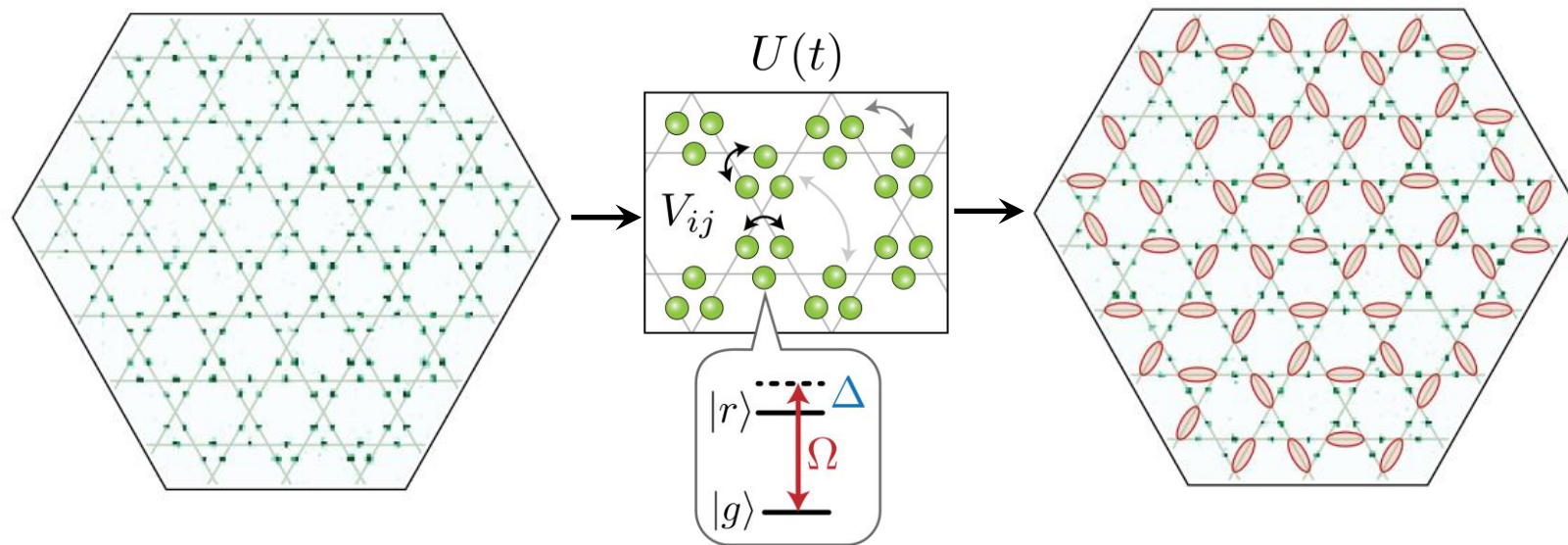


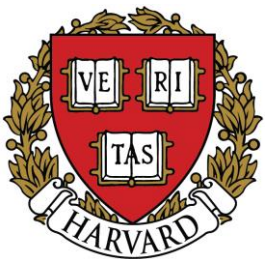
Exploring new scientific frontiers with programmable atom arrays



May 3rd, 2021

Giulia Semeghini

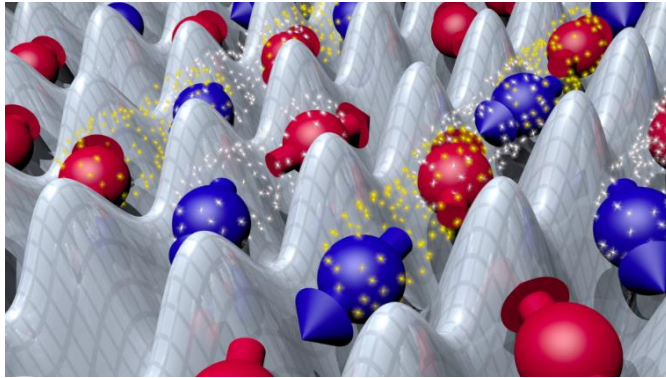
Lukin group - Harvard University



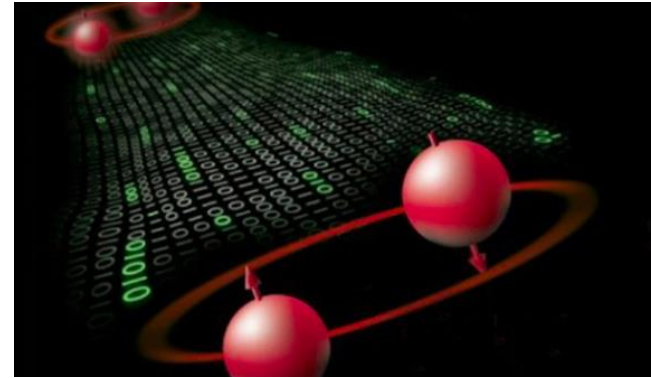
Frontiers of Quantum Science

Can we harness the power of quantum systems to learn new things about **nature** and develop new **technology**?

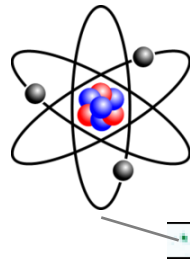
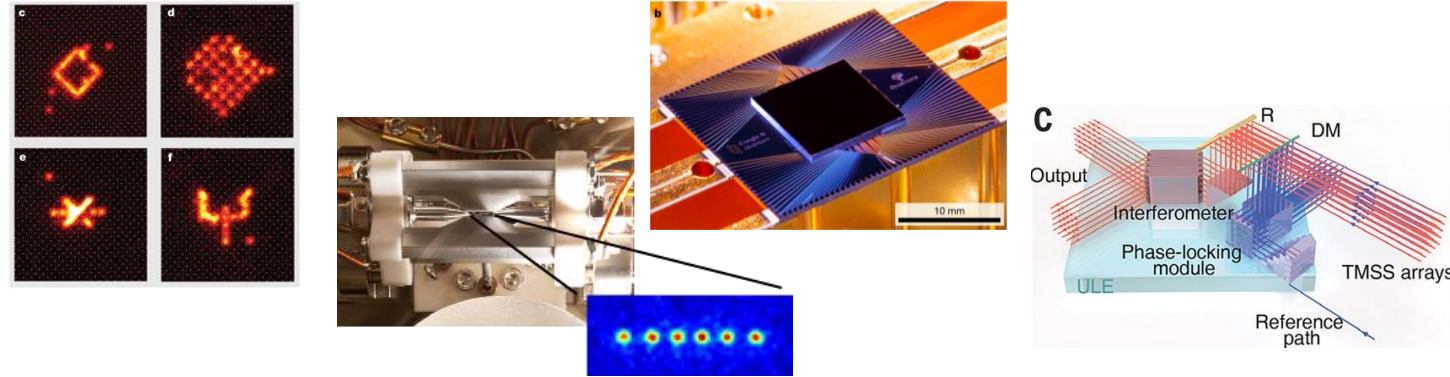
Quantum simulation



Quantum computing



Several platforms: ions, atoms, superconductors, photons, defects...

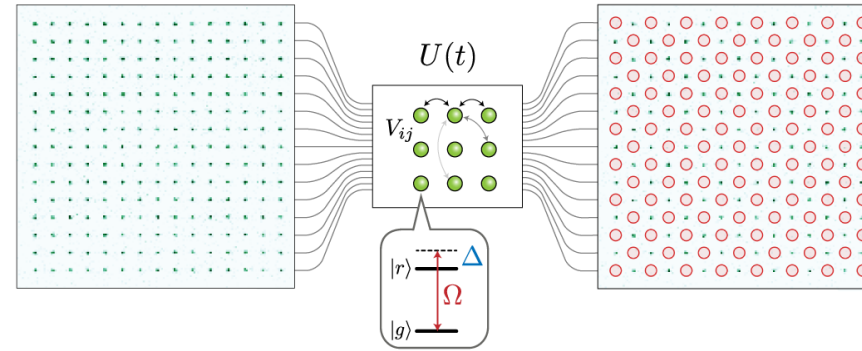


Our approach: **individual neutral atoms**

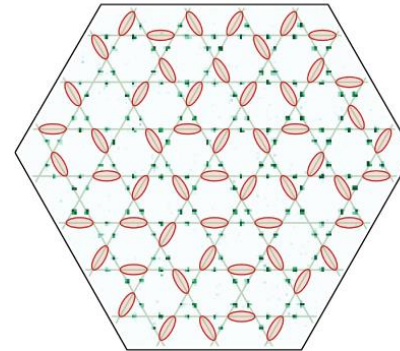
- Excellent isolation from the environment
- Well-developed toolbox:
 - Initialization, readout
 - Strong, switchable Rydberg interactions
- Highly scalable defect-free arrays (hundreds of atoms)
- Tunable system parameters

Outline

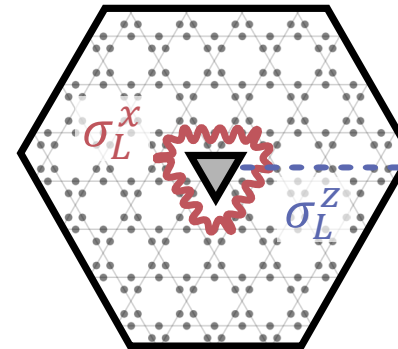
Programmable arrays
of Rydberg atoms



Quantum spin liquid phase
on a frustrated lattice

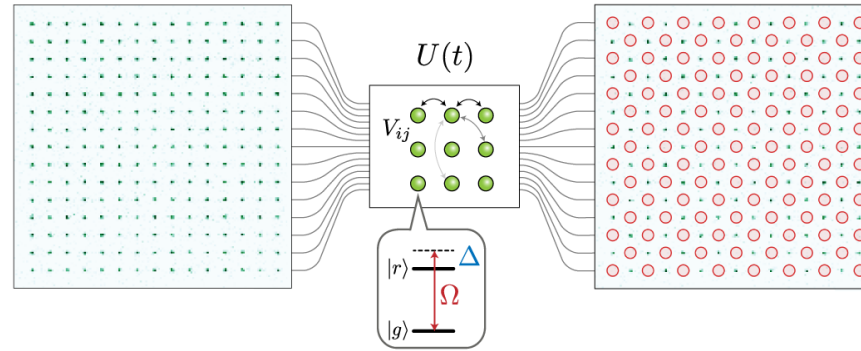


First steps towards a
topological qubit

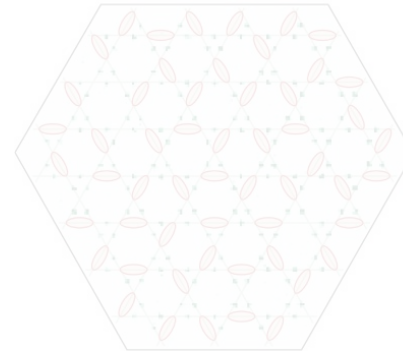


Outline

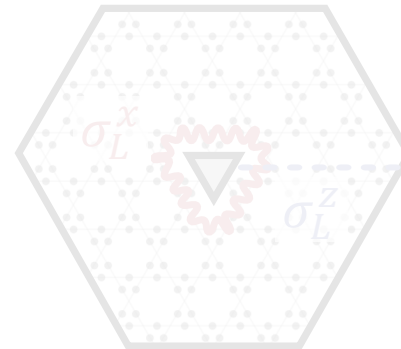
Programmable arrays of Rydberg atoms



Quantum spin liquid phase on a frustrated lattice



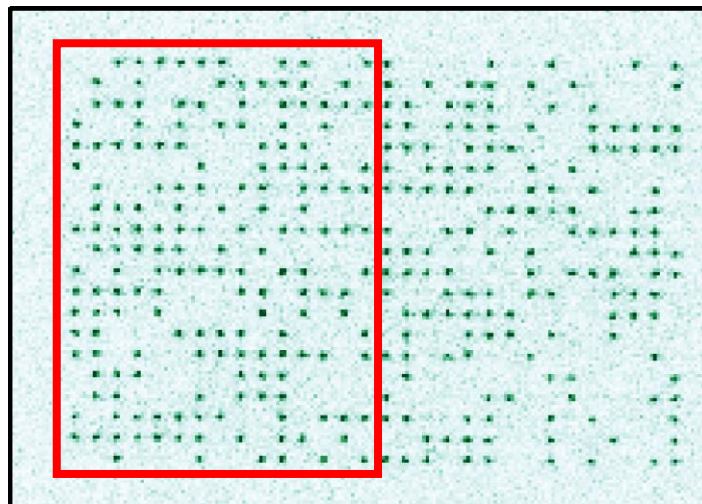
First steps towards a topological qubit



2D array of optical tweezers

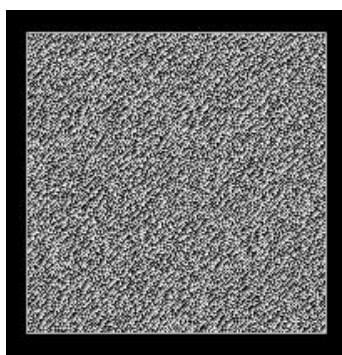


Atoms:



