



New Spin on Molecular
Quantum Materials

Wasem in Mainz
May 24-26th, 2022

Randomness effect on charge glass formation in θ -type BEDT-TTF compounds



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University of Tokyo, Japan*

Collaborators & Acknowledgements

Germany



Noise & Thermal expansion

J. Müller & M. Lang
Groups in Frankfurt Univ.

FT-IR

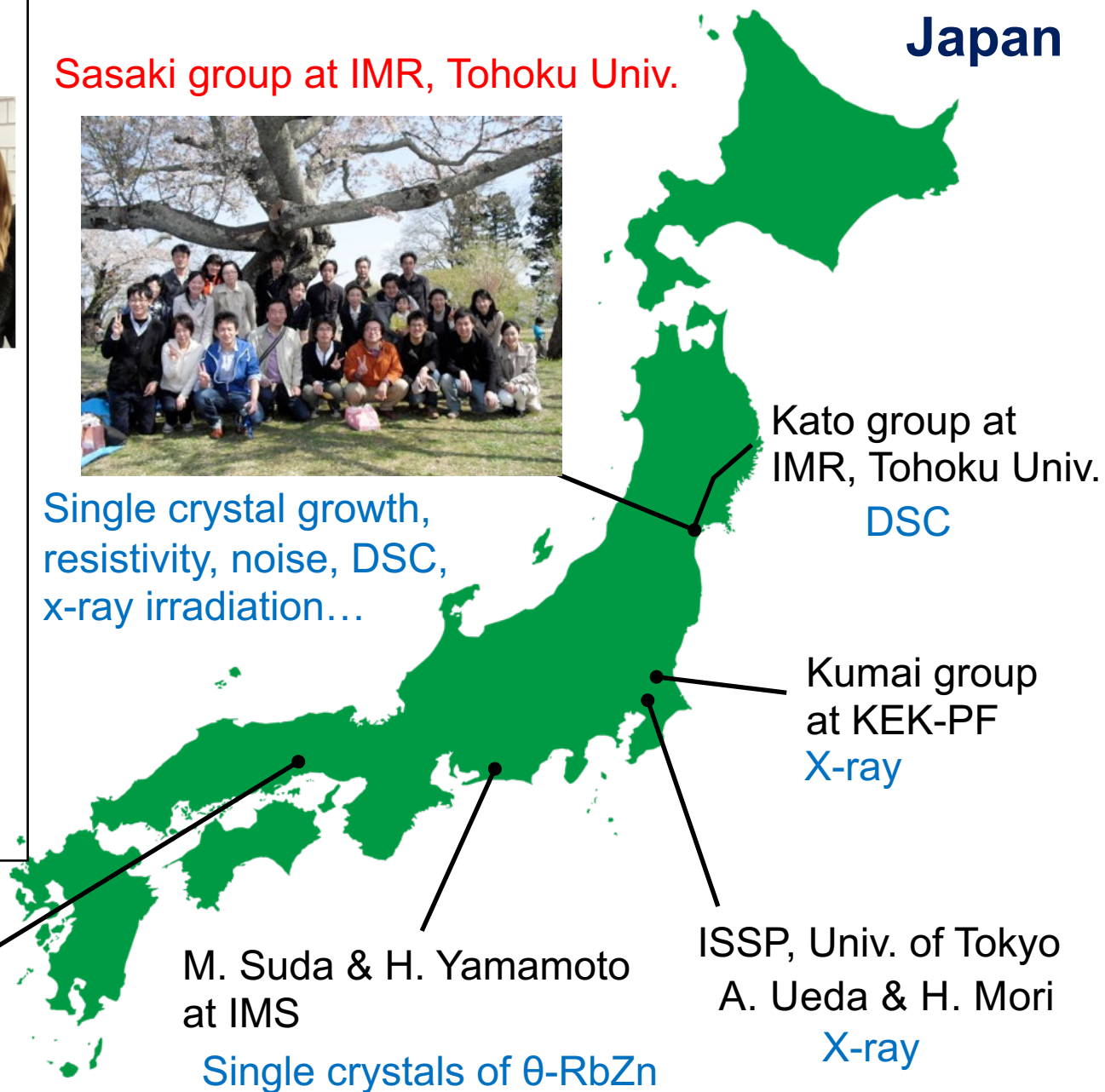
Y. Ikemoto & T. Moriwaki
at SPring-8 (BL43IR)

Sasaki group at IMR, Tohoku Univ.



Single crystal growth,
resistivity, noise, DSC,
x-ray irradiation...

Japan



1. Introduction

- Brief introduction of glass formation phenomena and charge liquid/crystal/glass states in θ -BEDT-TTF compounds

2. Comparison of CC and CG states in three different θ -type salts

- Resistivity and optical conductivity spectra in the orthorhombic and monoclinic systems

3. Charge crystallization process probed by thermodynamic quantity

- TTT diagram constructed by DSC measurements

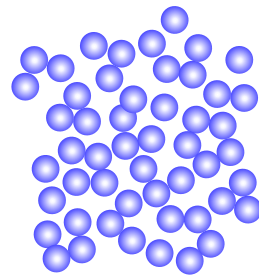
4. Randomness effects on charge crystallization and vitrification

- Reflectivity and DSC measurements for x-ray irradiated samples

5. Summary

Crystallization vs Vitrification

Glass-forming Liquid

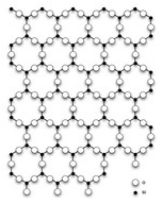


Slow Cooling

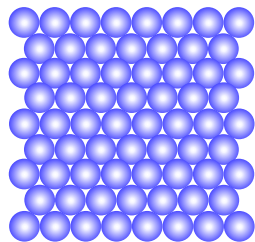
Rapid Cooling

Crystallization

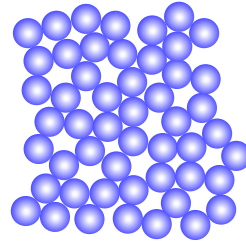
Vitrification



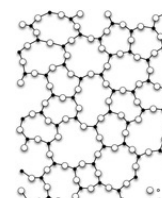
SiO_2



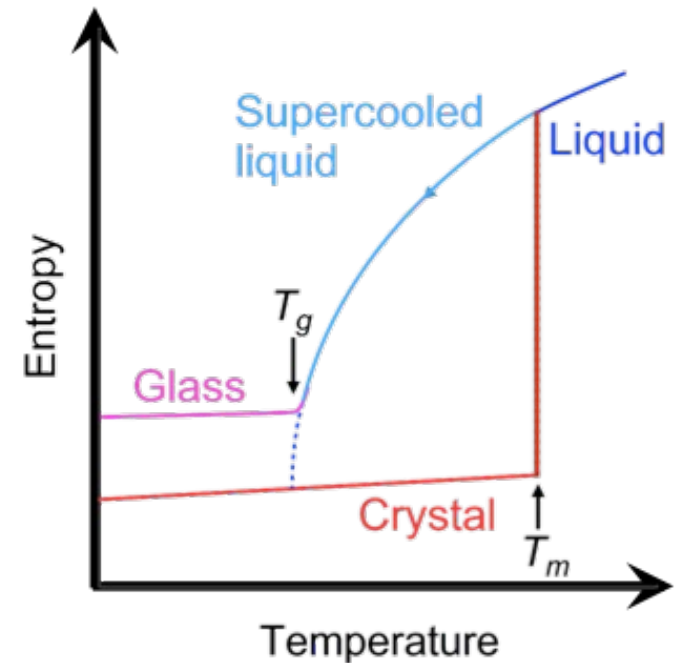
Crystal



Glass



SiO_2

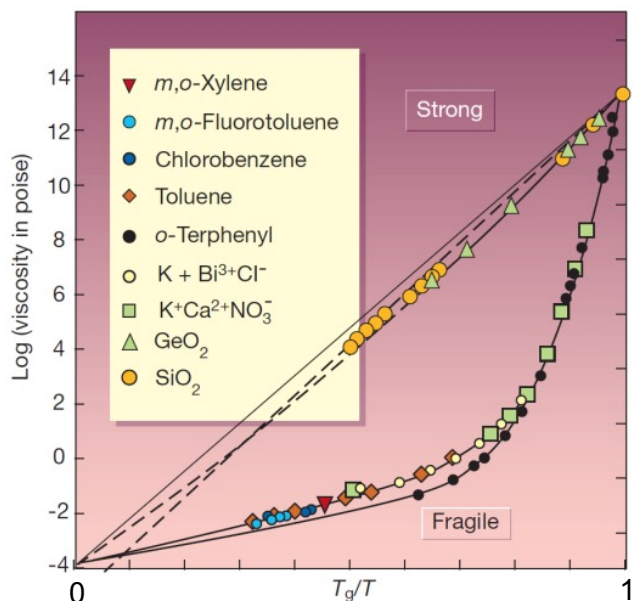


What is the mechanism of glass formation?

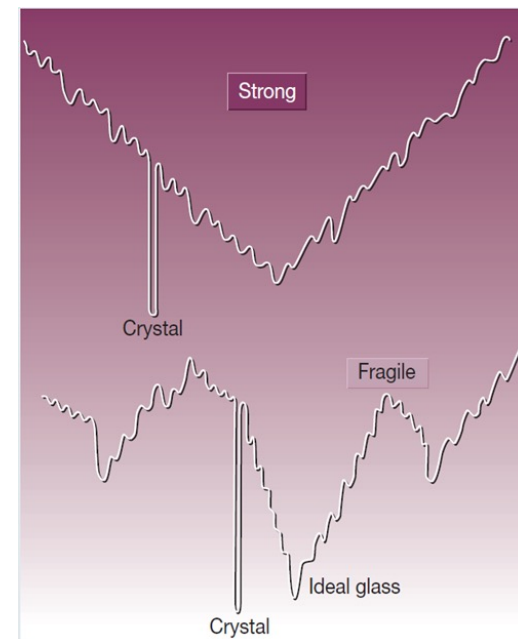
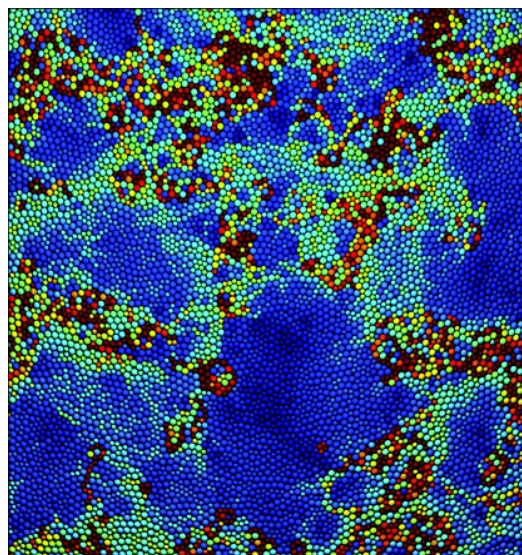
- Science 125 Questions: *What is the nature of the glassy state?* 125th Anniversary Issue of Science (2005)
- The 70 Wonders of Physics: *Is Glass a solid or liquid?* 75th Anniversary Issue of JPSJ (2017)



✓ **Slowing down of dynamics**
Divergence of viscosity



✓ **Dynamical heterogeneity**
Heterogeneity in space and time



Energy landscape

✓ **Energy landscape with multiple local minima**

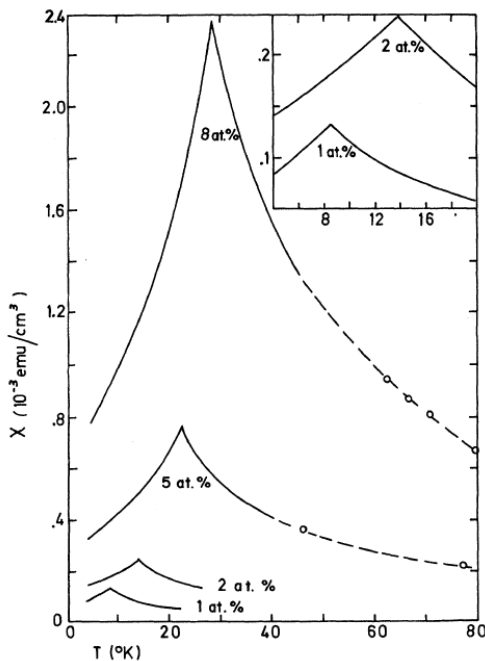
P. G. Debenedetti *et al.*, Nature **410**, 259 (2001).

From the cover of PNAS, Sep. 8th, 2009

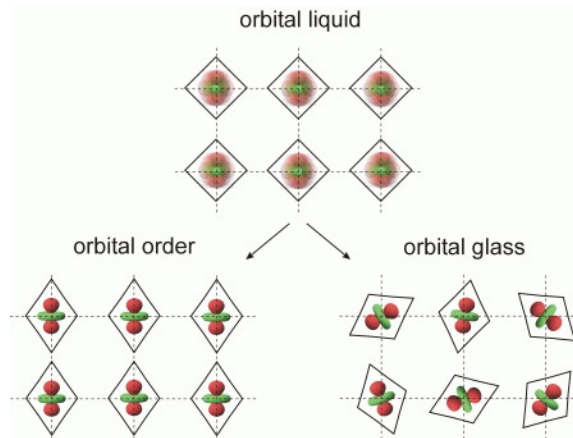
Glass formation phenomena have been reported in various degrees of freedom in strongly correlated electron systems:

spins, charges, orbitals, and superconducting vortices...

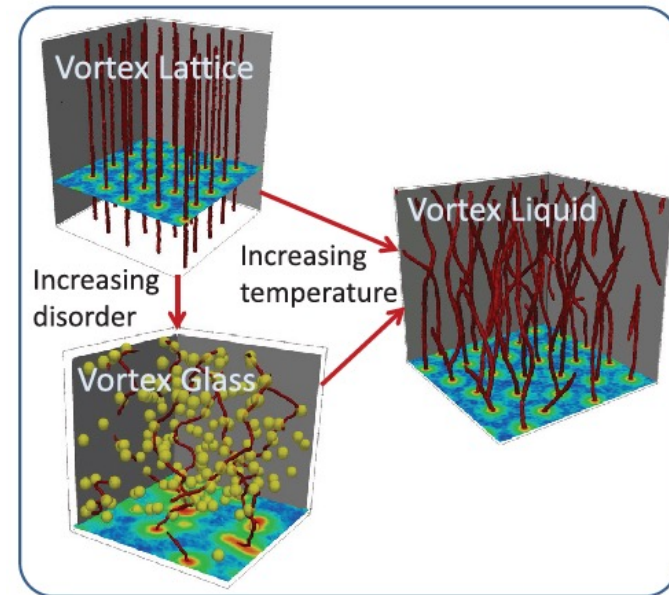
Spin glass



Orbital glass



Superconducting vortex glass



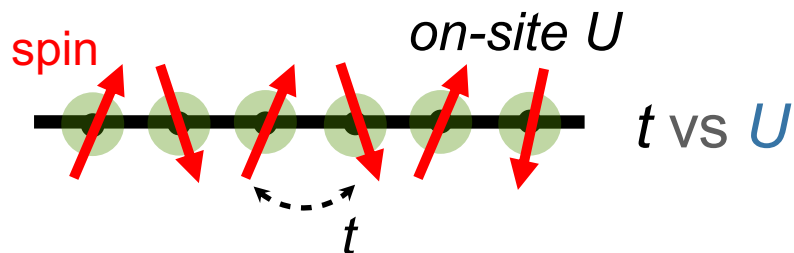
However, the fundamental understanding of the glass formation mechanism is still an open question.

