Seifert, Kläui, TK et al., APL (2017)
Fülöp, Tzortzakis, TK, Adv. Opt. Mat. (2020)



Tailor the magnetization
landscape to generate
structured THz beams

Hibberd, Graham et al., APL (2019)



The magnetization can also be temporally modulated

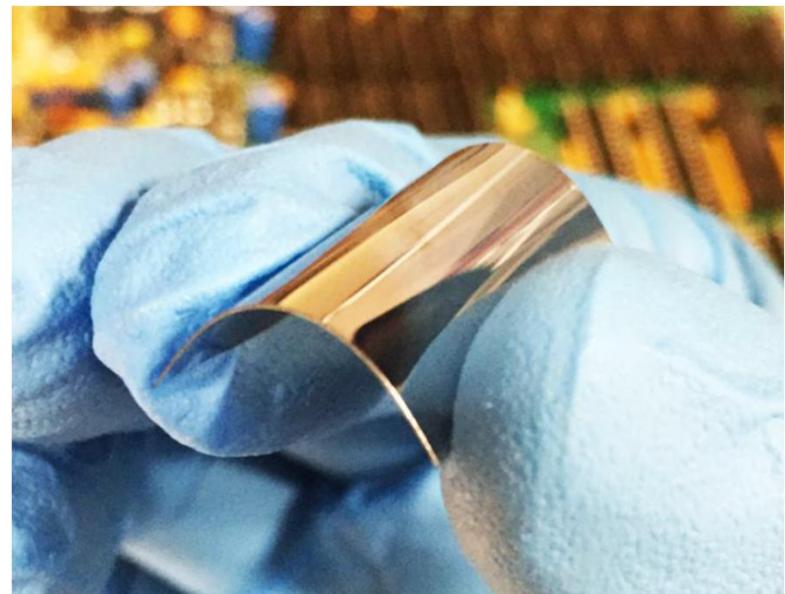
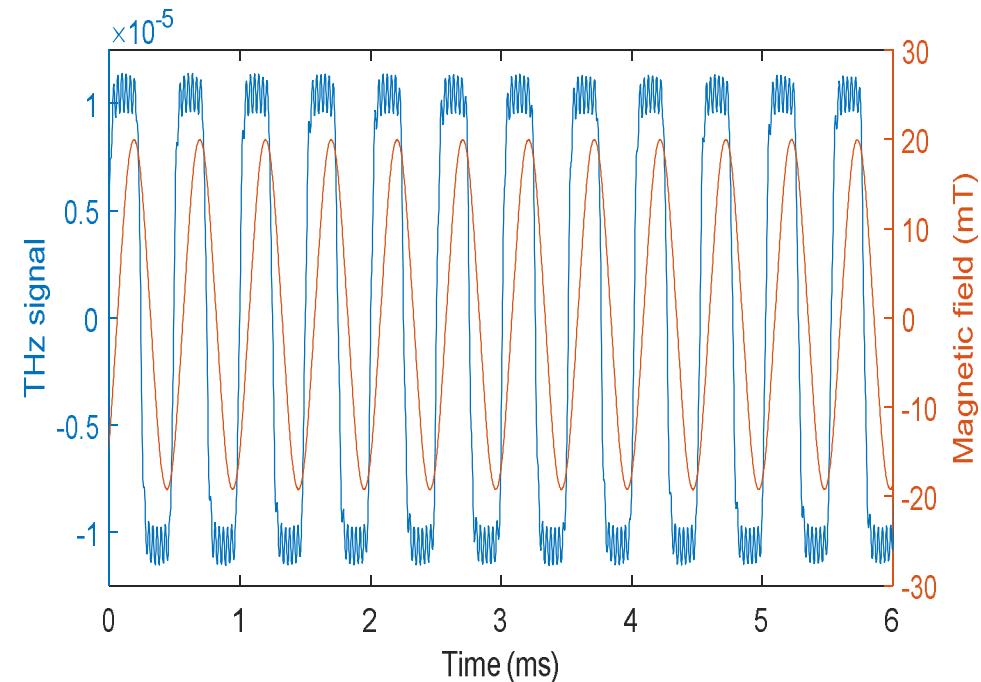
More features of spintronic THz emitters

Polarity modulation of THz field up to 20 kHz

Gückstock, Nadvornik, TK *et al.* (2020)

Bendable substrates

Wu, Yang *et al.*, Adv. Mat. (2016)

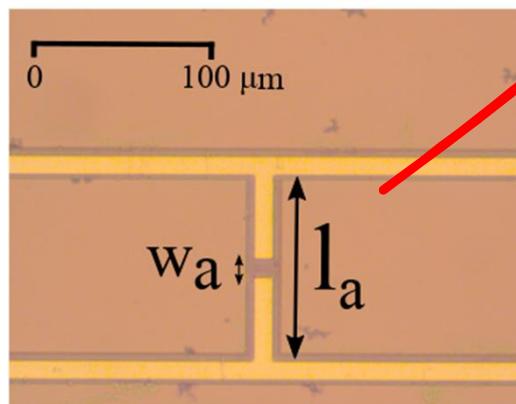
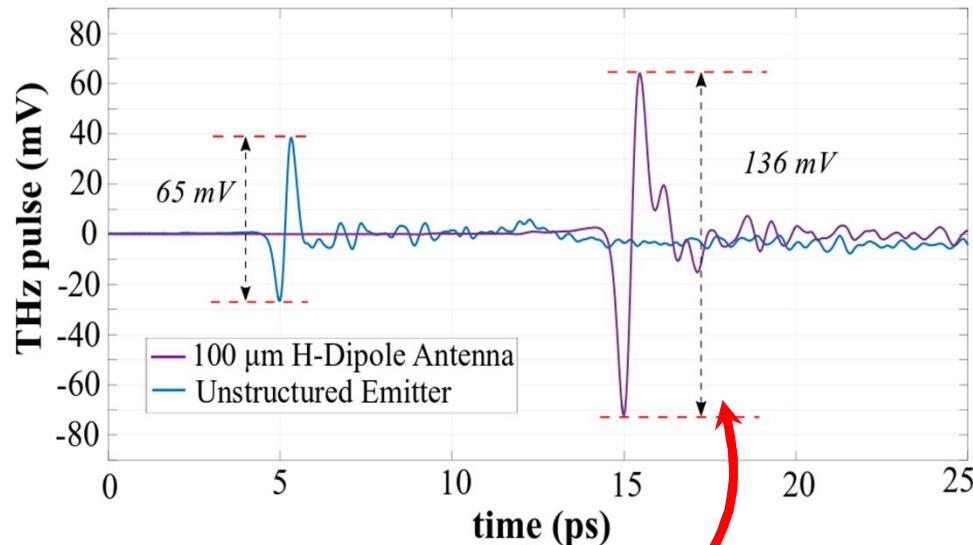


