

# Seeing or listening: magnetoelastic effects in antiferromagnetic textures

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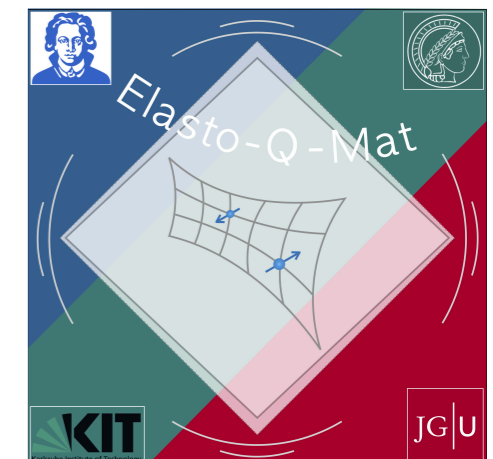


**SPIN+X**  
SFB/TRR 173  
Kaiserslautern • Mainz

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SPICE Seminar



European  
Research  
Council

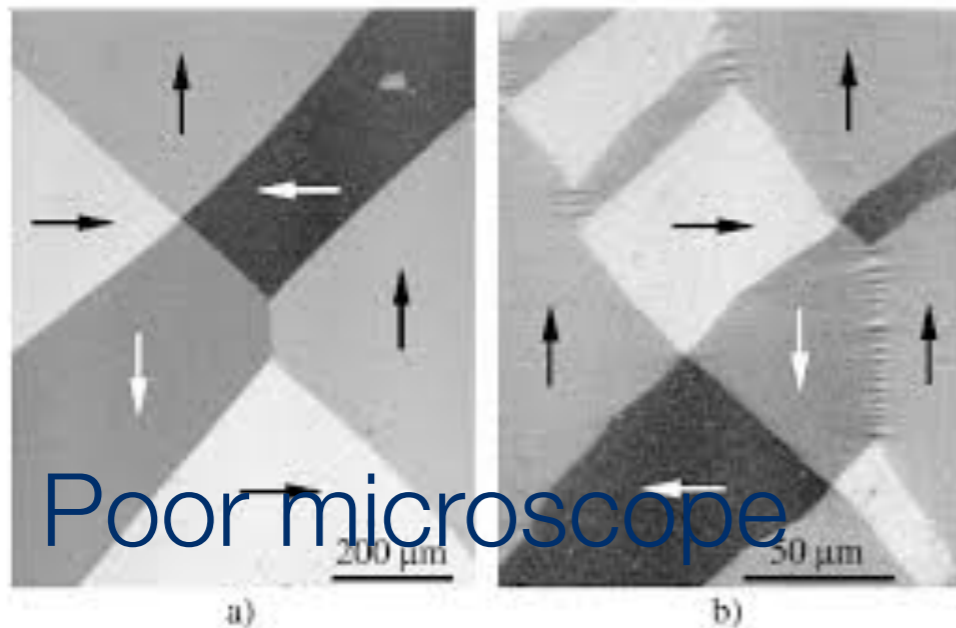


# Ashe Wednesday: the end of having fun

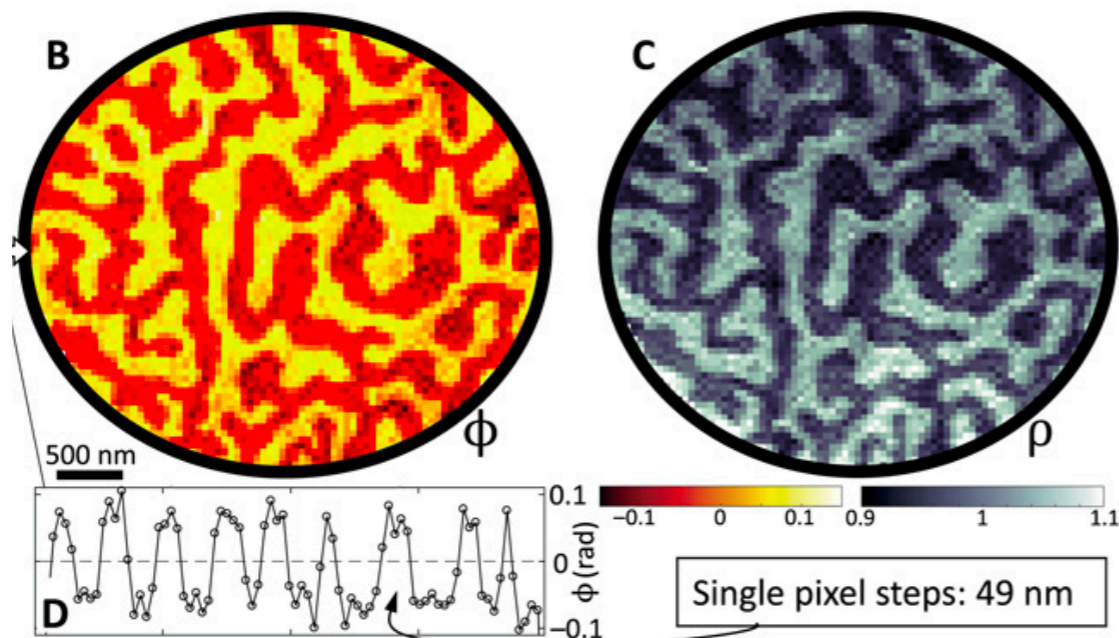




Visualization of  
ferromagnetic domains

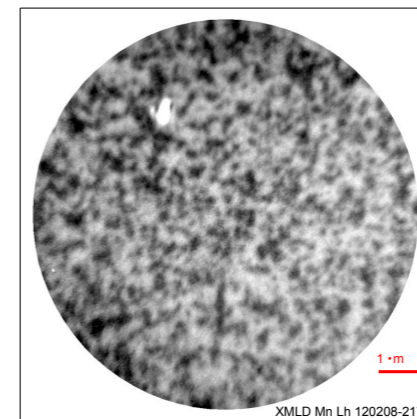


Poor microscope



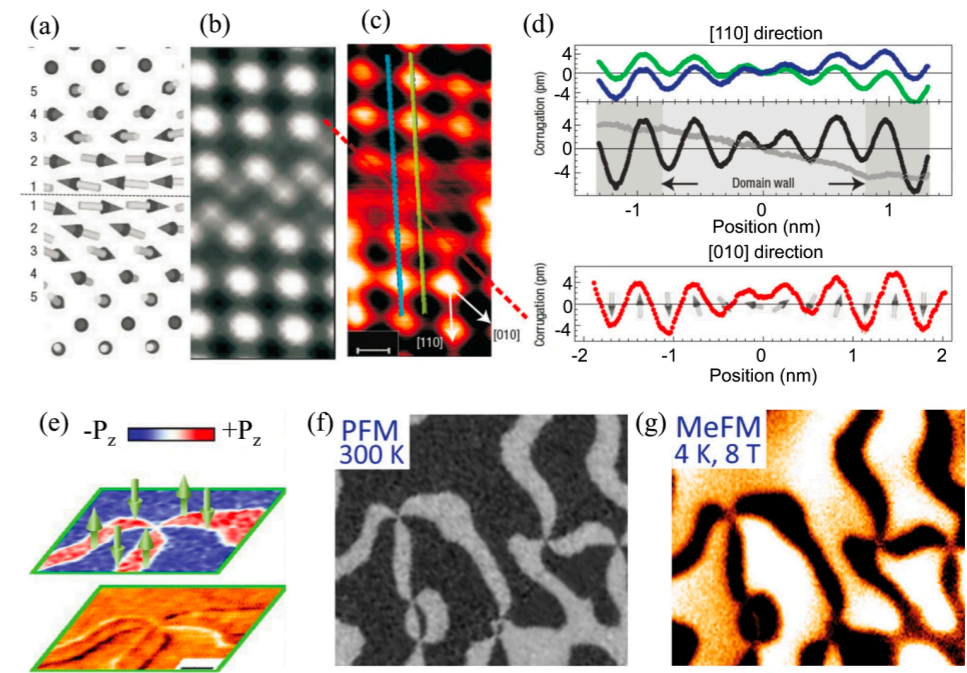
Good microscope

Visualization of  
**anti**ferromagnetic domains



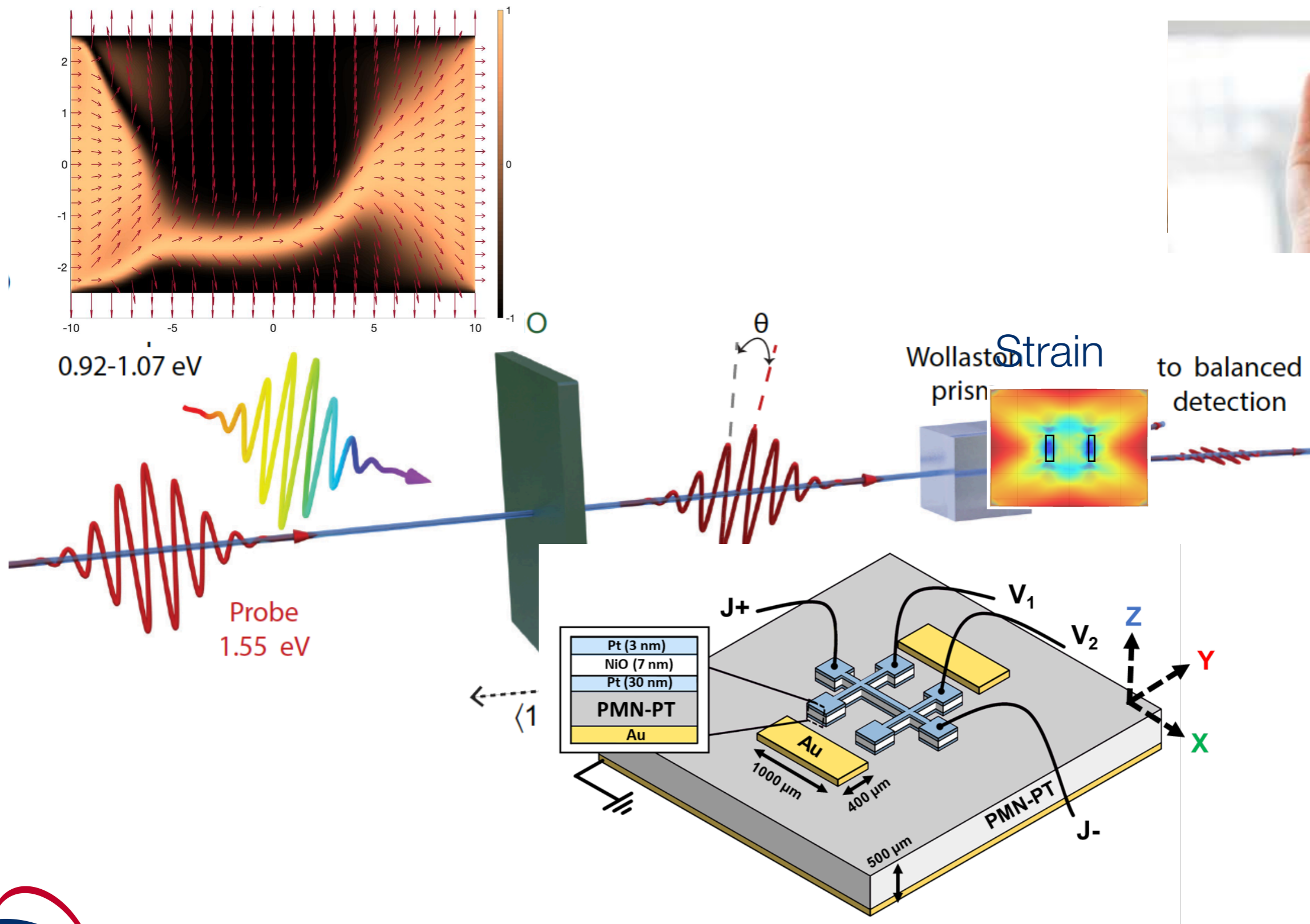
XMLD

SP-STM



S.W. Cheong et al, Quantum Mat. (2020)

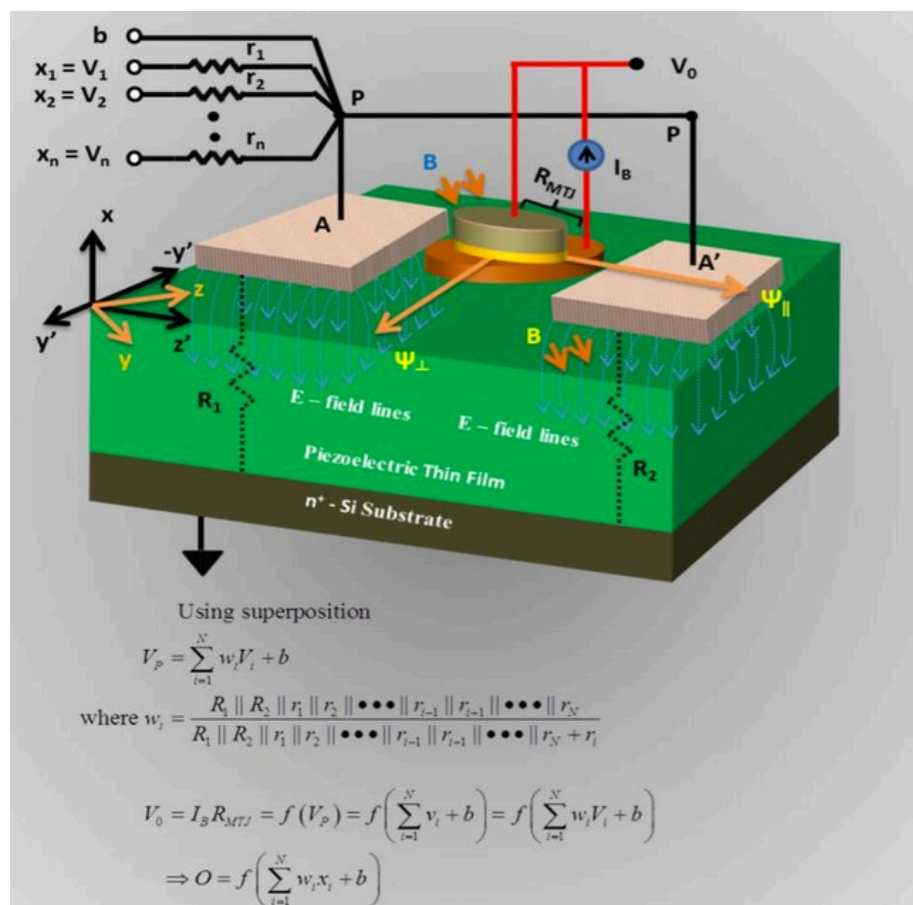
# Seeing or listening?





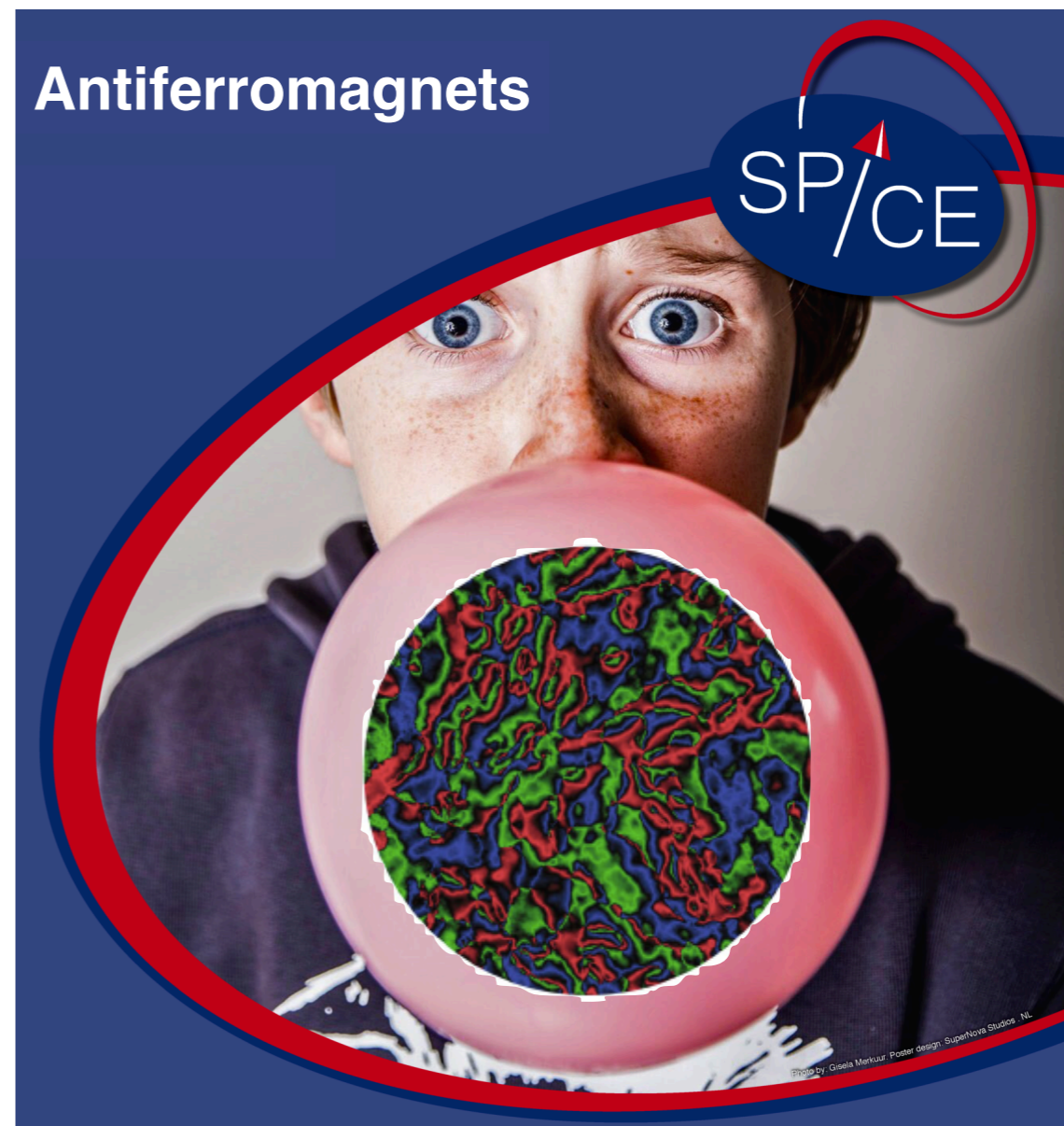
Supriyo Bandyopadhyay

Straintronics and neuromorphic computing



DOI: [10.1109/OJNANO.2020.3011637](https://doi.org/10.1109/OJNANO.2020.3011637)

## Antiferromagnets



F. P. Chmiel et al, Nature Mat. (2018)

- **Strains** affect morphology and dynamics of **antiferromagnetic** textures
- **Antiferromagnets** can be effectively manipulated by **strains**

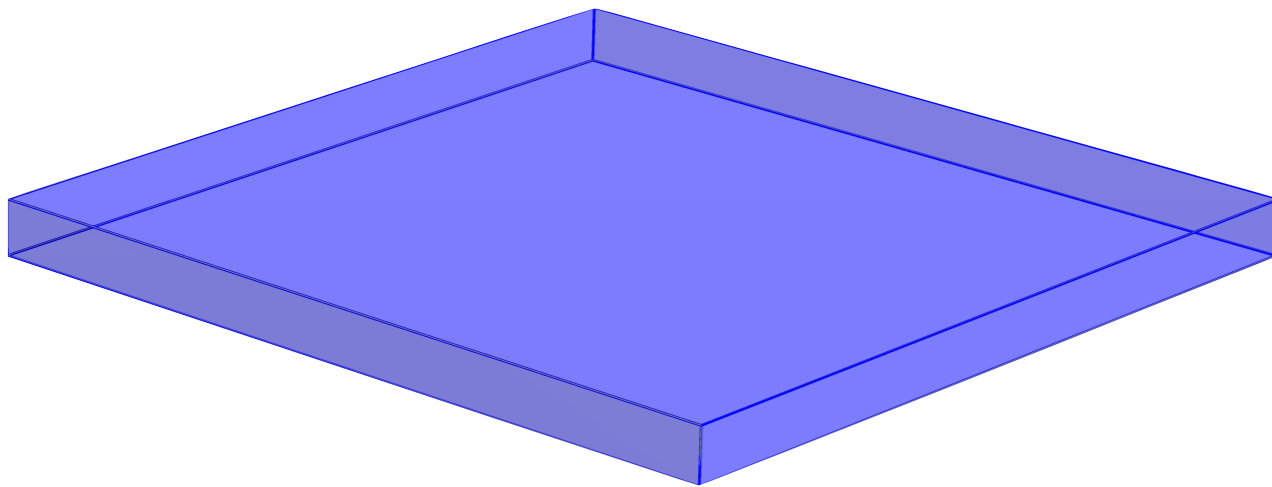




- **Magnetoelasticity: frozen spontaneous strains and domain walls**
- **Incompatibility of strains**
- **Nonequilibrium strains: switching**
- **Equilibrium domain structure: Micr-*a*-magnetics**
- **Conclusions**



Paramagnetic state



$$\hat{u}^{\text{spon}} \propto H_{\text{m-e}} \mathbf{n} \otimes \mathbf{n}$$

$$\hat{\sigma} \equiv - \frac{\partial}{\partial \hat{u}} (w_{\text{elas}} + w_{\text{me}}) = 0$$