Extended Hubbard model

$$H = t \sum_{\langle i,j \rangle \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + \text{h.c.}) + \underbrace{U \sum_i}_{\text{On-site}} n_{i\uparrow} n_{i\downarrow} + \underbrace{V \sum_{\langle i,j \rangle}}_{\text{Inter-site}} n_i n_j$$

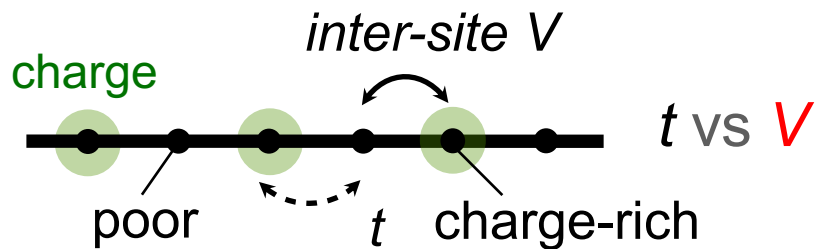
1/2-filled

→ Mott transition



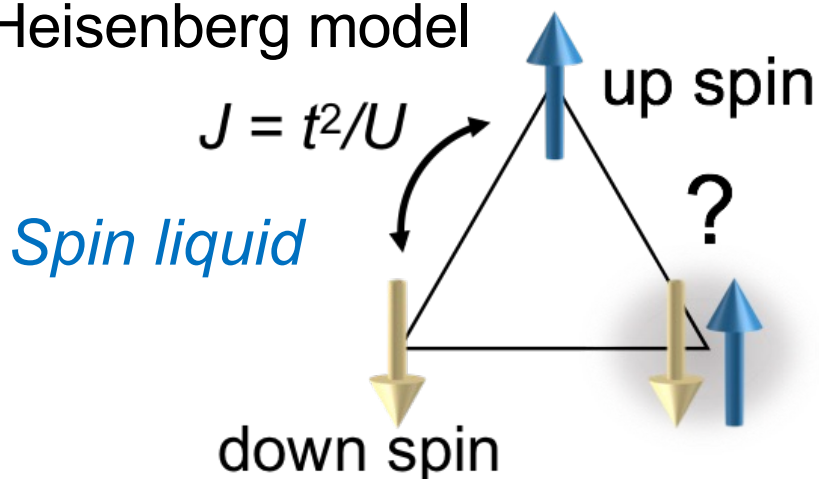
1/4-filled

→ Charge order (CO)



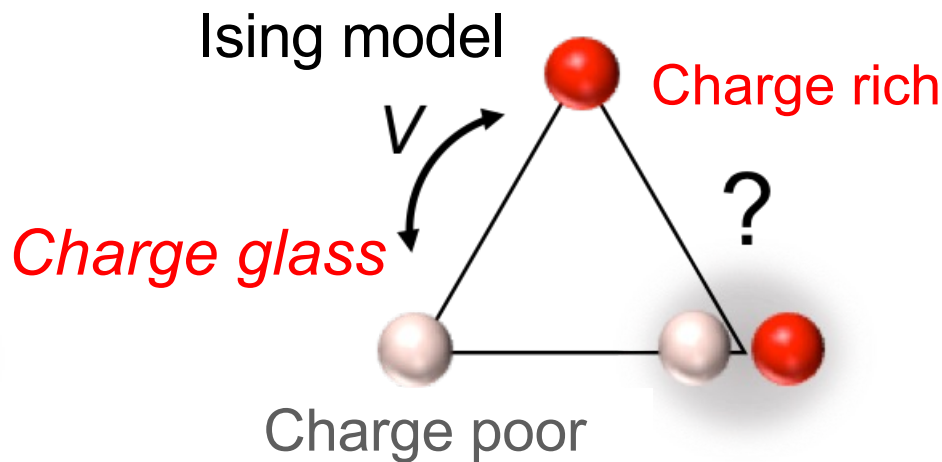
Spin frustration system

Heisenberg model



Charge frustration system

Ising model



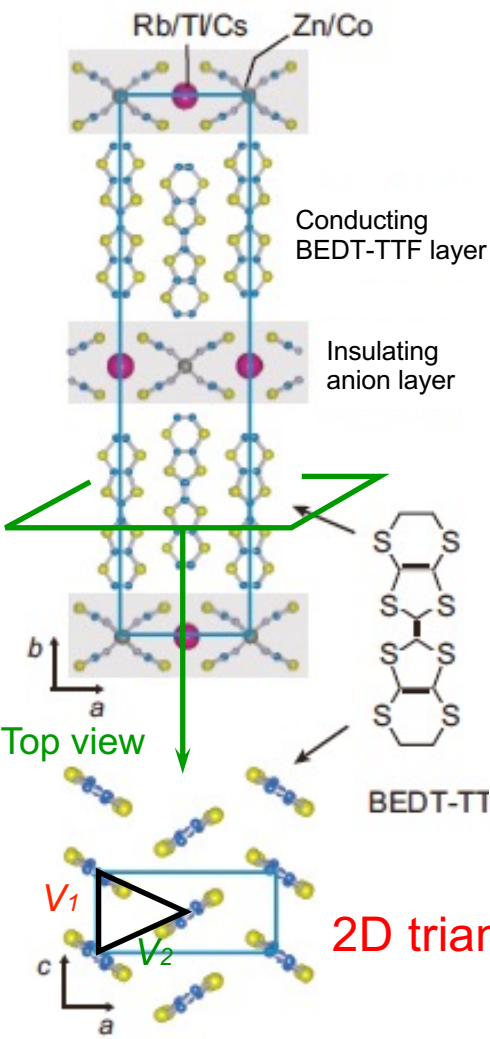
θ -(BEDT-TTF) $_2$ RbZn(SCN) $_4$ (abbreviated as θ -RbZn)

+0.5

-1

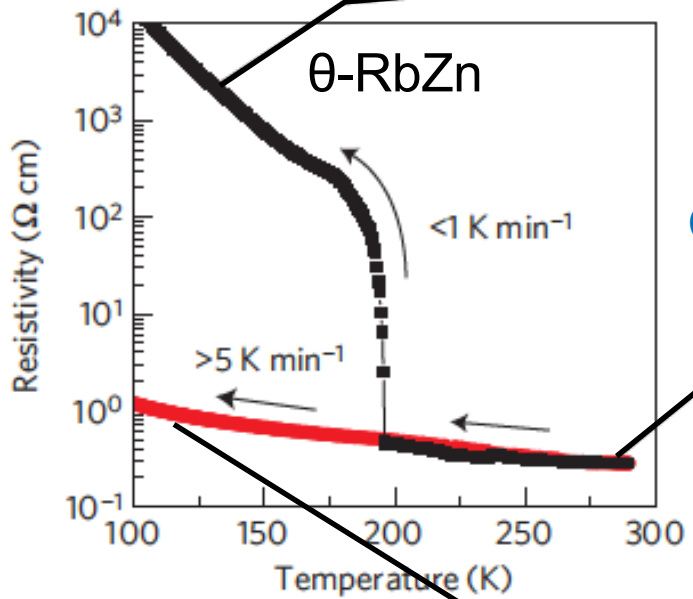
H. Mori *et al.*, PRB **91**, 041101 (1998).

$\frac{1}{4}$ filled hole band system with a 2D triangular lattice



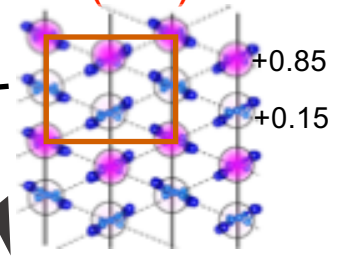
Strong inter-site Coulomb repulsion V

Charge ordering (CO)



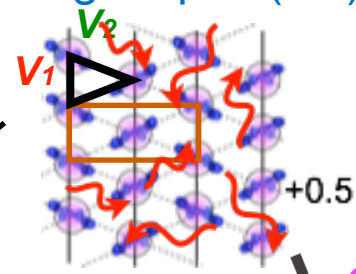
F. Kagawa *et al.*, Nature Phys. **9**, 419 (2013).

Charge crystal (CC)

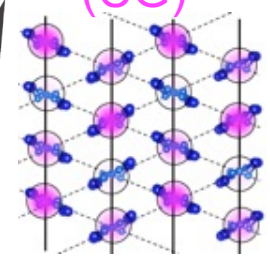


Slow cool

Charge liquid (CL)



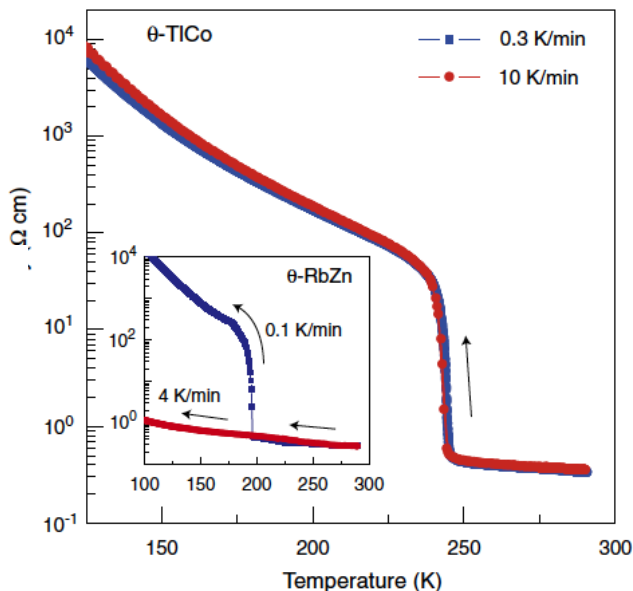
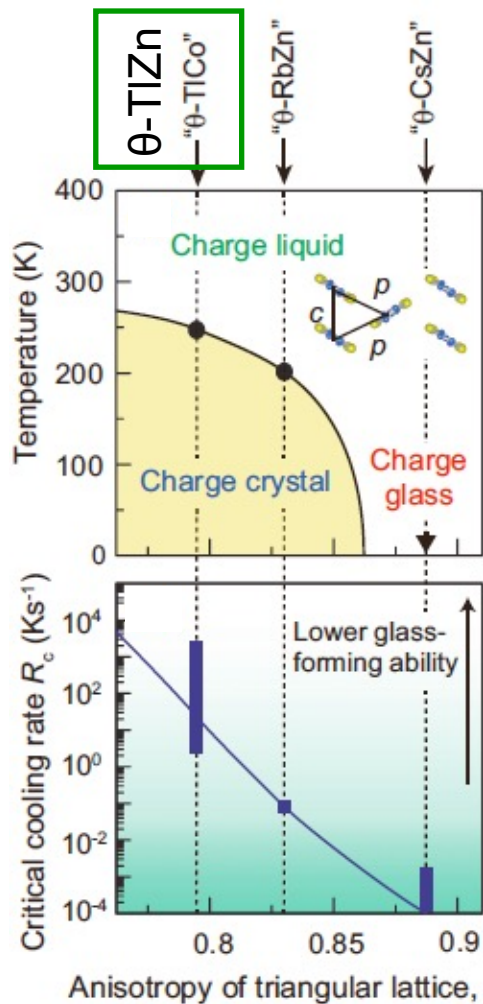
Charge glass (CG)



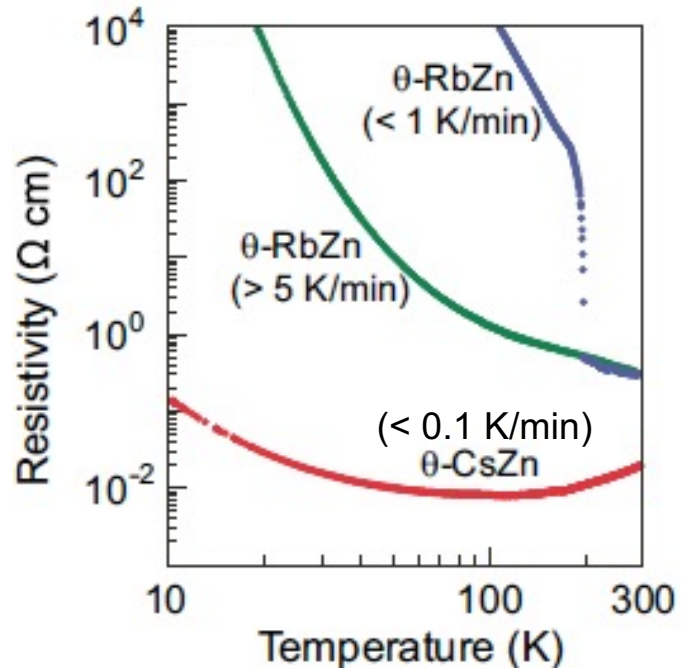
Randomly distributed

H. Oike *et al.*, PRB **91**, 041101(R) (2015).

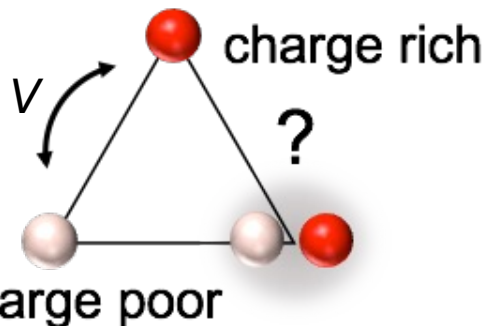
θ -(BEDT-TTF)₂MM'(SCN)₄ (abbreviated as θ -MM')



T. Sato *et al.*, JPSJ (2014).



T. Sato *et al.*, PRB (2014).



Geometrical frustration between V_1 and V_2

\rightarrow Charge glass formation

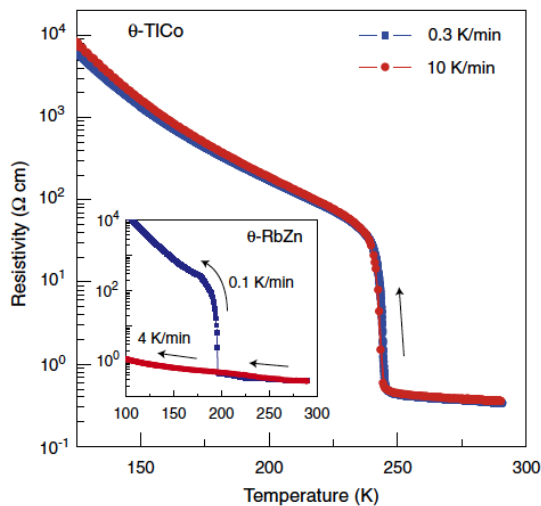
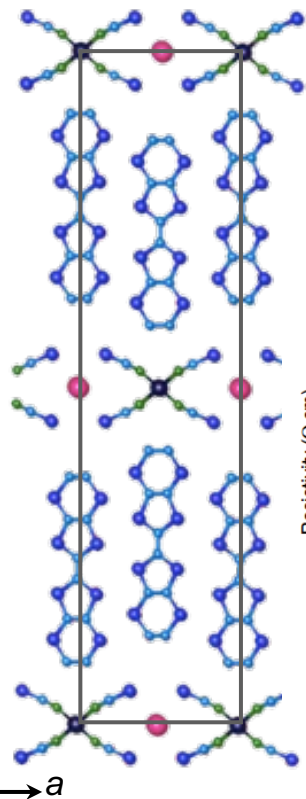
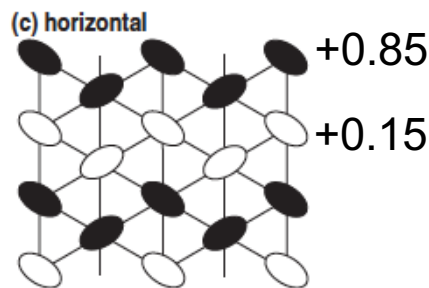


H. Oike *et al.*, PRB **91**, 041101(R) (2015).

Orthorhombic



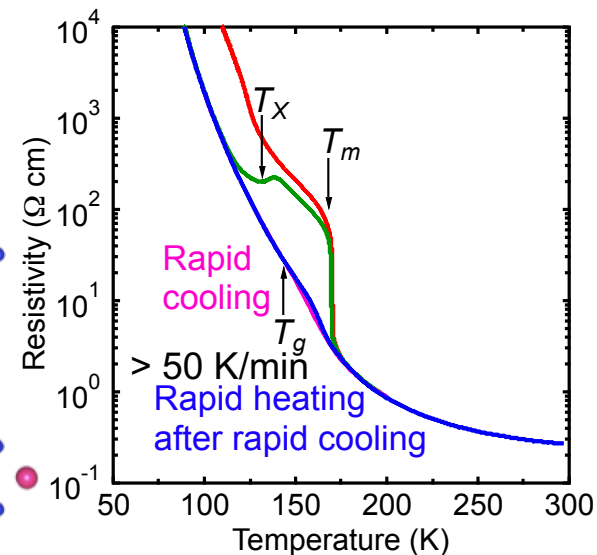
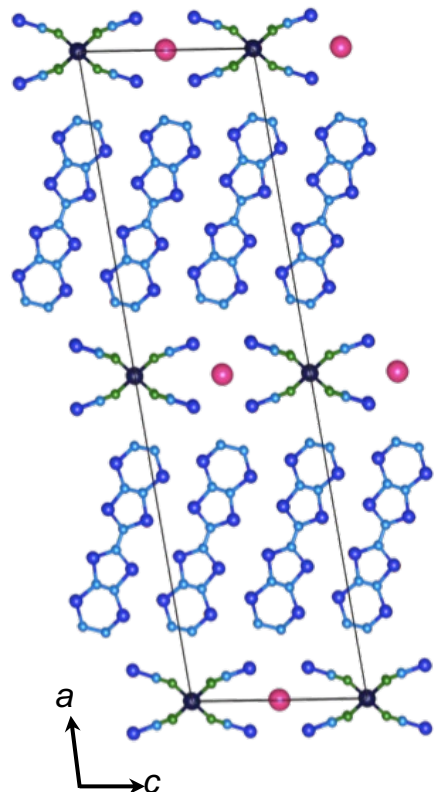
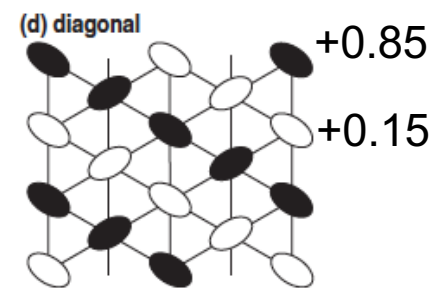
- Horizontal CO
- $T_m = 240$ K
- $V_2/V_1 = 0.835$



