

# Ab initio theory for coherent magnetic switching

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## *In collaboration with*



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Marco Berritta



Ashis Nandy



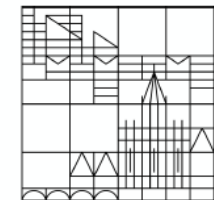
Ritwik Mondal



UPPSALA  
UNIVERSITET

Tobias Dannegger, Severin Selzer, Ulrike Ritzmann, Ulrich Nowak

Universität  
Konstanz



Eszter Simon, András Deák, László Szunyogh



Karel Carva



# General ways to switch AFMs or ferrimagnets

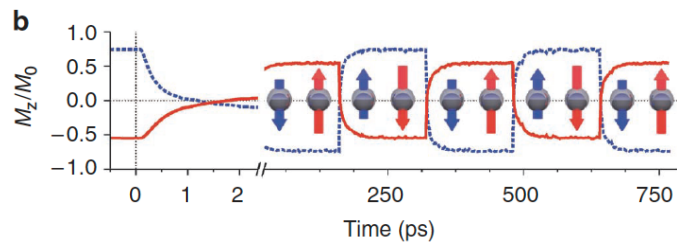
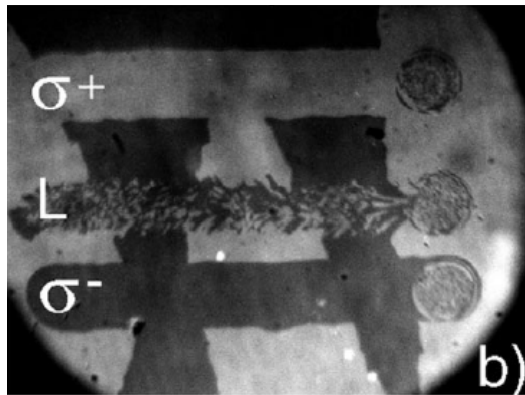
Incoherent processes



Coherent processes

Heating, "destroy & rebuild"

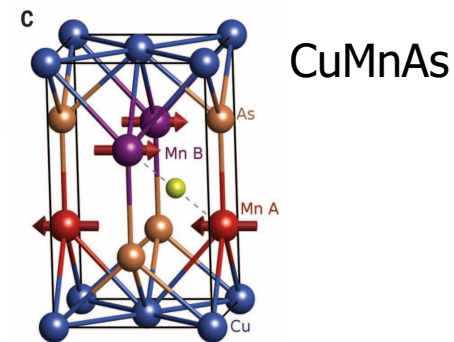
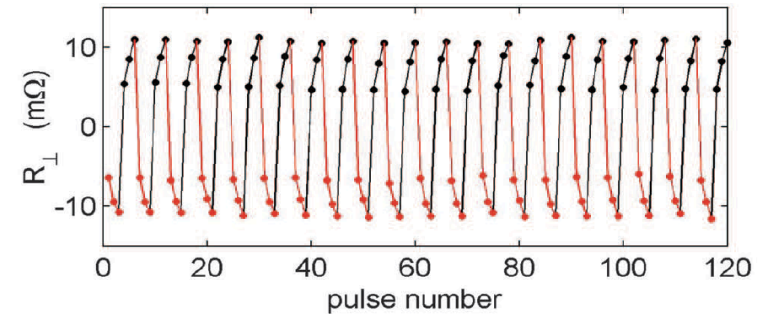
E.g., GdFeCo



Stanciu et al, PRL **99**, 047601 (2007)

Ostler et al, Nat. Comm. **3**, 666 (2012)

E.g. (Rashba - ) Edelstein effect  
"staggered torque in AFMs"



Zelezny et al, PRL **113**, 157201 (2014)

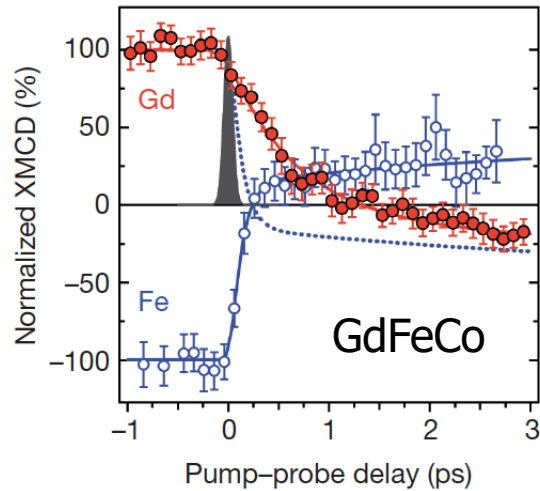
Wadley et al, Science **351**, 587 (2016)

Zelezny et al, PRB **95**, 014403 (2017)

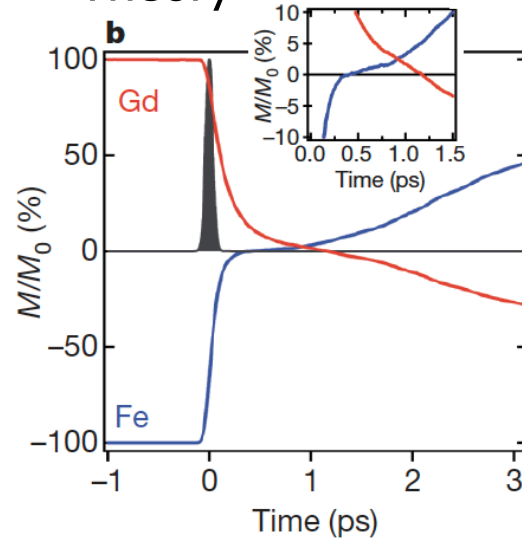


# Switching in ferrimagnets

## Experiment



## Theory



“Destroy & rebuild”