55-60% filling
fraction

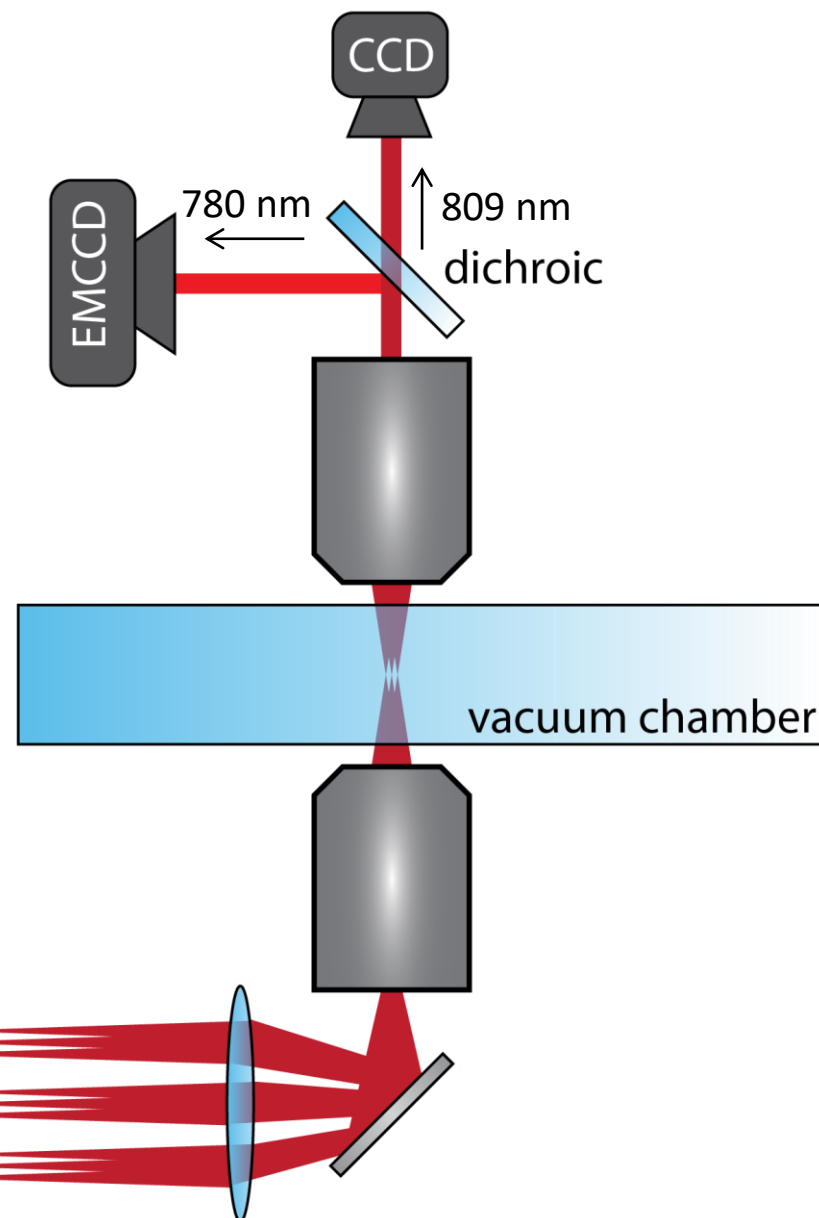
Deterministically
fill 300 sites?



phase
profile

Spatial Light
Modulator

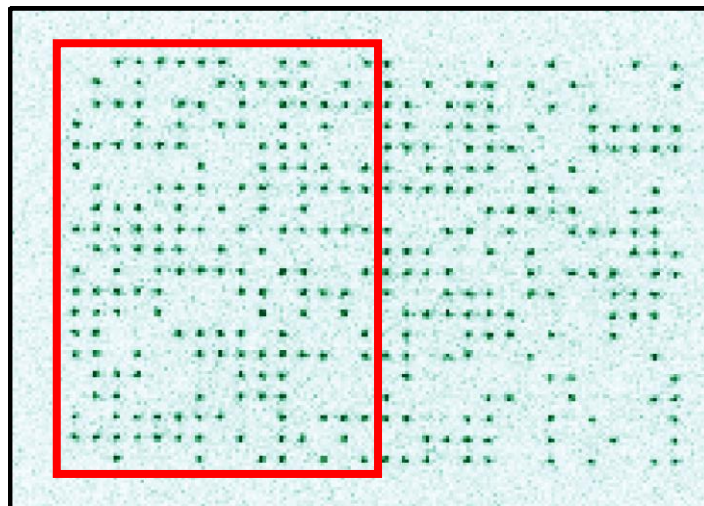
SLM



2D array of optical tweezers

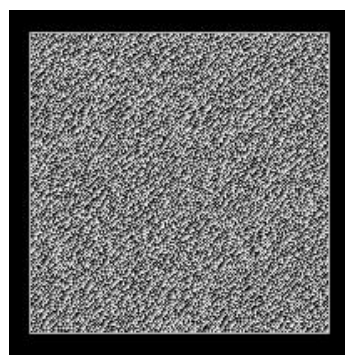


Atoms:



55-60% filling
fraction

Deterministically
fill 300 sites?



phase
profile

Spatial Light
Modulator

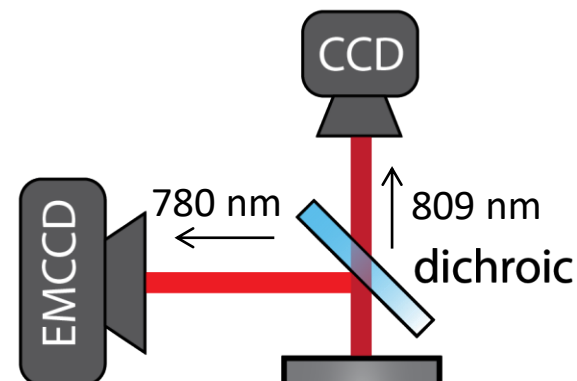
SLM

movable
tweezers

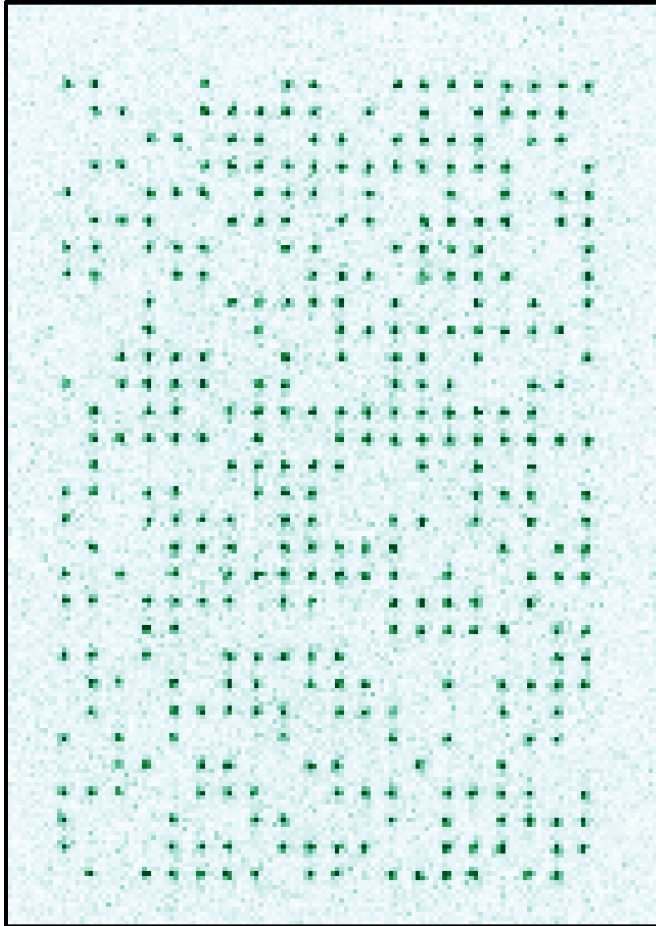
AOD

Acousto-Optical
Deflectors

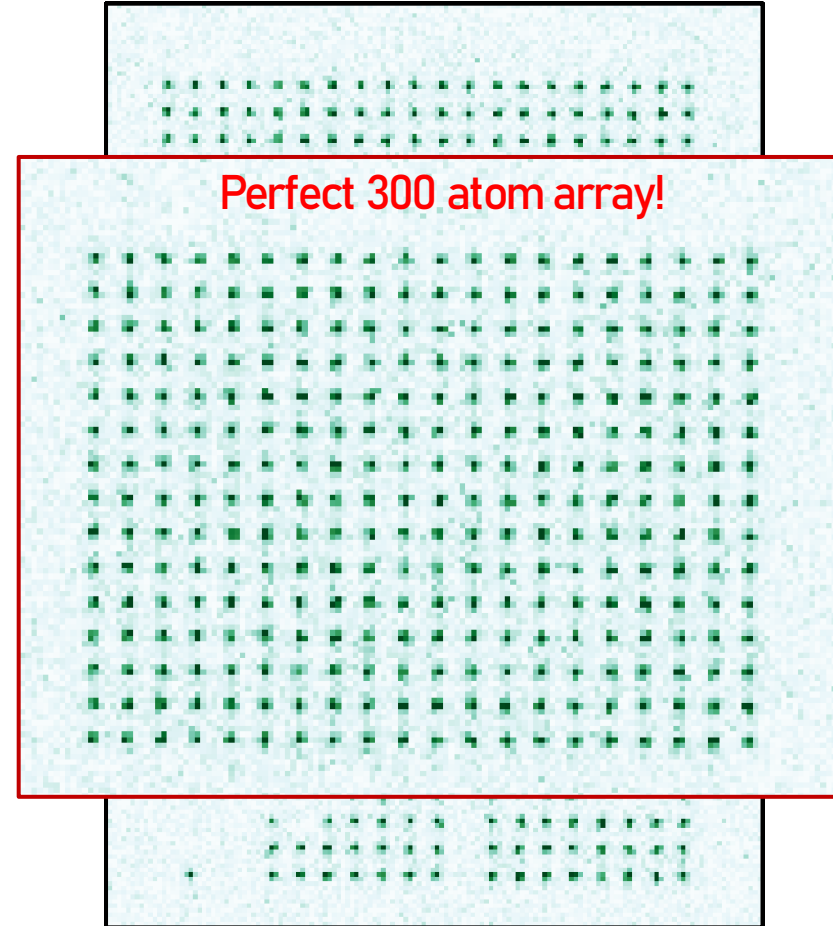
vacuum chamber



Initial loading:



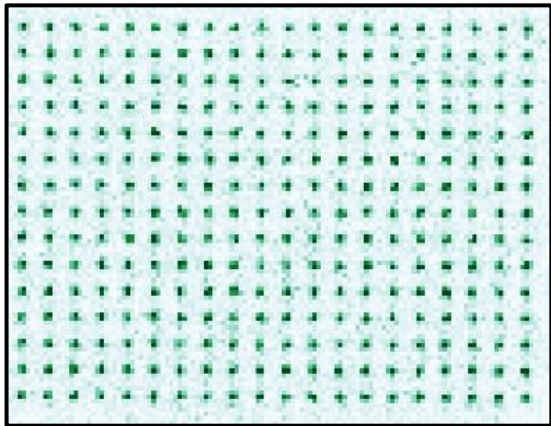
After sorting:



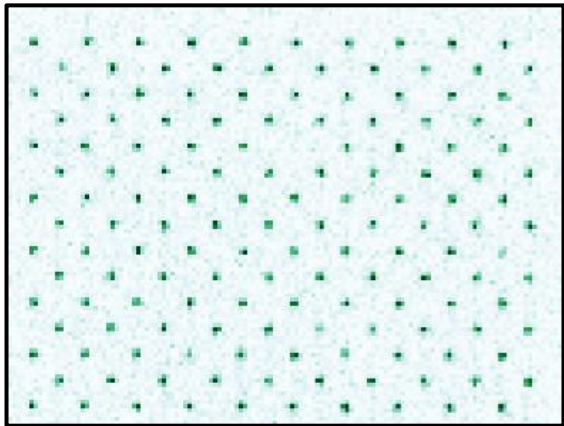
> 99% filling

Programmable 2D arrays

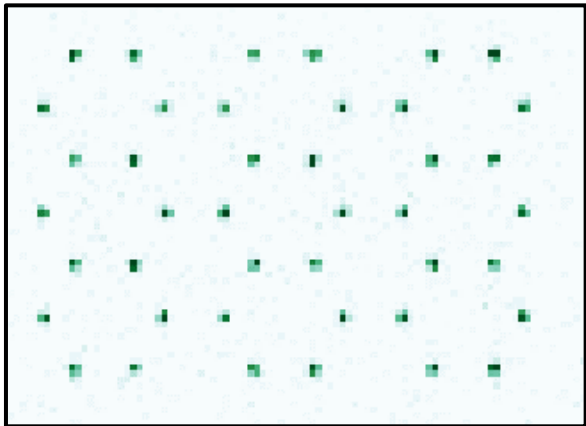
Square



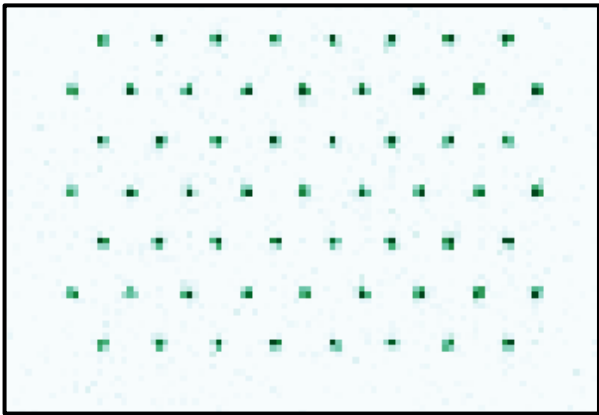
Tilted Square



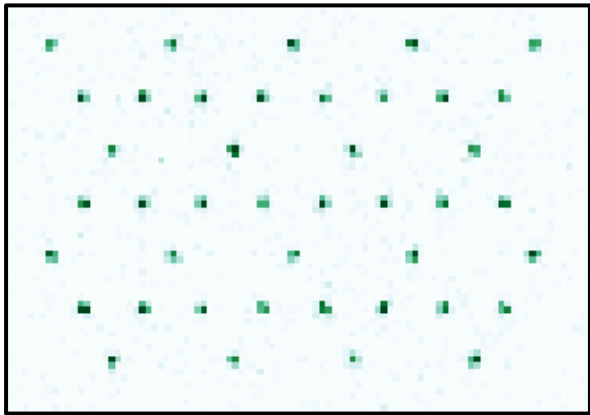
Honeycomb



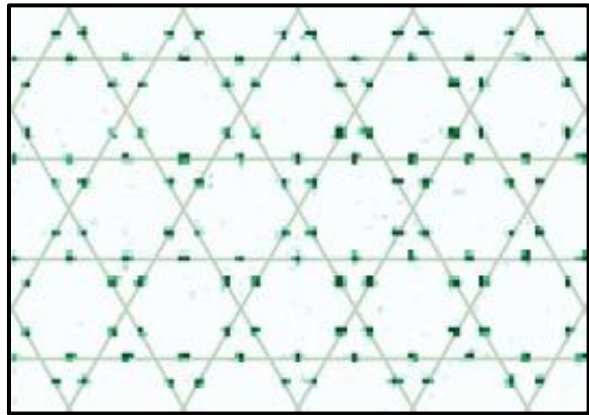
Triangular



Kagome



Link-kagome



Programmable 2D arrays

Square



Tilted Square

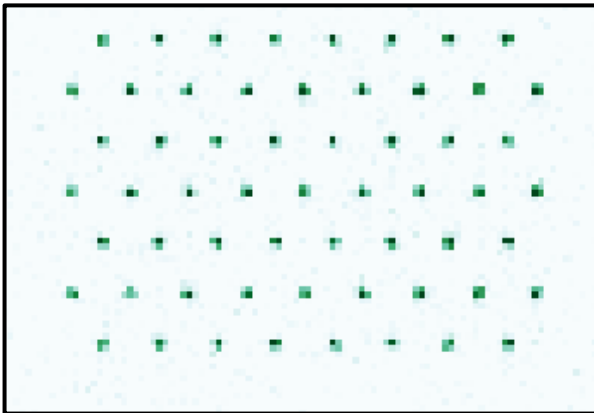


Honeycomb

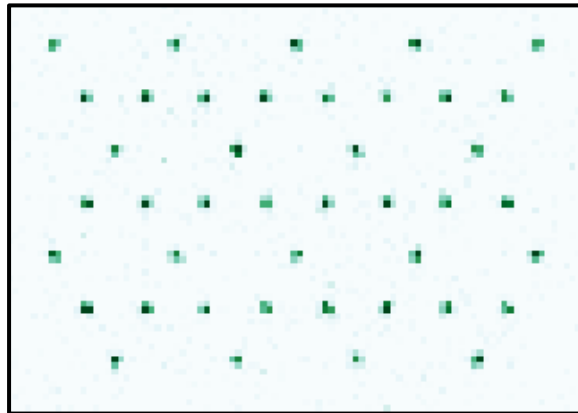


1) Programmable 2D arrays of neutral atoms

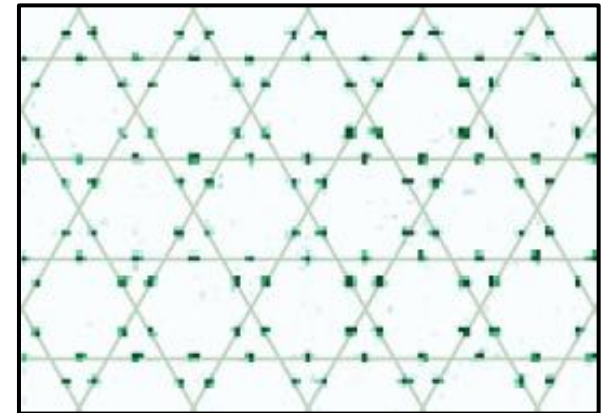
Triangular



Kagome

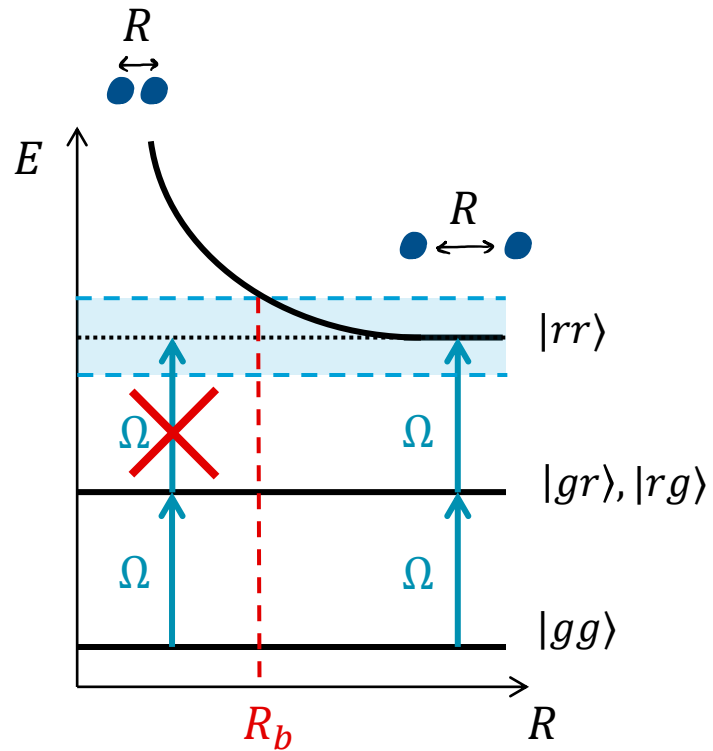


Link-kagome



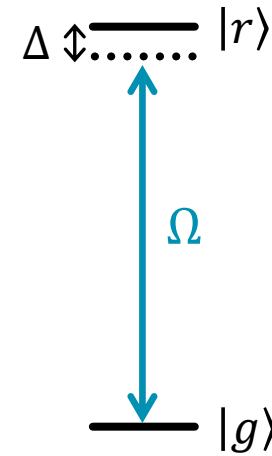
Rydberg states and long-range interactions

Rydberg blockade



blockade radius R_b :
 $V(R_b) = \Omega$

excitation to Rydberg state:



size ~ 0.2 nm

van der Waals interactions $\propto 1/R^6$

size > 200 nm

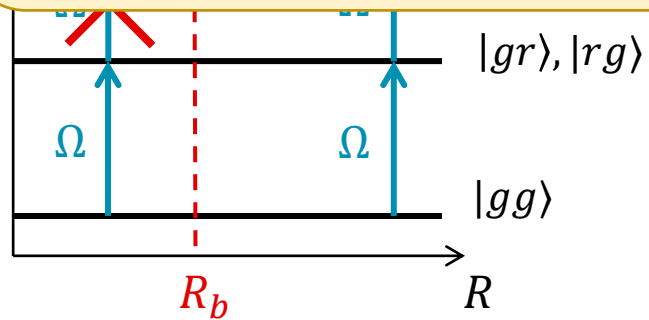
Rydberg states and long-range interactions

Rydberg blockade

excitation to Rydberg state:

size > 200nm

- 1) Programmable 2D arrays of neutral atoms
- 2) with strong long-range interactions



blockade radius R_b :
 $V(R_b) = \Omega$

$\Lambda \uparrow$ — $|r\rangle$

— $|g\rangle$

size ~ 0.2 nm

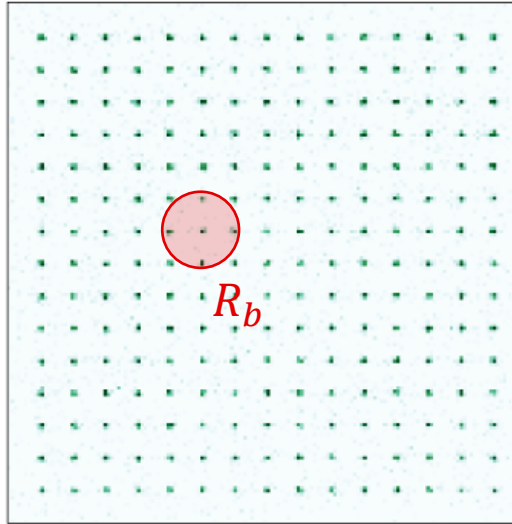
van der Waals interactions $\propto 1/R^6$

Rydberg blockade and quantum many-body phases

Many-Body ground state

$$\Delta > 0$$

Maximize number of
disconnected Rydberg atoms



antiferromagnetic
state

$$\mathcal{H} = \underbrace{\frac{1}{2}\Omega(t) \sum_i \sigma_x^{(i)}}_{\text{drive term}} - \underbrace{\sum_i \Delta(t) n_i}_{\text{detuning}} + \underbrace{\sum_{i<j} V_{ij} n_i n_j}_{\text{interaction}} \quad n_i = |r_i\rangle\langle r_i|$$

drive term

detuning

interaction

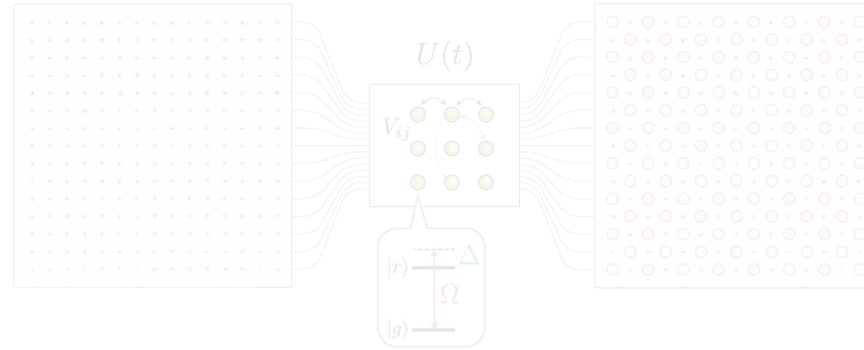
S. Ebadi et al., Quantum Phases of Matter on a 256-Atom Programmable Quantum Simulator, arXiv:2012.12281 (2020)

P. Scholl et al., Programmable quantum simulation of 2D antiferromagnets with hundreds of Rydberg atoms, arXiv:2012.12268 (2020)

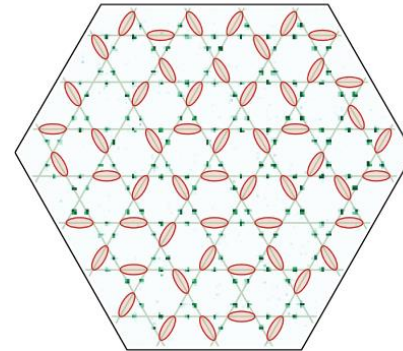
D. Bluvstein et al., Controlling quantum many-body dynamics in driven Rydberg atom arrays, Science 371, 1355 (2021)

Outline

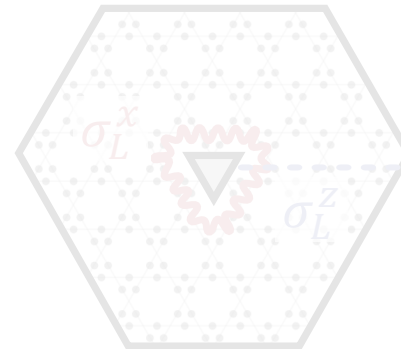
Programmable arrays
of Rydberg atoms



Quantum spin liquid phase
on a frustrated lattice



First steps towards a
topological qubit



Quantum spin liquids

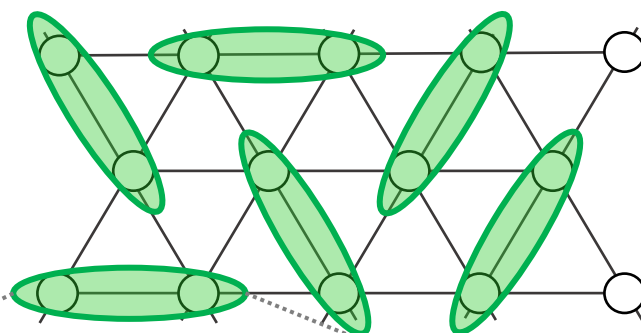
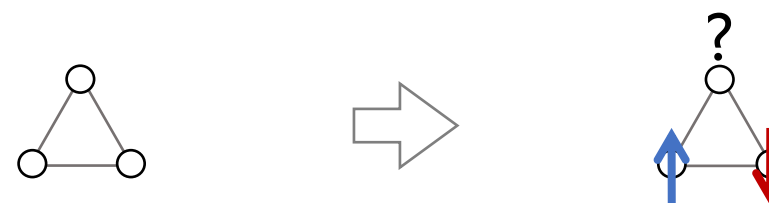
P. W. Anderson, *Science* 235 (1987)

The Resonating Valence Bond State in La_2CuO_4 and Superconductivity

P. W. ANDERSON

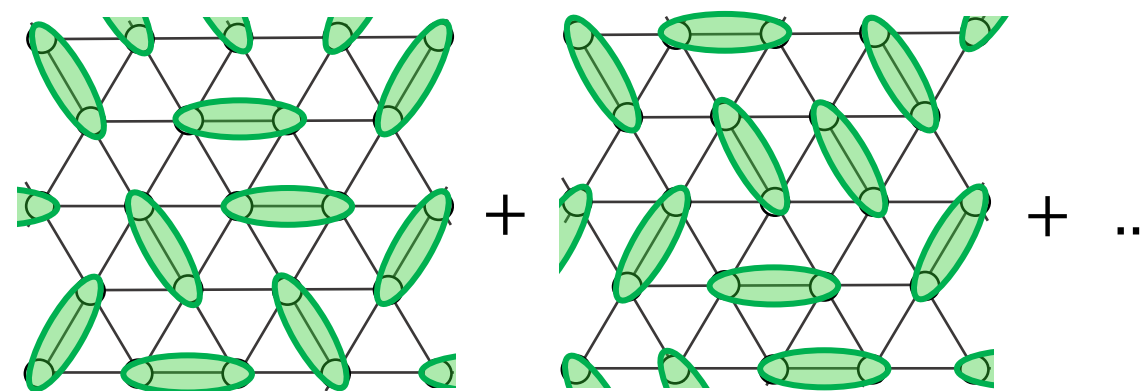
The oxide superconductors, particularly those recently discovered that are based on La_2CuO_4 , have a set of peculiarities that suggest a common, unique mechanism: they tend in every case to occur near a metal-insulator transition into an odd-electron insulator with peculiar magnetic properties. This insulating phase is proposed to be the long-sought “resonating-valence-bond” state or “quantum spin liquid” hypothesized in 1973. This insulating magnetic phase is favored by low spin, low dimensionality, and magnetic frustration. The preexisting magnetic singlet pairs of the insulating state become charged superconducting pairs when the insulator is doped sufficiently strongly. The mechanism for superconductivity is hence predominantly electronic and magnetic, although weak phonon interactions may favor the state. Many unusual properties are predicted, especially of the insulating state.

spin 1/2 particles with AF interactions
on a frustrated lattice



$$\frac{1}{\sqrt{2}} \left(\begin{array}{c} \uparrow \quad \downarrow \\ \downarrow \quad \uparrow \end{array} \right) - \begin{array}{c} \downarrow \quad \uparrow \\ \uparrow \quad \downarrow \end{array} \right)$$

entangled pair
(dimer)



resonating valence bond (RVB) state

Quantum spin liquids

P. W. Anderson, *Science* 235 (1987)

The Resonating Valence Bond State in La_2CuO_4 and Superconductivity

9000+ citations

P. W. ANDERSON

The oxide superconductors, particularly those recently discovered that are based on La_2CuO_4 , have a set of peculiarities that suggest a common, unique mechanism: they tend in every case to occur near a metal-insulator transition into an odd-electron insulator with peculiar magnetic properties. This insulating phase is proposed to be the long-sought “resonating-valence-bond” state or “quantum spin liquid” hypothesized in 1973. This insulating magnetic phase is favored by low spin, low dimensionality, and magnetic frustration. The preexisting magnetic singlet pairs of the insulating state become charged superconducting pairs when the insulator is doped sufficiently strongly. The mechanism for superconductivity is hence predominantly electronic and magnetic, although weak phonon interactions may favor the state. Many unusual properties are predicted, especially of the insulating state.