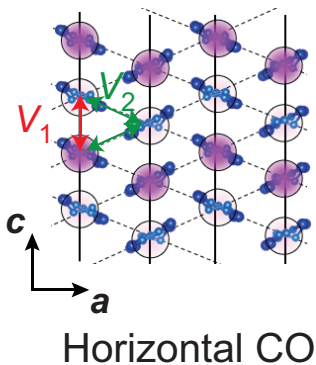
Monoclinic



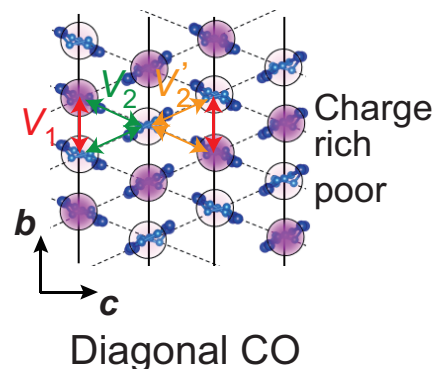
- Diagonal CO
- $T_m = 170$ K
- $V_2/V_1 = 0.816$



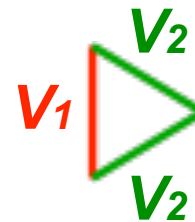
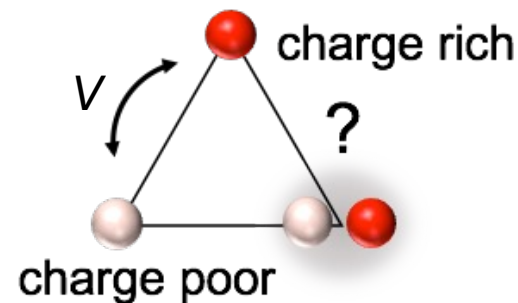
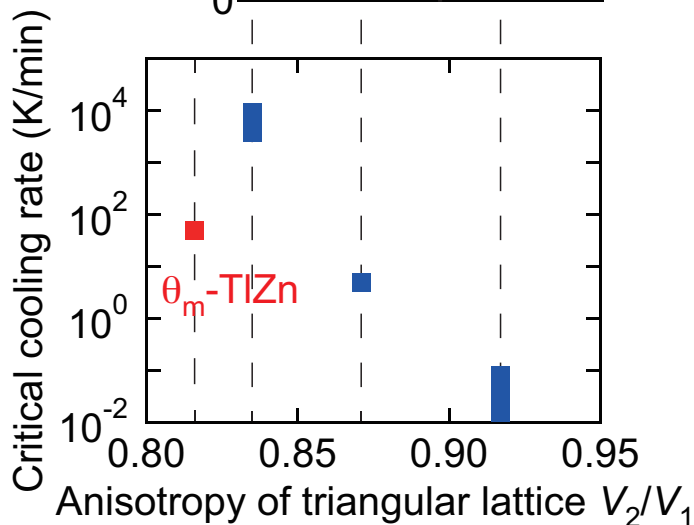
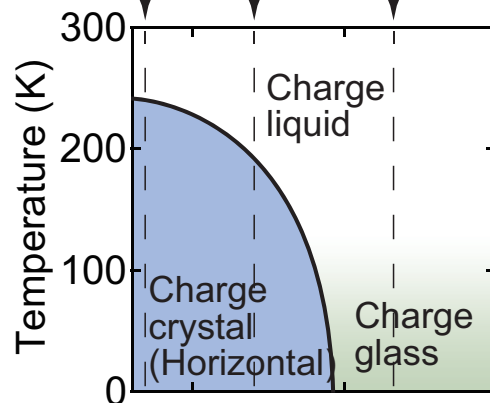
θ_o -type



θ_m -type



θ_o -TiZn/TiCo θ_o -RbZn θ_o -CsZn



V_1 vs V_2

Geometrical frustration of equilateral triangular lattice

The critical cooling rate of θ_m -TiZn is slower than that of θ_o -TiCo, suggesting different charge-glass formation mechanisms between these two systems.

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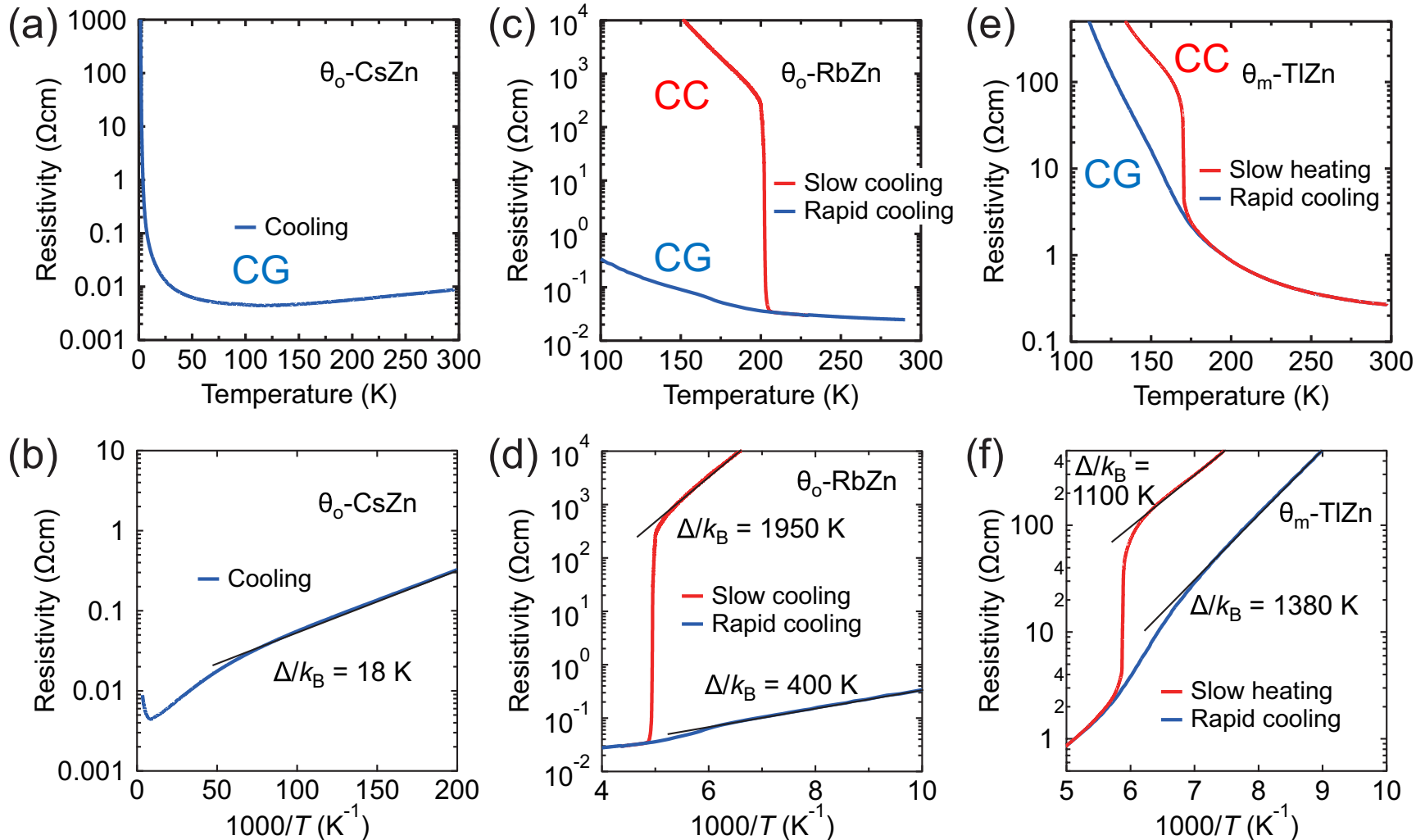
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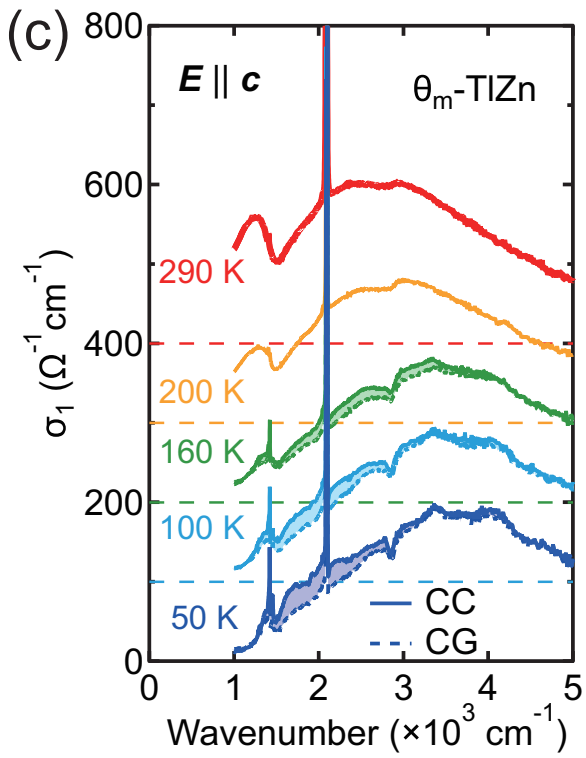
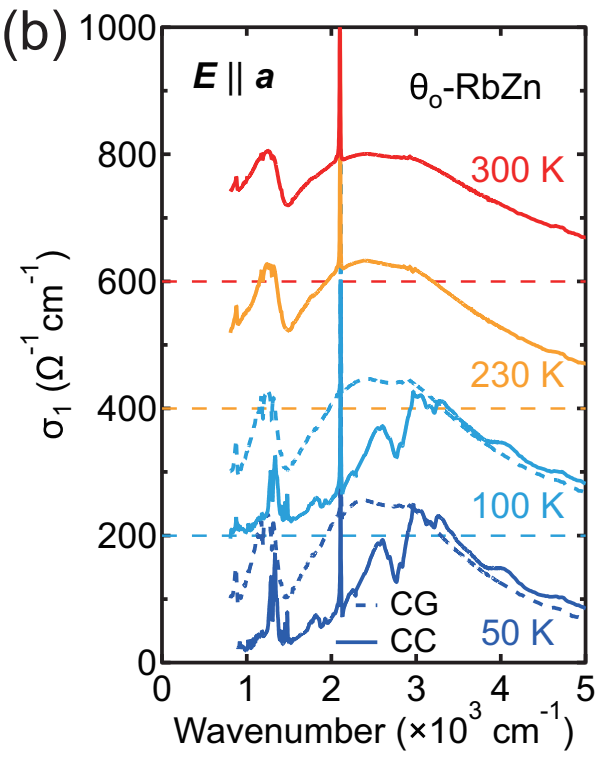
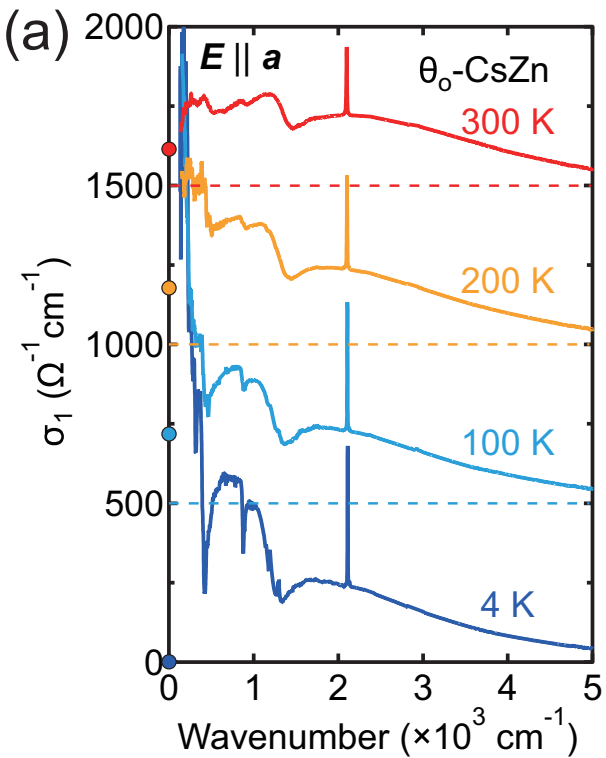
5. Summary

KH *et al.*, submitted.

θ_o -type: large difference between CC and CG states

θ_m -type: small difference between CC and CG states

KH *et al.*, submitted.

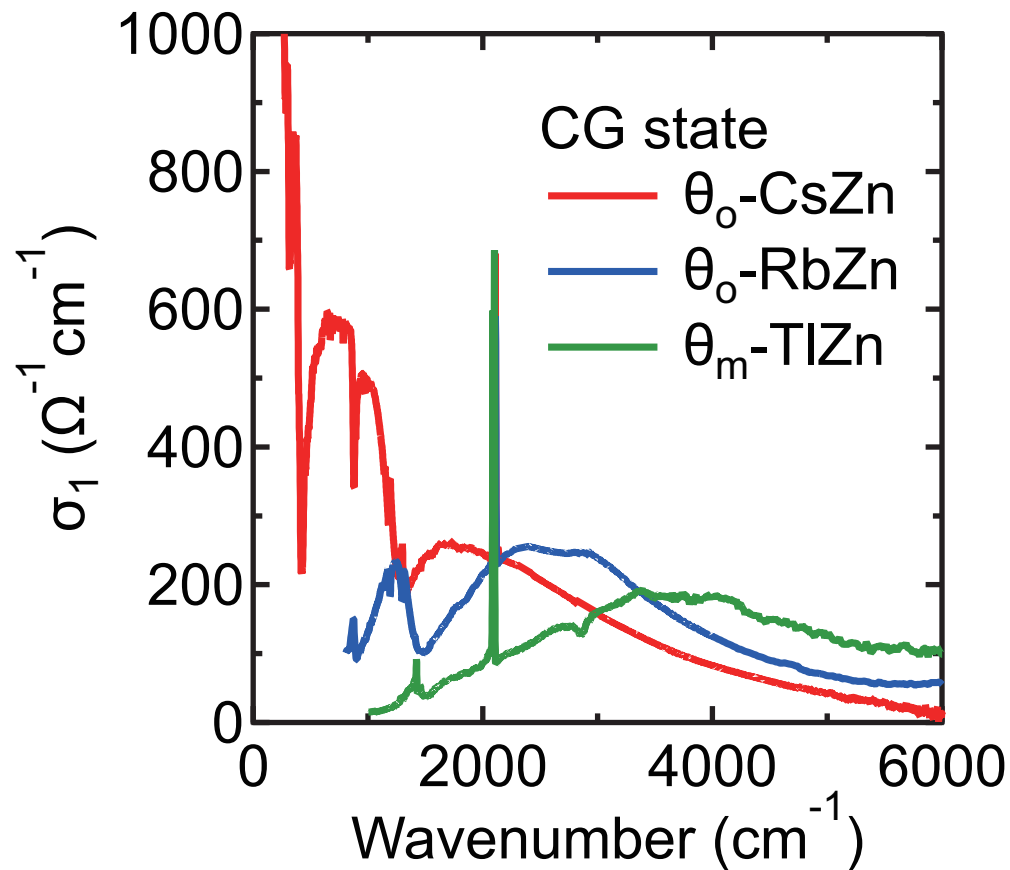


θ_o -type

The CG state exhibits larger low-energy excitations than the CC state and becomes more metallic with increasing V_1/V_2 .

θ_m -type

The optical gaps in the CG and CC states are almost identical.



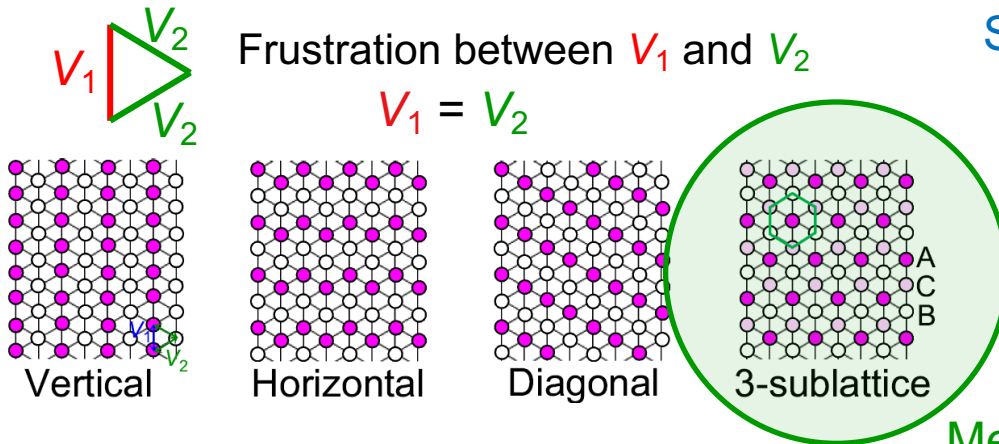
KH et al., submitted.

