

CRC TRR288

Slow dynamics in the Charge-Glass Forming Organic Conductors θ -(BEDT-TTF)₂MM'(SCN)₄

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SPICE Workshop, May 24th-26th 2022, Ingelheim

Outline

1) Quasi-2D Organic Charge-Transfer Salts

- Charge-glass state in θ -(ET)₂MM'(SCN)₄

2) Fluctuation Spectroscopy

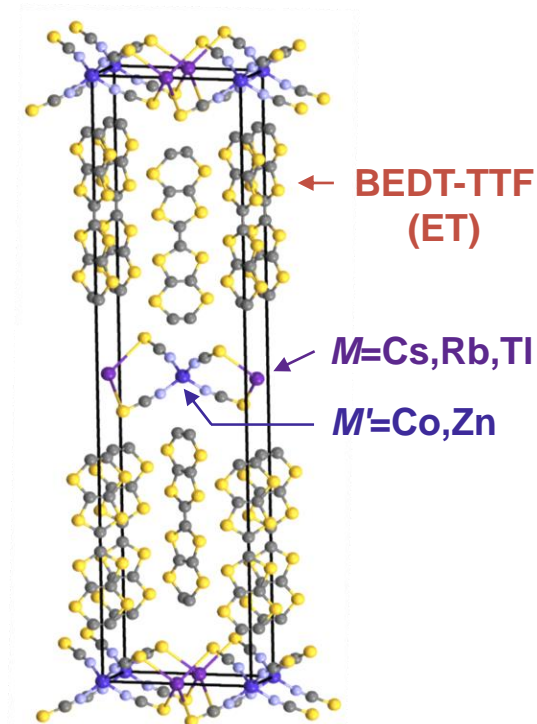
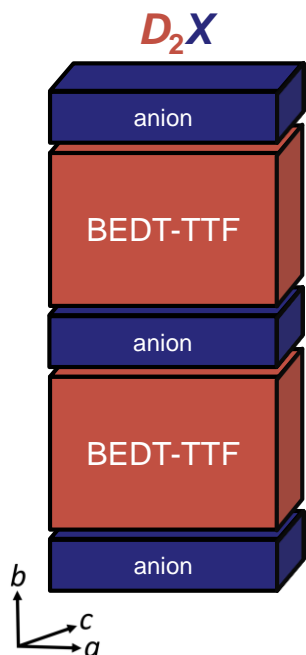
3) Structural Transition in θ -(ET)₂CsM'(SCN)₄

- Resistance
- Thermal expansion
- Resistance noise

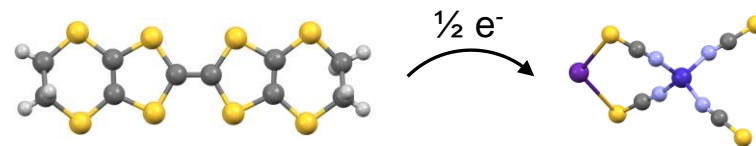
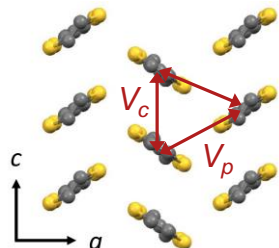
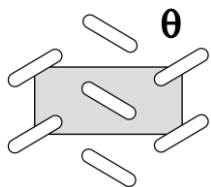
4) Comparison of the Charge-Crystal and Charge-Glass State

- Resistance noise in θ -(ET)₂RbZn(SCN)₄
- Resistance noise in θ_m -(ET)₂TlZn(SCN)₄

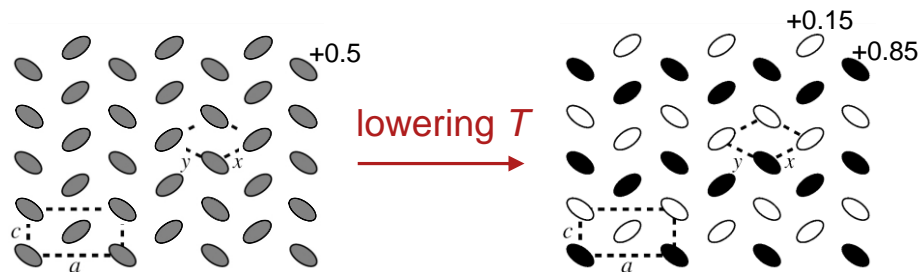
Quasi-2D Organic Charge-Transfer Salts



Packing motif:



- Charge transfer:
 $[D_2] + [X] \rightarrow [D_2]^{+1} + [X]^{-1}$
 $\rightarrow 1/4$ -filled conduction band
- Strong electronic correlations
 \rightarrow Charge ordering (CO)