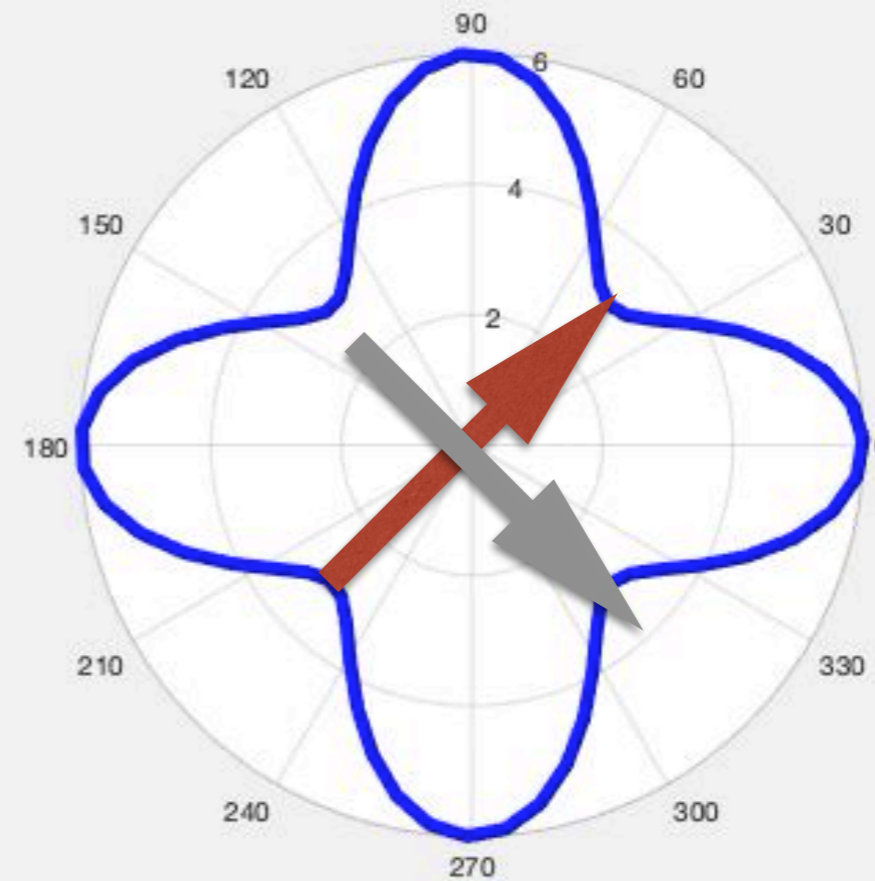
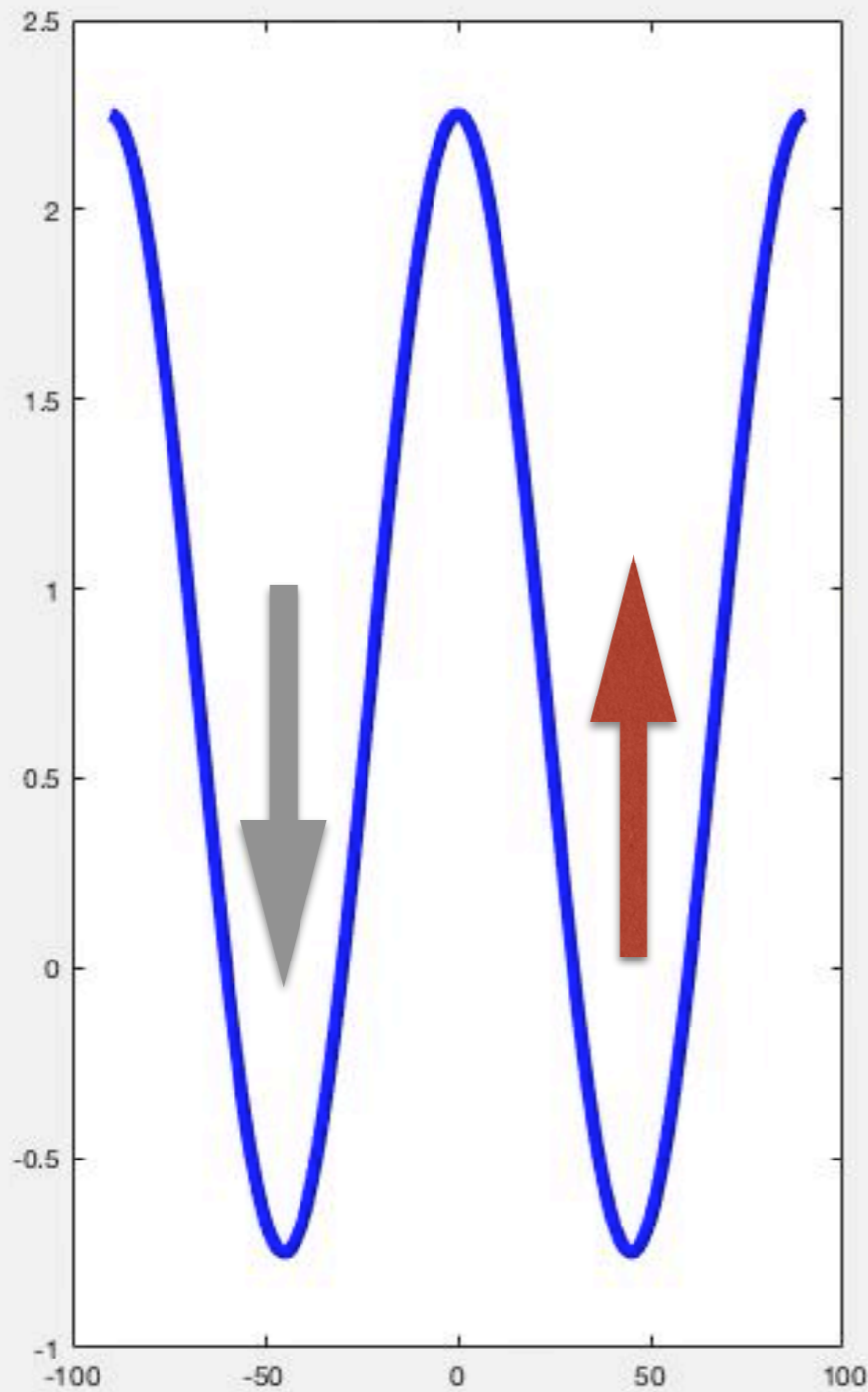
$$\tau_{\text{elas}} \propto L/v_{\text{sound}} \geq 10^{-9} \text{ s}^{-1} \quad \tau_{\text{mag}} \propto 1/\nu_{\text{AFMR}} \propto 10^{-12} \text{ s}^{-1}$$

Frozen spontaneous strains

$$\varepsilon^{\text{spon}} = - \frac{H_{\text{me}} M_s}{\mu} (n_x^2 - n_y^2) |_0$$

$$w_{\text{me}} = H_{\text{me}} M_s \varepsilon^{\text{spon}} (n_x^2 - n_y^2)$$

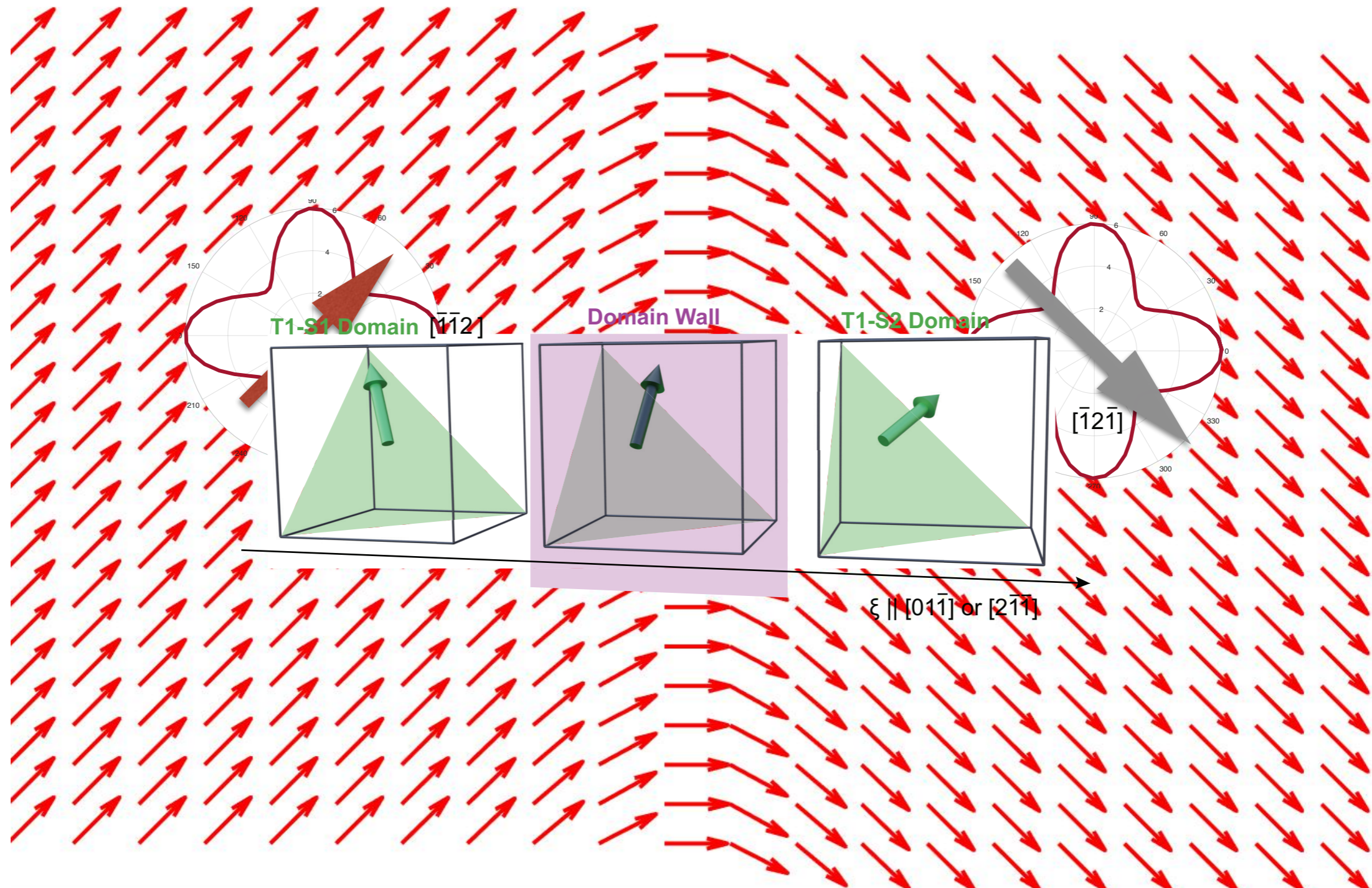
A Borovik-Romanov, 1965  
V. Ozhogin, 1980



$$\omega_{\text{AFMR}} \propto \sqrt{H_{\text{an}} + H_{\text{me}} \epsilon^{\text{spon}}}$$

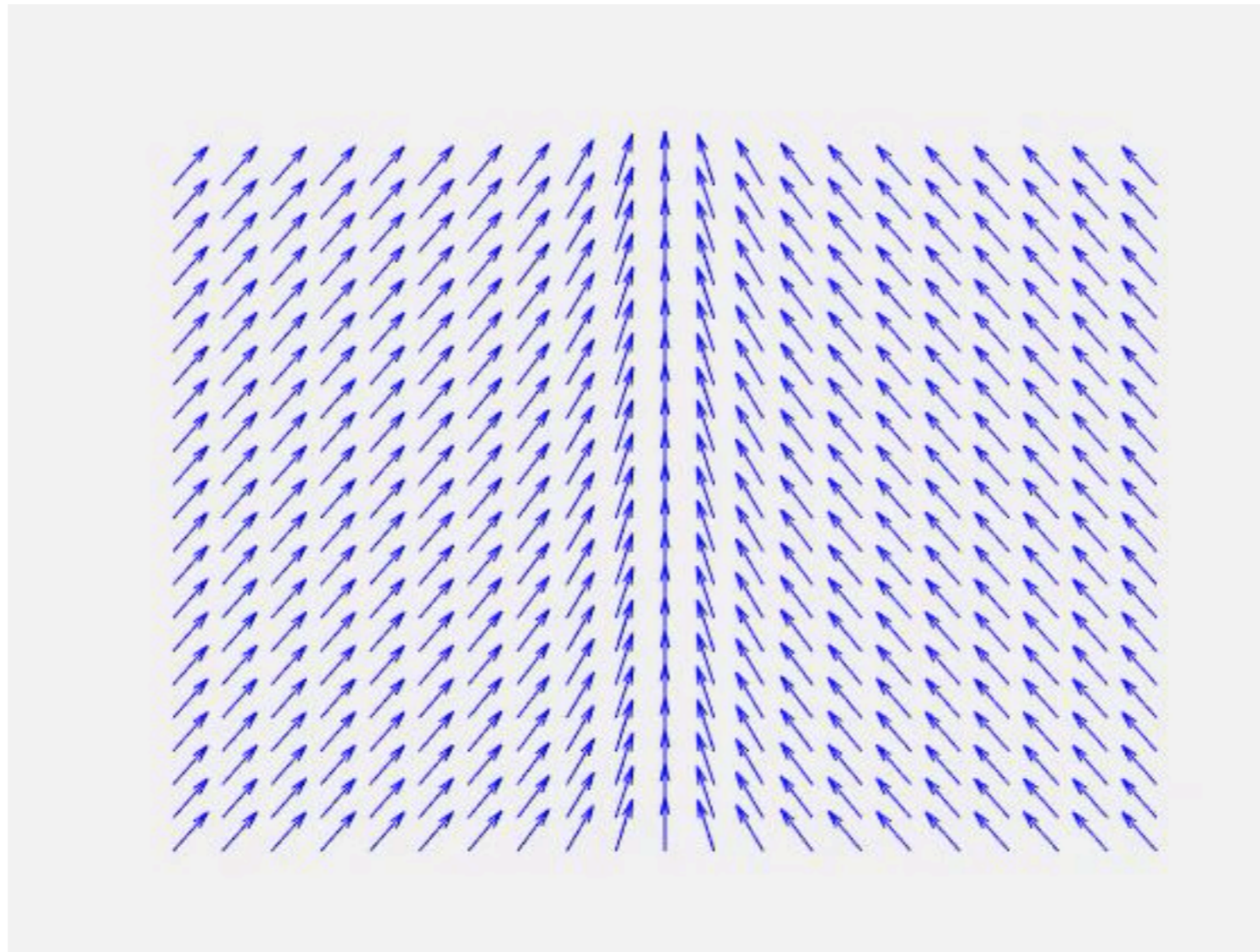


# Magnetic domain wall





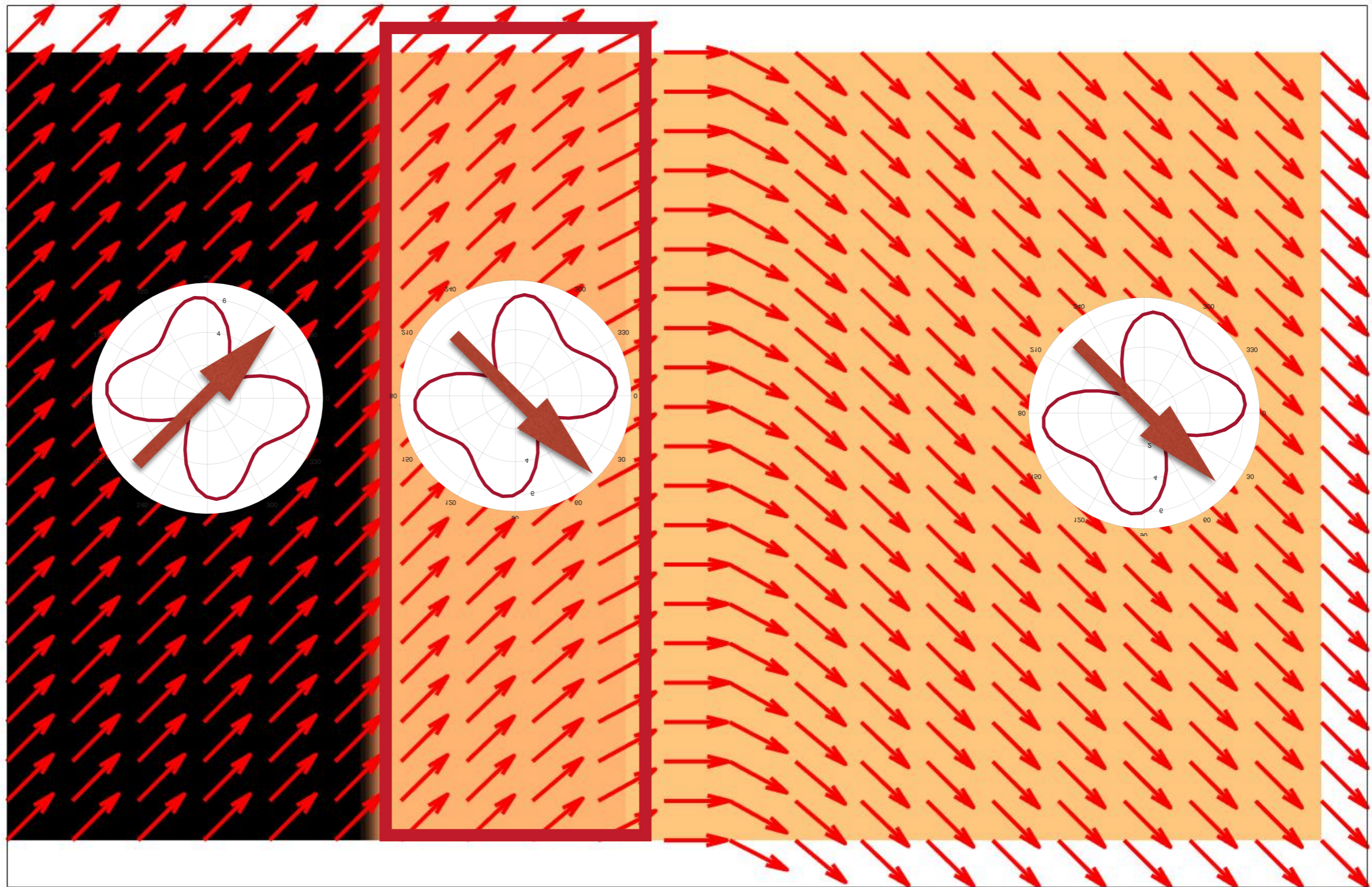
$$\mathbf{n} \times (\ddot{\mathbf{n}} - c^2 \Delta \mathbf{n} + \gamma^2 H_{\text{ex}} \mathbf{H}_{\text{an}}) = 0$$



$$\omega_{\text{DW}} = 0$$

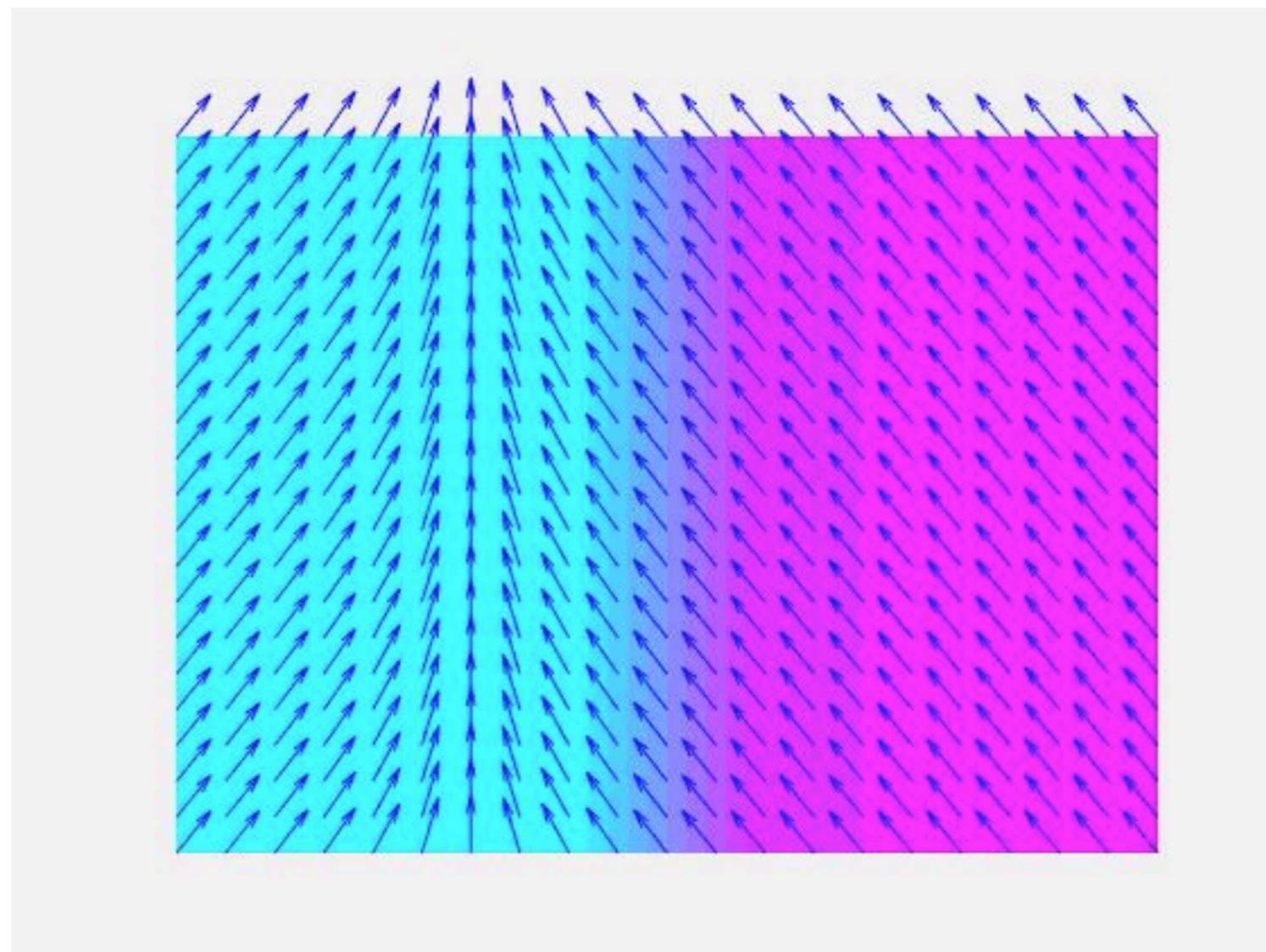
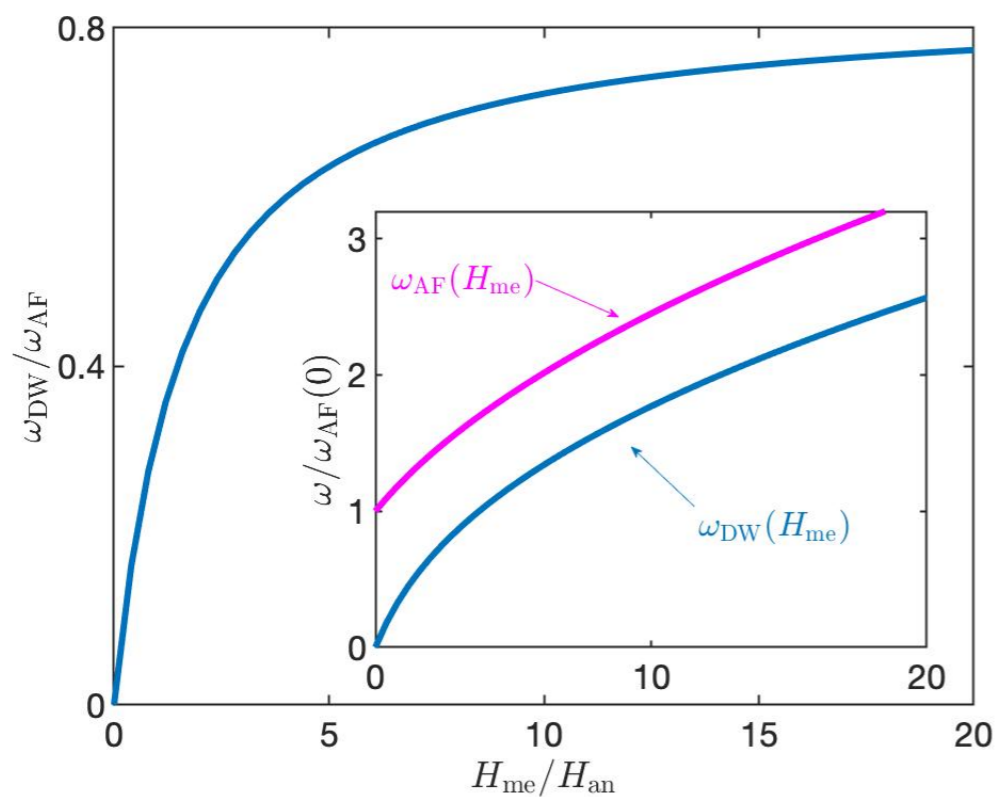