Metal films can be nanostructured easily

More features by nanostructuring

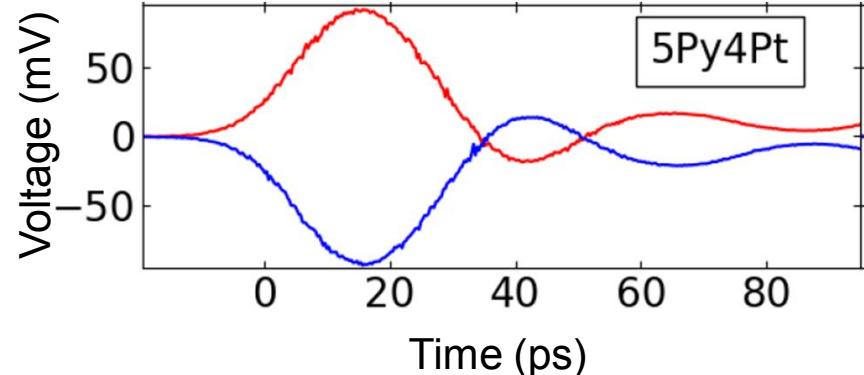
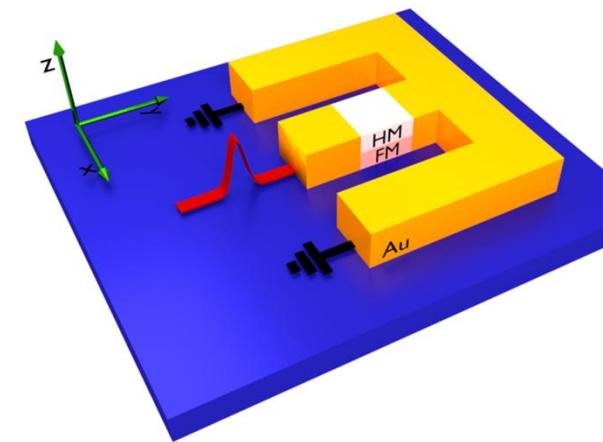
Antenna-coupled spintronic emitters

Nandi, Kläui, TK, Preu *et al.*, APL (2019)



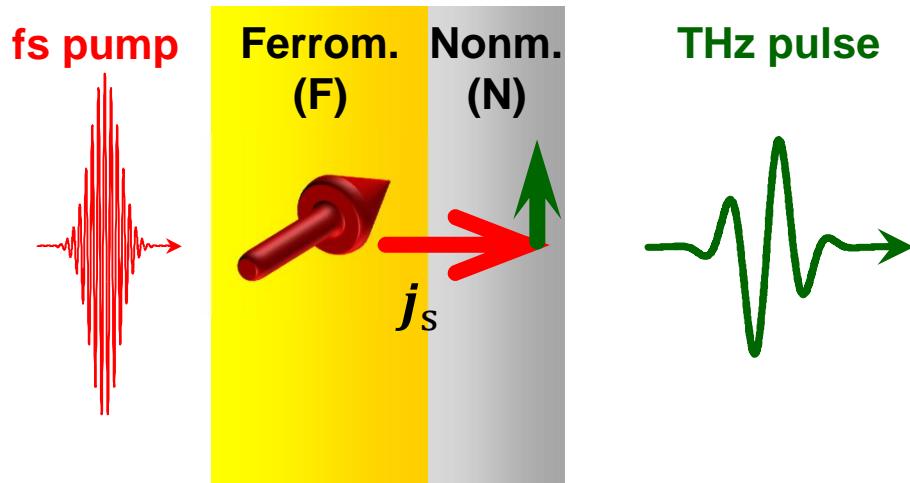
On-chip THz source

Weber,
Hoppe, TK,
Woltersdorf
et al. (2020)



Challenge:
How to further increase the THz field strength?

Challenge: More THz amplitude



Maximize spin-to-charge-current conversion

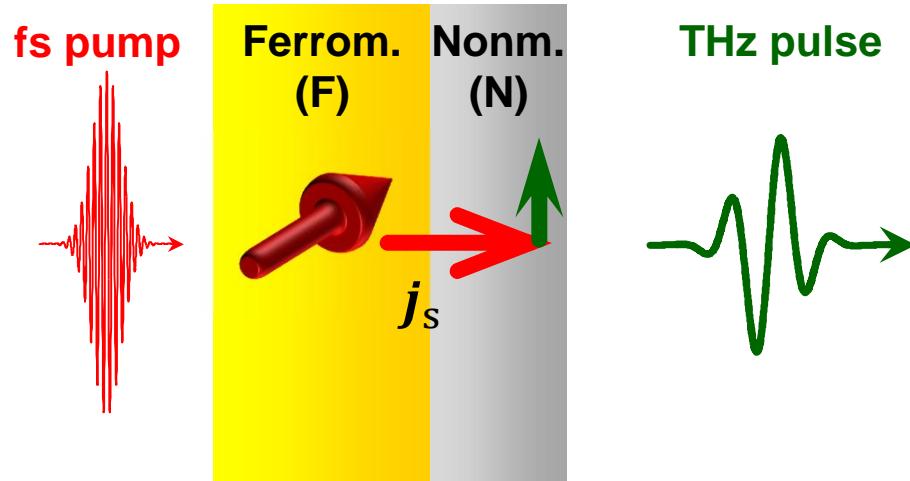
- N should have large ISHE and large spin diffusion length
- Interface engineering

Maximize spin current for given pump energy

- Understand how the spin current is induced

What does the driving THz spin current look like?