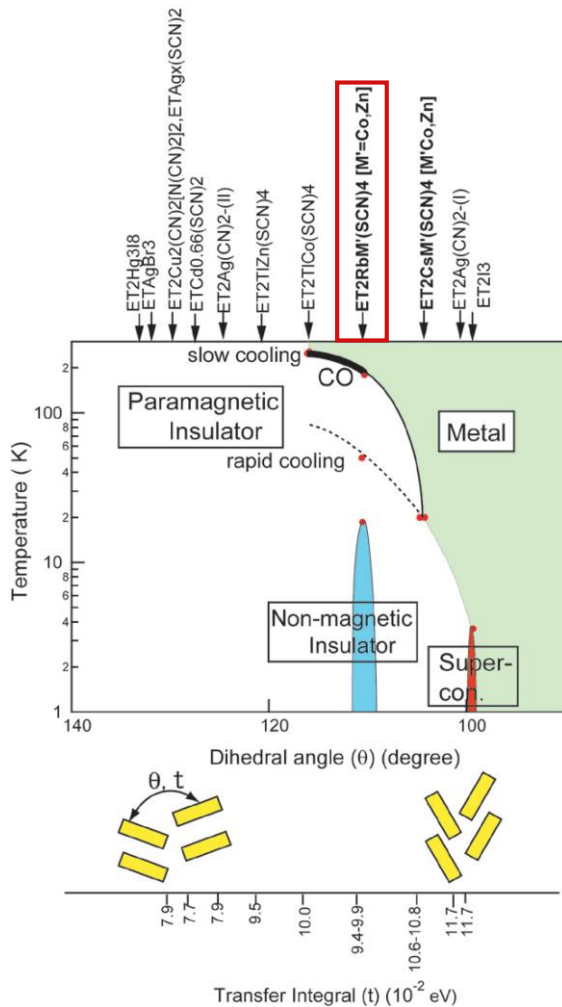
Kuroki, *Sci. Technol. Adv. Mater.* **10**, 024312 (2009)

Extended Hubbard model:

$$\mathcal{H} = \sum_{\langle ij \rangle, \sigma} t_{ij} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + \sum_i U n_{i\uparrow} n_{i\downarrow} + \sum_{\langle ij \rangle} V_{ij} n_i n_j$$

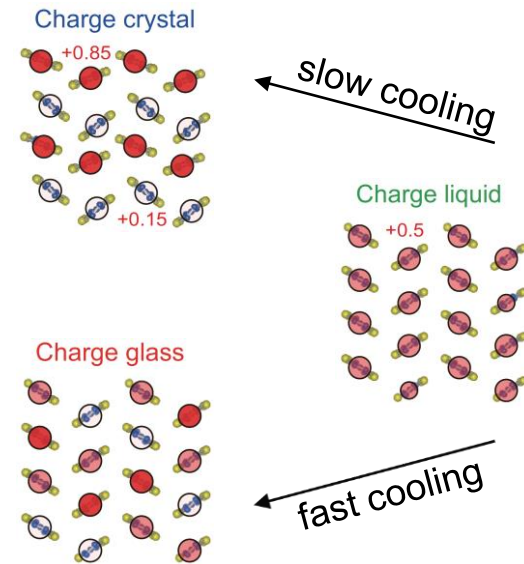
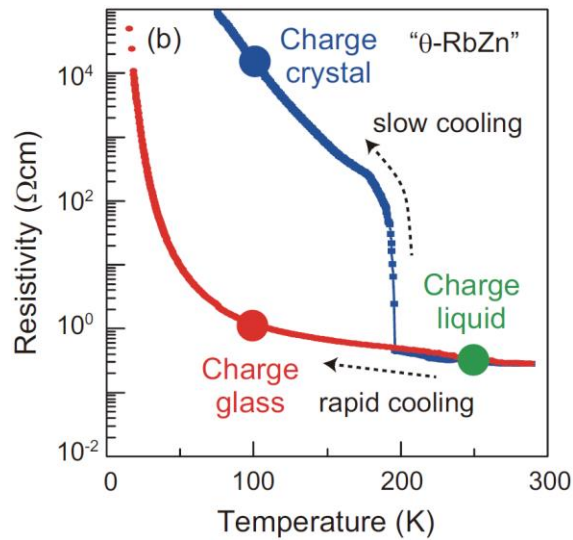
Seo, *JPSJ* **69**, 805 (2000)