# Magnetoelastic domain wall





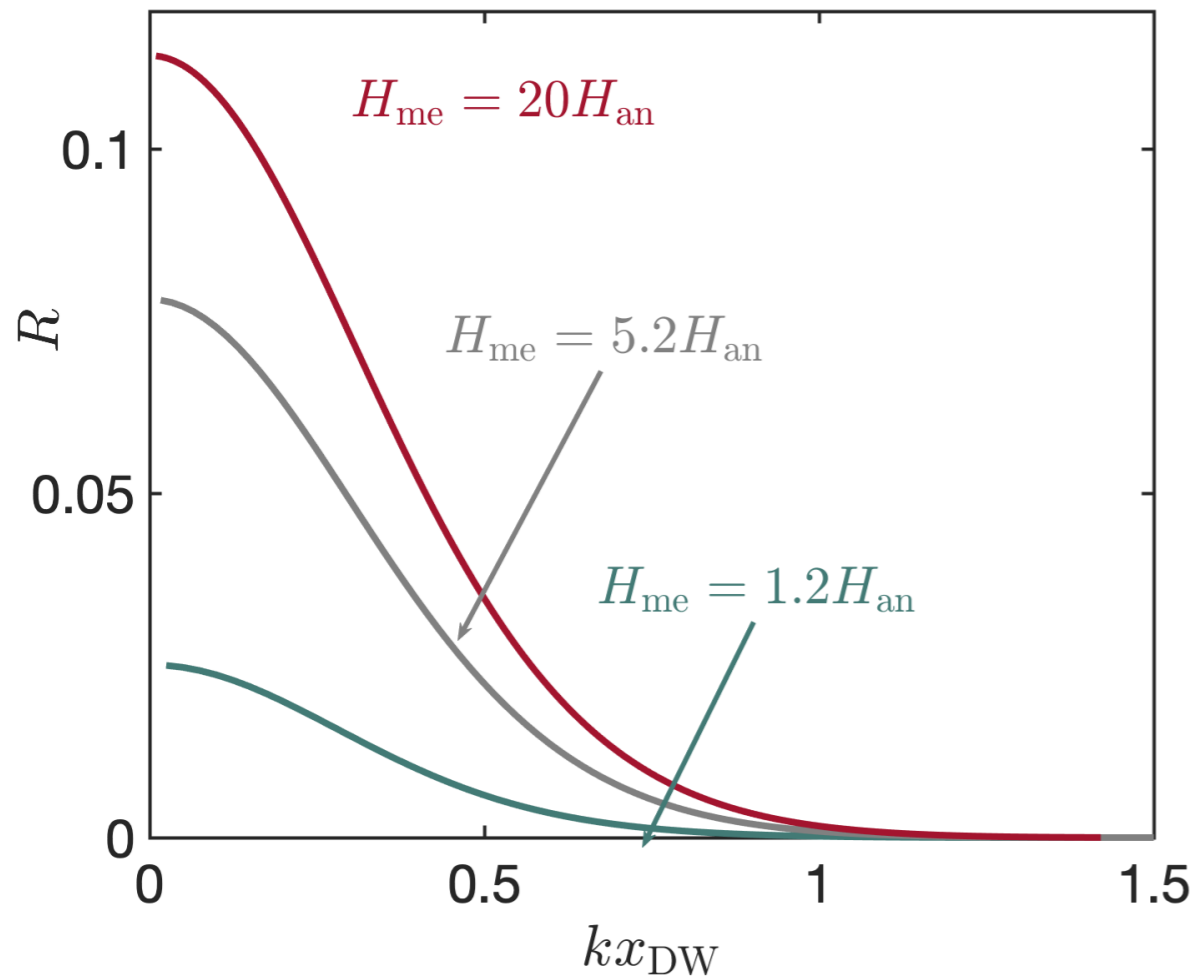
$$\mathbf{n} \times (\ddot{\mathbf{n}} - c^2 \Delta \mathbf{n} + \gamma^2 H_{\text{ex}} \mathbf{H}_{\text{an}}(\xi)) = 0$$



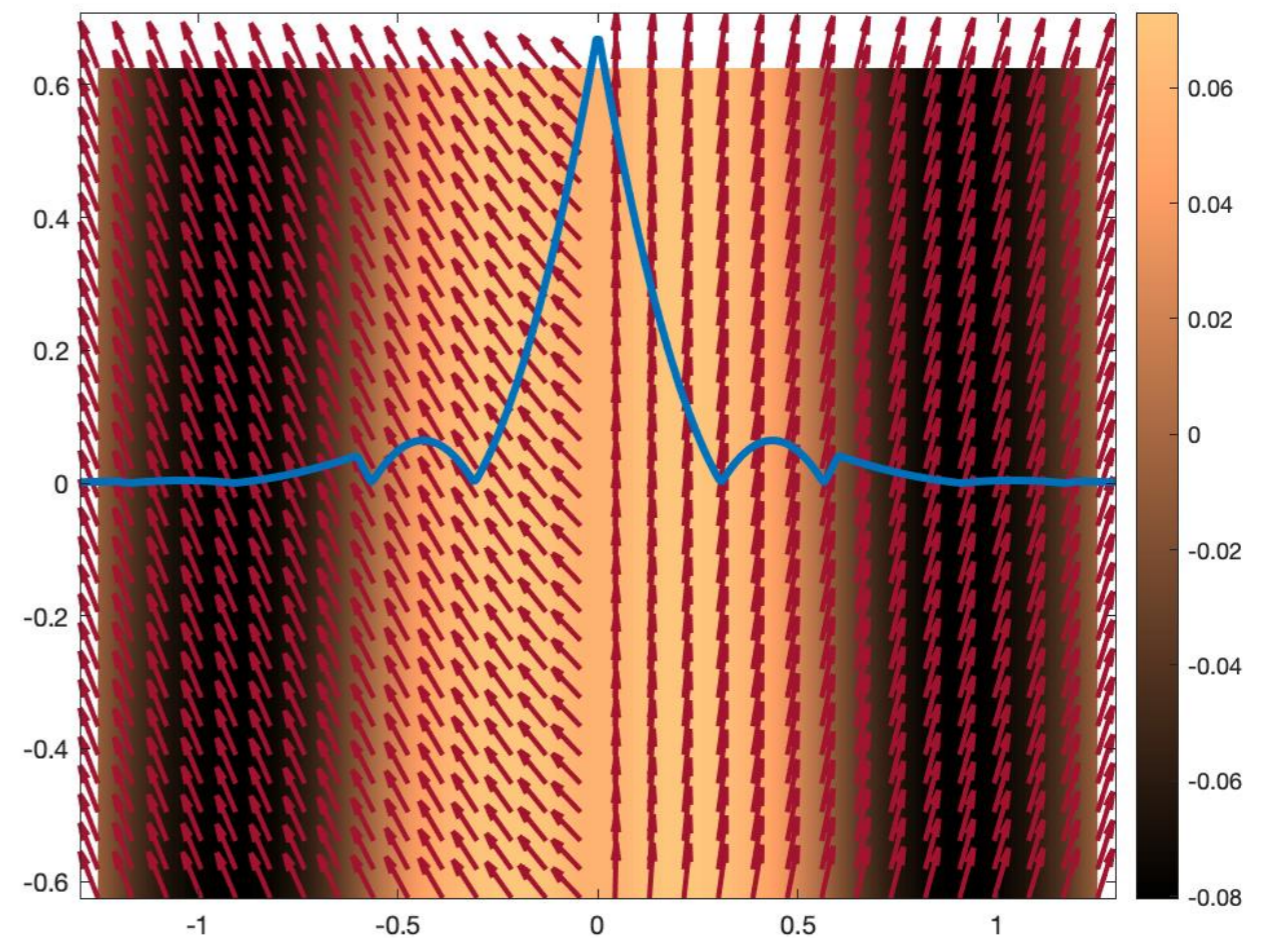
$$\omega_{\text{DW}} \propto \sqrt{H_{\text{me}}}$$



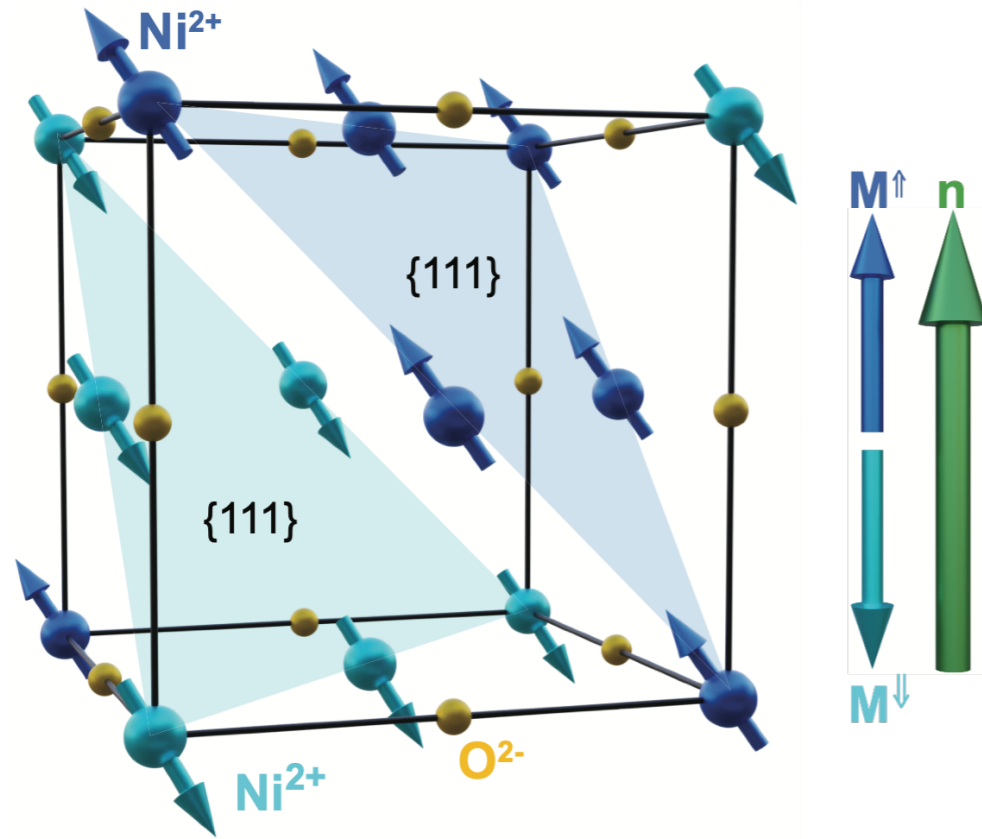
$$\mathbf{n} \times (\ddot{\mathbf{n}} - c^2 \Delta \mathbf{n} + \gamma^2 H_{\text{ex}} \mathbf{H}_{\text{an}}(\xi)) = 0$$



Reflection of magnons

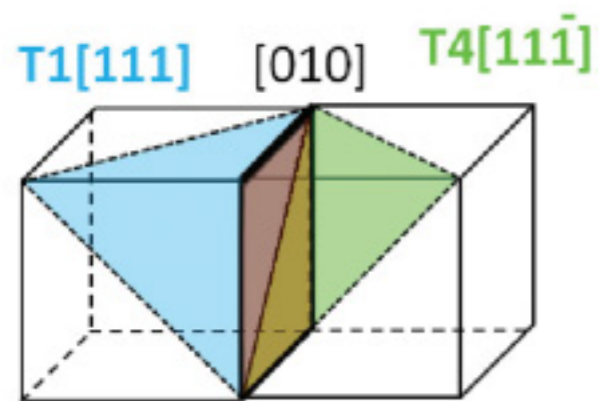


Localised mode

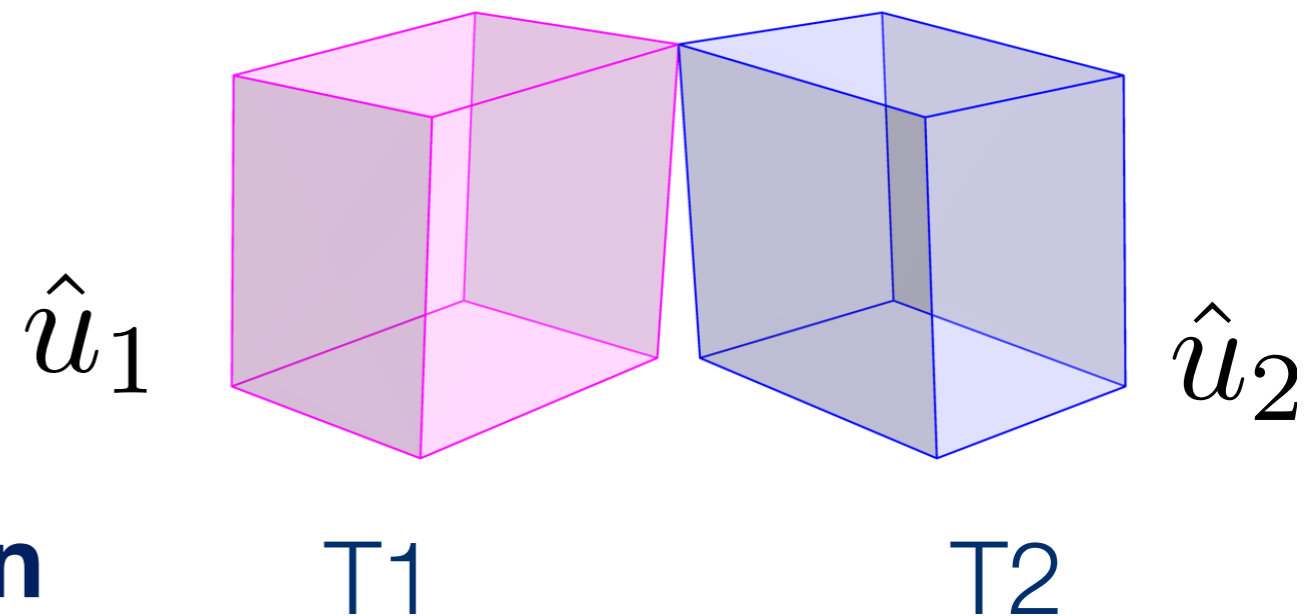


(011)T-wall 5  $\mu\text{m}$

K. Arai et al, PRB, 85, 104418 (2012)

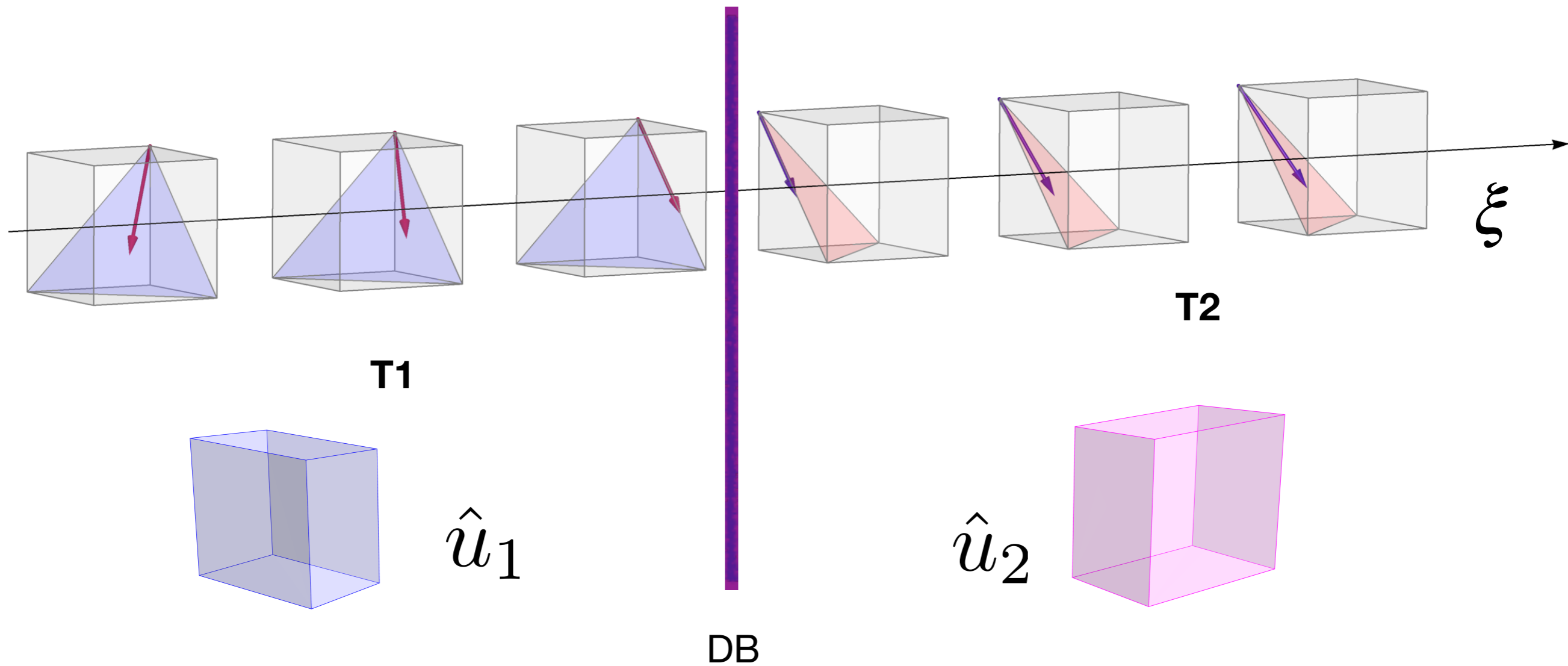


## Distortion



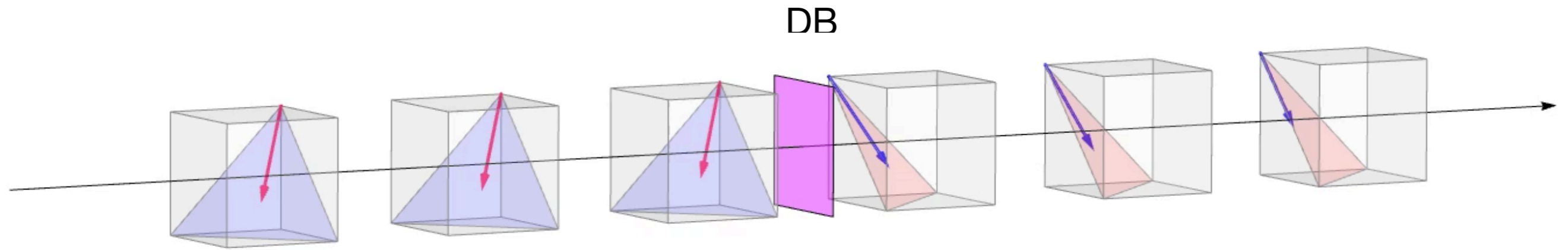
## Exchange striction

# T domain wall



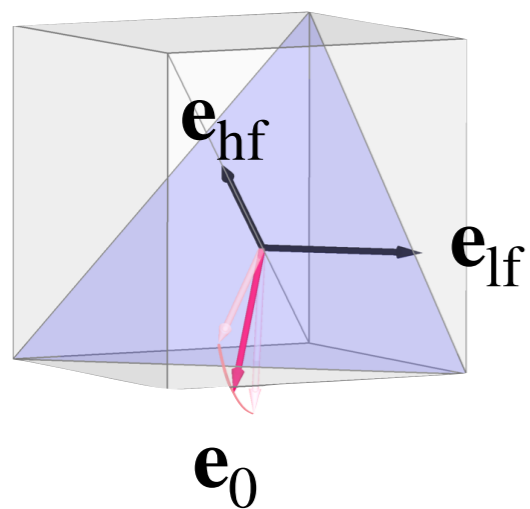
$$w_{\text{an}}(\xi) = w_{\text{mag}} + \begin{cases} M_s H_{\text{me}} \mathbf{n} \cdot \hat{u}_1 \cdot \mathbf{n}, & \xi < 0, \\ M_s H_{\text{me}} \mathbf{n} \cdot \hat{u}_2 \cdot \mathbf{n}, & \xi > 0 \end{cases}$$

# Polarization of eigen modes



**T1**

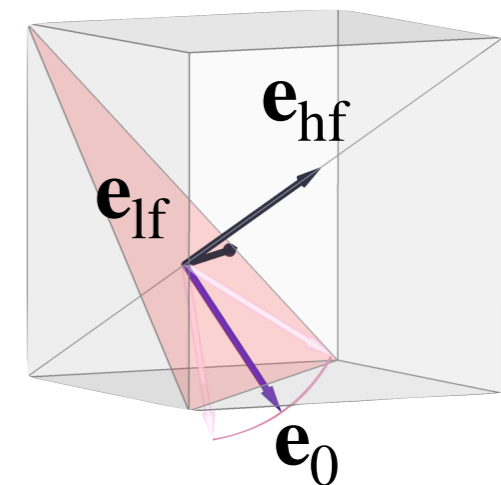
$[111]$



HF

**T2**

$[\bar{1}11]$



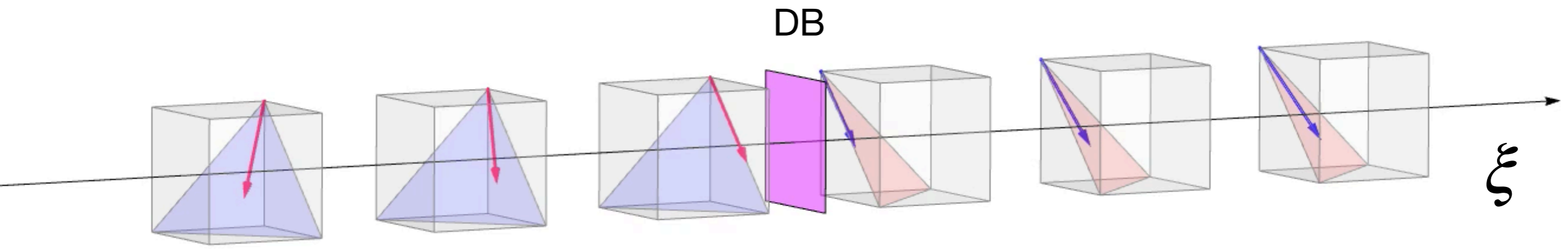
HF

$$\mathbf{e}_{\text{hf}}(\xi) = \frac{1}{\sqrt{3}} \begin{cases} [111] & \xi < 0, \\ [\bar{1}11] & \xi > 0, \end{cases}$$

$$\mathbf{e}_0(\xi) = \frac{1}{\sqrt{6}} \begin{cases} [\bar{1}12] & \xi < 0, \\ [1\bar{1}2] & \xi > 0, \end{cases}$$



# Polarization of eigen modes

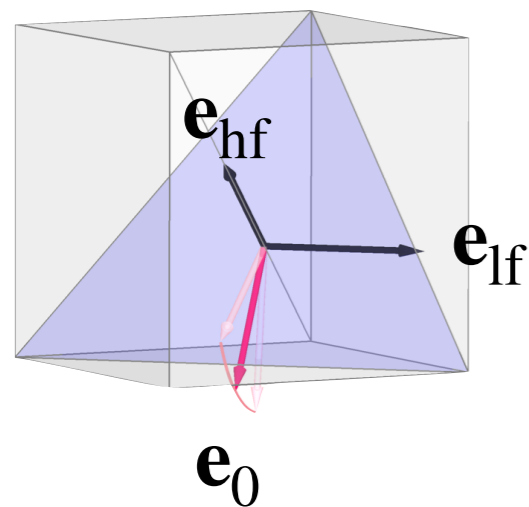


T1

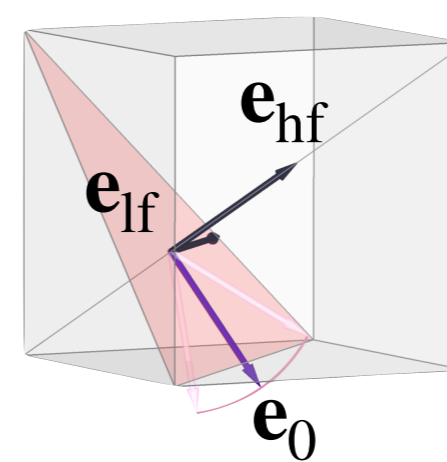
T2

[111]