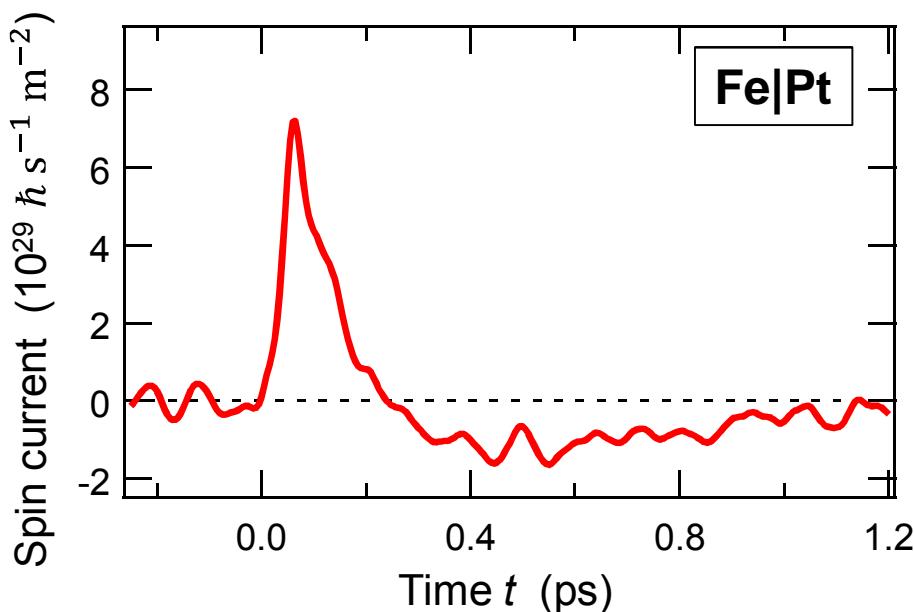
Dynamics of the spin current



Approach: “Calculate back”

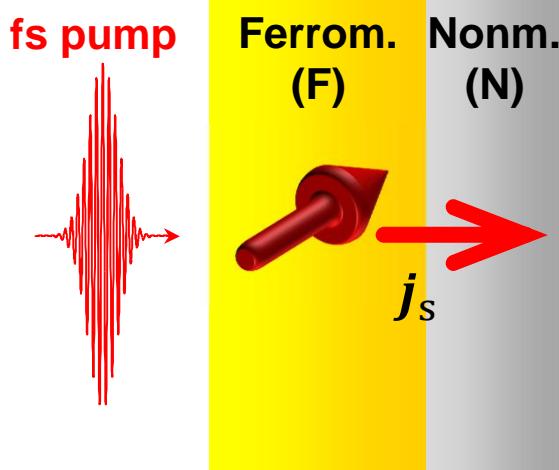
THz field $\rightarrow j_s(t)$

Seifert, Barker, TK et al.,
Nature Commun. (2018)



- Extremely fast bipolar response
- What is the driving force?

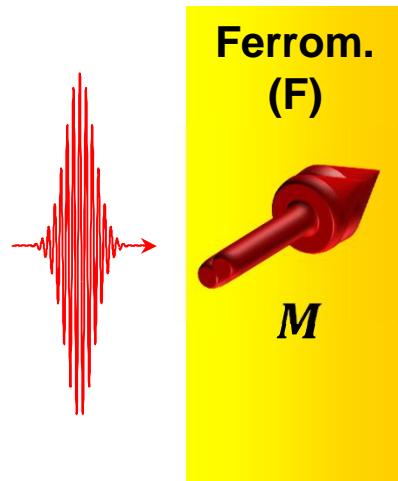
Transport vs demagnetization



What is the origin
of the ultrafast spin transport?

Idea:

Compare to dynamics of another
fundamental process



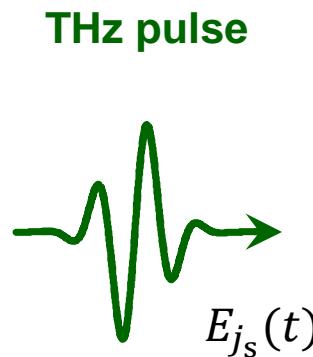
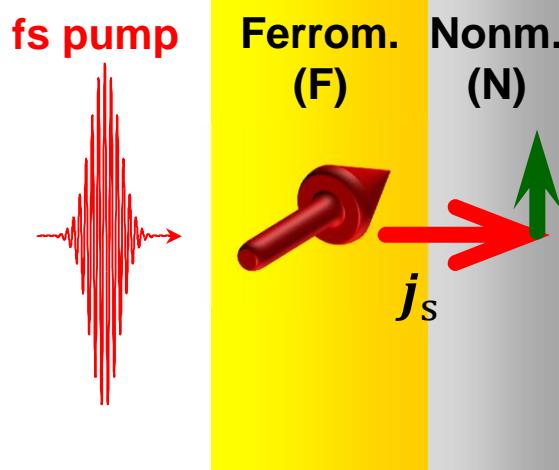
Ultrafast local magnetization quenching
of single ferromagnetic layers

Beaurepaire *et al.*, PRL (1996)

Is there a correlation with
ultrafast transport in bilayers?

Need to measure both processes in one experiment

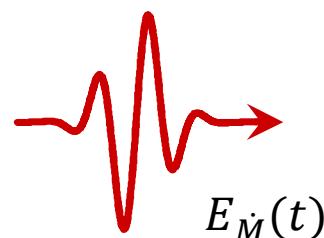
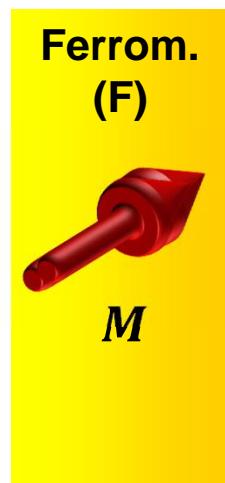
Transport vs demagnetization



Spin-to-charge current conversion
⇒ Electric-dipole (ED) radiation

Seifert *et al.*, Nat. Phot. (2016)

