

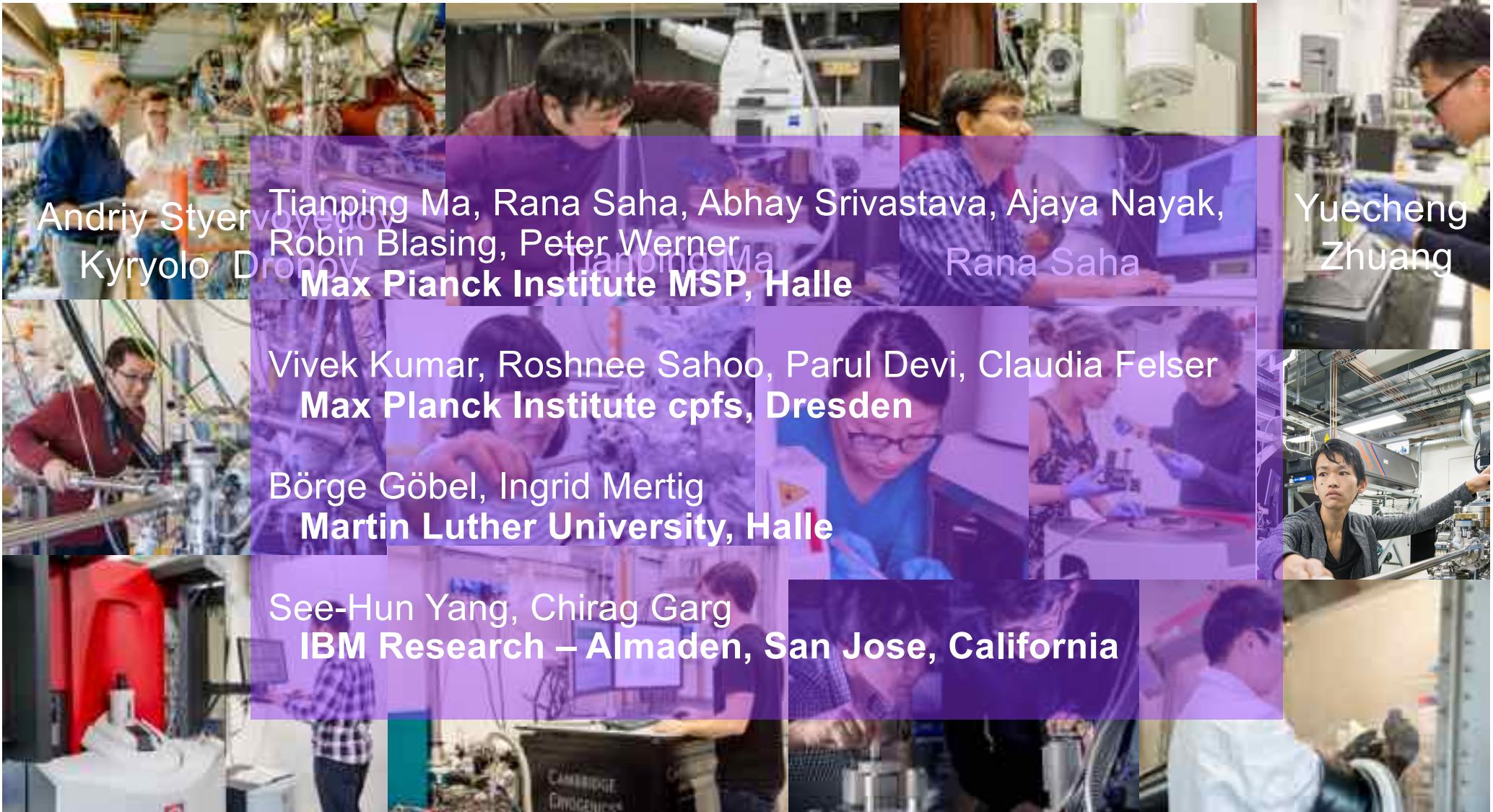
Chiral spintronics: non collinear spin textures with application to Racetrack Memory

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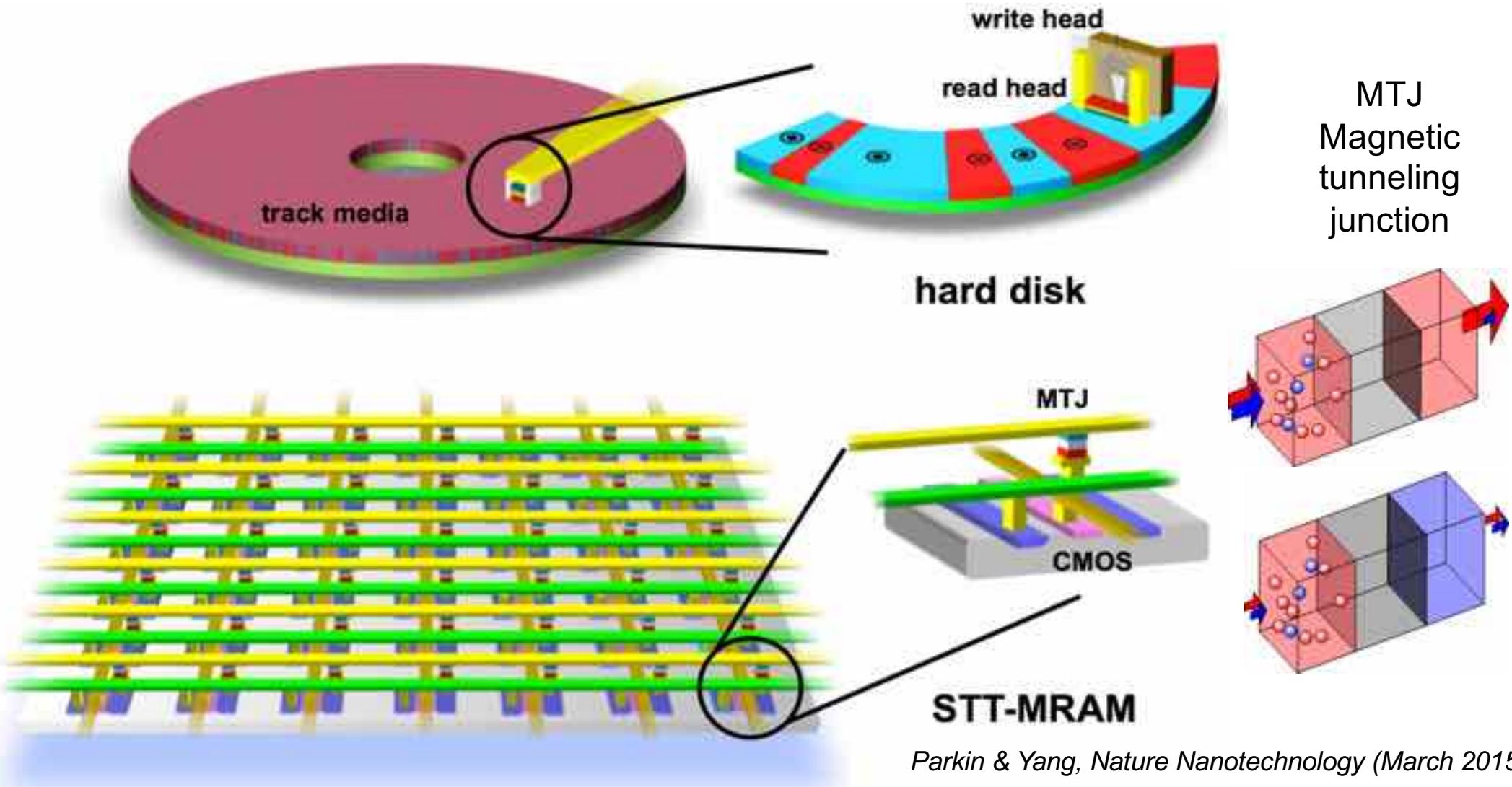
Yuecheng Zhuang

Vivek Kumar, Roshnee Sahoo, Parul Devi, Claudia Felser
Max Planck Institute cpfs, Dresden

Börge Göbel, Ingrid Mertig
Martin Luther University, Halle

See-Hun Yang, Chirag Garg
IBM Research – Almaden, San Jose, California

Spintronics – from materials and phenomena to applications



Spintronic technologies evolution

HDD Spin Valve & MTJ read heads

1988-1991

GMR
Spin Engineering
- Oscillatory interlayer coupling
- Co/Cu
- Interface spin-dependent coupling
- Dusting layers

1997

Spin-valve recording heads (mass-production)



2007

Magnetic Tunnel Junction recording heads (mass production)

2015

Slow-down of increased HDD capacity

MRAM



1995

MTJ based MRAM proposal

1996

Spin transfer torque (STT) switching

1999

1 kbit array
High MR
Low RA
 $0.25 \mu\text{m}$
2.3 ns R/W
Giant TMR using MgO

2005

16 Mbit demonstrator (180nm CMOS)
Field switched

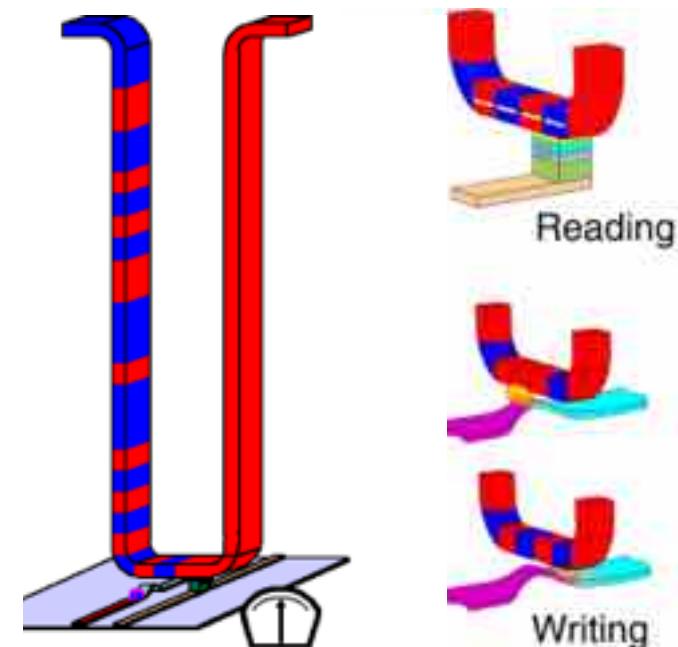
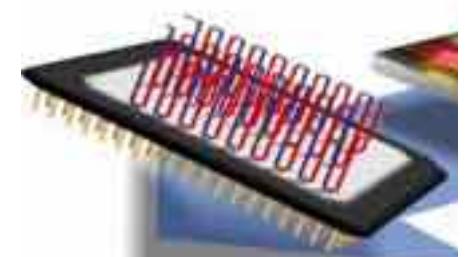
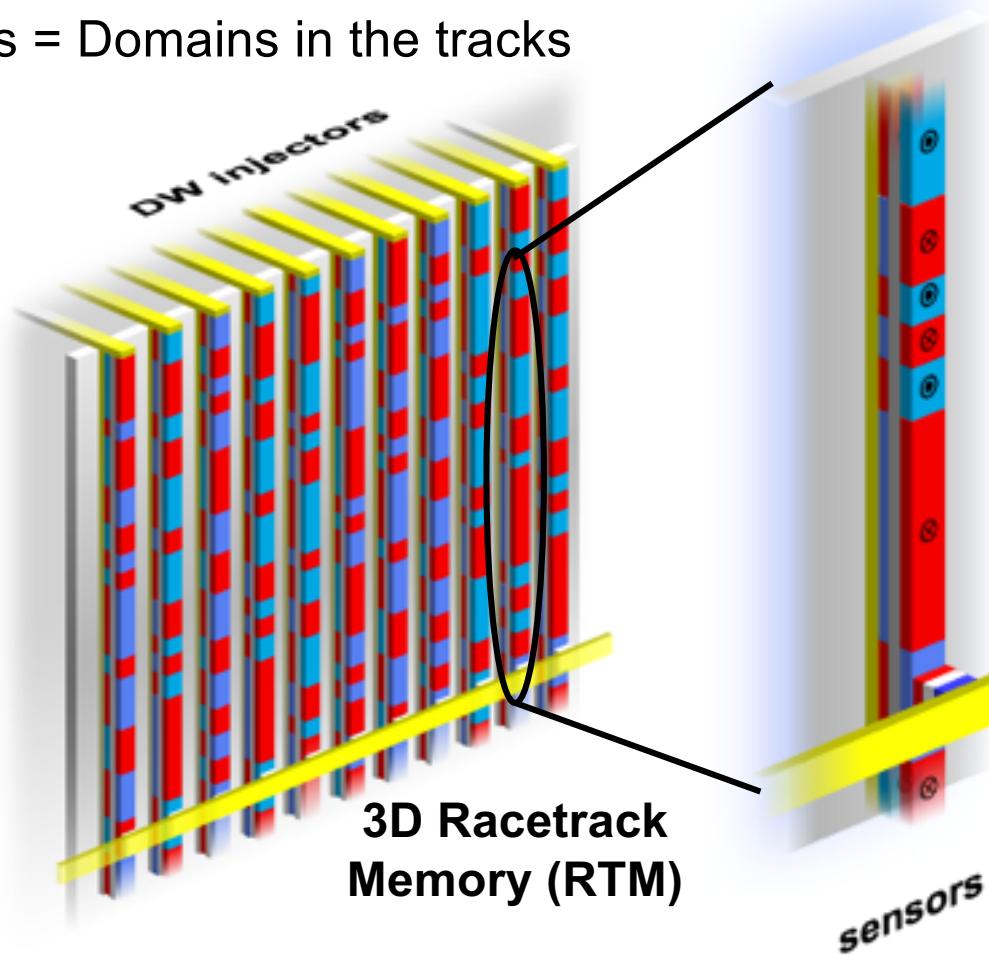


2019

Samsung announces shipment of high-performance foundry 28nm CMOS STT e-MRAM

Magnetic Racetrack Memory

- Bits = Domains in the tracks

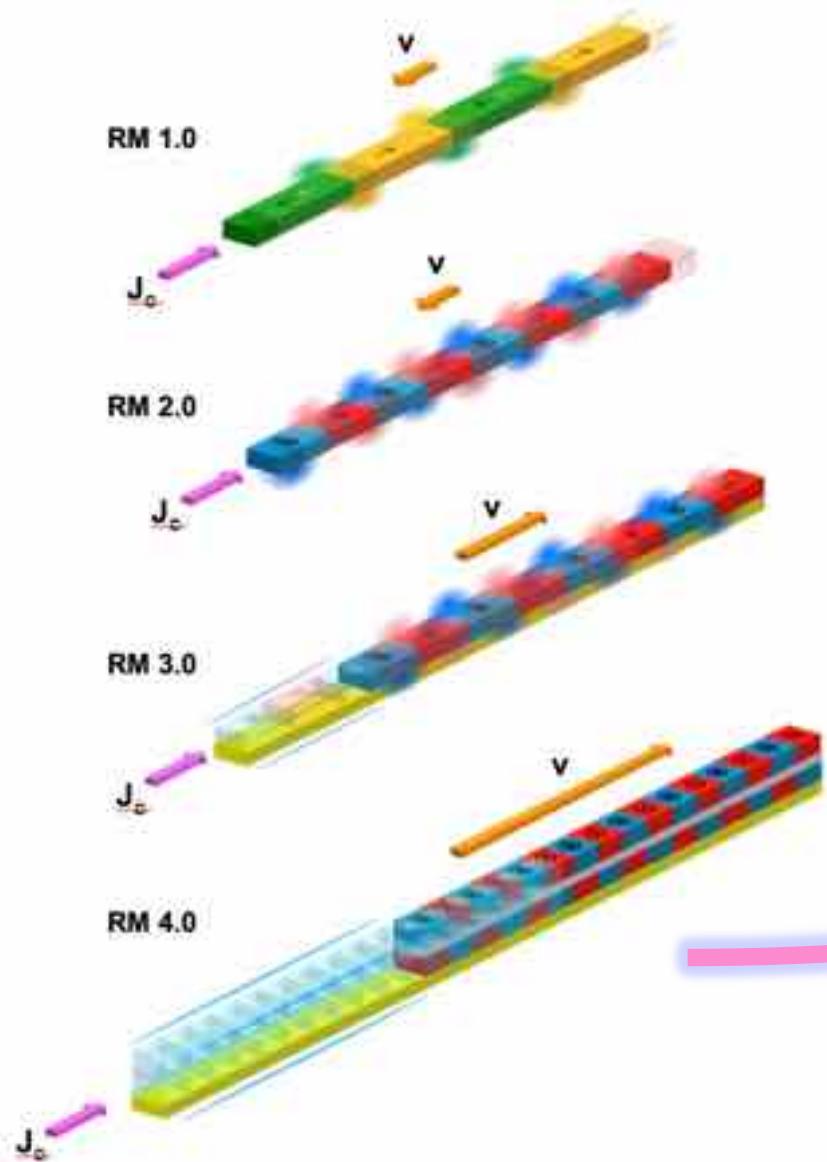


- A novel *three-dimensional* storage-class memory
- The capacity of a hard disk drive

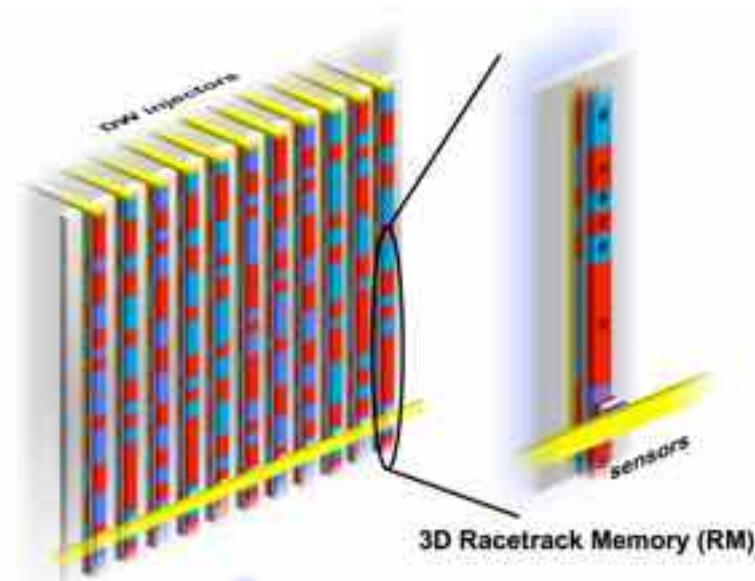
Parkin, US patents 6834005, 6898132 (2004)