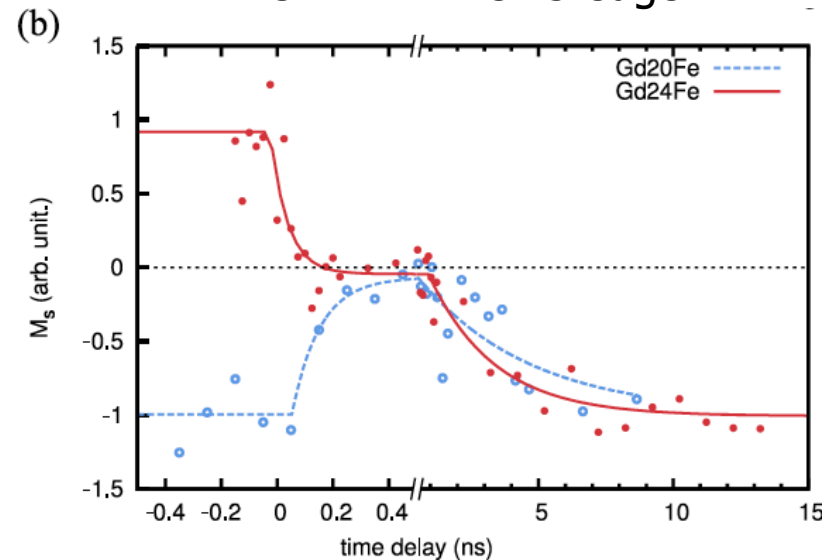
Nearly fully demagnetized  $\sim 1$  ps

Radu et al, Nature **472**, 205 (2011)

Very long *restoring* time  $\sim 1$  ns  
of incoherent process

Re-writable at  $> 0.5$  ns

## XMCD PEEM Fe L3 edge

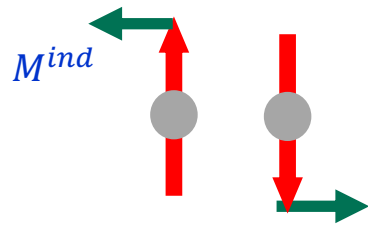


LeGuyader et al, PRB **93**, 134402 (2016)

# Possible *coherent* processes

- (Rashba - ) Edelstein effect

$$\vec{M}^{ind}(\omega) = \chi(\omega)\vec{E}(\omega) \quad \text{Current or THz field (ps)}$$



Staggered torques

Requires local symmetry breaking

Zelezny et al, PRL **113**, 157201 (2014)

Wadley et al, Science **351**, 587 (2016)

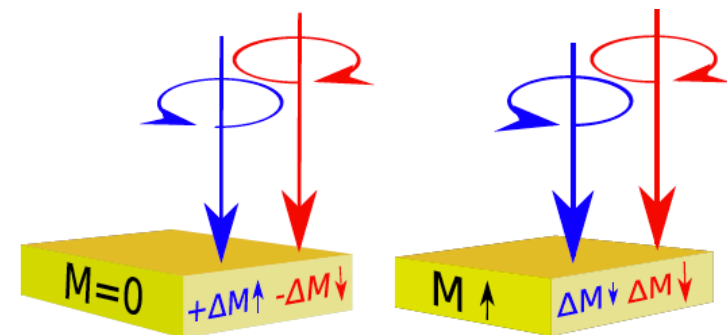
- Inverse Faraday effect

$$\vec{M}^{ind}(\omega) = \kappa^{IFE}(\omega)\vec{E}^* \times \vec{E} = K^{IFE}(\omega)E^2(\omega) \quad \text{fs-laser field}$$

Kimel et al, Nature **435**, 655 (2005)

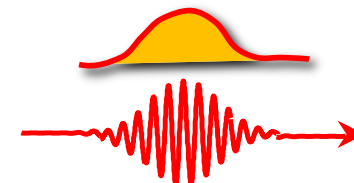
Berritta, Mondal, Carva, PMO, PRL **117**, 137203 (2016)

John et al, Sci.Rep. **7**, 4114 (2017)



