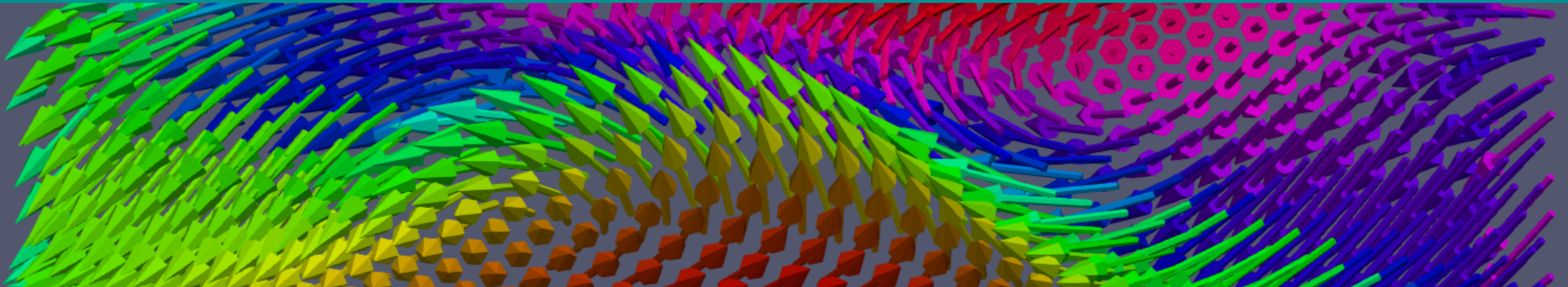




# Quantum Magnonics: quantum optics with magnons



**Silvia Viola Kusminskiy**



**MAX PLANCK INSTITUTE**  
FOR THE SCIENCE OF LIGHT



**FRIEDRICH-ALEXANDER**  
UNIVERSITÄT  
ERLANGEN-NÜRNBERG

# Max Planck Research Group



MAX PLANCK INSTITUTE  
for the science of light

Erlangen, Germany



# Theory of hybrid quantum systems



Sanchar



Silvia Viola Kusminski



Victor Bittencourt



Anna-Luisa Römling



Fabian Engelhardt



Maximilian Hollendonner



Kai



jgraf



Helene

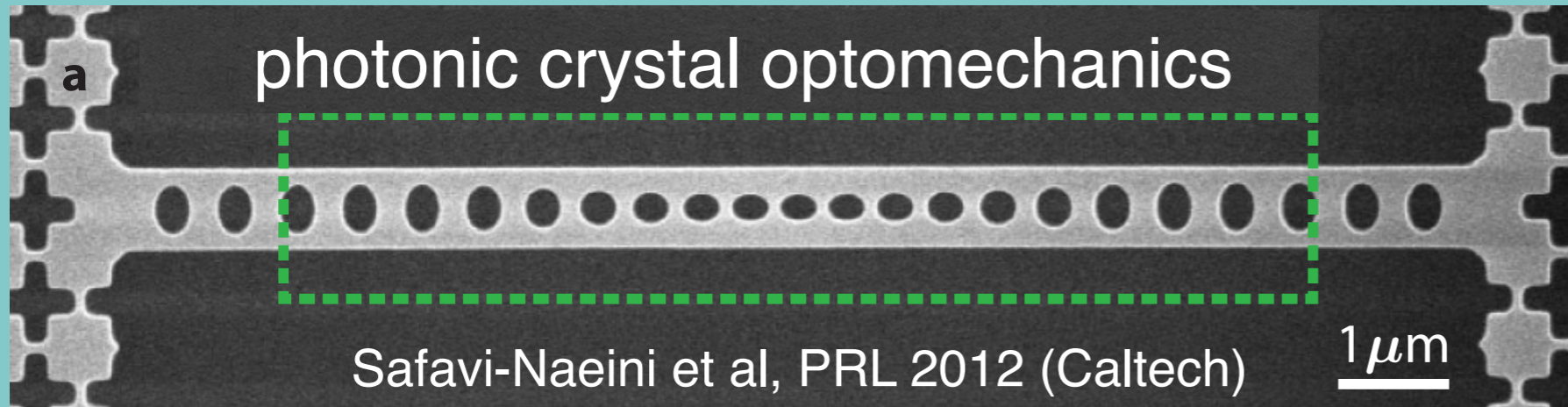


Vanessa Wachter

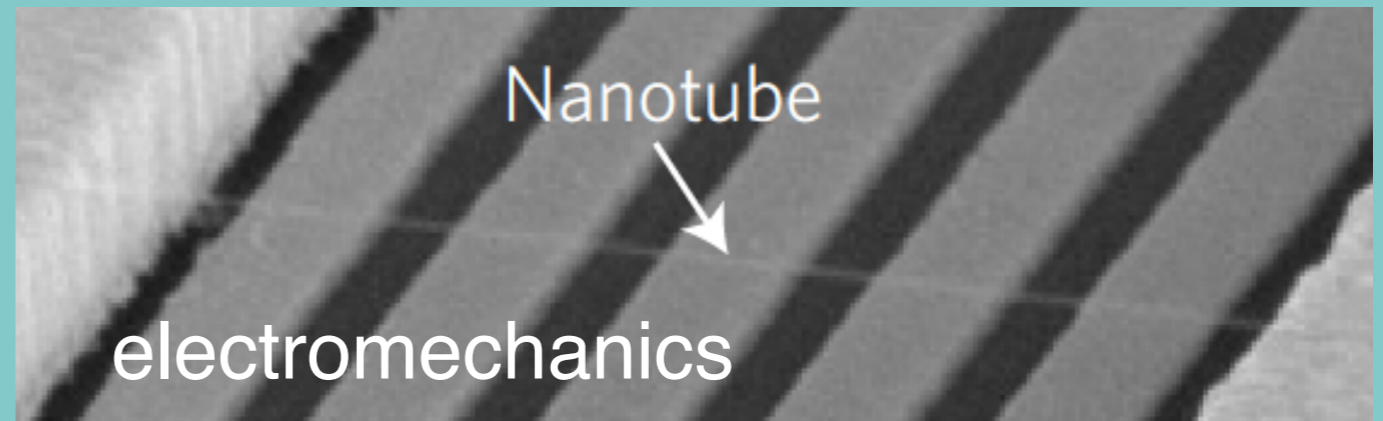
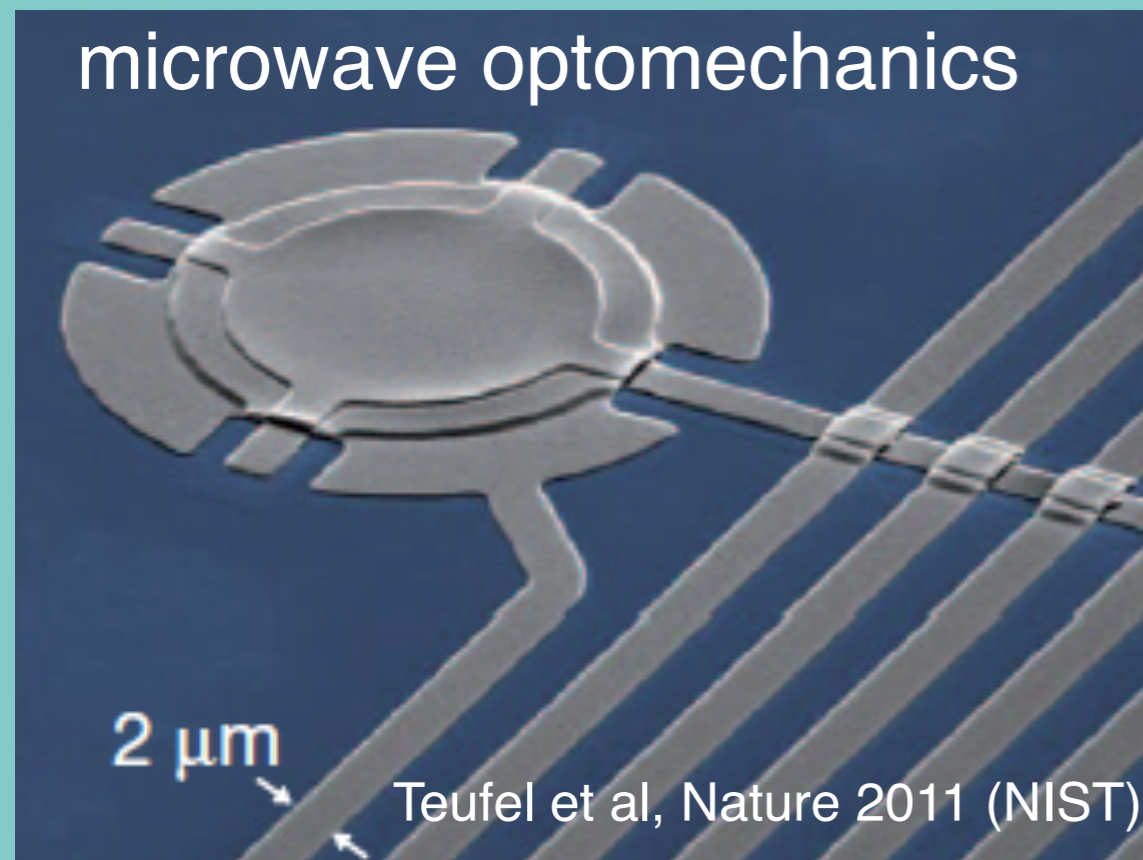


tparvin

# Examples of Hybrid Quantum Systems



nano/micrometer  
scale systems



Benyamini et al, Nature Physics 10, 151 (2014)



Osada et. al PRL 116, 223601 (2016)

use collective excitations



# Applications: Quantum Technologies

**Process**



**Communicate**

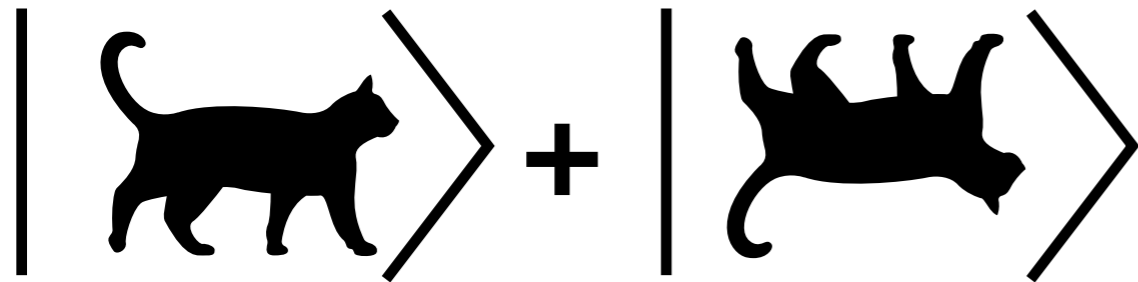


**Quantum Information**

**Store**



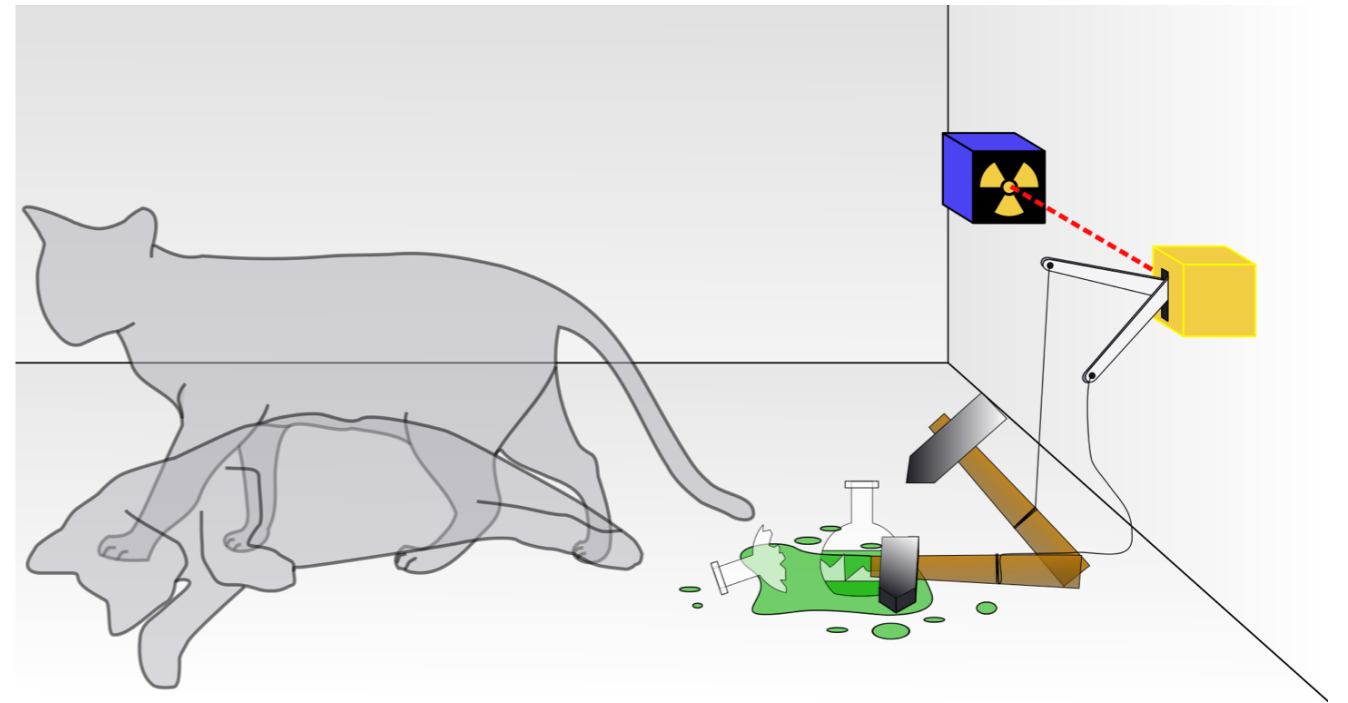
# Fundamentals: How macroscopic can a quantum state be?



**Schrödinger's cat**



## Ultimate hybrid system!



### Schrödinger, Gedanken experiment (1935)

*“Man kann burleske Fälle konstruieren.”*

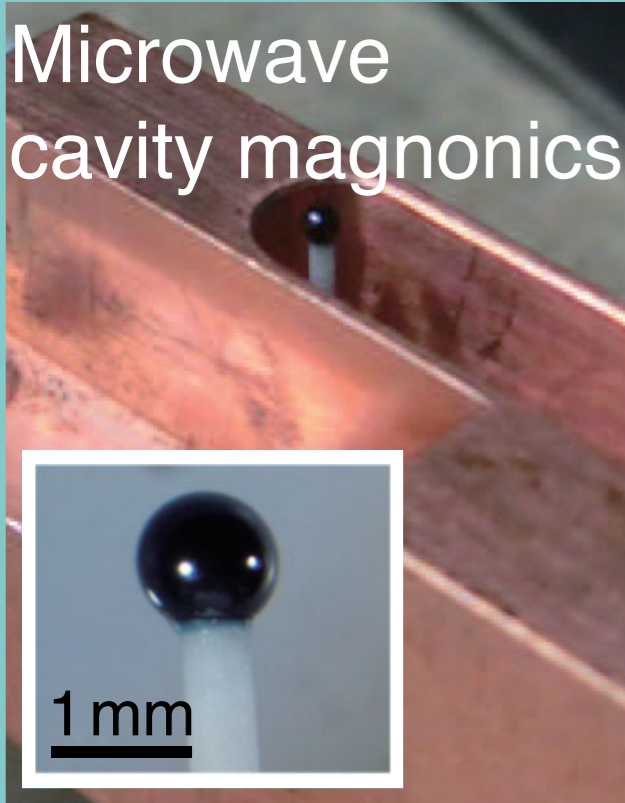
*“One can construct absurd cases.”*

“Die Naturwissenschaften”, 1935

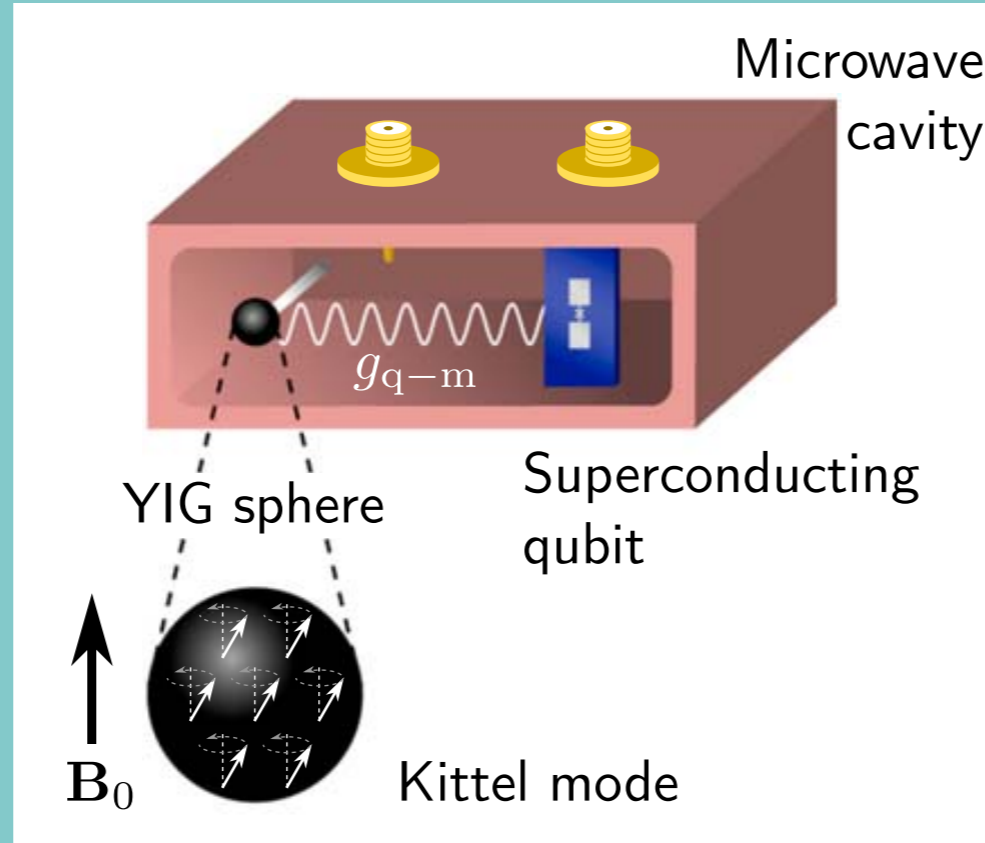


# Hybrid quantum systems based on magnetic elements

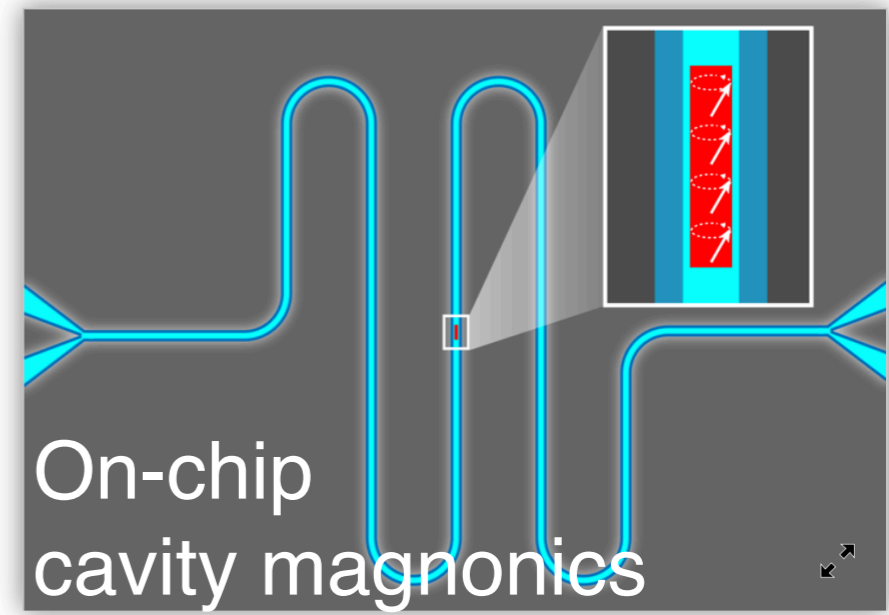
Microwave cavity magnonics



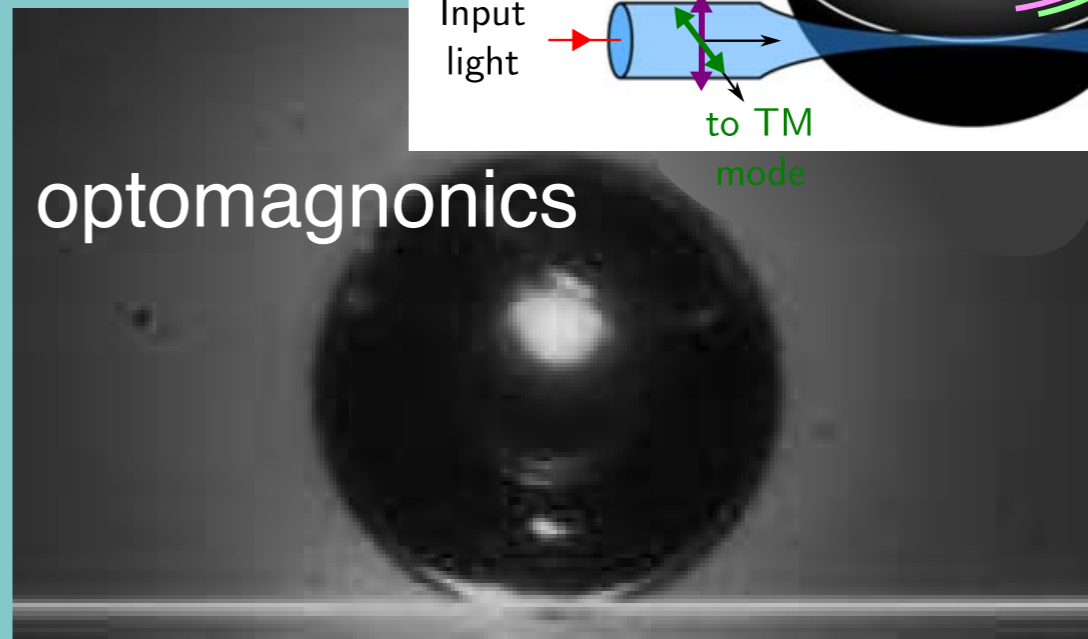
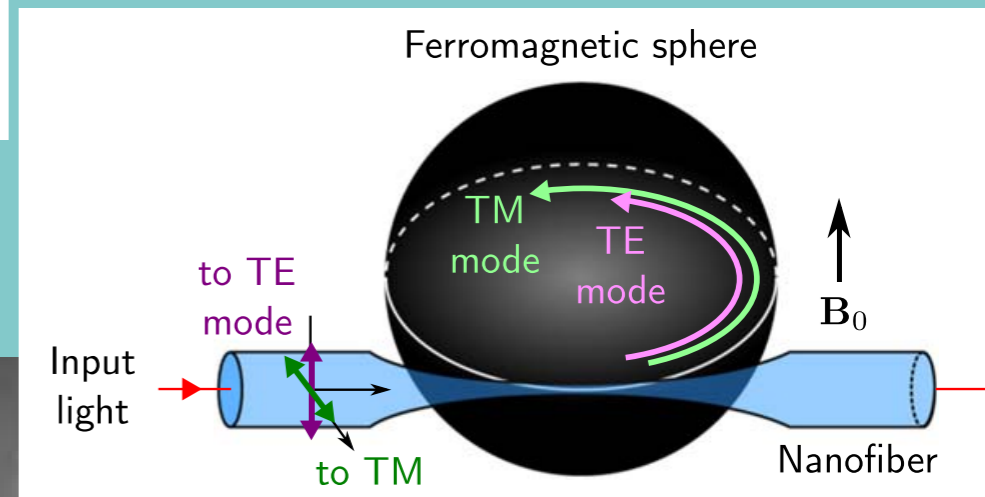
Tabuchi et. al PRL 2014