Parkin et al., Science 320, 190 (2008)

Parkin, Scientific American (2009)



Memory on Racetrack! 4+ stages



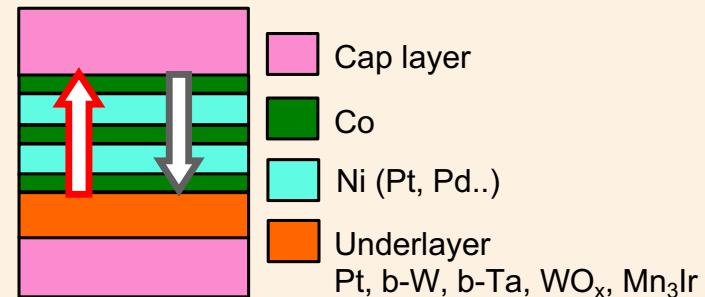
- Yang et al. *Nature Phys.* (2019)
Bläsing et al. *Nature Commun.* (2018)
Garg et al. *Nature Commun.* (2018)
Filippou et al. *Nature Commun.* (2018)
Garg et al. *Science Adv.* (2017)
Parkin & Yang, *Nature Nano.* (2015)
Yang, Ryu and Parkin, *Nature Nano.* (2015)
Ryu et al. *Nature Nano.* (2013)
Parkin et al. *Science* (2008)

Chiral domain walls from Dzyaloshinskii-Moriya Interaction

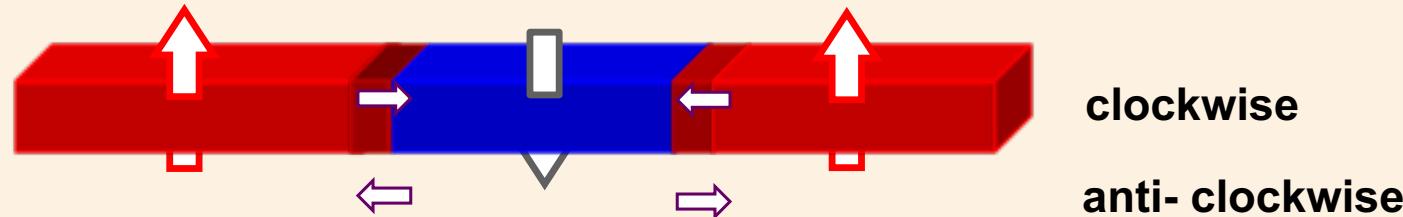
(#1) Perpendicular Magnetic anisotropy (PMA)

- materials: Co/Ni, Co/Pt, Co/Pd, RE/TM multilayers, low symmetry magnetic materials
- racetracks magnetized perpendicular to the plane of the racetrack

- *Bloch* and *Neel* domain walls
→ narrow domain walls
- Dzyaloshinskii-Moriya Interaction
→ *chiral* domain walls



Neel domain walls

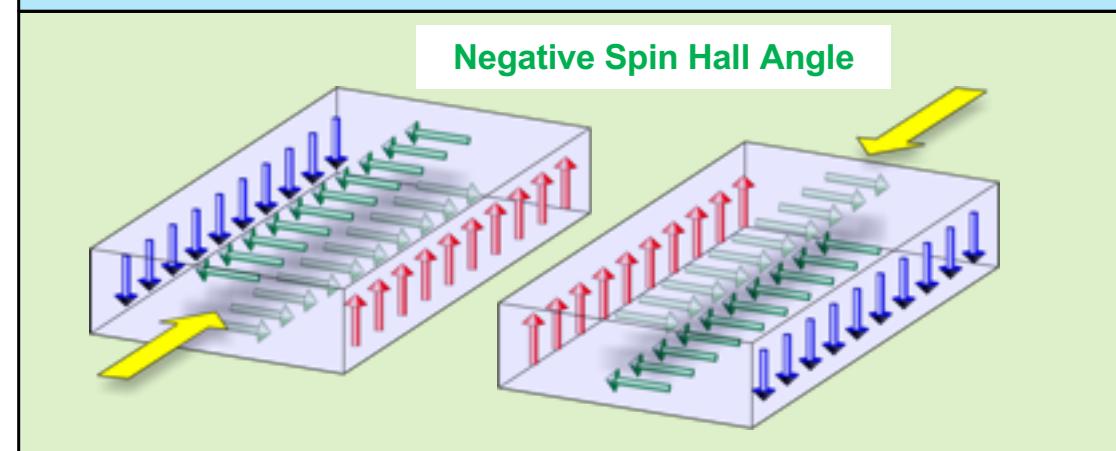
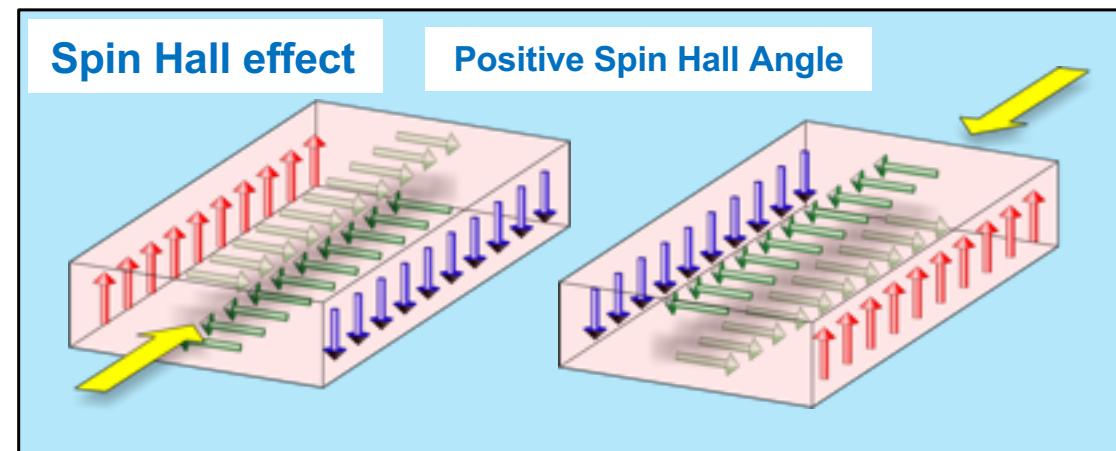
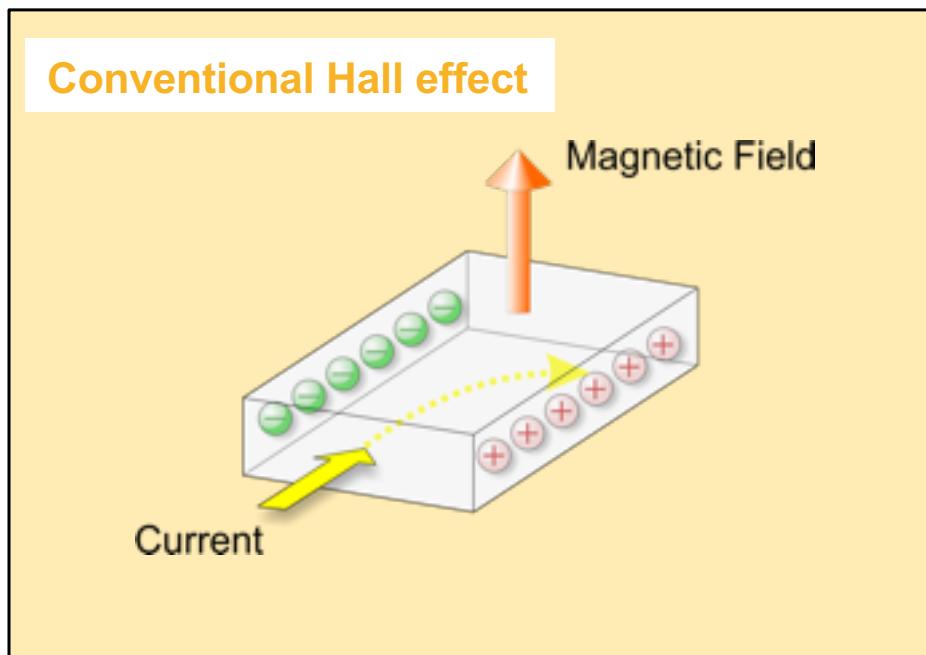


(#2) Dzyaloshinskii–Moriya interaction (DMI)

- *chirality* of domain walls set by interface(s) vector exchange interaction

(#3) Proximity Induced Magnetization (PIM)

(#4) Spin Hall effect: conversion of charge to spin current

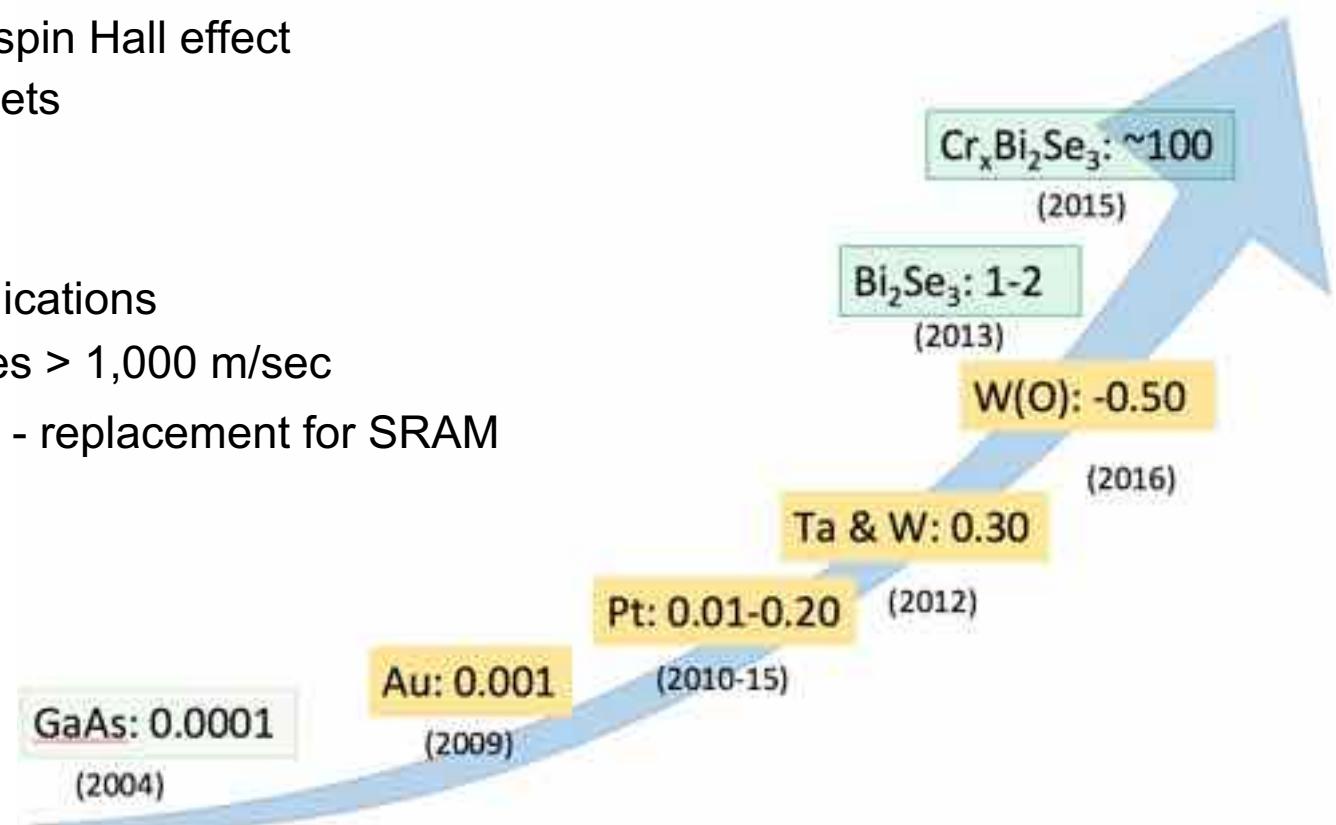


SHA: Spin Hall Angle
charge to spin conversion efficiency: ~10-30 %

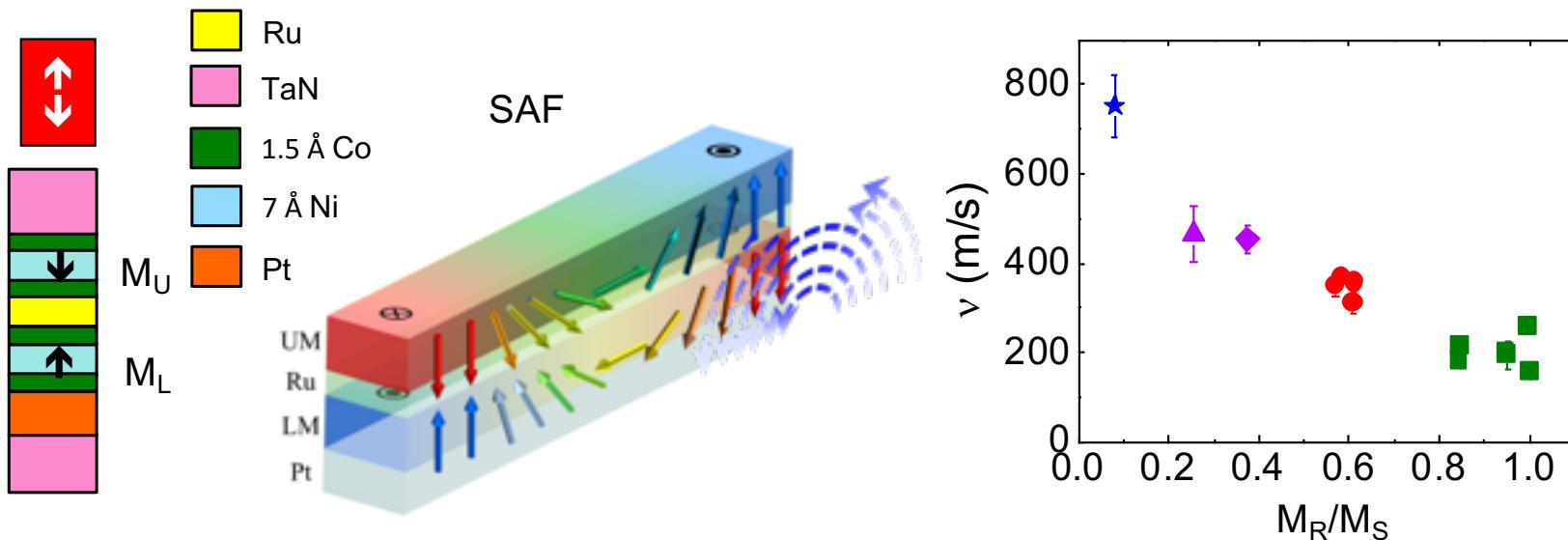
Charge to spin conversion via the Spin Hall Effect

- Spin Hall effect: dramatic evolution over past 5 years
- Novel materials show large spin Hall effect
 - Triangular antiferromagnets
 - Novel crystalline phases
- Important technological applications
 - Racetrack – DW velocities > 1,000 m/sec
 - 3T Single DW Racetrack - replacement for SRAM

Zhang *et al.* Nat. Phys. (2015)
Demasius *et al.*, Nat. Comm. (2016)
Nayak *et al.* Sci. Adv. (2016)
Zhang *et al.* Sci. Adv. (2016)



RM 4.0: Very High DW Speeds in Synthetic Antiferromagnet (SAF) racetracks



- SAF: Synthetic antiferromagnet: upper racetrack = exact mirror image of lower
- DW velocity increases as degree of compensation of moments in upper and lower racetracks increased
- 4x highest DW velocity yet reported → Speeds exceed 1.5- 5 km/sec
- AF exchange field > DMI field

Yang et al. *Nature Nanotechnology* (2015)

Racetrack Memory 4.0



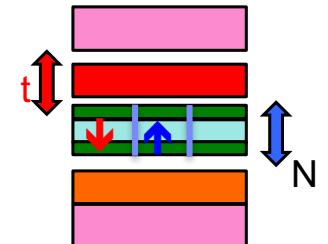
- 20 domain walls moved in lock step with current pulses
- High velocity at low current densities
- Narrow domain walls (~6 nm)
- Very thin racetracks (~1 nm)



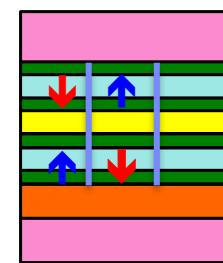
- Giant domain wall velocities in Synthetic Antiferromagnet racetracks > 1 km/sec

Domain wall motion: complex interplay of 4 spin-orbit derived phenomena

1. Perpendicular magnetic anisotropy
2. Proximity induced magnetization
3. Chiral domain walls – DMI
4. Spin currents – from spin Hall effect (SHE)