$$E_{j_s} \propto j_s$$



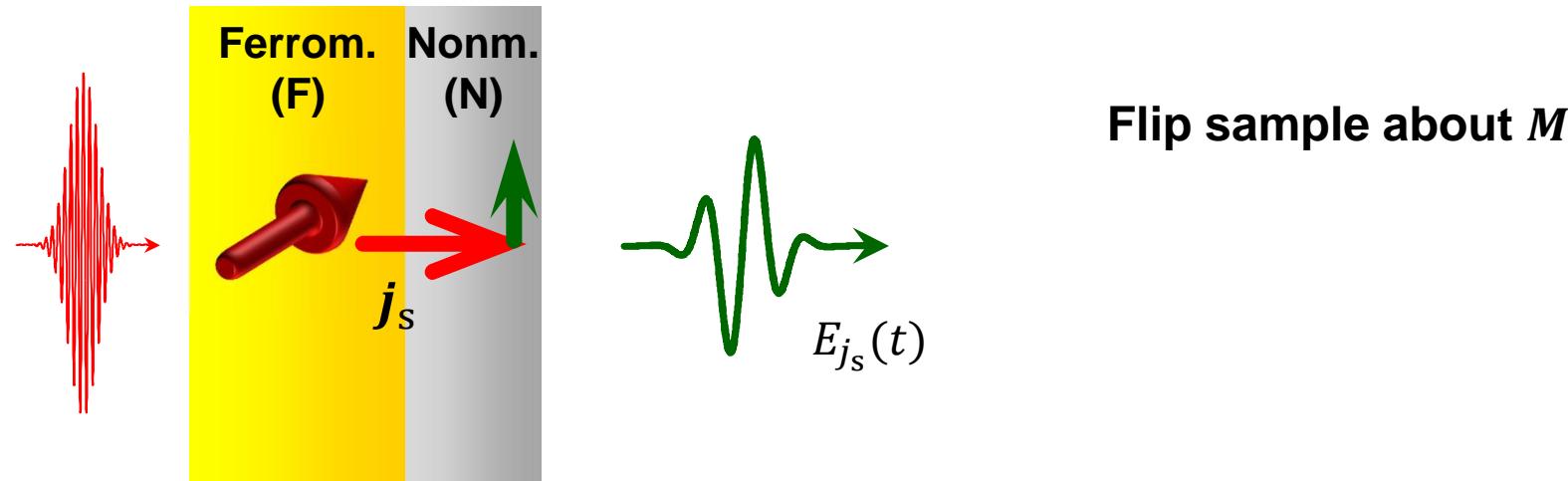
Ultrafast magnetization quenching
⇒ Magnetic-dipole (MD) radiation

Beaurepaire *et al.*, APL (2004)

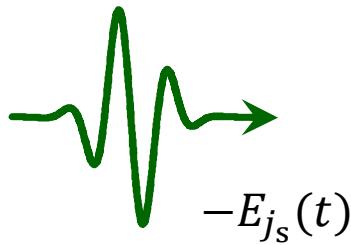
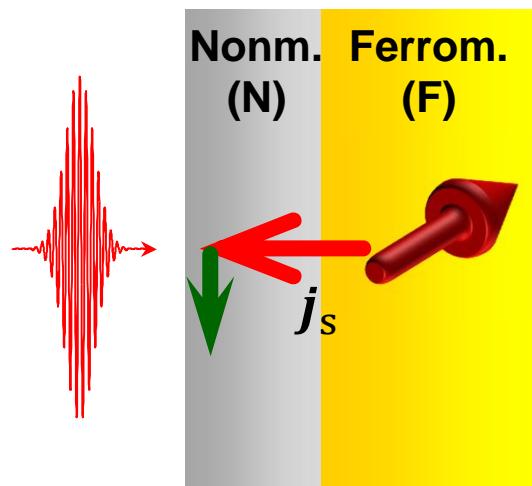
$$E_M \propto \dot{M} = \partial M / \partial t$$

How to isolate the weak \dot{M} signal?

j_s vs \dot{M} : Exploit symmetry



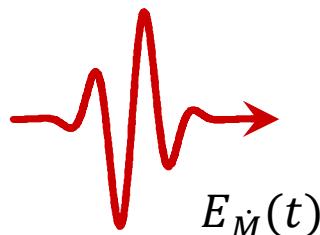
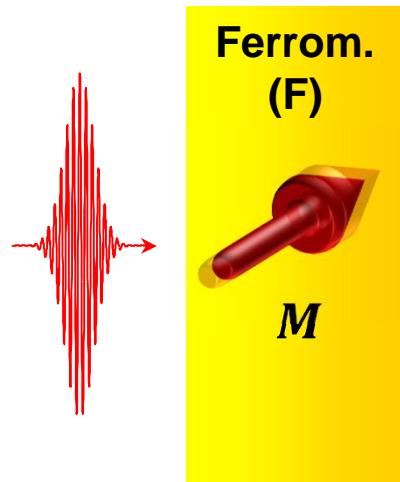
j_s vs \dot{M} : Exploit symmetry



Flip sample about M

- THz field due to j_s changes sign
- Has “odd” symmetry

Beaurepaire *et al.*, Ultrafast Phen. (2004)
Huisman, Kimel *et al.*, PRB (2015)

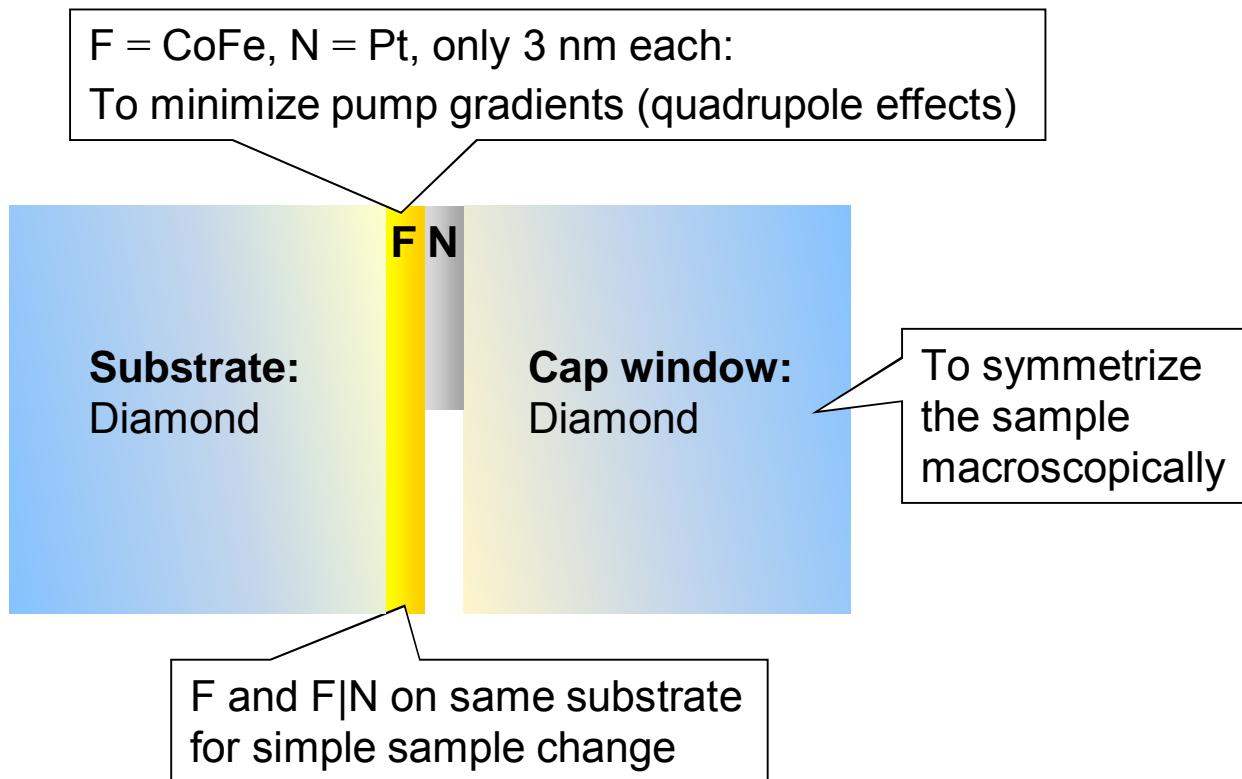


Flip sample about M

- \dot{M} emission does not change
- Has “even” symmetry
- A potential even signal is hopefully small