- no spatial order
- topological order
- long-range quantum entanglement
- anyonic non-local excitations
- emergent gauge fields
- robust ground state degeneracy
- link with high- T_c superconductivity
- application to fault-tolerant quantum computing → toric code

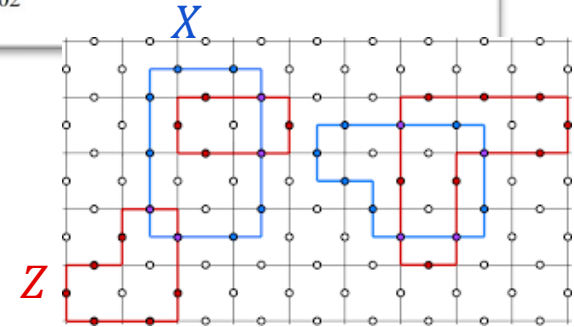
Fault-tolerant quantum computation by anyons

A.Yu. Kitaev*

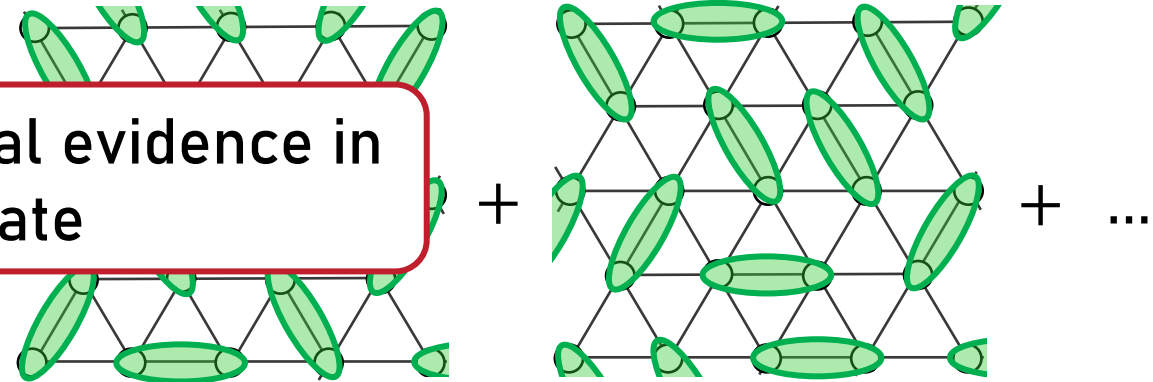
5200+ citations

L.D. Landau Institute for Theoretical Physics, 117940, Kosygina St. 2, Germany

Received 20 May 2002



No conclusive experimental evidence in any system to date



resonating valence bond (RVB) state

Towards topological phases in frustrated lattices of Rydberg atoms

Subir Sachdev



Rhine Samajdar



Wen Wei Ho

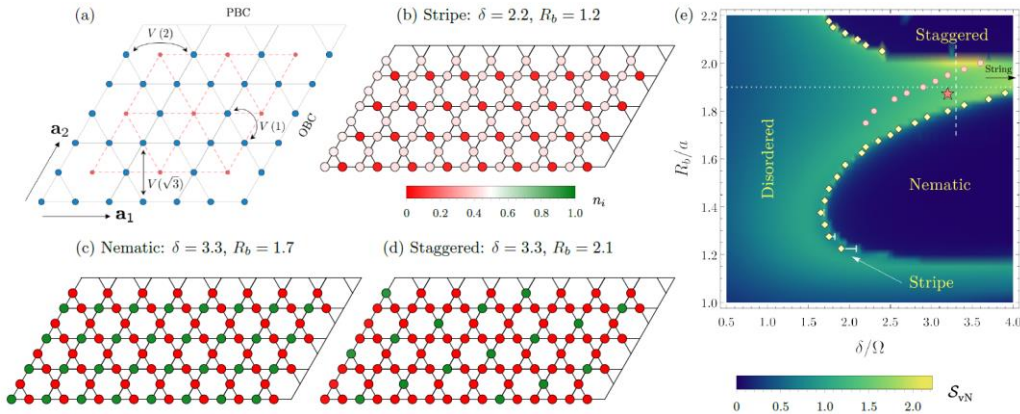


Hannes Pichler

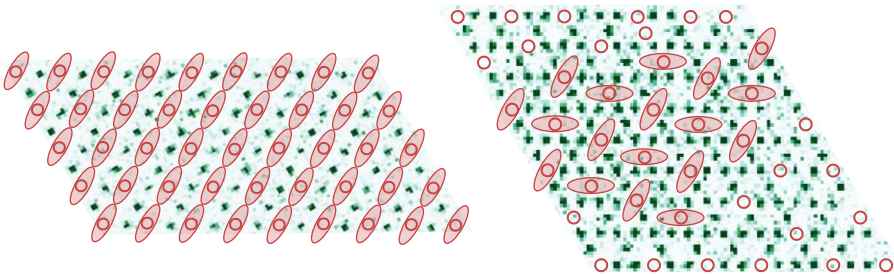


Quantum phases of Rydberg atoms on a kagome lattice

Rhine Samajdar^{a,1}, Wen Wei Ho^{a,b}, Hannes Pichler^{c,d}, Mikhail D. Lukin^a, and Subir Sachdev^a



Emergent solid and liquid 'dimer' phases in frustrated lattices



Ashvin Vishwanath

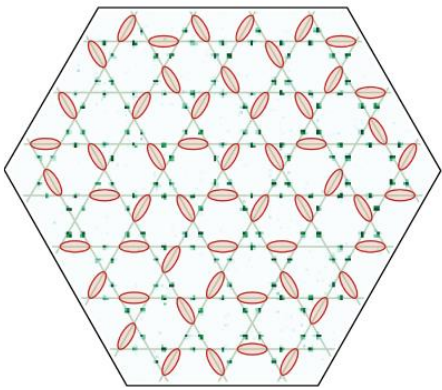
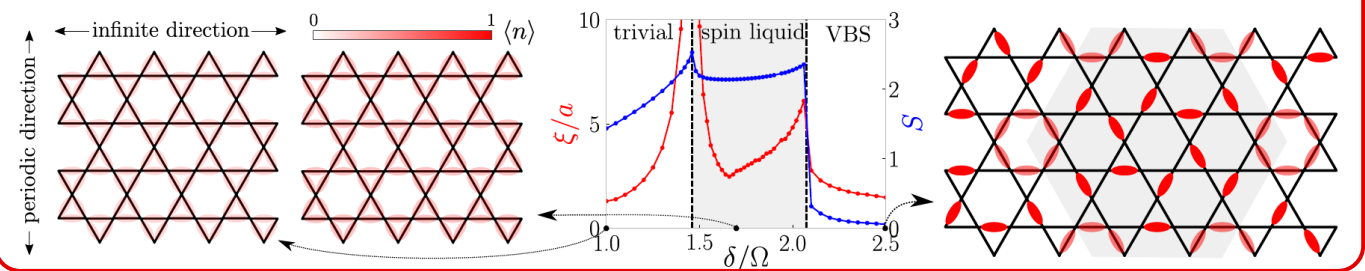


Ruben Verresen

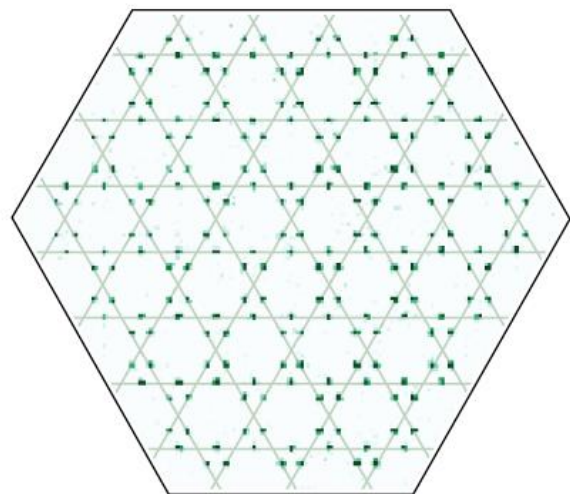


Prediction of Toric Code Topological Order from Rydberg Blockade

Ruben Verresen, Mikhail D. Lukin, and Ashvin Vishwanath
Department of Physics, Harvard University, Cambridge, MA 02138, USA
(Dated: November 26, 2020)



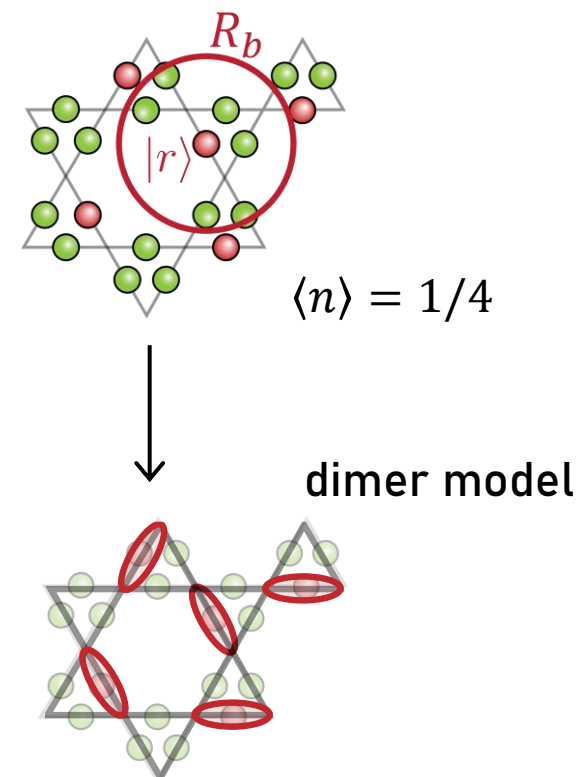
Rydberg atoms on a ruby lattice – dimer model



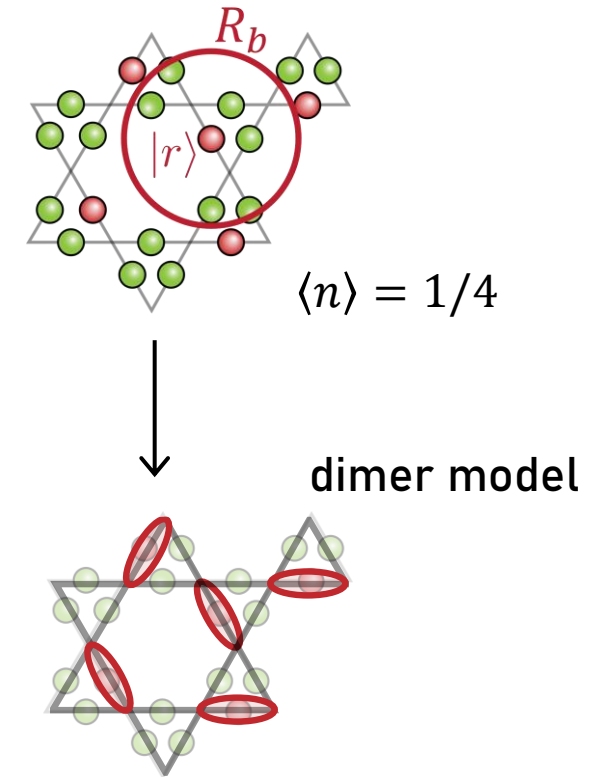
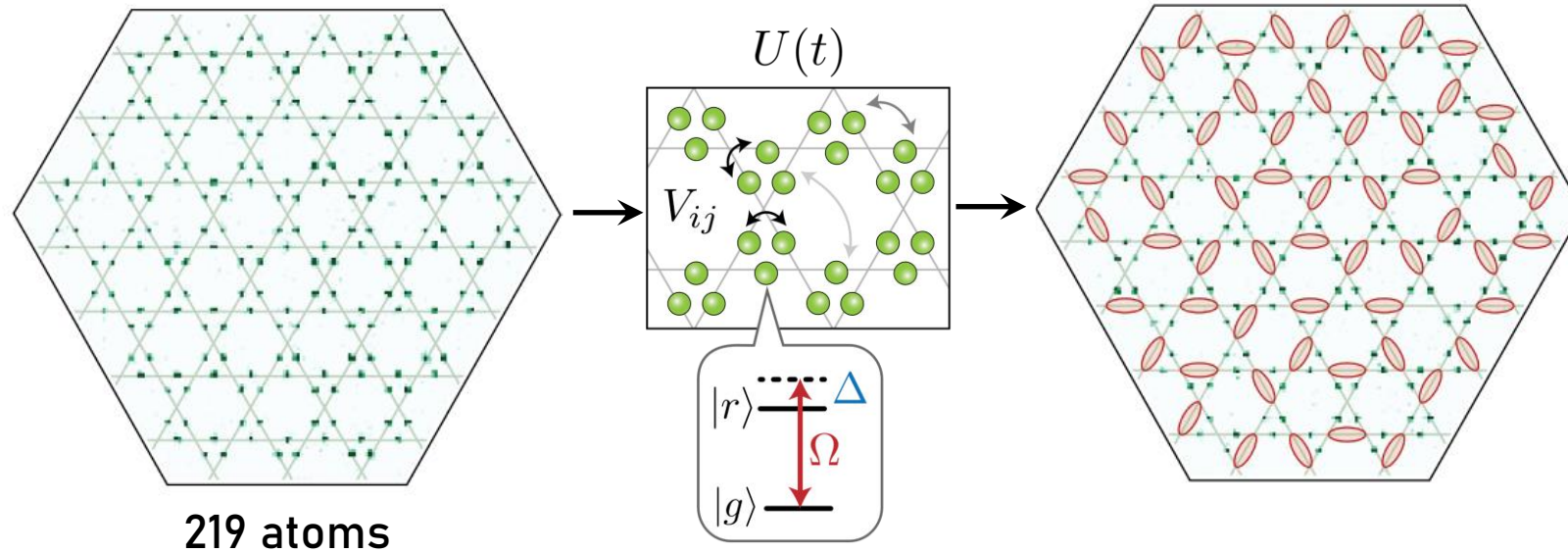
219 atoms

quantum spin liquid state: superposition of all dimer coverings

$$|\psi_{QSL}\rangle = \left| \begin{array}{c} \text{dimer covering 1} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 2} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 3} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 4} \end{array} \right\rangle + \dots \rightarrow \begin{array}{c} \text{RVB state} \\ \mathbb{Z}_2 \text{ topological order} \end{array}$$