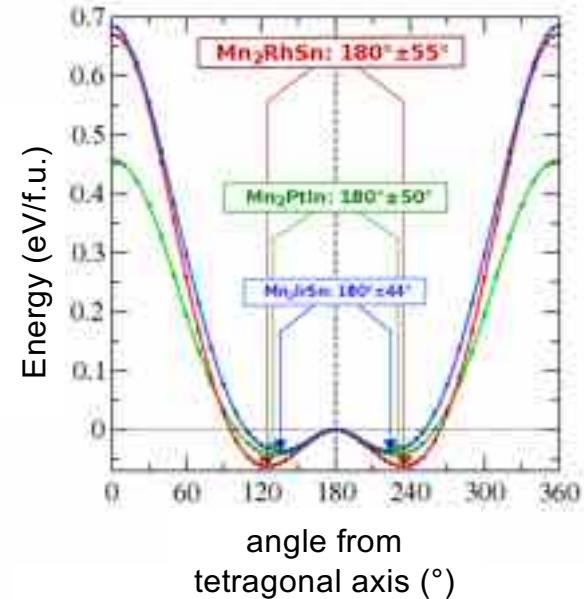
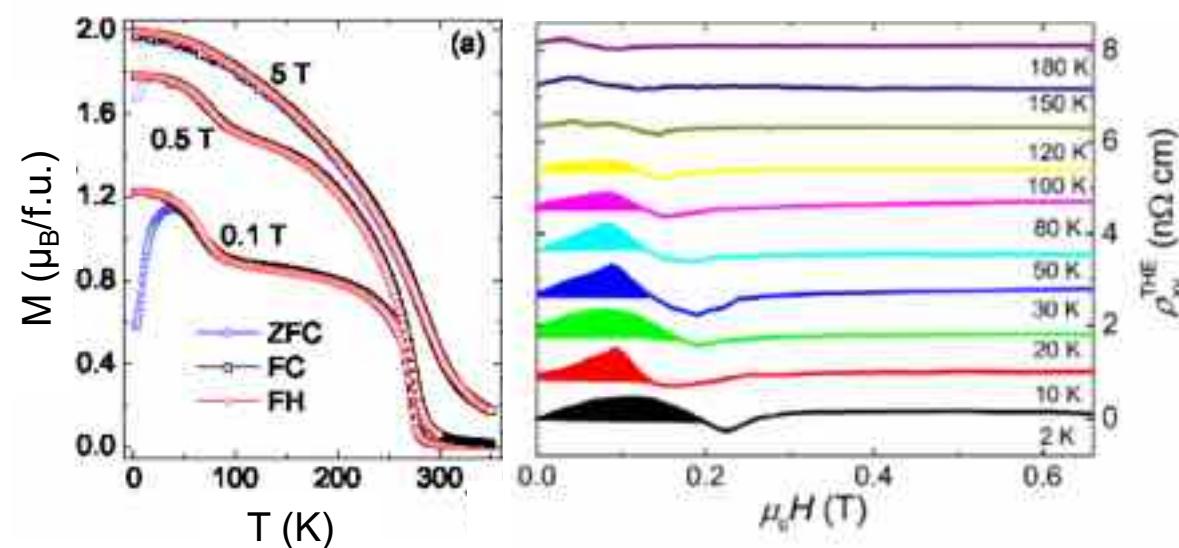
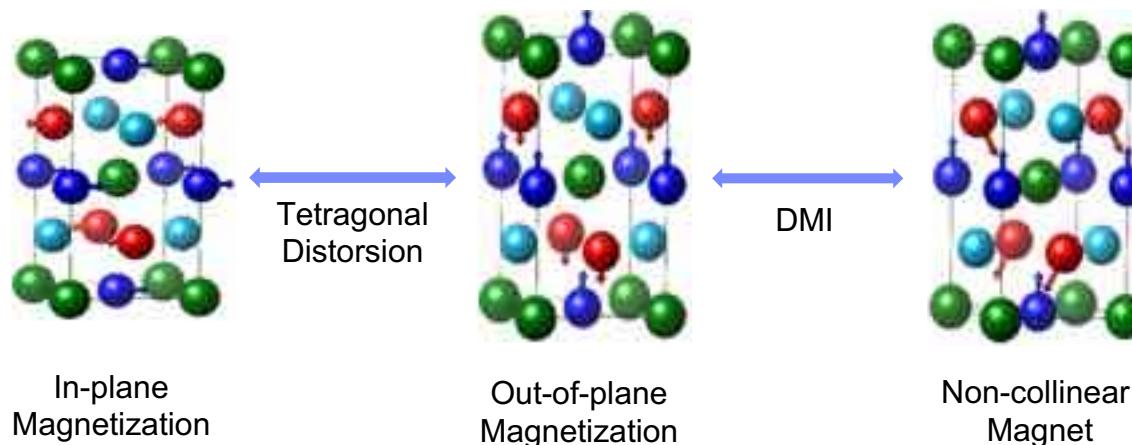


TaN
1.5 Å Co
7 Å Ni
Pt
Ru



Ryu et al. *Nature Nanotechnology* (2013)
Yang et al. *Nature Nanotechnology* (2015)
Parkin et al. *Nature Nanotechnology* (2015)
Garg et al. *Sci. Adv.* (2017); *Nat. Comm.* (2018)
Yang et al. *Nat. Phys.* (2019)

Heusler Family of compounds



Mn_2RhSn
Also Mn_2IrSn , **Mn_2PtSn**

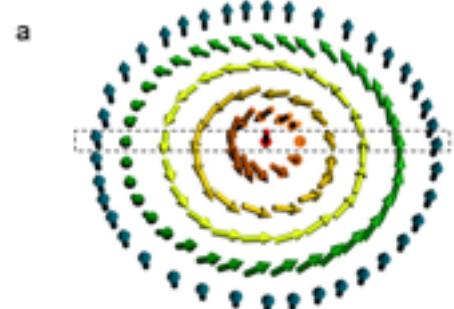
Claudia Felser et al.

Meshcheriakova, et al.
Phys. Rev. Lett. **113**, 087203 (2014)

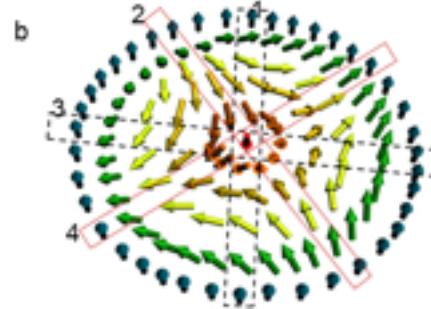
Rana et al. New J. Phys. **18**, 85007 (2016)

Skyrmions and anti-skyrmions

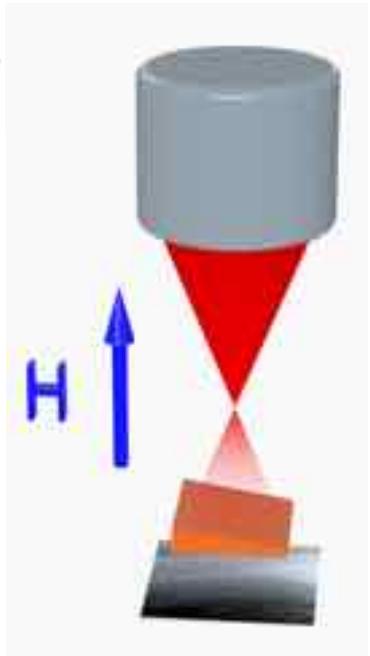
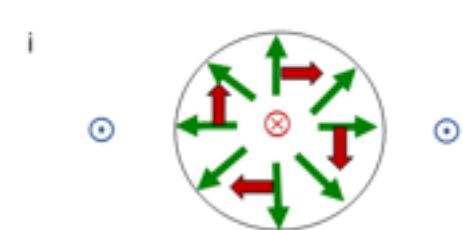
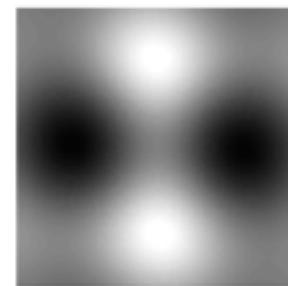
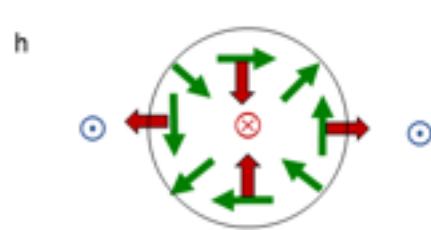
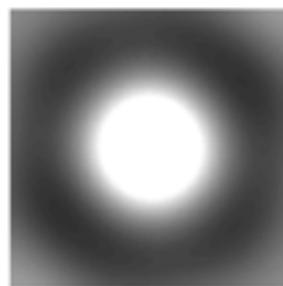
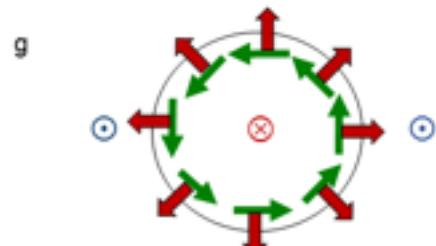
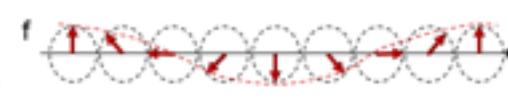
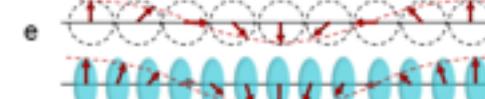
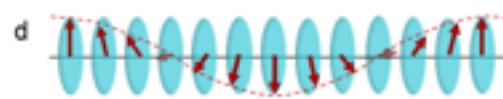
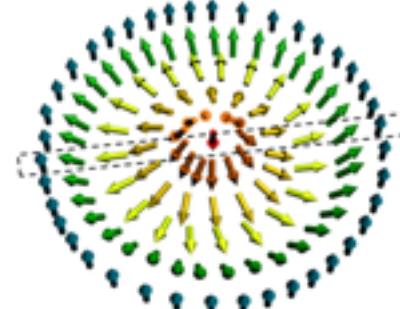
Bloch skyrmion



antiskyrmion



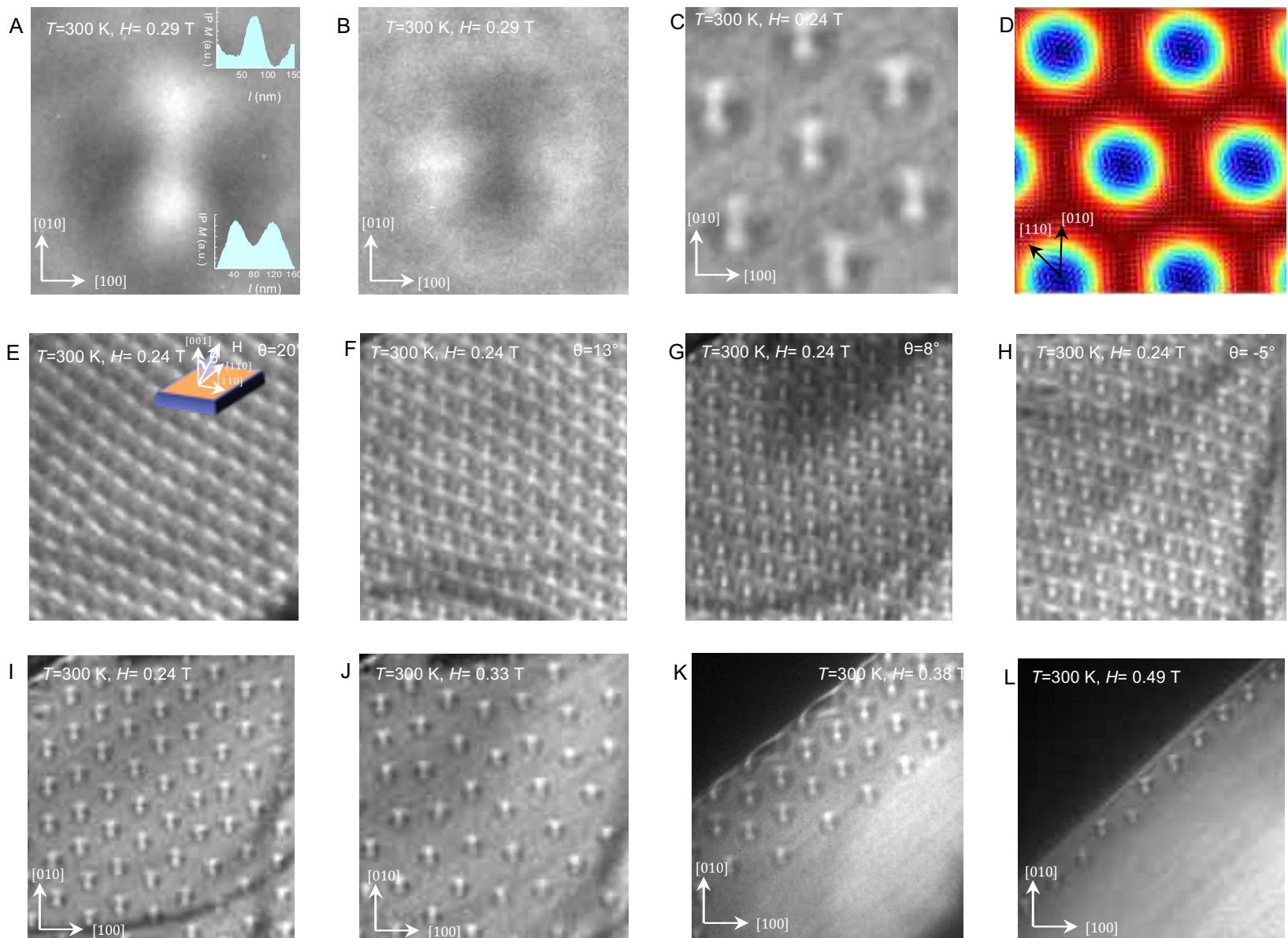
Néel skyrmion



$\text{Mn}_{1.4}\text{Pt}(\text{Pd})\text{Sn}$

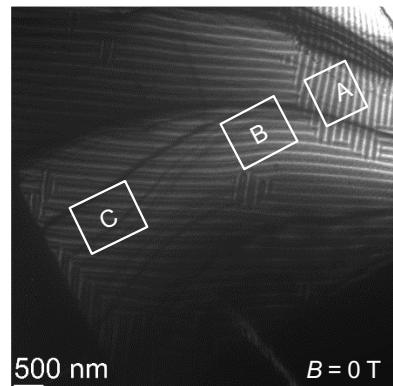
Anti-Skyrmions

- stable above 300 K



Nayak *et al.*
Nature (2017)

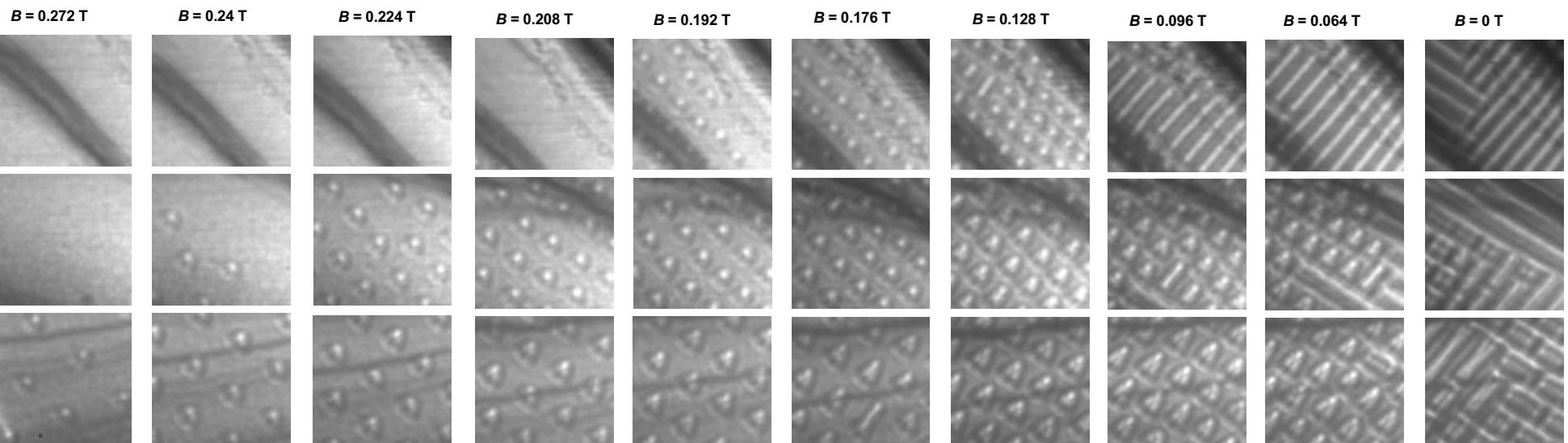
Stabilization of anti-skyrmion lattice in wedge-shaped lamella of $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$



C : $t \sim 247 \text{ nm}$

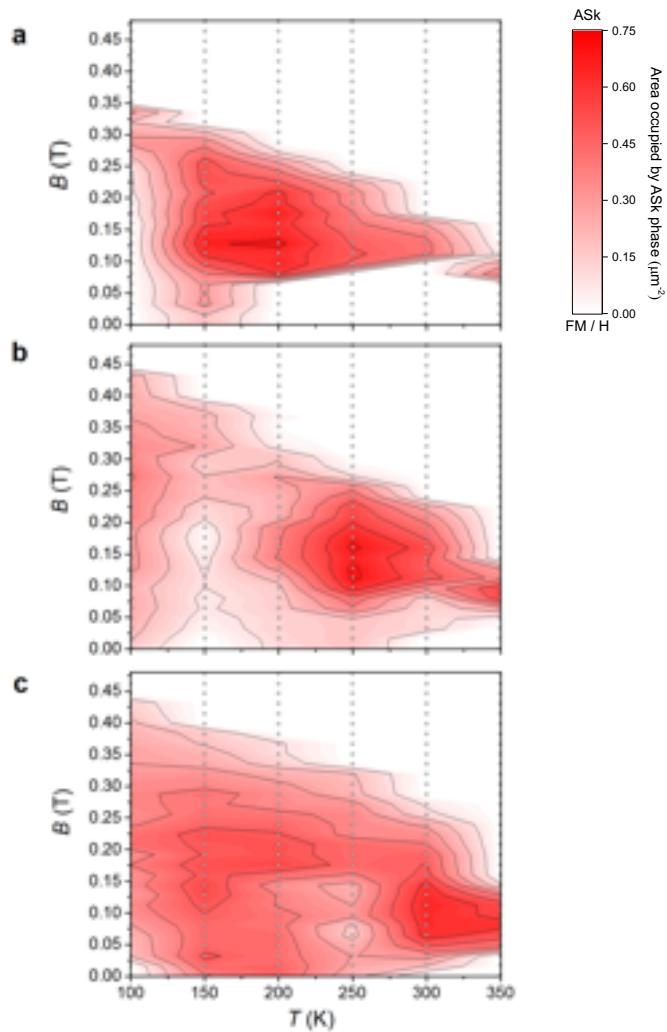
B : $205 \text{ nm} < t < 225 \text{ nm}$