Lachance Quirion et. al Appl. Phys. Express 2019

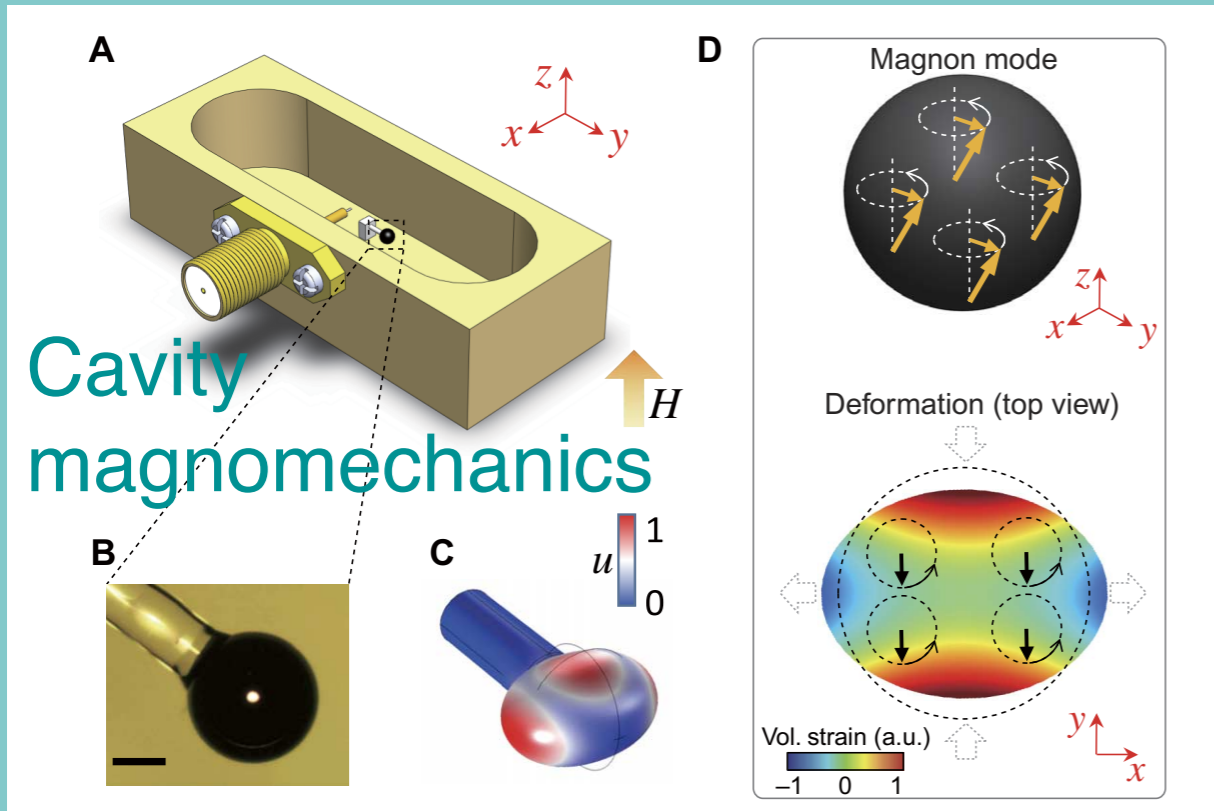


APS/Alan Stonebraker



Osada et. al PRL 116 2016

Cavity magnomechanics



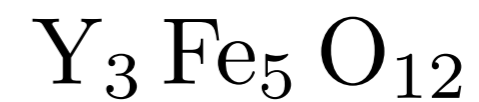
Zhang et. al Science Advances 2016

# YIG



## YIG

Yttrium Iron Garnet



- ferrimagnetic
- insulator
- transparent in the infrared



# What is a magnon?

**magnon**

elementary magnetic  
excitation  
(quantum of spin wave)

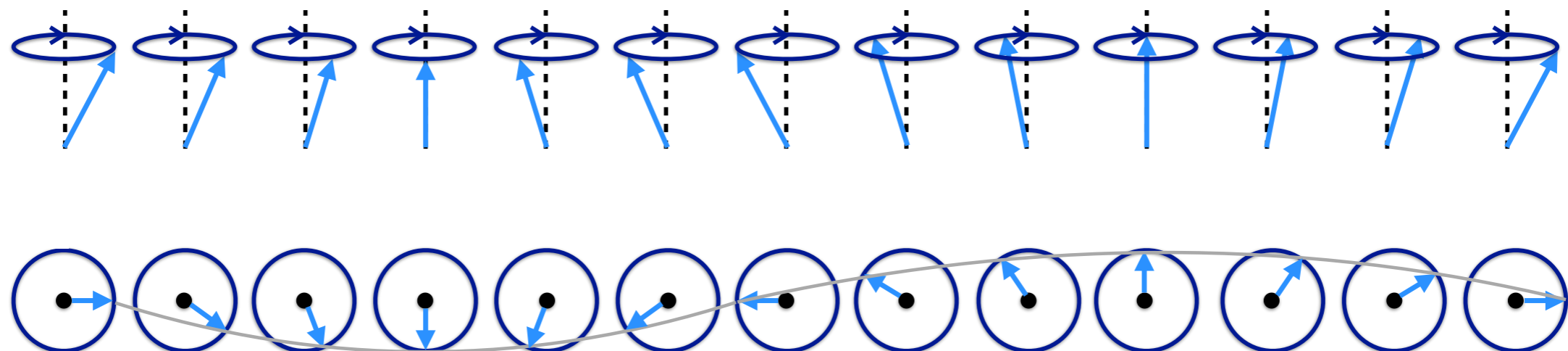


image: Jasmin Graf

# Why do we like magnons?



**magnon**

Rich physics

elementary magnetic  
excitation  
(quantum of spin wave)

Amenable  
to miniaturization

Robust

Low Power

Tunable



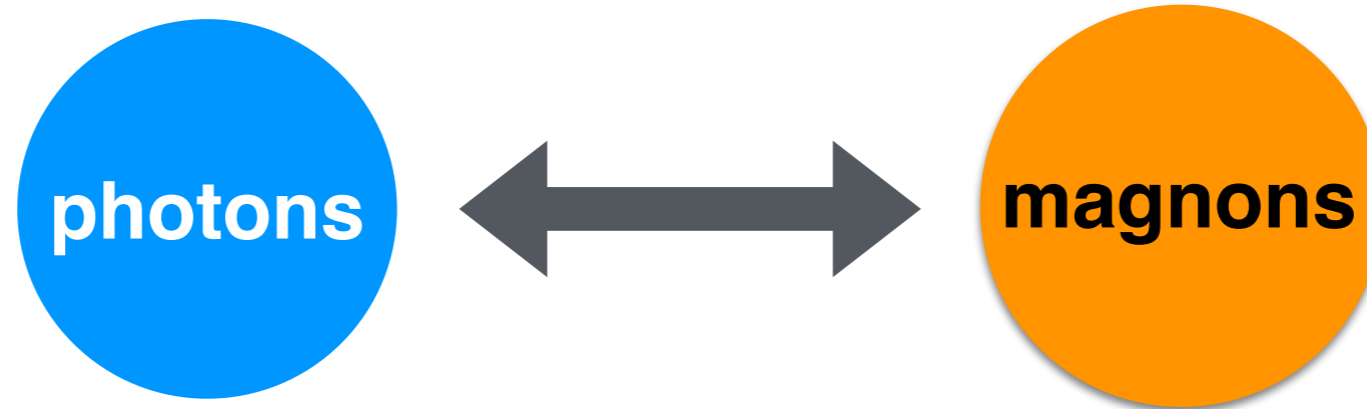
# Why do we like magnons?



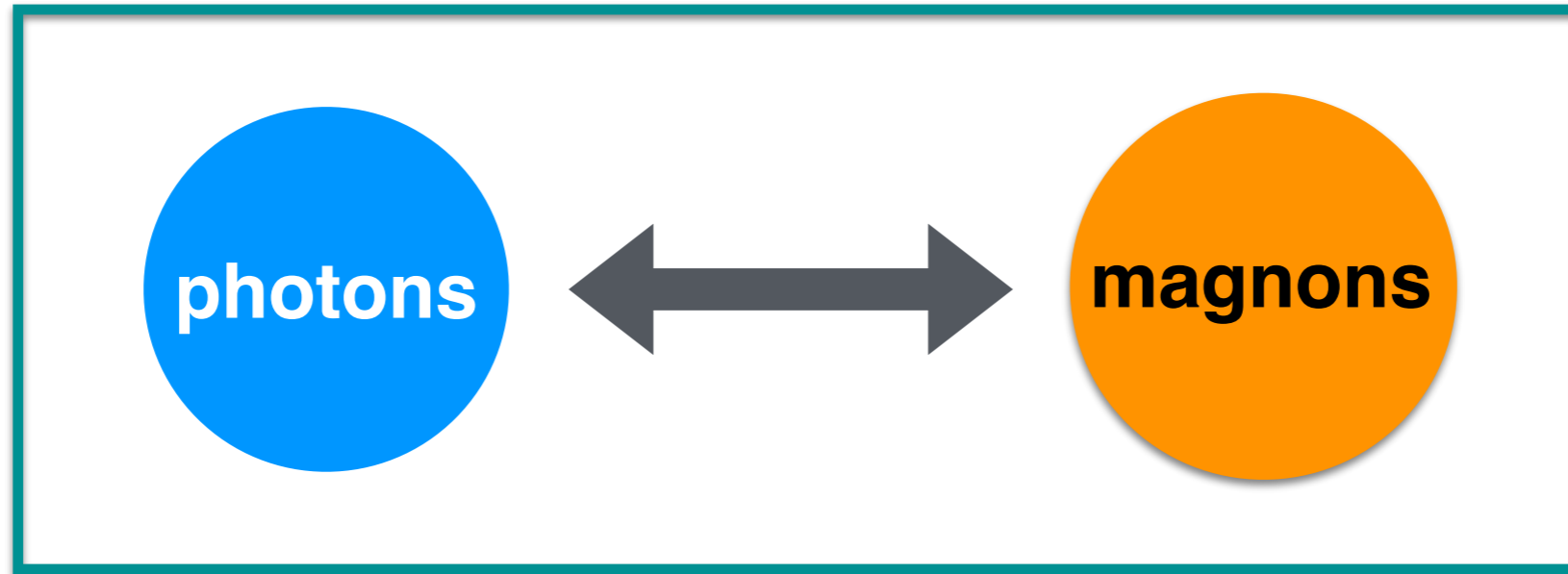
Figure: m&m's

Magnons are friendly: coupling to electrons, phonons, photons...

# Optomagnonics



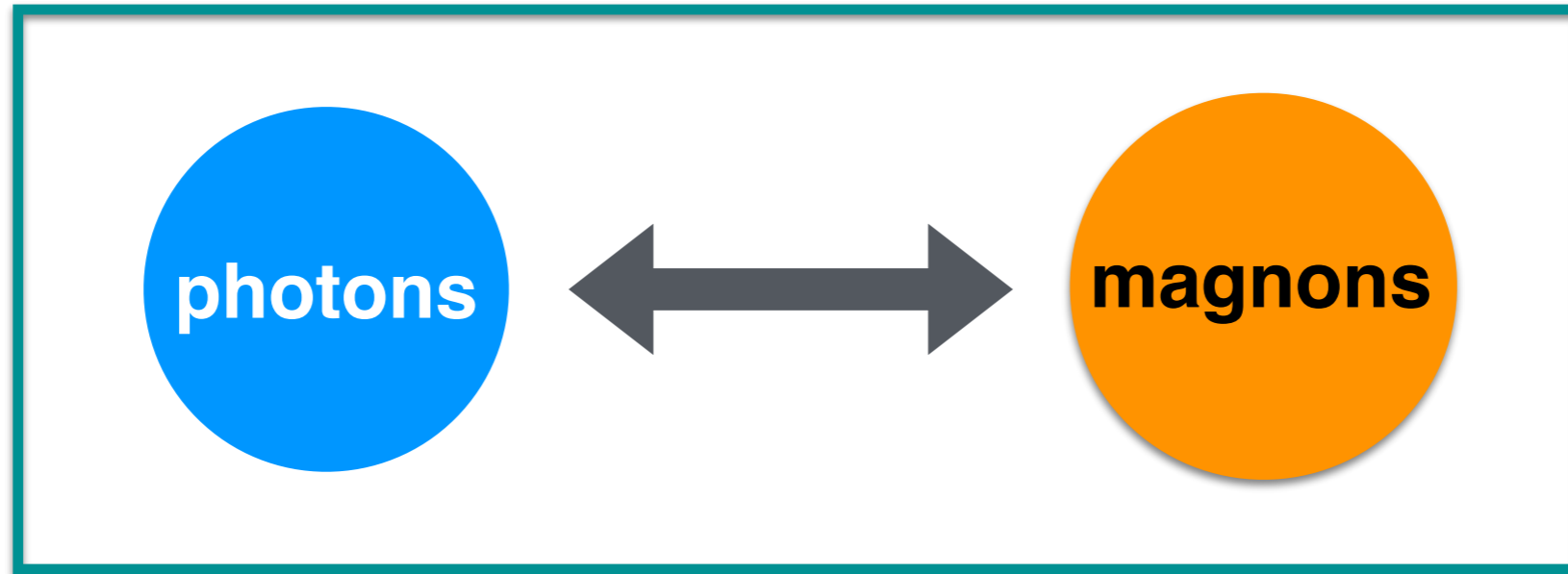
# Cavity Optomagnonics



cavity-enhanced spin-photon interaction



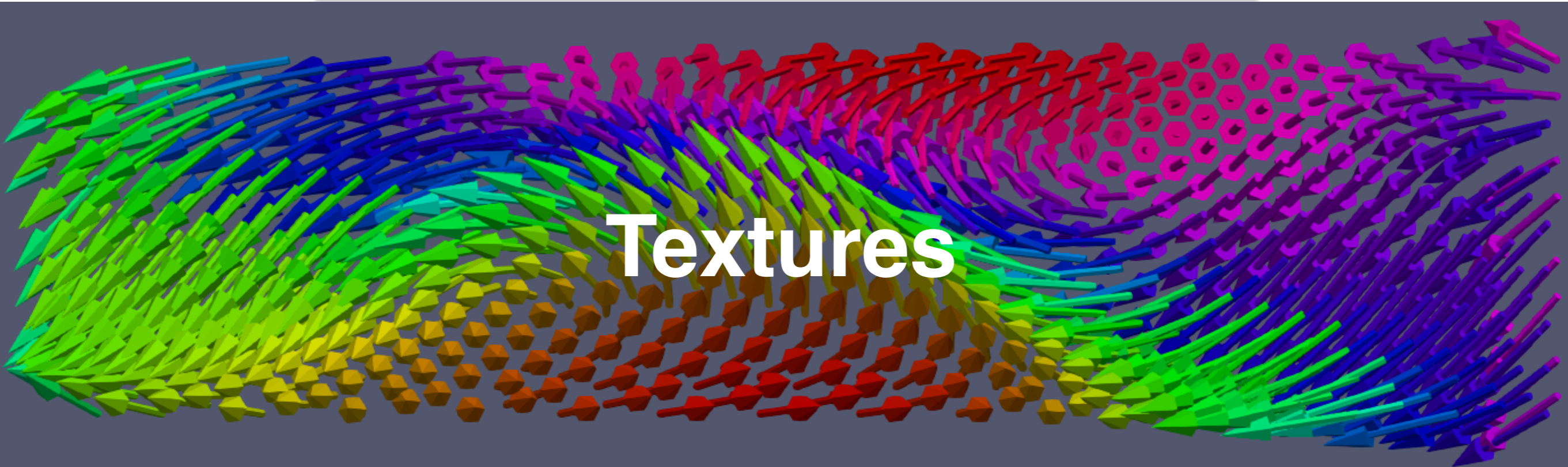
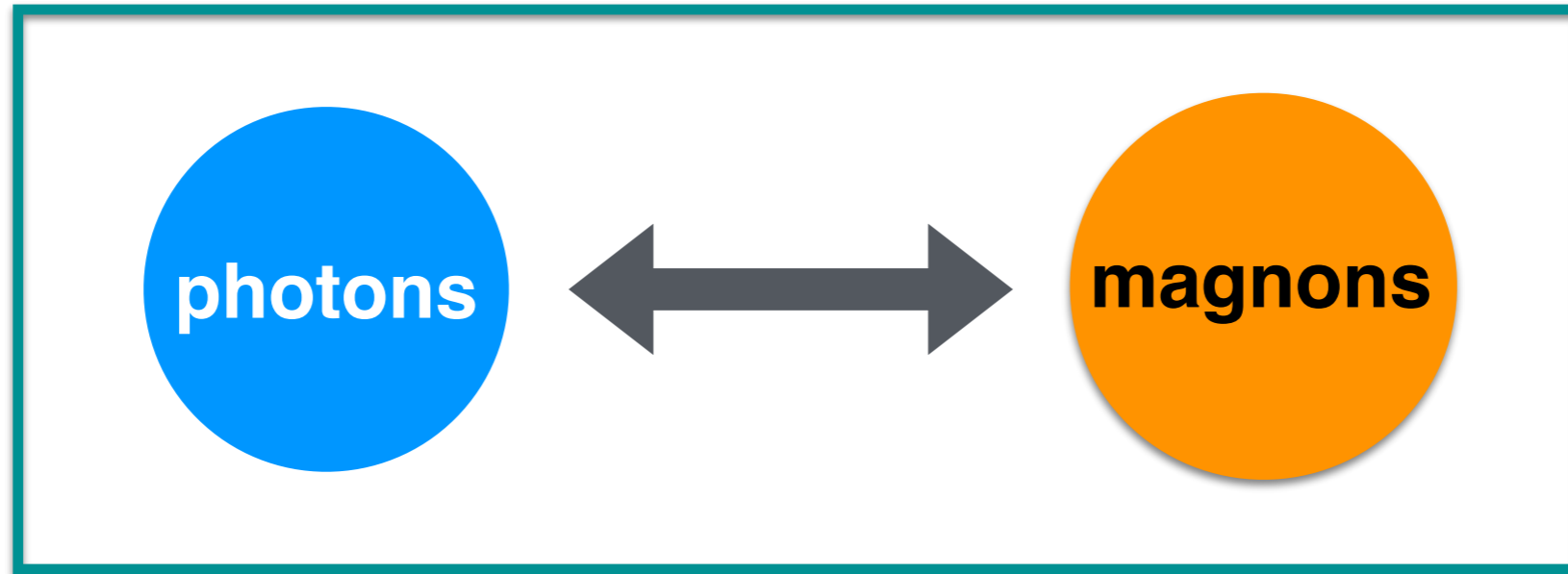
# Cavity Optomagnonics



cavity-enhanced spin-photon interaction

cavity QED + magnetism

# Cavity Optomagnonics



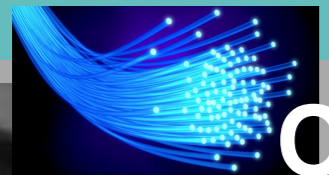
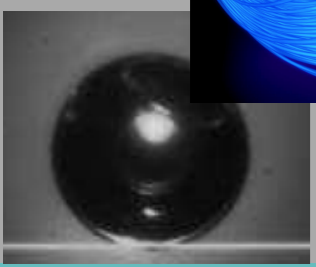
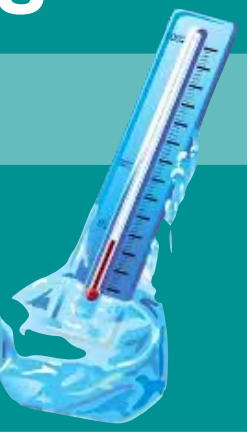


**Microwave regime**

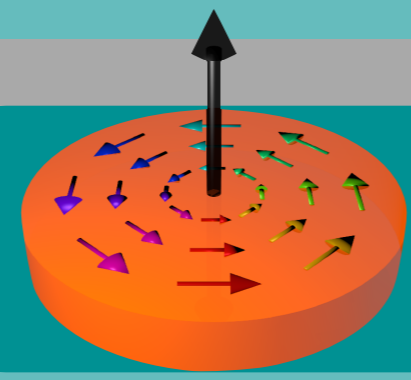


**Spin cat states in ferromagnetic insulators**

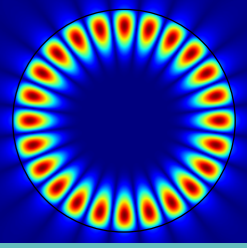
**Magnon-phonon quantum thermometry**



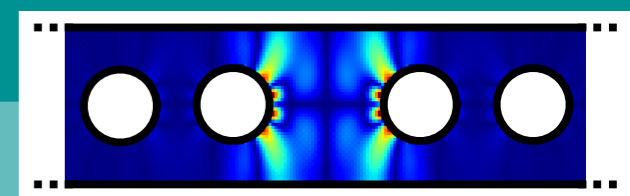
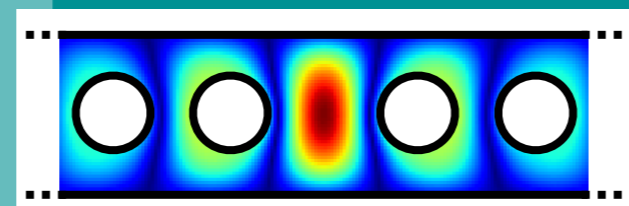
**Optical regime**



**Optomagnonics with Magnetic textures**



**Optomagnonic crystals**



**Outlook**





# Microwave regime

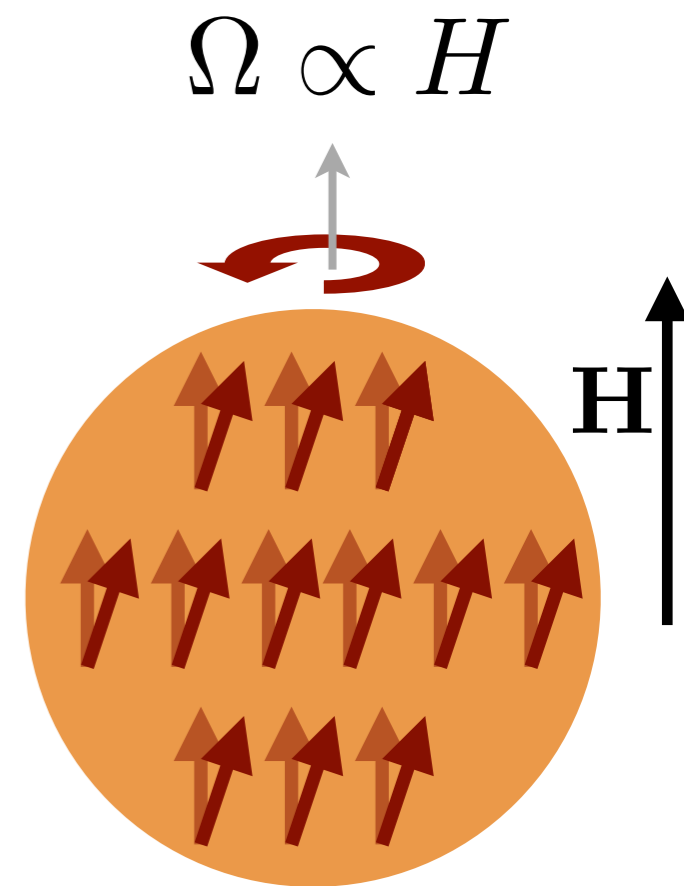
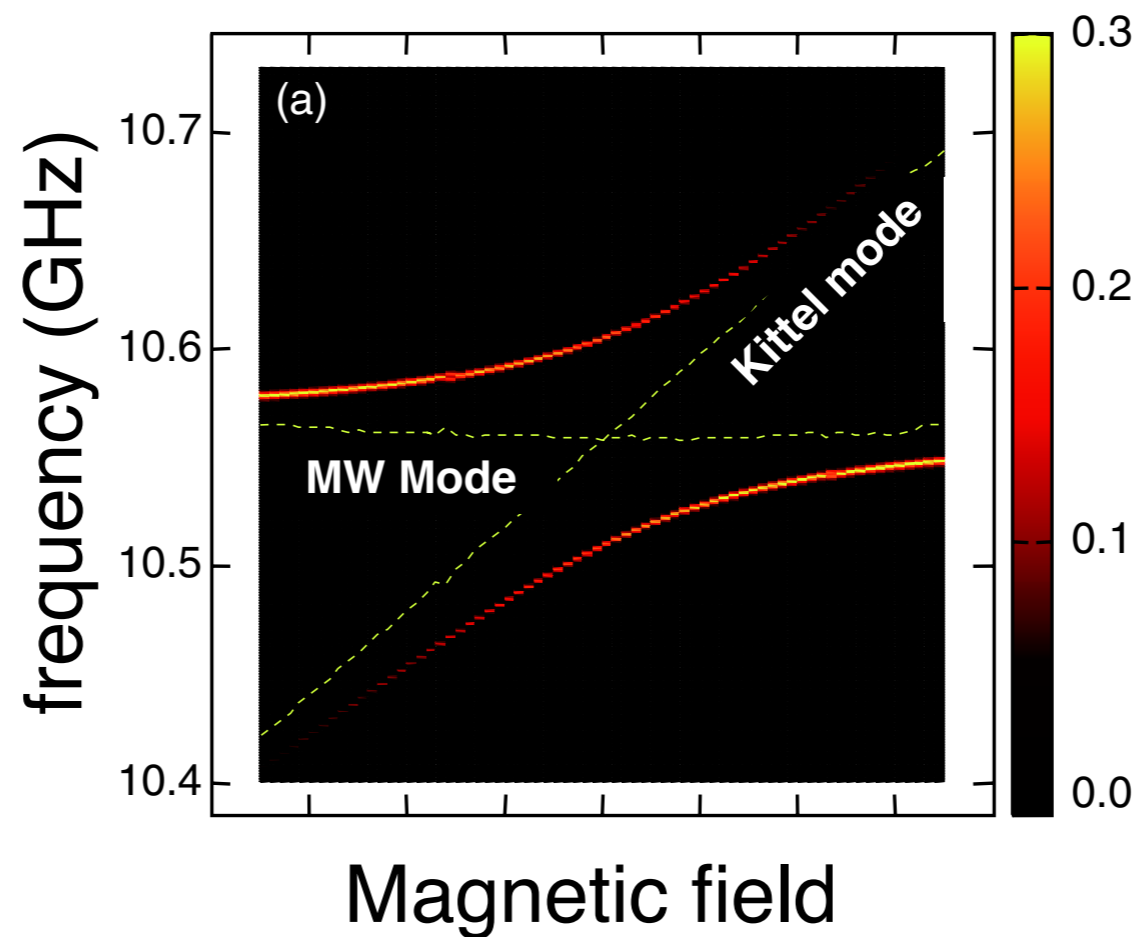
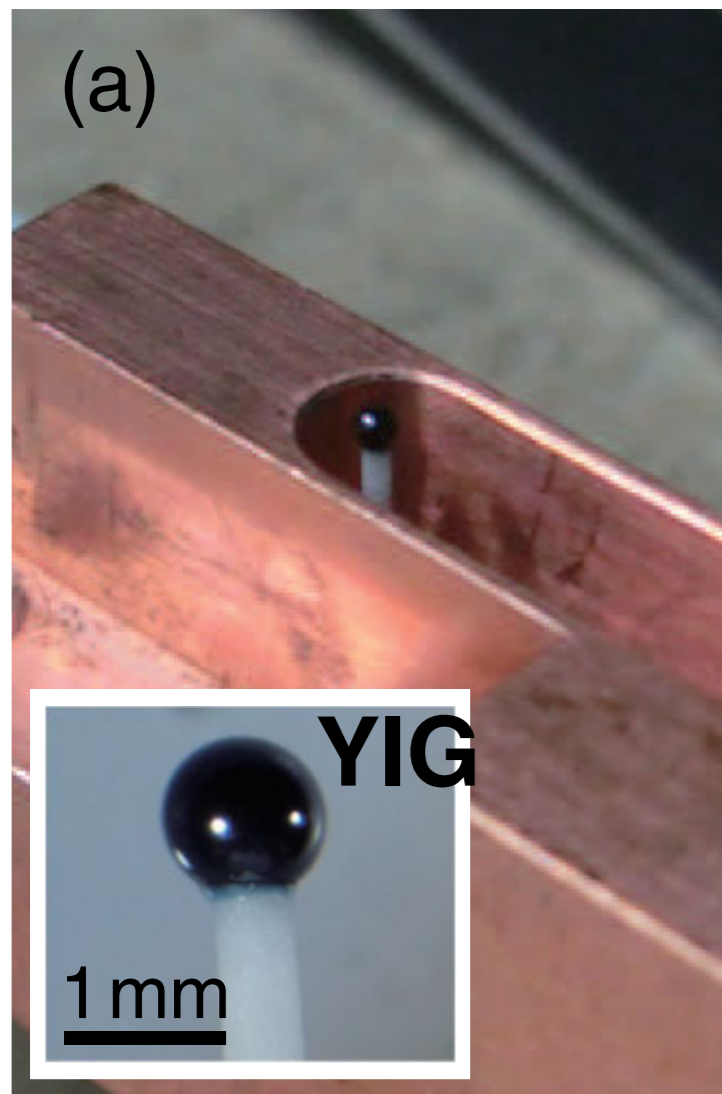
# Microwave Regime