Rydberg atoms on a ruby lattice – dimer model

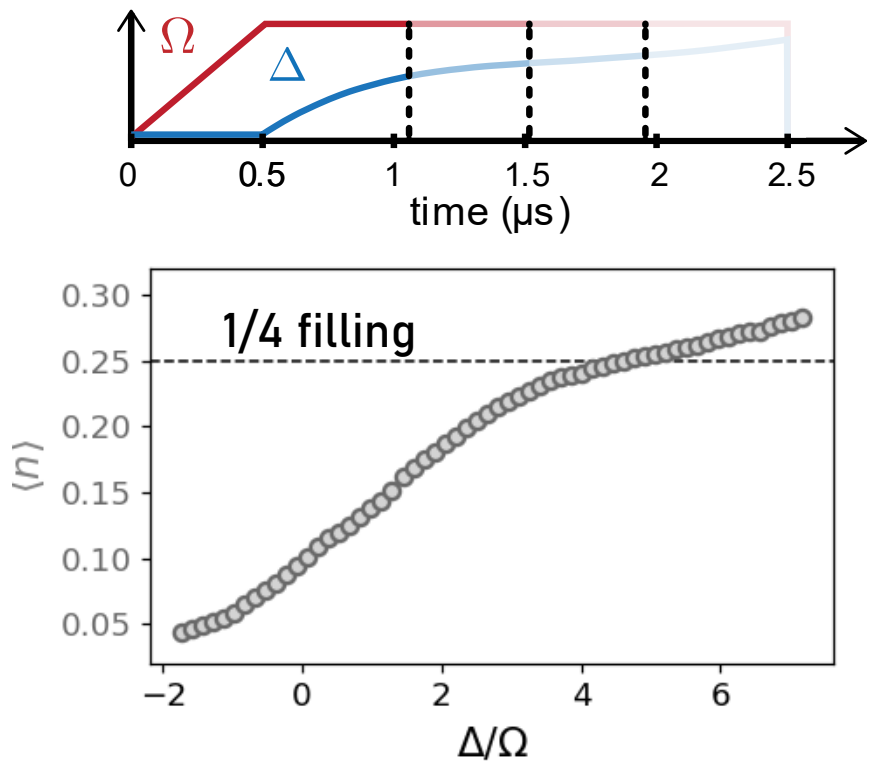


quantum spin liquid state: superposition of all dimer coverings

$$|\psi_{QSL}\rangle = \left| \begin{array}{c} \text{dimer covering 1} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 2} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 3} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 4} \end{array} \right\rangle + \dots \rightarrow \text{RVB state}$$

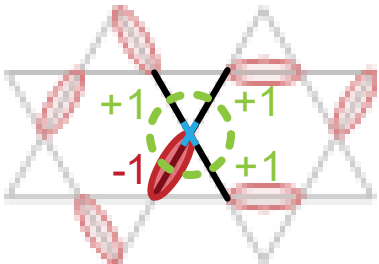
$\mathbb{Z}_2$ topological order

Quasi-adiabatic preparation of a dimer phase



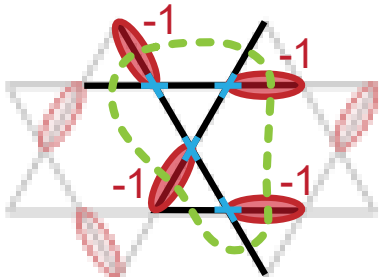
diagonal string operator Z :

$$\langle Z \rangle = -1$$

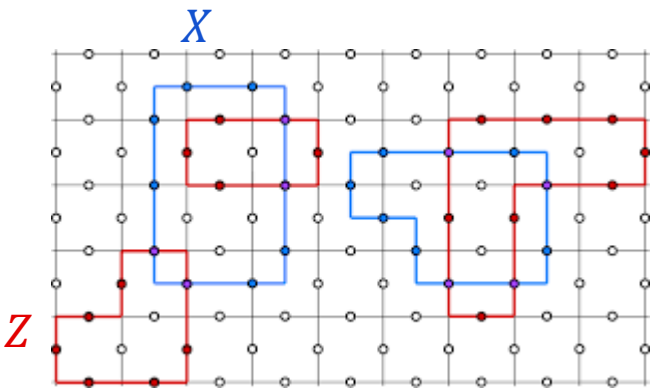


parity of dimers along a string

$$\langle Z \rangle = (-1)^{\# \text{ enclosed vertices}}$$

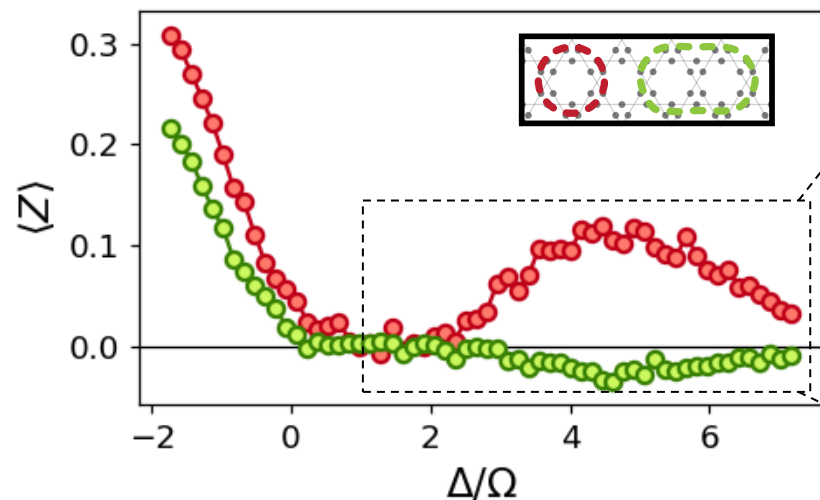
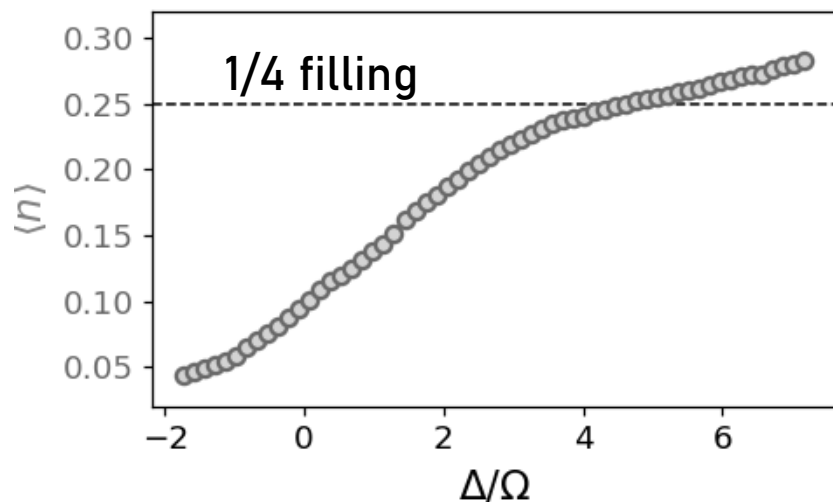
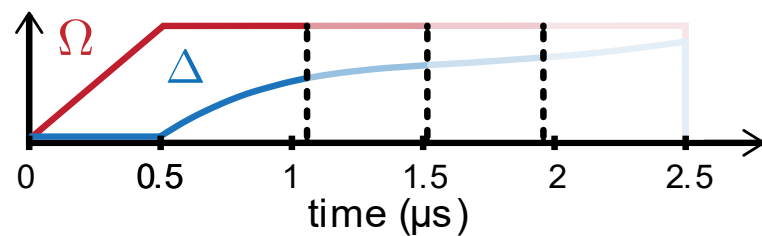


topological string operators associated with a $\mathbb{Z}_2$ quantum spin liquid (toric code)

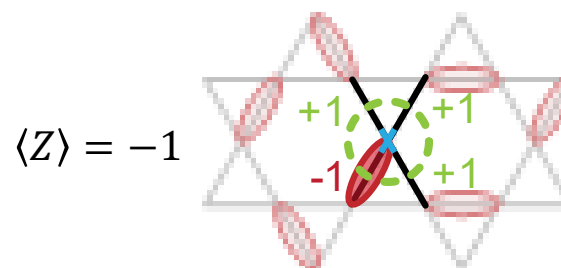


Quasi-adiabatic preparation of a dimer phase

G. Semeghini et al., arXiv:2104.04119 (2021)

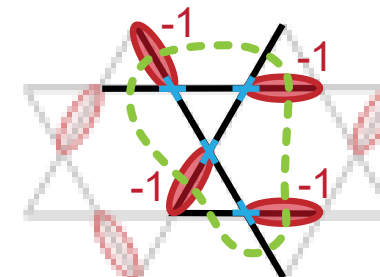


diagonal string operator Z :

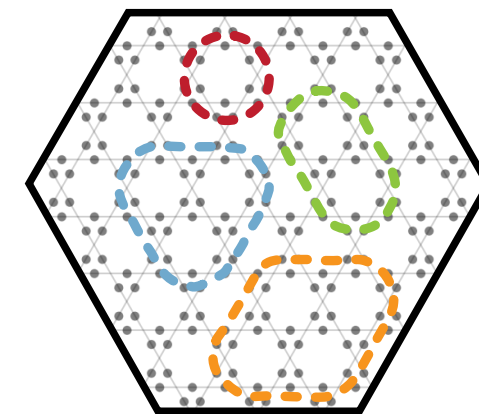
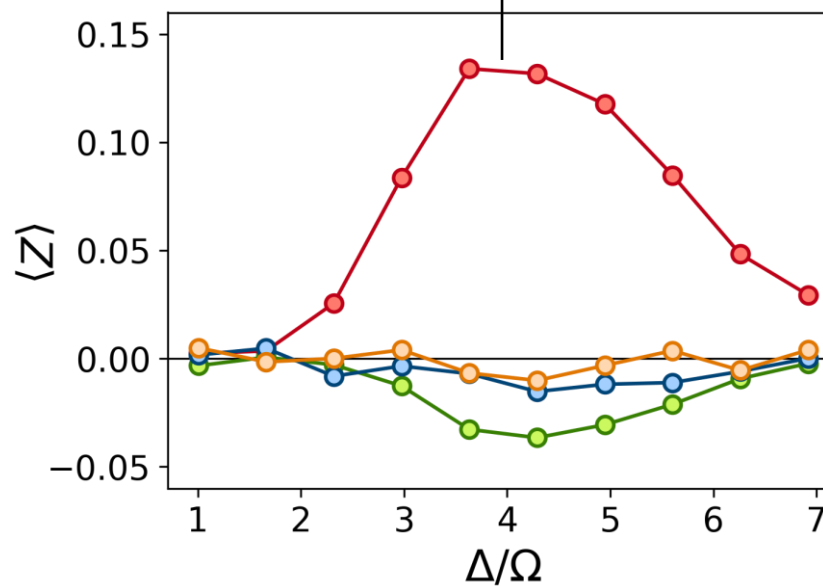


parity of dimers along a string

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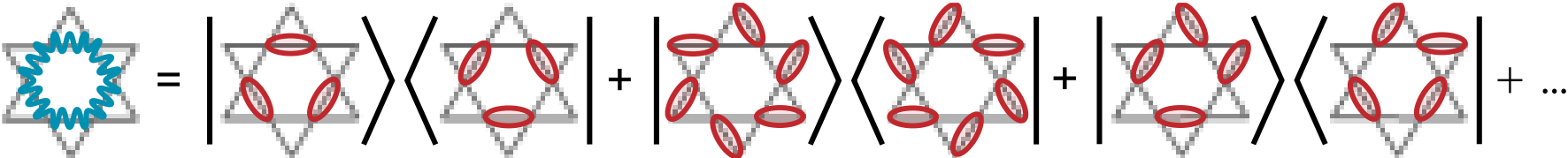
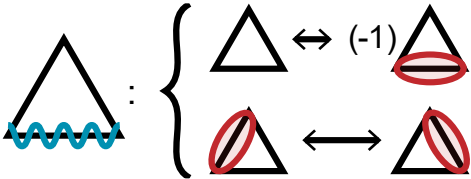


transition into an approximate dimer phase



Probing coherence: off-diagonal string operator

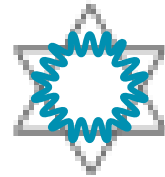
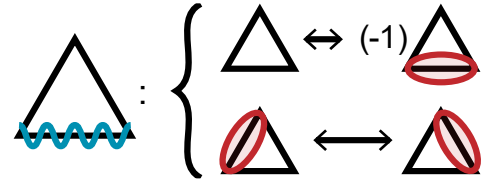
off-diagonal string operator X :



$$\langle \text{star with blue wavy line} \rangle > 0 \rightarrow \text{coherence between dimer coverings}$$