As the anisotropy of the triangular lattice increases, the spectral weight in the low-energy region shifts to a higher-energy region.

Extended Hubbard model (EHM)

$$H = t \sum_{\langle i,j \rangle \sigma} (c_{i\sigma}^\dagger c_{j\sigma} + \text{h.c.}) + U \sum_i n_{i\uparrow} n_{i\downarrow} + V \sum_{\langle i,j \rangle} n_i n_j$$

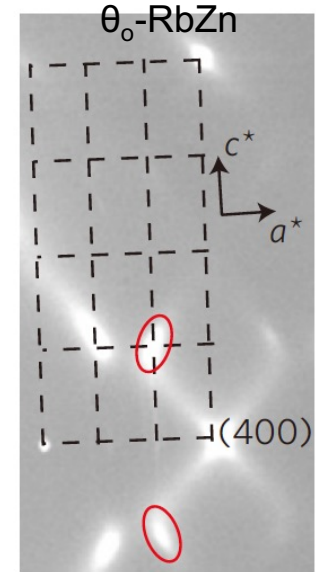
θ_o -type: equilateral triangular lattice



Short-range 3-fold COs

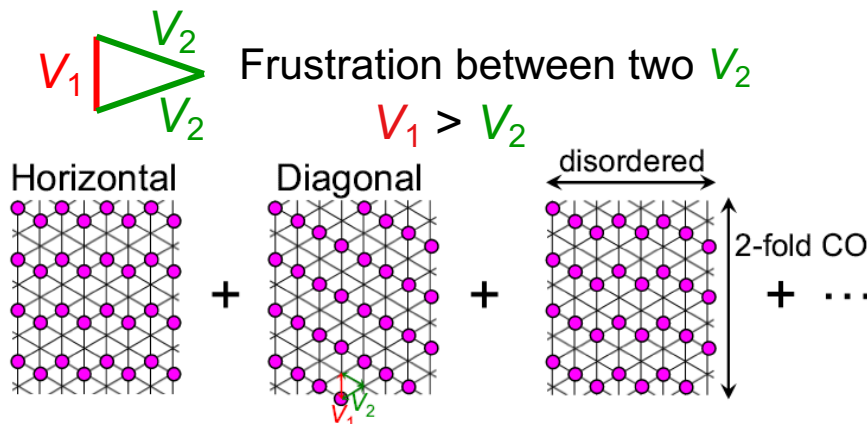
θ_o -CsZn: $q_d = (2/3 \ 1/3)$

θ_o -RbZn: $q_d = (1/4 \ 1/3)$



F. Kagawa *et al.*,
Nature Phys. **9**, 419 (2013).

θ_m -type: isosceles triangle lattice

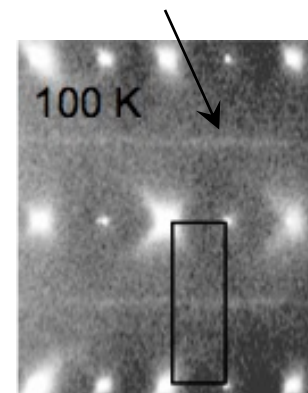


Chain-striped COs along the V_1 direction

Superpositions of insulating 2-fold stripe COs

θ_m -TiZn: $q_d = (1/2 \ l)$

Diffuse lines



S. Sasaki, KH *et al.*, Science **357**, 1381 (2017).

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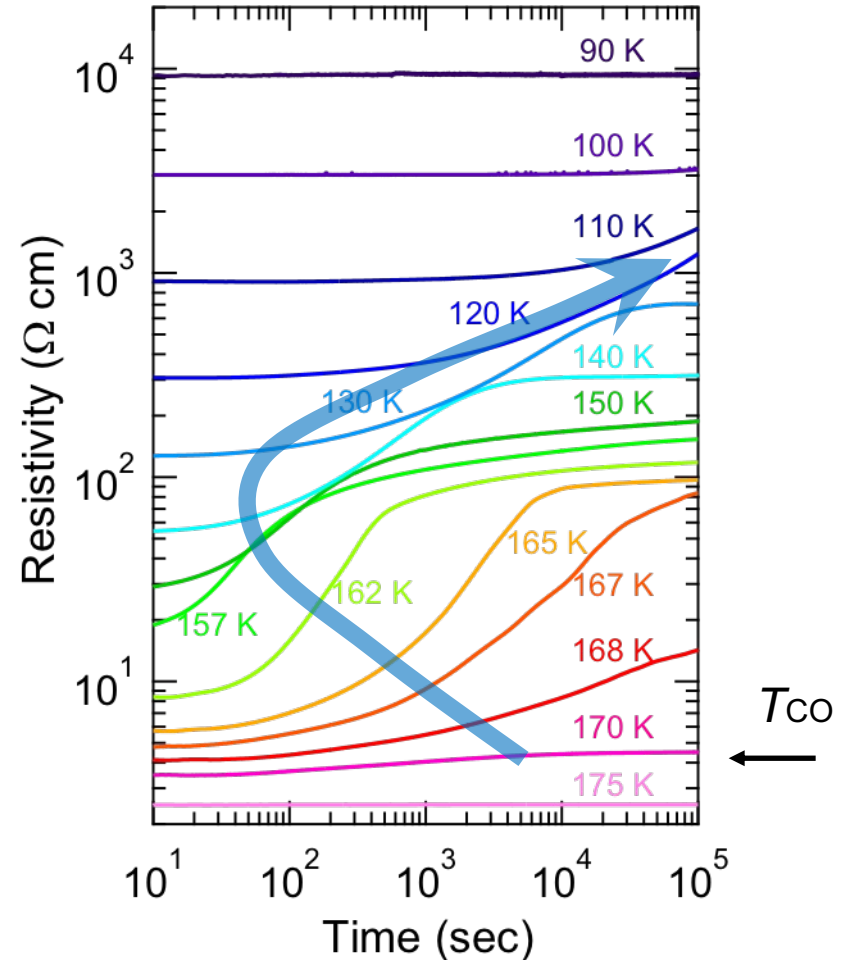
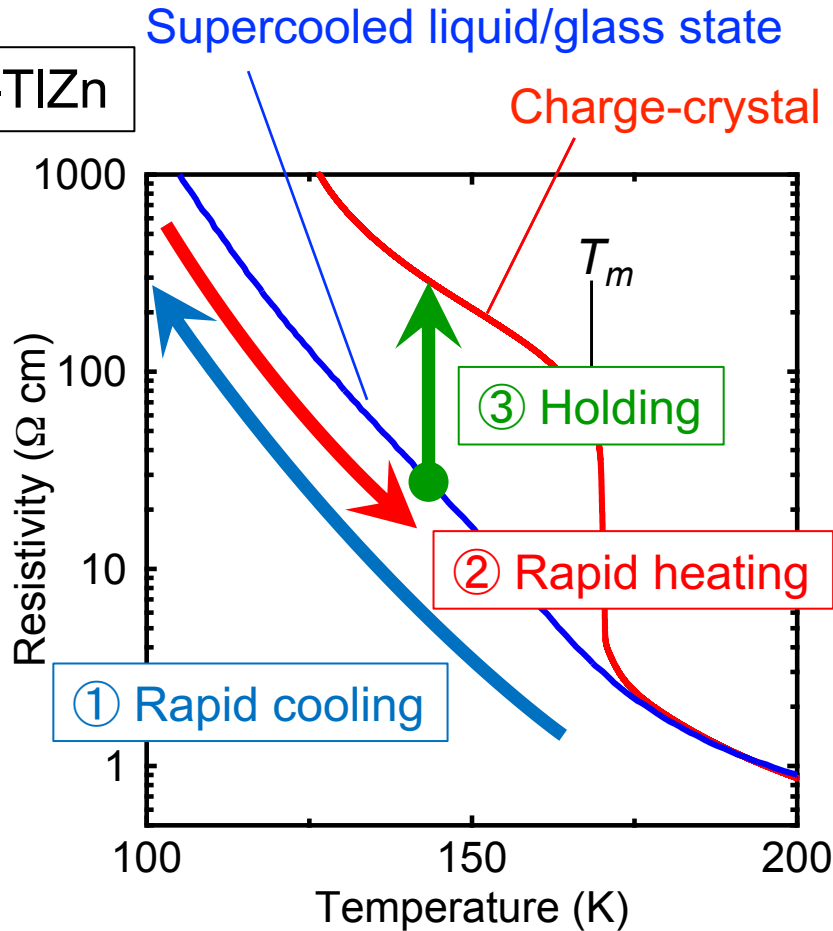
- TTT diagram constructed by DSC measurements

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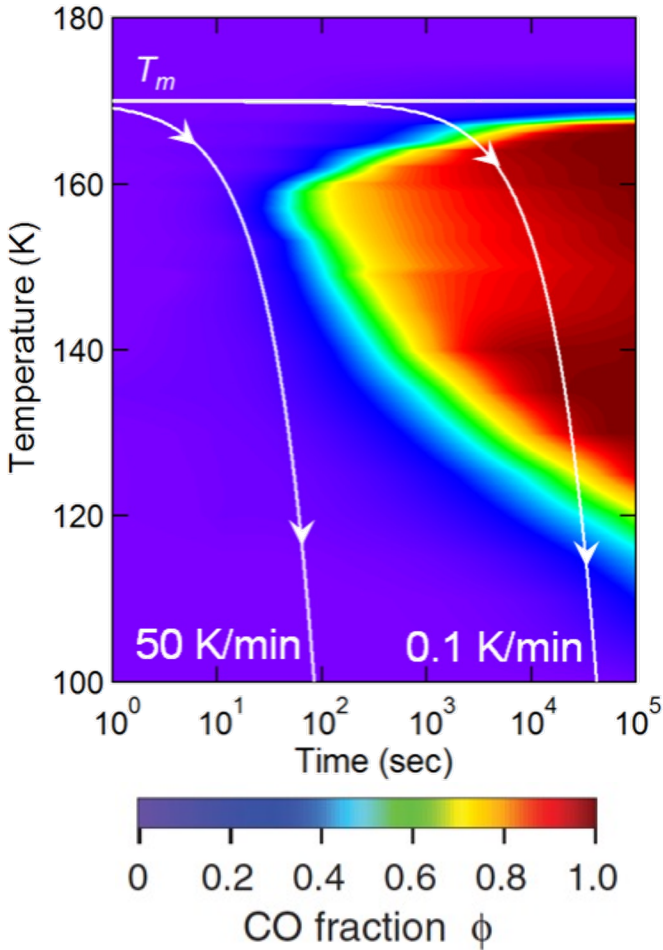
S. Sasaki, KH *et al.*,
 Science **357**, 1381 (2017).



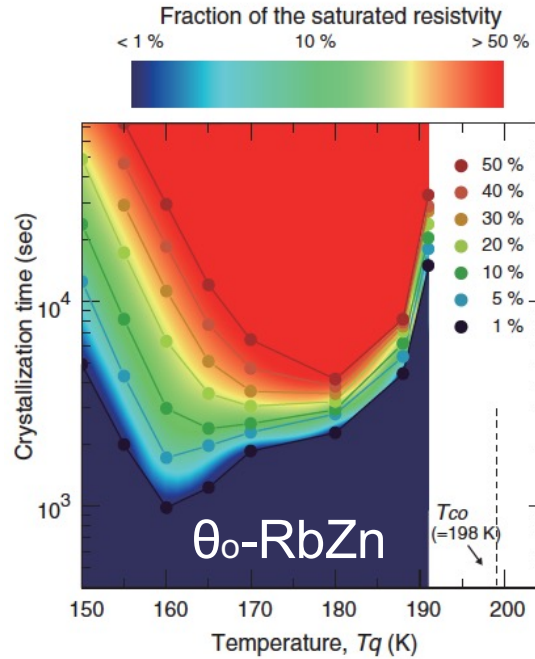
- ① Rapid cooling (> 50 K/min)
- ② Rapid heating (> 50 K/min)
- ③ Holding

The crystallization time becomes faster with decreasing temperature, and then, it becomes slower below ~ 160 K.

θ_m -TiZn Dome-shaped

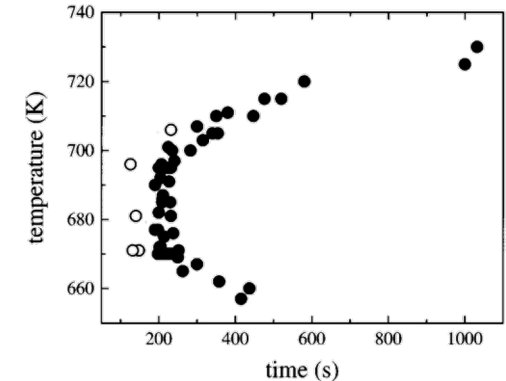


S. Sasaki, KH *et al.*,
 Science **357**, 1381 (2017).



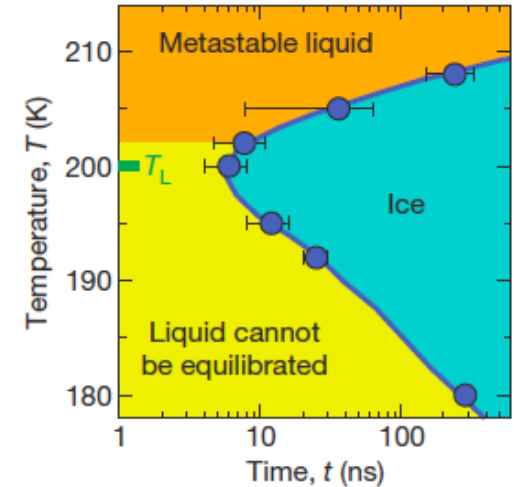
T. Sato *et al.*,
 Science **357**, 1378 (2017).

Metallic glass



J. Schroers, *et al.*, APL **77**, 1158 (2000).

H₂O(simulation)



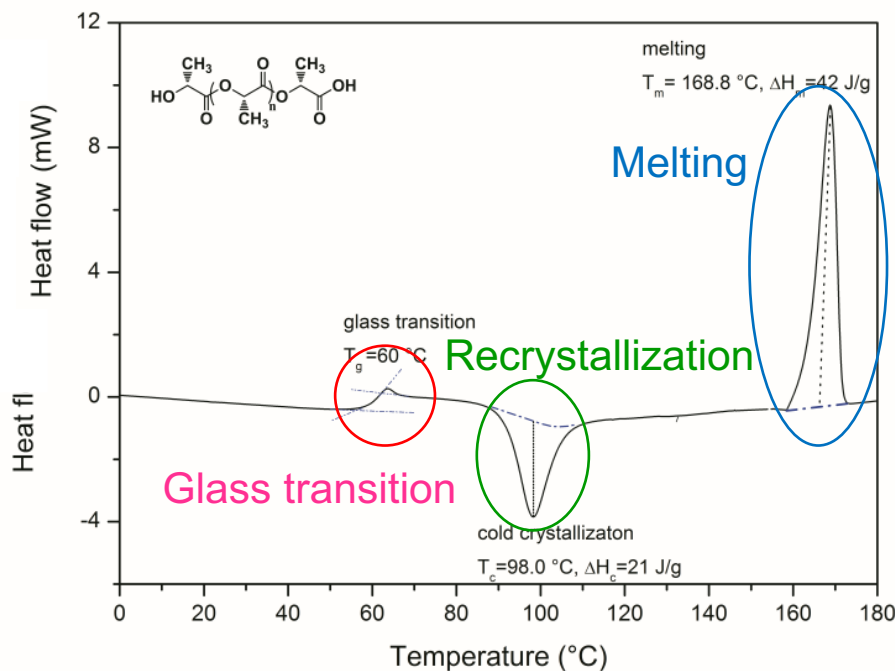
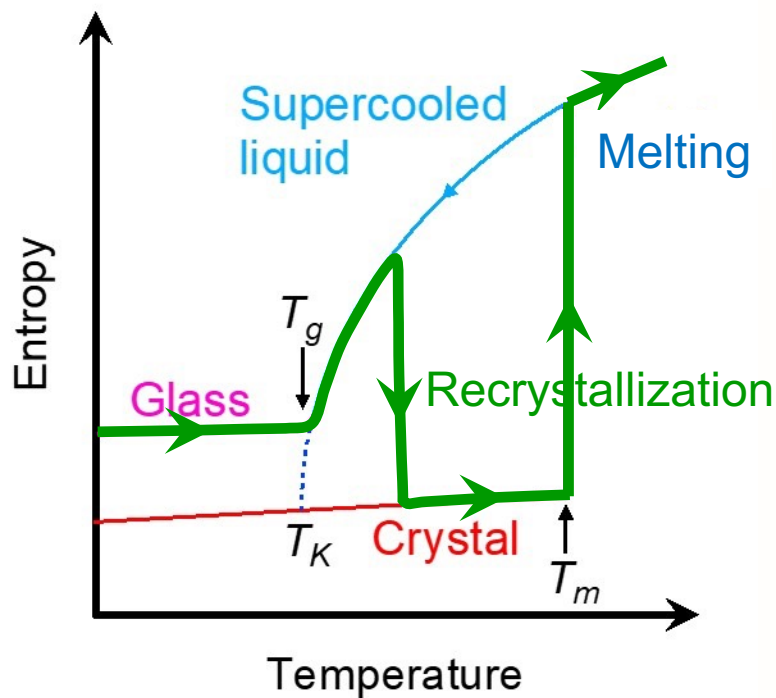
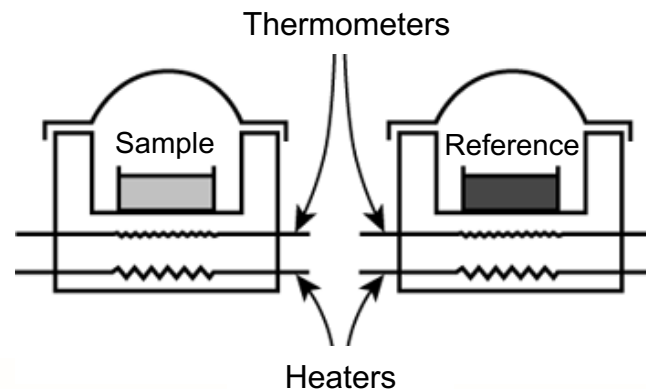
E. B. Moore, *et al.*, Nature **479**, 506 (2011).