Magnons



Microwaves

Strong coupling regime



$\Omega \sim \text{GHz}$   
for 30mT

- Huebl et. al, PRL 111, 127003 (2013)
- Zhang et. al PRL 113, 156401 (2014)
- Tabuchi et. al PRL 113, 083603 (2014)

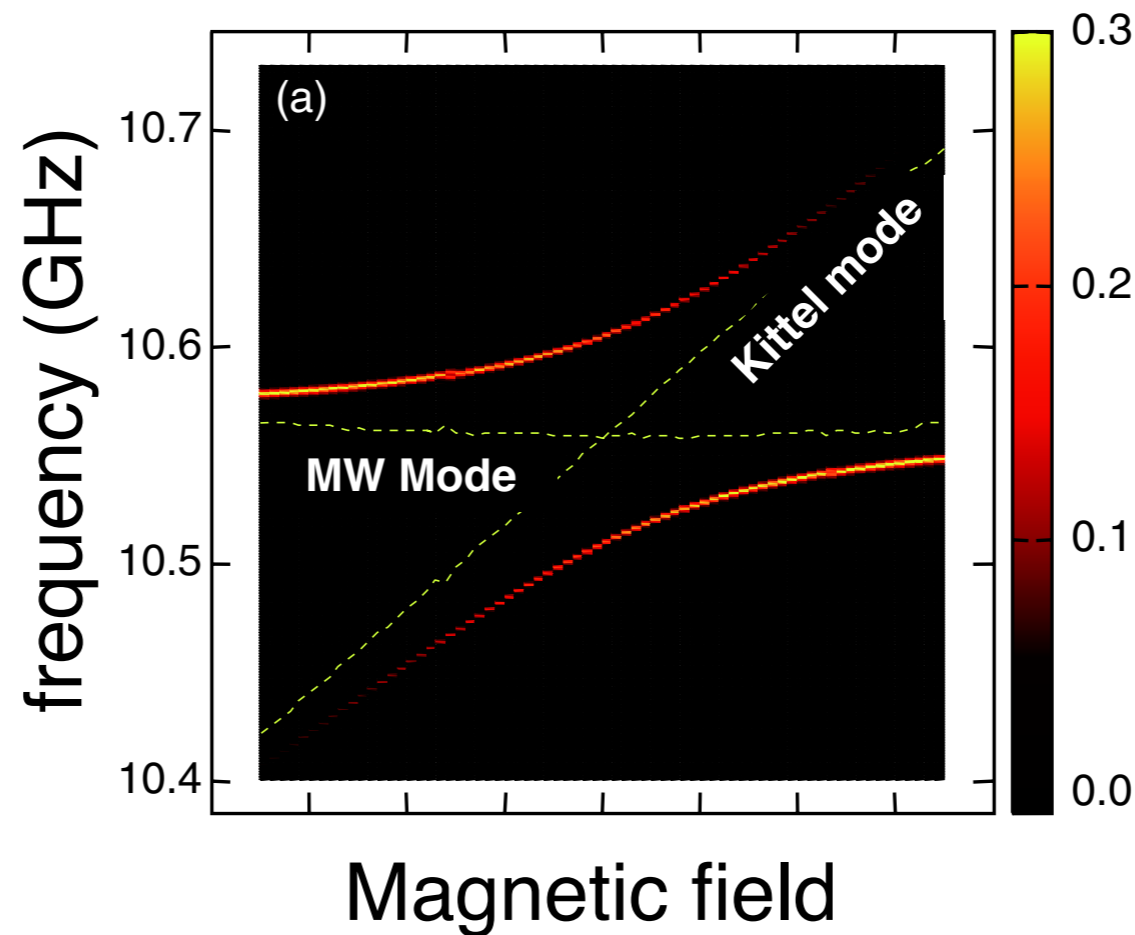
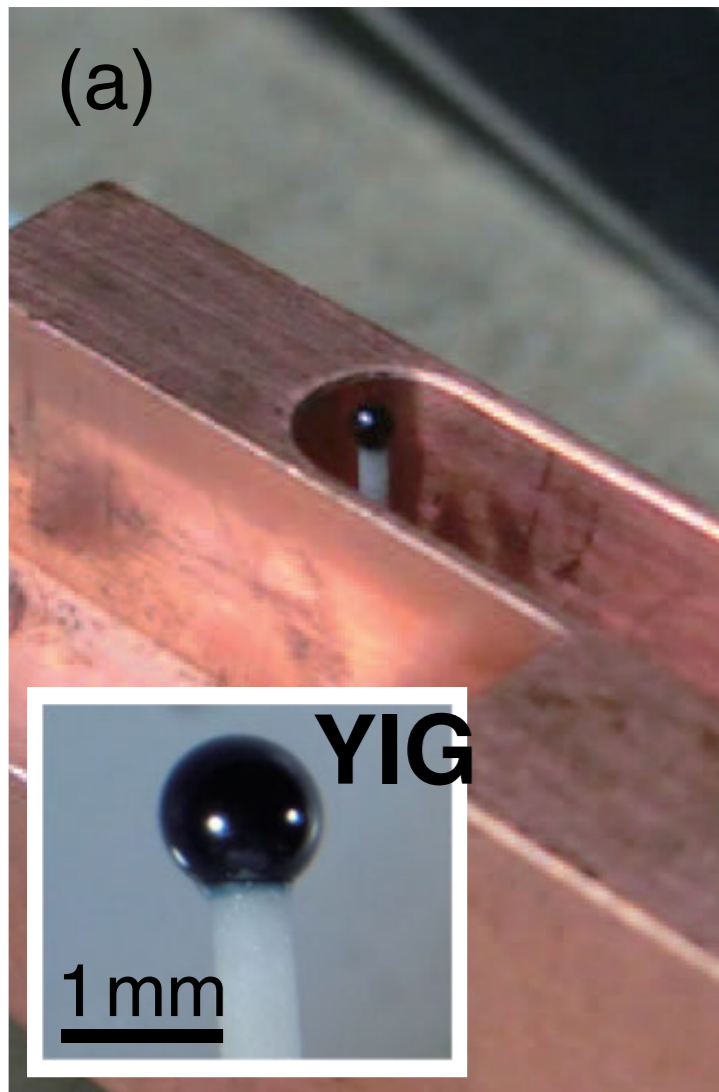
# Microwave Regime

Magnons



Microwaves

Strong coupling regime

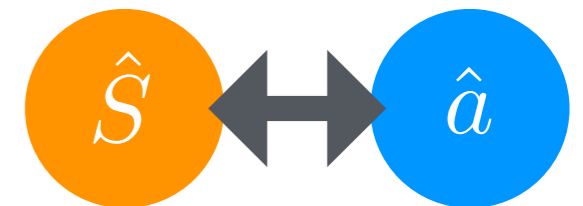


Resonant coupling

$$-\frac{\mu_0}{2} \int d^3r \mathbf{M} \cdot \mathbf{H}_{\text{MW}}$$



$$\hat{S}^- \hat{a} + \hat{S}^+ \hat{a}^\dagger$$



- Huebl et. al, PRL 111, 127003 (2013)
- Zhang et. al PRL 113, 156401 (2014)
- Tabuchi et. al PRL 113, 083603 (2014)

Soykal and M. E. Flatte  
PRL 104, 077202 (2010)



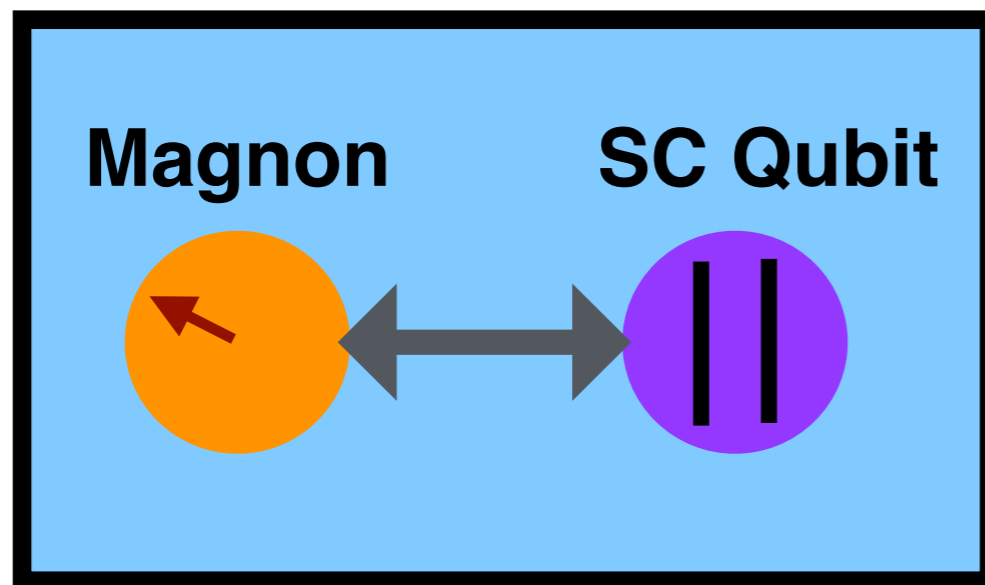
# Microwave Regime

QUANTUM INFORMATION

(Science 2015)

## Coherent coupling between a ferromagnetic magnon and a superconducting qubit

Yutaka Tabuchi,<sup>1\*</sup> Seiichiro Ishino,<sup>1</sup> Atsushi Noguchi,<sup>1</sup> Toyofumi Ishikawa,<sup>1</sup> Rekishu Yamazaki,<sup>1</sup> Koji Usami,<sup>1</sup> Yasunobu Nakamura<sup>1,2</sup>



MW Cavity

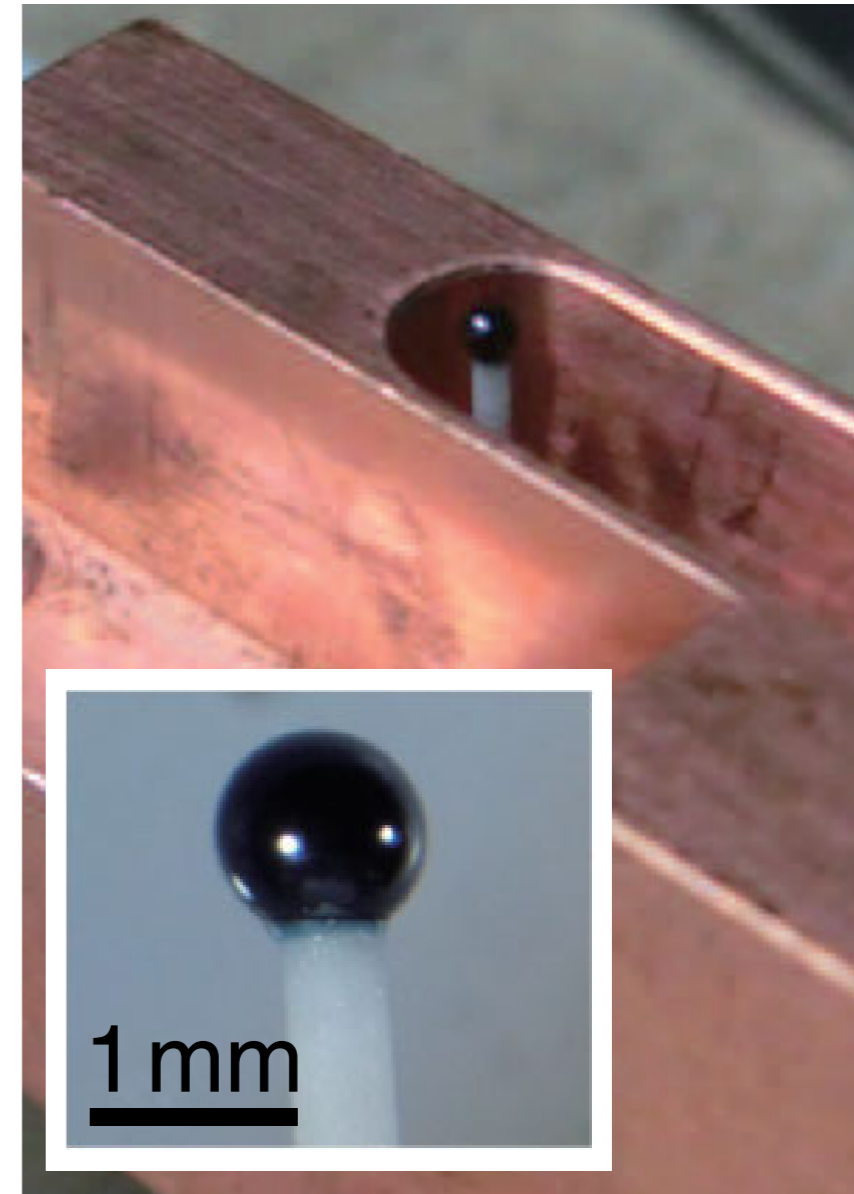
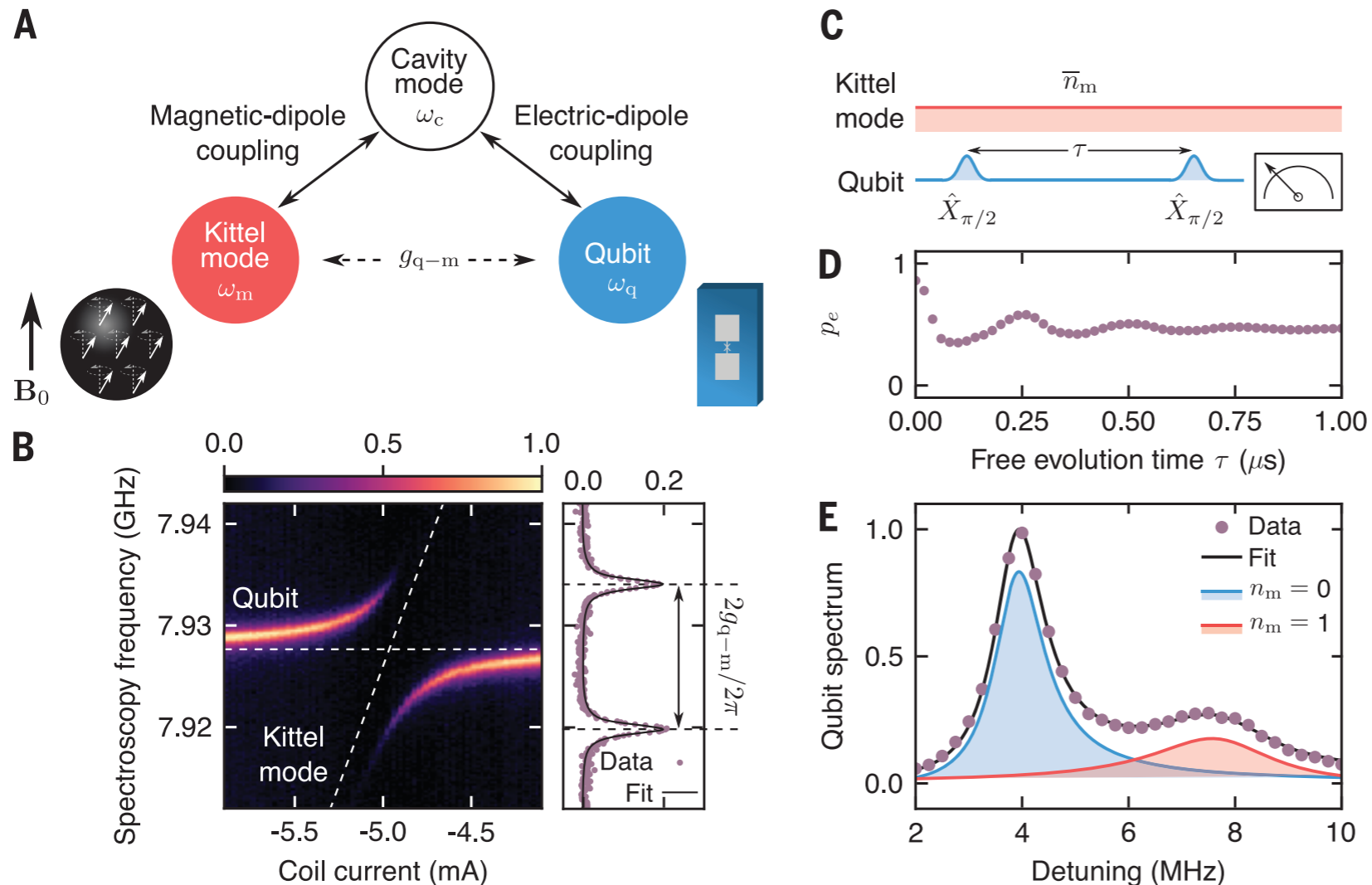


# Microwave Regime: Single Magnon Detector

QUANTUM SENSING (Science 2020)

## Entanglement-based single-shot detection of a single magnon with a superconducting qubit

Dany Lachance-Quirion<sup>1</sup>, Samuel Piotr Wolski<sup>1</sup>, Yutaka Tabuchi<sup>1</sup>, Shingo Kono<sup>1</sup>, Koji Usami<sup>1</sup>, Yasunobu Nakamura<sup>1,2\*</sup>





Microwave  
regime

# Spin cat states in ferromagnetic insulators



S. Sharma, V. Bittencourt, A. Karenowska, SVK;  
PRB 103, L100403 (2021)



MAX PLANCK INSTITUTE  
FOR THE SCIENCE OF LIGHT



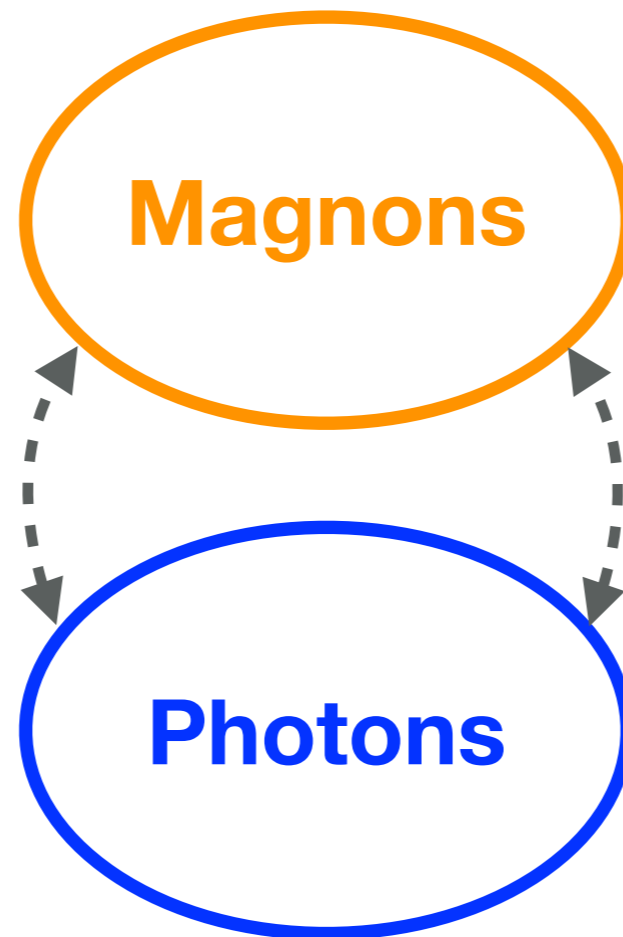
FRIEDRICH-ALEXANDER  
UNIVERSITÄT  
ERLANGEN-NÜRNBERG



UNIVERSITY OF  
OXFORD

# Proposal: Heralding Magnetic Cat States

**Heralding**



**Generation of entangled state**

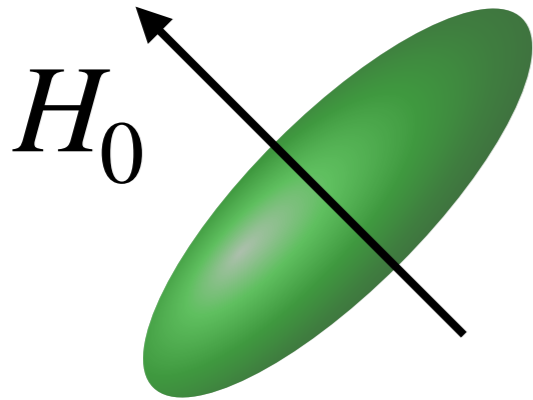
**+**

**Measurement**



# Heralding Magnetic Cat States

Anisotropic Nanomagnet



$$\mathcal{H}_{\text{mag}} = \frac{\mu_0}{2} \mathbf{M} \tilde{\mathbf{N}} \mathbf{M} - \mu_0 M_z H_0$$

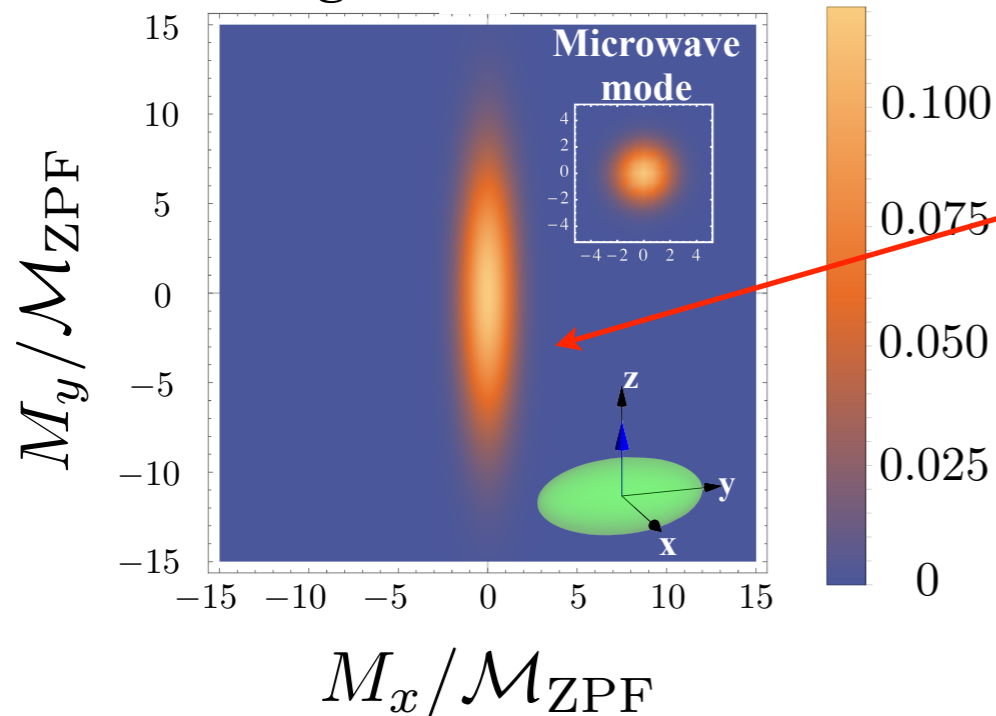


isotropic  
magnet

anisotropy

$$\frac{\hat{H}}{\hbar} = \omega_0 \hat{s}^\dagger \hat{s} + \frac{\omega_s}{2} (\hat{s}^2 + \hat{s}^{\dagger 2})$$