symmetry breaking not required

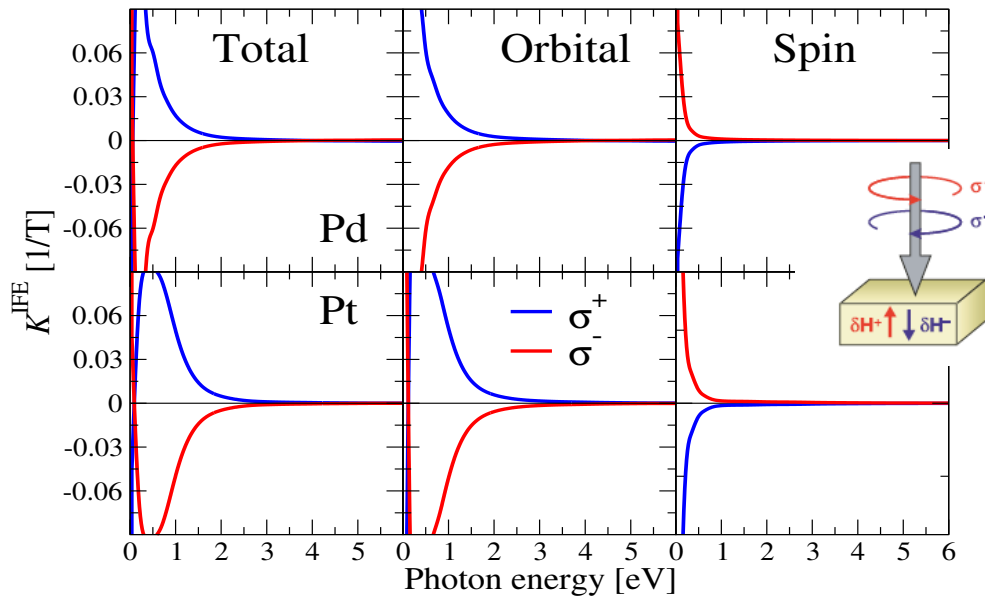
Disadvantage: not much known about it



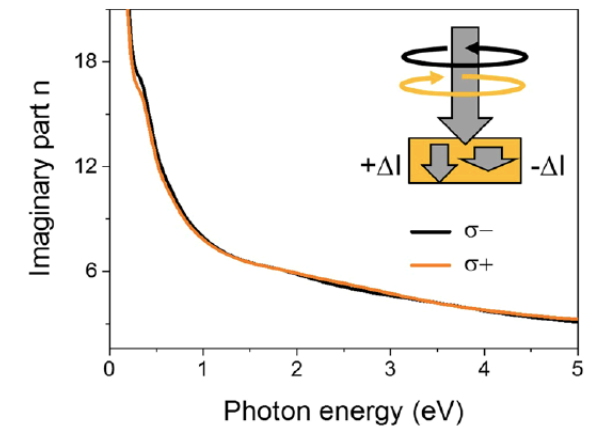
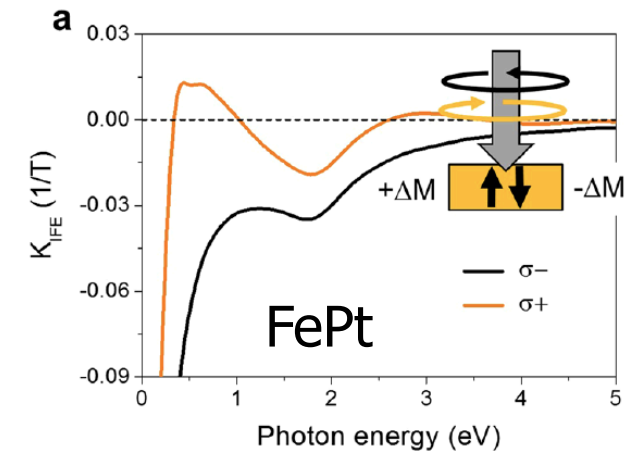
# Nonmagnetic metals / ferromagnets

$$M^{ind}(0) = [K_L^{IFE}(\omega) + K_S^{IFE}(\omega)] I/c, I = \varepsilon_0 c E^2/2 \quad (\hbar\Gamma = 0.03 Ry)$$

2nd order response theory



- *Antisymmetric* in helicity for nonmagnetic materials
- Opposite effect of spin and orbital IFE contributions possible
- Asymmetric for ferromagnets



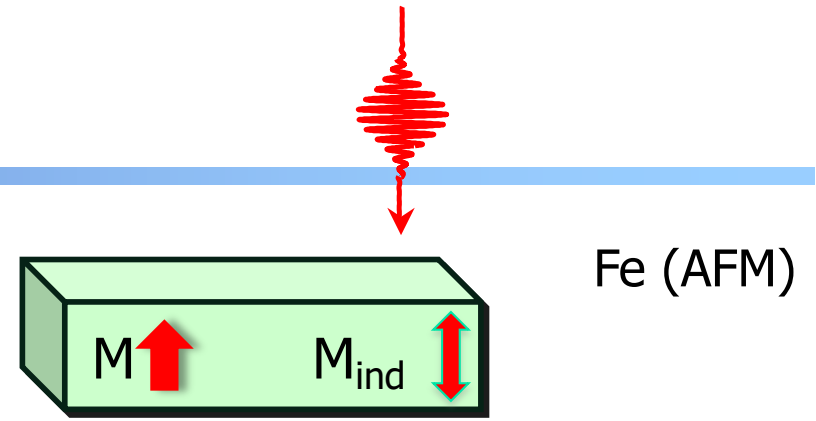
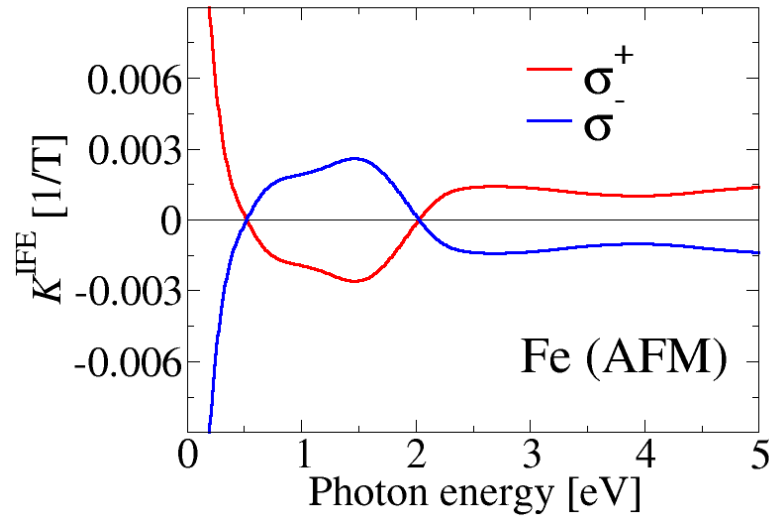
Dichroic absorption

John et al, Sci.Rep. **7**, 4114 (2017)

Berritta, Mondal, Carva, PMO, PRL **117**, 137203 (2016)