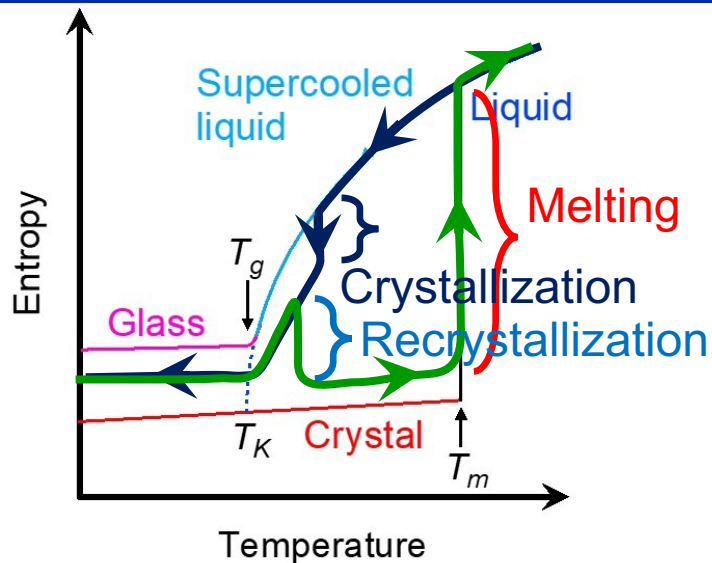
The dome-shaped TTT diagram is universally observed in other glass-forming liquids such as metallic glasses and water as well as θ_o -RbZn.

DSC is a very powerful probe to observe crystallization, melting, and glass transitions.

$$\Delta S(T) = \left(\int \frac{\Delta \dot{q}(T) dt}{T} \right)$$

Heat flow



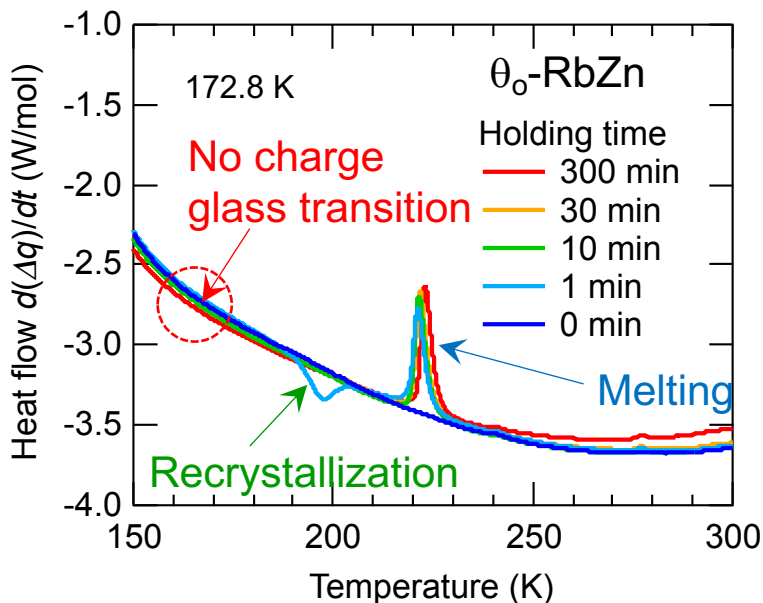


DSC-8500 (Perkin Elmer Corp.) at IMR, Tohoku Univ.

- N₂ liquid (Lowest temperature: 135 K)
- He gas atmosphere
- Cooling/heating rate: 100 K/min

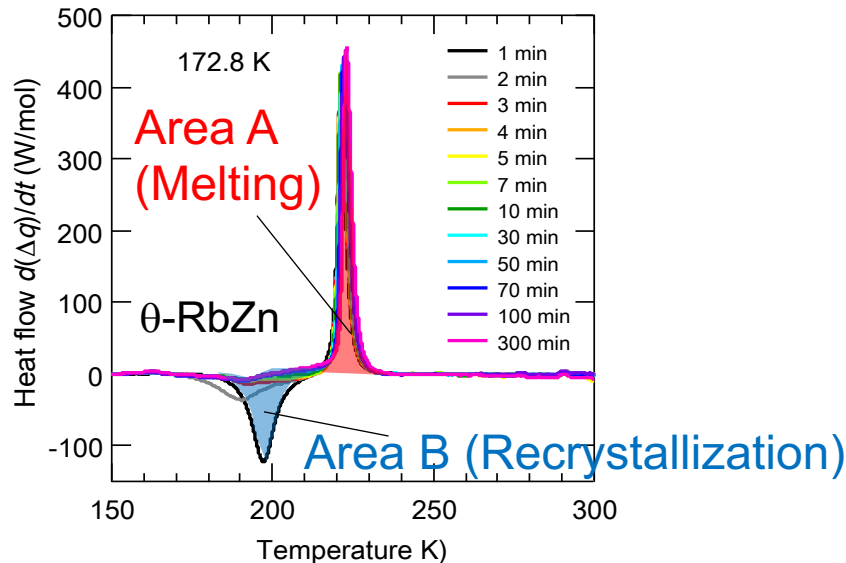
$$\Delta H(T) = \int \Delta \dot{q}(T) dt$$

$$\Delta S = \Delta H/T$$



- No clear charge glass transition
- Clear charge recrystallization and melting

Entropy change = $\frac{\text{Area A}}{\text{Crystallization}}$ - $\frac{\text{Area B}}{\text{Melting Recrystallization}}$



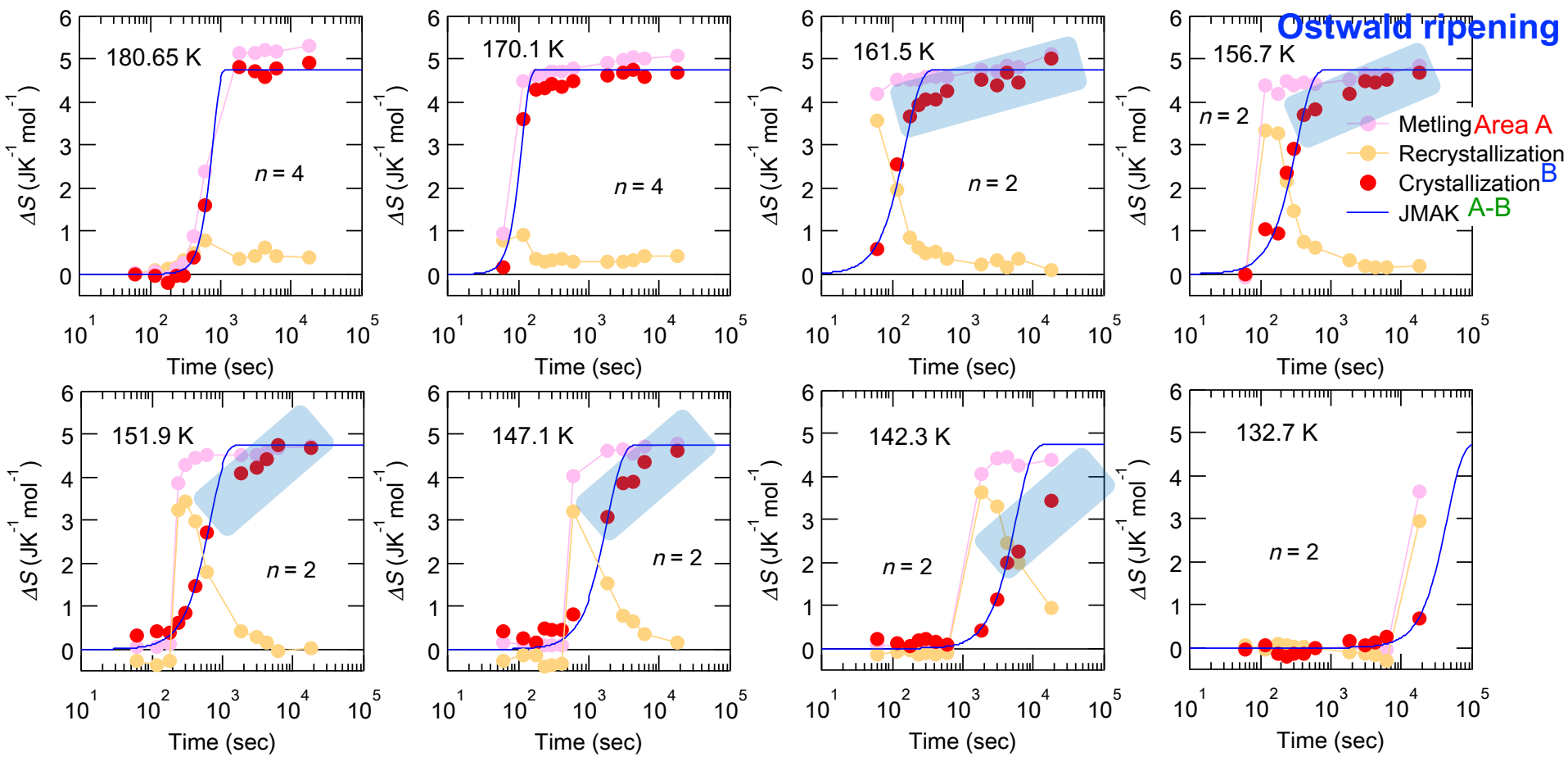
Johnson–Mehl–Avrami–Kolmogorov (JMAK) model

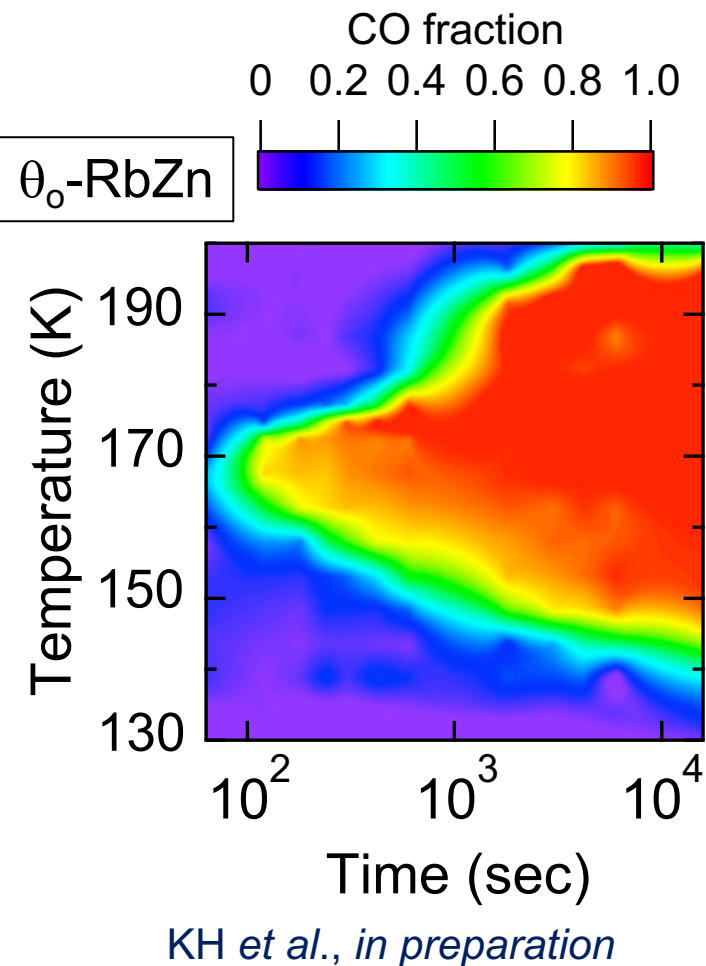
$$\phi(t) = 1 - \exp(-bt^n) \quad 3D: n = 3-4, 2D: n = 2-3, 1D: n = 1-2$$

- At high T, 3D crystal growth ($n = 4$)
- At low T, 1-2D crystal growth ($n = 2$)
→ Ostwald ripening

Consistent with the results of NMR

T. Sato *et al.*, Science **357**, 1378 (2017).





JMAK model

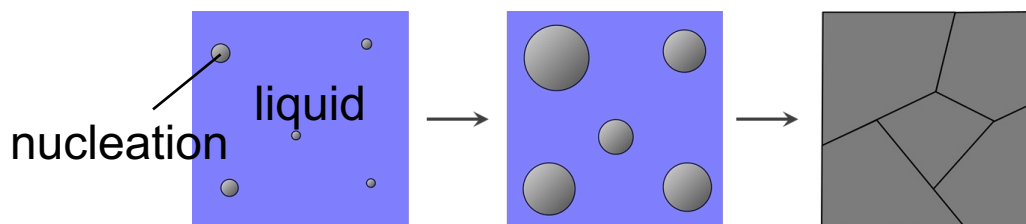
$$\phi(t) = 1 - \exp(-bt^n)$$

3D: $n = 3-4$

2D: $n = 2-3$

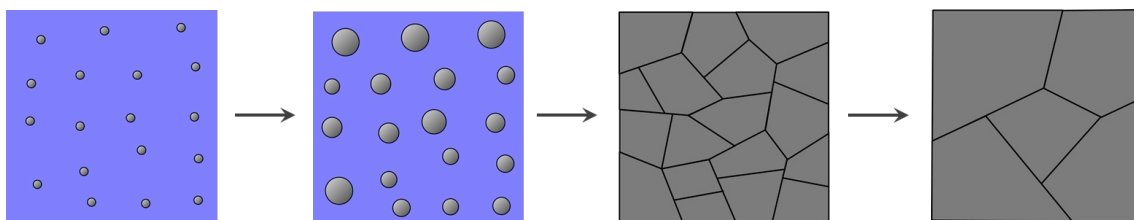
1D: $n = 1-2$

High T : Nucleus growth is dominant



$n = 4$ 3D growth

Low T : Nucleation is dominant



$n = 2$

1 or 2D growth + Ostwald ripening

Dome shaped TTT diagram constructed by thermodynamic quantity

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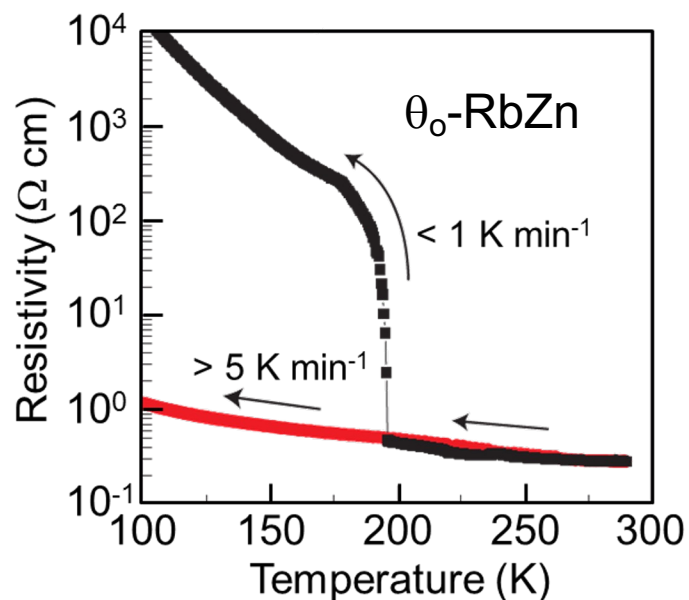
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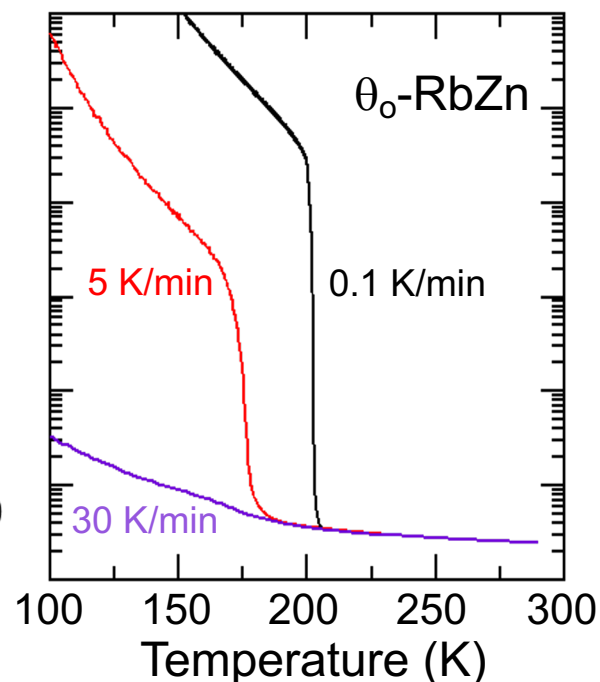
5. Summary