$[\bar{1}11]$

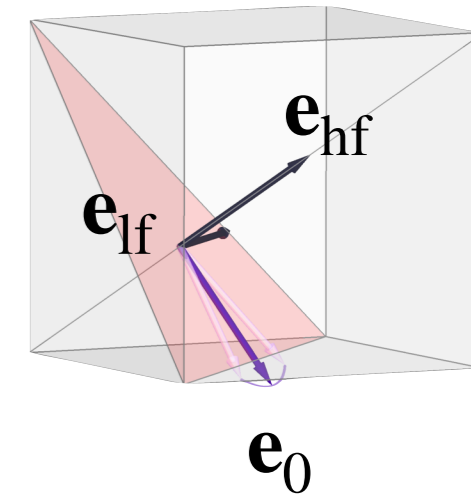


HF



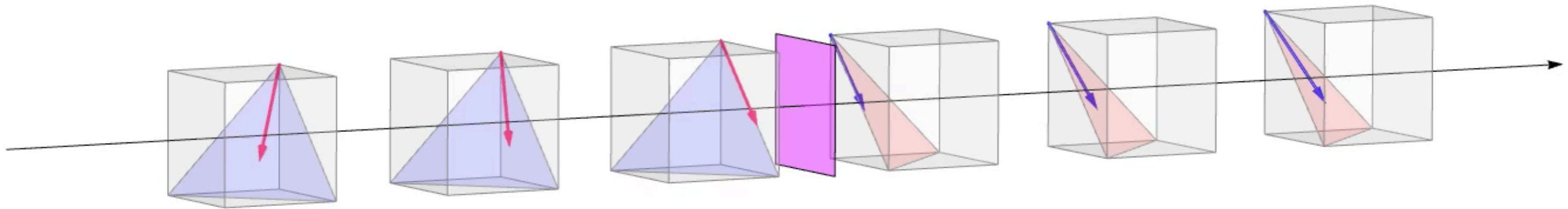
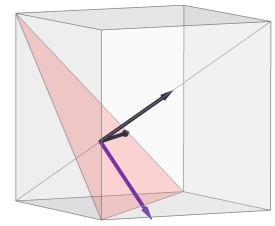
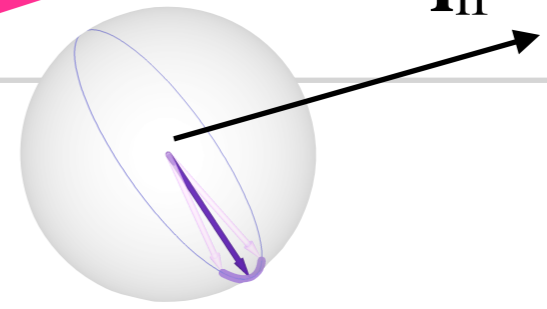
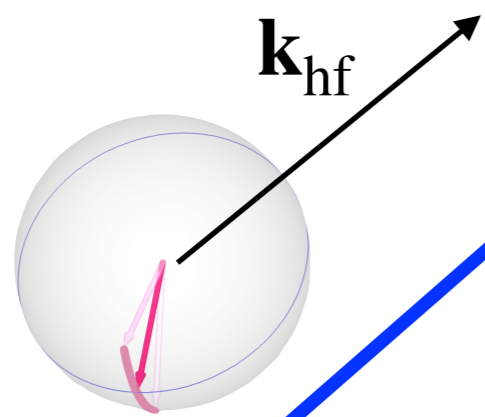
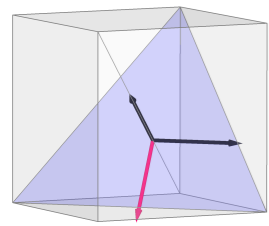
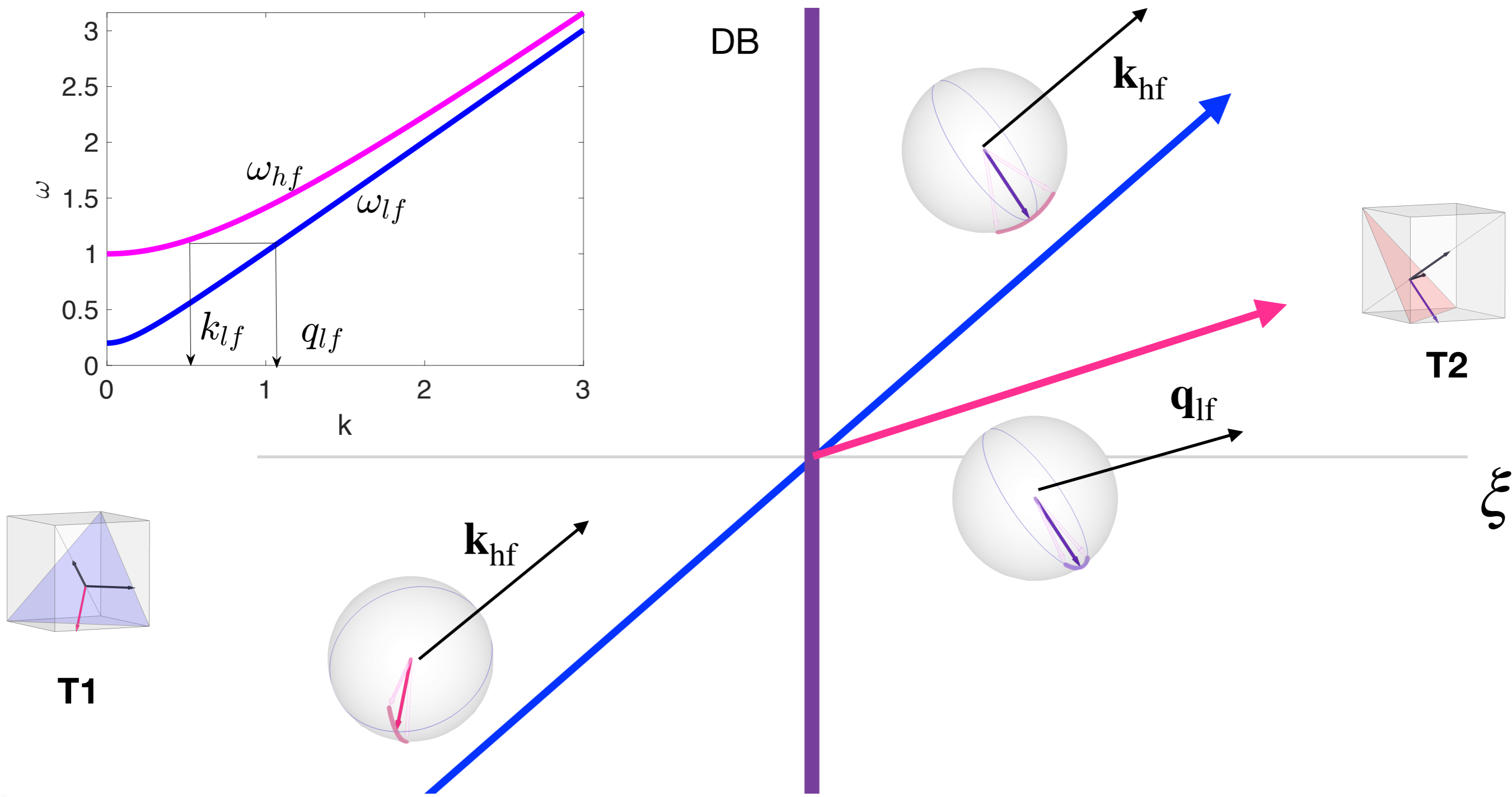
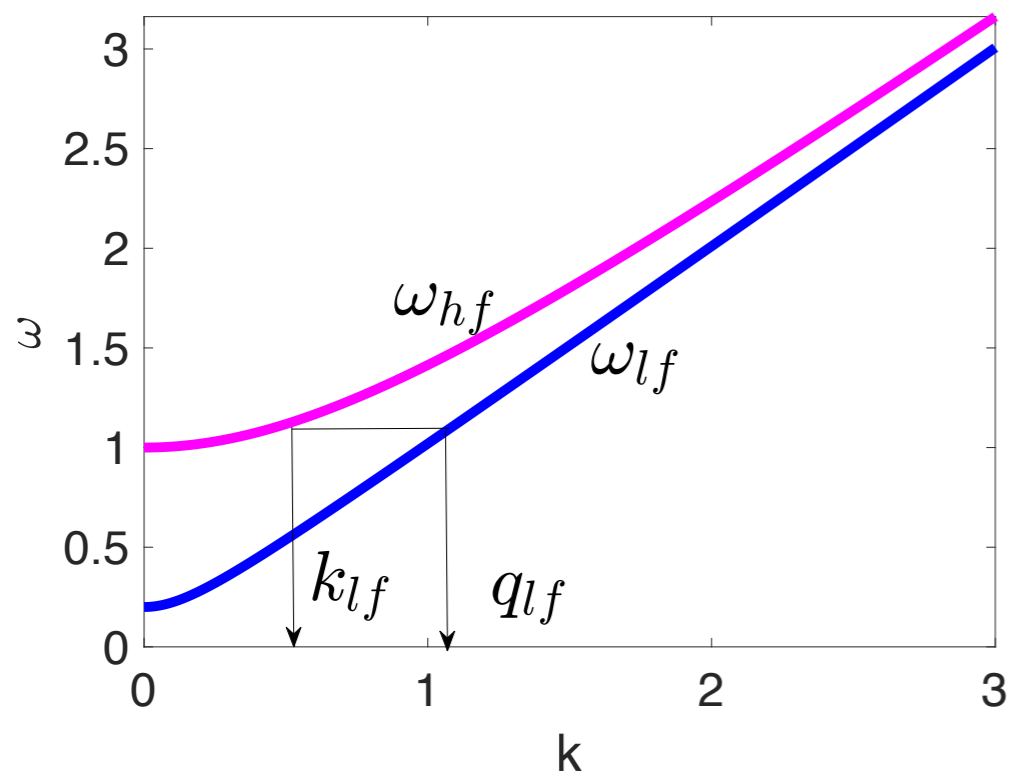
HF

+

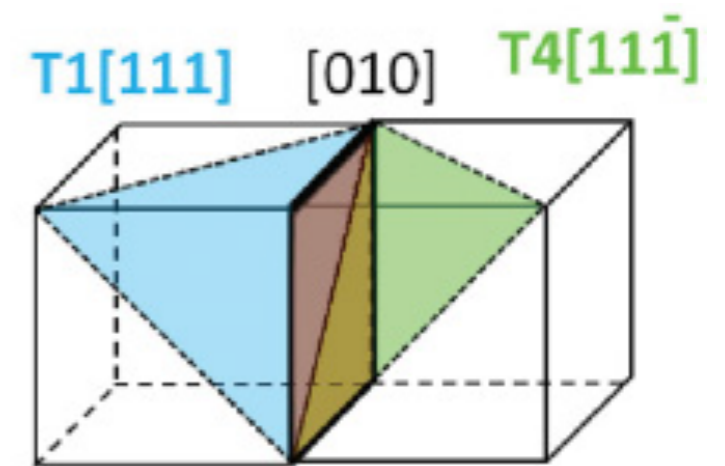
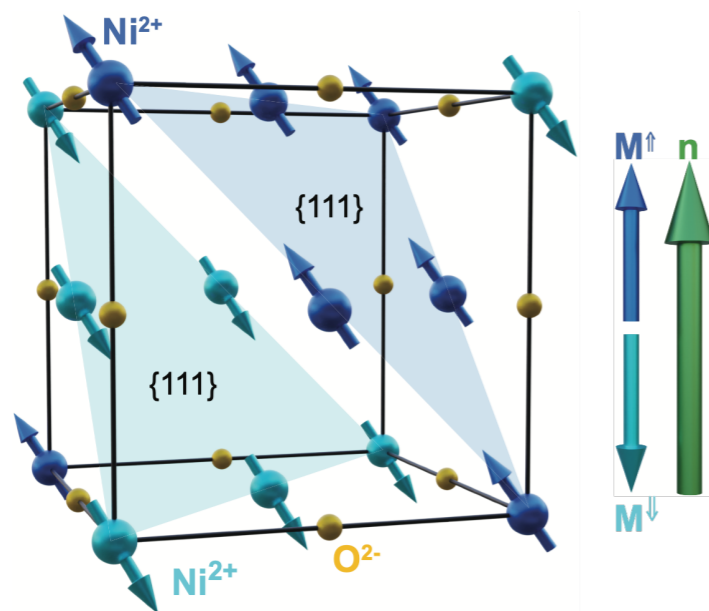
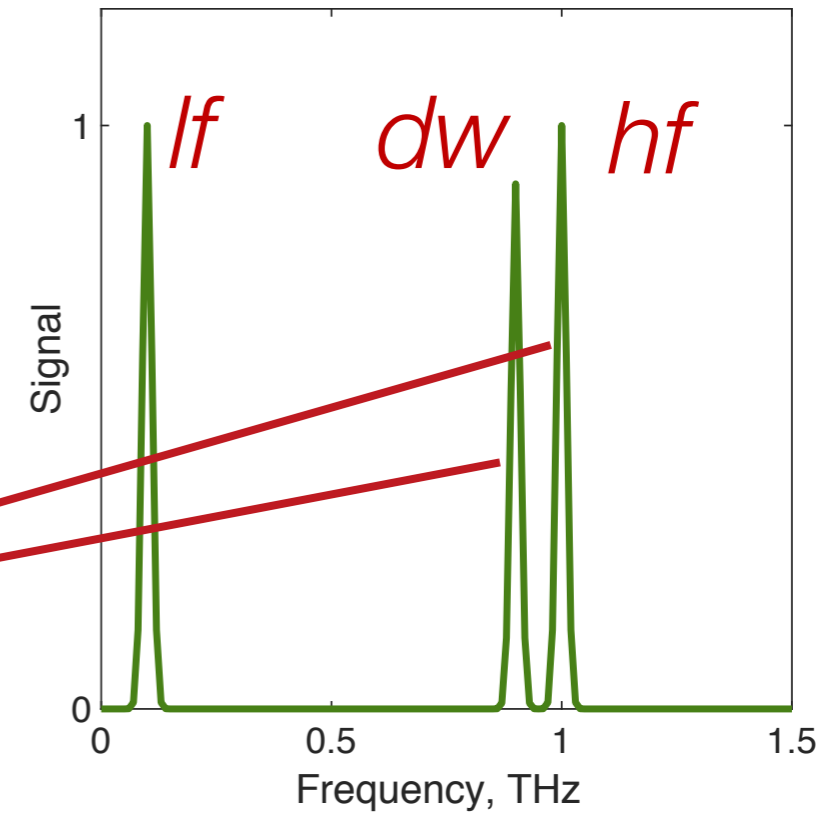
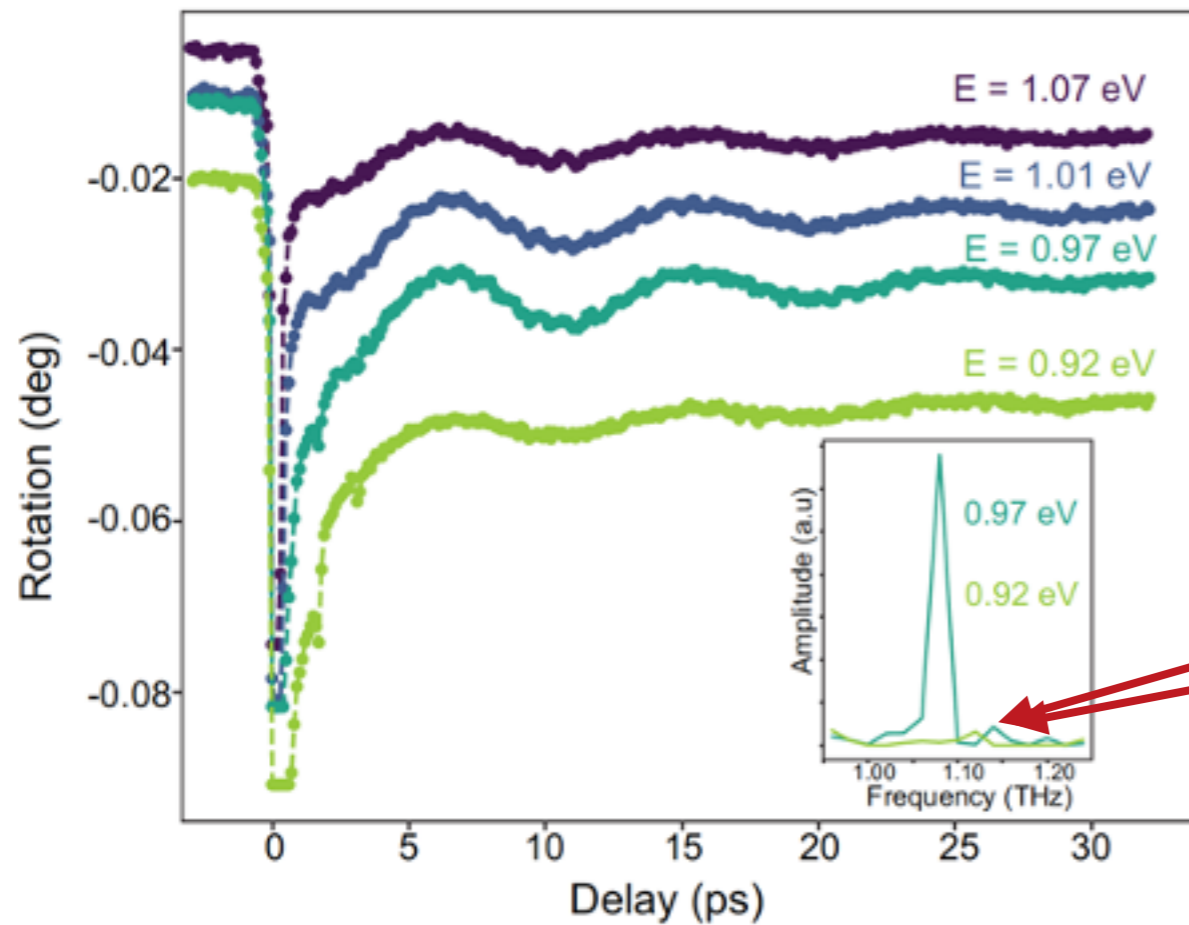


LF

# Magnon birefringence



# Parametric downconversion, NiO



Bossini, O.G, et, submitted to PRL

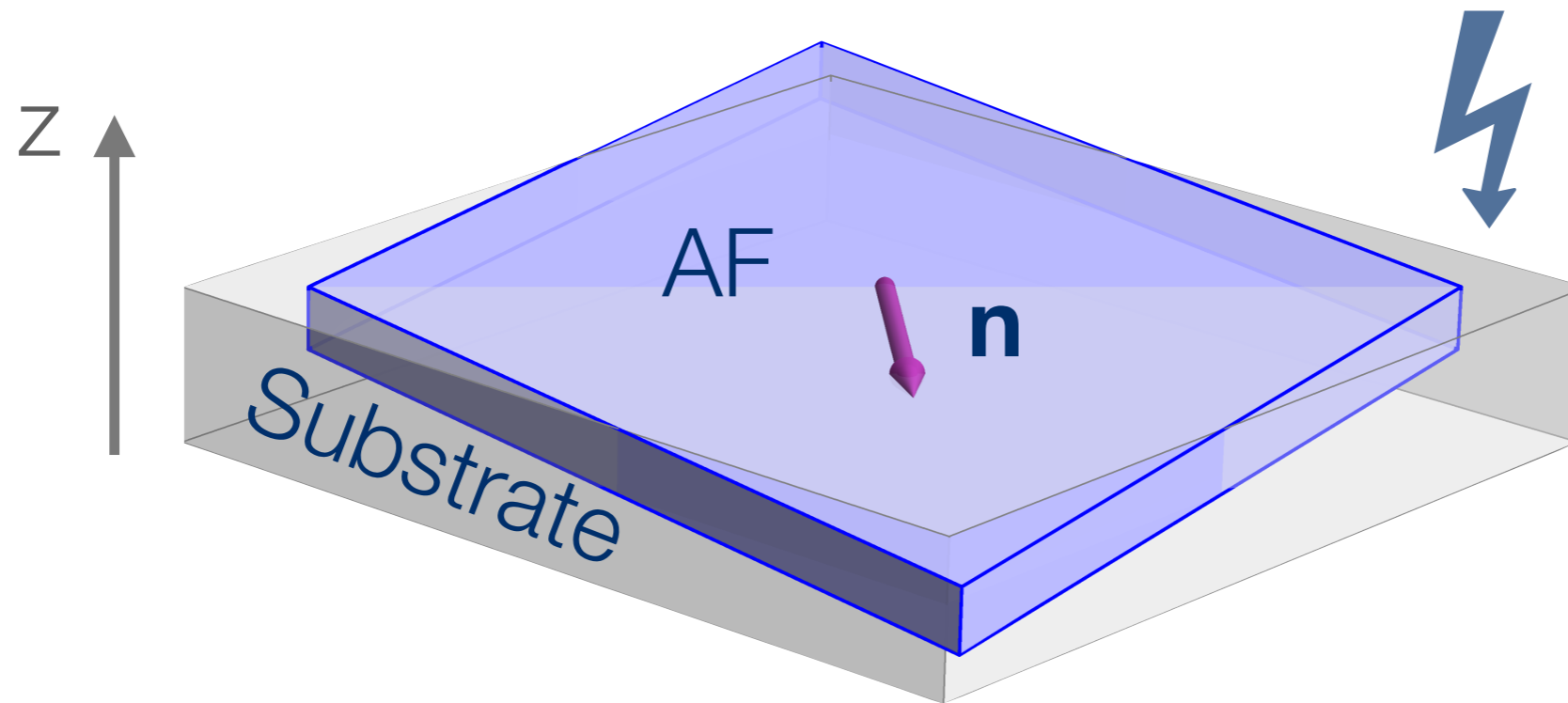
- Magnetoelasticity: frozen spontaneous strains and domain walls
- **Incompatibility of strains**
- Nonequilibrium strains: switching
- Equilibrium domain structure: *Micr-a*-magnetics
- Conclusions





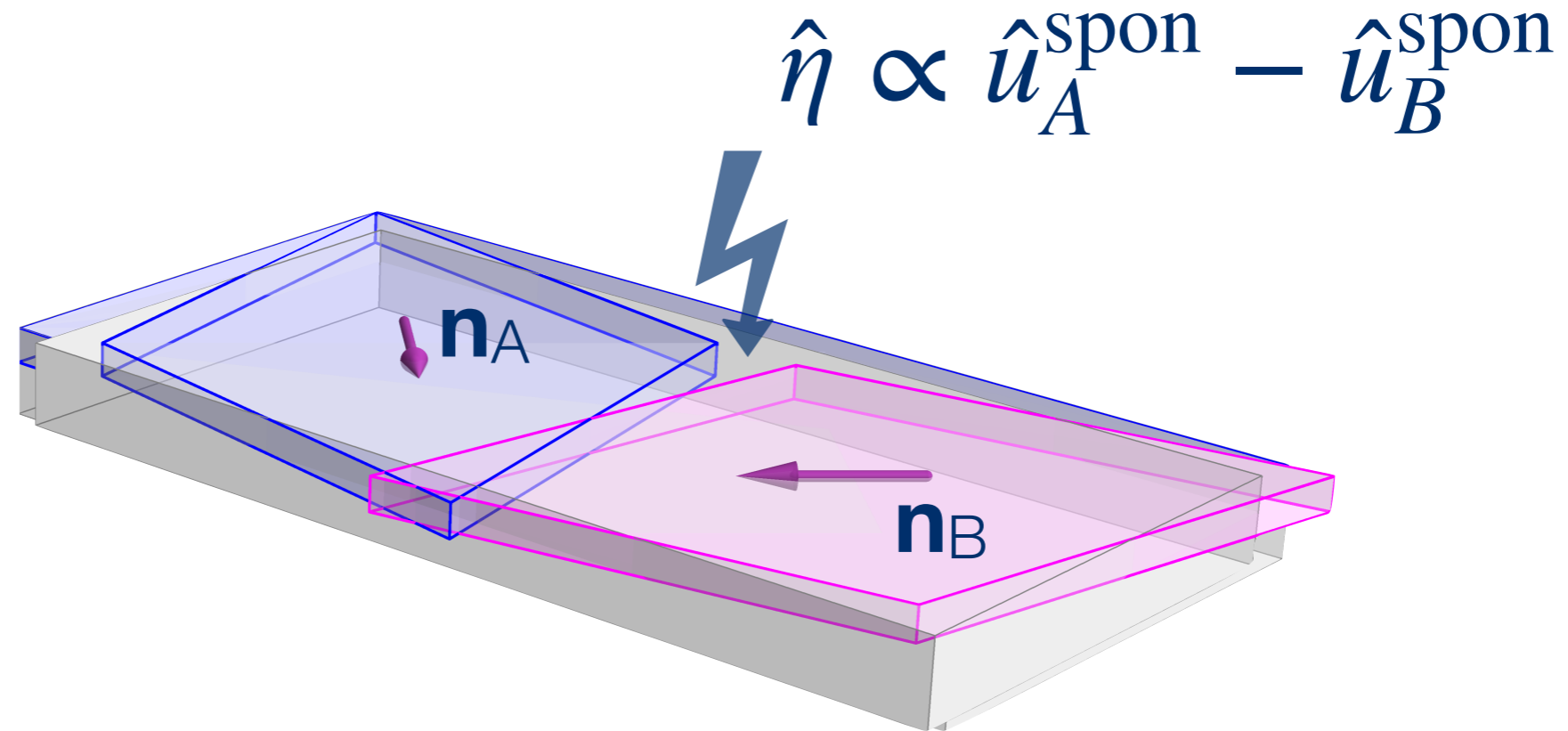
# Example I: AF + substrate

$$\hat{\eta} \propto \hat{u}^{\text{spon}} - \hat{u}^{\text{sub}}$$



Incompatibility at the interface

# Example II: domain A + domain B



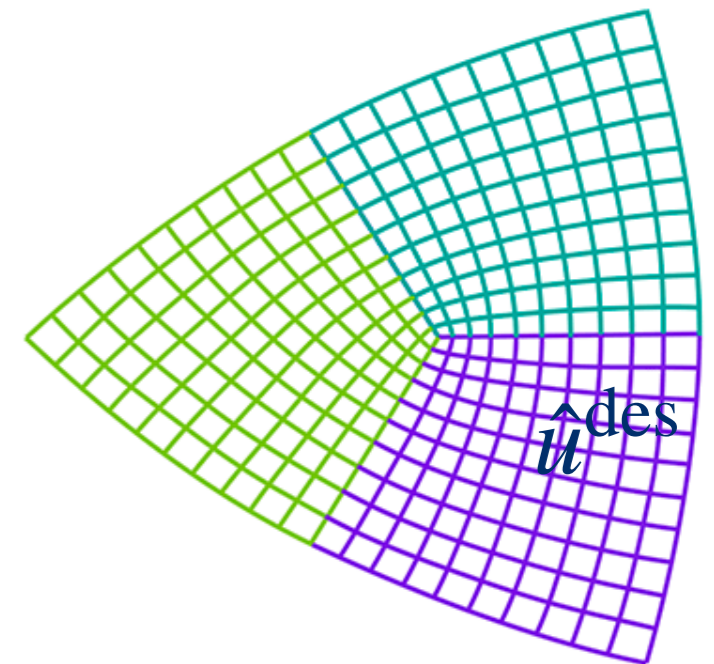
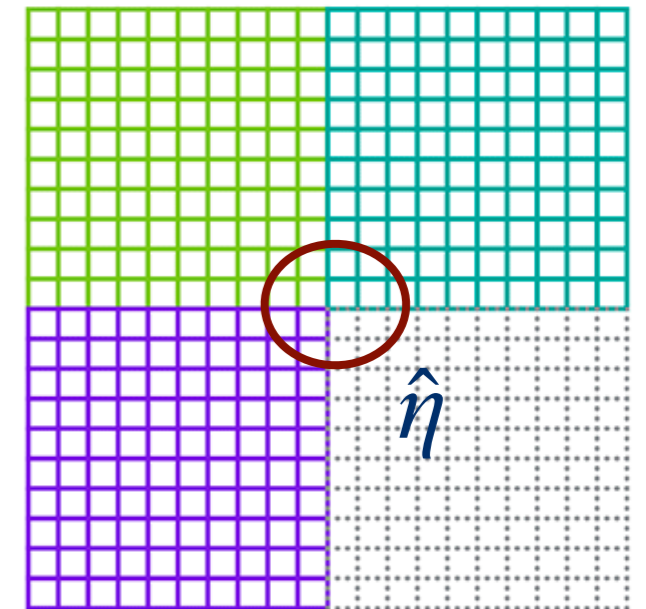
Incompatibility between deformed regions