Probing coherence: **off-diagonal string operator**

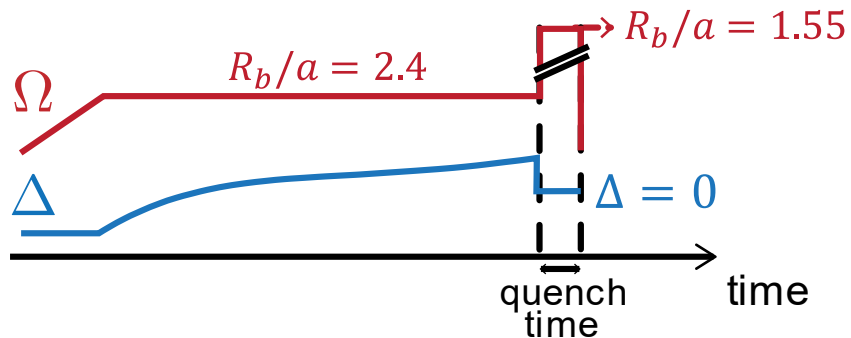
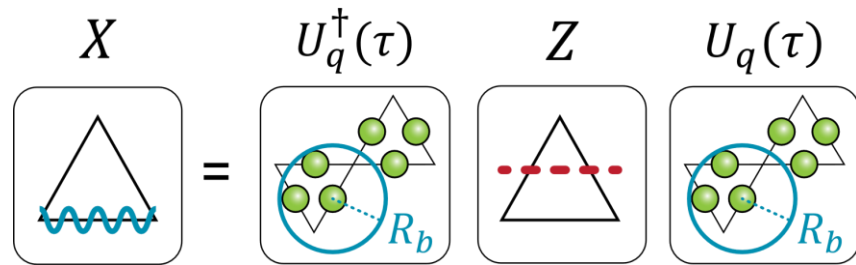
off-diagonal string operator X :



$$= \left| \begin{array}{c} \text{Star with 3 red ellipses} \end{array} \right\rangle \left\langle \begin{array}{c} \text{Star with 3 red ellipses} \end{array} \right| + \left| \begin{array}{c} \text{Star with 4 red ellipses} \end{array} \right\rangle \left\langle \begin{array}{c} \text{Star with 4 red ellipses} \end{array} \right| + \left| \begin{array}{c} \text{Star with 5 red ellipses} \end{array} \right\rangle \left\langle \begin{array}{c} \text{Star with 5 red ellipses} \end{array} \right| + \dots$$

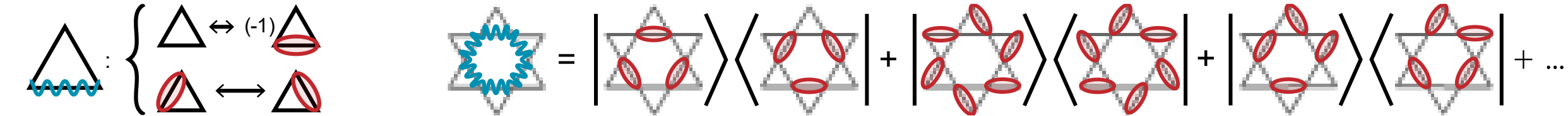
$$\left\langle \begin{array}{c} \text{Star with blue wavy line} \end{array} \right\rangle > 0 \rightarrow \text{coherence between dimer coverings}$$

basis rotation to measure X :

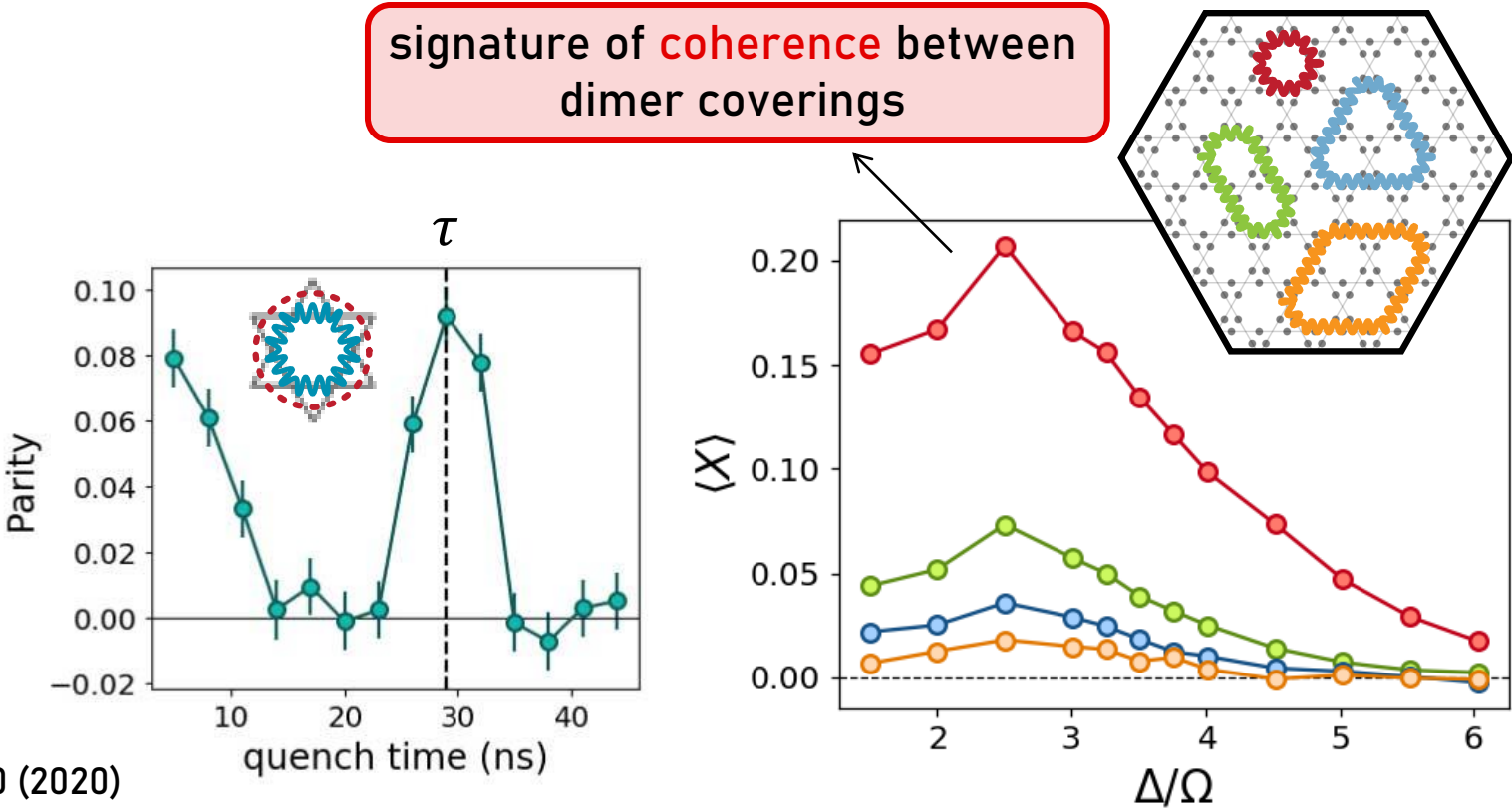
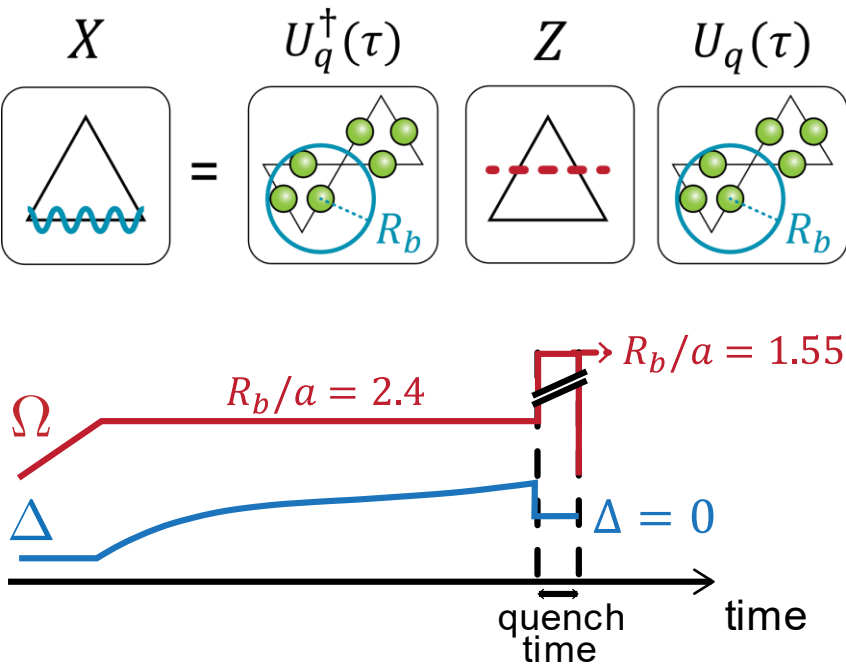


Probing coherence: off-diagonal string operator

off-diagonal string operator X :



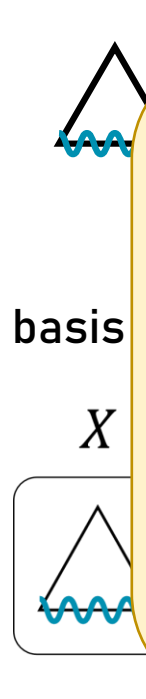
basis rotation to measure X :



Probing coherence: **off-diagonal string operator**

$$\langle \text{star diagram} \rangle > 0 \rightarrow \text{coherence between dimer coverings}$$

off-diagonal string operator X :



basis X

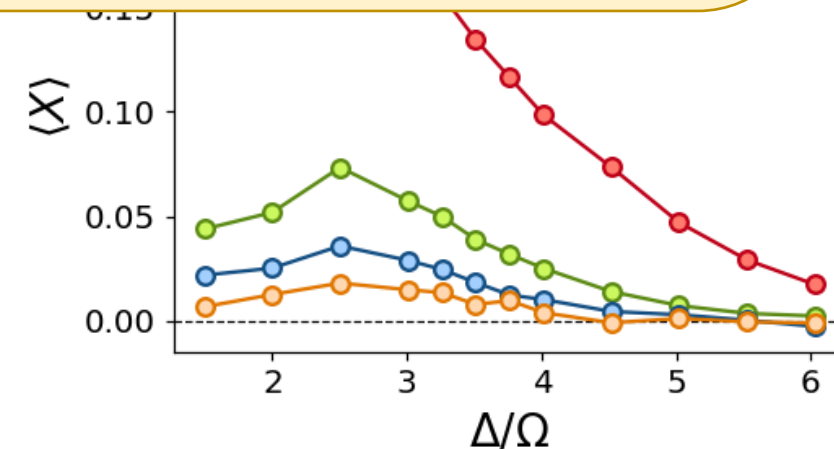
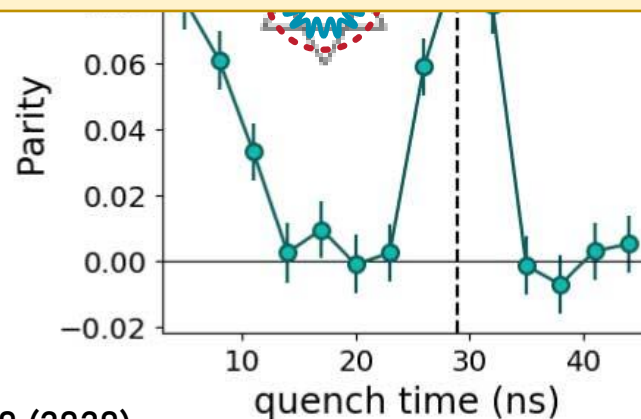
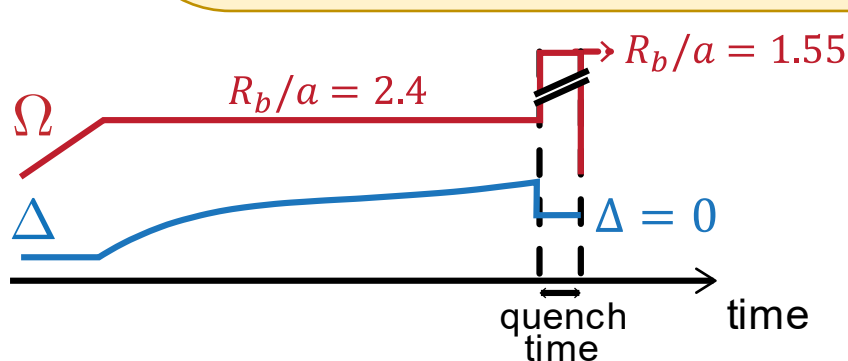
- $|\langle Z_{loop} \rangle| > 0 \rightarrow$ dimer phase
- $\langle X_{loop} \rangle > 0 \rightarrow$ coherent superposition

onset of a **quantum spin liquid** phase!

$$|\psi_{QSL}\rangle = \left| \text{dimer covering 1} \right\rangle + \left| \text{dimer covering 2} \right\rangle + \left| \text{dimer covering 3} \right\rangle + \left| \text{dimer covering 4} \right\rangle + \dots$$

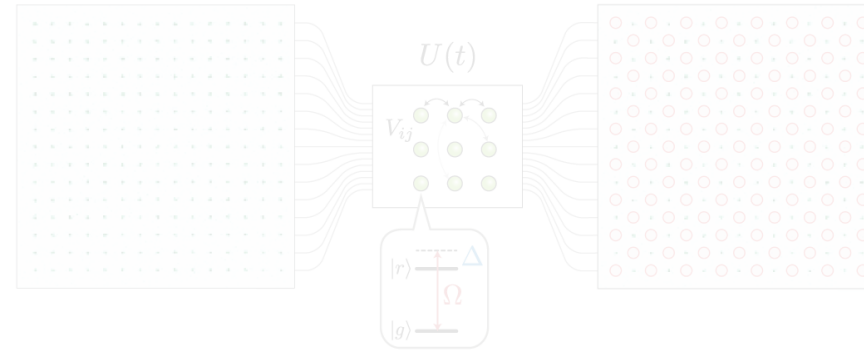


+ ...