Large sample dependence in the charge glass forming ability

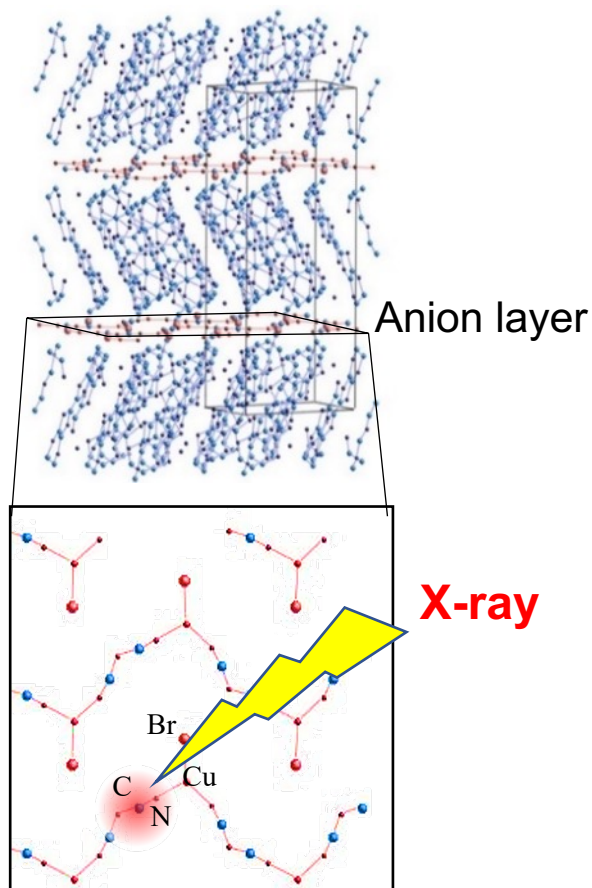
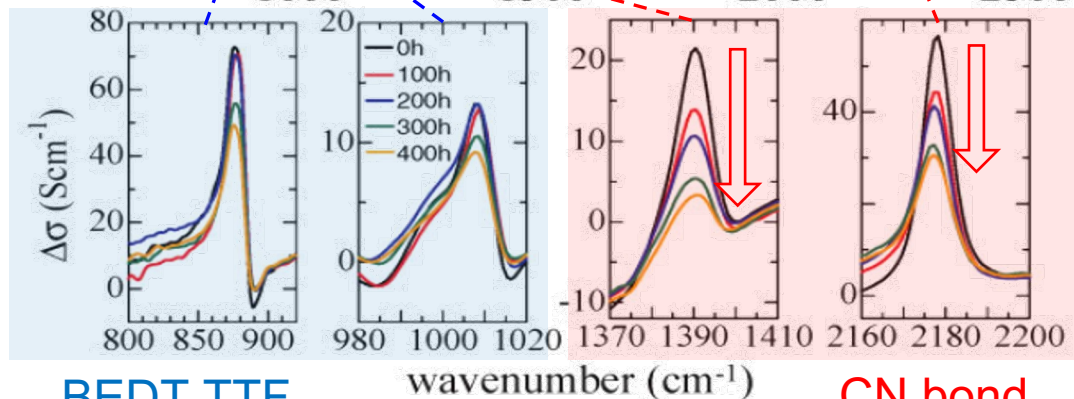
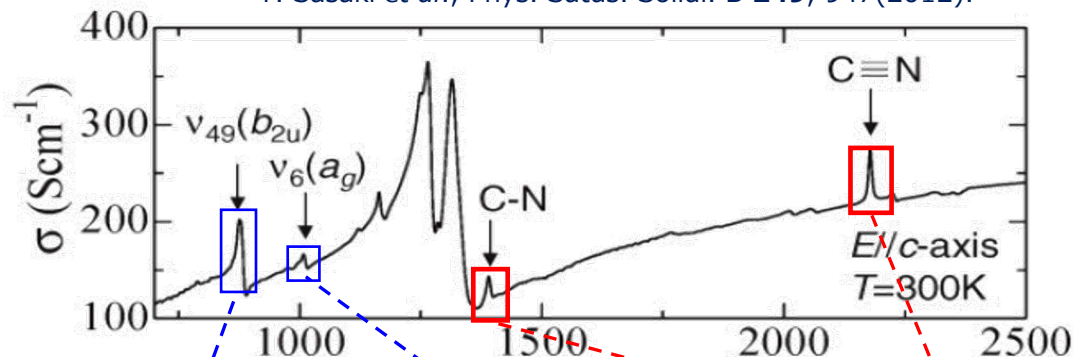
F. Kagawa *et al.*,
Nature Physics **9**, 2642 (2013)



Our crystals (from Yamamoto Group
in the Institute for Molecular Science)



We examined how the randomness effect affects the charge-glass forming ability by using molecular defects introduced by x-ray irradiation.

Infrared molecular vibrations in $\kappa\text{-(ET)}_2\text{Cu}[\text{N}(\text{CN})_2]\text{Br}$ T. Sasaki et al., Phys. Status Solidi. B **249**, 947(2012).

Defects are mainly introduced within the anion layers (CN bonds).

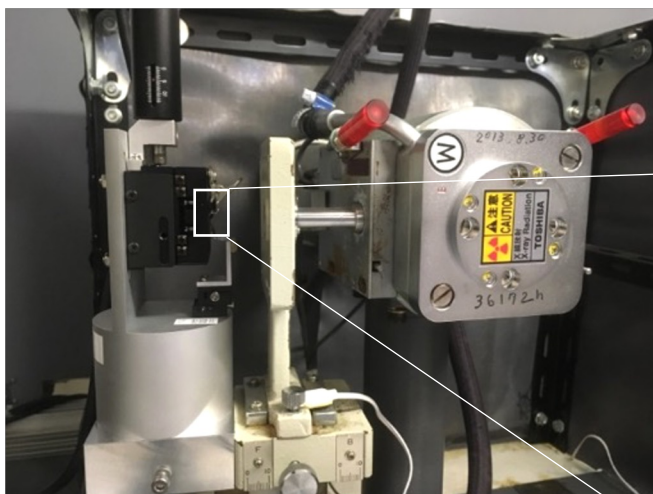
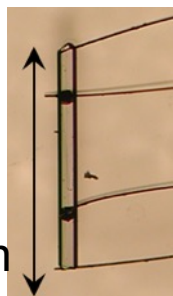


X-ray irradiation does not largely affect the physical properties of π electrons such as charge carriers.

- θ -(ET)₂RbZn(SCN)₄ single crystals
 - Electrochemical redox method
 - Provided by S. Suda and H. Yamamoto

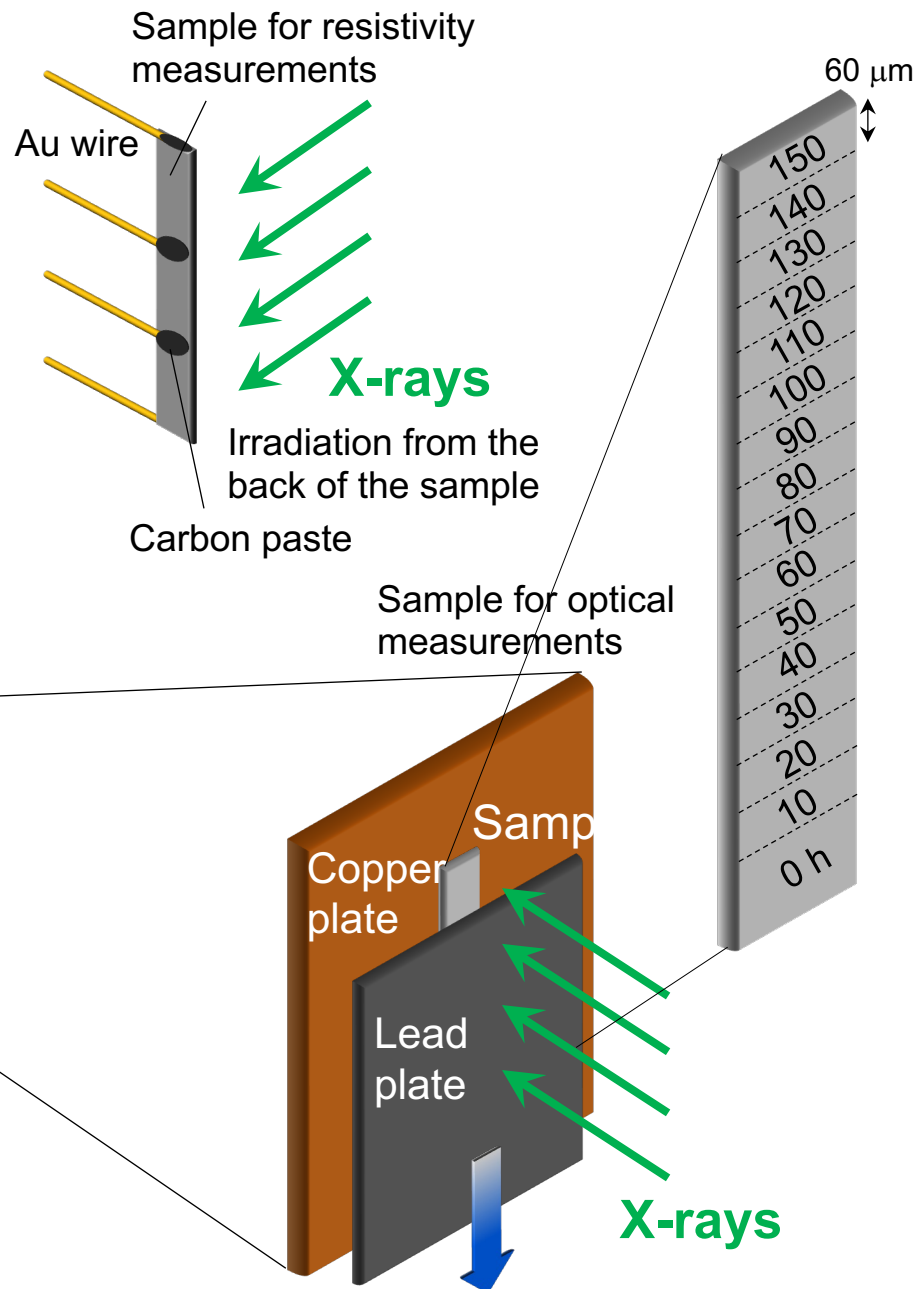
• X-ray irradiation

- Room temperature
- Tungsten tube (non-filtered, 40 kV, 20 mA)
- White x-rays

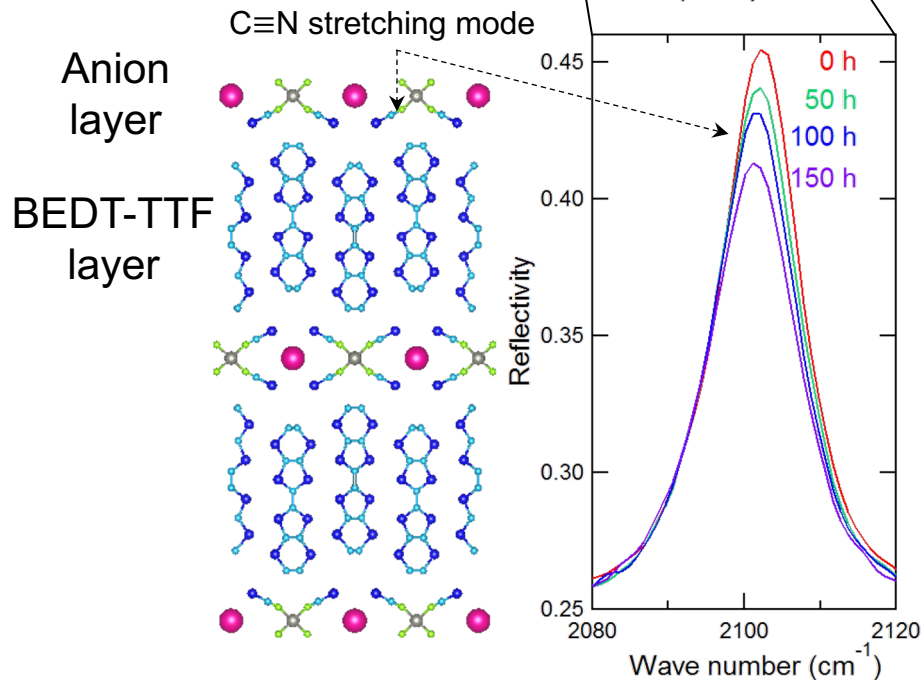
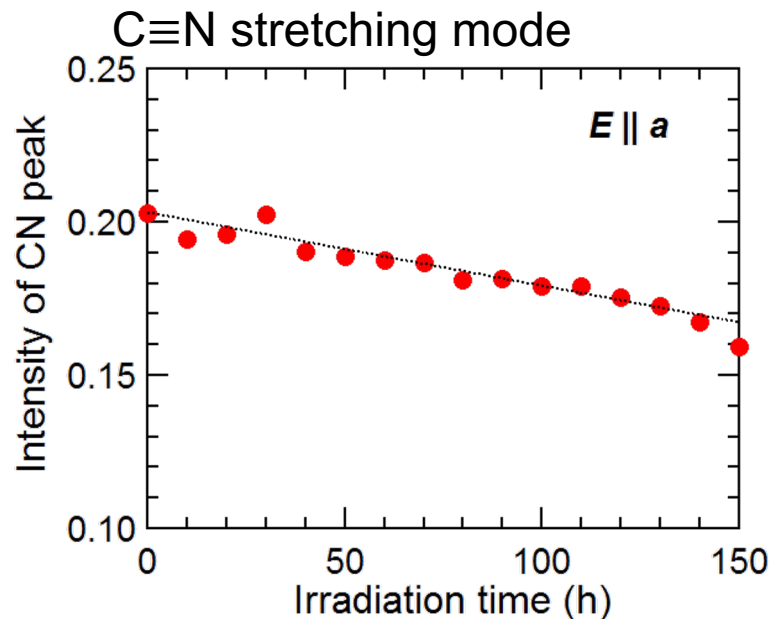
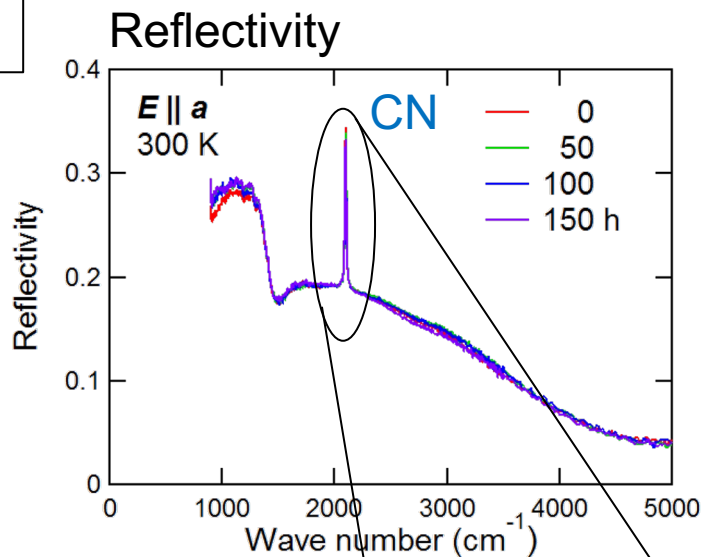


• Experiments

- Resistivity
- Optical conductivity
- Differential scanning calorimetry (DSC)



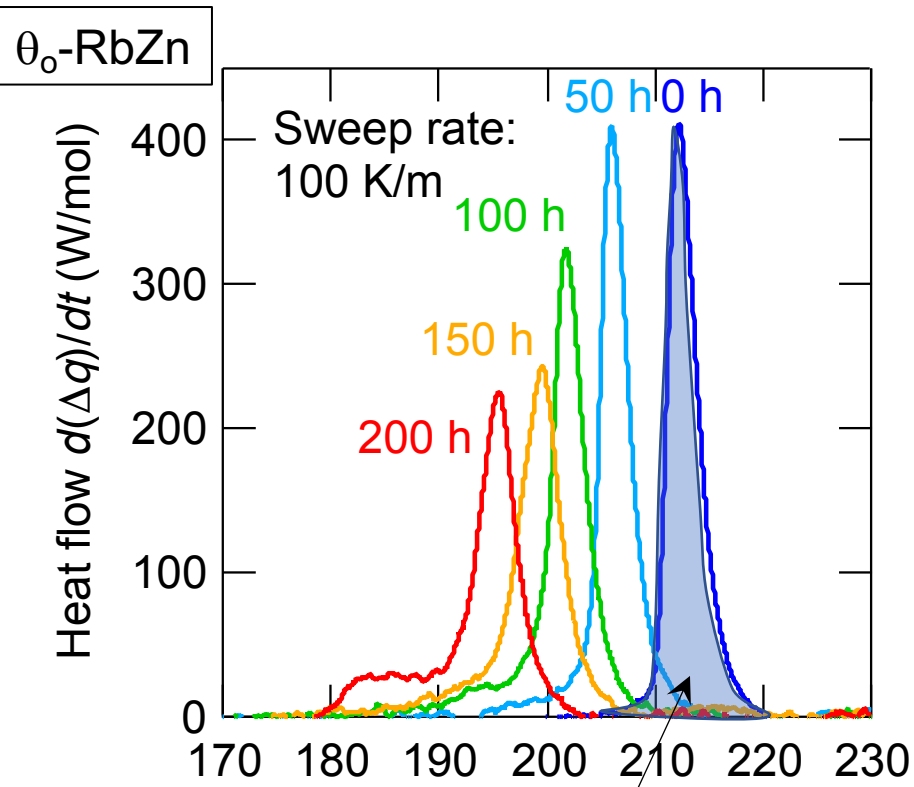
θ_0 -RbZn



Reduction by 15% after 150 h irradiation

Defects are mainly introduced within the anion layers (CN bonds).