Direct way to measure j_s and \dot{M} in same setup with ~10 fs resolution

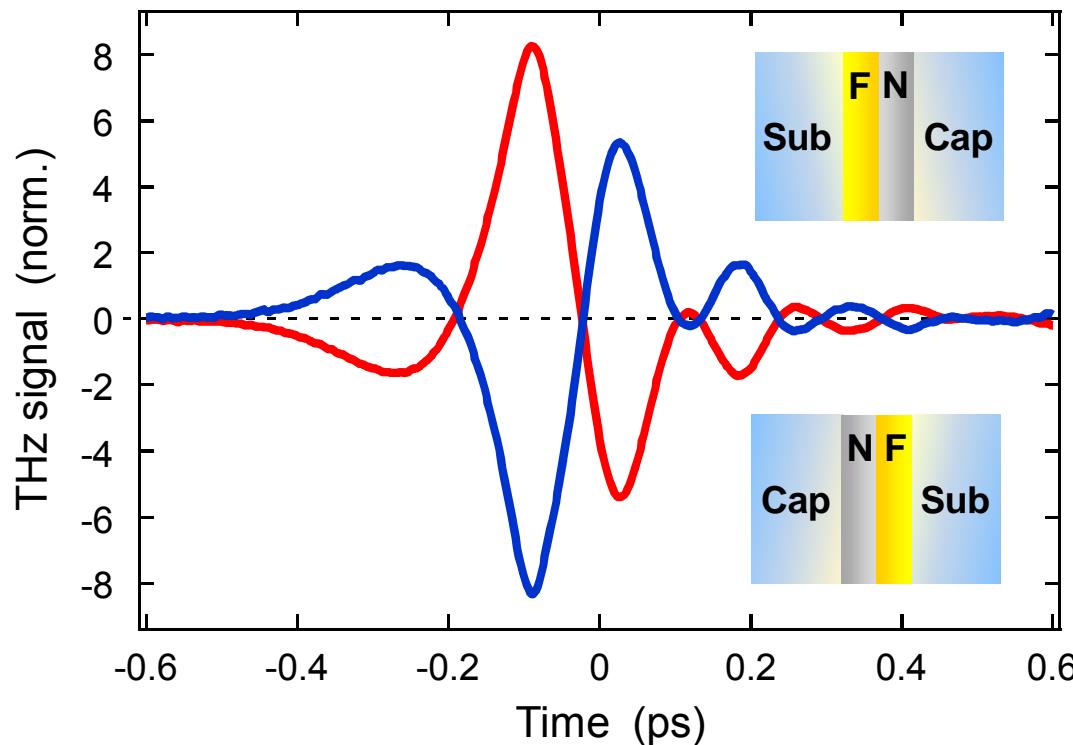
Sample details



Pump pulses: From Ti:sapphire oscillator (10 fs, 800 nm, 1 nJ, 80 MHz)

Start with F|N stack

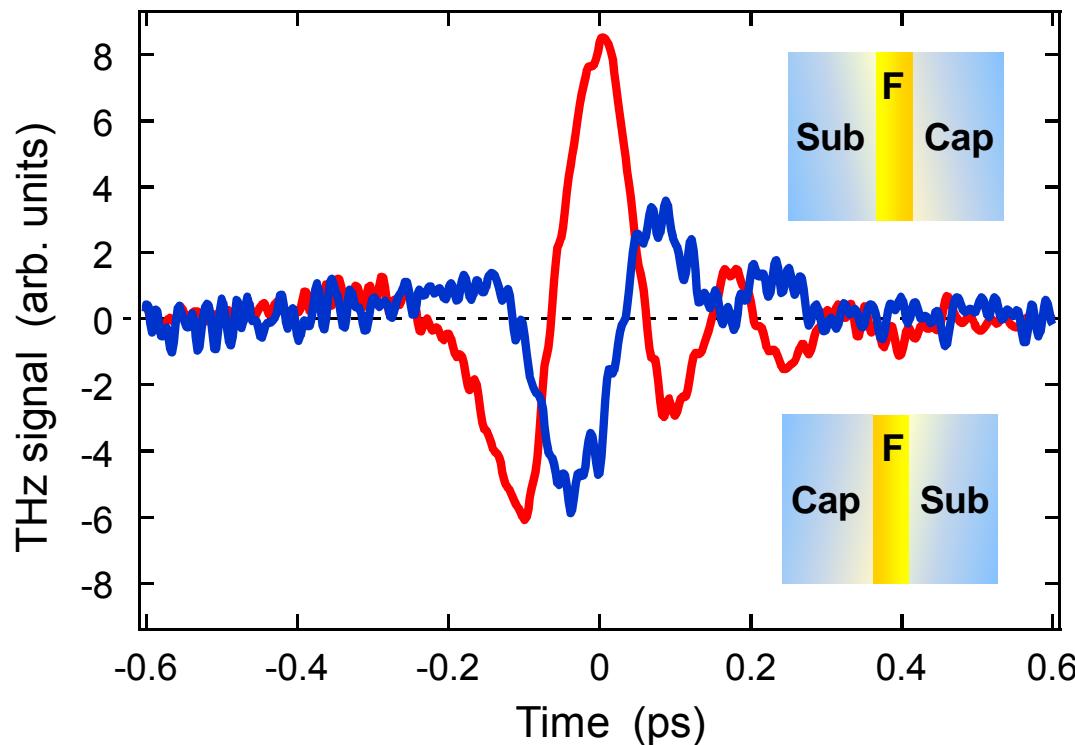
THz emission from F|N stack



- $\approx 100\%$ signal reversal \Rightarrow Even symmetry
- Thus: Emission $E_{js}(t) \propto j_s(t)$ is dominant
- Use as reference for proper alignment