A : $160 \text{ nm} < t < 185 \text{ nm}$

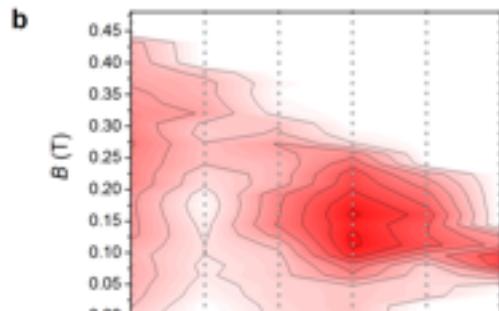


$\text{Mn}_{1.4}(\text{Pt}_{0.9}\text{Pd}_{0.1})\text{Sn}$

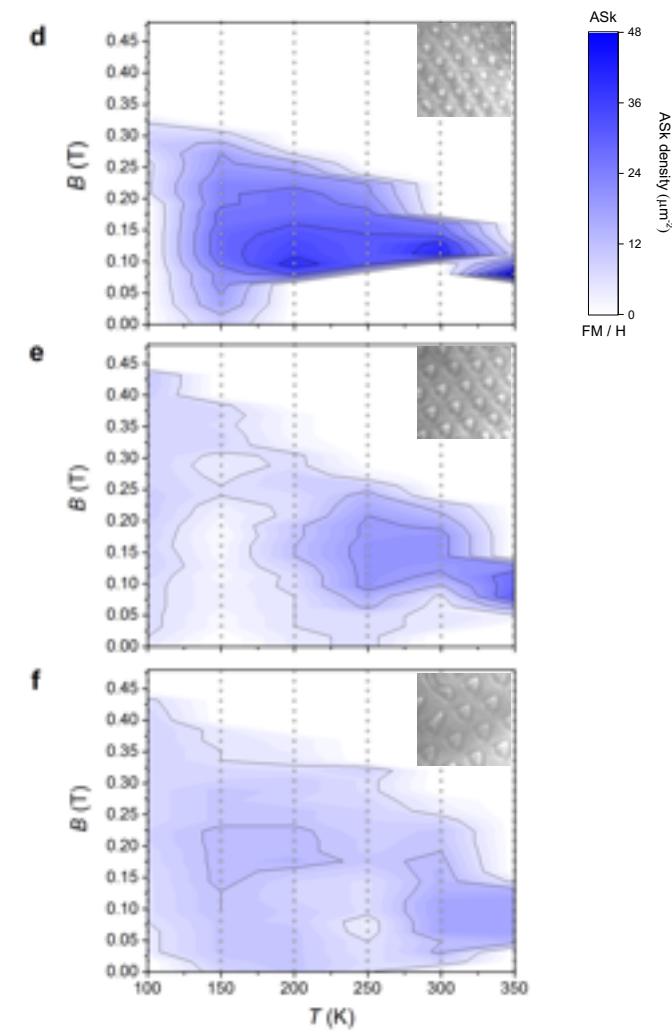
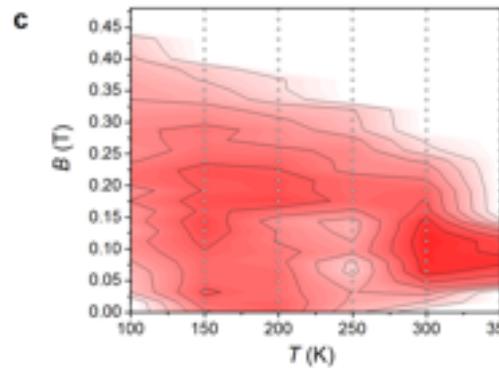
A : $160 \text{ nm} < t < 185 \text{ nm}$



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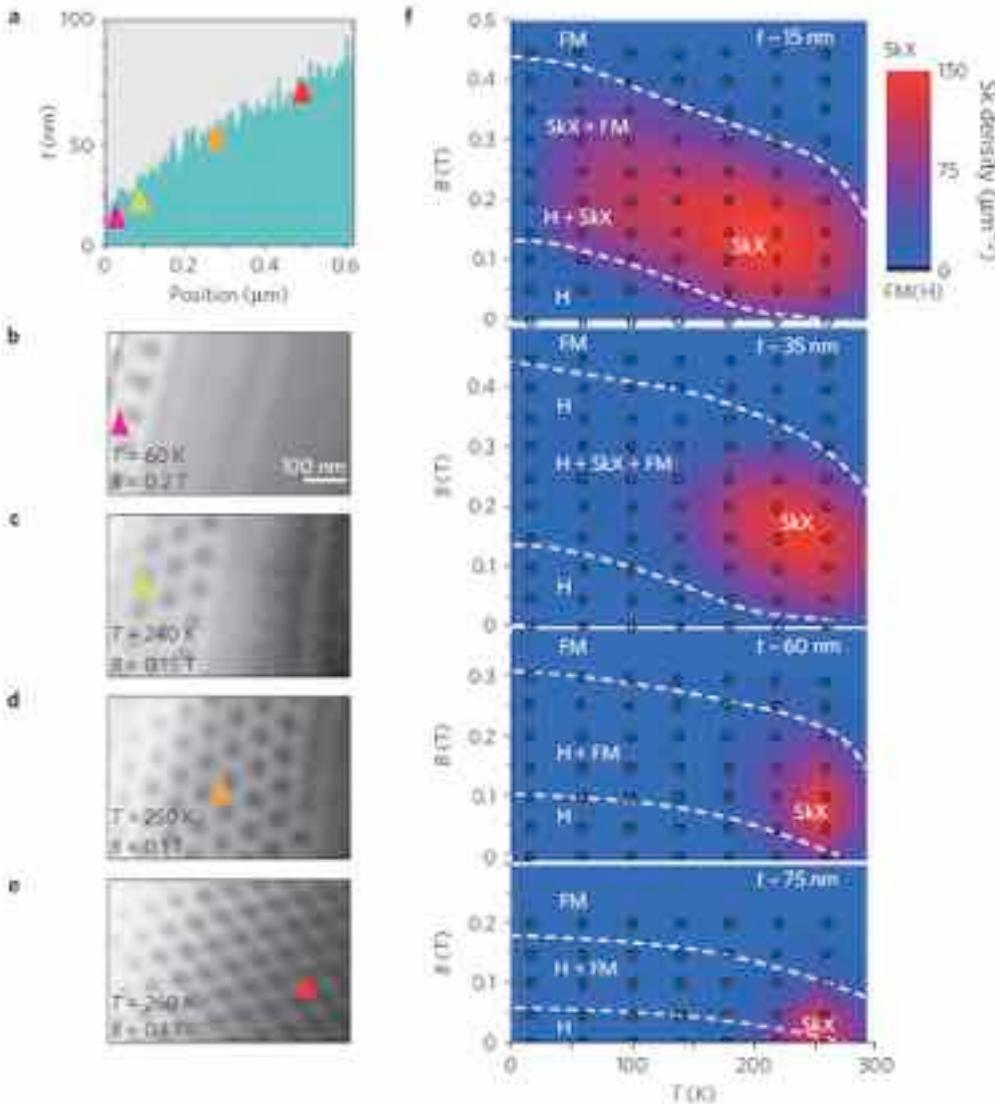


C : $t \sim 247 \text{ nm}$



Magnetic phase diagram of wedge-shaped lamella of single crystalline
 $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ after field-cooling (FC) process with 0.2 T applied at 365 K.

Magnetic phase diagram of skyrmions: FeGe (B20, chiral)

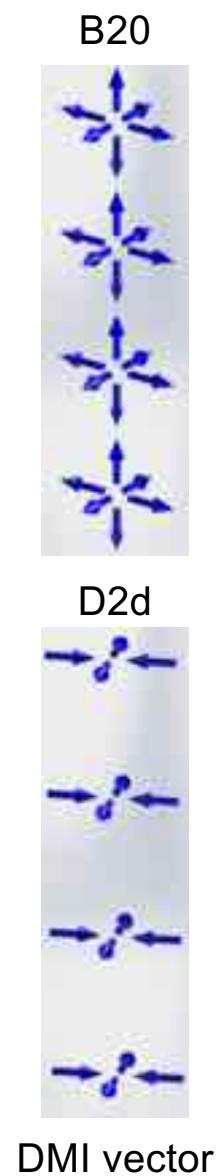
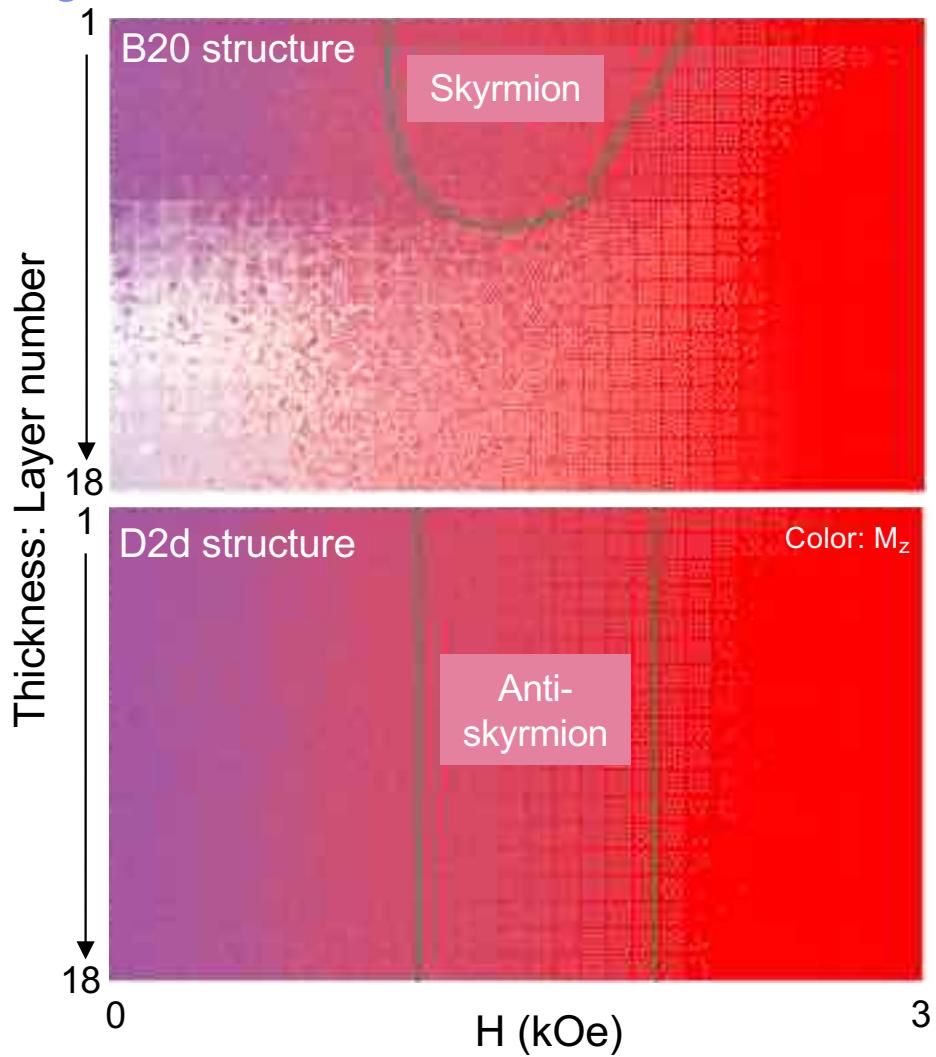


→ Anti-skyrmions are stable over wide range of temperature and magnetic fields as compared to skyrmions

→ Skyrmions are more stable at lower thicknesses

X. Z. Yu , Y. Tokura *et al.*,
Nat. Mater. 10, 106, (2011)

Thickness and field dependent phase diagram of B20 and D2d structures



$$H_{B20} = -J \sum_r \vec{S}_r \cdot (\vec{S}_{r+a\hat{x}} + \vec{S}_{r+a\hat{y}} + \vec{S}_{r+a\hat{z}}) - \vec{H} \cdot \sum_r \vec{S}_r \\ - D_{B20} \sum_r (\vec{S}_r \times \vec{S}_{r+a\hat{x}} \cdot \hat{x} + \vec{S}_r \times \vec{S}_{r+a\hat{y}} \cdot \hat{y} + \vec{S}_r \times \vec{S}_{r+a\hat{z}} \cdot \hat{z})$$

Magnetization **modulated** along thickness, due to z component of DMI.

→ Twisted Skyrmion tube preferred.

→ Skymions **not stable** when thickness is too large (twist from surface to center larger than $\pi/2$).

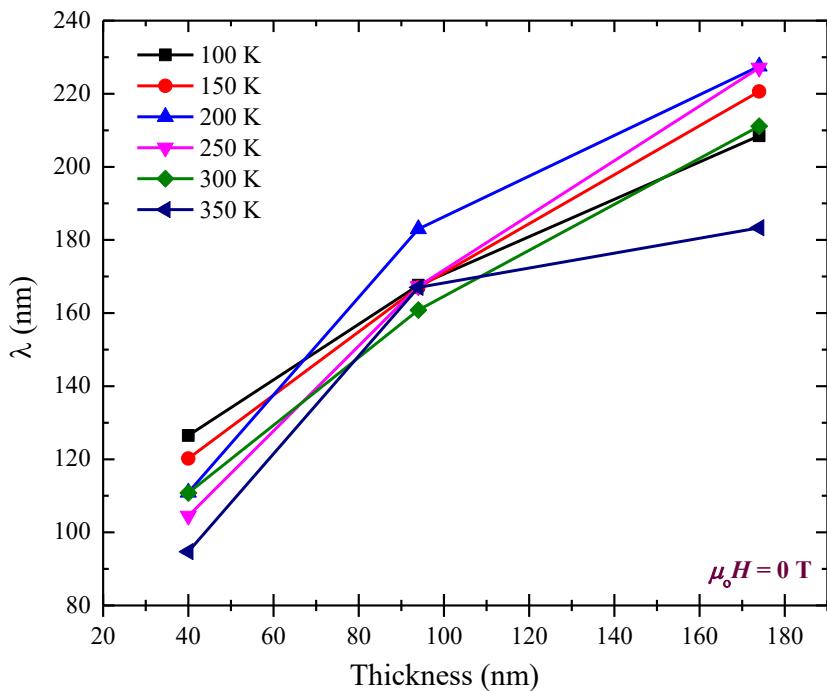
$$H_{D2d} = -J \sum_r \vec{S}_r \cdot (\vec{S}_{r+a\hat{x}} + \vec{S}_{r+a\hat{y}} + \vec{S}_{r+a\hat{z}}) - \vec{H} \cdot \sum_r \vec{S}_r \\ - D_{D2d} \sum_r (-\vec{S}_r \times \vec{S}_{r+a\hat{x}} \cdot \hat{x} + \vec{S}_r \times \vec{S}_{r+a\hat{y}} \cdot \hat{y})$$

Magnetization **unchanged** along thickness, due to Heisenberg exchange.

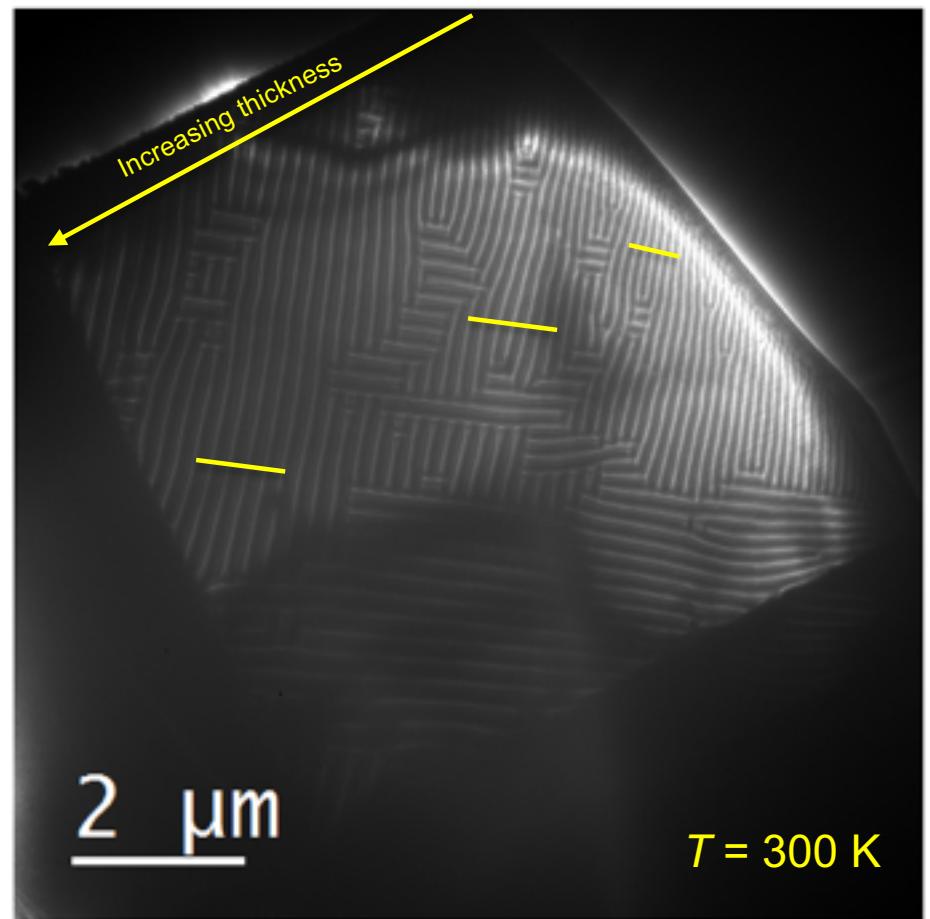
→ Untwisted Anti-Skyrmion tube preferred.