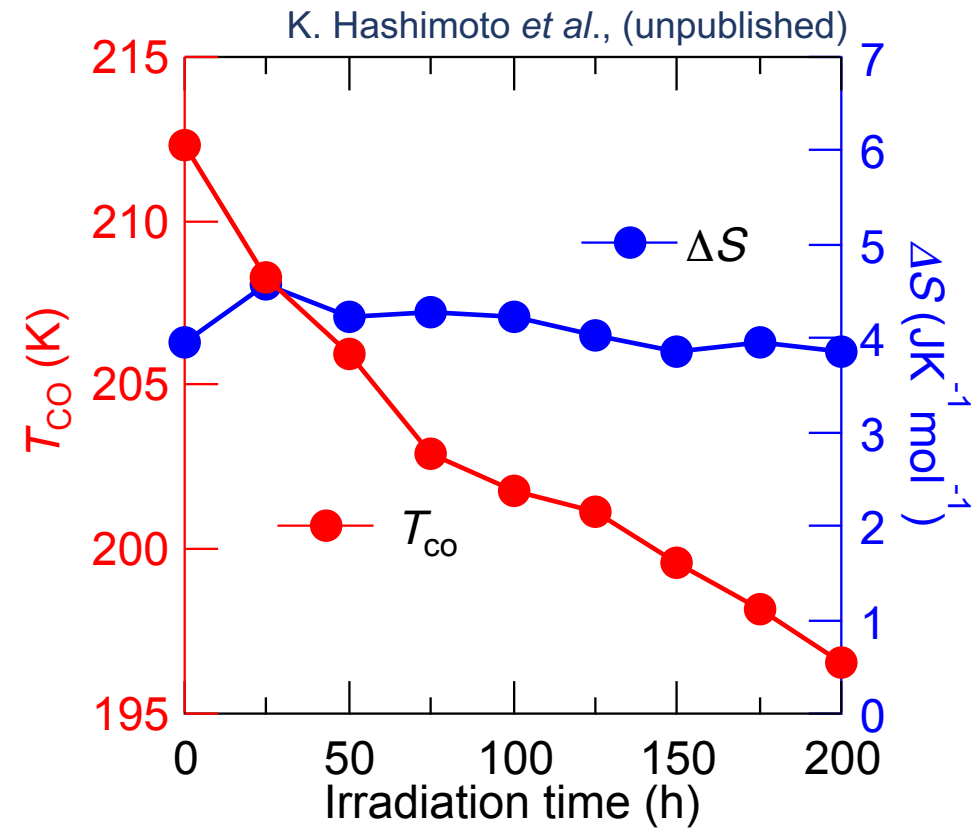
- T_{CO} is systematically suppressed by x-ray irradiation.
- Almost no change in the total entropy at the CO transition



$$\Delta H(T) = \int \Delta \dot{q}(T) dt$$

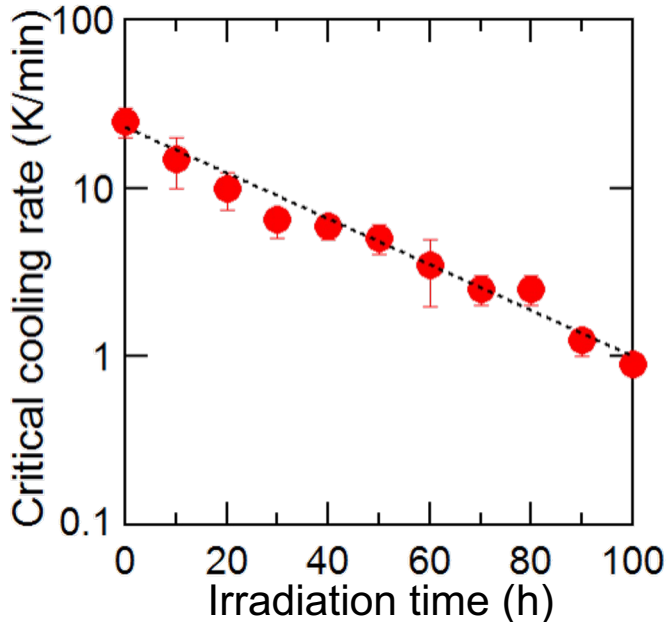
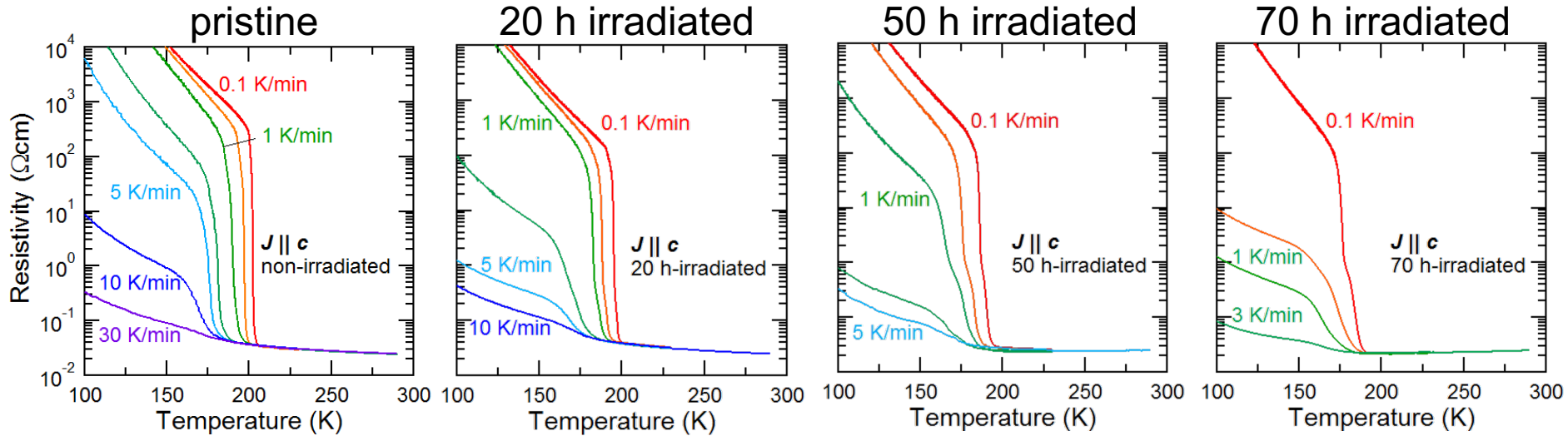
$$\Delta S = \Delta H / T$$

The entropy change at CO is obtained by integrating the heat flow over time and divided by T_{CO} .



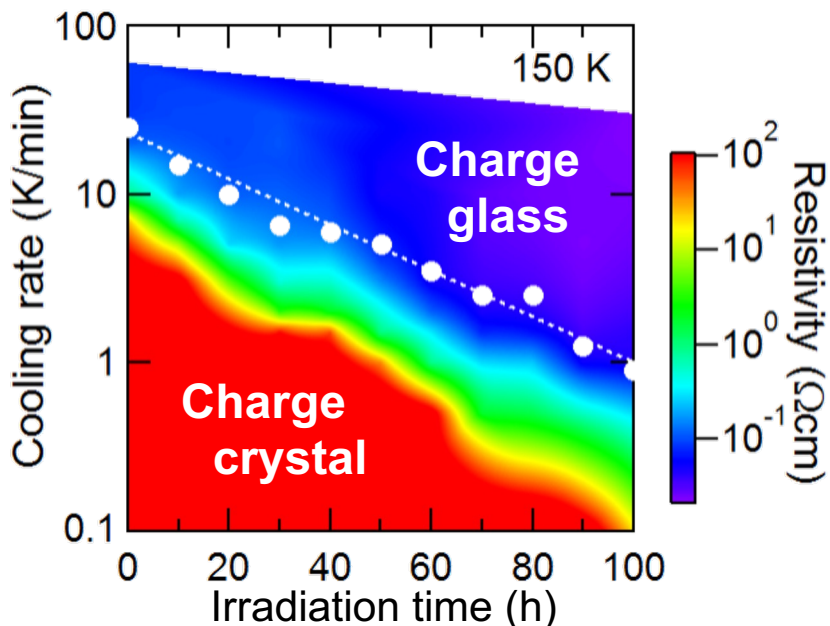
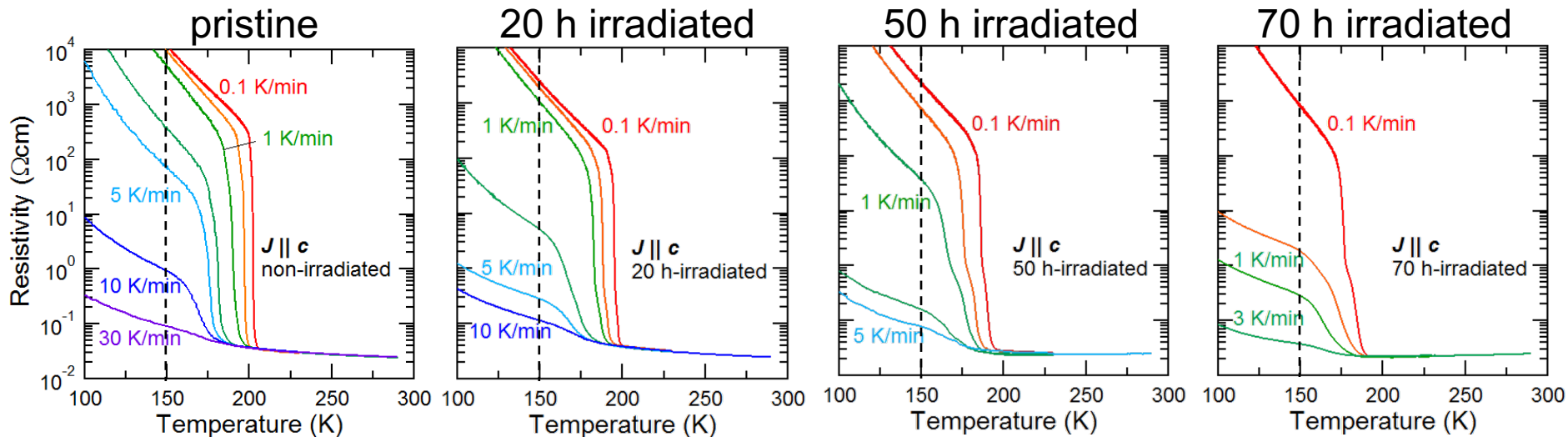
The molecular defects introduced by x-ray irradiation act as impurities.

Cooling-rate dependence of the resistivity curve



The critical cooling rate for the charge glass formation becomes slower with increasing x-ray irradiation time.

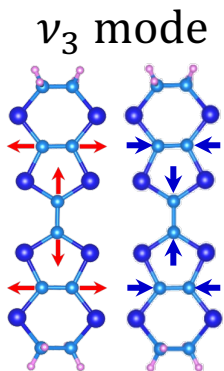
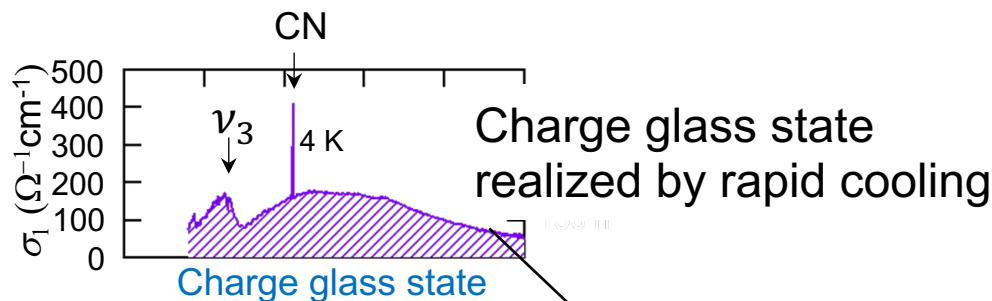
Cooling-rate dependence of the resistivity curve



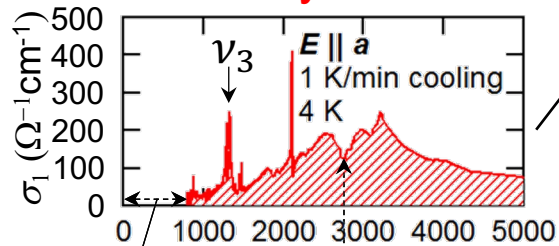
The critical cooling rate for the charge glass formation becomes slower with increasing x-ray irradiation time.