$$\nabla \times \hat{u}^{\text{des}}(\mathbf{r}) \times \nabla = -\hat{\eta}$$

$$\hat{\eta} = \nabla \times \hat{u}^{\text{spon}}(\mathbf{r}) \times \nabla$$

$$\hat{u} = \underbrace{\hat{u}^{\text{spon}}(\mathbf{r})}_{\text{dislocations}} + \underbrace{\hat{u}^{\text{des}}(\mathbf{r})}_{\text{plastic}}$$

$$\hat{u}^{\text{spon}} \propto \mathbf{n} \otimes \mathbf{n}$$

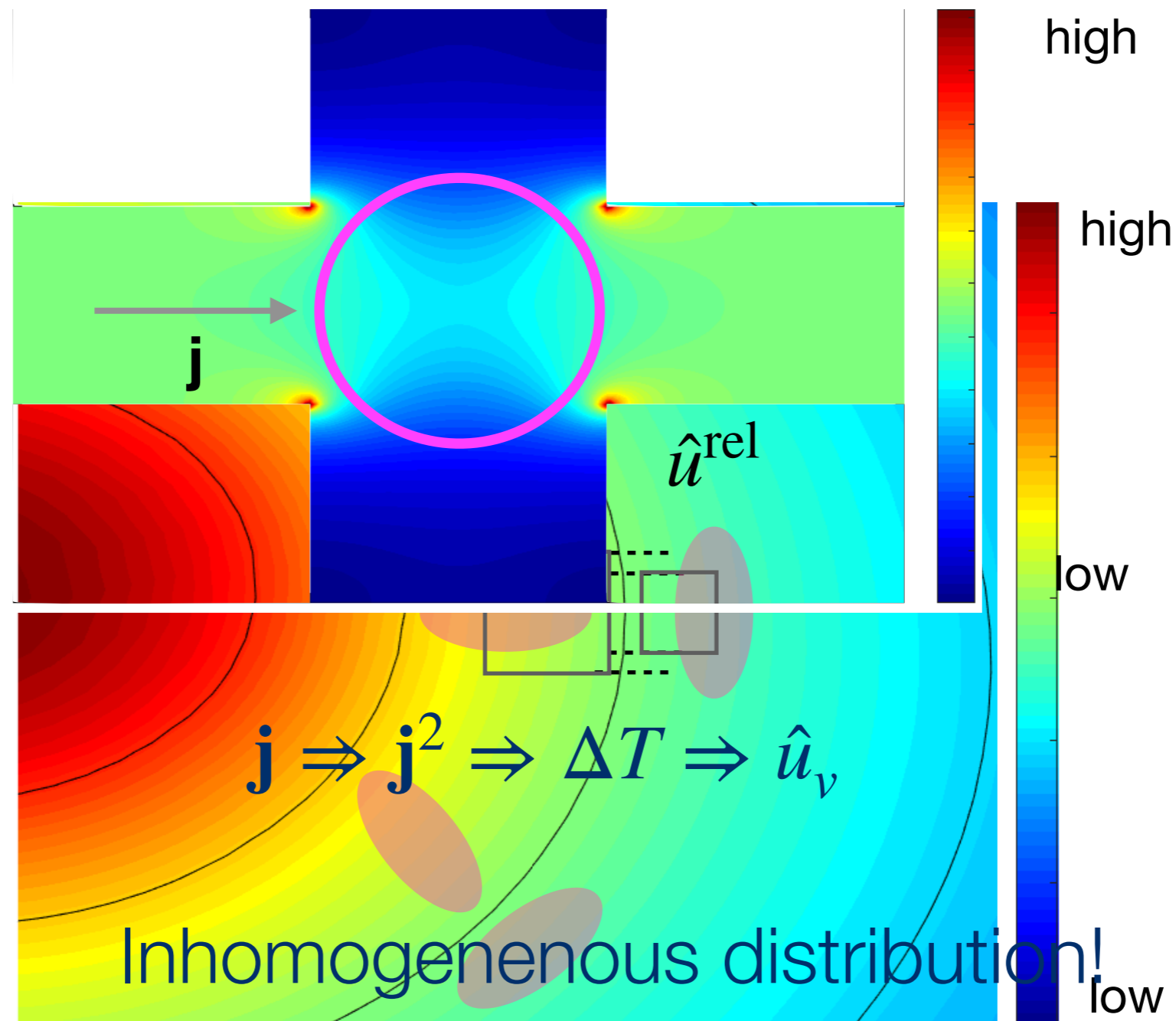


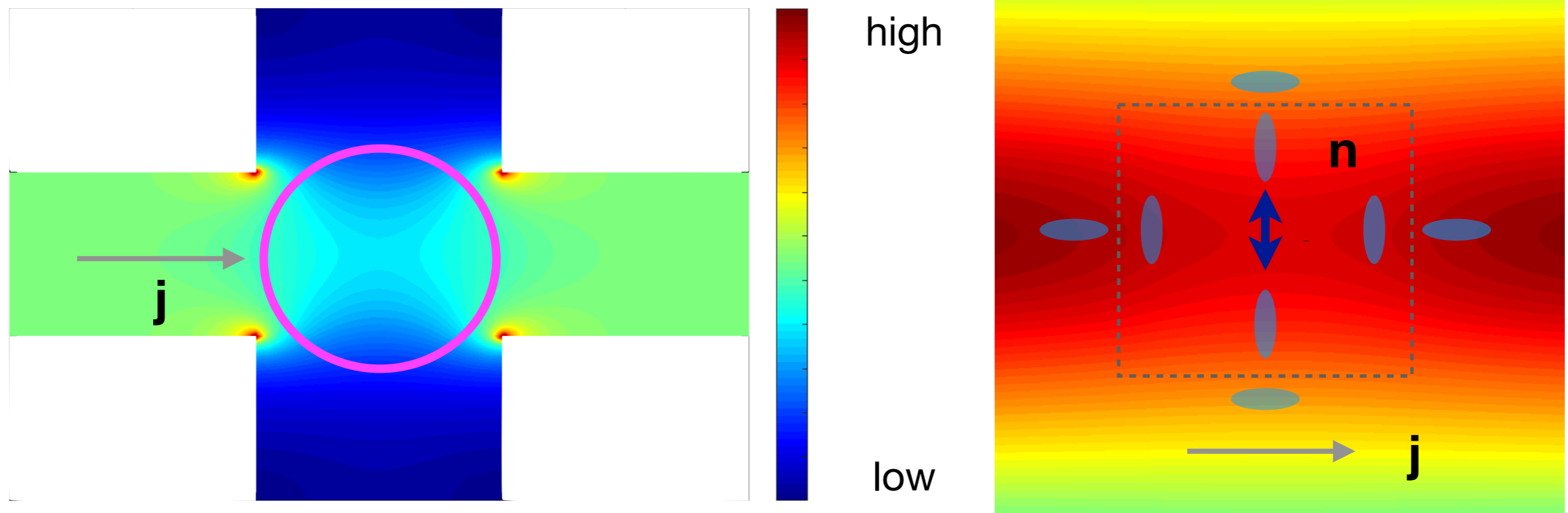


- Magnetoelasticity: frozen spontaneous strains and domain walls
- Incompatibility of strains
- **Nonequilibrium strains: switching**
- Equilibrium domain structure: *Micr-a*-magnetics
- Conclusions



# Thermo-magneto-elastic effect





$$U_{\text{ext}} \propto \lambda(\mathbf{j} \cdot \mathbf{n})^2$$

Zhang et al Phys. Rev. Lett. 123, 247206 (2019)

Baldrati, OG, et al. Phys. Rev. Lett. 125, 77201 (2020)

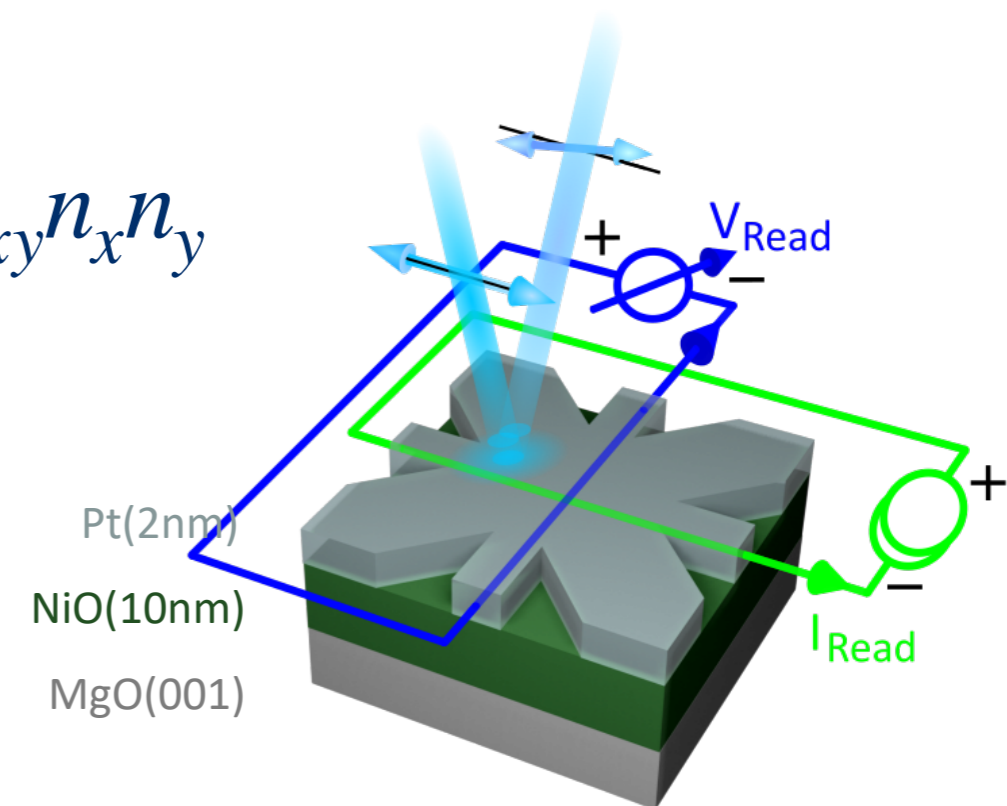


## Zeeman energy

$$U_{\text{ext}} \propto \left( H_x^2 - H_y^2 \right) \left( n_x^2 - n_y^2 \right) + 4H_x H_y n_x n_y = (\mathbf{H} \cdot \mathbf{n})^2$$

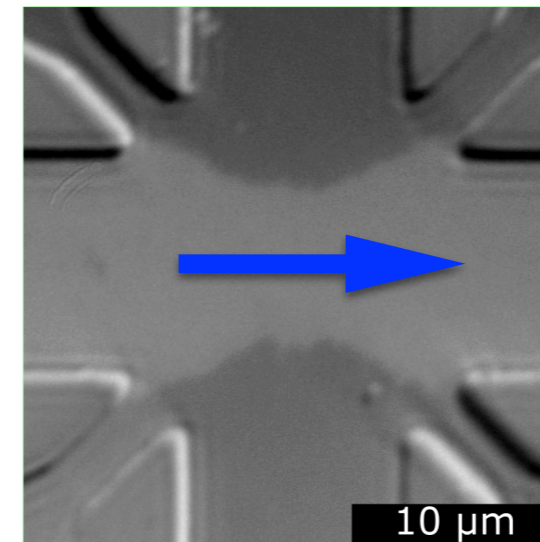
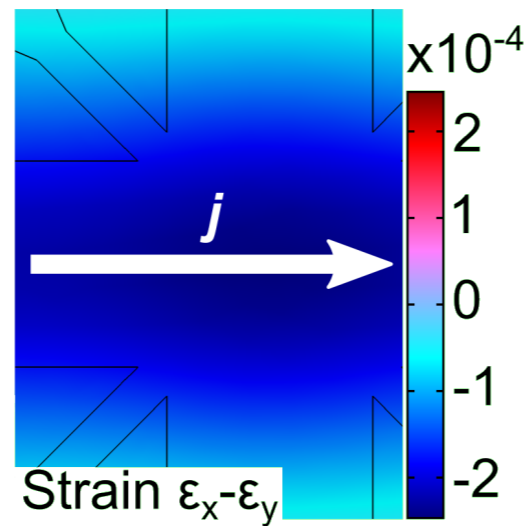
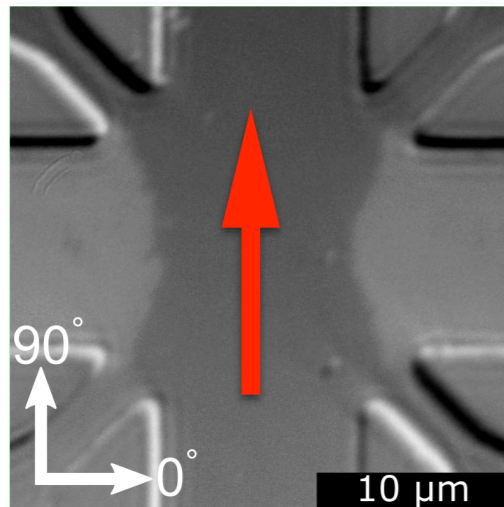
## Magnetoelastic energy, external strains

$$U_{\text{ext}} \propto \left( u_{xx} - u_{yy} \right) \left( n_x^2 - n_y^2 \right) + 2u_{xy} n_x n_y$$



Courtesy F. Schreiber

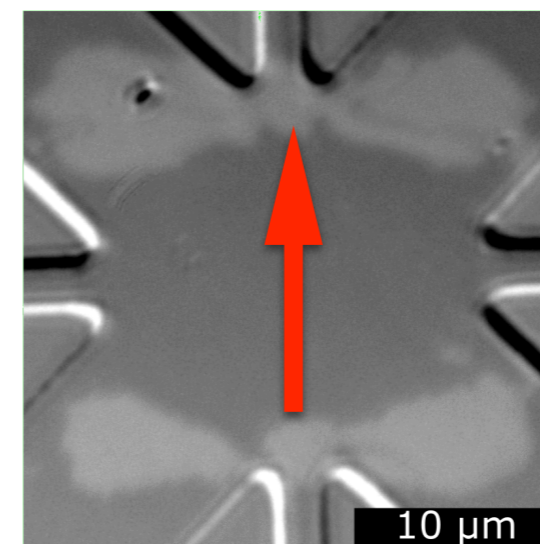
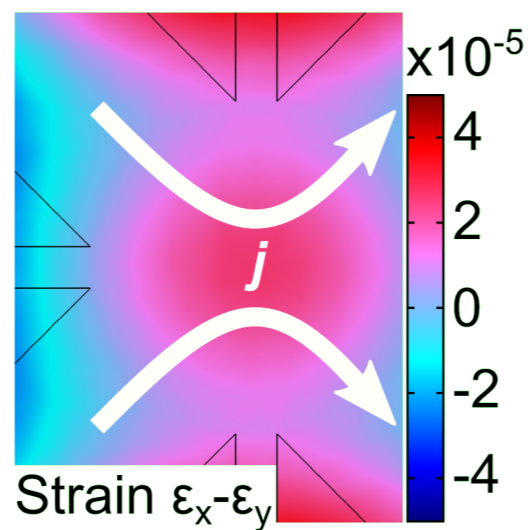
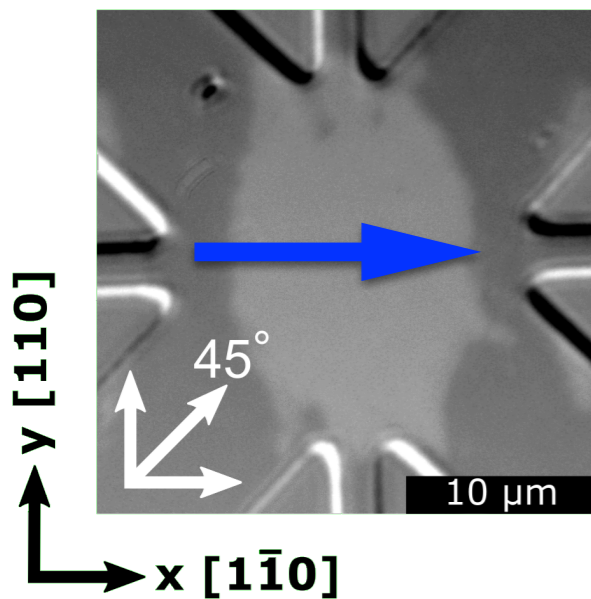
## From SMR



$j \parallel n$

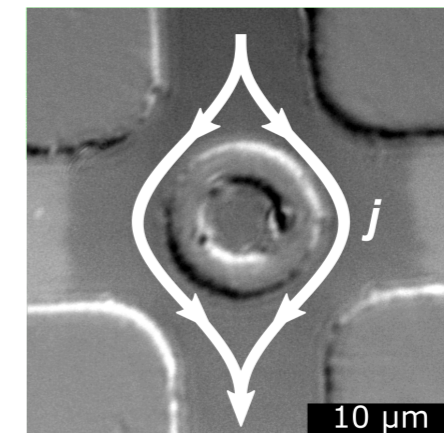
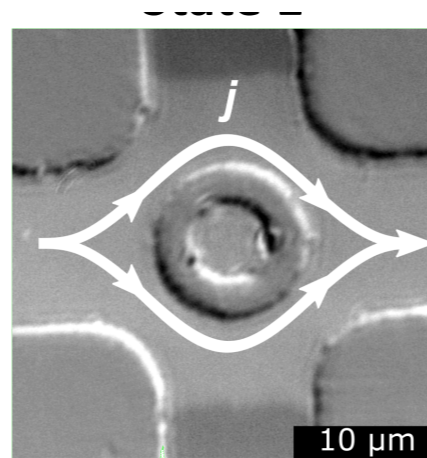
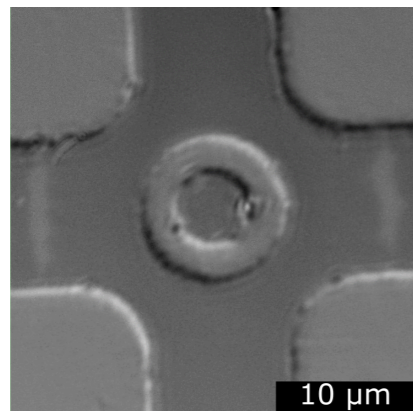
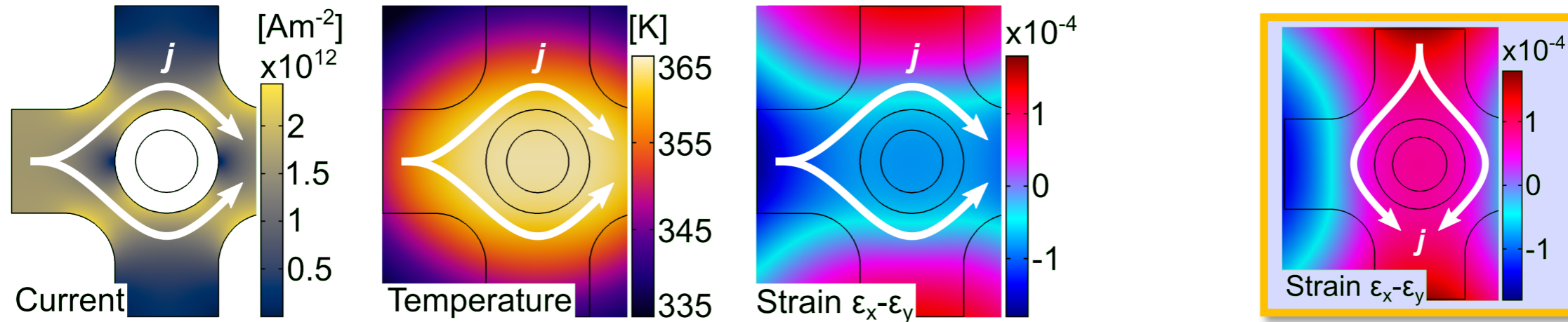
## Initial state