→ Anti-Skyrmion remain **stable** to large thicknesses.

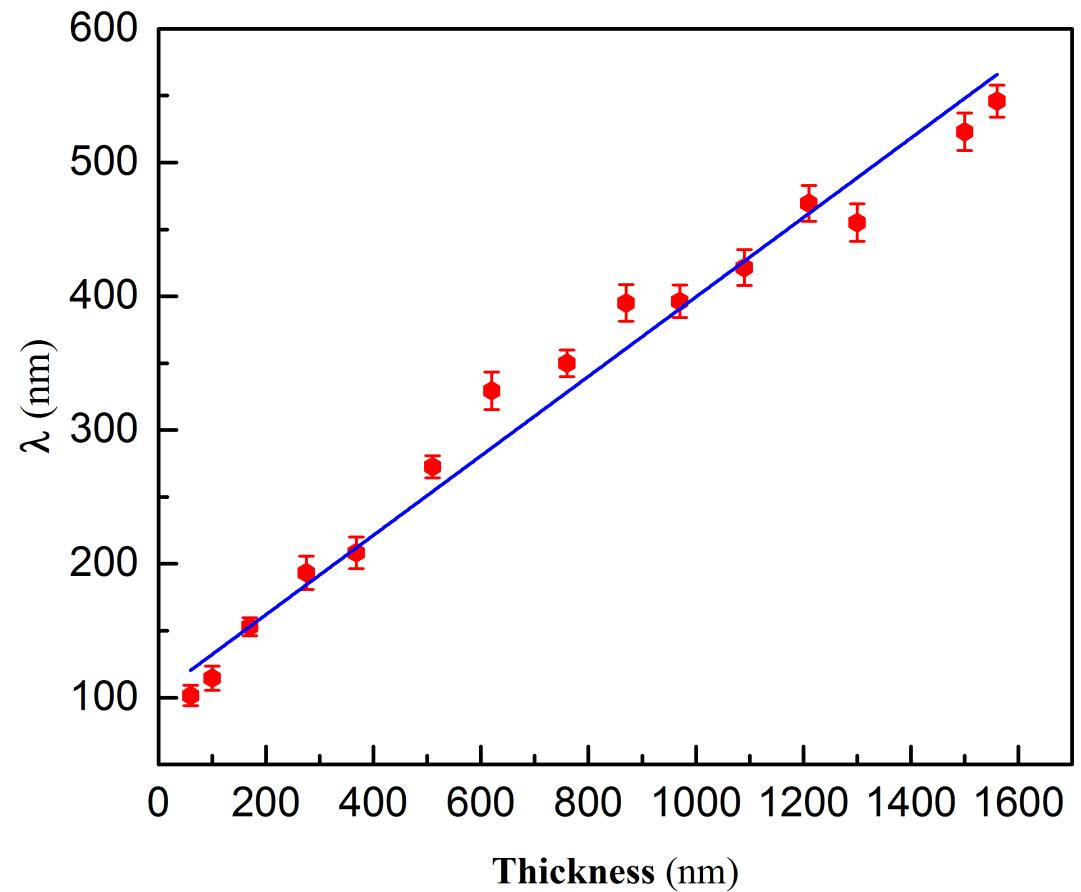
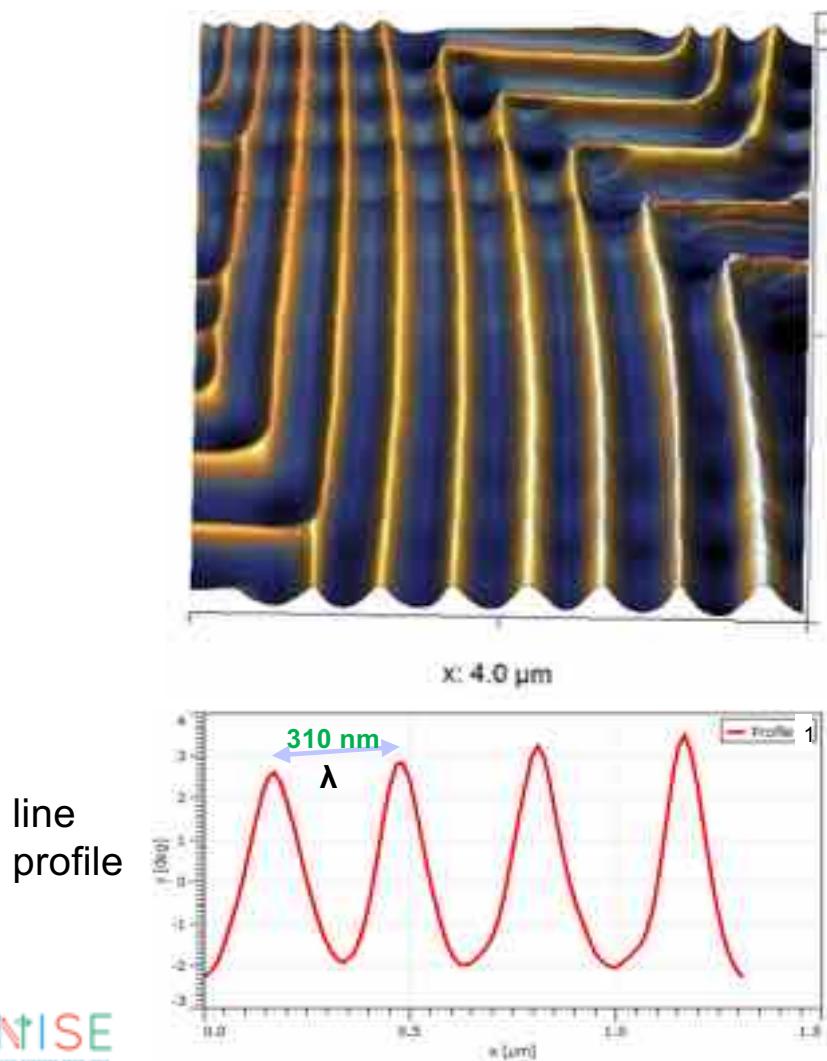
Thickness dependence of helical periodicity in wedge-shaped lamella (Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn)



Helical periodicity increases with thickness

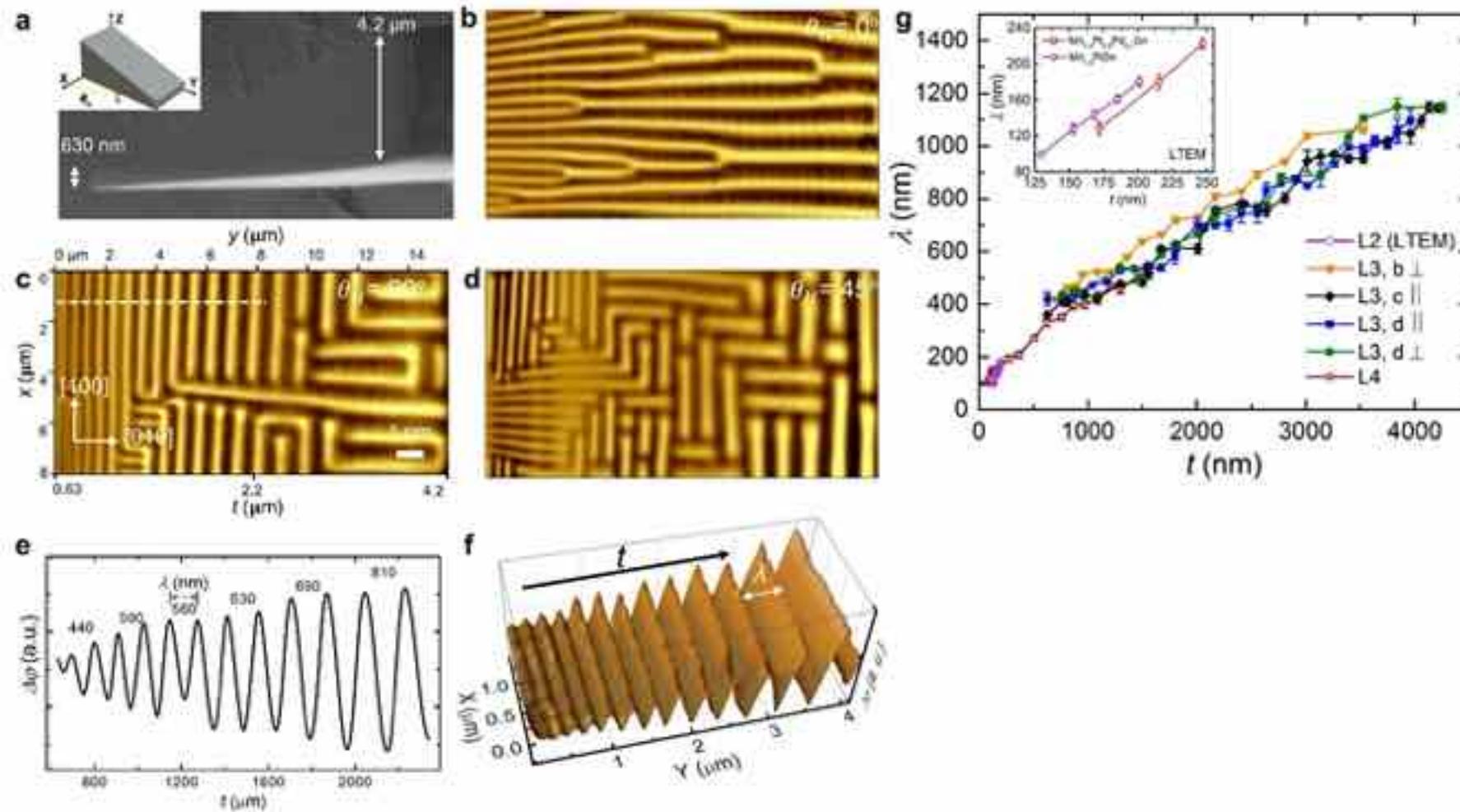


Variable temperature magnetic force microscopy imaging of helical phase



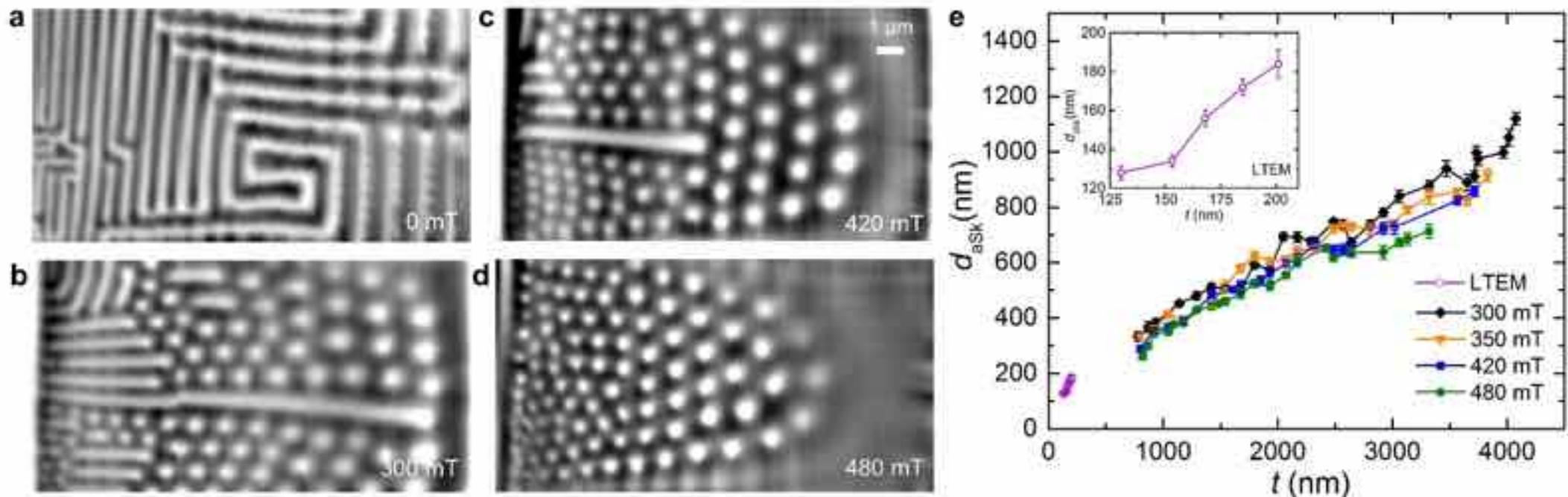
→ Helical wavelength λ increases linearly with thickness from 80 nm to 1600 nm

Variable temperature MFM: helical wavelength vs thickness up to 4 microns



Magnetic force microscopy (MFM) image of helical magnetic phase in a wedge-shaped lamella of single crystalline Mn_{1.4}PtSn at 300 K in zero magnetic fields and Helical wavelength vs. thickness of the lamella crystal.

Variable temperature MFM: aSk size vs thickness up to 4 microns

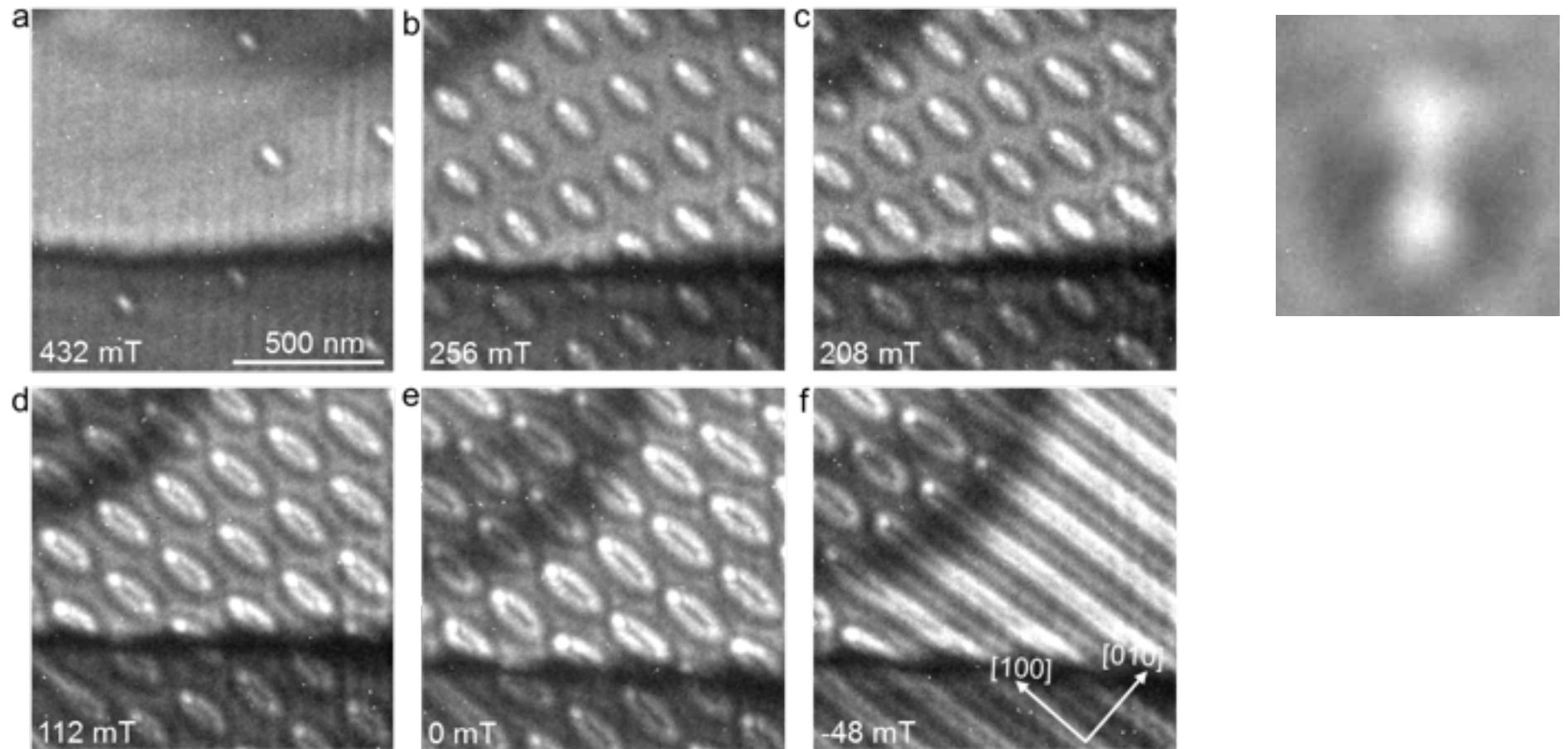


Magnetic force microscopy (MFM) image of magnetic field dependence of Antiskyrmion

- Thickness of the lamella increases from left end to right. The size of the images are $8 \times 16 \mu\text{m}$.
- Maximum size of the antiskyrmion is found to be $1.16 \mu\text{m}$.