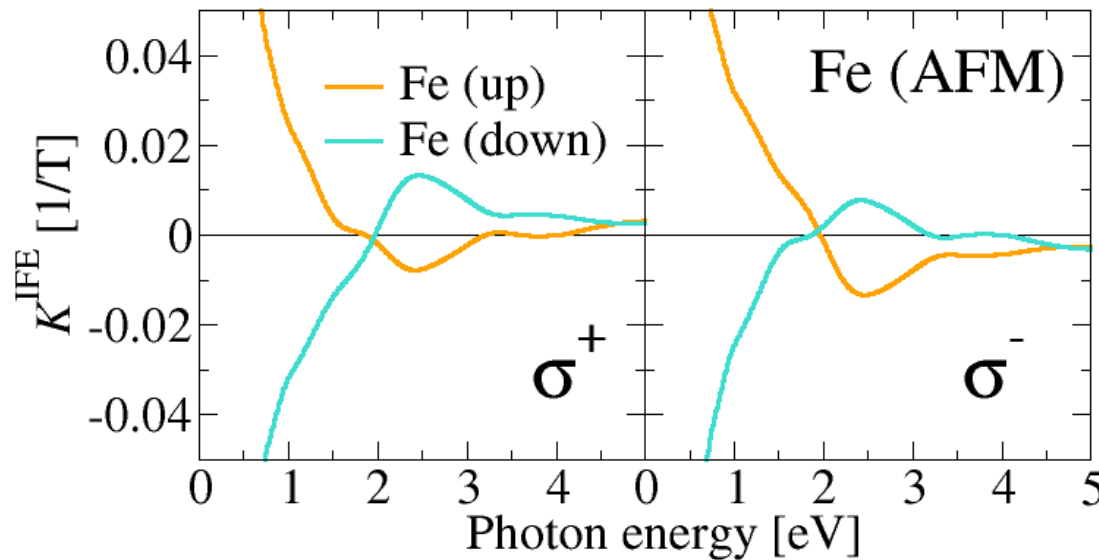
Freimuth, Blügel, Mokrousov, PRB **94**, 144432 (2016)

# Antiferromagnetic materials



$$M^{Fe} = \pm 1.635 \mu_B$$

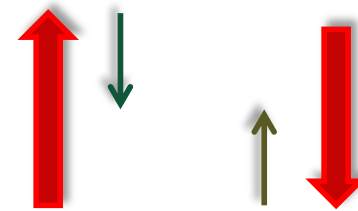
Small effect, antisymmetric?



$$I=1 \text{ GW/cm}^2 \quad \omega=2.5 \text{ eV}$$

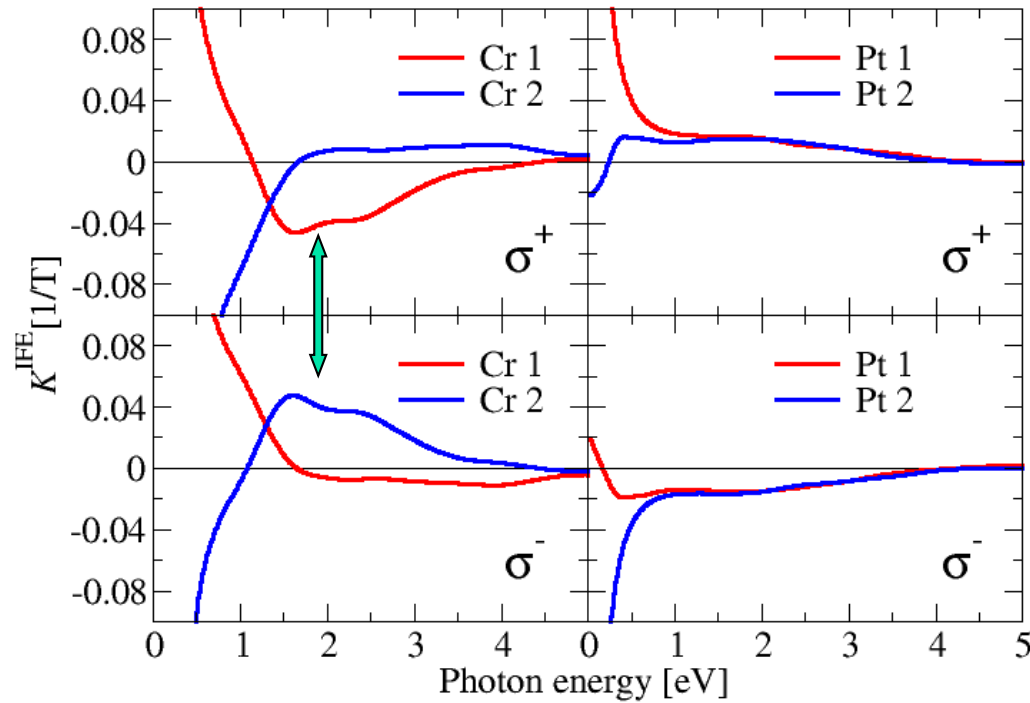
$$M_{ind}^{\uparrow}(\sigma^+) = -6.4 \times 10^{-4} \mu_B$$

$$M_{ind}^{\downarrow}(\sigma^+) = 1.1 \times 10^{-3} \mu_B$$



Gives staggered induced moments

# Antiferromagnetic CrPt



$\omega = 2.26 \text{ eV}$

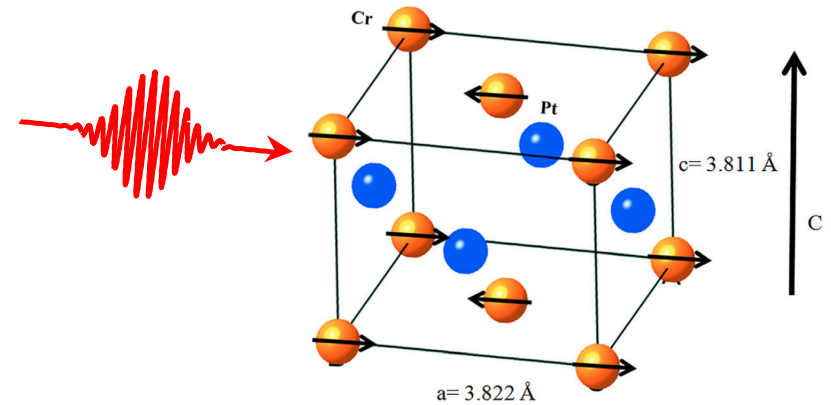
$I = 1 \text{ GW/cm}^2$

Cr<sub>1</sub>  $M^{\text{ind}}(\sigma^+) = -0.08\mu_B$

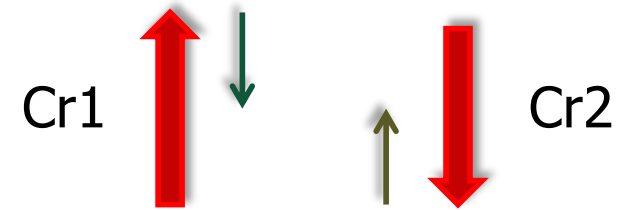
Cr<sub>2</sub>  $M^{\text{ind}}(\sigma^+) = 0.014\mu_B$