## Final state



$j \perp n$

# Thermo-magneto-elastic effect, NiO

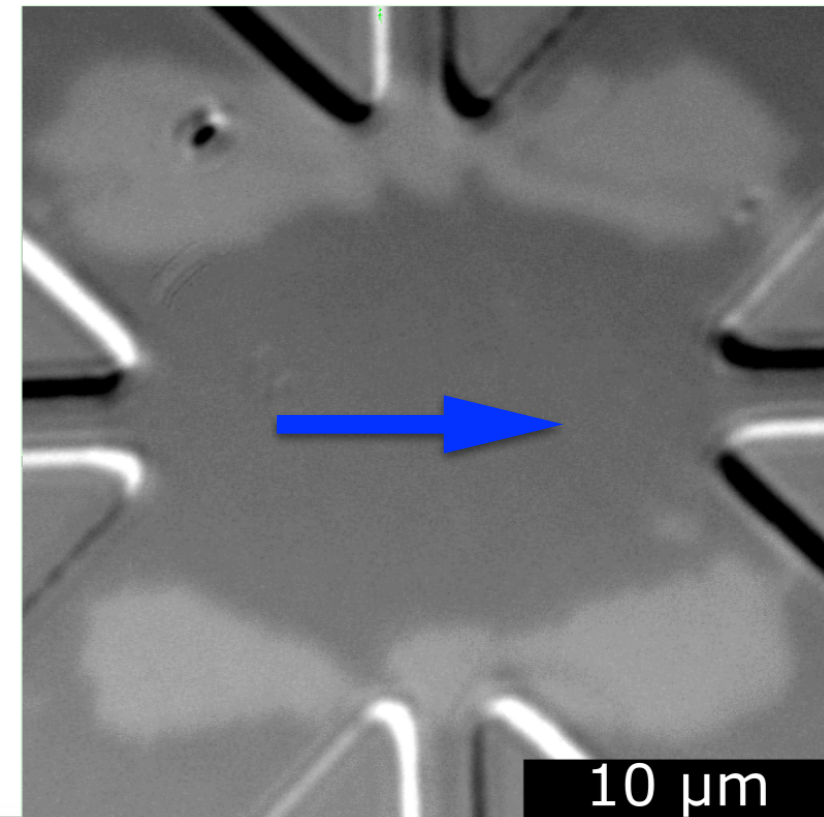
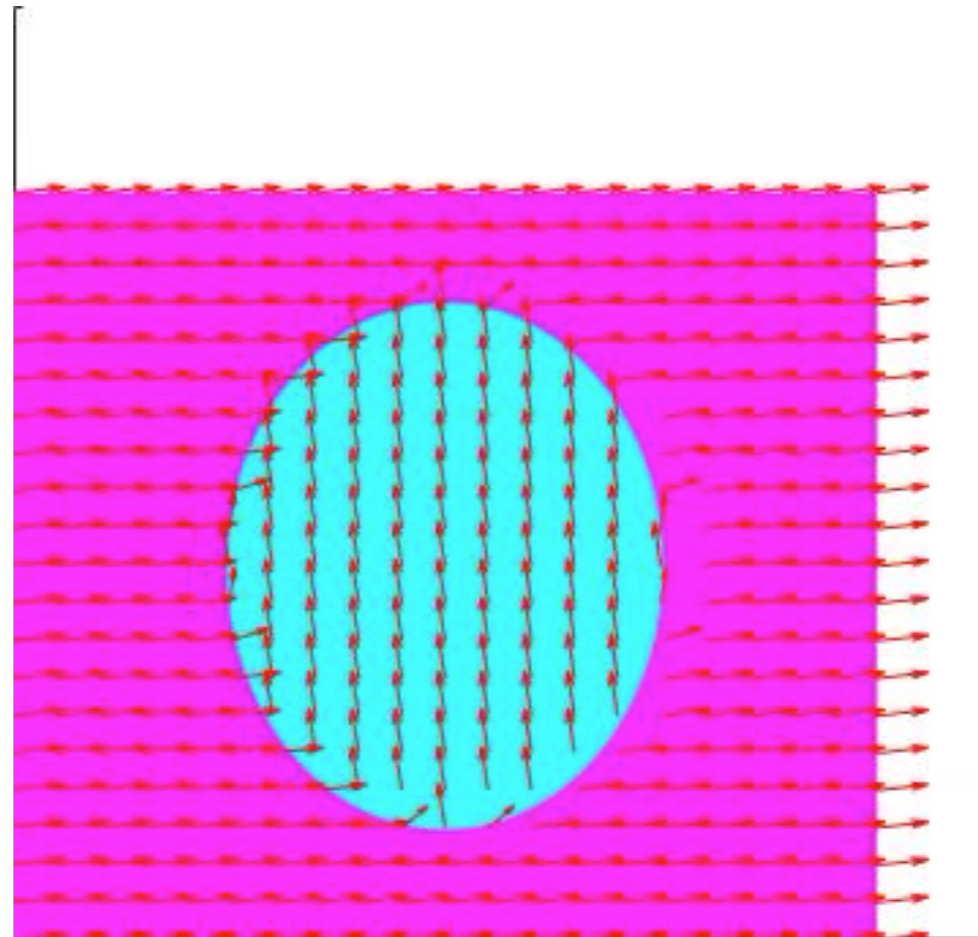
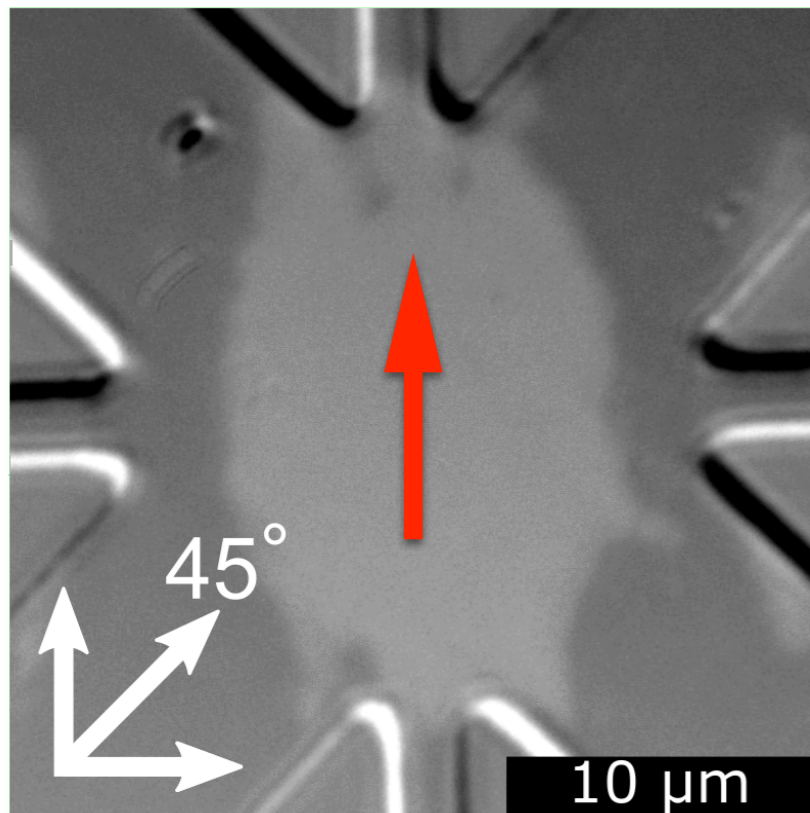


Meer, OG, et al. , Nano Lett. 21, 114 (2020)

Courtesy H. Meer



# Dynamics of the domain switching



- Magnetoelasticity: frozen spontaneous strains and domain walls
- Incompatibility of strains
- Nonequilibrium strains: switching
- **Equilibrium domain structure: Micr-*a*-magnetics**
- Conclusions



## Elastostatic

$$\hat{\eta} = \nabla \times \hat{u}^{\text{spon}}(\mathbf{r}) \times \nabla$$

$$\nabla \times \hat{u}^{\text{des}}(\mathbf{r}) \times \nabla = -\hat{\eta}$$

$$\hat{I}(\mathbf{r}) \propto \nabla_j \nabla_k |\mathbf{r}| \propto \frac{1}{|\mathbf{r}|}$$

$$\hat{u}^{\text{spon}} \propto \mathbf{n} \otimes \mathbf{n}$$

## Magnetostatics

$$\rho_m = \text{div} \mathbf{M}$$

$$\nabla \cdot \nabla \psi = -\rho_m$$

$$G(\mathbf{r}) \propto \nabla \cdot \nabla |\mathbf{r}| = \frac{1}{|\mathbf{r}|}$$

$$\mathbf{M} \propto \mathbf{M}$$

$$\hat{u} = \underbrace{\hat{u}^{\text{spon}}(\mathbf{r})}_{\text{dislocations}} + \underbrace{\hat{u}^{\text{des}}(\mathbf{r})}_{\text{plastic}}$$

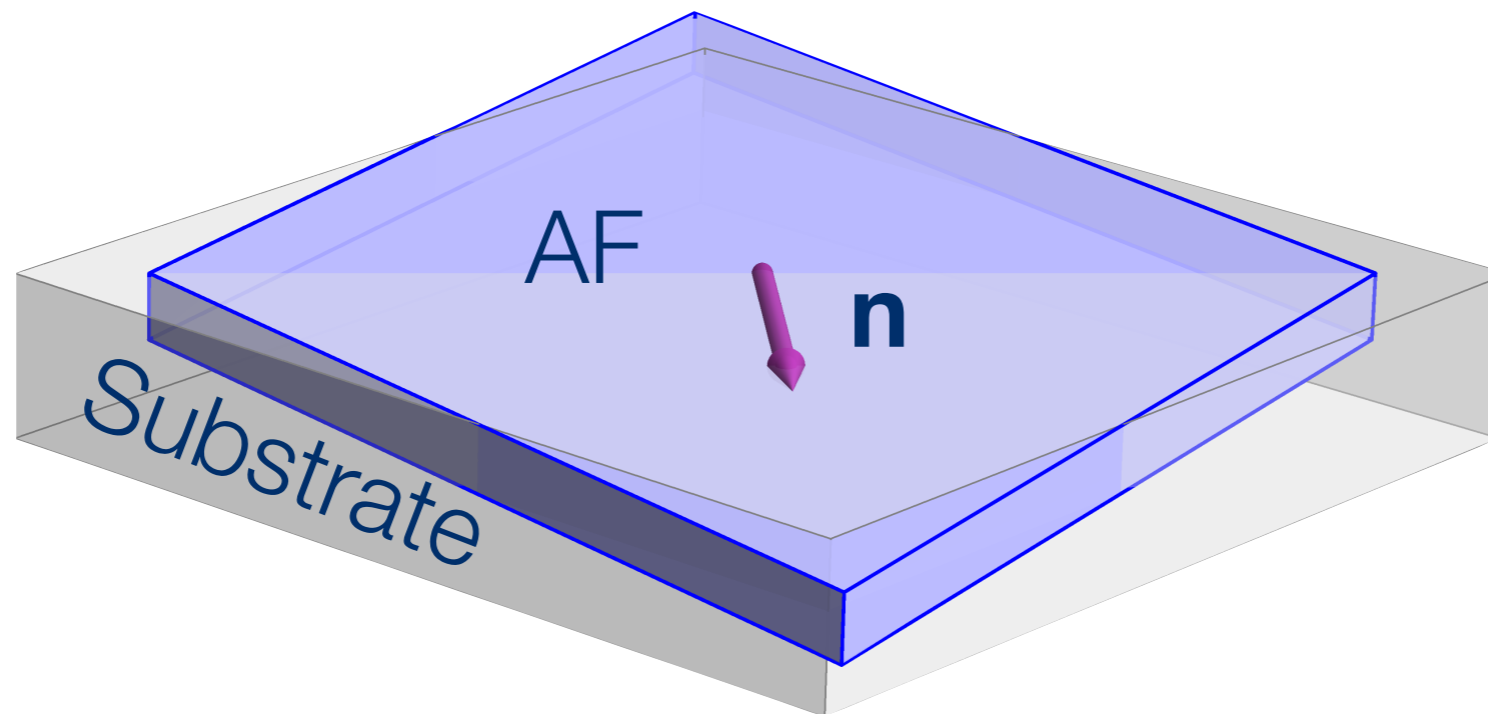
dislocations

plastic





Demagnetising  $\rightarrow$  de**stressing**

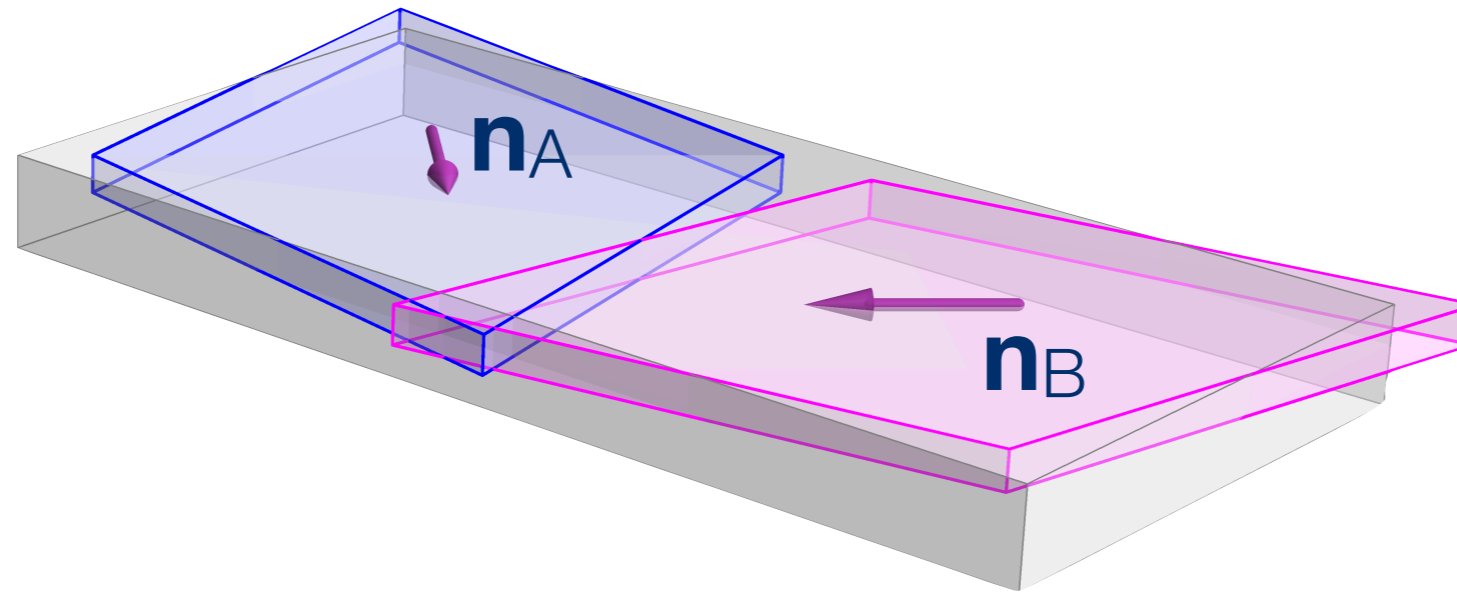


$$E_{\text{des}} = \langle \mathbf{n} \otimes \mathbf{n} \rangle \hat{N} \langle \mathbf{n} \otimes \mathbf{n} \rangle V$$

$$\hat{N} = \int_V \hat{K}(\mathbf{r}_1 - \mathbf{r}_2) d\mathbf{r}_2$$

Gomonay, Loktev, 2002

Bloch type DW  $\rightarrow$  Nye type DW



Flat domain wall

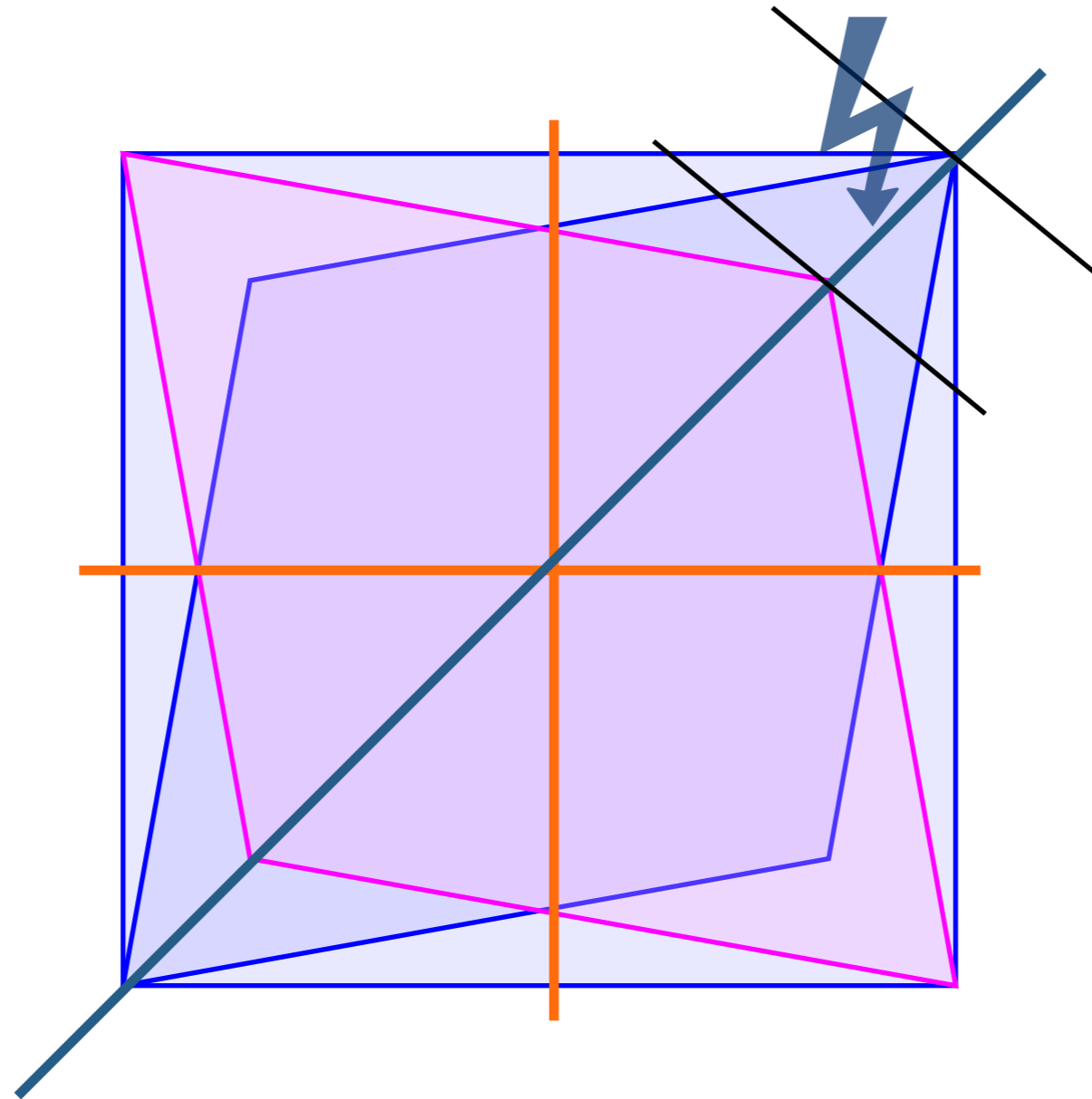
$$\hat{u}^{\text{des}}(\mathbf{r}) = \mathbf{N} \times \left[ \hat{u}_A^{\text{spon}} - \hat{u}_B^{\text{spon}} \right] \times \mathbf{N} \int \nabla \hat{I}(\mathbf{r} - \mathbf{r}_{\text{DW}}) d\mathbf{r}_{\text{DW}}$$

$$E_{\text{des}} \propto V$$

Kleman, Miltat, 1972, for FM

# Nye domain wall

$$\hat{\eta} \propto \mathbf{N} \times (\hat{u}_A^{\text{spon}} - \hat{u}_B^{\text{spon}}) \times \mathbf{N}$$



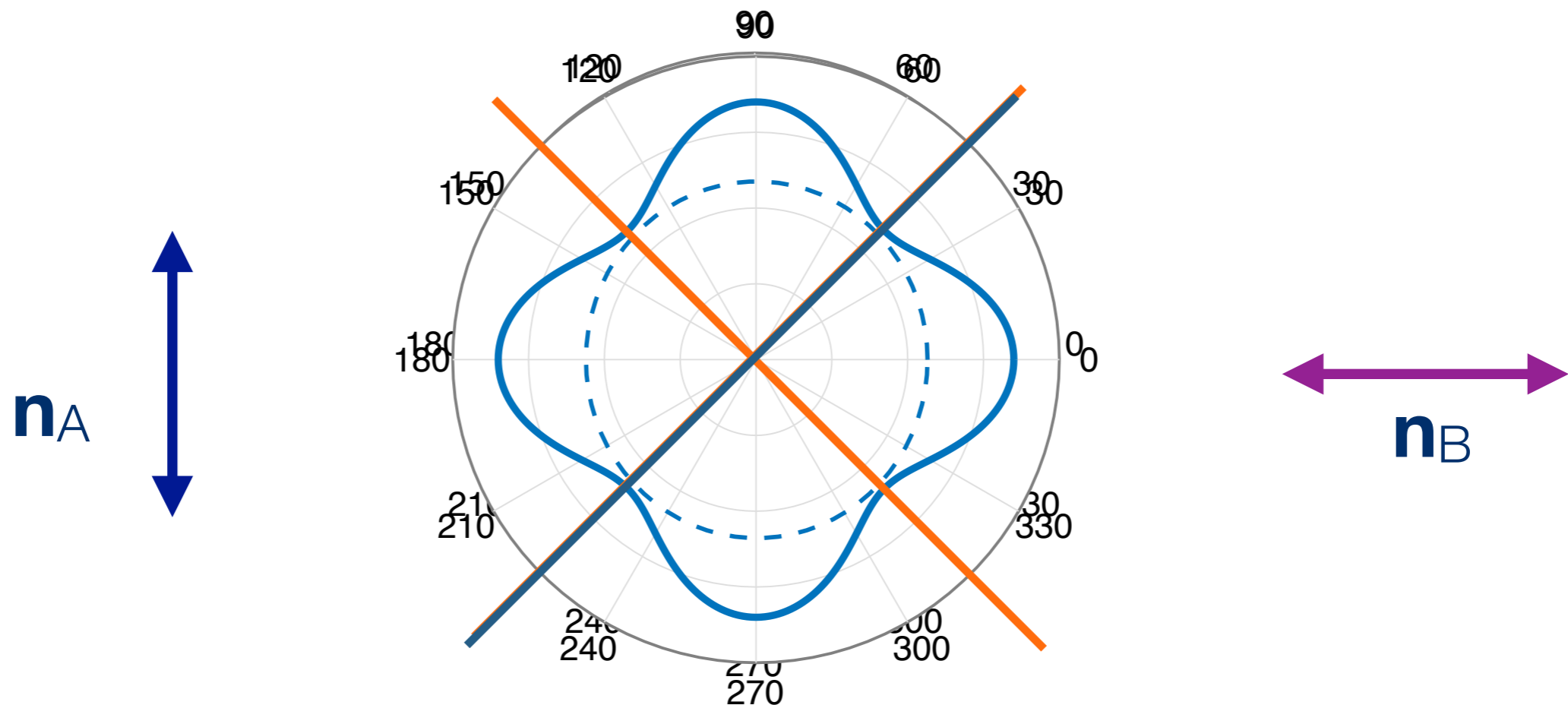
$$\hat{\eta}(\varphi) \neq 0$$



# Domain wall energy

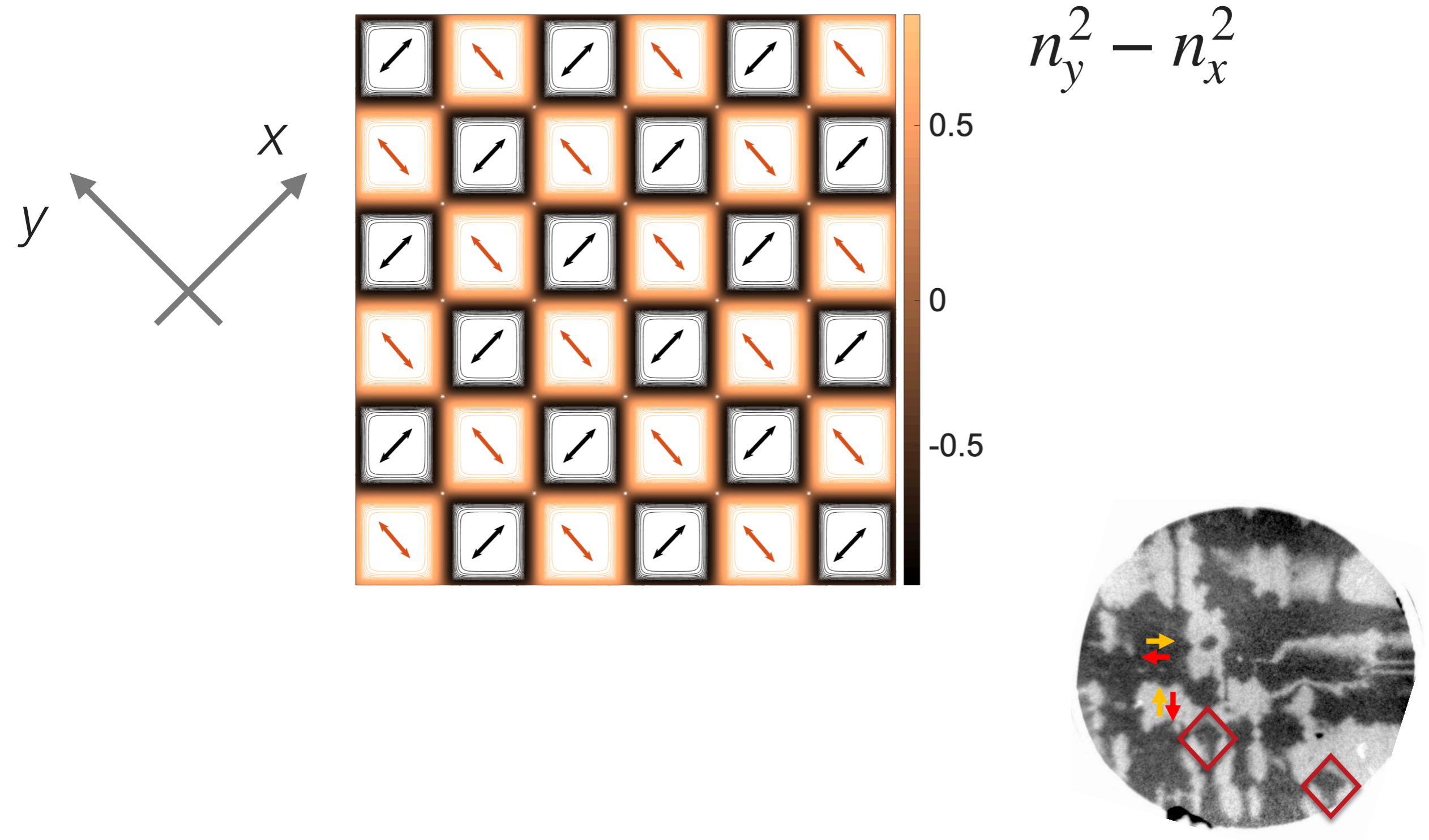
$$E_{\text{mag}} = \sqrt{AH_{\text{an}}M_s}$$

$$E_{\text{m-el}} \propto Cu_{\text{spon}}^2 \cos^2(2\theta)$$



$$E_{\text{DW}} = E_{\text{mag}} + E_{\text{m-el}}$$

# Domain structure in the blanket film



Wadley, Reimers, et al, Nature Nanotechnology

Elastic degrees of freedom => integrated out

$$\hat{u}^{\text{des}}(\mathbf{r}) = \int d\mathbf{r}' \hat{I}(\mathbf{r} - \mathbf{r}') \hat{\eta}(\mathbf{r}')$$