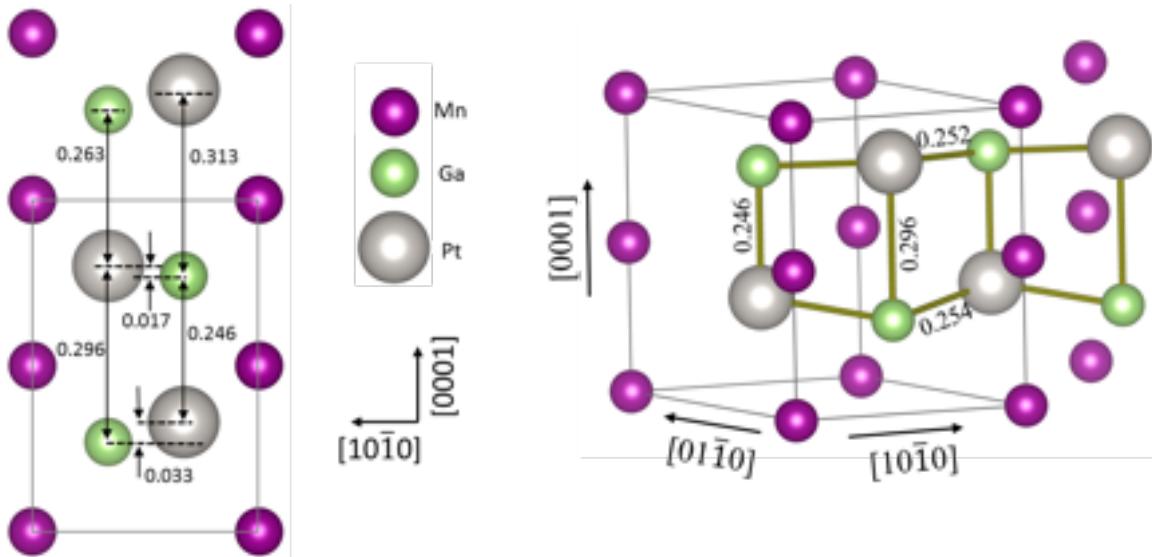
Antiskyrmion size vs. thickness 300 K

Anti-skymion → elliptical skymion due to dipole-dipole interaction



→ Metastable elliptical Bloch skyrmions can be stabilized in same material

PtMnGa: crystal structure (buckled layered)



Crystal structure:

trigonal (P3m1, space group no. 156)

Lattice Parameters:

$a=b=4.35\text{\AA}$, $c=5.59\text{\AA}$

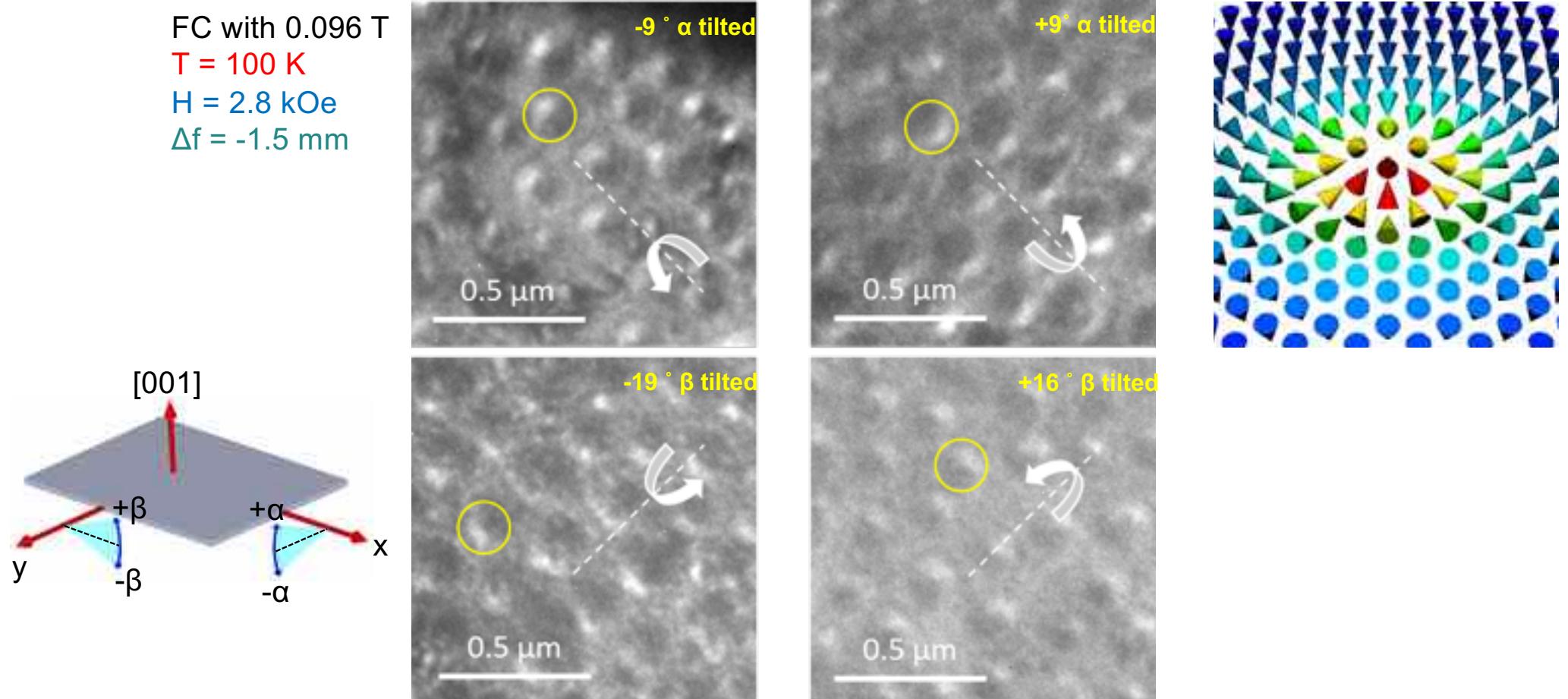
$\alpha=\beta=90^\circ$, $\gamma=120^\circ$

→ Isostructural to hourglass Fermion material KHgSb

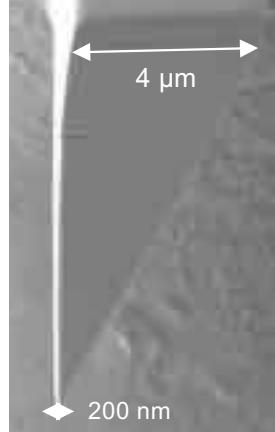
H K L	$ F_{\text{obs}}(HKL) ^2$	$ F_{\text{calc}}(HKL) ^2$
0 0 1	196.59	180.80
0 1 1	6840.50	5999.50
0 2 1	2750.74	2577.01
-1 1 0	2646.92	2518.25
-1 1 1	3647.66	4349.99
-1 1 2	16313.61	11653.70
-1 2 0	29582.05	28119.56
-1 2 1	124.15	85.47
0 0 3	518.51	654.83
-1 3 0	649.65	614.80
-2 2 1	2081.43	3479.31
-2 2 2	3779.58	2752.53
0 2 2	4448.76	5720.66
-1 2 2	6201.97	8002.20
-2 2 0	1419.61	1180.43
-2 3 1	2616.53	1623.58
-1 2 4	2441.26	2914.10

Bold letters indicate observed reflections, which are forbidden in space group P6₃/mmc

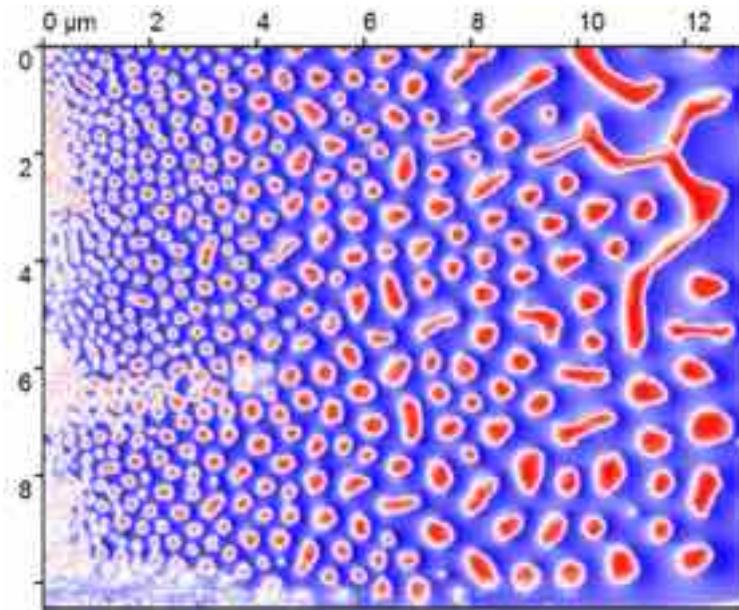
PtMnGa: Lorentz TEM observation of Néel skyrmions



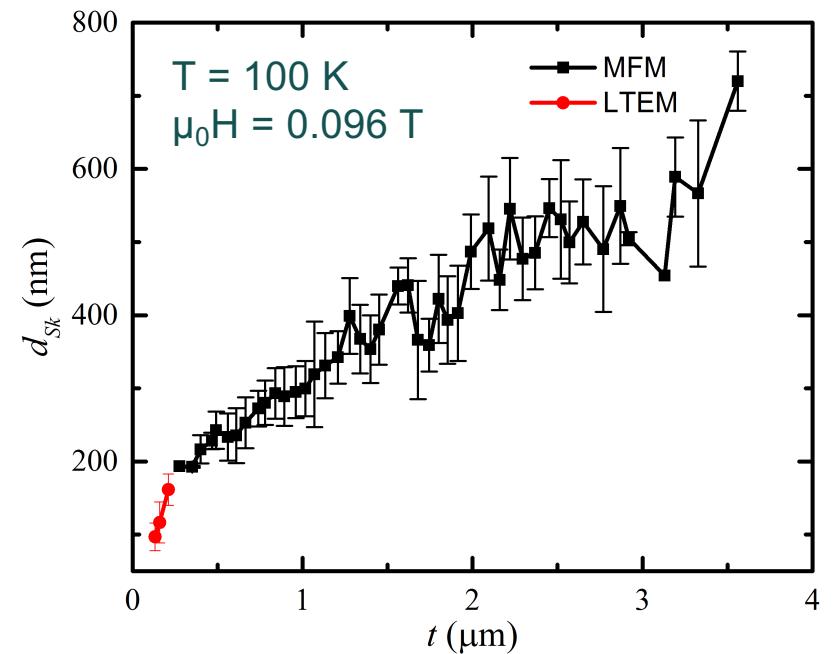
Stable Skyrmions for thicker layers observed using MFM



SEM side view image of the wedge-shaped lamella

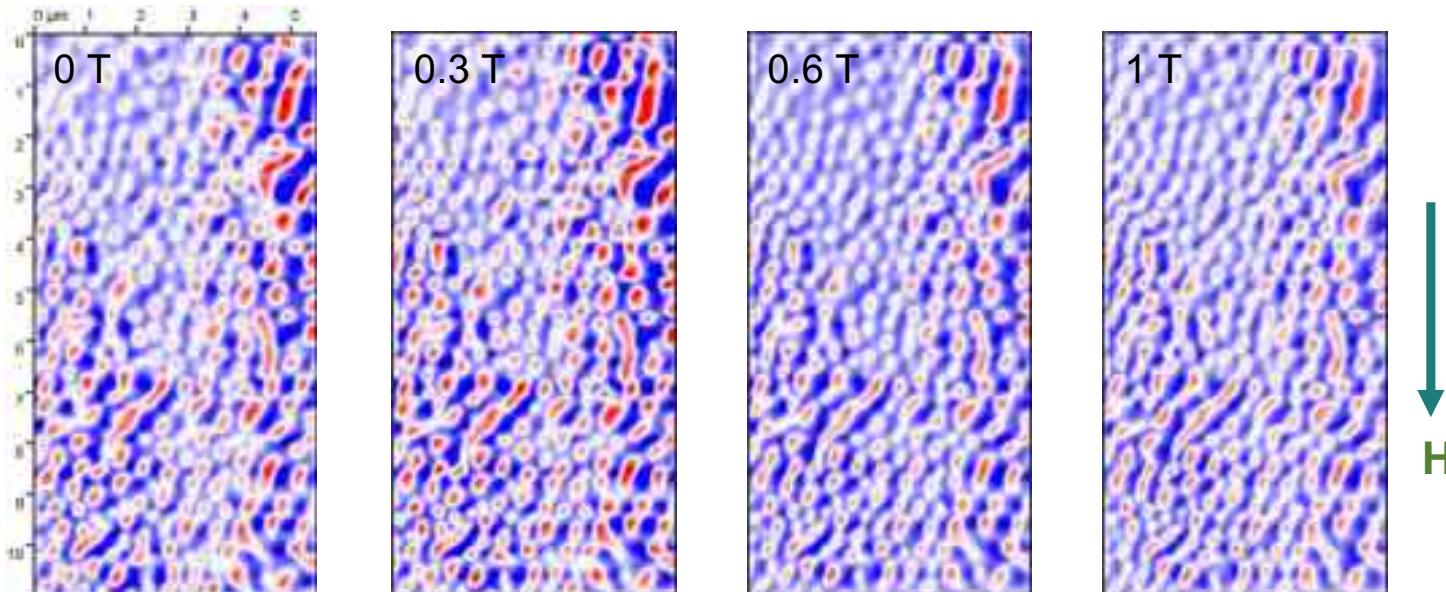


MFM image of Néel skyrmions

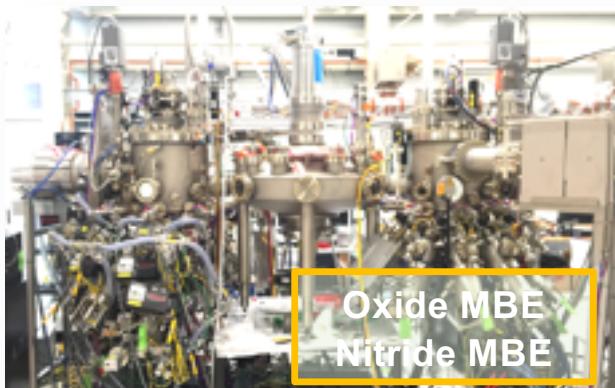


Skrymion size vs thickness

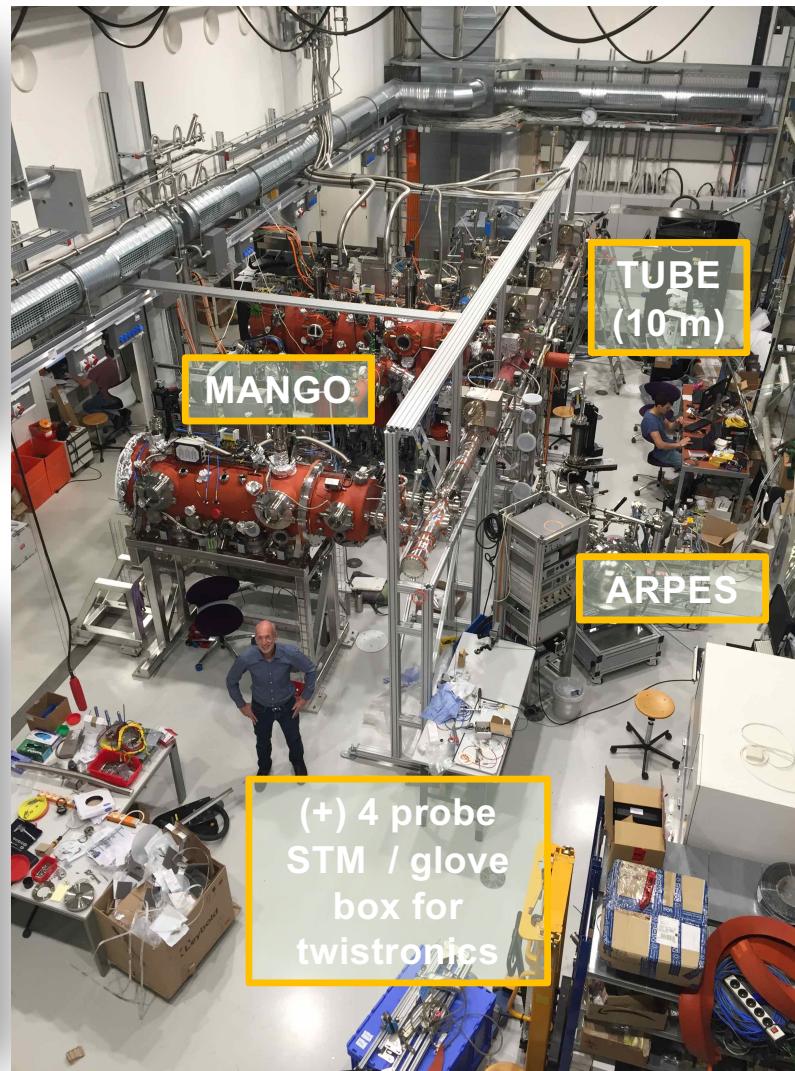
Robustness against in-plane magnetic field