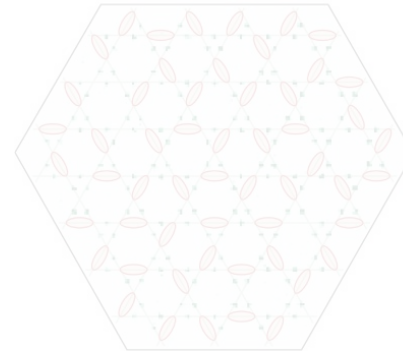


Outline

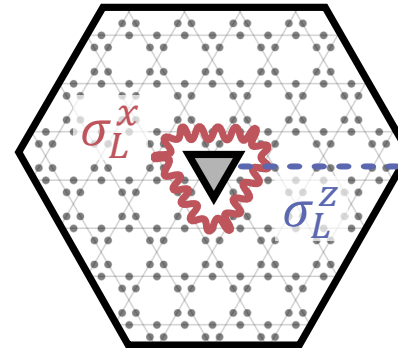
Programmable arrays
of Rydberg atoms



Quantum spin liquid phase
on a frustrated lattice



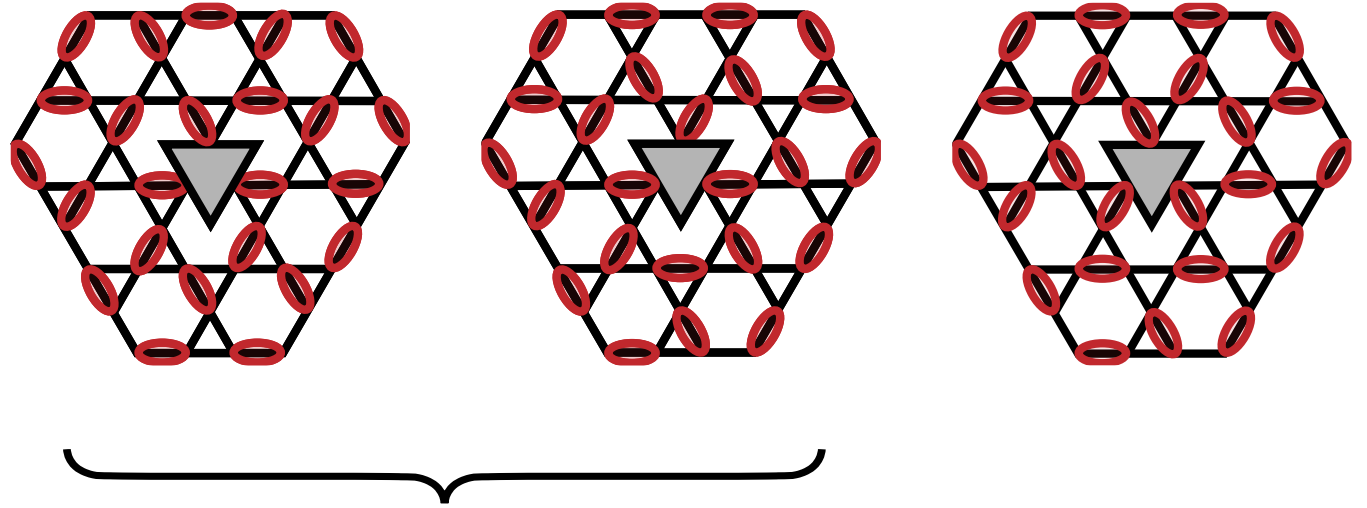
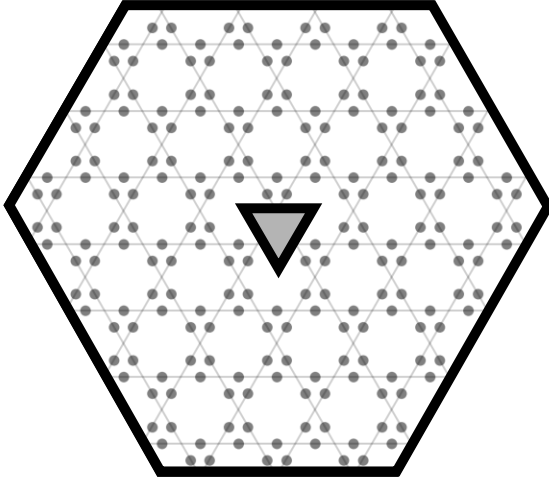
First steps towards a
topological qubit



Towards topological states – array with a hole

two distinct topological sectors

non-trivial topology

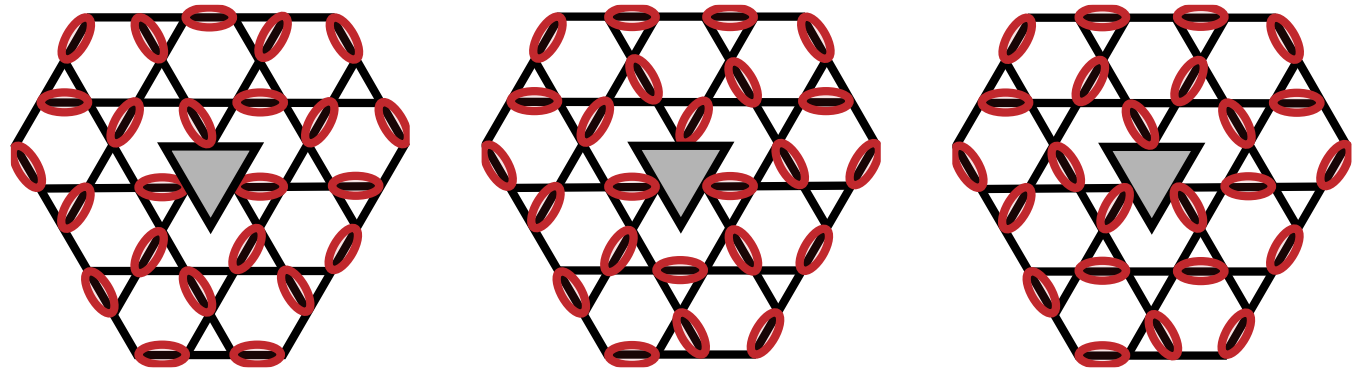
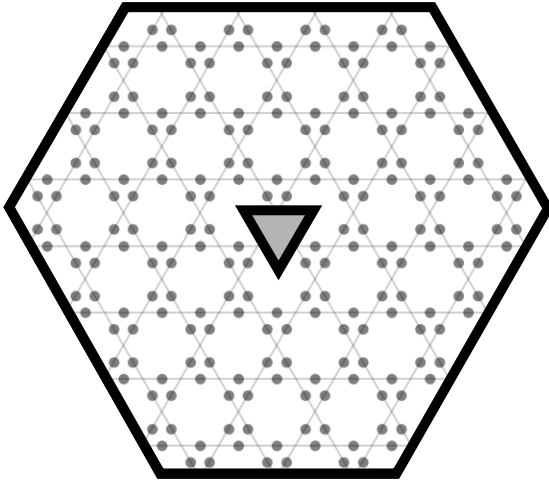


transition graph:
 $|A\rangle \cup |B\rangle$

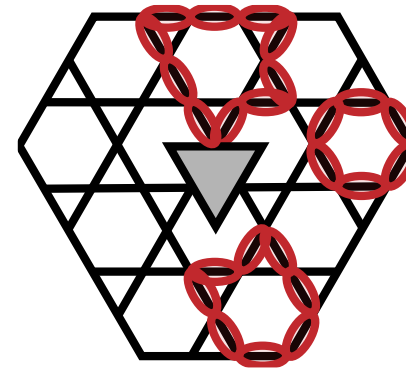
Towards topological states – array with a hole

two distinct topological sectors

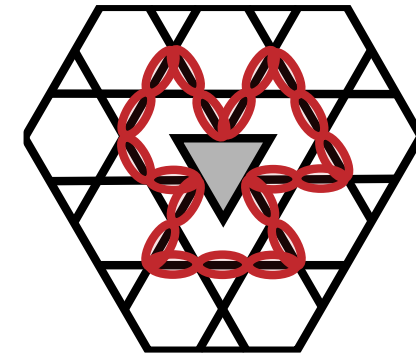
non-trivial topology



transition graph:
 $|A\rangle \cup |B\rangle$

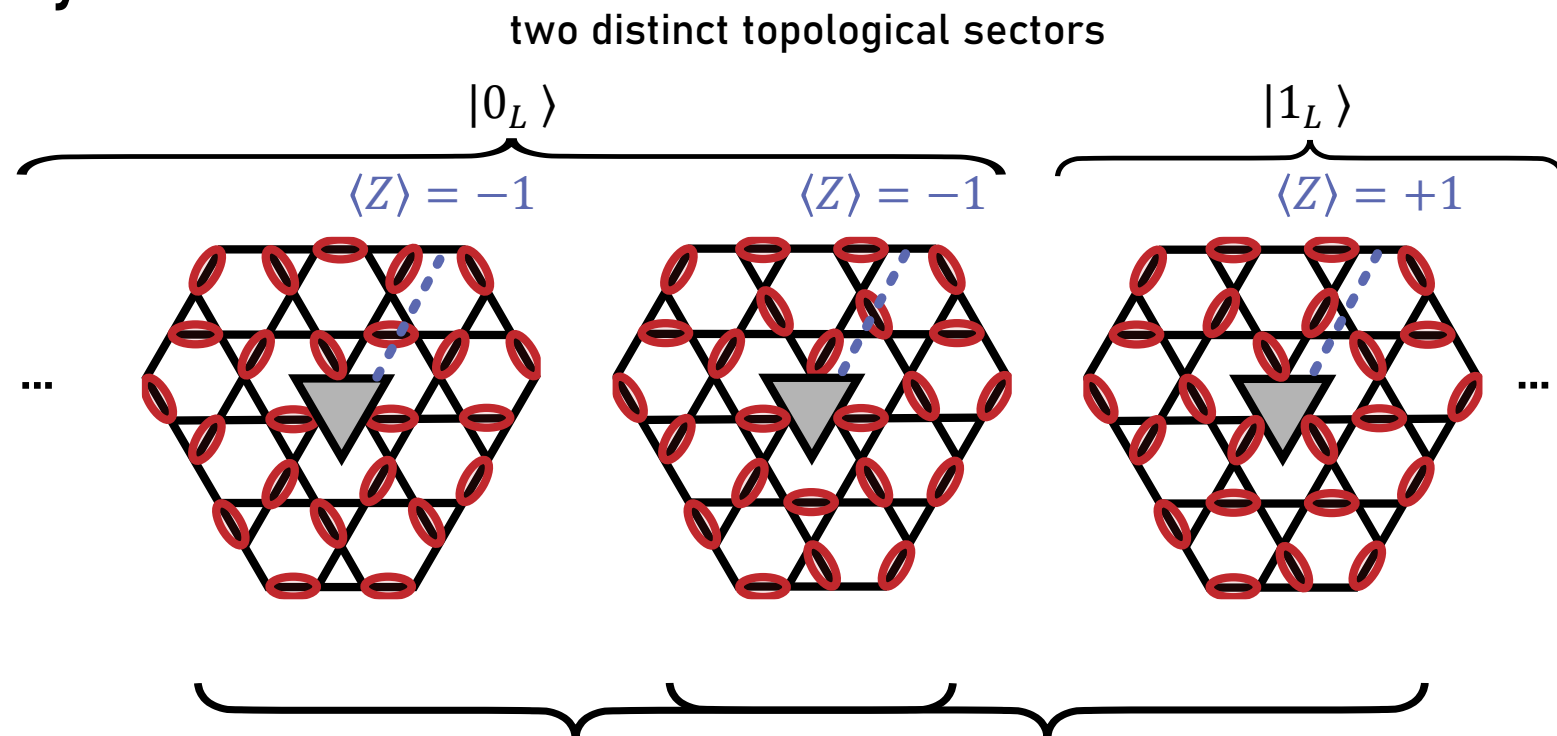
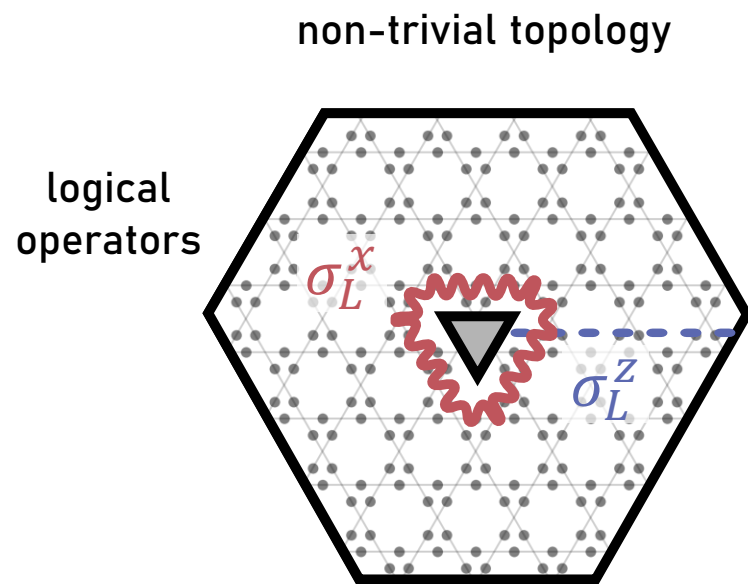


same
topological sector



opposite
topological sectors

Towards topological states – array with a hole



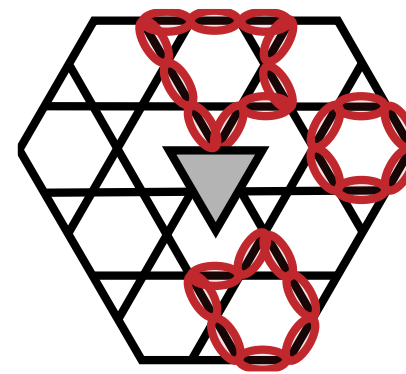
$$|0_L\rangle \longrightarrow \langle \sigma_L^z \rangle = -1$$

$$|1_L\rangle \longrightarrow \langle \sigma_L^z \rangle = +1$$

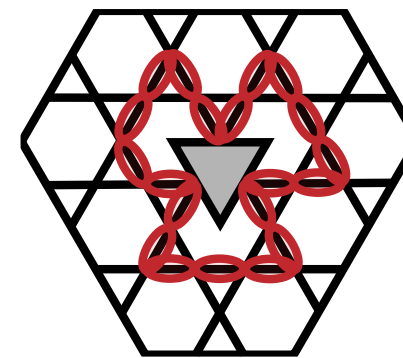
$$|+\rangle \sim \frac{|0_L\rangle + |1_L\rangle}{\sqrt{2}} \longrightarrow \langle \sigma_L^x \rangle = +1$$

$$|-\rangle \sim \frac{|0_L\rangle - |1_L\rangle}{\sqrt{2}} \longrightarrow \langle \sigma_L^x \rangle = -1$$