Let's go for the single F layer

THz emission from single F layer



- Signal is ~ 170 times smaller than from F|N
- Nontrivial signal change upon sample reversal
- Indicates both even (\dot{M}) and odd signals contribute

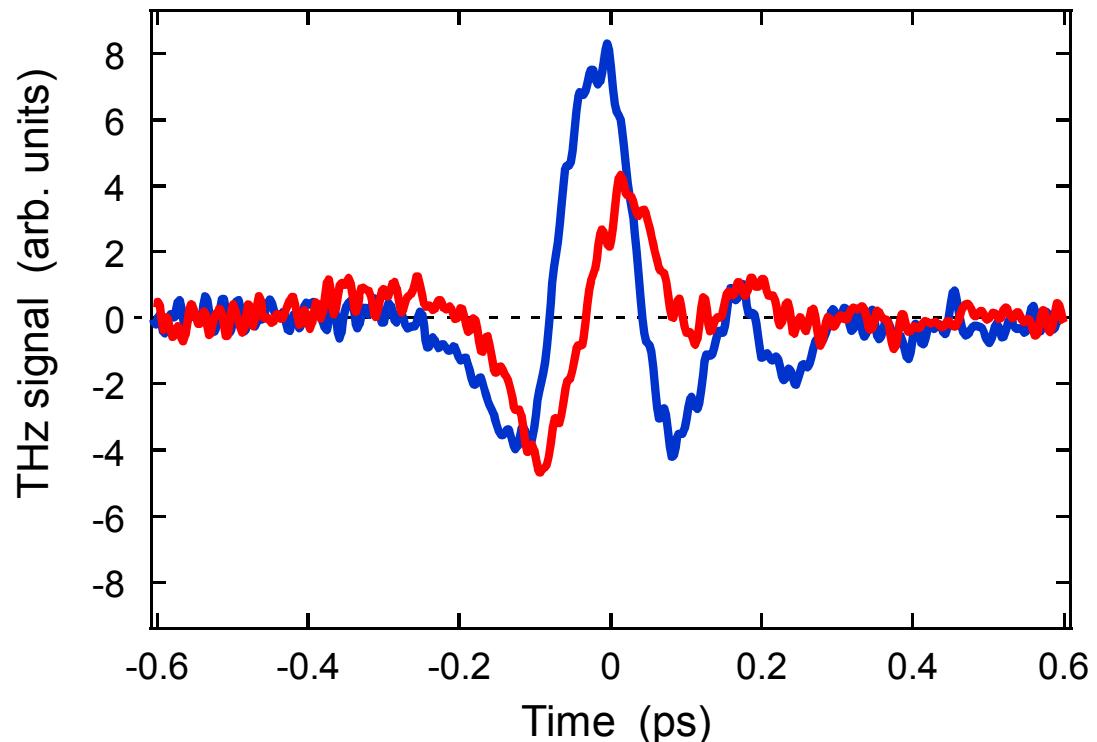
Separate even and odd contributions

F monolayer: Odd and even emission

$$\begin{array}{c} \text{S F C} \quad \text{C F S} \\ \hline 2 \\ \text{S F C} + \text{C F S} \\ \hline 2 \end{array}$$

= Odd

= Even



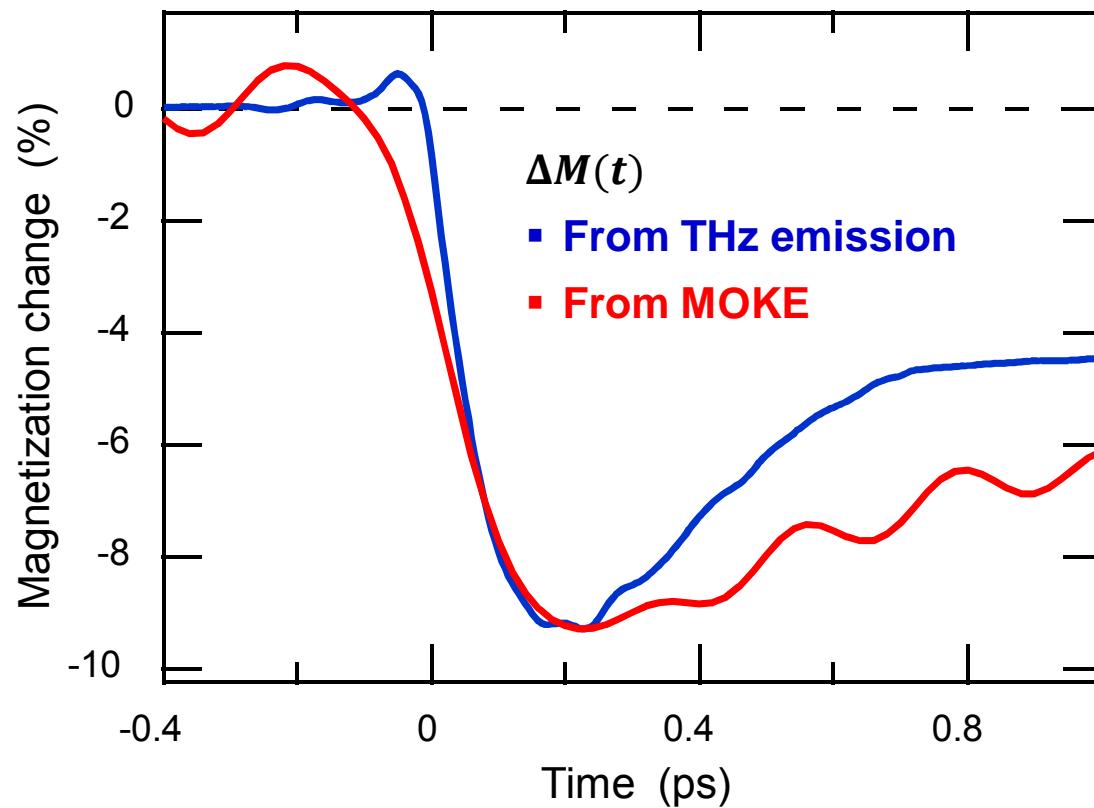
Odd signal: Considerable

- Origin not fully clear: Due likely to film inhomogeneities along growth direction

Even signal: We assign it to ultrafast demagnetization, $E_M \propto \dot{M}$

- Signal has right polarity and order of magnitude
- Extract $M(t)$ and check temporal dynamics