Pt<sub>1</sub>  $M^{\text{ind}}(\sigma^+) = 0.03\mu_B$

Pt<sub>2</sub>  $M^{\text{ind}}(\sigma^+) = 0.03\mu_B$



$T_N \sim 760 \text{ K}$

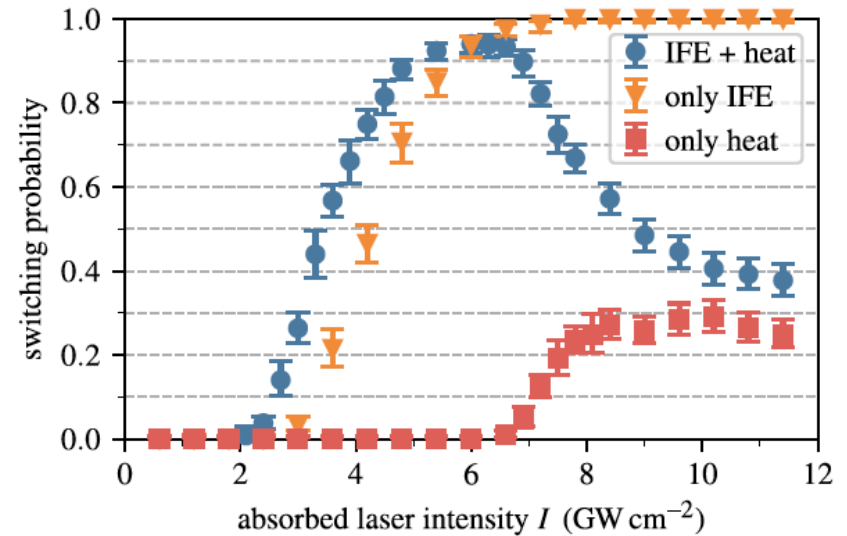
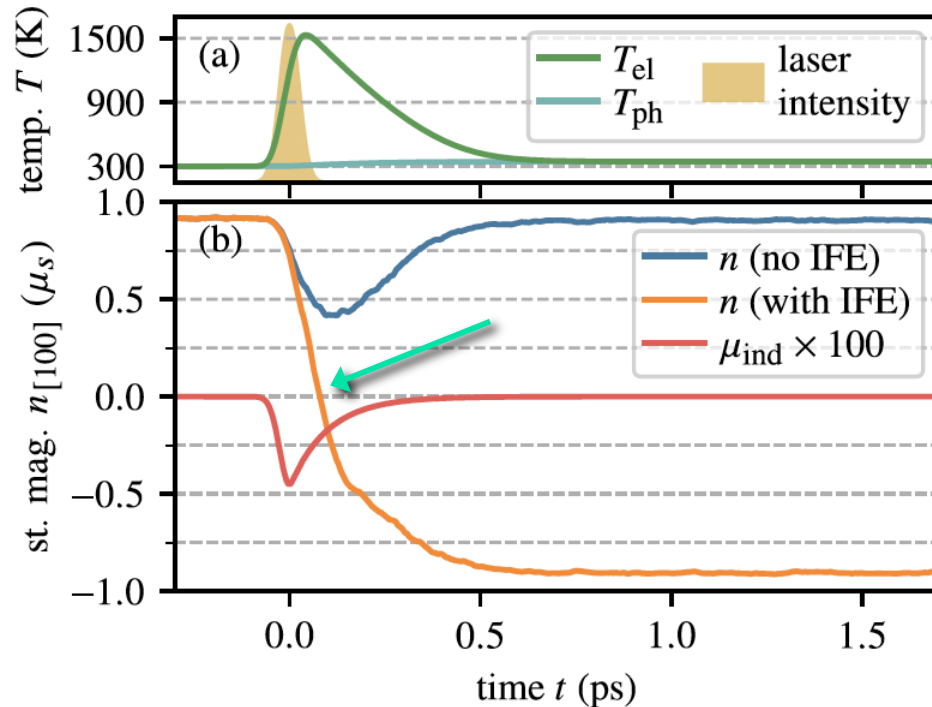