Spin fluctuations  
ground state



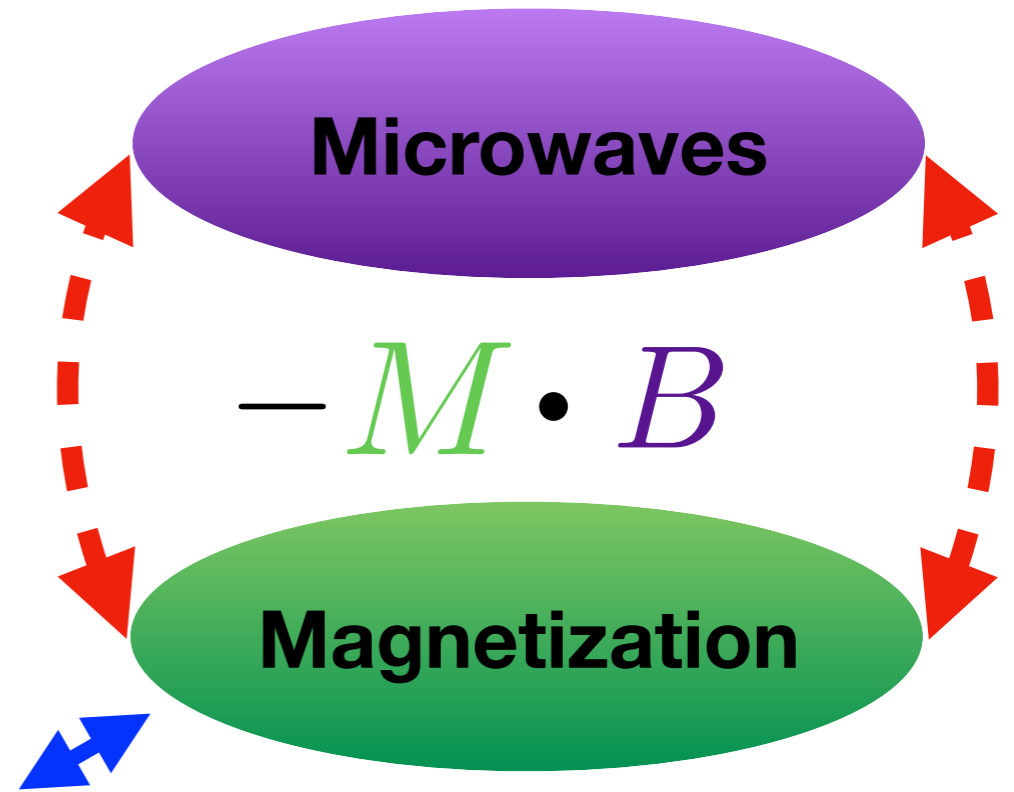
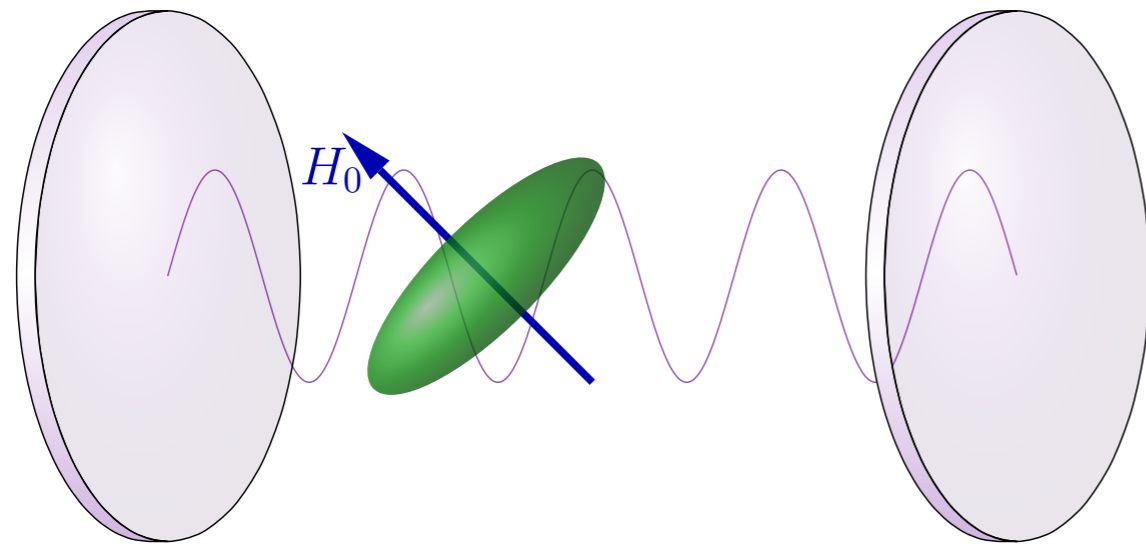
**Squeezed magnetic ground state  
due to anisotropy**

Kamra et al, PRL 2016

Anisotropic quantum fluctuations

# Heralding Magnetic Cat States

MW Cavity + Anisotropic Nanomagnet



Shape anisotropy

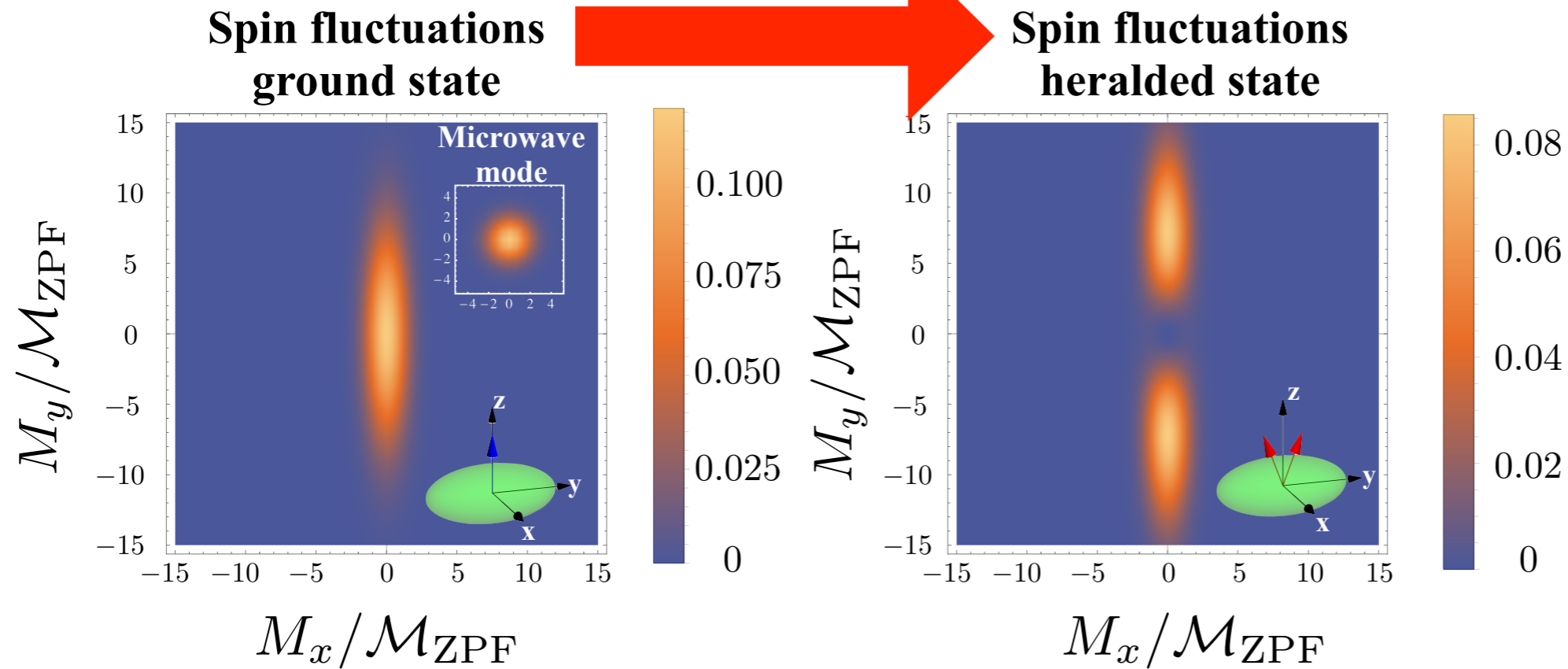
$$\frac{\hat{H}}{\hbar} = \omega_0 \hat{s}^\dagger \hat{s} + \frac{\omega_s}{2} (\hat{s}^2 + \hat{s}^{\dagger 2}) + \omega_a \hat{a}^\dagger \hat{a} + g(\hat{s} \hat{a}^\dagger + \hat{s}^\dagger \hat{a})$$

cavity photons
cavity-magnet interaction

**The ground state of the coupled system is an entangled state of photons and flipped spins**

# Heralding Magnetic Cat States

measurement of the state of the cavity



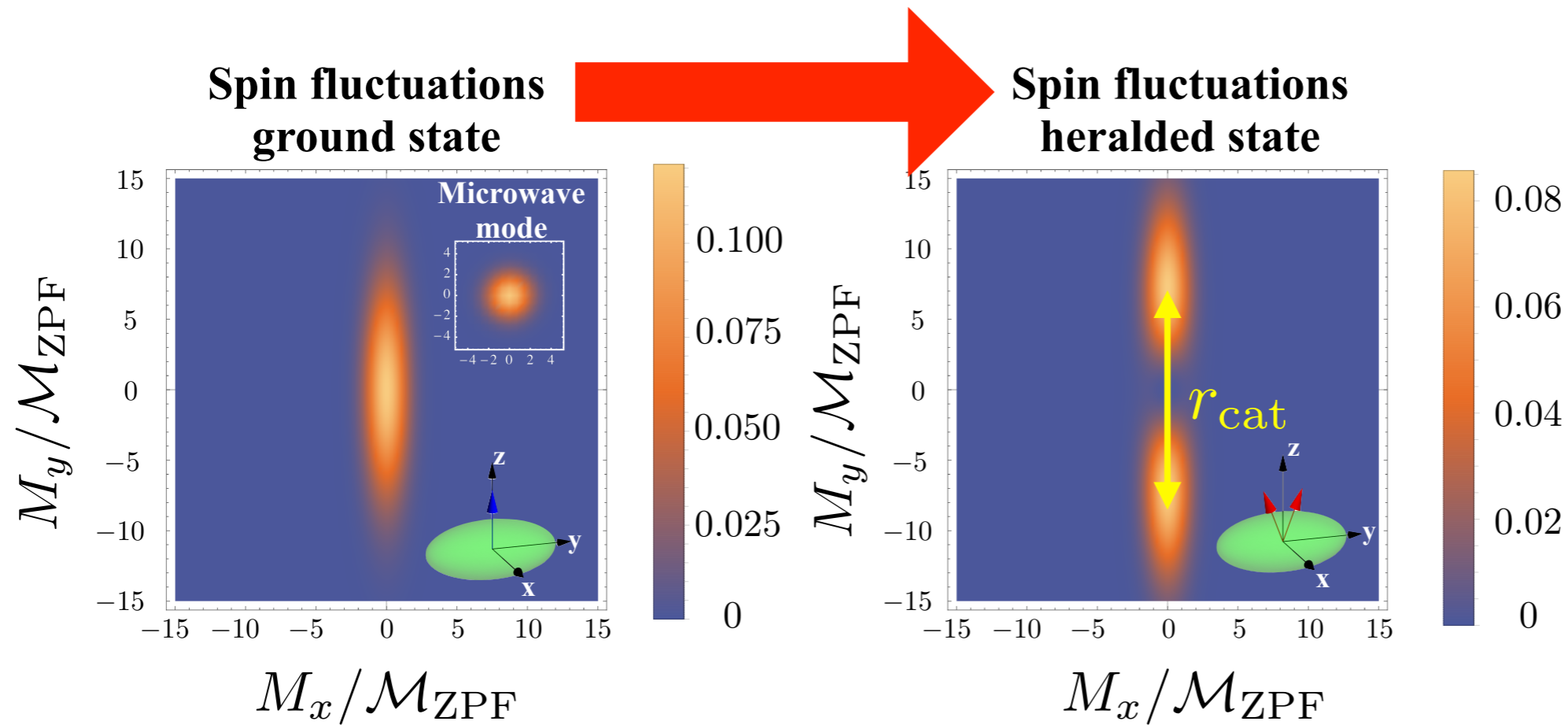
squeezed  
magnetic  
cat state!



The magnetization points simultaneously in two directions

# Creating Magnetic Cat States

measurement of the state of the cavity



squeezed  
magnetic  
cat state!



The magnetization points simultaneously in two directions

Experimental  
limitation

Low temperatures needed for a "large cat"

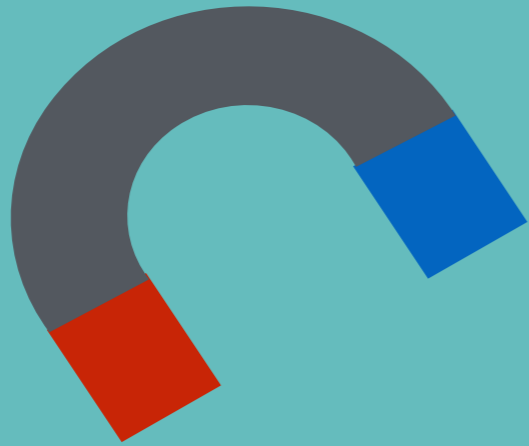
$$T < 5 \text{ mK}$$





Microwave  
regime

# Magnon-phonon quantum thermometry



C. Potts, V. Bittencourt, SVK, John P. Davis;  
Phys. Rev. Applied 13, 064001 (2020)



MAX PLANCK INSTITUTE  
FOR THE SCIENCE OF LIGHT

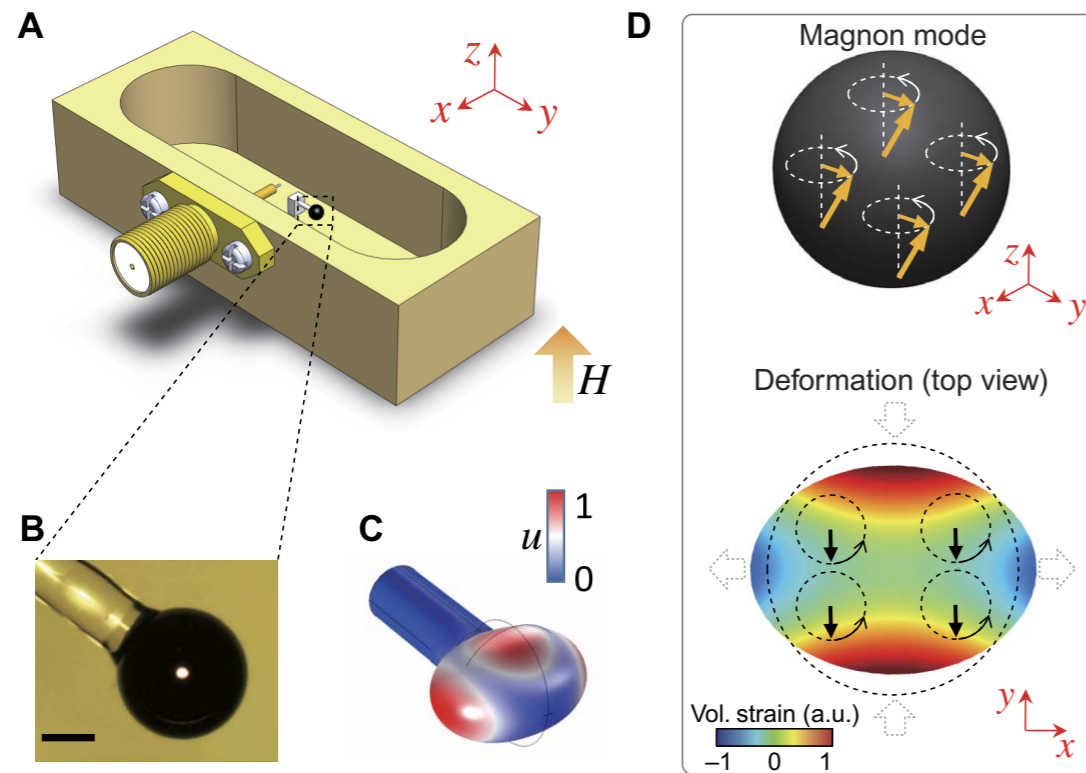
FAU  
FRIEDRICH-ALEXANDER  
UNIVERSITÄT  
ERLANGEN-NÜRNBERG



UNIVERSITY OF  
ALBERTA

# Magnons can couple coherently to phonons

## Cavity magnomechanics



Zhang et. al Science Advances 2016

## Magnetoelastic energy

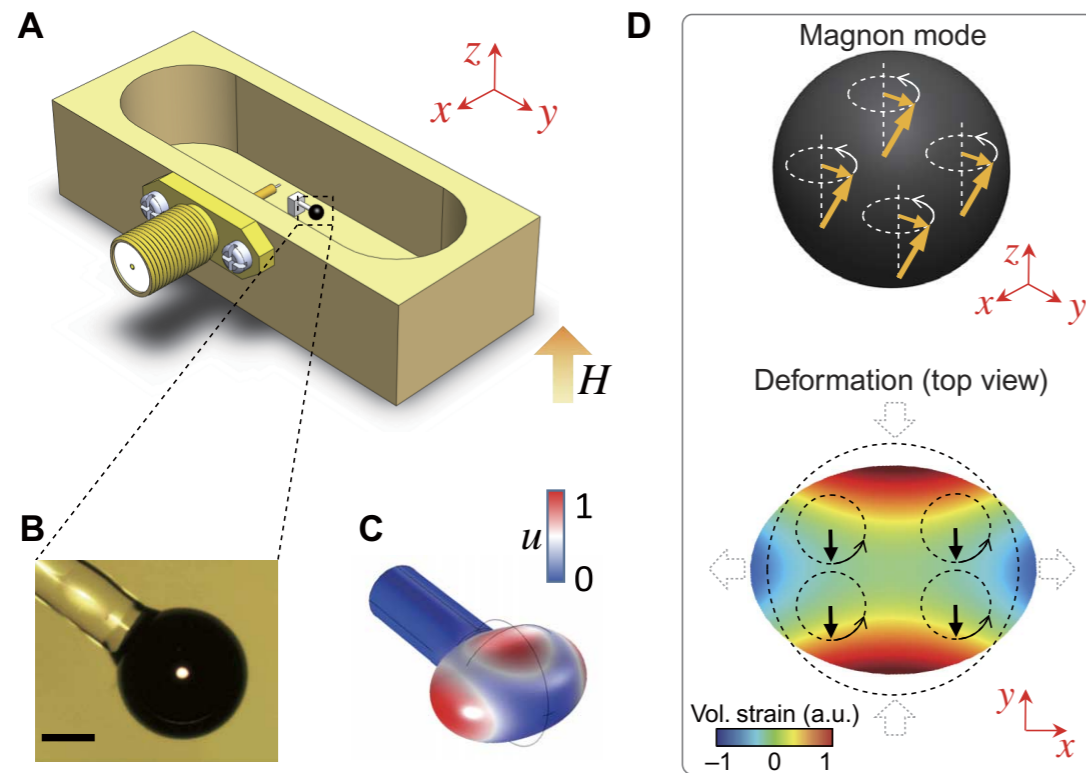
magnetoelastic constants

strain tensor

$$E_{mel} = \frac{B_1}{M_S^2} \int d^3r [M_x^2 u_{xx} + M_y^2 u_{yy} + M_z^2 u_{zz}] + \frac{2B_2}{M_S^2} \int d^3r [M_x M_y u_{xy} + M_y M_z u_{yz} + M_z M_x u_{zx}]$$

# Magnons can couple coherently to phonons

## Cavity magnomechanics



Zhang et. al Science Advances 2016

## Magnetoelastic energy

magnetoelastic constants / strain tensor

$$E_{mel} = \frac{B_1}{M_S^2} \int d^3r [M_x^2 u_{xx} + M_y^2 u_{yy} + M_z^2 u_{zz}] + \frac{2B_2}{M_S^2} \int d^3r [M_x M_y u_{xy} + M_y M_z u_{yz} + M_z M_x u_{zx}]$$

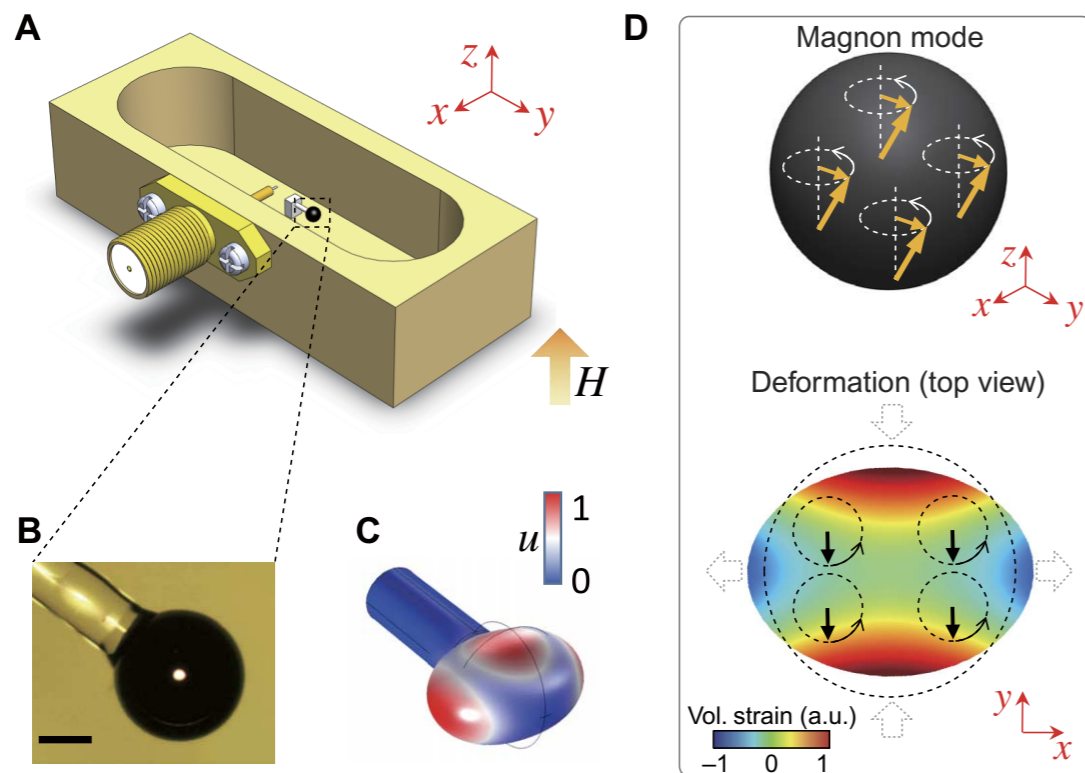
Quantizing in terms of magnon  $\mathbf{m}$  and phonon  $\mathbf{b}$  operators for an arbitrary magnetic texture

$$\hat{H}_{ME} = \sum_{\alpha\beta\gamma ij} \left[ G_{\alpha\gamma}^{ij} \hat{m}_\alpha (\hat{b}_\gamma + \hat{b}_\gamma^\dagger) + T_{\alpha\beta\gamma}^{ij} \hat{m}_\alpha \hat{m}_\beta (\hat{b}_\gamma + \hat{b}_\gamma^\dagger) + P_{\alpha\beta\gamma}^{ij} \hat{m}_\beta^\dagger \hat{m}_\alpha (\hat{b}_\gamma + \hat{b}_\gamma^\dagger) + h.c. \right]$$

F. Engelhardt et al, in preparation

# Magnons can couple to phonons

## Cavity magnomechanics



Zhang et. al Science Advances 2016

## Magnetoelastic energy

magnetoelastic constants / strain tensor

$$E_{mel} = \frac{B_1}{M_S^2} \int d^3r [M_x^2 u_{xx} + M_y^2 u_{yy} + M_z^2 u_{zz}] + \frac{2B_2}{M_S^2} \int d^3r [M_x M_y u_{xy} + M_y M_z u_{yz} + M_z M_x u_{zx}]$$

Quantizing in terms of magnon  $\mathbf{m}$  and phonon  $\mathbf{b}$  operators for an arbitrary magnetic texture

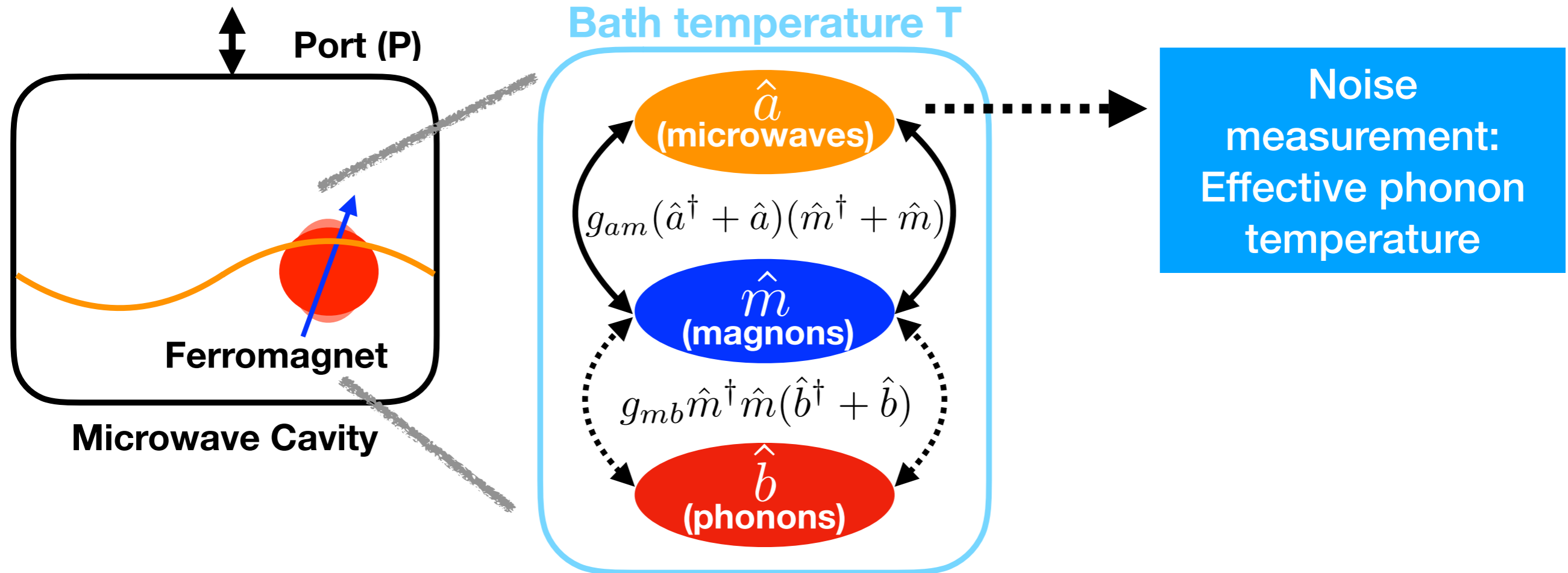
$$\hat{H}_{ME} = \sum_{\alpha\beta\gamma ij} [G_{\alpha\gamma}^{ij} \hat{m}_\alpha (\hat{b}_\gamma + \hat{b}_\gamma^\dagger) + T_{\alpha\beta\gamma}^{ij} \hat{m}_\alpha \hat{m}_\beta (\hat{b}_\gamma + \hat{b}_\gamma^\dagger) + P_{\alpha\beta\gamma}^{ij} \hat{m}_\beta^\dagger \hat{m}_\alpha (\hat{b}_\gamma + \hat{b}_\gamma^\dagger) + h.c.]$$