MFM images of metastable Néel skyrmions in a uniform lamella of thickness 1 μm

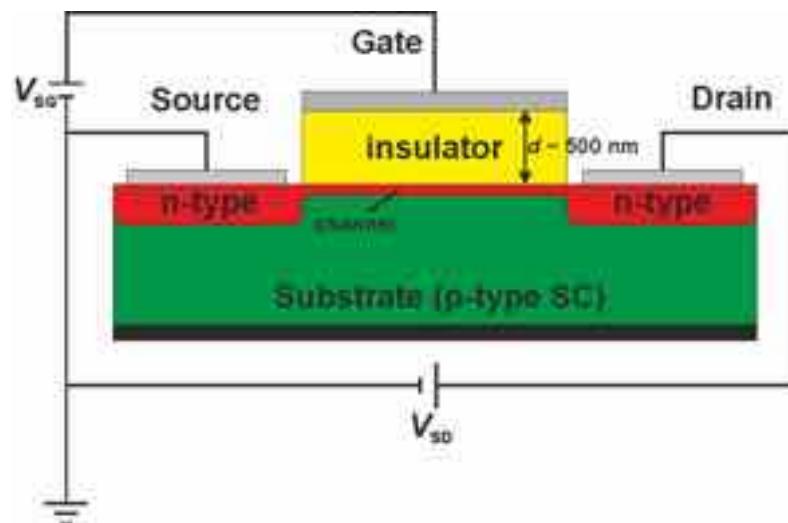


Thin film deposition



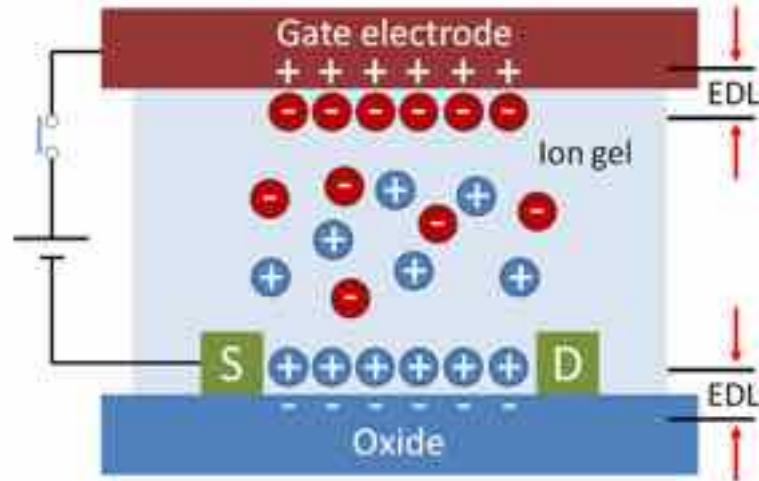
Conventional versus ionic liquid gating

**Metal-Insulator-Semiconductor
(MIS) FET structure**



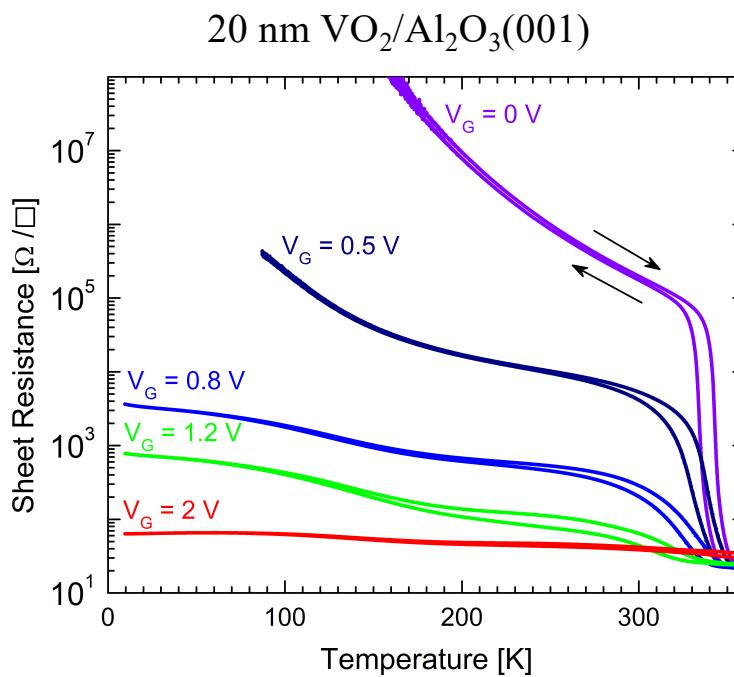
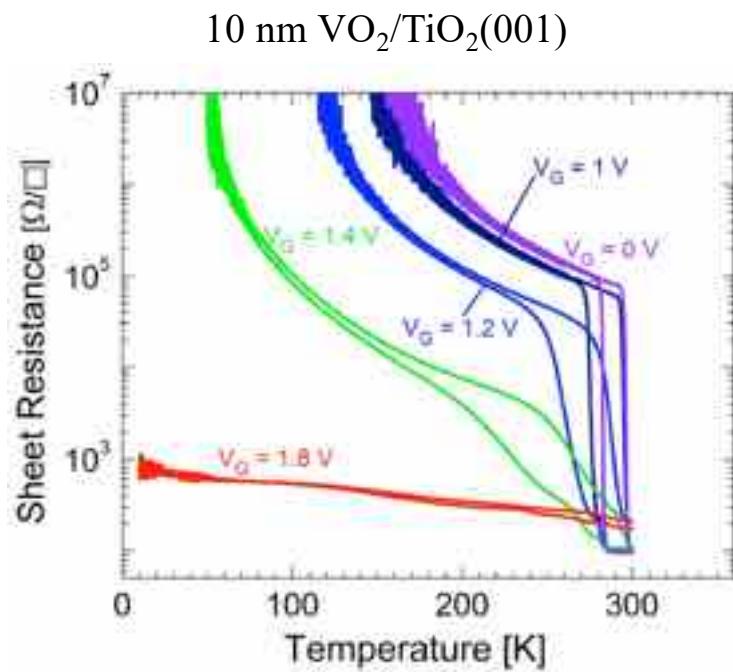
- Insulating layer ≈ 100 nm
- Capacitance $\approx 10 - 50 \text{ nF cm}^{-2}$
- Charge carrier density $\approx 10^{13} \text{ cm}^{-2}$
- High gate voltages necessary

**Electrochemical Double Layer
(EDL) FET structure**



- EDL ≈ 1 nm
- Capacitance $\approx 1 \mu\text{F cm}^{-2}$
- **Charge carrier density $\approx 10^{15} \text{ cm}^{-2}$**
- Gate voltage $< 3V$

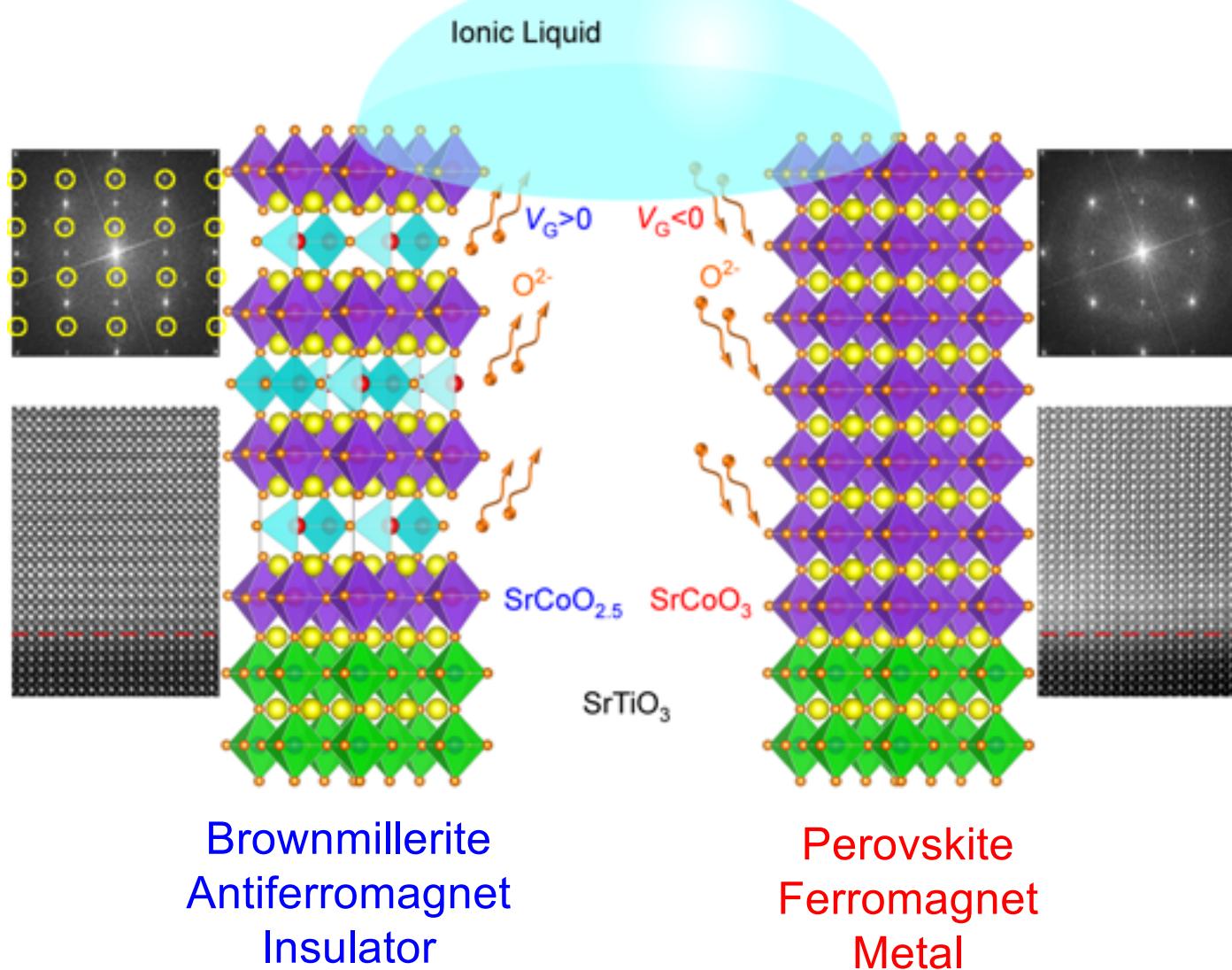
Ionic Liquid Gate induced Suppression of MIT



- MIT is suppressed down to 10K with the application of gate voltage.
- No signature of residual MIT
→ The entire film is metallized.

Science (2013), PNAS (2015, 2016), Nano Lett. (2013, 2016), Adv. Mater. (2016, 2017)
Phys. Rev. Lett, (2016, 2017), Nano Lett. (2017), DRC (2018), Nat. Commun. (2018)

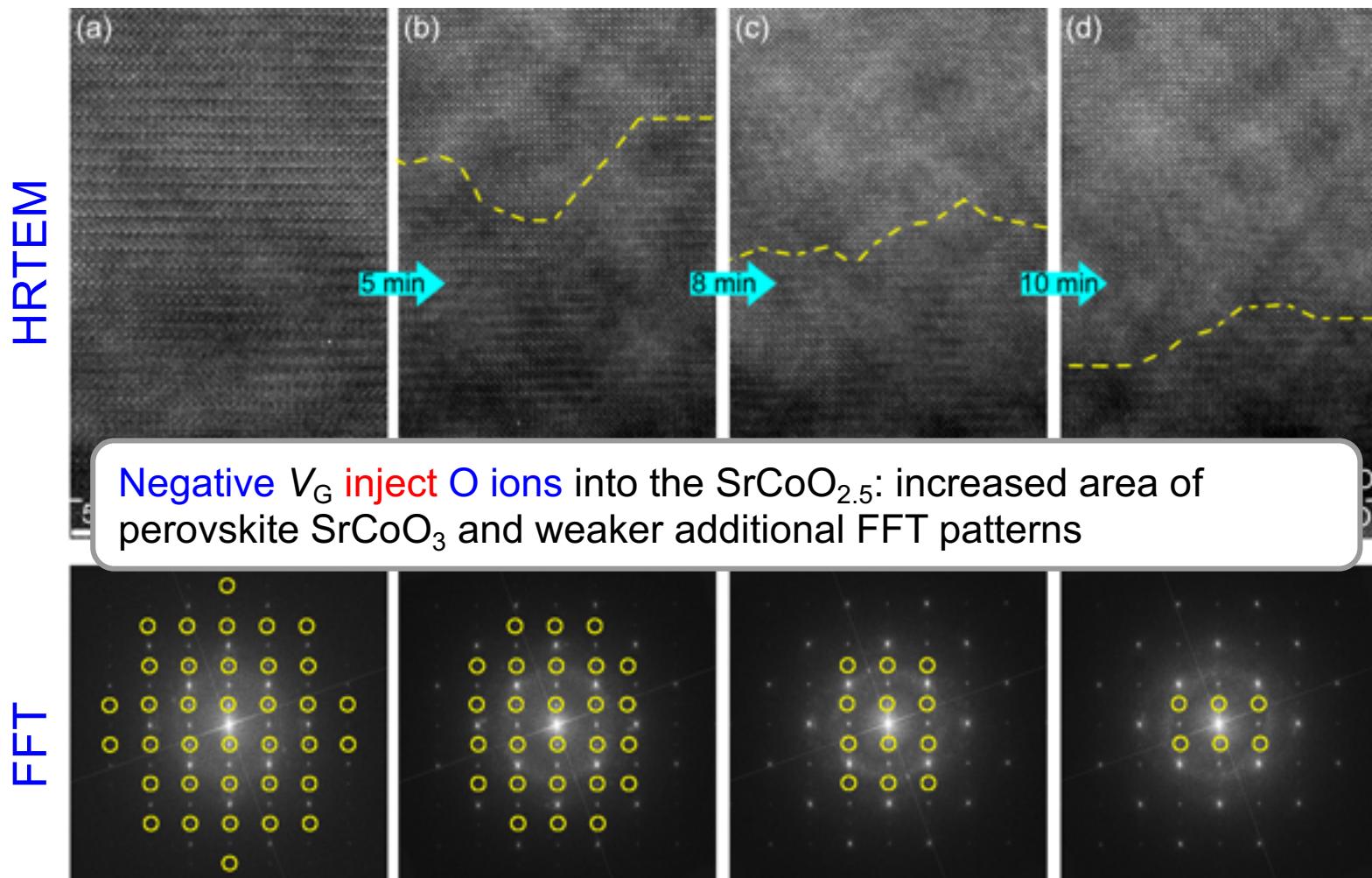
Brownmillerite – perovskite: $\text{SrCoO}_{2.5}$ – SrCoO_3



Positive and negative gate voltages applied through ionic liquid will extract oxygen from and inject oxygen into SCO, respectively, resulting in the brownmillerite $\text{SrCoO}_{2.5}$ and perovskite SrCoO_3 .

In-situ TEM: Brownmillerite \rightarrow perovskite: $\text{SrCoO}_{2.5} - \text{SrCoO}_3$

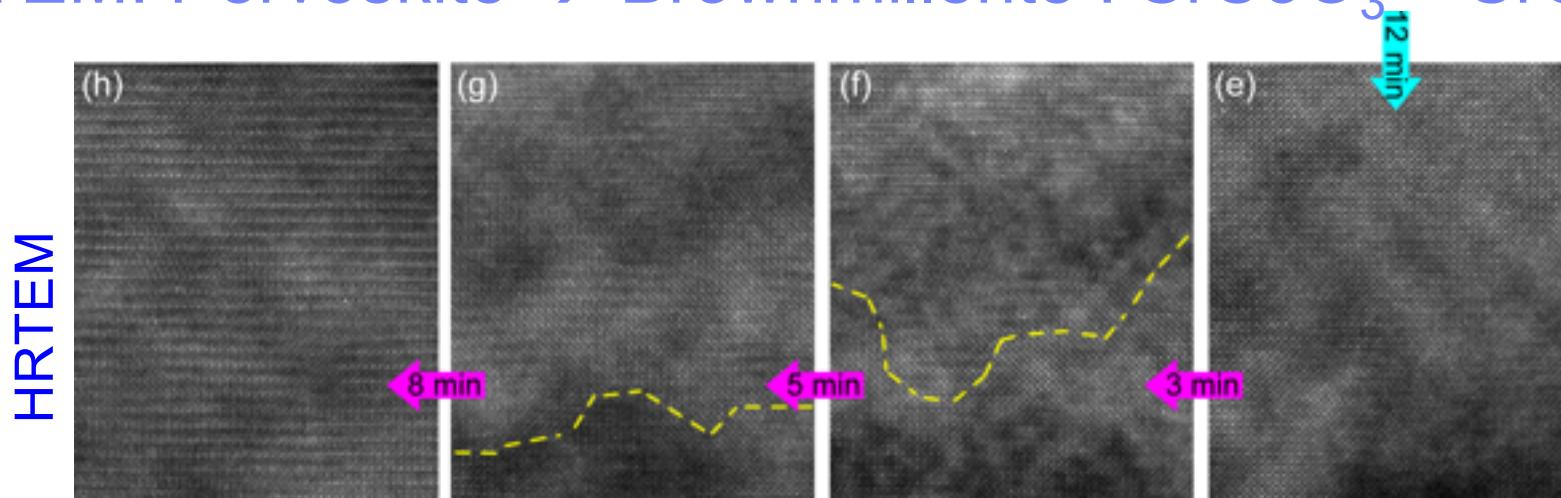
$V_G = -3V$



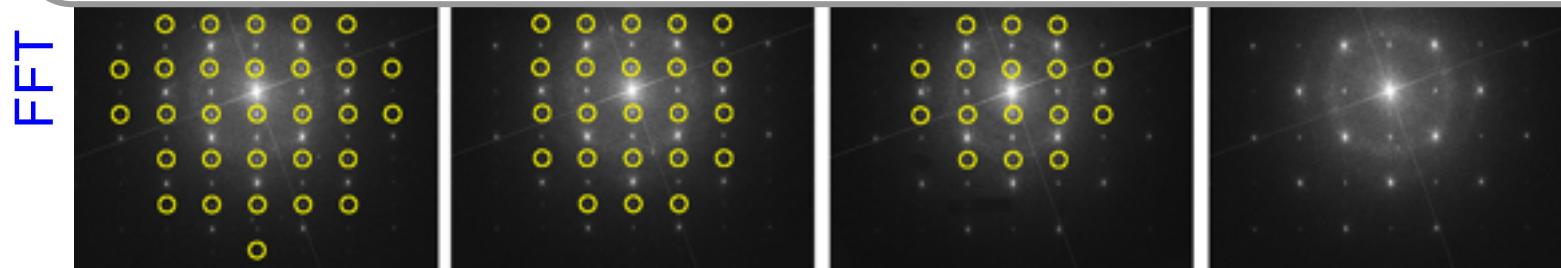
Time-dependent phase transition between $\text{SrCoO}_{2.5}$ and SrCoO_3 with ionic liquid gating. Ionic liquid on the left side.