Staggered induced  
moments, not equal size

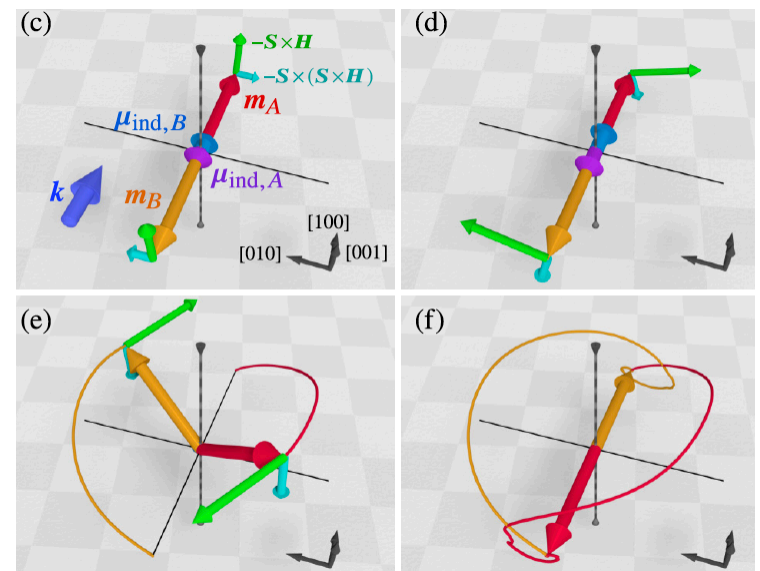


# Atomistic spin dynamics simulations

➔ Talk of Uli Nowak, poster Tobias Dannegger

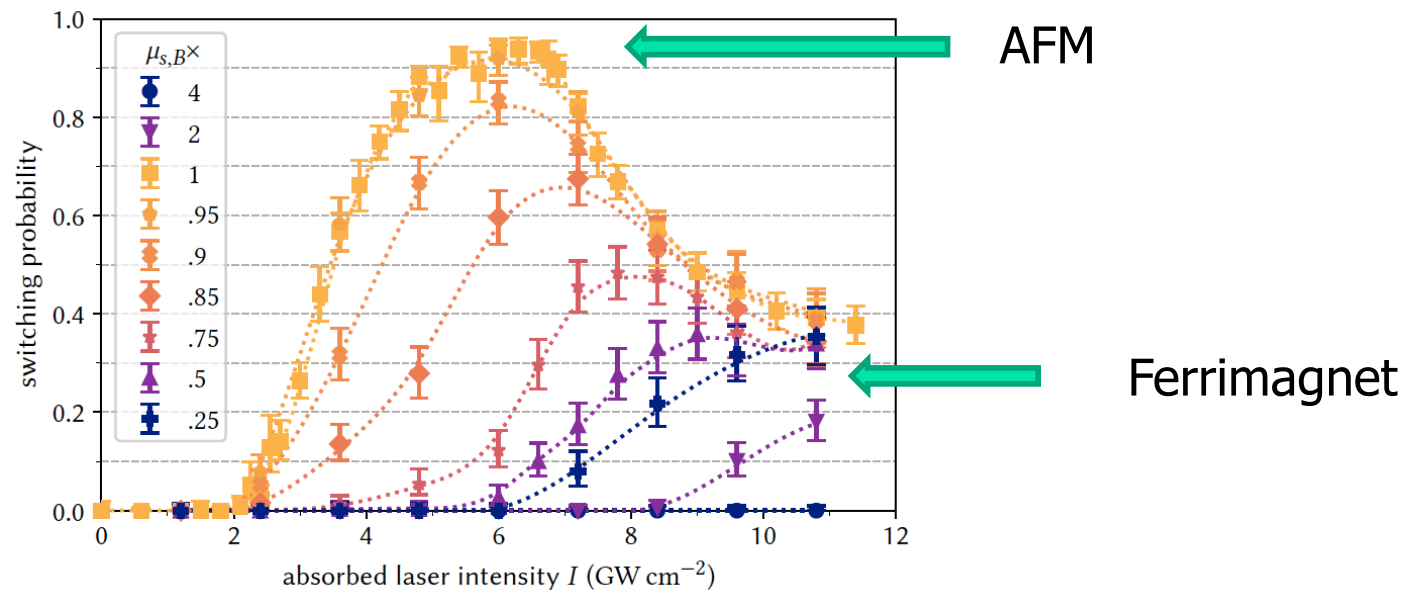


- Switches in  $\sim 200$  fs, completed at 500 fs
- Nonthermal switching, much heat works against switching probability
- $90^\circ$  coherent switching possible
- AFM exchange enhanced switching



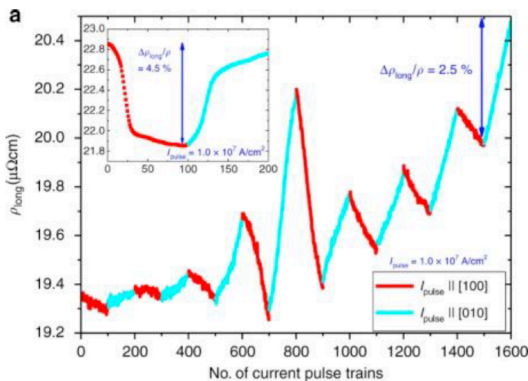
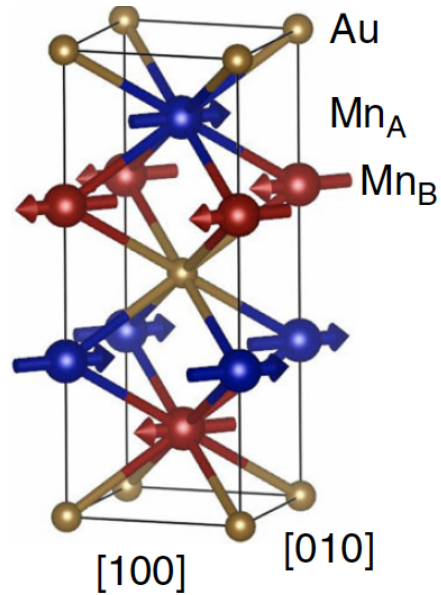
# From AFM to ferrimagnet

Ferrimagnetic order – unequal sublat. moments



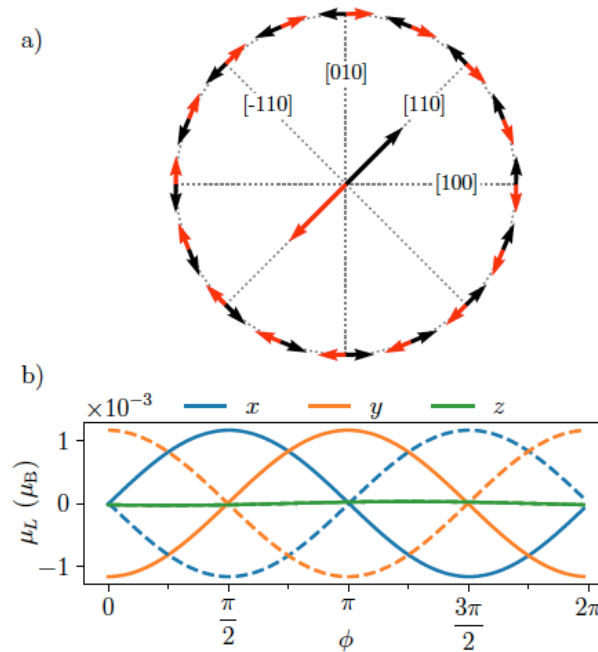


# Current induced switching in Mn<sub>2</sub>Au



Bodnar et al, Nat. Commun. **9**, 348 (2018)