F monolayer: $M(t)$ dynamics



Approximate agreement of

- THz-extracted $M(t)$
- $M(t)$ from MOKE

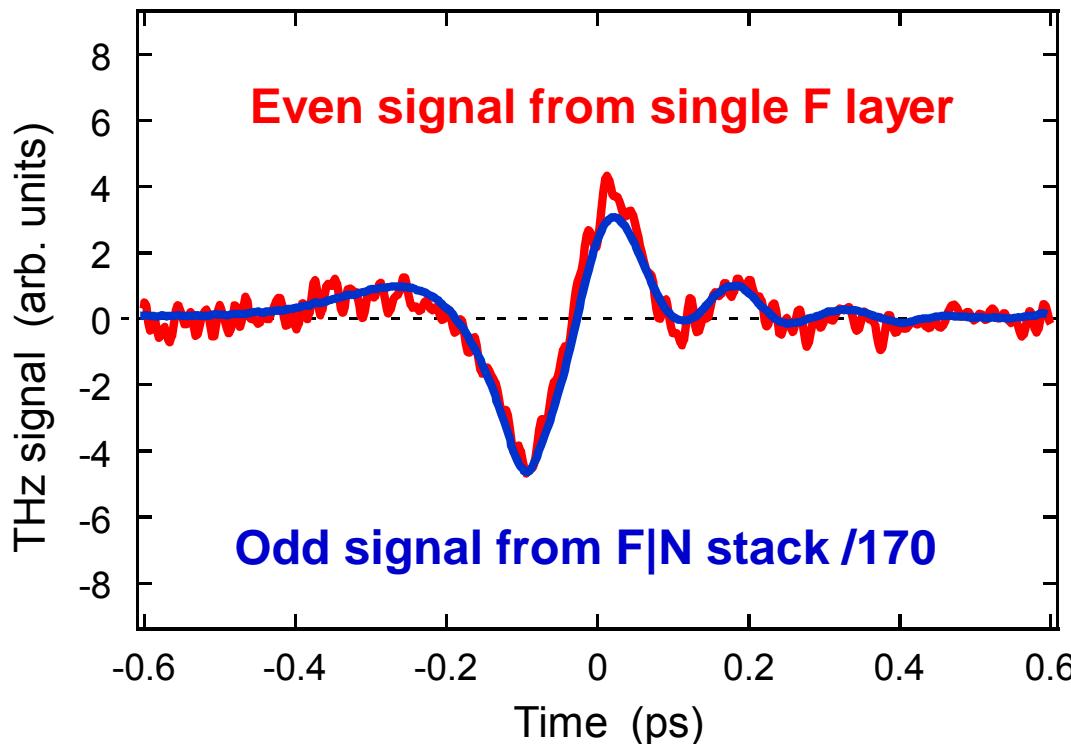
Huisman, Kimel *et al.*, PRB (2015)

Razdolski, Bovensiepen, Melnikov *et al.*,
J. Phys. D (2017)

The even signal from the F single layer arises from ultrafast demagnetization

Now compare: Even signal from F \leftrightarrow Odd signal from F|N

Even from F vs odd from F|N



Rouzegar, Brandt,
Reiß, Brouwer,
Woltersdorf, TK
et al. (2020)

Even emission from single F layer \propto Odd emission from F|N stack

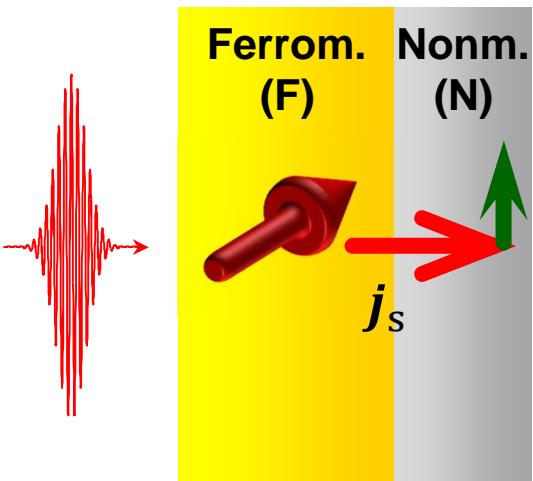


$$E_{\text{even}}^{\text{F}}(t) \propto E_{\text{odd}}^{\text{F|N}}(t)$$

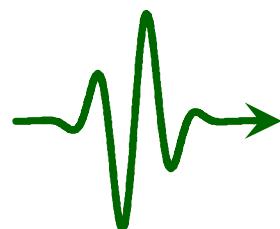


What does this imply?

Transport vs demagnetization



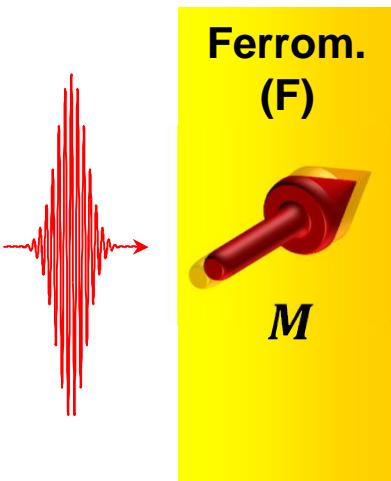
$$E_{\text{odd}}^{F|N}(t) \propto j_s^{F|N}$$



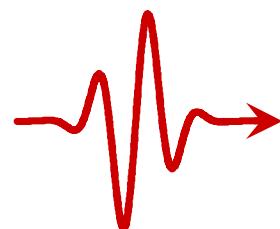
Our observations indicate:

Demagnetization rate of F
 \propto Spin current F \rightarrow N in F|N

$$j_s^{F|N} \propto \dot{M}^F$$



$$E_{\text{even}}^F \propto \dot{M}^F$$



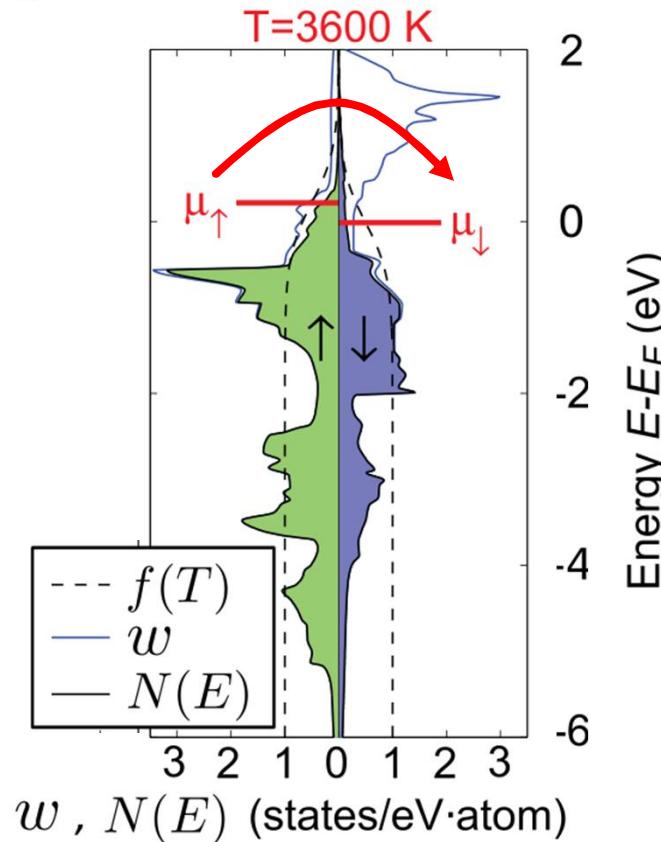
Confirmed for

- N = Pt and W
- F = Co₇₀Fe₃₀ and Ni₈₀Fe₂₀

1) Ultrafast spin transport and
2) Ultrafast demagnetization
are driven by same force

How to explain?

Interpretation: Driven by $\mu_{\uparrow} - \mu_{\downarrow}$



Acremann et al.:

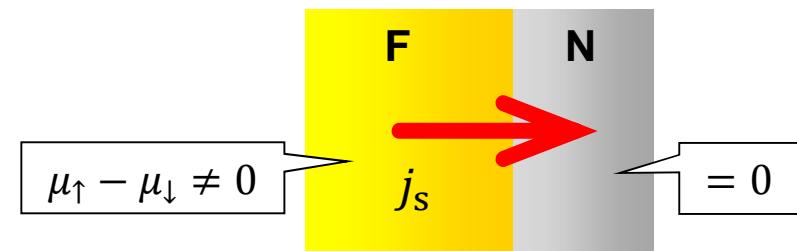
Pump-induced heating of F layer causes $\mu_{\uparrow} \neq \mu_{\downarrow}$

J. Phys. Cond. Mat. (2017)

Spin-dependent Seebeck effect

Bauer, Saitoh, Wees,
Nature Mat. (2013)

$$j_s^{F|N} \propto \mu_{\uparrow} - \mu_{\downarrow}$$



Rethfeld et al.:

$\mu_{\uparrow} - \mu_{\downarrow}$ drives spin flow
from \uparrow to \downarrow channel

PRB 90, 144420 (2014)

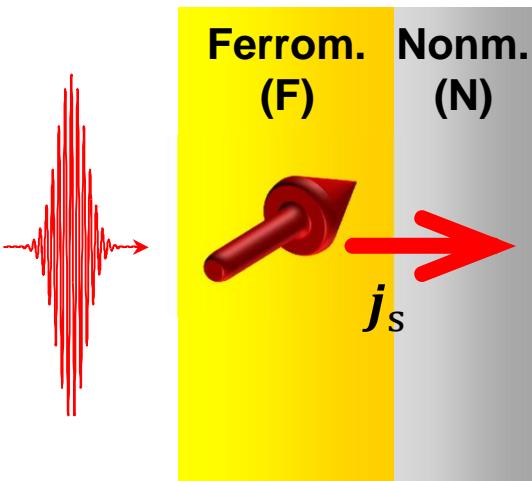
$$\dot{M}^F \propto \mu_{\uparrow} - \mu_{\downarrow}$$

Consistent with
our observations

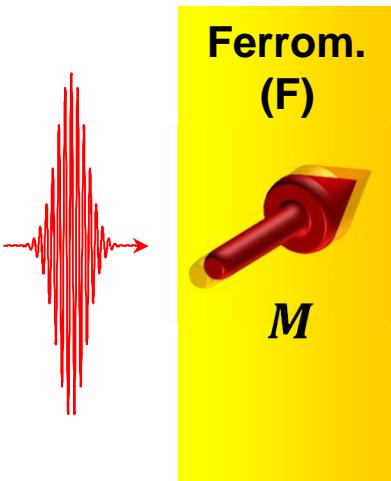
$$\dot{M}^F \propto j_s^{F|N}$$

Can be extended to nonthermal states

Summary



- 1) Ultrafast spin transport and
 - 2) Ultrafast demagnetization
- are driven by the same force**



Significance:

- We can apply all knowledge of 2) to 1)
- Will guide us to maximize spin current from F to N
- Can use N layer as an ultrafast detector of $\mu_\uparrow - \mu_\downarrow$

Summary

THz spin transport

- Enables interesting applications, e.g. THz wave generation, THz spin torque
- First temporal steps of spintronic effects

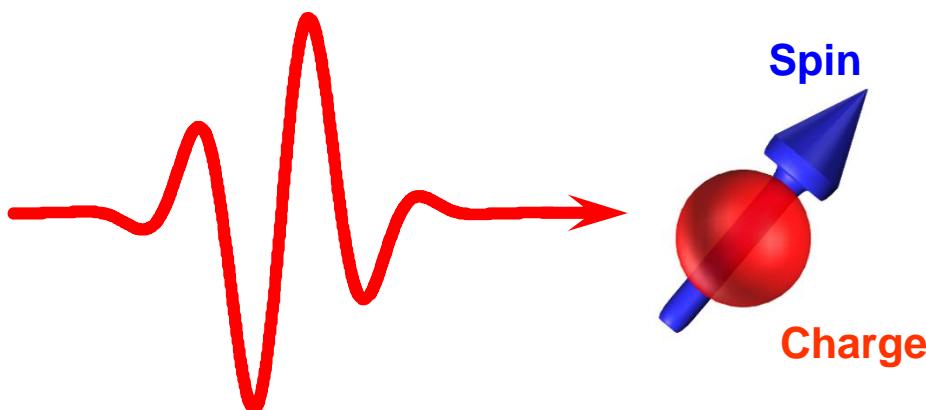
New insights

- Ultrafast demagnetization and spin transport are driven by same force

Future directions

- Use THz currents to apply spin torque

THz radiation is a useful tool to probe and control spin dynamics



Interested in a project in
in THz spintronics?

→ Contact Tobias Kampfrath