Destressing energy

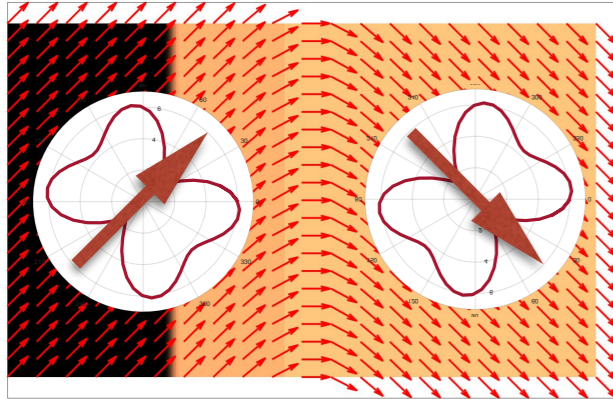
$$E_{\text{des}} = \iint d\mathbf{r}_1 d\mathbf{r}_2 \mathbf{n}_1 \otimes \mathbf{n}_1 \hat{K}(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{n}_2 \otimes \mathbf{n}_2$$

Kernel => dipole

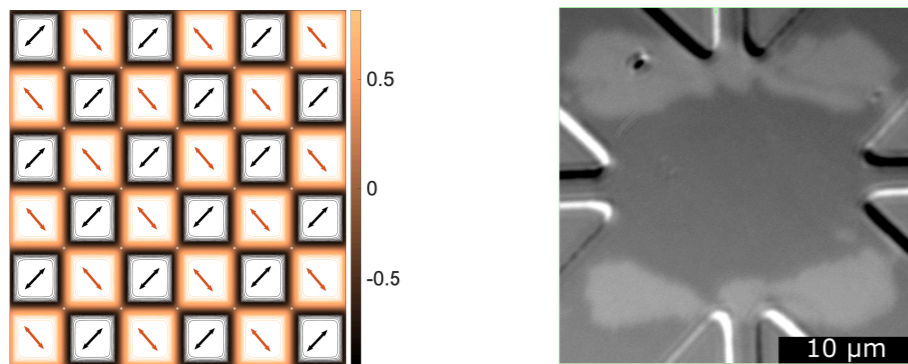
$$\hat{K}(\mathbf{r}) \propto \nabla_j \nabla_k \nabla_l \nabla_m |\mathbf{r}| \propto \frac{1}{|\mathbf{r}|^3}$$



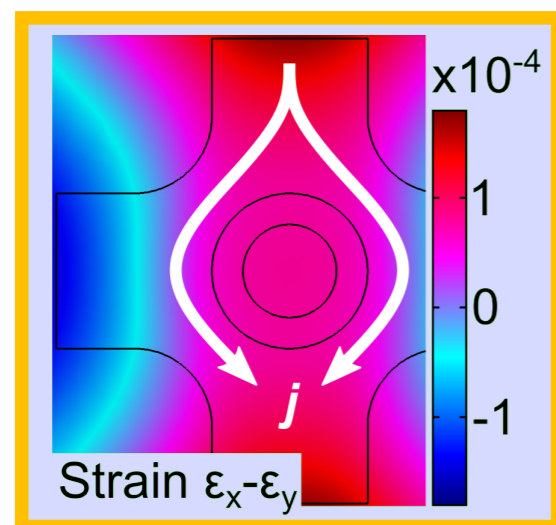




Domain wall dynamics



Domain structure



Domain switching

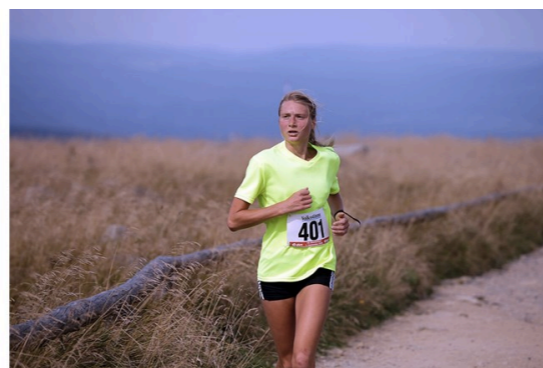
# Collaborators



Pete Wadley



Kevin Edmonds

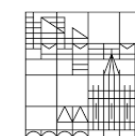


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**THANK YOU!**



Magneto+elastic:

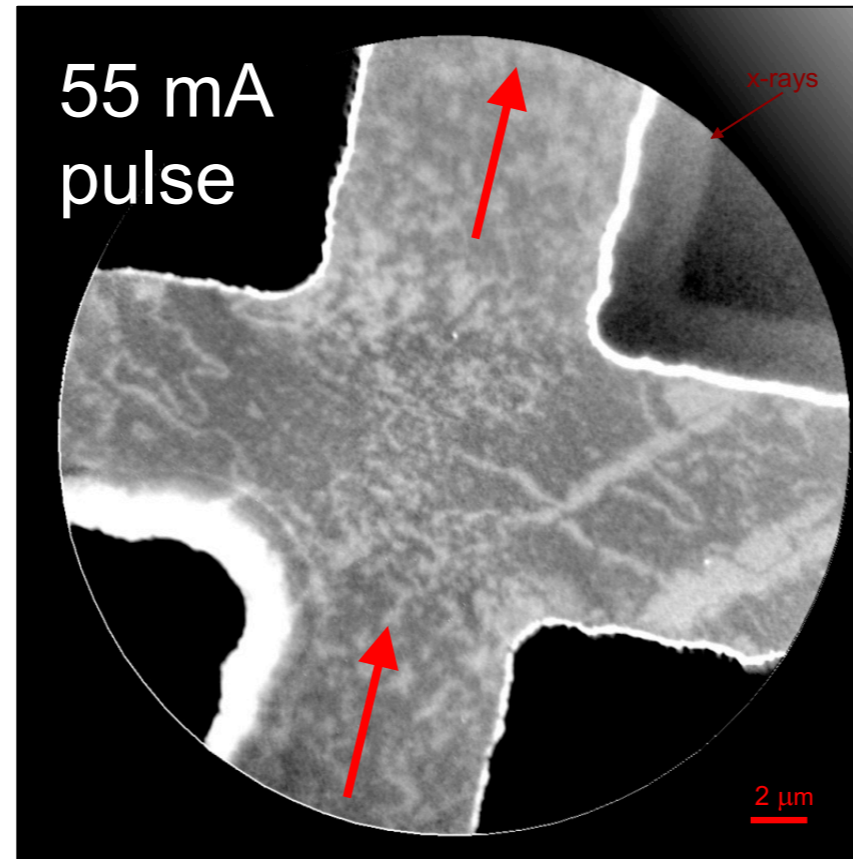
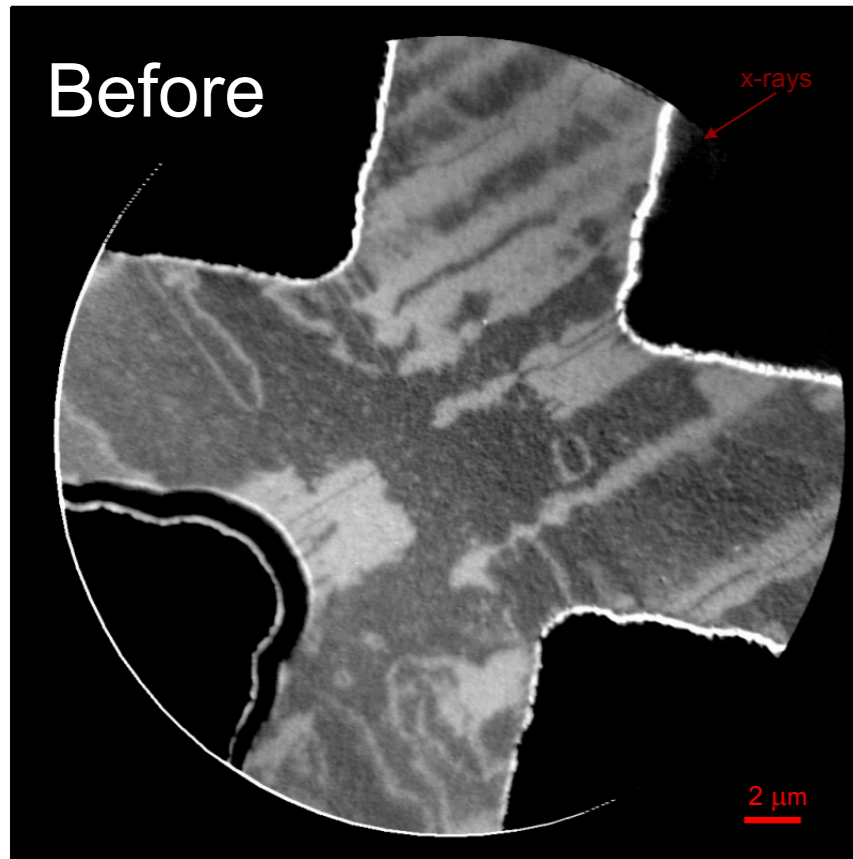
$$A\Delta\mathbf{n} + \mathbf{H}_n = -H_{me}\hat{u}\mathbf{n}$$

$$\nabla(\hat{c}\hat{u}) = -H_{me}\nabla(\mathbf{n} \otimes \mathbf{n})$$

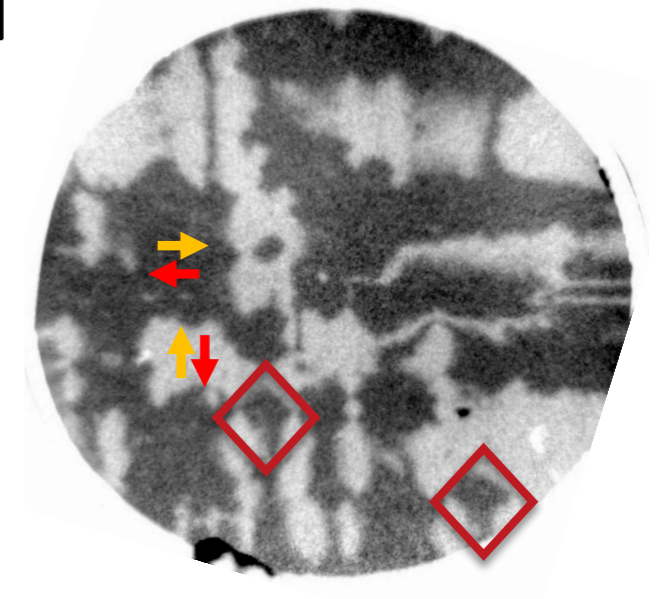
Magnetic plasticity (Kleman, Miltat, 1972)

$$\hat{u} = \underbrace{\hat{u}^{\text{spon}}(\mathbf{r})}_{\text{dislocations}} + \underbrace{\hat{u}^{\text{des}}(\mathbf{r})}_{\text{plastic}}$$

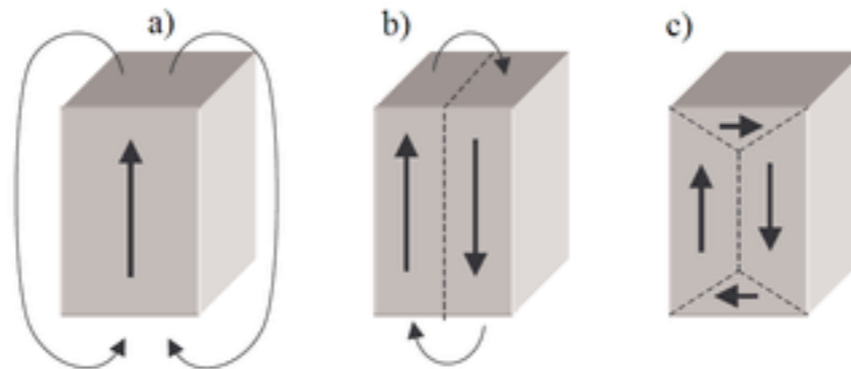




Wadley, Reimers, et al, Nature Nanotechnology,



$$E_{\text{demag}} = \iint d\mathbf{r}_1 d\mathbf{r}_2 \frac{\nabla \cdot \mathbf{M}(\mathbf{r}_1) \nabla \cdot \mathbf{M}(\mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|}$$



	FM	AF
Order parameter	<b>M</b>	<b>N</b>
Conjugated field	<b>H</b>	—
“Charge”	div <b>M</b>	—
Long range field	<b>H</b> <sub>dip</sub>	—

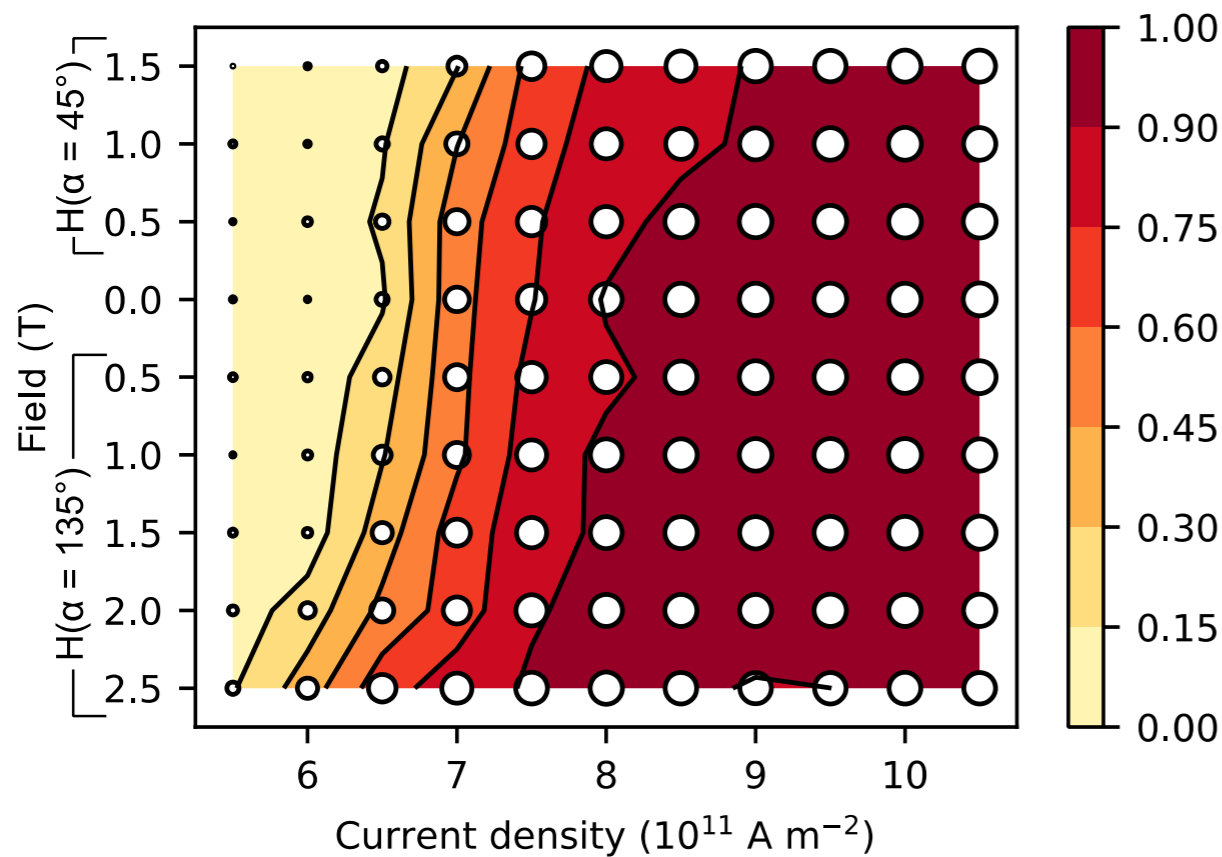


# Possible solution

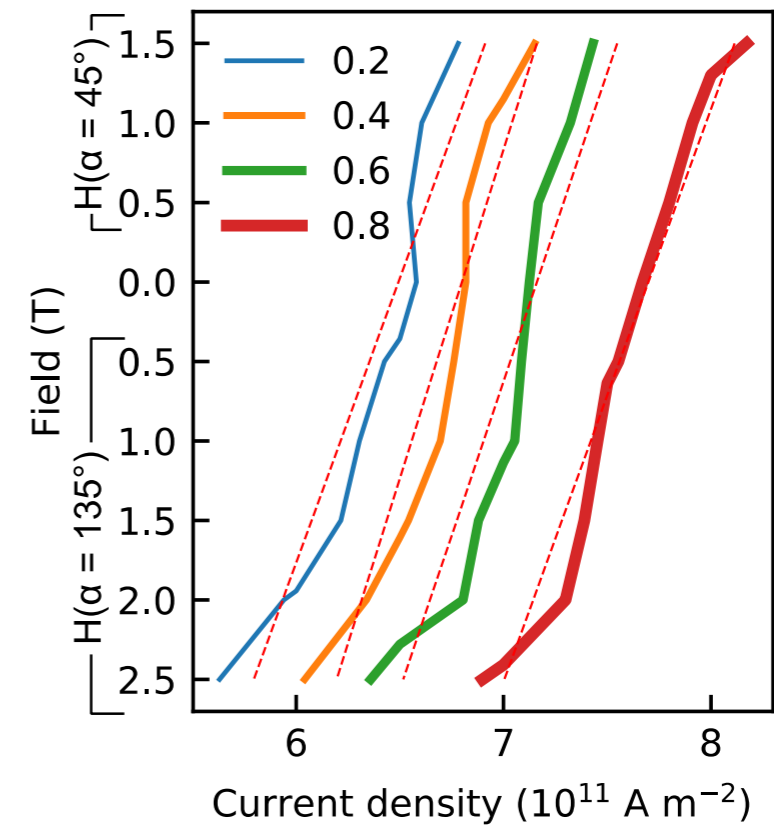
	FM	AF+ ME
Order parameter	$\mathbf{M}$	$\mathbf{N} \otimes \mathbf{N} \propto \hat{u}^{\text{spon}}$
Conjugated field	$\mathbf{H}$	$\hat{\sigma}$
“Charge”	$\text{div } \mathbf{M}$	$\text{inc } \hat{u}^{\text{spon}}$
Long range field	$\mathbf{H}_{\text{dip}}$	$\hat{\sigma}^{\text{destr}}$



# Field-current equivalence in CoO



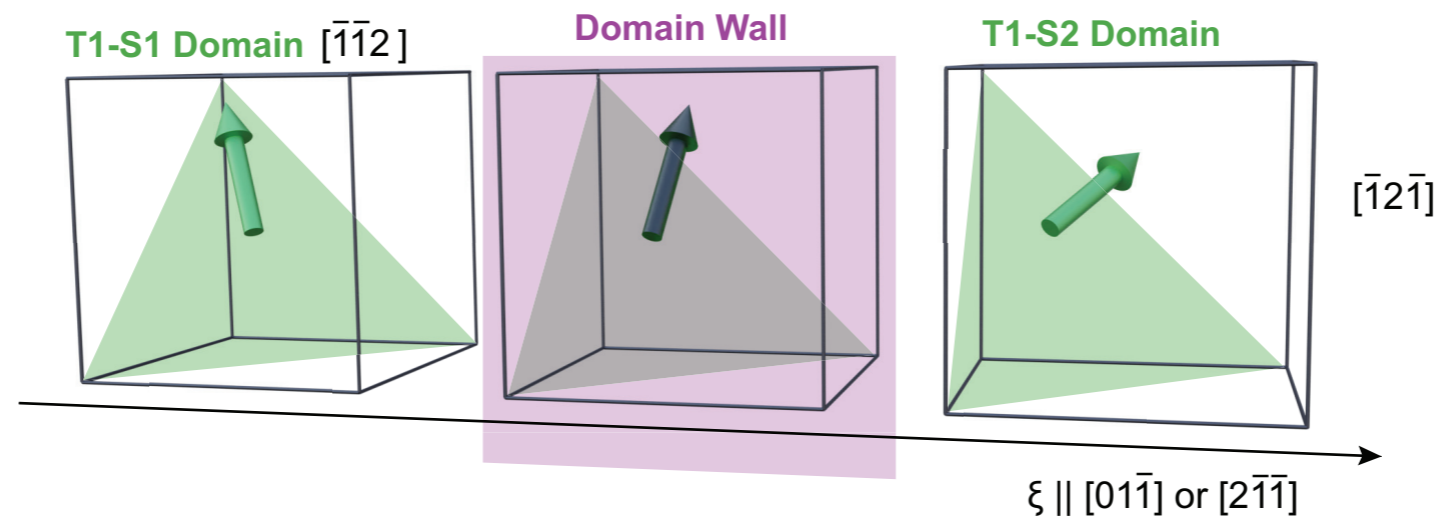
$$U_{\text{ext}} \propto (\mathbf{H} \cdot \mathbf{n})^2$$