Orb. Edelstein effect



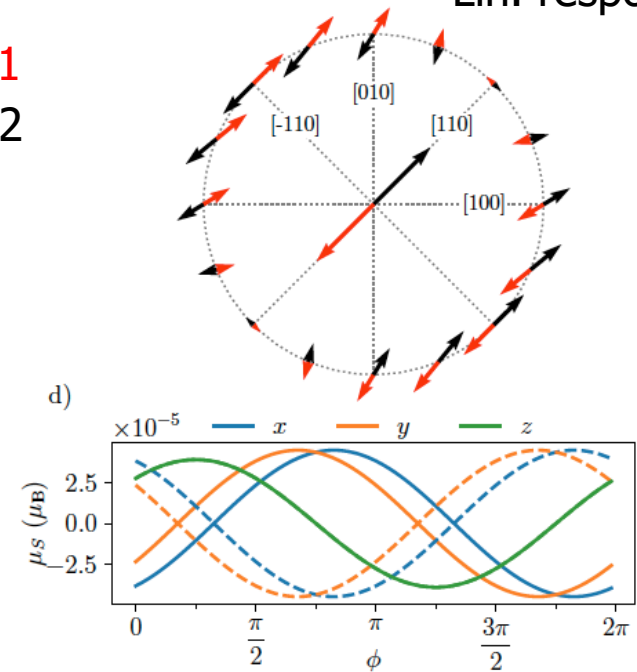
$10^{-3} \mu_B$  "Rashba"

Spin Edelstein effect

$$M_i^{\text{ind}} = \chi_{ij} \cdot E_j$$

Lin. response

Mn1  
Mn2



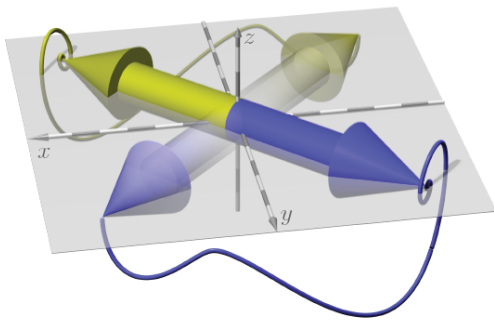
$10^{-5} \mu_B$  "Dresselhaus"

- Orbital REE dominant (not due to SOC)
- Large non-Néel elements present (spin)
- Orbital polarization is staggered

# Atomistic spin dynamics simulations

$$\mathcal{H} = -\frac{1}{2} \sum_{i \neq j} J_{ij} (\mathbf{S}_i + \mathbf{s}_i) \cdot (\mathbf{S}_j + \mathbf{s}_j) - \sum_i J^{\text{sd}} \mathbf{S}_i \cdot \mathbf{s}_i + \sum_i \xi \mathbf{S}_i \cdot \mathbf{l}_i - \sum_i d_z S_{i,z}^2 - \sum_i d_{xy} S_{i,x}^2 S_{i,y}^2,$$

With *ab initio* input

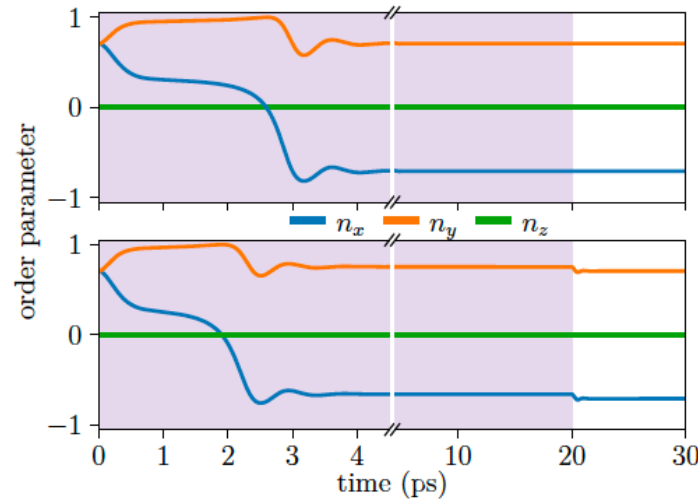


Switching path

AFM exchange enhanced

$$T \sim (E_{ani} E_{xch})^{-1/2}$$

Roy, Otxoa, Wunderlich, PRB **94**, 014439 (2016)

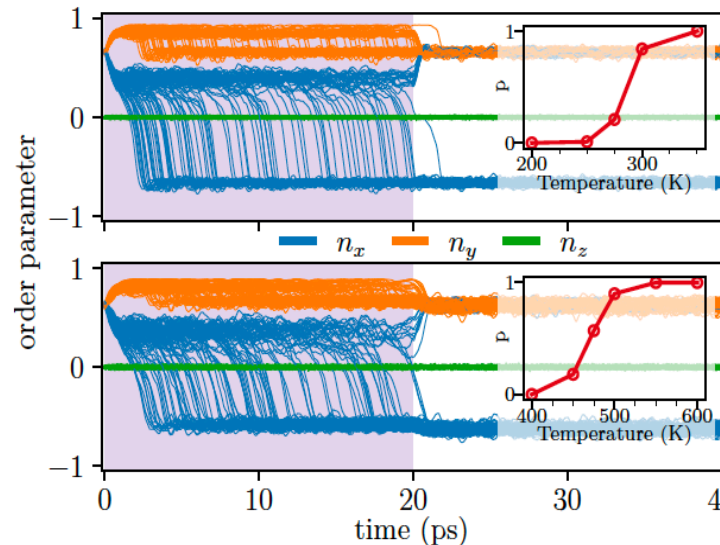


$E \parallel [110]$   
 $E = 1.9 \cdot 10^7 \text{ V/m}$

$E \parallel [100]$   
 $E = 3.1 \cdot 10^7 \text{ V/m}$

90° switching simulation for 20 ps pulse (T=0 K, high E field)

=> Fast switching ~4 ps



$E \parallel [110]$   
 $E = 10^7 \text{ V/m}, T = 300 \text{ K}$

Switching assisted by temperature !