Parametric coupling

Magnons ~ GHz  
Phonons ~ kHz

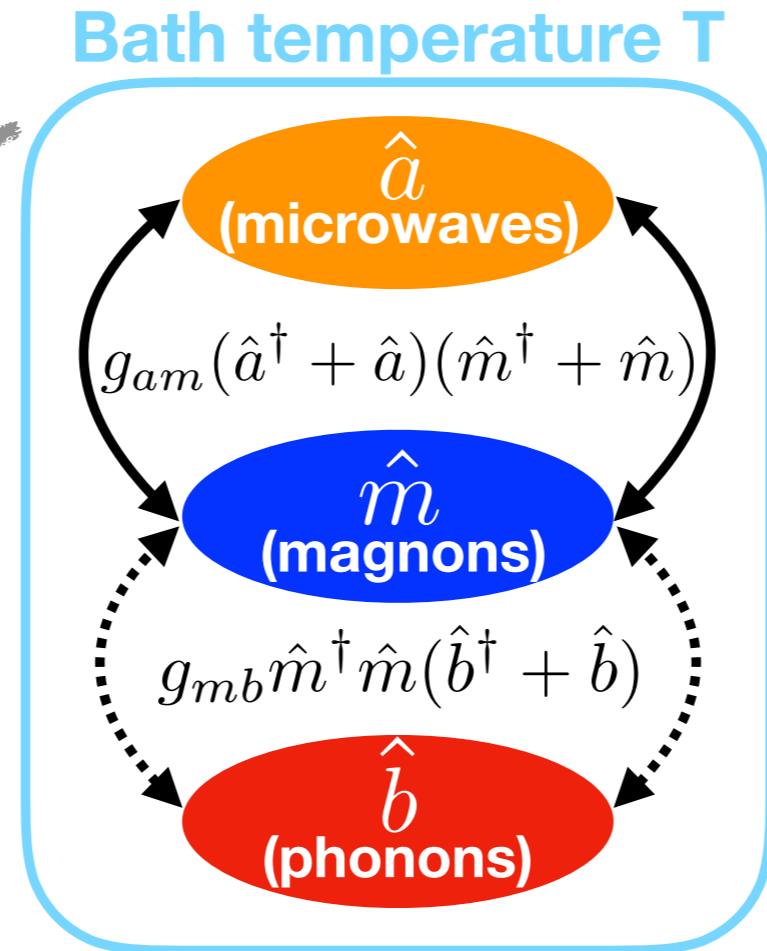
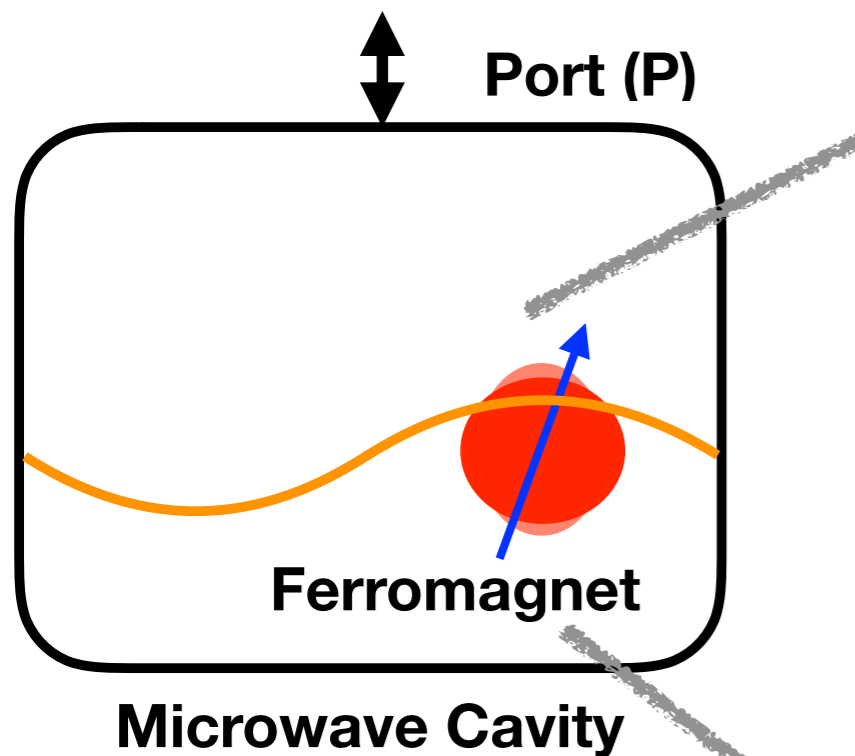


# Magnon - Phonon Quantum Correlation Thermometry



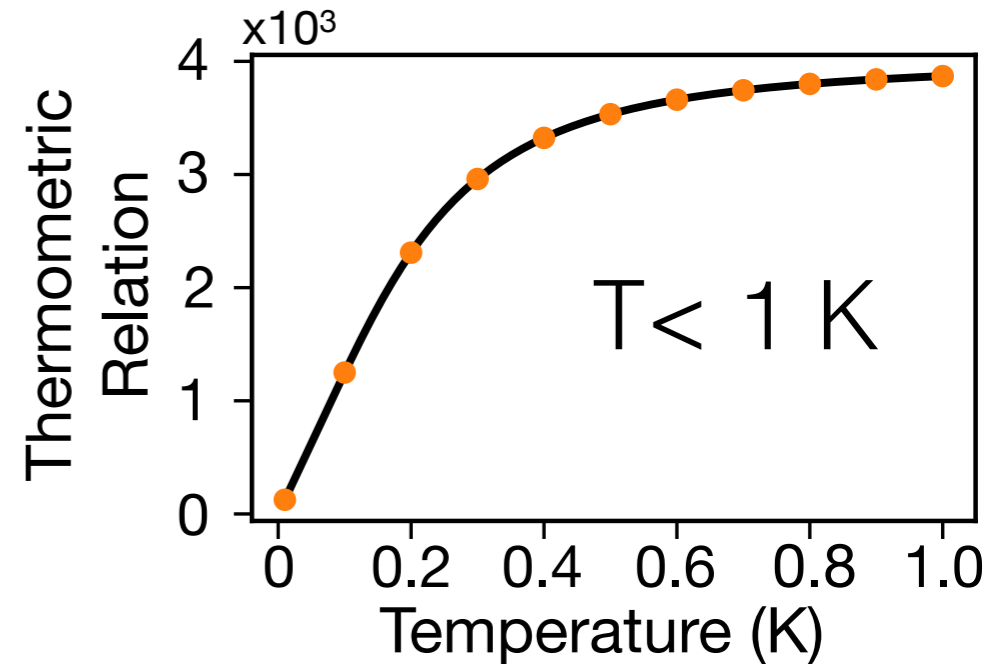
- Thermal vibrations imprinted on the magnetization
- Drive the MW at resonance and measure noise correlations

# Magnon - Phonon Quantum Correlation Thermometry



Noise measurement:  
Effective phonon temperature

Primary thermometer



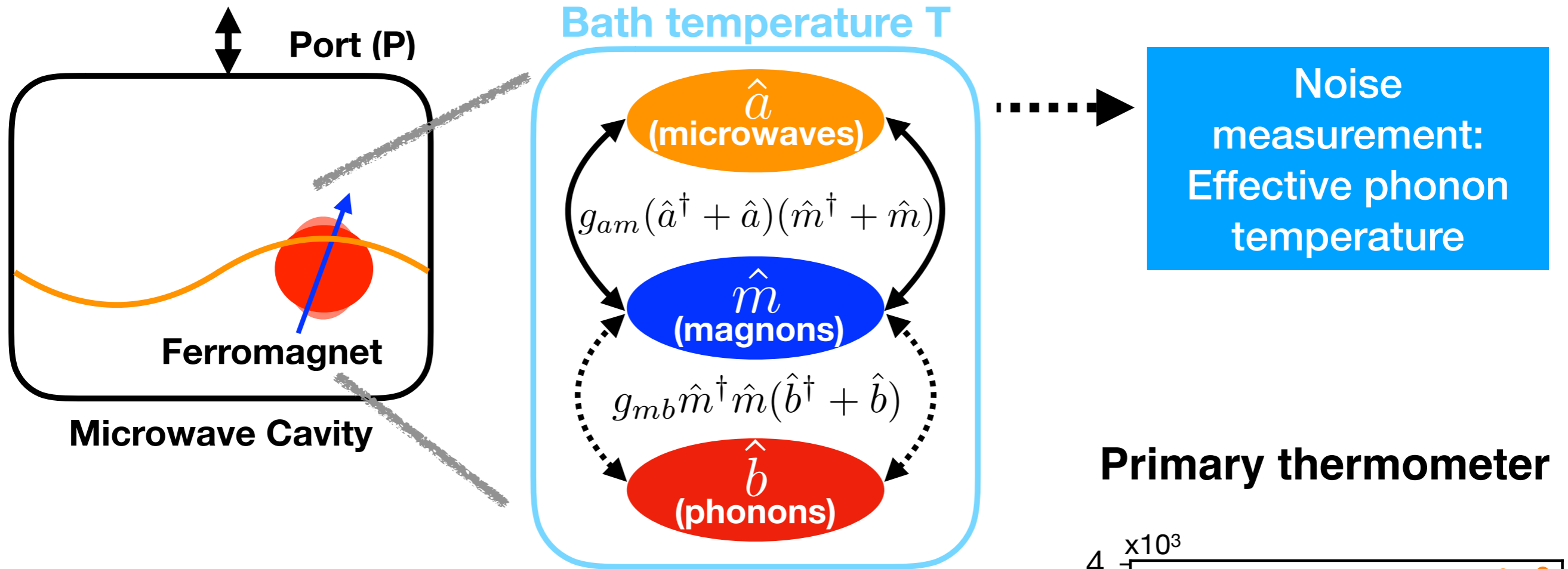
MW output noise correlation

$$\frac{\text{Re}\{S_{\frac{\pi}{2}, \frac{\pi}{2}}[\omega]\}}{\text{Im}\{S_{0, \frac{\pi}{2}}[\omega]\}} = \frac{4 \coth\left(\frac{\hbar\omega}{2k_B T}\right)}{2n_{\text{th}} + 1}$$

Obtained by solving the quantum Langevin equations (input-output theory)

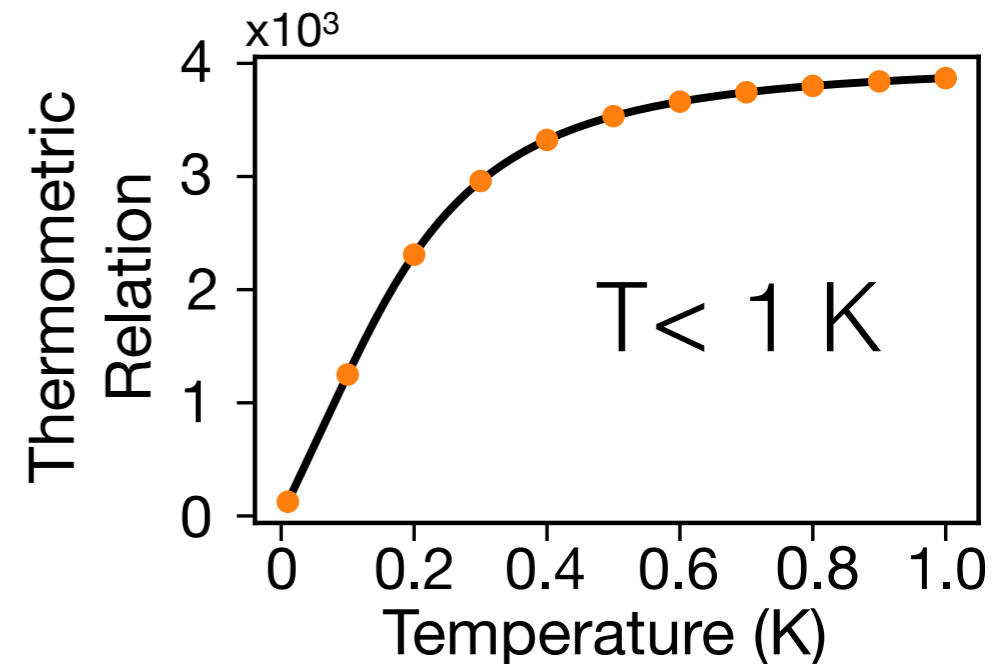
Measurement: two quadratures

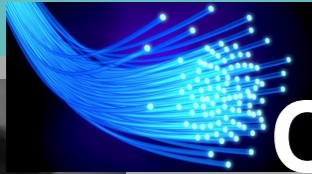
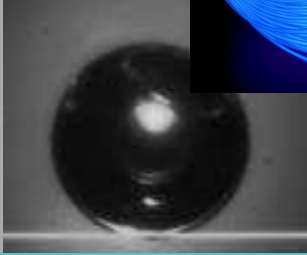
# Magnon - Phonon Quantum Correlation Thermometry



- No-calibration thermometer for cryogenic temperatures with low heating effects
- Compatible with current experimental setups
- Experimental imperfections sources: Finite detuning and deviations from the optimal driving condition

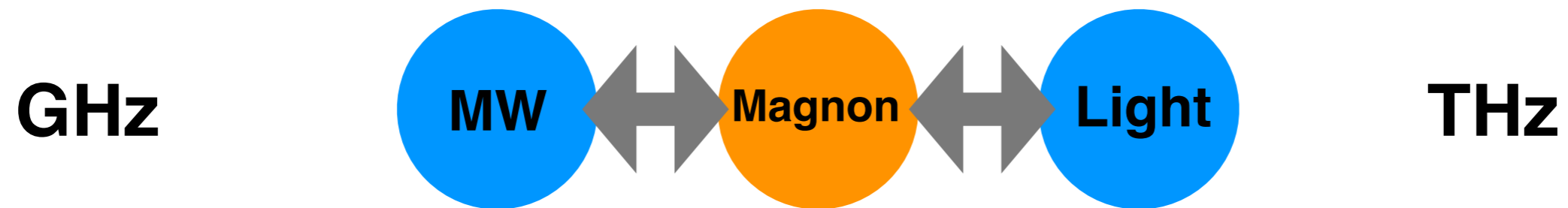
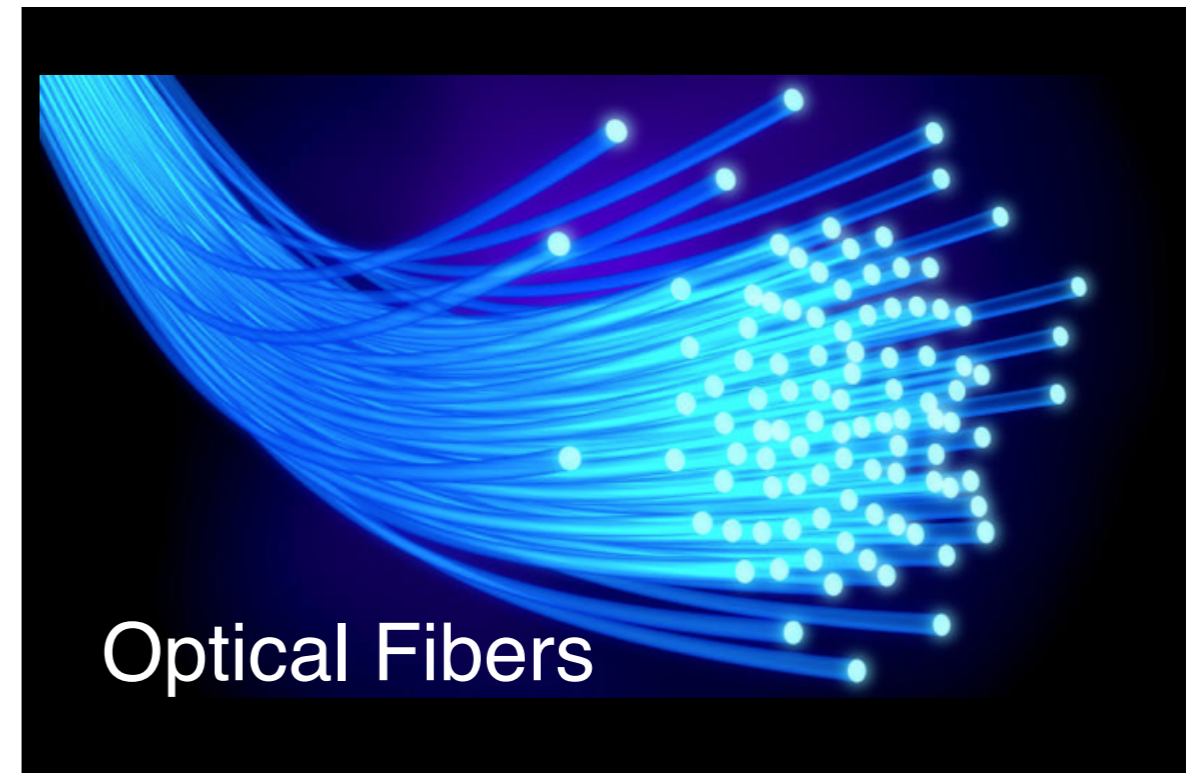
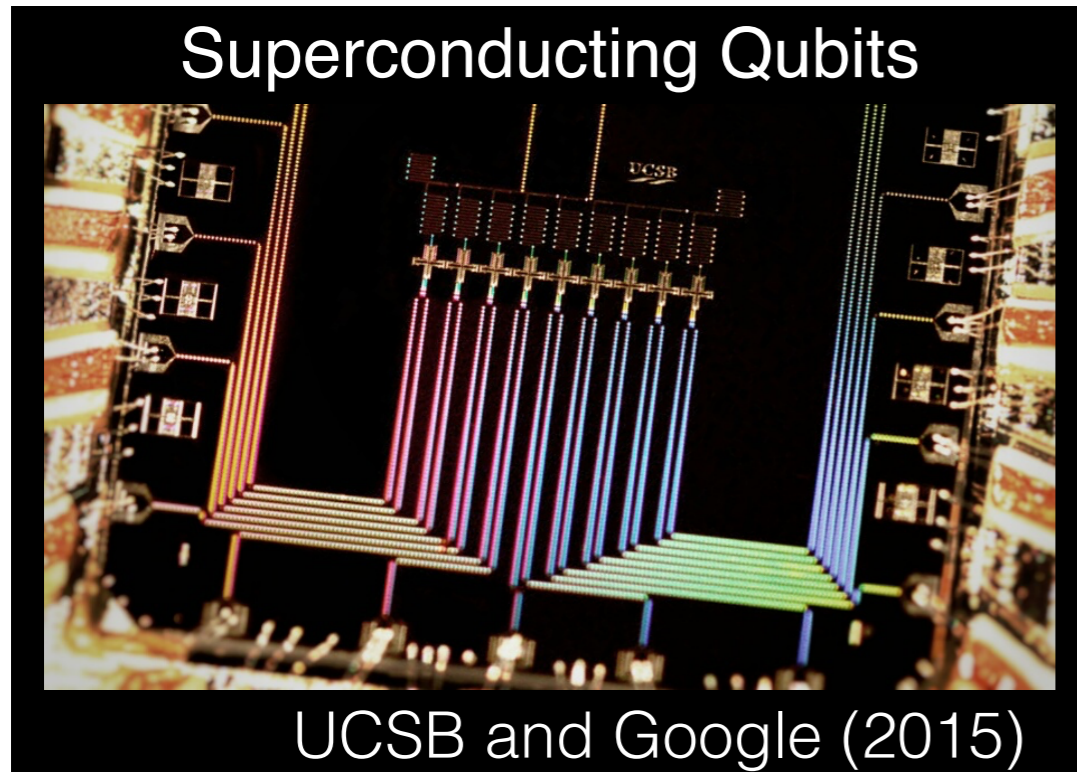
## Primary thermometer





# Optical regime

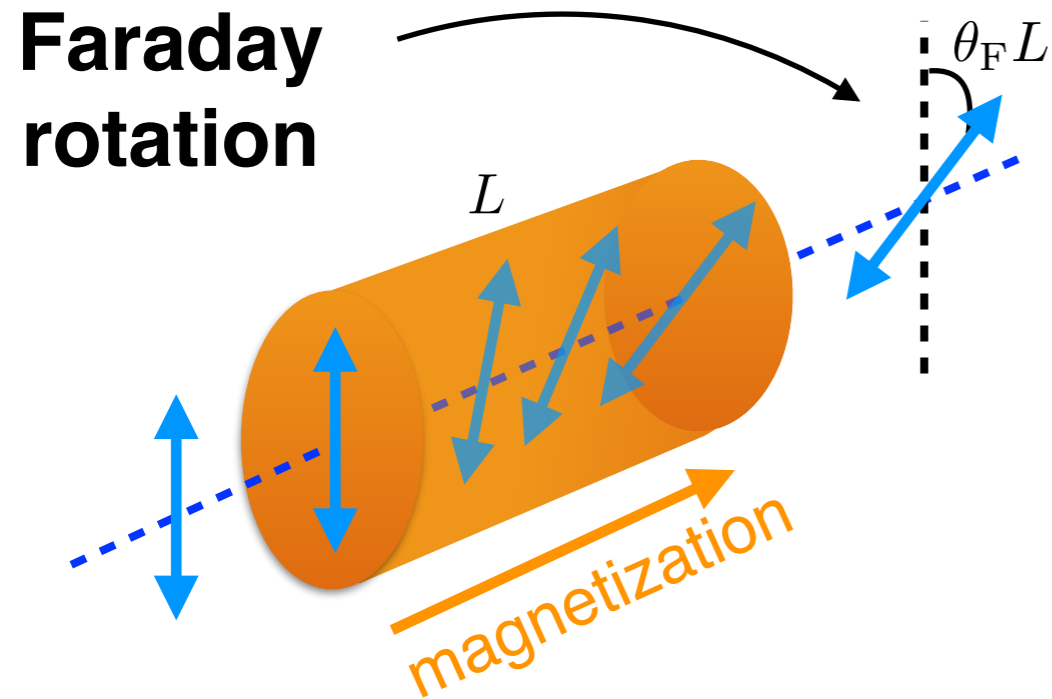
# Coupling to Optics?