In-situ TEM: Pervoskite \rightarrow Brownmillerite : $\text{SrCoO}_3 - \text{SrCoO}_{2.5}$

$V_G = +3V$

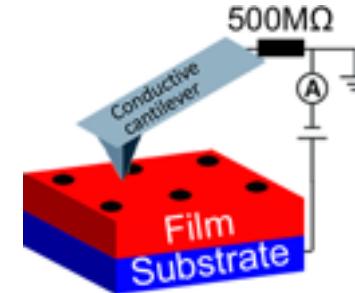
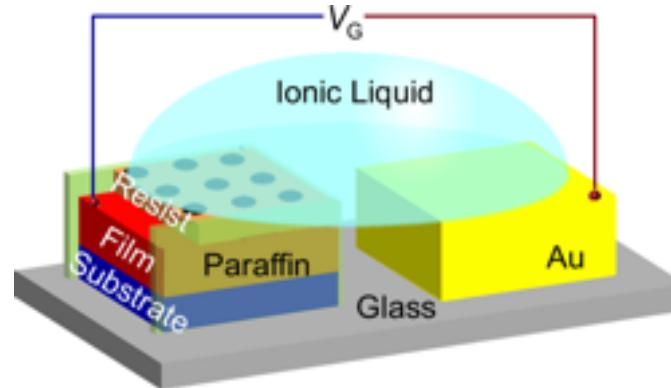
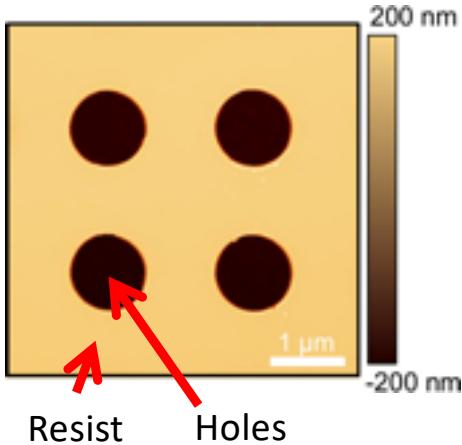


1. V_G reversibly manipulates phase transition of $\text{SrCoO}_3\text{-}\text{SrCoO}_{2.5}$ by injecting or extracting oxygen ions;
2. Phase transition always starts from the surface and then extends into the bulk of thin film.



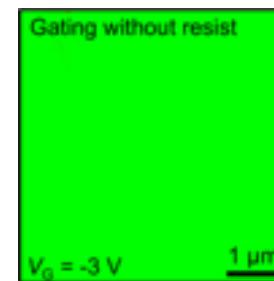
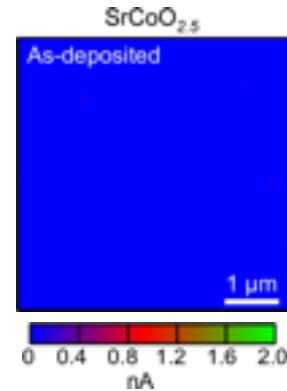
Top HRTEM, bottom FFT; yellow dashed line: rough phase boundary

IL gating controlled meso-structures

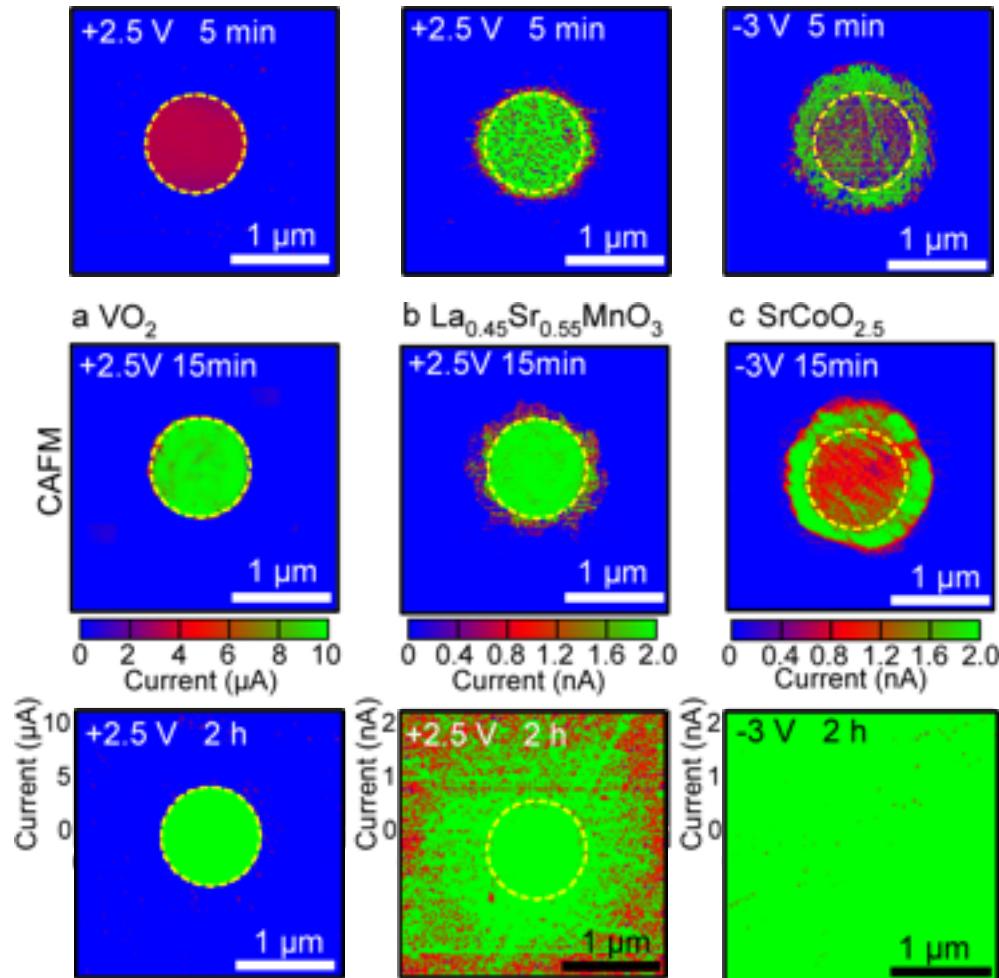


Resist can effectively obstruct the gating effect right above

CAFM



IL gated patterned oxide film → 3D meso-structures



➤ VO_2

Strictly limited in the resist
hole: [001] O transport channel

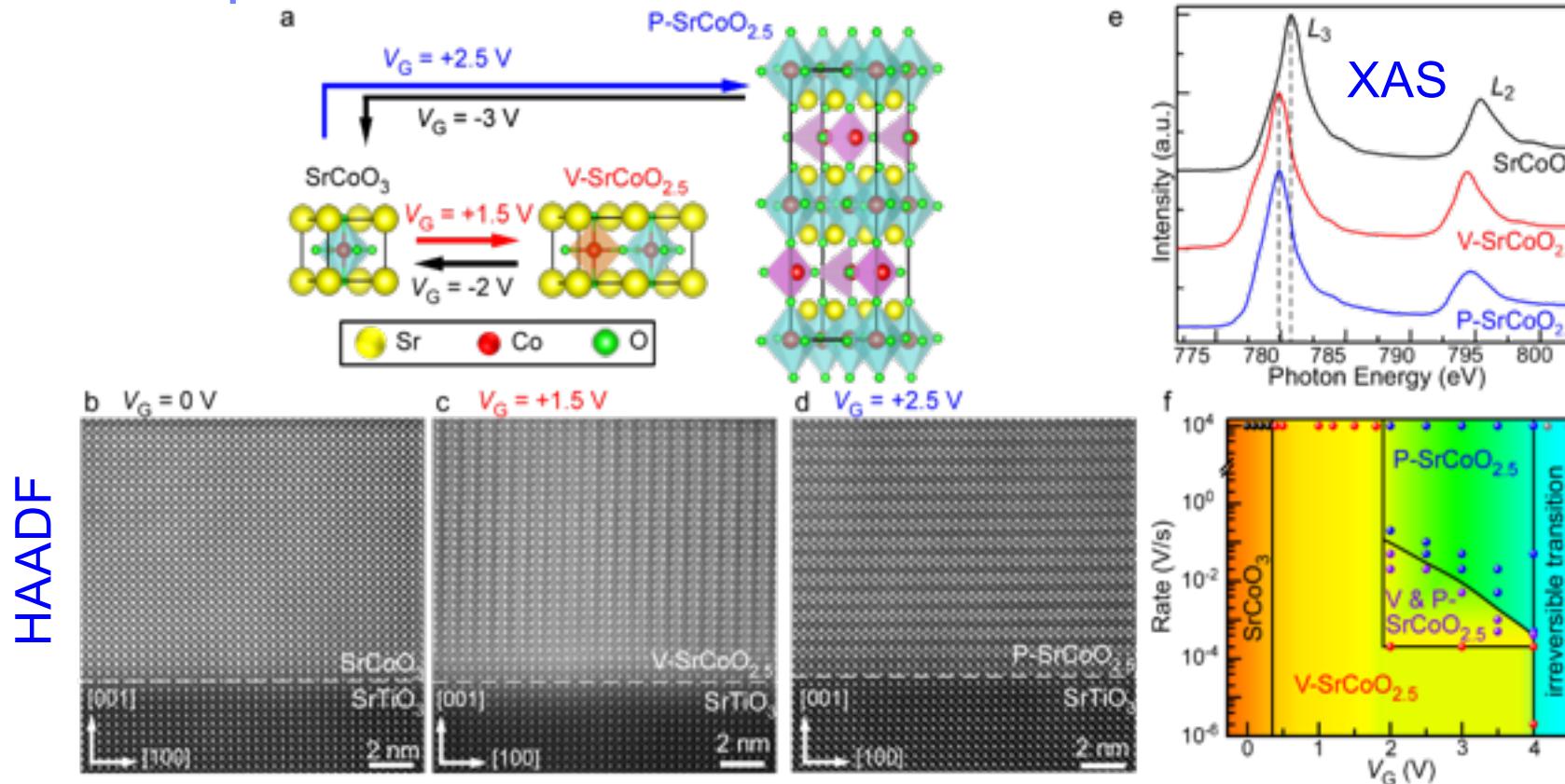
➤ $\text{La}_{0.45}\text{Sr}_{0.55}\text{MnO}_3$

Emanative pattern beyond the
hole: multidirectional O
transport

➤ $\text{SrCoO}_{2.5}$

A more conductive ring around
the hole: multidirectional O
transport &
perovskite/brownmilleriate
interface

Multi-state phase transitions

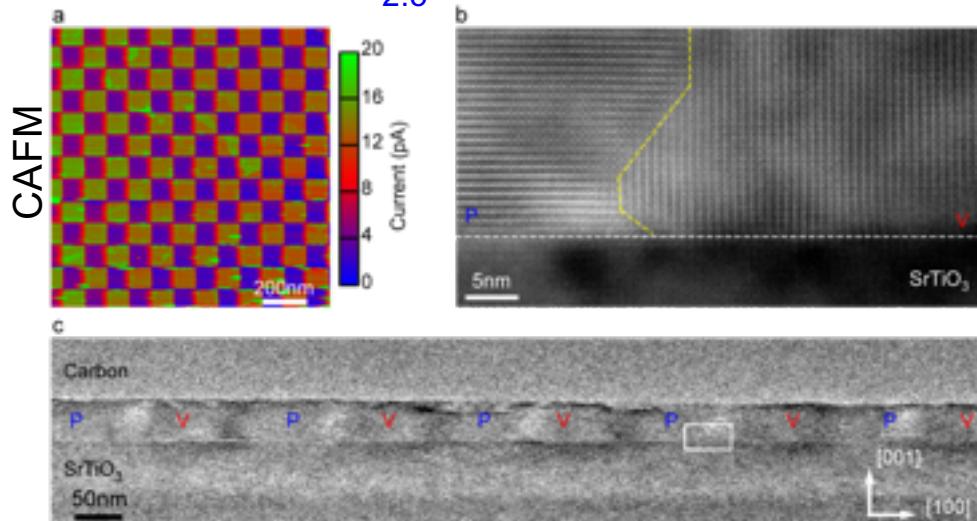


Beyond two states phase change

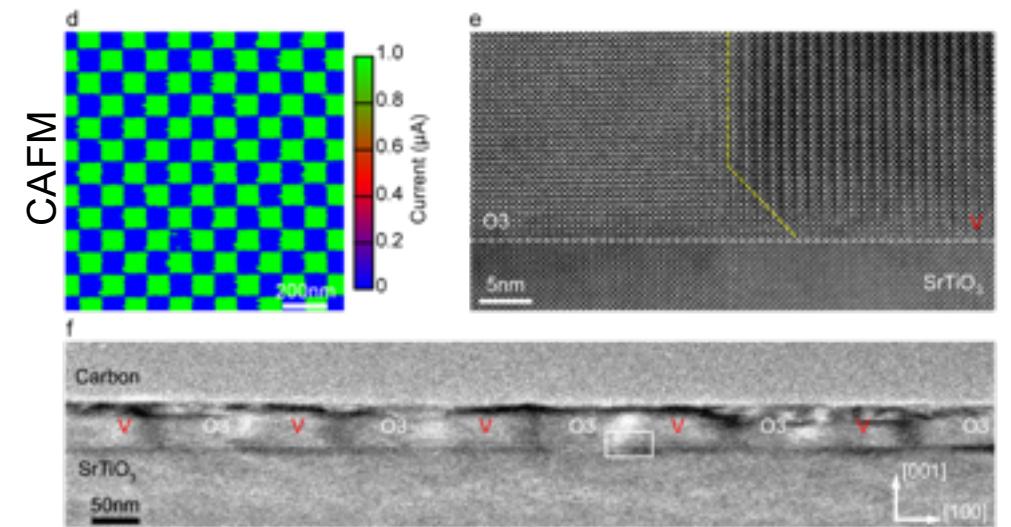
SrCoO_3 could be reversibly changed to $\text{SrCoO}_{2.5}$ with superstructure vertical and parallel to the surface using small and large gate voltages, respectively

Creation of magnetic meso-structures

P-SrCoO_{2.5} & V-SrCoO_{2.5}



SrCoO₃ & V-SrCoO_{2.5}

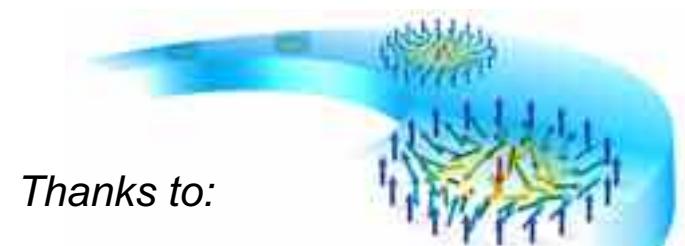


- Nanoscale gridding (~100nm) made by patterned ionic liquid gating

Chiral spintronics: Chiral and spatial spin textures

■ Anti-skyrmions

- Observed in several tetragonal inverse Heuslers
- Anti-skyrmion size & helical wavelength increase with thickness
- "elliptical" twinned skyrmions found in same system



Thanks to:

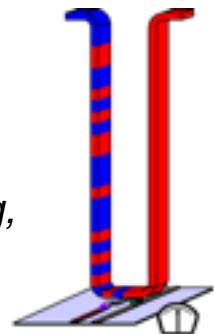
Tianping Ma, Rana Saha, Abhay Srivastava, Ajaya Nayak, Peter Werner, Robin Blasing, Kai-Uwe Demasius, MPI-Halle

■ Important technological applications – Racetrack Memory

- DW velocities > 1,000 m/sec in synthetic antiferromagnets
- 3T Single DW Racetrack - replacement for SRAM
- Promises 3T device with SRAM performance but increased density, much lower energy consumption & non-volatility (fast start-up)

Vivek Kumar, Roshnee Sahoo, Parul Devi, Claudia Felser
MPI-cpfs

Börge Göbel, Ingrid Mertig
MLU, Halle



See-Hun Yang, Chirag Garg,
IBM Research - Almaden

■ Artificial spin textures

- Using ionic liquid gate induced FM and AFM (2x) phases can create wide variety of spatial spin textures – racetrack, spin liquid...

Nayak et al. Nature (2017)
Yang et al. Nat. Phys. (2019)
Sana et al. Nat Comm. (2019)
Srivastava et al. Adv. Mater. (2019)

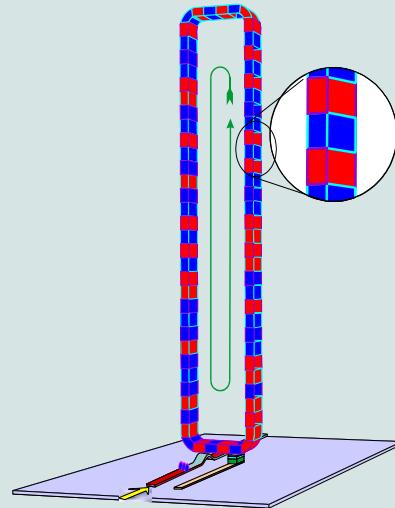
Innately 3D memory and logic devices



Today: 2D

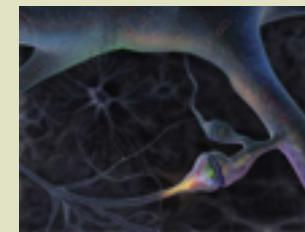


Future: innately 3D



Racetrack Memory

10 to 100 times the storage capacity of conventional solid state memory
→ Could displace flash memory and hard disk drives



Cognitive Devices
emulating synaptic functions in a solid state device
→ *Million times more energy efficient than charge based computers*