$E \parallel [100]$   
 $E = 10^7 \text{ V/m}, T = 500 \text{ K}$

# Other possible coherent torques?

$$\frac{\partial \mathbf{M}}{\partial t} = -\gamma \mathbf{M} \times \mathbf{H}_{\text{eff}} + \alpha \mathbf{M} \times \frac{\partial \mathbf{M}}{\partial t} + \mathbf{T}$$

■ Inertial torque  $T^{\text{inert}} = \mathbf{M} \times \left[ I \cdot \frac{\partial^2 \mathbf{M}}{\partial t^2} \right]$

Short time scale (fs-ps), but intrinsic  $I$  – *how to steer it?*

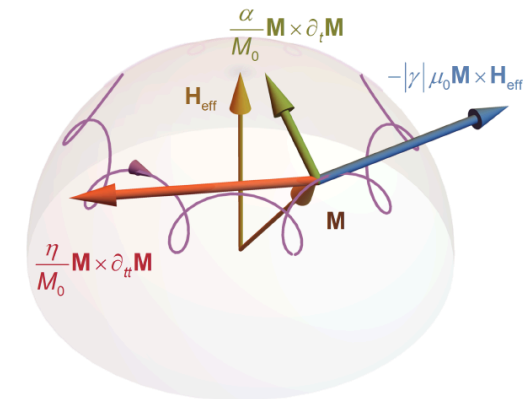
Mondal, Berritta, Nandy, PMO, PRB **96**, 024425 (2017)

Cherkasskii et al, PRB **102**, 184432 (2020)

Neeraj et al, Nat. Phys. **17**, 245 (2021)

Mondal, Großenbach, Rozsa, Nowak, PRB **103**, 104404 (2021)

$$I \propto -\bar{\tau} \alpha \quad \bar{\tau} \approx 700 \text{ fs}$$



■ Optical spin-orbit torque  $T^{\text{OSOT}} = -\frac{e^2}{2m^2 c^2 \epsilon_0} \mathbf{M} \times \mathbf{j}_s$   $\mathbf{j}_s = -2\mathbf{E} \times \mathbf{A} (\epsilon_0 = 1)$

Photon spin angular momentum

Mondal, Berritta, Paillard et al, PRB **92**, 100402R (2015)

Mondal, Berritta, Oppeneer, JPCM **29**, 194002 (2017)

Mondal, Donges, Nowak, PRRes **3**, 023116 (20121)

Some possible observations (?)

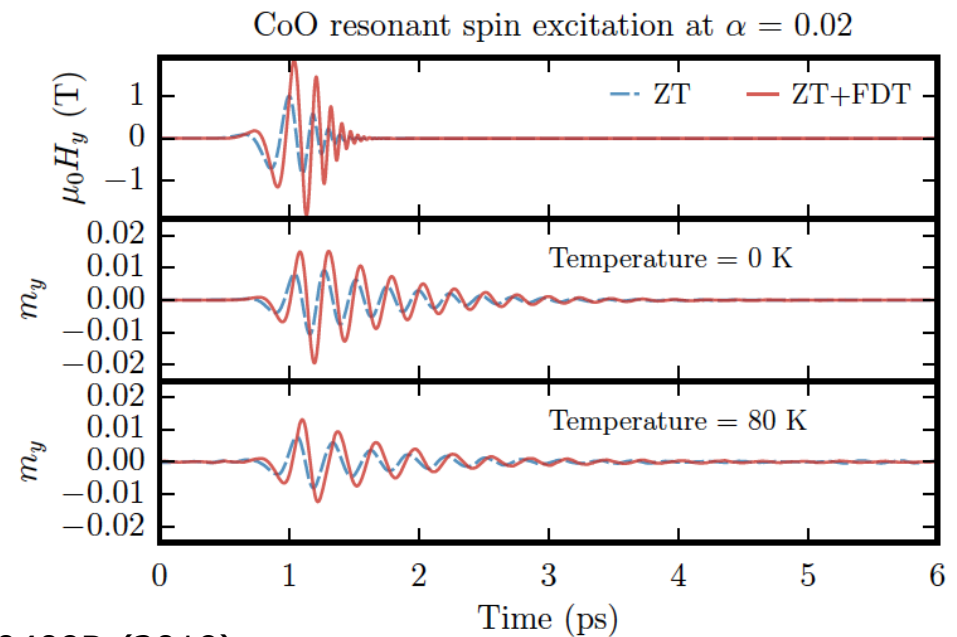
# Other possible torques?

- Field-derivative torque  $T^{FDT} \propto M \times \frac{\partial H}{\partial t}$  (non-relativistic)

- for THz pulses FDT could be important
- phase difference between Zeeman torque and FDT
- No experimental observation so far

Mondal, Berritta, PMO, PRB **94**, 144419 (2016)

Mondal, Donges, Ritzmann, PMO, Nowak, PRB **100**, 060409R (2019)





# Conclusions

- Inverse Faraday effect can give staggered induced moments in AFMs (no symmetry-breaking needed)
- Can initiate fast AFM switching process ( $\sim 200$ fs) in CrPt *w/o incoherent heating*
- Electric-field induced switching in  $\text{Mn}_2\text{Au}$  possible in  $\sim 4$  ps
- Heating by current pulses strongly assists the AFM switching in  $\text{Mn}_2\text{Au}$
- Other possible torques for ultrafast switching:  
Inertial torque, FDT, and OSOT – not enough known and possibly too small