transition graph:
 $|A\rangle \cup |B\rangle$

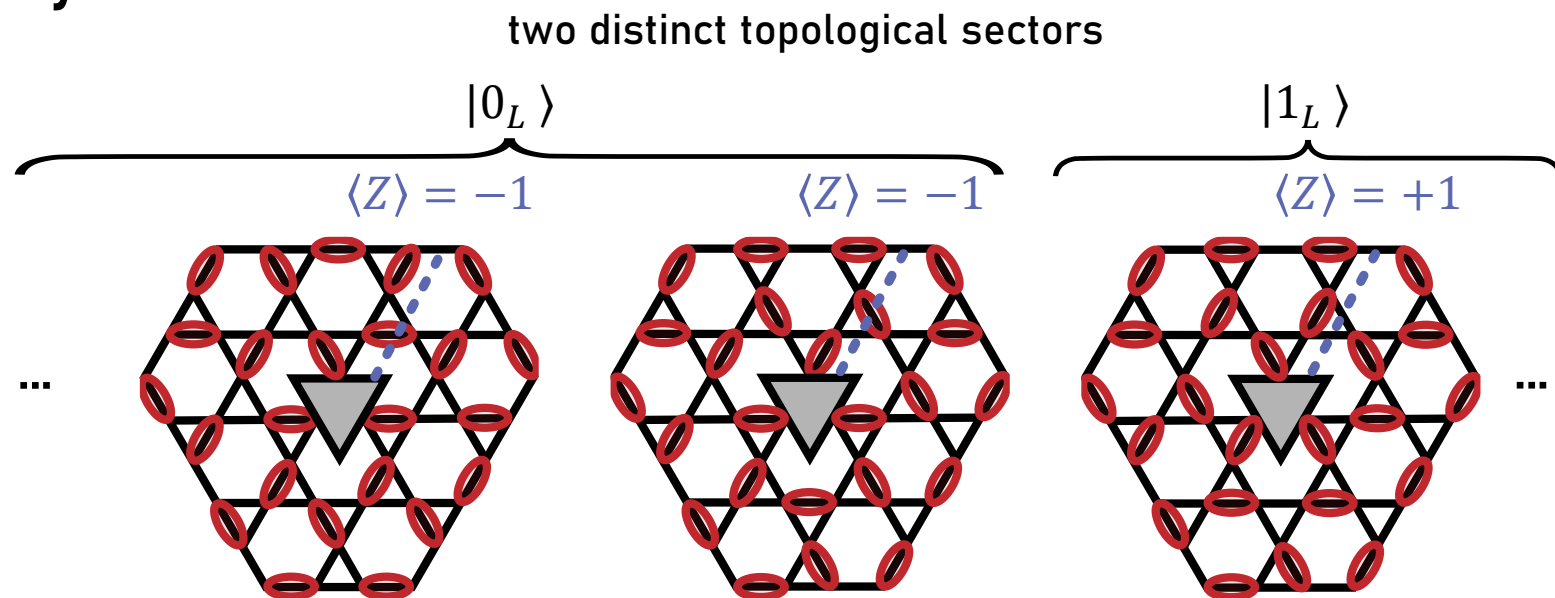
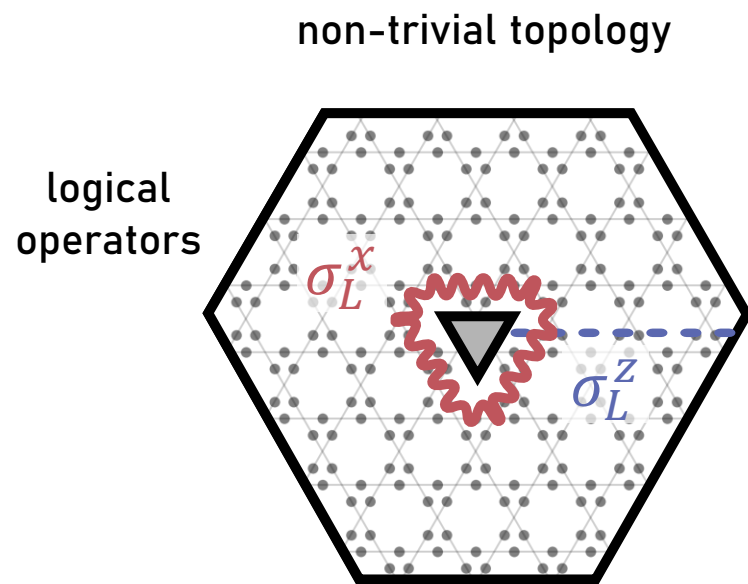


same
topological sector



opposite
topological sectors

Towards topological states – array with a hole

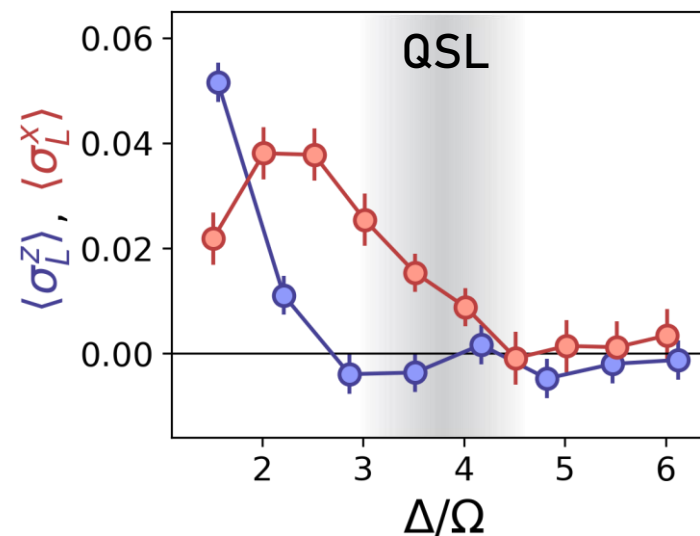


$$|0_L\rangle \longrightarrow \langle \sigma_L^z \rangle = -1$$

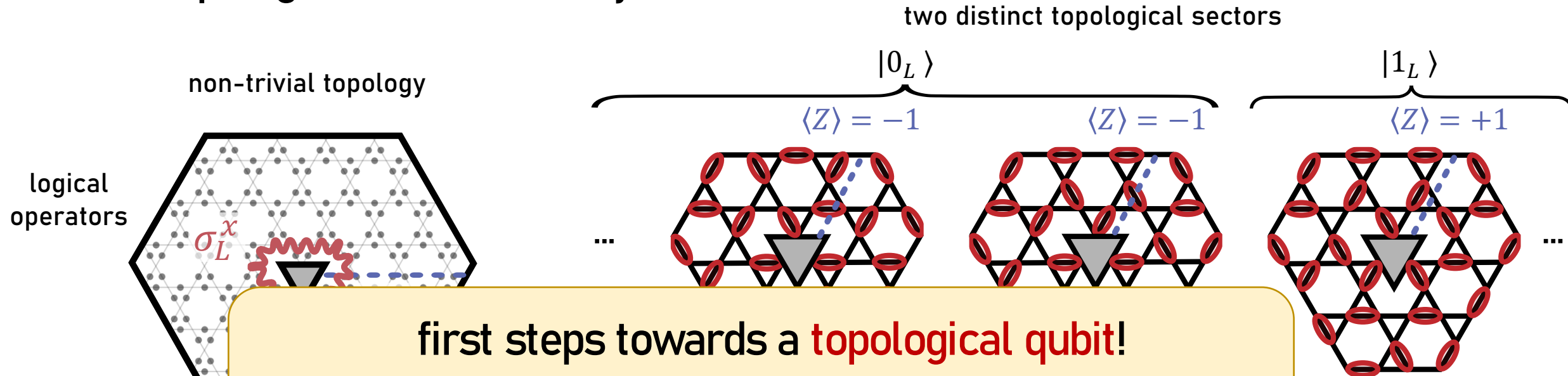
$$|1_L\rangle \longrightarrow \langle \sigma_L^z \rangle = +1$$

$$|+\rangle \sim \frac{|0_L\rangle + |1_L\rangle}{\sqrt{2}} \longrightarrow \langle \sigma_L^x \rangle = +1$$

$$|-\rangle \sim \frac{|0_L\rangle - |1_L\rangle}{\sqrt{2}} \longrightarrow \langle \sigma_L^x \rangle = -1$$



Towards topological states – array with a hole



first steps towards a **topological qubit!**

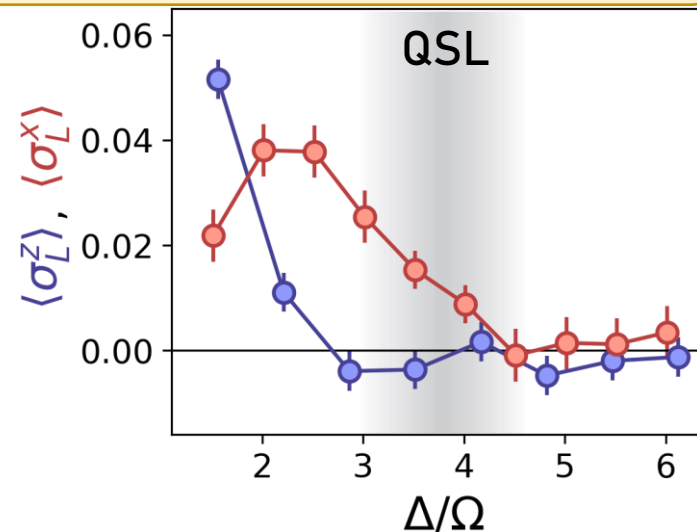
related work from Google Quantum AI:
K. J. Satzinger et al., arXiv:2104.01180 (2021)

$$|0_L\rangle \rightarrow \langle \sigma_L^z \rangle = -1$$

$$|1_L\rangle \rightarrow \langle \sigma_L^z \rangle = +1$$

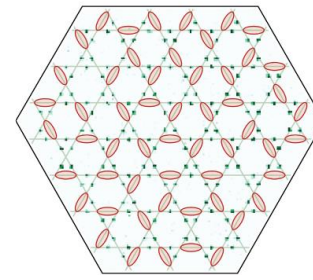
$$|+\rangle \sim \frac{|0_L\rangle + |1_L\rangle}{\sqrt{2}} \rightarrow \langle \sigma_L^x \rangle = +1$$

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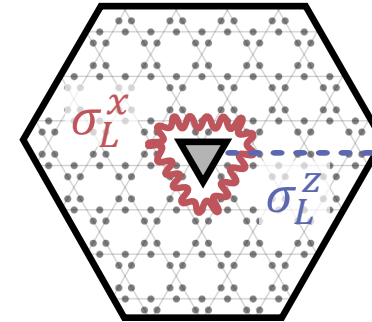


Summary

- 1) Transition into a **dimer phase** with no local order
- 2) Experimental evidence for **coherent superposition** of dimer coverings
- 3) First signatures of **non-trivial topological order**
- 4) Exciting opportunities for experimental exploration of topological states of matter!



$$\text{Star} = \left| \begin{array}{c} \text{triangle} \\ \text{with 3 dimers} \end{array} \right\rangle \left\langle \begin{array}{c} \text{triangle} \\ \text{with 3 dimers} \end{array} \right| + \left| \begin{array}{c} \text{triangle} \\ \text{with 3 dimers} \end{array} \right\rangle \left\langle \begin{array}{c} \text{triangle} \\ \text{with 3 dimers} \end{array} \right| + \dots$$

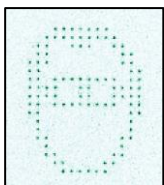


Open questions

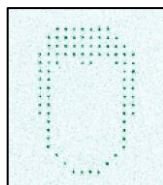
- 1) Origin of the observed QSL: metastable state?
- 2) Improving state preparation and measurement of string operators
- 3) Encoding and manipulating topological qubits
- 4) Implications for topological matter in other systems, e.g. 2D materials



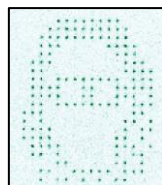
Image: Sci Am (2006)



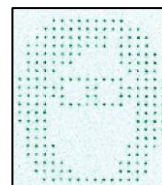
Harry Levine



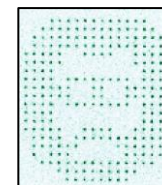
Tout Wang



Sepehr Ebadi



GS



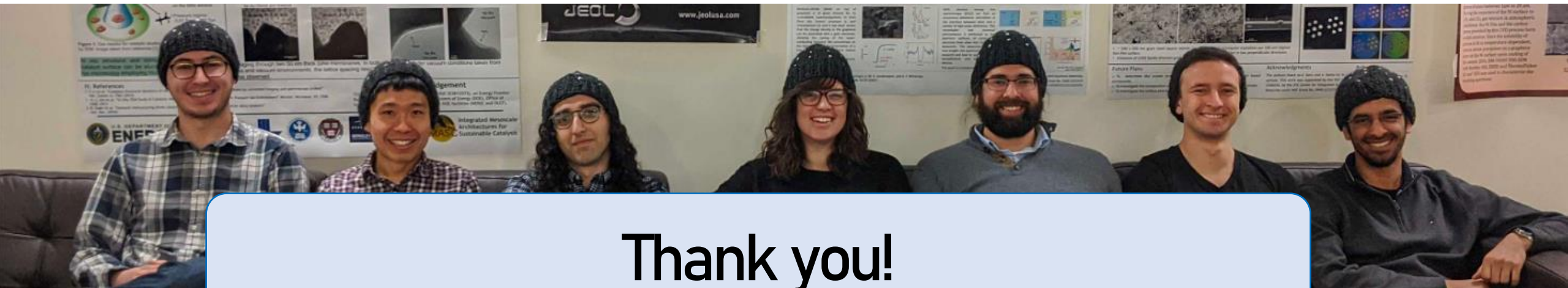
Alex Keesling



Dolev Bluvstein



Ahmed Omran



Thank you!

Pls: M. Lukin



V. Vuletic



M. Greiner



R. Verresen



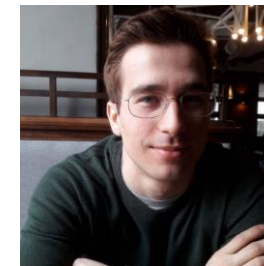
R. Samajdar



W. W. Ho



M. Kalinowski



A. Vishwanath



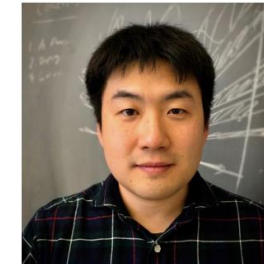
S. Sachdev



H. Pichler



S. Choi



THEORY COLLABORATORS