The charge crystal regime is significantly suppressed by randomness introduced by x-rays

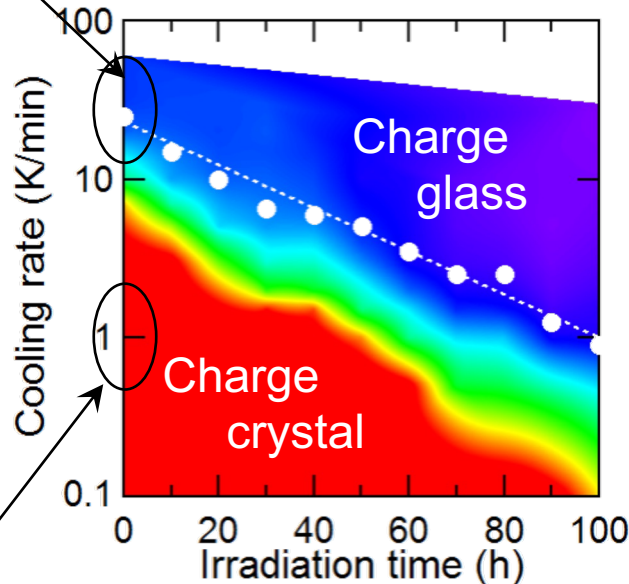


Charge crystal state realized by slow cooling



Optical gap

Overtone of ν_3 mode (2ω)

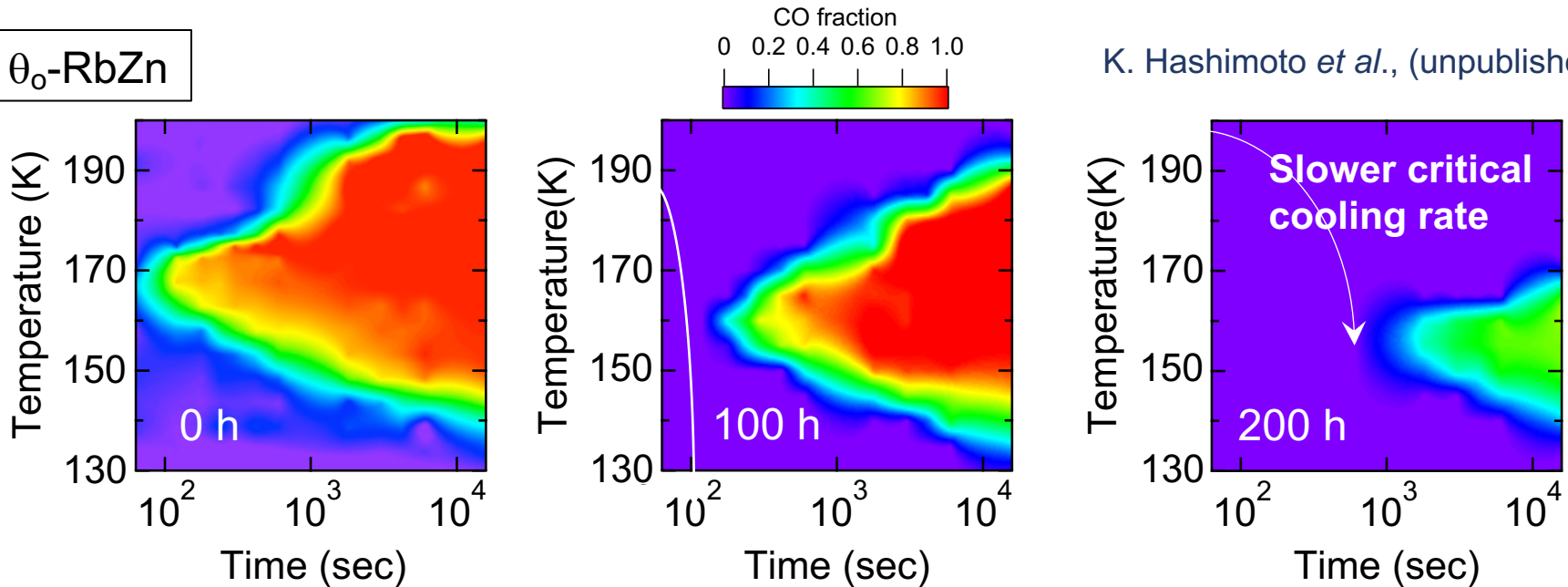


Features of CO state

- Clear optical gap
- Sharpening of ν_3 mode
- Overtone of ν_3 mode

TTT diagram conducted by DSC measurements

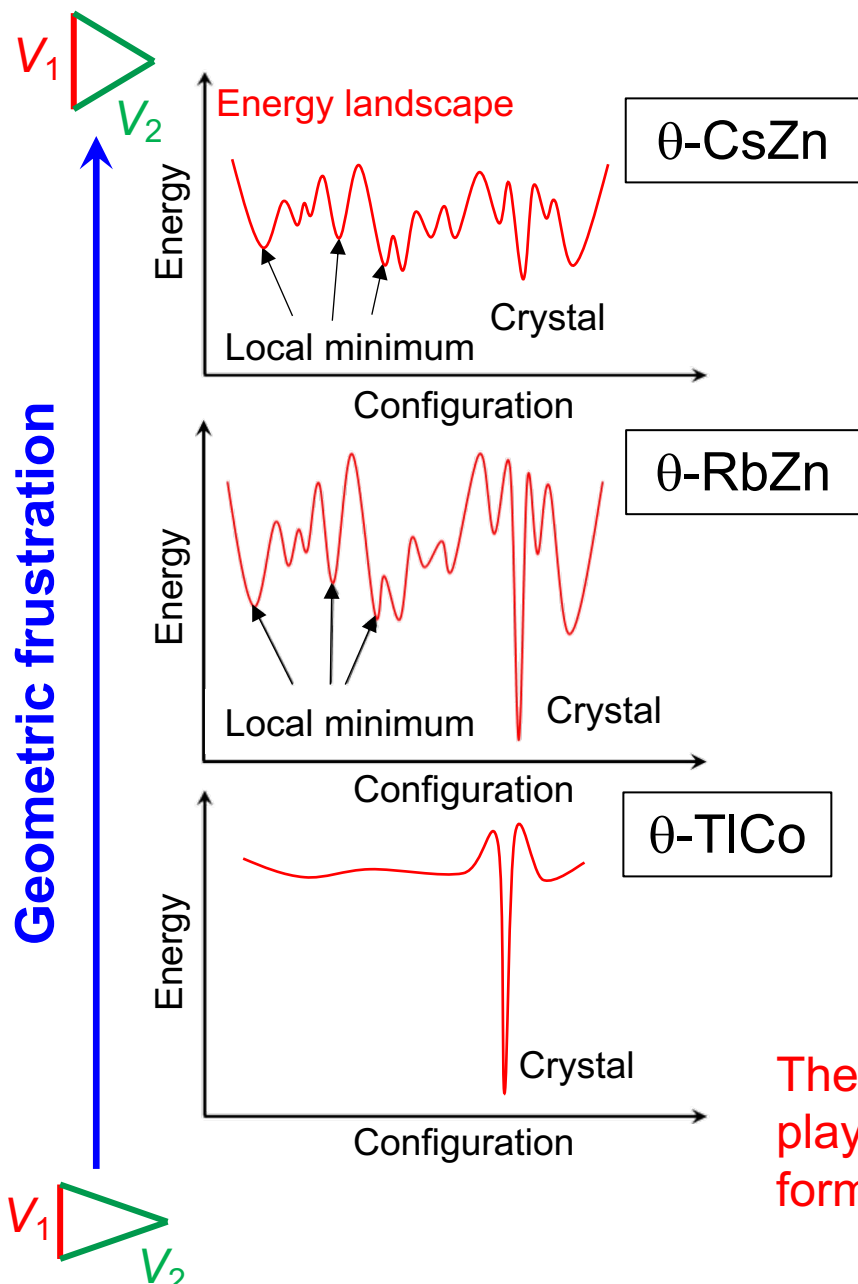
θ_0 -RbZn



- The charge crystallization regime shifts to a longer-time region.
- The nose temperature decreases.

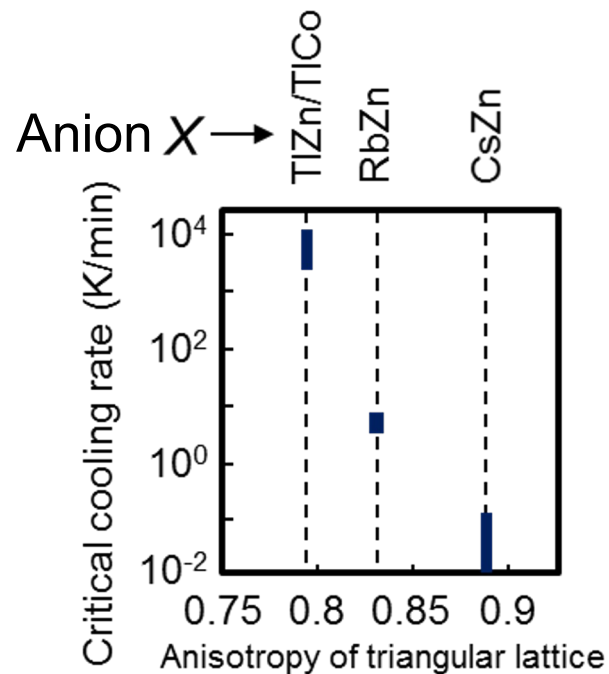
Charge crystallization is significantly suppressed by x-ray irradiation, which leads to a slower critical cooling rate compared to that of the pristine sample.

➤ Even in θ_0 -RbZn, one can make a charge glass state in a very slow cooling rate.

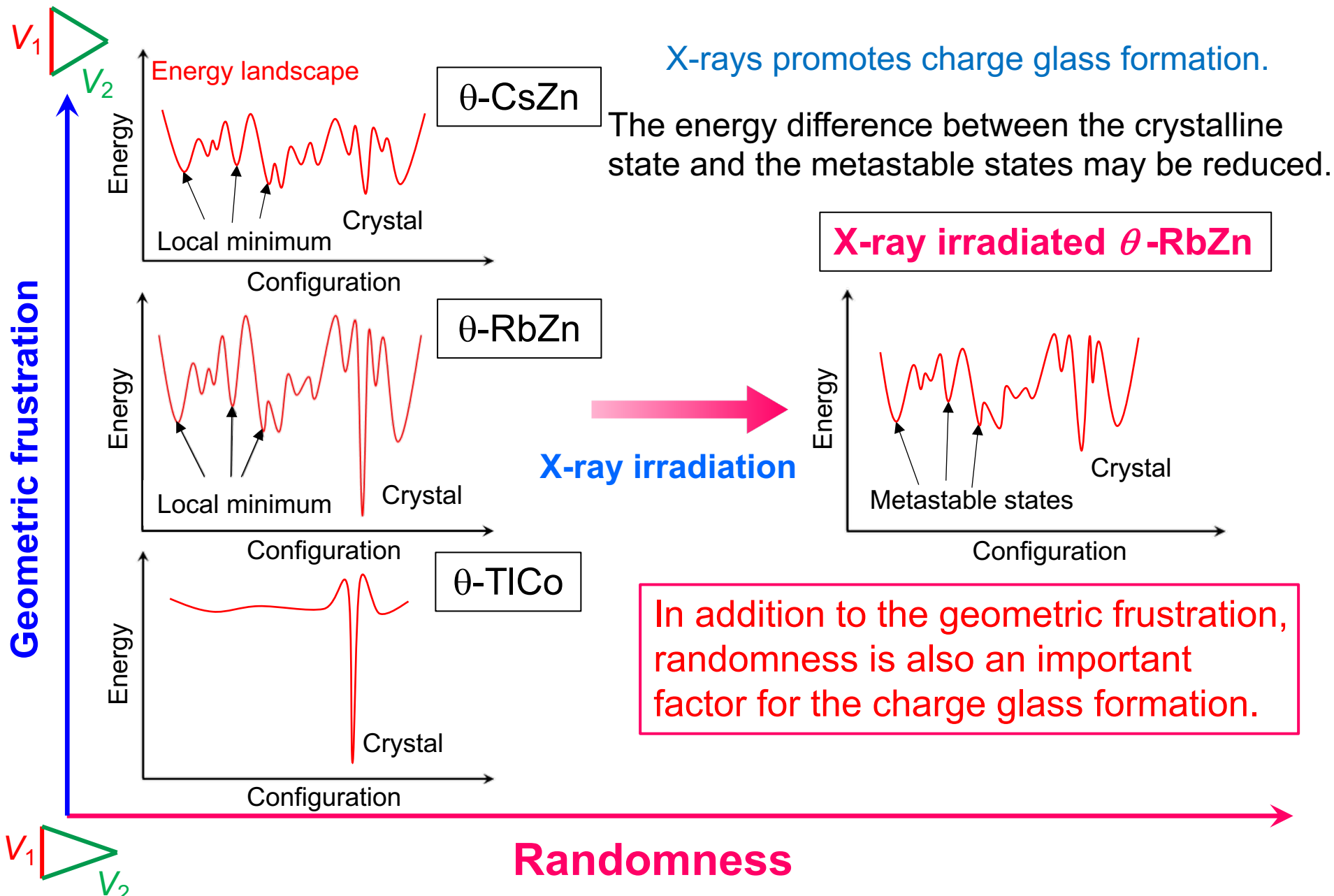


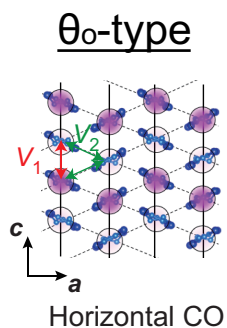
T. Sato *et al.*, JPSJ **83**, 083602 (2014).

H. Oike *et al.*, PRB **91**, 041101 (2015).

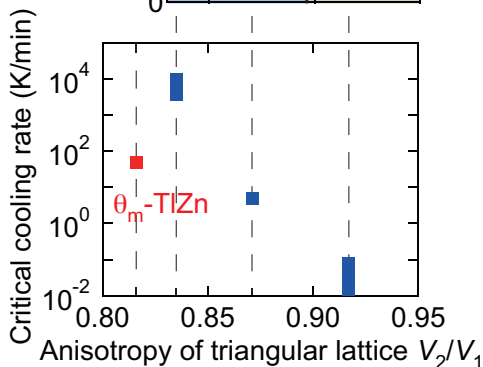
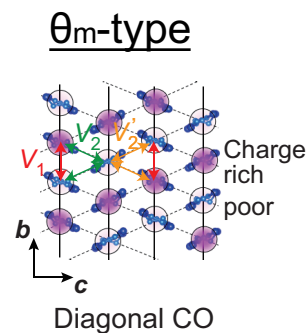
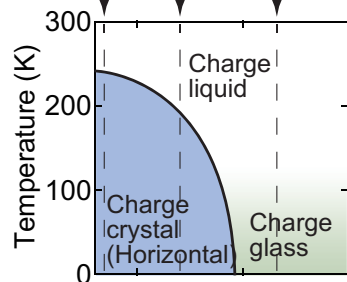


The geometrical frustration of the triangular lattice plays an important role for the charge glass formation in this system.



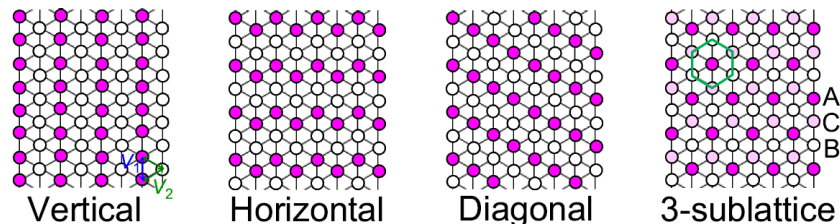


θ_o -TiZn/TiCo θ_o -RbZn θ_o -CsZn



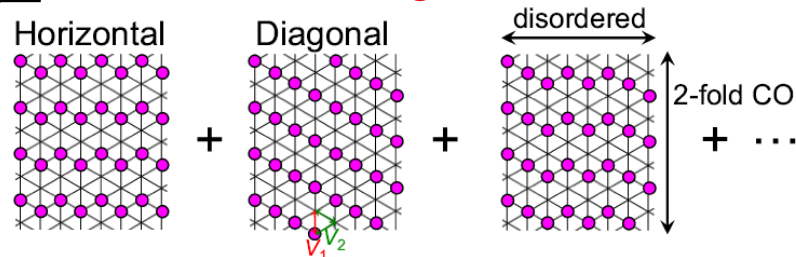
θ_o -type

Metallic 3-fold COs



θ_m -type

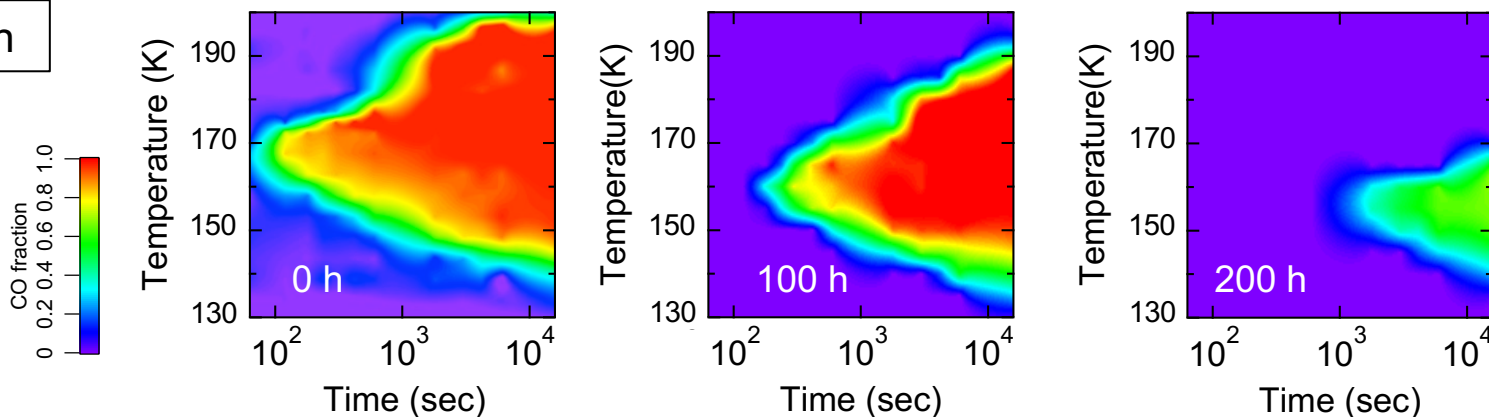
Insulating 2-fold COs



Different charge-glass formation mechanisms in the orthorhombic and monoclinic systems.

TTT diagram conducted by DSC

θ_o -RbZn



Randomness effect is also an important factor for the charge glass formation.