**Motivation:**  
**magnon as a transducer**



# Coupling to Optics?: Faraday Effect



permittivity

$$\epsilon_{ij}(\mathbf{M}) = \epsilon_0 (\epsilon \delta_{ij} - i f \epsilon_{ijk} M_k)$$

broken time-reversal symmetry

optical  
spin density

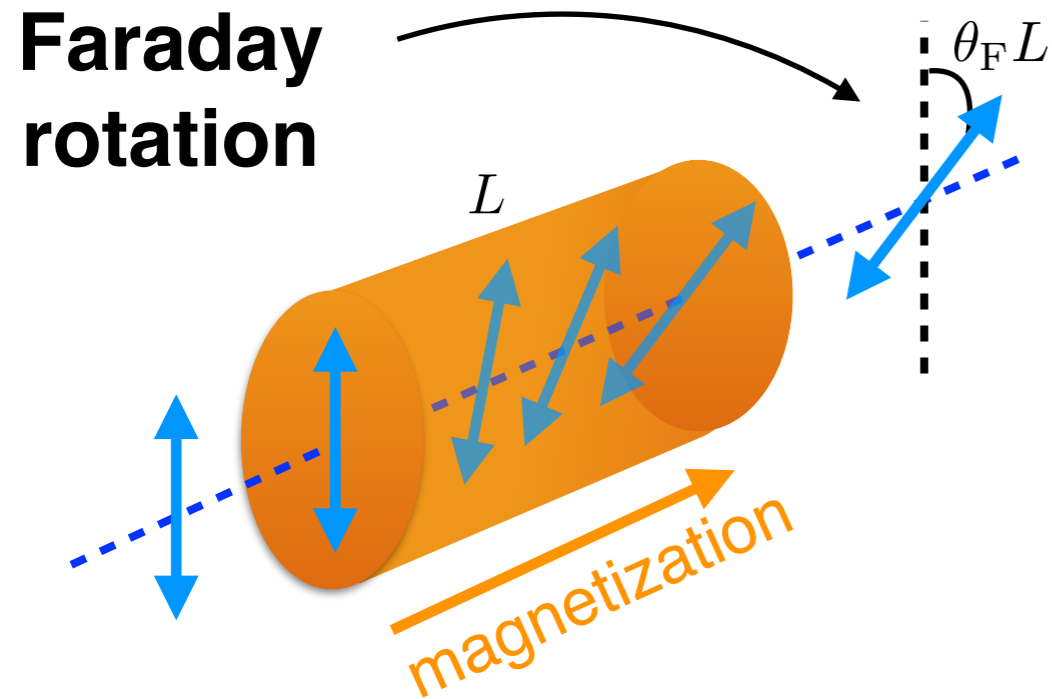


$$\bar{U}_{\text{MO}} = \theta_F \sqrt{\frac{\epsilon}{\epsilon_0}} \int d\mathbf{r} \frac{\mathbf{M}(\mathbf{r})}{M_s} \cdot \frac{\epsilon_0}{2i\omega} [\mathbf{E}^*(\mathbf{r}) \times \mathbf{E}(\mathbf{r})]$$



magnetization  
density

# Coupling to Optics?: Faraday Effect



permittivity

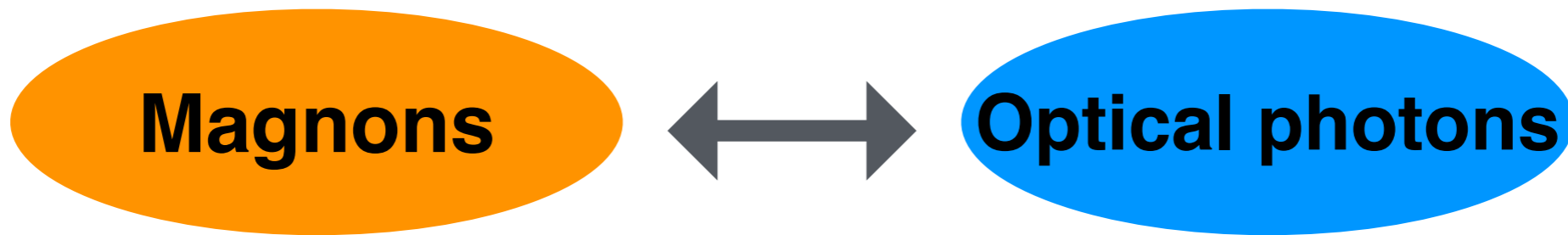
$$\epsilon_{ij}(\mathbf{M}) = \epsilon_0 (\epsilon \delta_{ij} - i f \epsilon_{ijk} M_k)$$

↑  
broken time-reversal symmetry

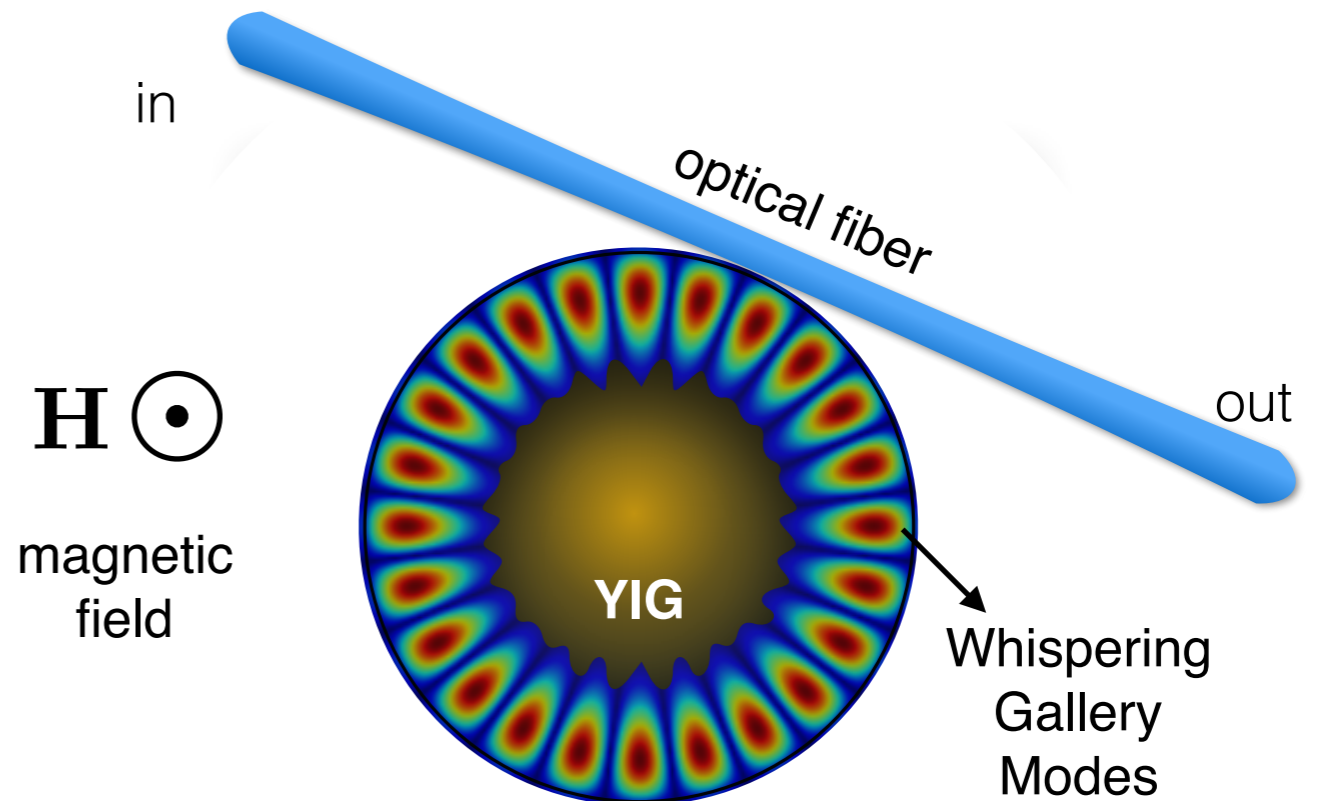
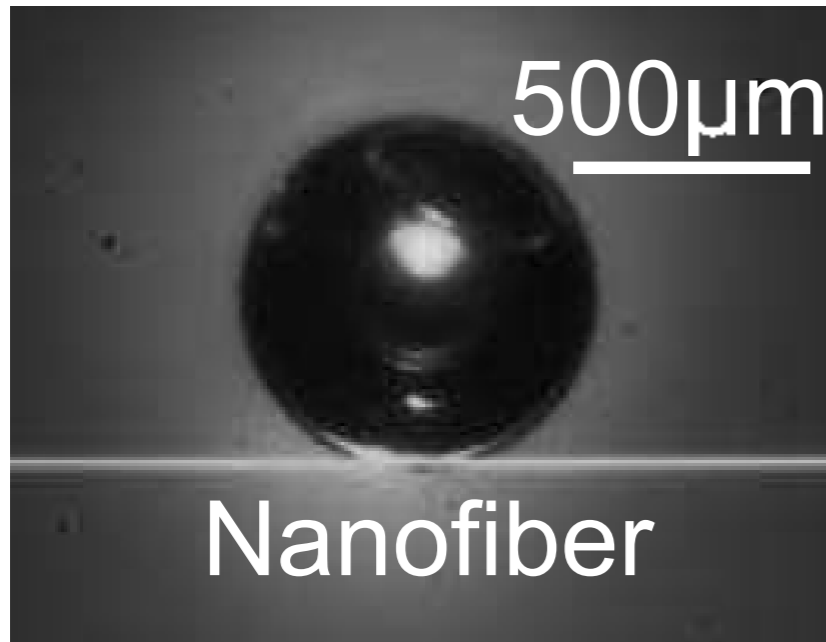
$$\bar{U}_{\text{MO}} = \theta_F \sqrt{\frac{\epsilon}{\epsilon_0}} \int d\mathbf{r} \frac{\mathbf{M}(\mathbf{r})}{M_s} \cdot \frac{\epsilon_0}{2i\omega} [\mathbf{E}^*(\mathbf{r}) \times \mathbf{E}(\mathbf{r})]$$



# Cavity Optomagnonics



Coupling demonstrated in 2016



- Osada et. al PRL 116, 223601 (Nakamura's group, Tokyo)
- Haigh et. al PRL 117, 133602 (Cambridge Univ / Hitachi)
- Zhang et. al PRL 117, 123605 (Hong Tang's group, Yale)

**Cavity Enhanced Coupling**

# But...

## Problem

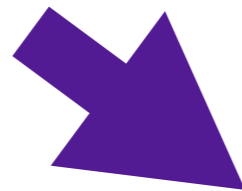
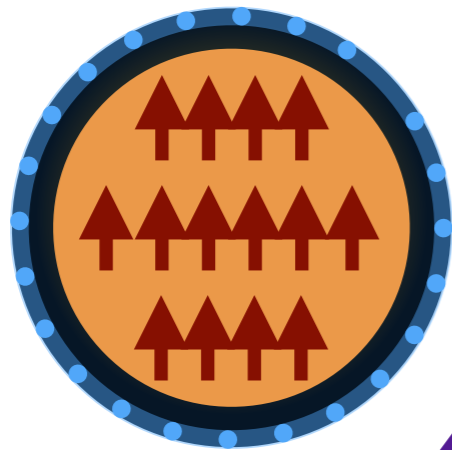
the state of the art optomagnonic coupling is small

**We have shown that the theoretical limit  
is much larger**

SVK, H. X. Tang, and F. Marquardt  
PRA 94, 033821 (2016)

## What we need:

better overlap of modes



?

smaller systems

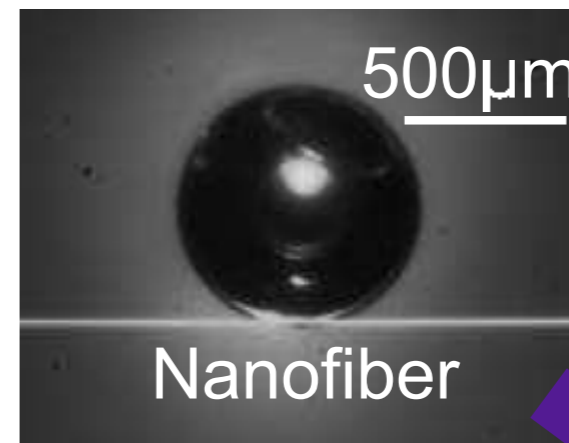
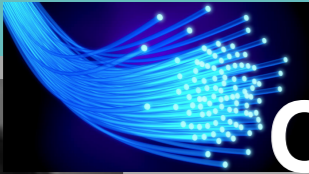
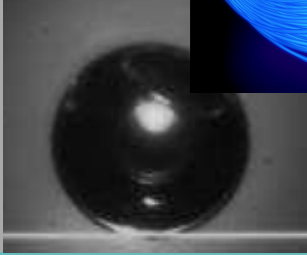


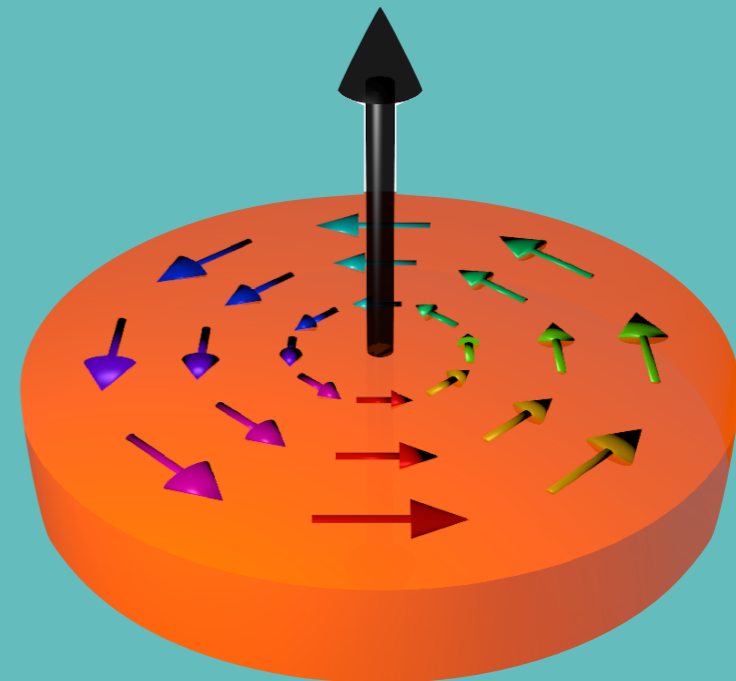
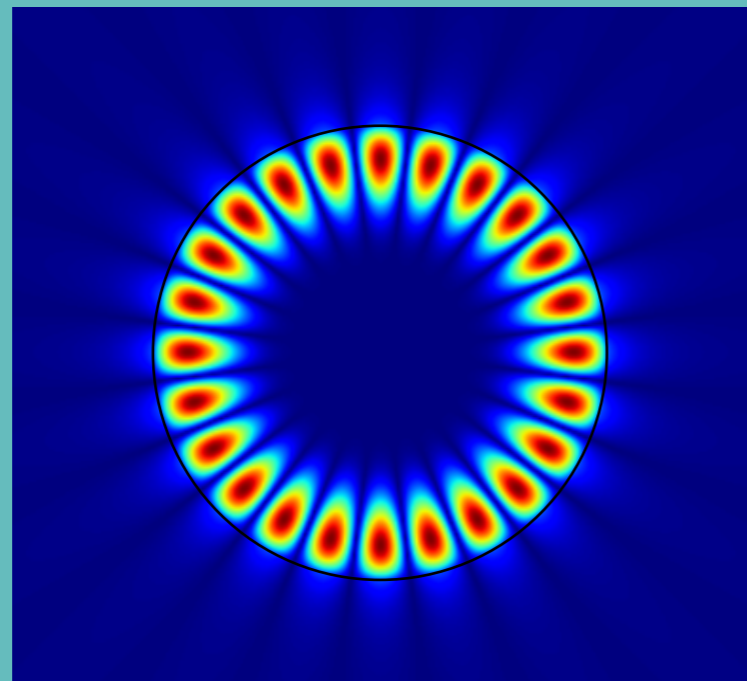
Fig: Osada et. al.  
PRL 116, 223601





Optical  
regime

# Optomagnonics with Magnetic textures



J. Graf, H. Pfeifer, F. Marquardt, SVK;  
PRB 98, 241406(R) (2018)



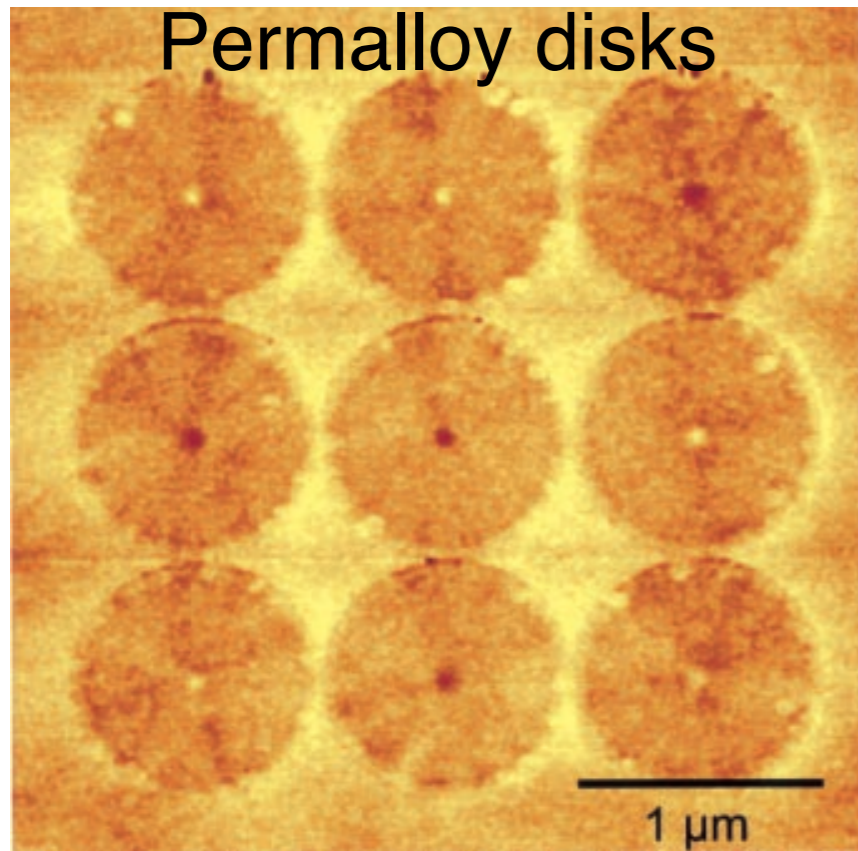
MAX PLANCK INSTITUTE  
FOR THE SCIENCE OF LIGHT



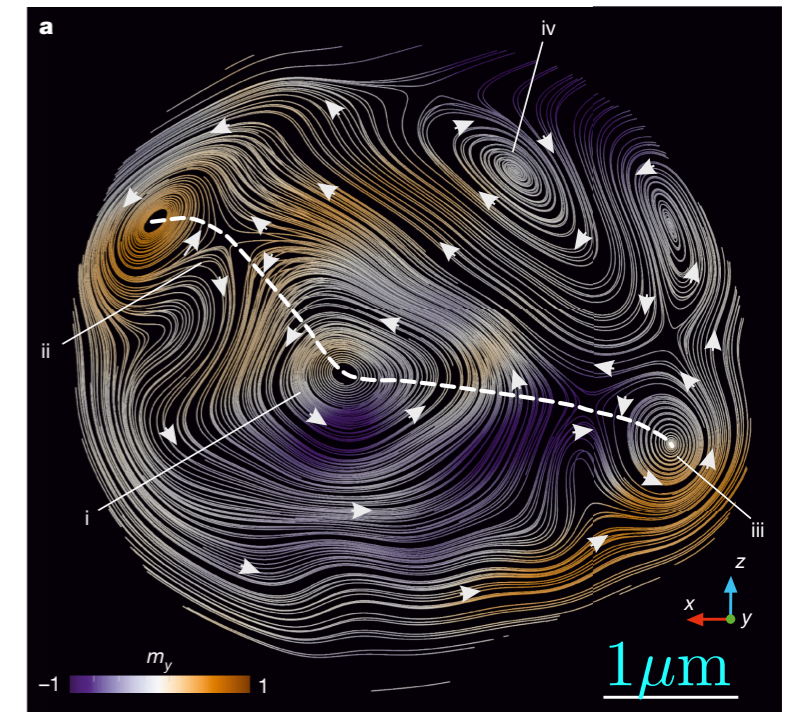
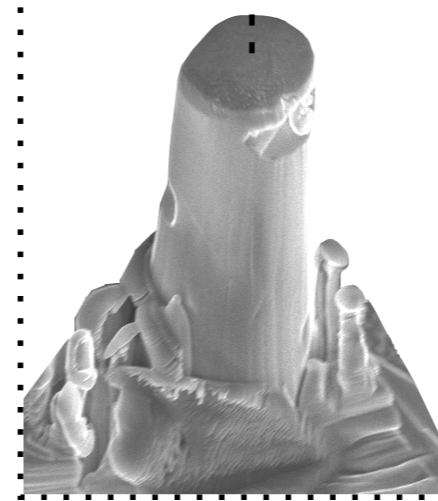
FRIEDRICH-ALEXANDER  
UNIVERSITÄT  
ERLANGEN-NÜRNBERG



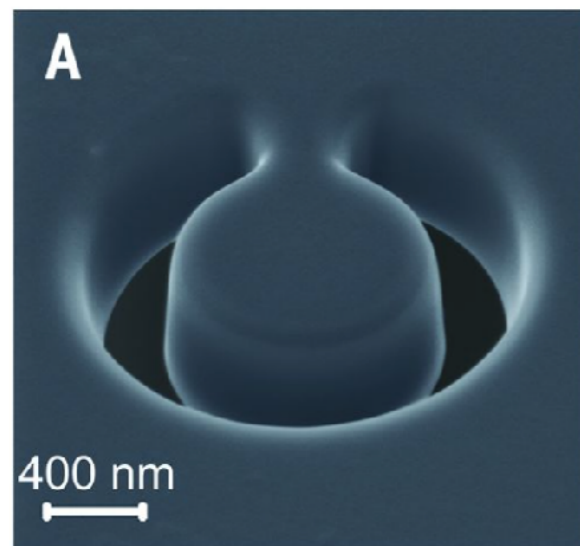
# Magnetic Textures: Vortex in Microdisks



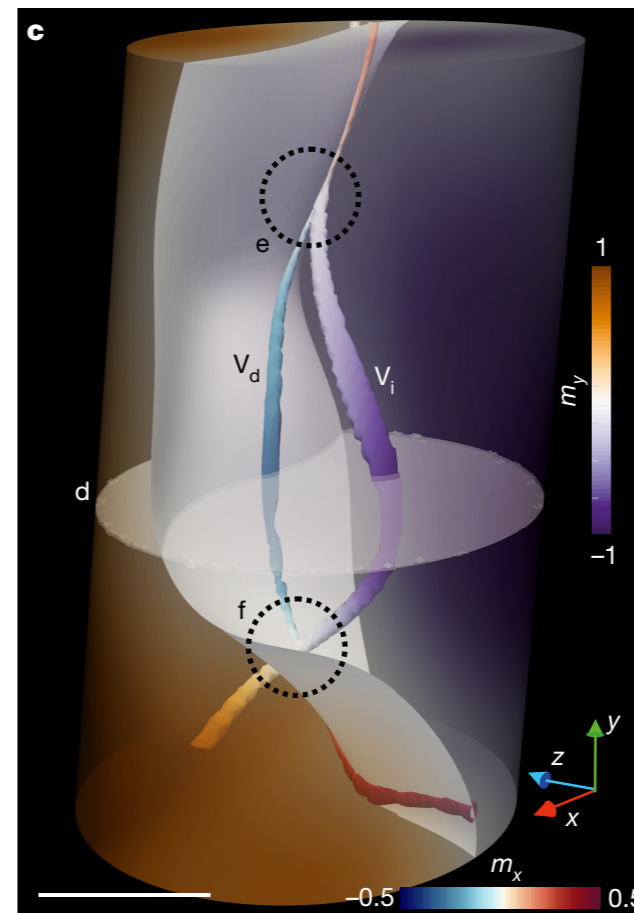
T. Shinjo et al, Science 289, 930 (2000)



## YIG disks



Losby et al, Science 350, 798 (2015)



## Cobalt Gadolinium pillars

C. Donally et al,  
Nature 547  
328 (2017)

# Optomagnonics beyond the Kittel mode

$$H_{\text{MO}} = -i \frac{\theta_{\text{F}} \lambda_n \varepsilon_0 \varepsilon}{2\pi} \int d\mathbf{r} \mathbf{m}(\mathbf{r}, t) \cdot [\mathbf{E}^*(\mathbf{r}, t) \times \mathbf{E}(\mathbf{r}, t)]$$

# Optomagnonics beyond the Kittel mode

$$H_{\text{MO}} = -i \frac{\theta_{\text{F}} \lambda_n \varepsilon_0 \varepsilon}{2\pi} \int d\mathbf{r} \mathbf{m}(\mathbf{r}, t) \cdot [\mathbf{E}^*(\mathbf{r}, t) \times \mathbf{E}(\mathbf{r}, t)]$$

$$\mathbf{m}(\mathbf{r}, t) = \mathbf{m}_0(\mathbf{r}) + \delta\mathbf{m}(\mathbf{r}, t)$$



small

# Optomagnonics beyond the Kittel mode

$$H_{\text{MO}} = -i \frac{\theta_{\text{F}} \lambda_n \varepsilon_0 \varepsilon}{2\pi} \int d\mathbf{r} \mathbf{m}(\mathbf{r}, t) \cdot [\mathbf{E}^*(\mathbf{r}, t) \times \mathbf{E}(\mathbf{r}, t)]$$

$$\mathbf{m}(\mathbf{r}, t) = \mathbf{m}_0(\mathbf{r}) + \delta\mathbf{m}(\mathbf{r}, t)$$

**Quantize: Holstein Primakoff to first order**

$$\delta\mathbf{m}(\mathbf{r}, t) \rightarrow \frac{1}{2} \sum_{\gamma} \left( \delta\mathbf{m}_{\gamma}(\mathbf{r}) \hat{b}_{\gamma} e^{-i\omega_{\gamma} t} + \delta\mathbf{m}_{\gamma}^*(\mathbf{r}) \hat{b}_{\gamma}^{\dagger} e^{i\omega_{\gamma} t} \right)$$

↓

magnon mode index      mode functions      bosonic operator

# Optomagnonics beyond the Kittel mode

$$H_{\text{MO}} = -i \frac{\theta_{\text{F}} \lambda_n \varepsilon_0 \varepsilon}{2\pi} \int d\mathbf{r} \mathbf{m}(\mathbf{r}, t) \cdot [\mathbf{E}^*(\mathbf{r}, t) \times \mathbf{E}(\mathbf{r}, t)]$$

$$\mathbf{m}(\mathbf{r}, t) = \mathbf{m}_0(\mathbf{r}) + \delta\mathbf{m}(\mathbf{r}, t)$$

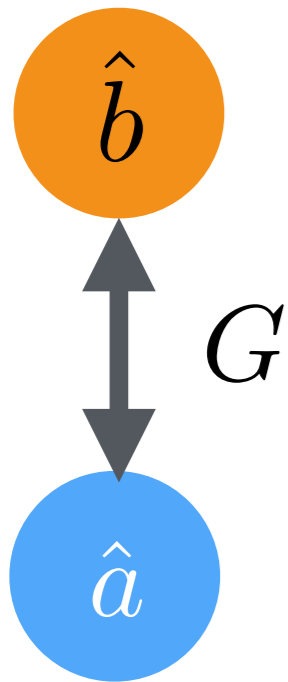
$$\delta\mathbf{m}(\mathbf{r}, t) \rightarrow \frac{1}{2} \sum_{\gamma} \left( \delta\mathbf{m}_{\gamma}(\mathbf{r}) \hat{b}_{\gamma} e^{-i\omega_{\gamma} t} + \delta\mathbf{m}_{\gamma}^*(\mathbf{r}) \hat{b}_{\gamma}^{\dagger} e^{i\omega_{\gamma} t} \right)$$

$$\mathbf{E}^{(*)}(\mathbf{r}, t) \rightarrow \sum_{\beta} \mathbf{E}_{\beta}^{(*)}(\mathbf{r}) \hat{a}_{\beta}^{(\dagger)} e^{-(+)\omega_{\beta} t}$$