$$U_{\text{ext}} \propto \lambda(\mathbf{j} \cdot \mathbf{n})^2$$

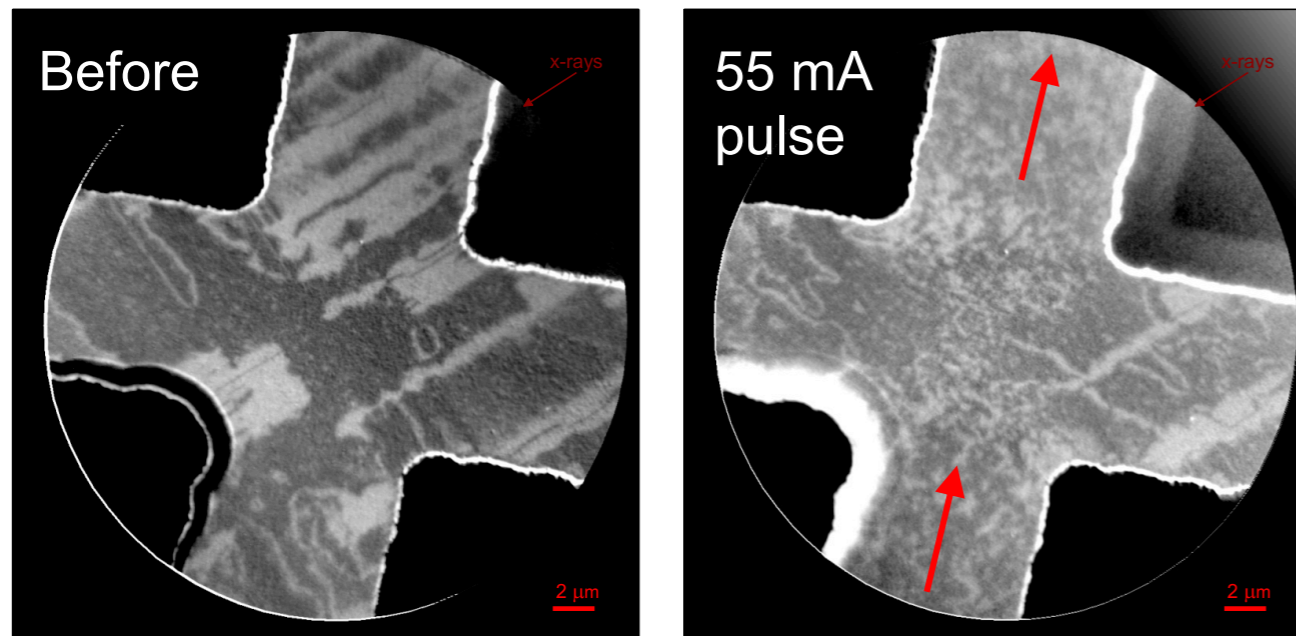
Baldrati, OG, et al. Phys. Rev. Lett. 125, 77201 (2020)

$$\mathbf{n} \times (\ddot{\mathbf{n}} - c^2 \Delta \mathbf{n} + \gamma^2 H_{\text{ex}} \mathbf{H}_{\text{an}}(\xi)) = 0$$



$$\omega_{\text{DW}} \propto \sqrt{H_{\text{me}}}$$

## CuMnAs

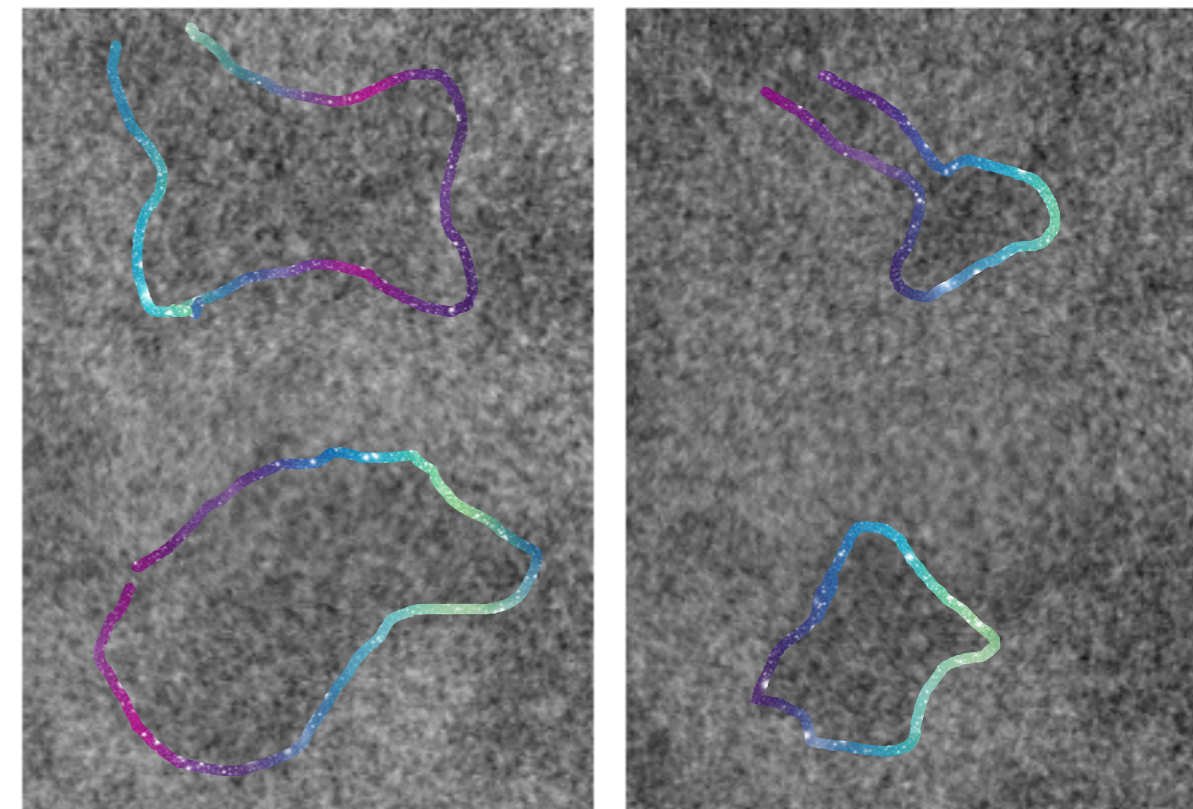


Wadley, Reimers et al *Nat. Nanotech.* 18  
Wadley, Reimers, et al

## NiO

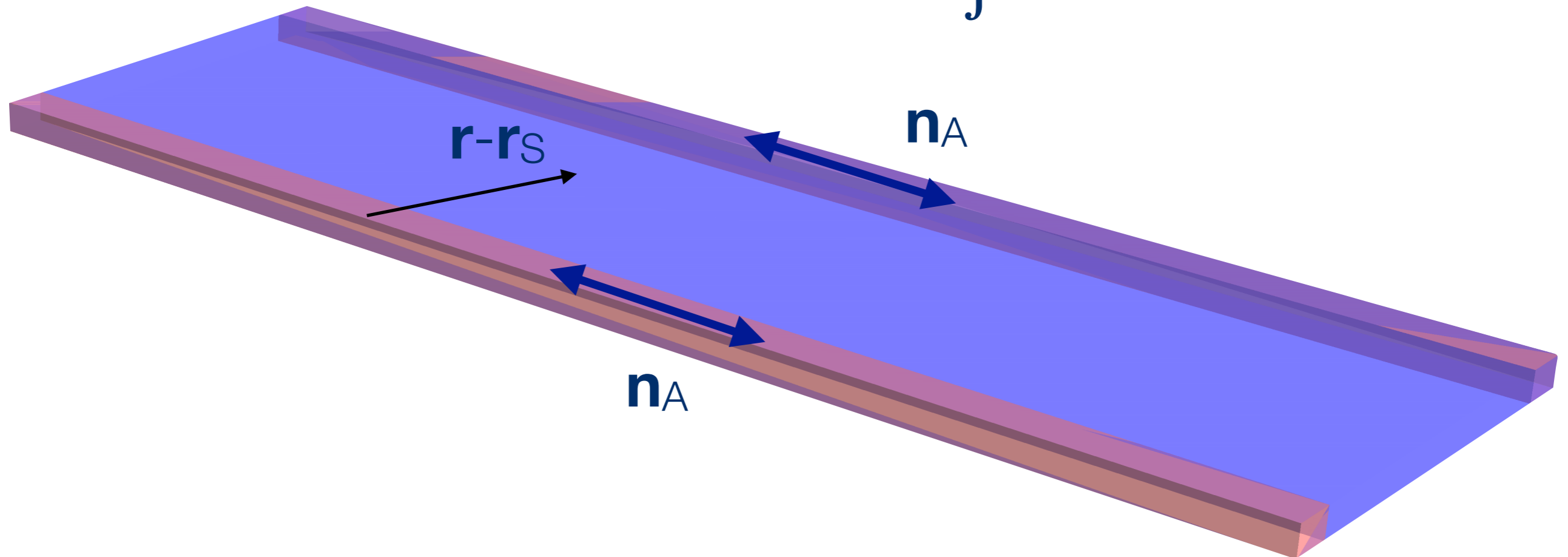
(a) Before pulse

(b) After 5x pulse



Baldrati, OG et al *PRL* 19

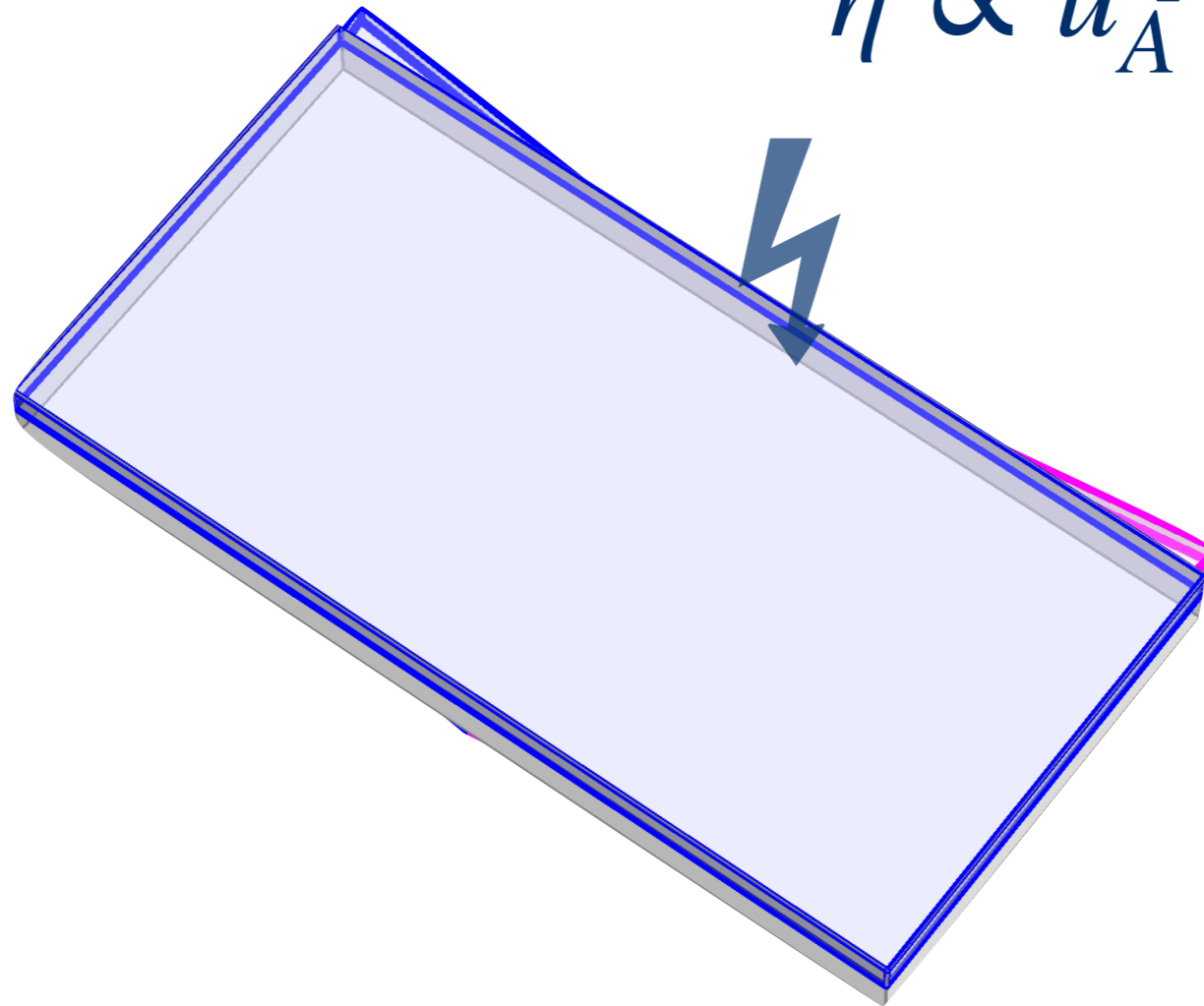
$$W_{\text{surf}} = K_S \oint dS (\mathbf{n} \cdot \mathbf{e})^2$$

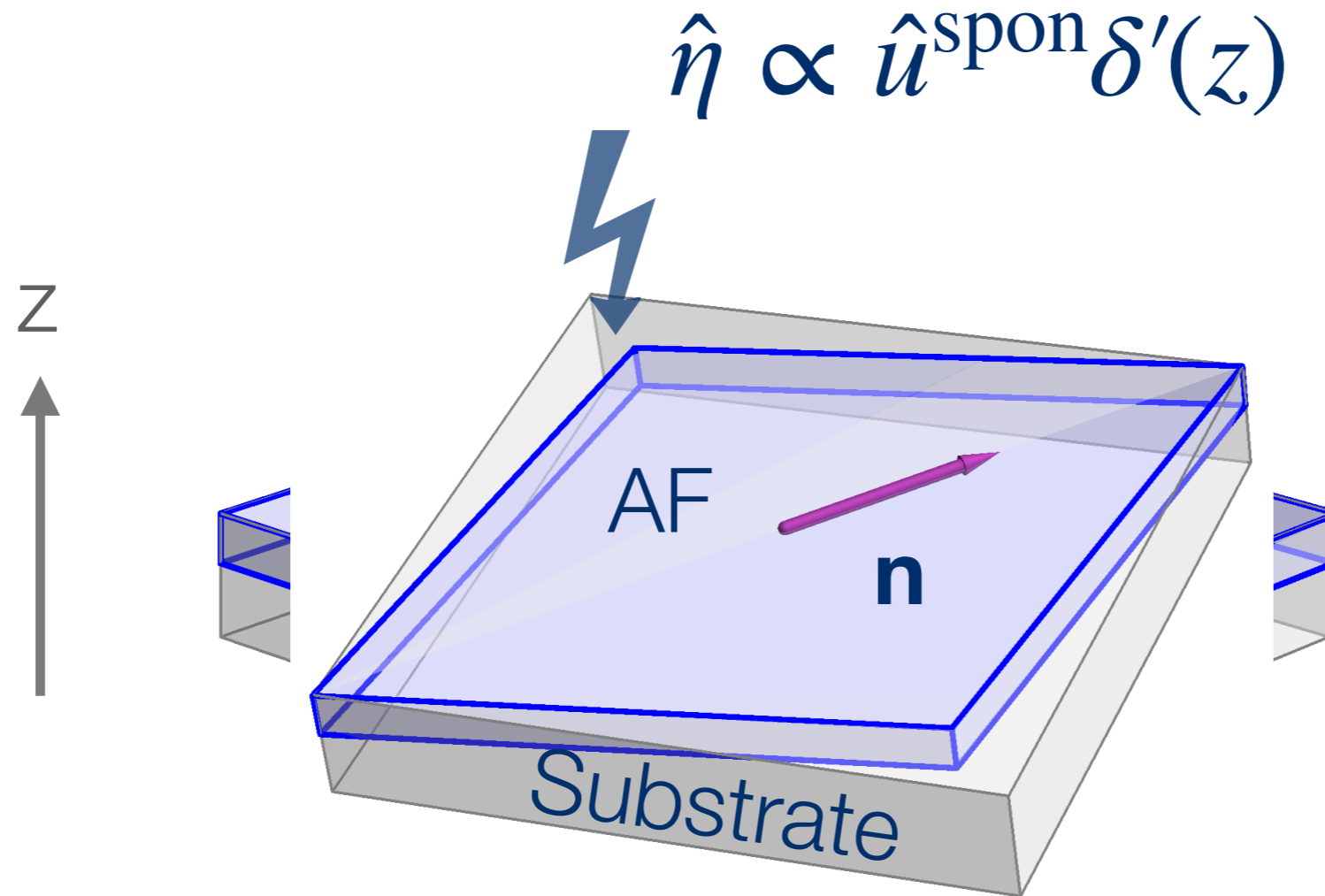


$$E_{\text{surf}} = \int_V (\mathbf{n} \otimes \mathbf{n})(\mathbf{e} \cdot \nabla) \int_S \hat{I}(\mathbf{r} - \mathbf{r}_S) \mathbf{n}_S \otimes \mathbf{n}_S d\mathbf{r}_S$$

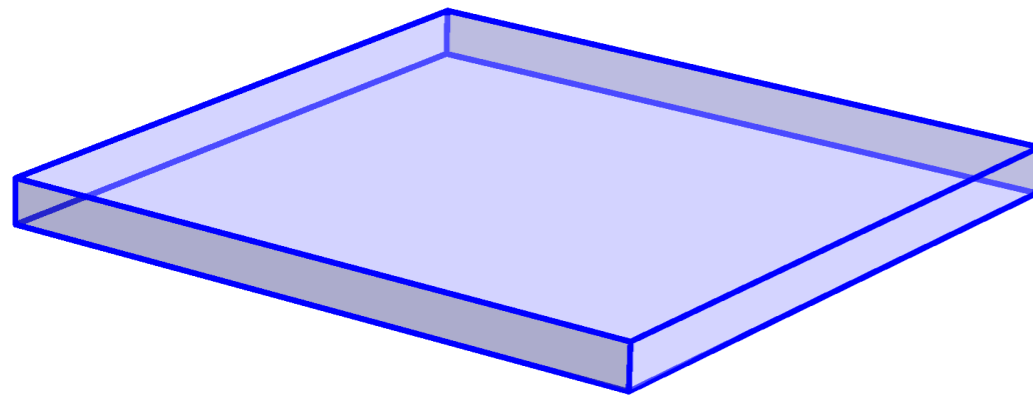


$$\hat{\eta} \propto \hat{u}_A^{\text{spon}} - \hat{u}_B^{\text{spon}}$$

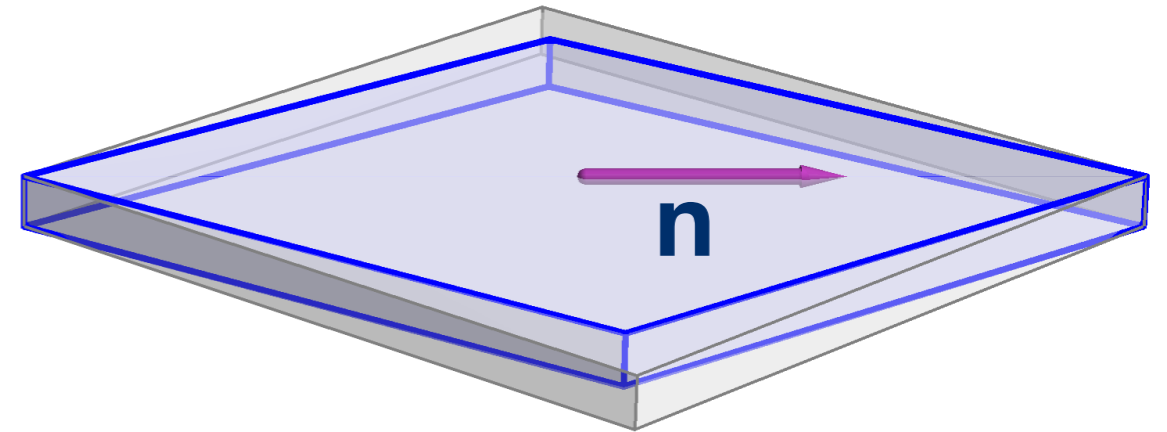




Paramagnetic state



AF state



$$\hat{u}^{\text{spon}} \propto H_{m-e} \mathbf{n} \otimes \mathbf{n}$$

Elastostatic in terms of incompatibility

$$\nabla \times \hat{u}^{\text{des}}(\mathbf{r}) \times \nabla = -\hat{\eta}$$

$$\hat{\eta}(\mathbf{r}) = \nabla \times \hat{u}^{\text{spon}}(\mathbf{r}) \times \nabla$$

$$\hat{I}(\mathbf{r}) \propto \nabla_j \nabla_k |\mathbf{r}| \propto \frac{1}{|\mathbf{r}|}$$

$$\hat{u}^{\text{spon}} \propto H_{m-e} \mathbf{n} \otimes \mathbf{n}$$





$$\hat{u}^{\text{des}} = \int d\mathbf{r}' \hat{I}(\mathbf{r} - \mathbf{r}') \hat{\eta}(\mathbf{r}')$$

$$\hat{\eta}(\mathbf{r}) = \nabla \times \hat{u}^{\text{spon}}(\mathbf{r}) \times \nabla$$

$$\hat{I}(\mathbf{r}) \propto \nabla_j \nabla_k |\mathbf{r}| \propto \frac{1}{|\mathbf{r}|}$$

$$\hat{u}^{\text{spon}} \propto H_{\text{m-e}} \mathbf{n} \otimes \mathbf{n}$$



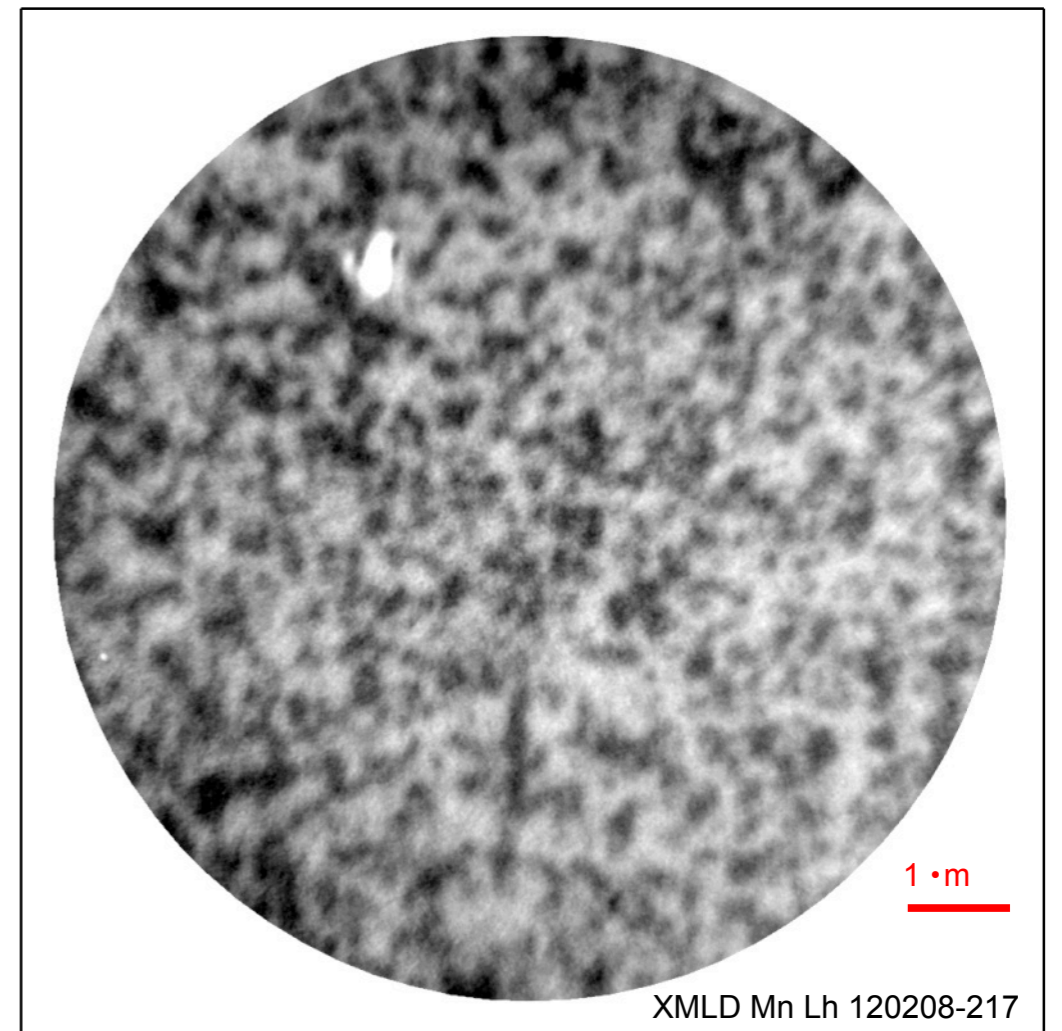
$$w_{\text{me}} = H_{\text{me}} M_s \left[ (u_{xx} - u_{yy})(n_x^2 - n_y^2) + 2u_{xy}n_xn_y \right]$$

Frozen spontaneous strains

$$\varepsilon^{\text{spon}} = - \frac{H_{\text{me}} M_s}{\mu} (n_x^2 - n_y^2) |_0$$

$$w_{\text{me}} = H_{\text{me}} M_s \varepsilon^{\text{spon}} (n_x^2 - n_y^2)$$

# Motivation



15 μm field view



**b**