# Optomagnonics beyond the Kittel mode

## Optomagnonic Hamiltonian in the spin-wave limit



$$\hat{H}_{MO} = \sum_{\alpha\beta\gamma} G_{\alpha\beta\gamma} \hat{a}_{\alpha}^{\dagger} \hat{a}_{\beta} \hat{b}_{\gamma} + \text{h.c.}$$

## Optomagnonic coupling

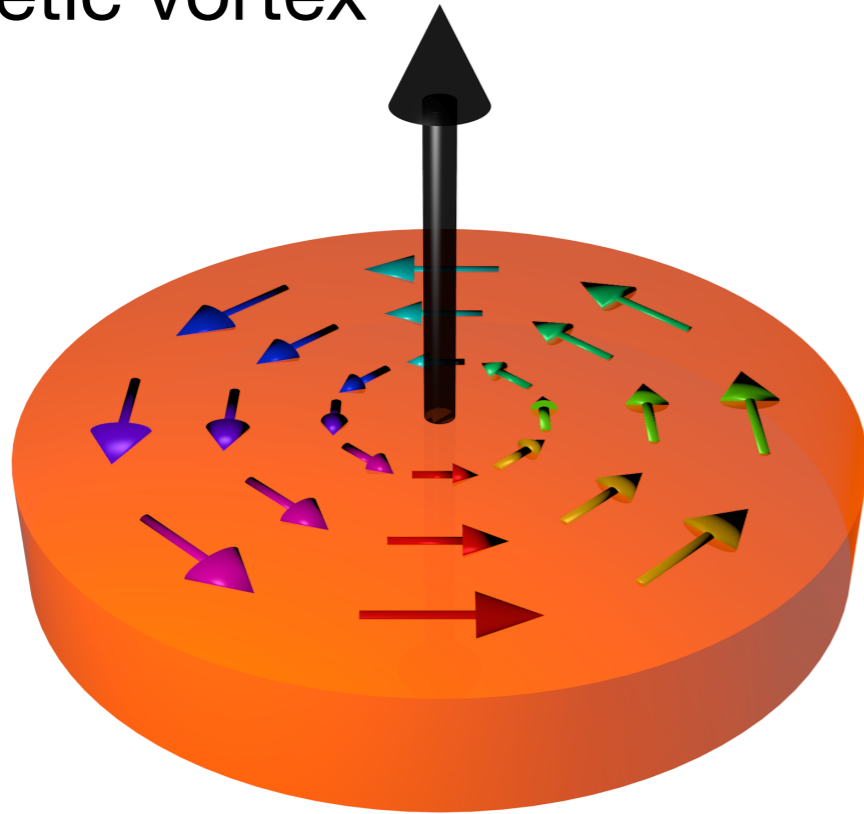
$$G_{\alpha\beta\gamma} = -i \frac{\theta_F \lambda_n \varepsilon_0 \varepsilon}{4\pi} \int d\mathbf{r} \delta \mathbf{m}_{\gamma}(\mathbf{r}) \cdot [\mathbf{E}_{\alpha}^*(\mathbf{r}) \times \mathbf{E}_{\beta}(\mathbf{r})]$$

magnon mode function

photon mode functions

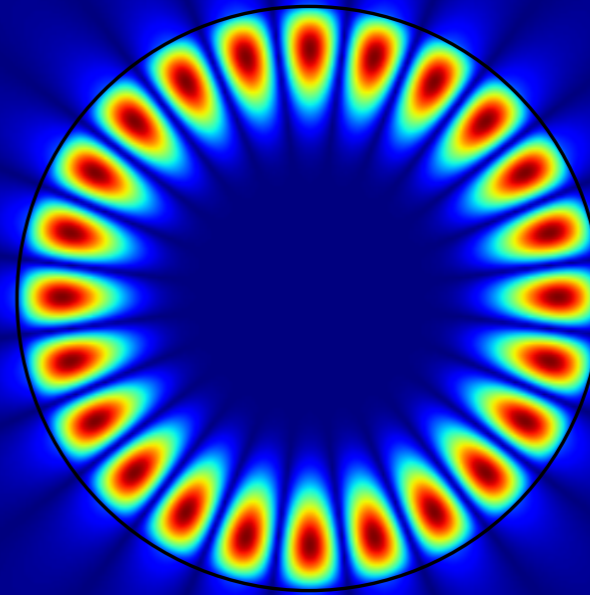
# Nanostructures: magnetic textures + light

Magnetic vortex

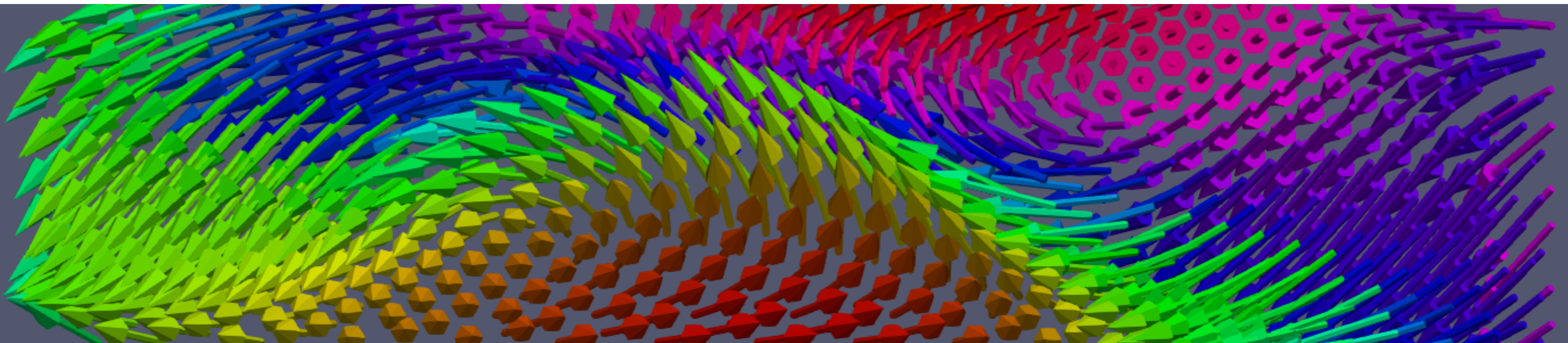


+

Optical

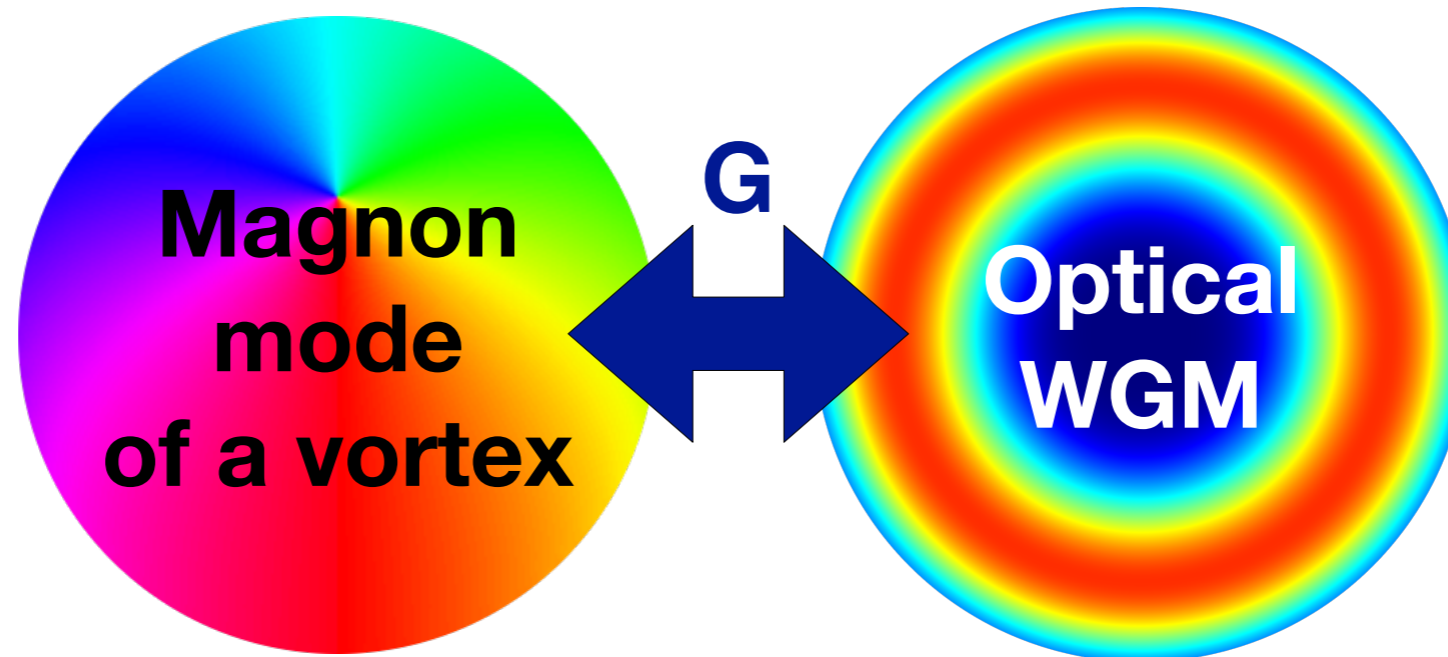


Whispering Gallery Modes



# Optomagnonic Coupling

$$G_{\alpha\beta\gamma} = -i \frac{\theta_F \lambda_n \varepsilon_0 \varepsilon}{4\pi} \int d\mathbf{r} \delta\mathbf{m}_\gamma(\mathbf{r}) \cdot [\mathbf{E}_\alpha^*(\mathbf{r}) \times \mathbf{E}_\beta(\mathbf{r})]$$



MuMax3

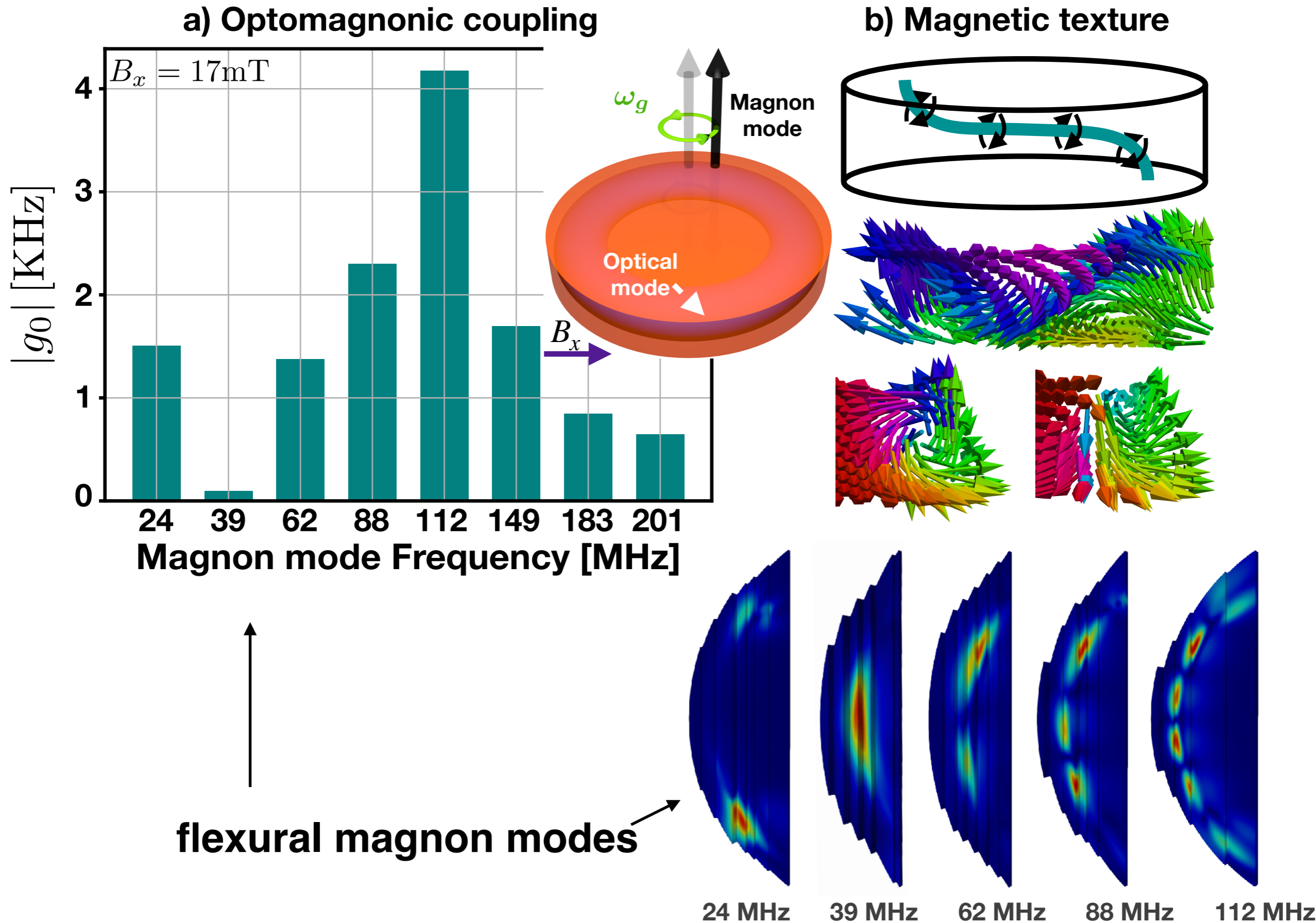
Vansteenkiste et. al  
AIP Advances 4,  
107133 (2014)

micromagnetics

finite element

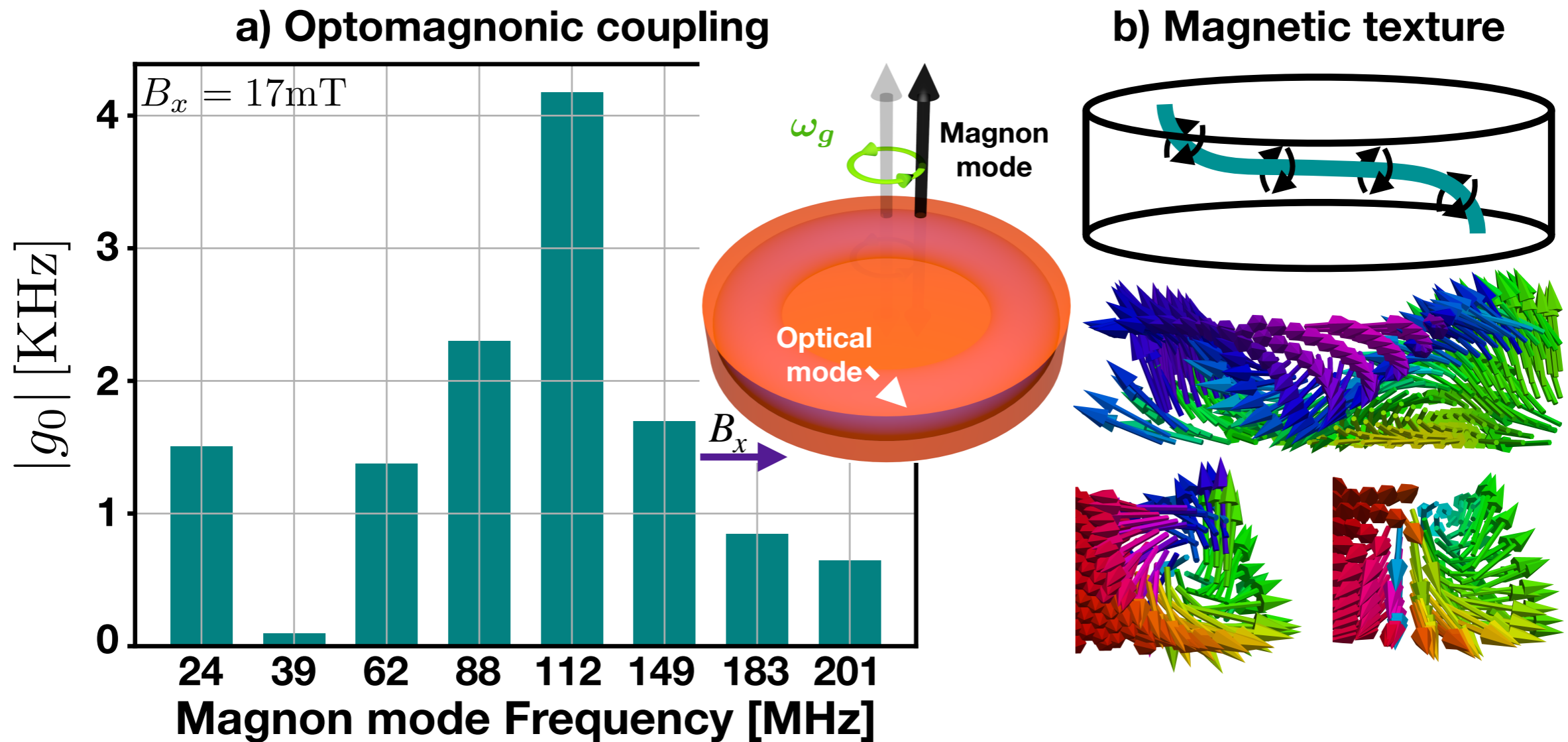
Simulation software

# Nanostructures: magnetic textures + light





# Nanostructures: magnetic textures + light



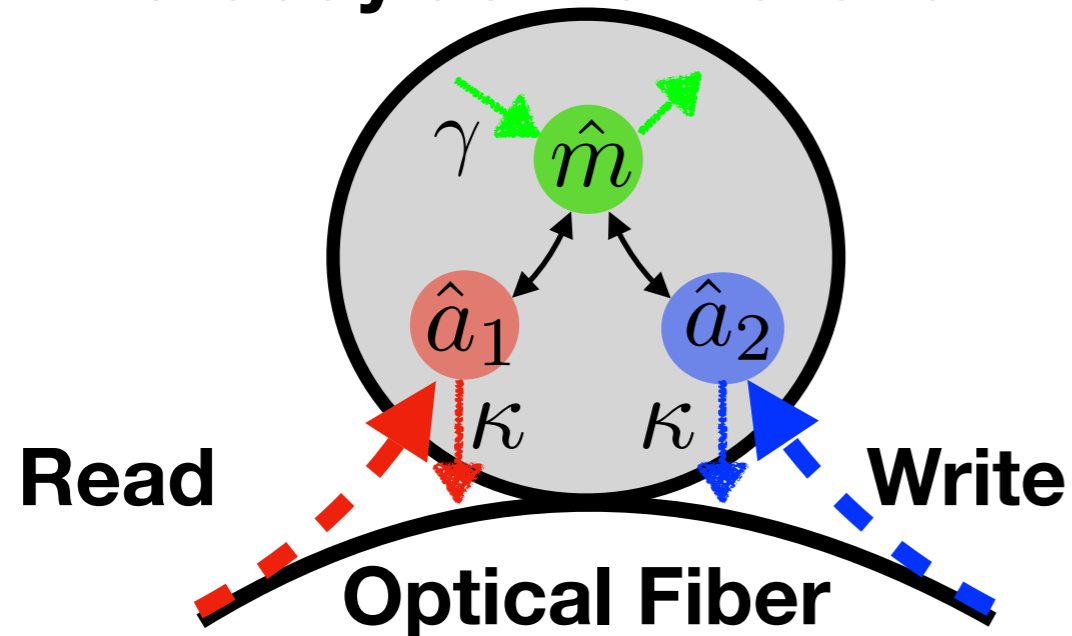
- Coupling for each mode tunable by a magnetic field

- Predicted cooperativity  $\mathcal{C} = 4n_{\text{ph}} \frac{g_0^2}{\Gamma \kappa} \approx 10^{-2}$

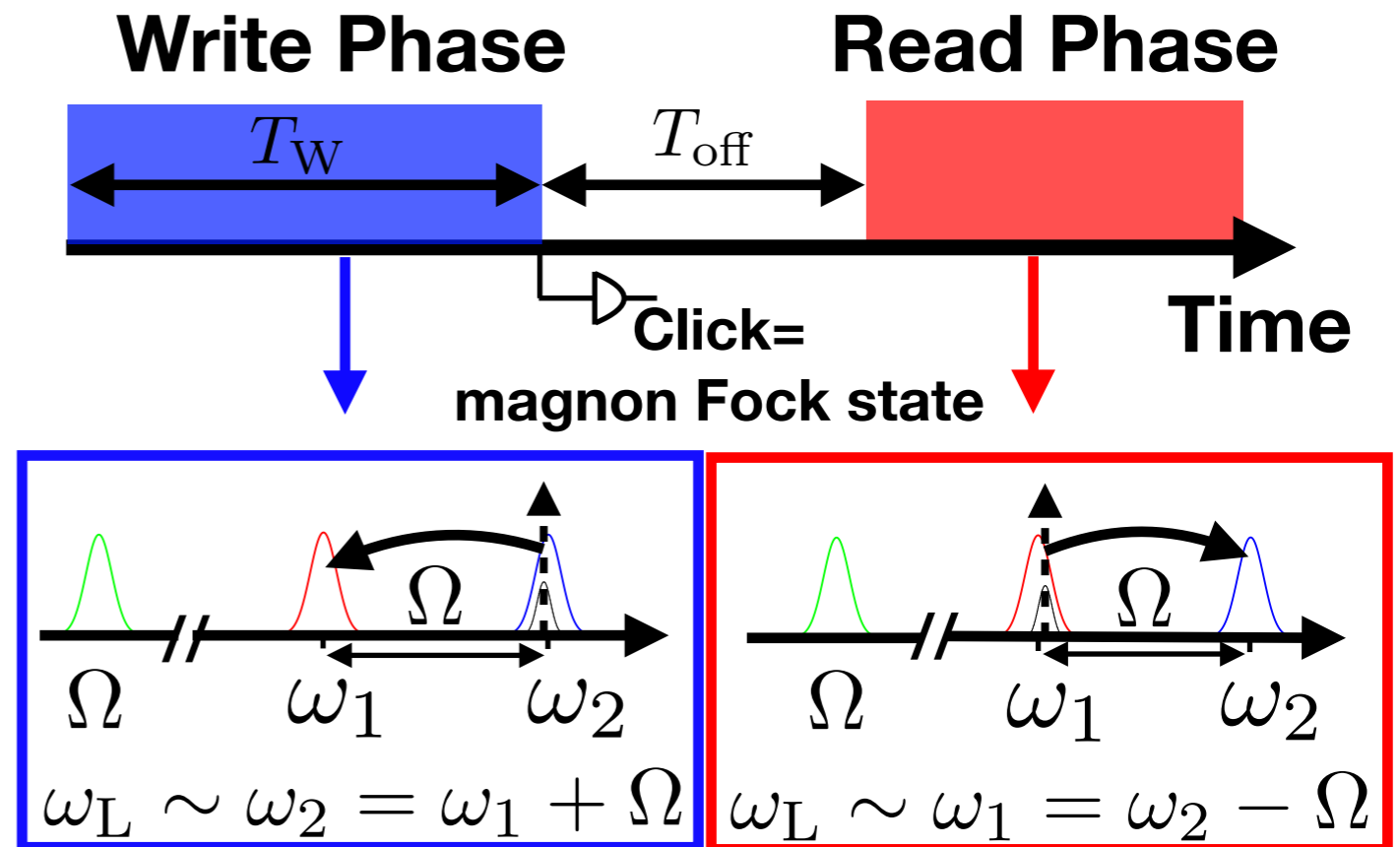


# Application: Magnon heralding

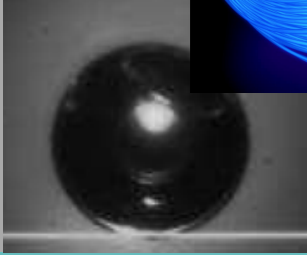
Faraday active material



## All-optical protocol for generating magnon Fock states

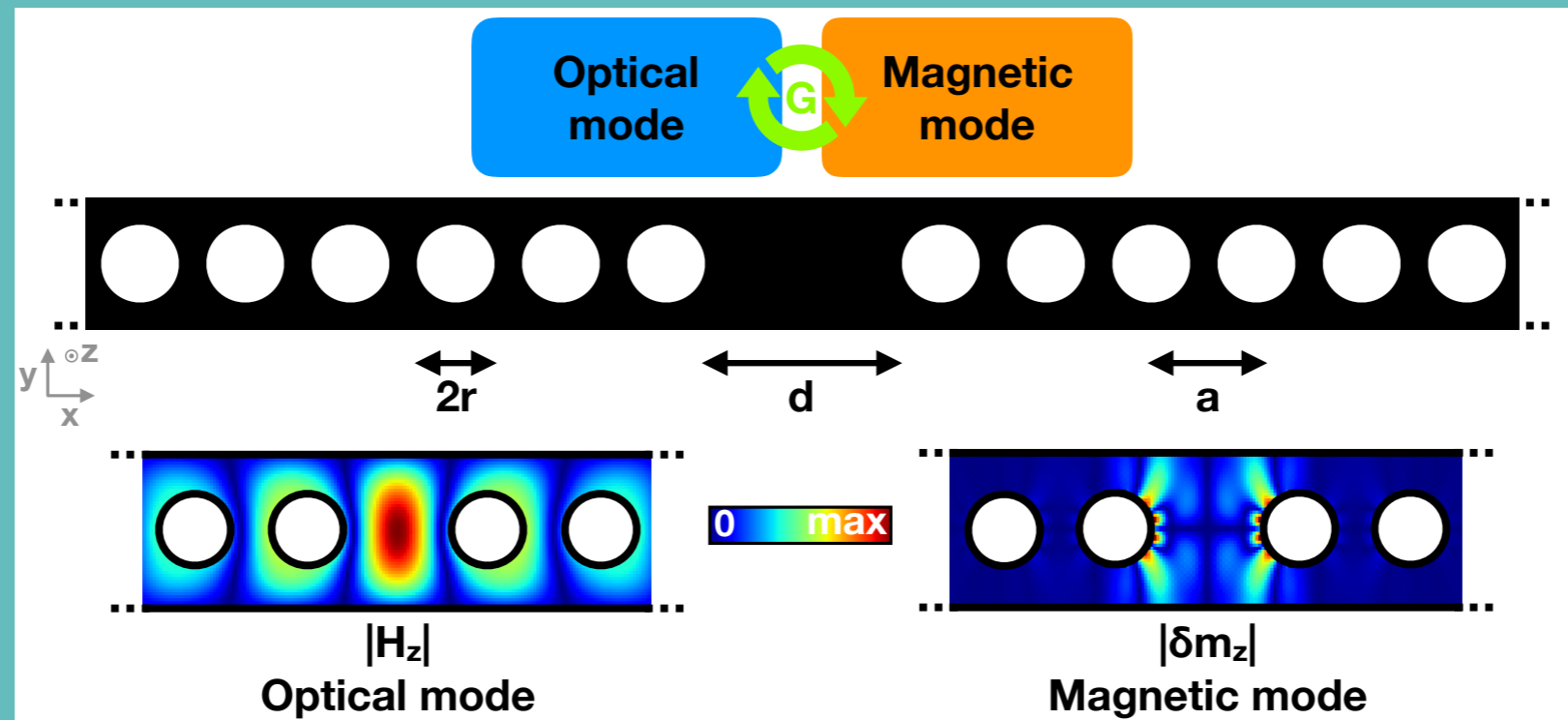


V. Bittencourt, V. Feulner, SVK;  
PRA **100**, 013810 (2019)



Optical regime

# Optomagnonic crystals



J. Graf, S. Sharma, H. Huebl, SVK;  
Physical Review Research 3, 013277 (2021)



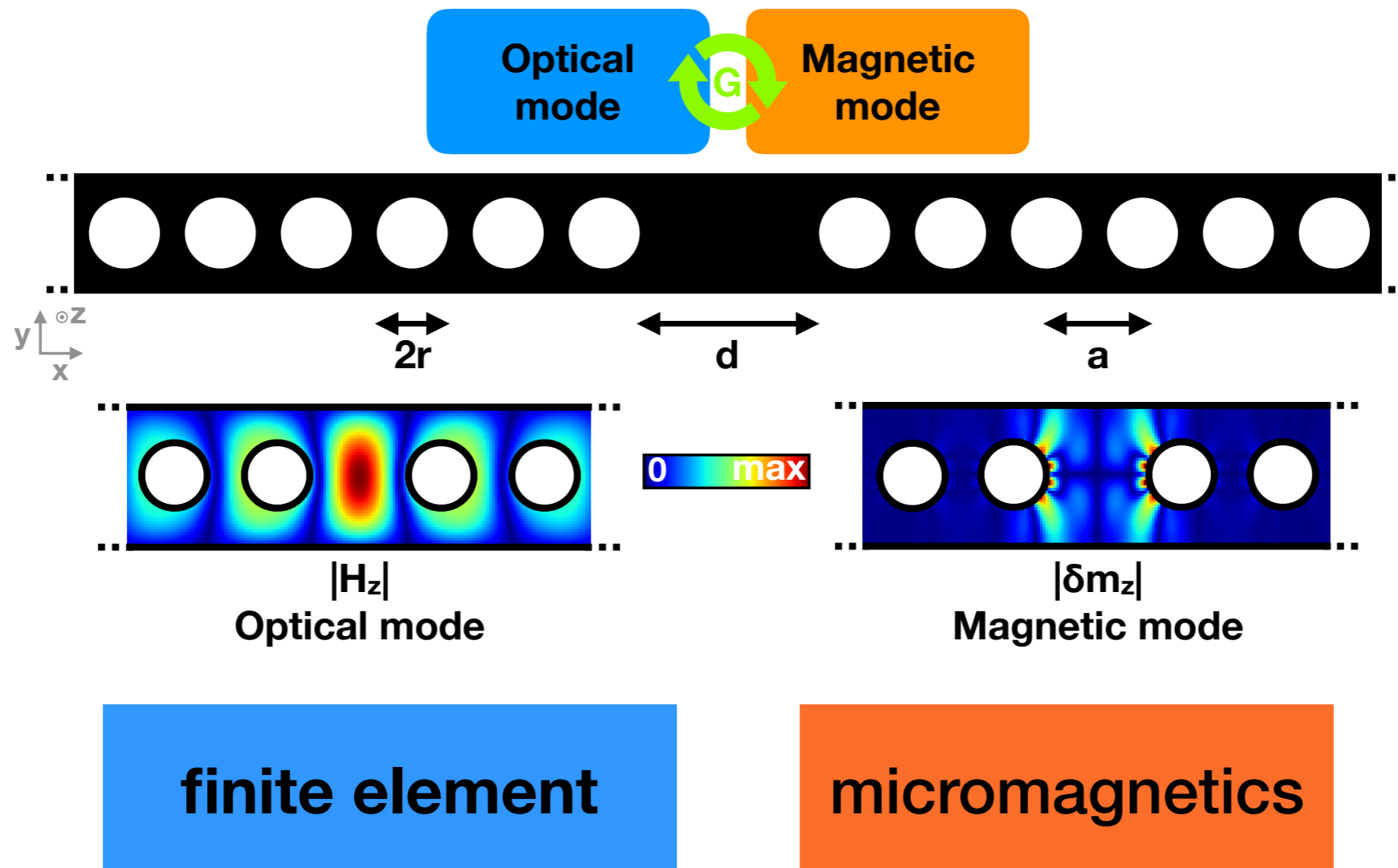
MAX PLANCK INSTITUTE  
FOR THE SCIENCE OF LIGHT

FAU  
FRIEDRICH-ALEXANDER  
UNIVERSITÄT  
ERLANGEN-NÜRNBERG



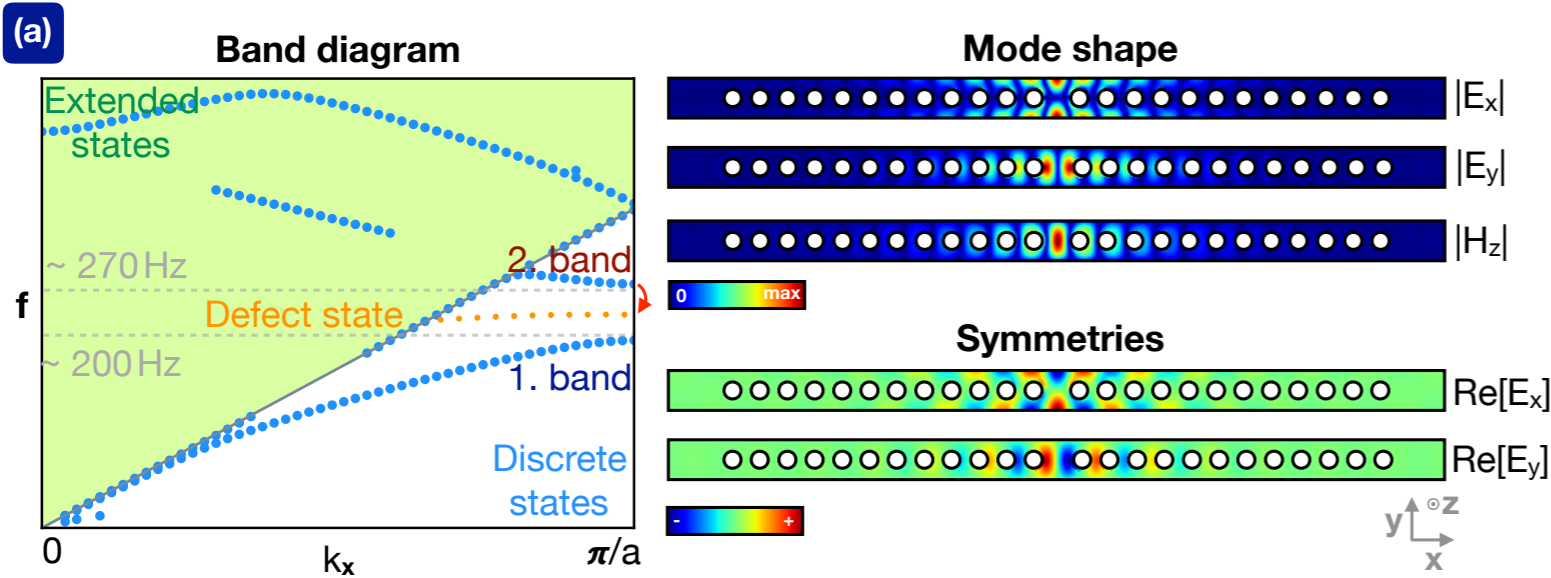
# Optomagnonic Coupling

$$G_{\alpha\beta\gamma} = -i \frac{\theta_F \lambda_n \varepsilon_0 \varepsilon}{4\pi} \frac{1}{2} \int d\mathbf{r} \delta\mathbf{m}_\gamma(\mathbf{r}) \cdot [\mathbf{E}_\alpha^*(\mathbf{r}) \times \mathbf{E}_\beta(\mathbf{r})]$$

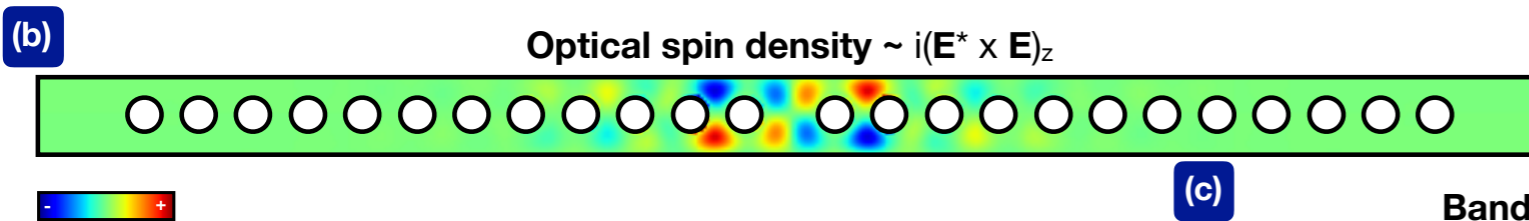


# Design: towards optimal mode matching

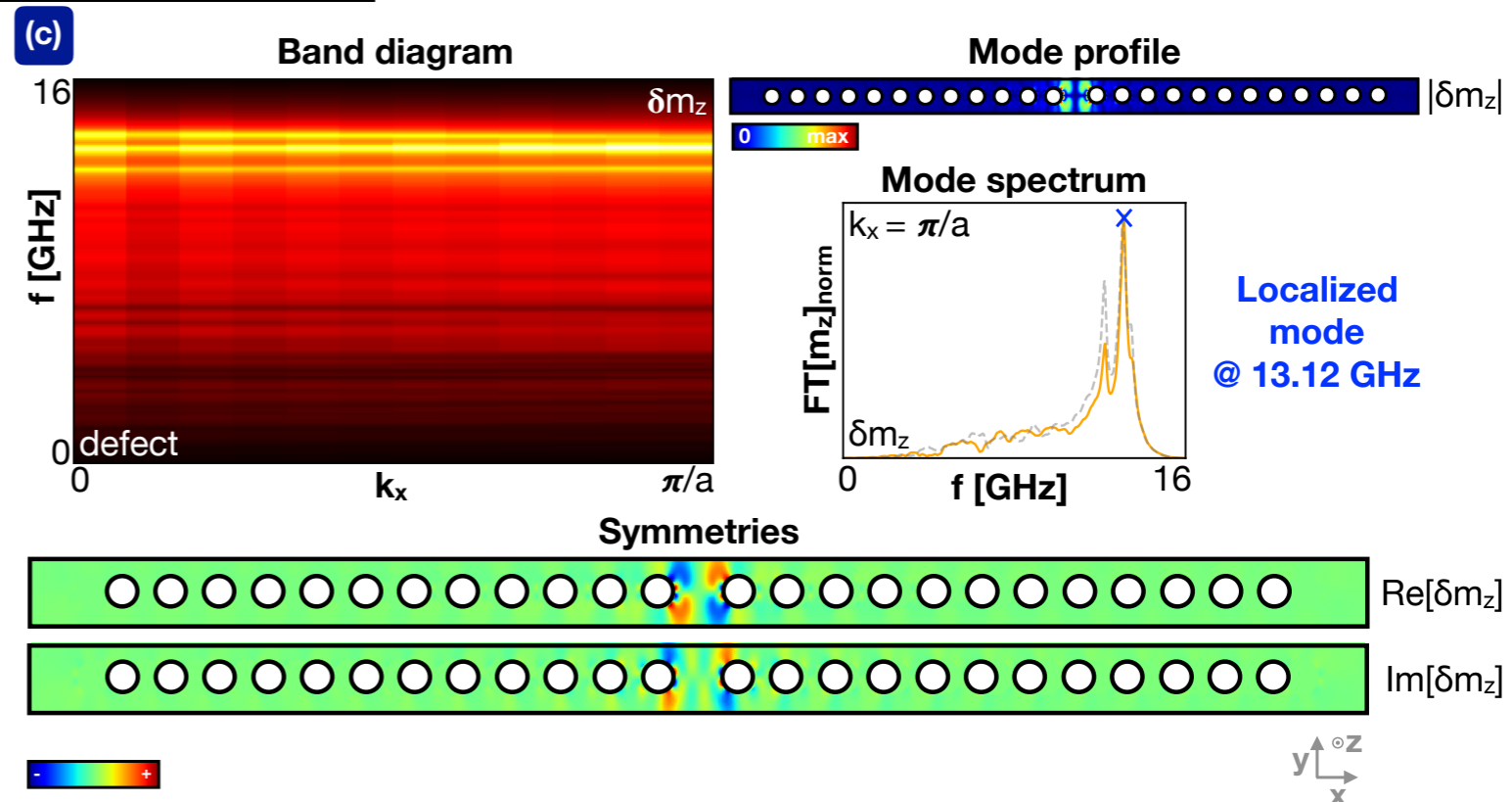
## Optics



→ Localized optical mode



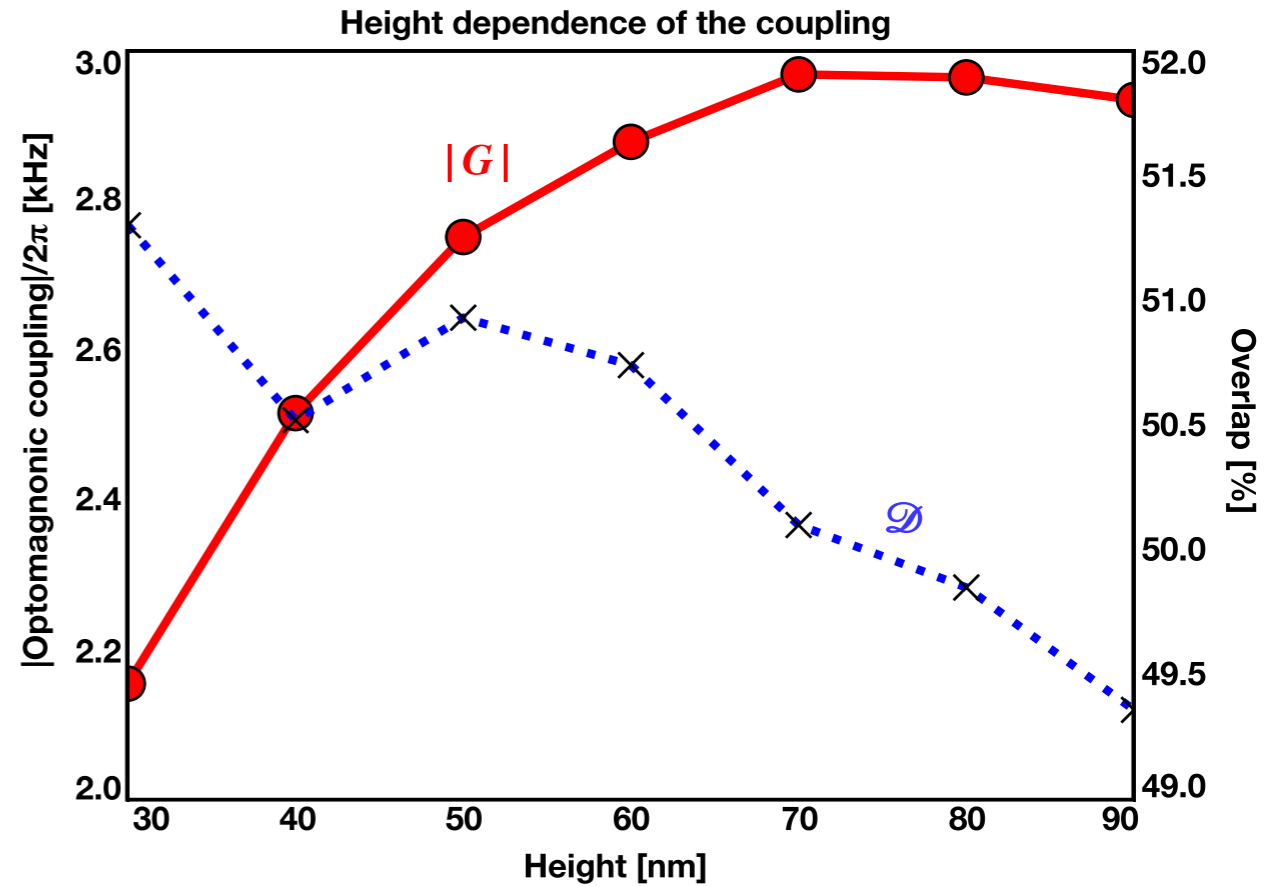
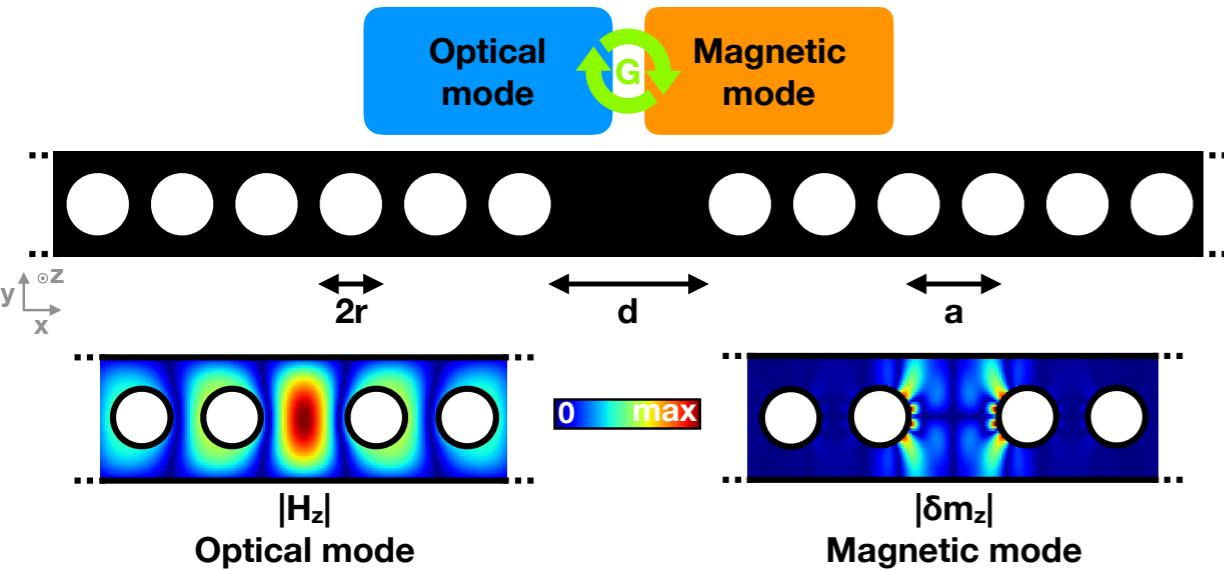
## Magnetics



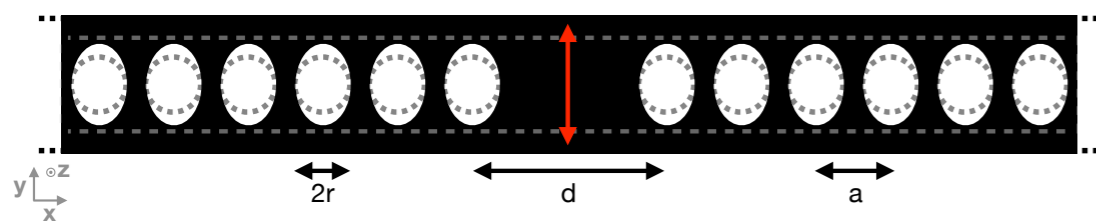
← Localized magnon mode

Match symmetries of magnon mode and optical spin density

# Design: towards optimal mode matching



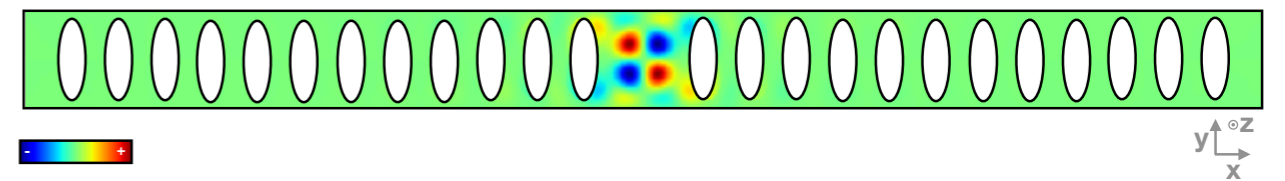
## Optimization



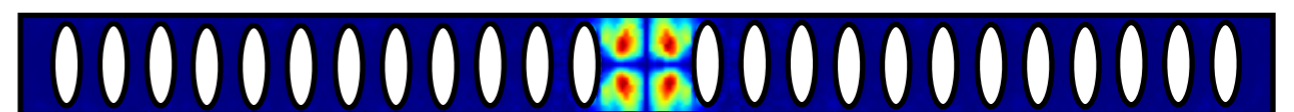
- One order of magnitude gain in coupling
- **However very lossy optical mode**

**further work needed**

## Optical spin density



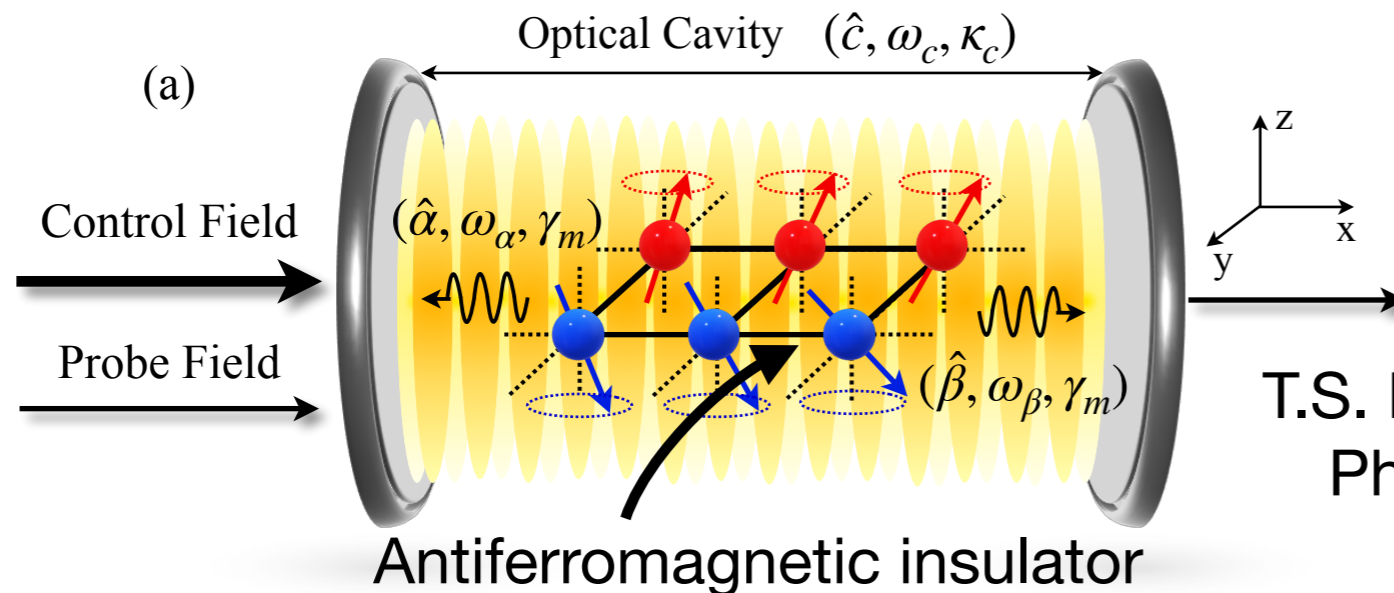
## Magnon mode





# What about Antiferromagnets?

## Antiferromagnetic cavity optomagnonics



T.S. Parvini, Victor A.S.V. Bittencourt, SVK;  
Phys. Rev. Research 2, 022027 (2020)

- Model for cavity optomagnonics with antiferromagnetic insulators  
(one-magnon scattering processes)
- Strong coupling requires large magneto-optical asymmetry  
(compared with known values for simple AFMs)
- Tunable coupling with a magnetic field  
(for finite hard-axis anisotropy)
- Dark-to-bright transition for magnon modes possible



# Outlook

- State of the art in the MW regime ripe for generation of macroscopic non-classical states of the magnetization: e.g. cat states
- Coupling to vibrations at the single quanta can be harnessed for applications: e.g. quantum thermometer
- Design of optomagnonic systems in the optical regime shows promise for improved coupling values
- Developed a method for calculating the optomagnonic coupling in the presence of magnetic textures: more to follow



# Thank you



MAX-PLANCK-  
GESELLSCHAFT



**DFG** Deutsche  
Forschungsgemeinschaft  
German Research Foundation