Different Charge States in θ -(BEDT-TTF)₂X



Mori et al., *PRB* **57**, 12023 (1998),
 Kuroki, *Sci. Technol. Adv. Mater.* **10**, 024312 (2009)

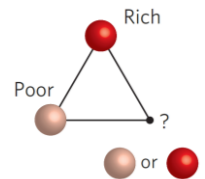
Cooling-rate dependence:



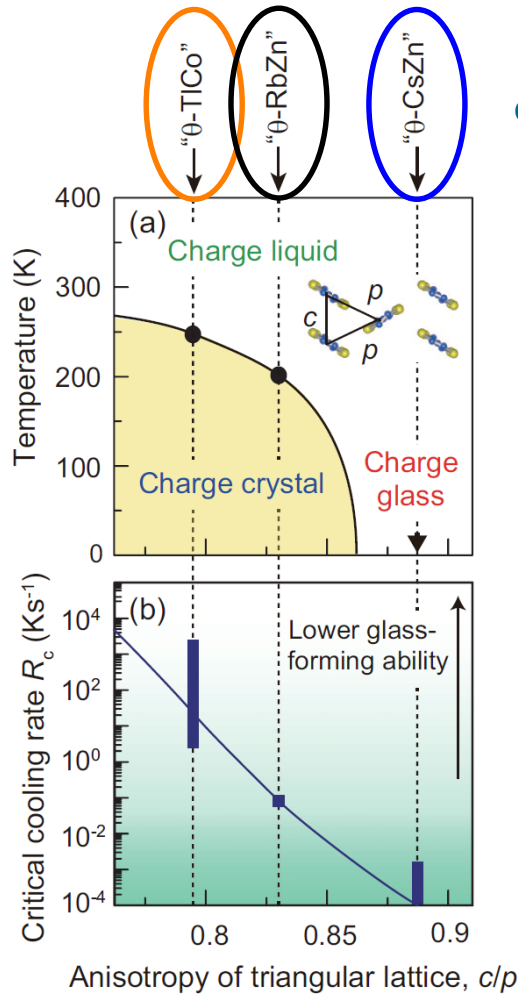
Oike et al., *PRB* **91**, 041101(R) (2015),
 Kagawa et al., *Nat. Phys.* **9**, 419 (2013)

- CO transition can be kinetically avoided
 → Charge-glass state

Geometric frustration:

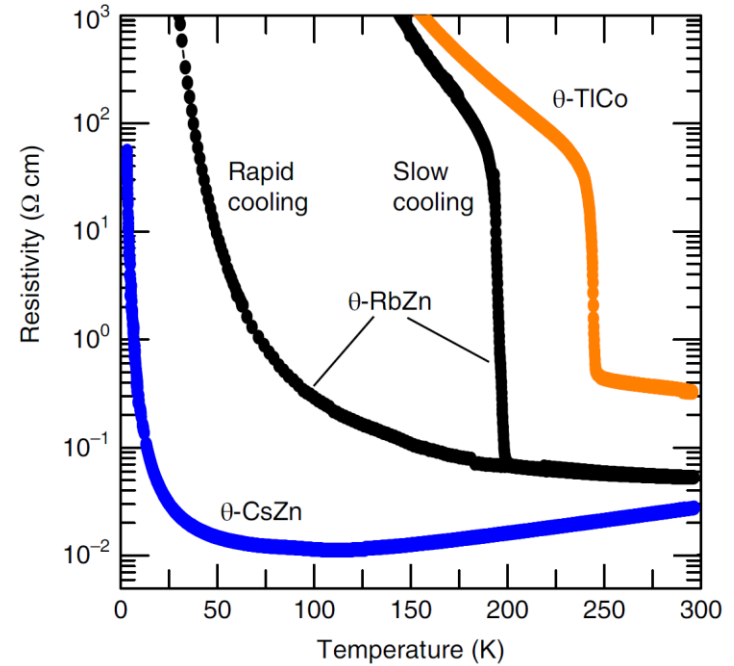
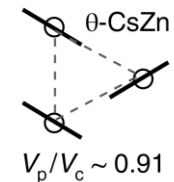
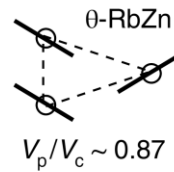
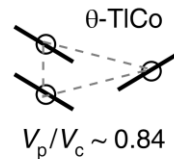


Geometric Frustration



Oike et al., *PRB* **91**, 041101(R) (2015)

orthorhombic θ - MM'



Sato et al., *Nat. Mater.* **18**, 229 (2019)

Stronger frustration

- reduces T_{CO}
- reduces critical cooling rate q_c
- enhances glass-forming ability

Charge-Glass State

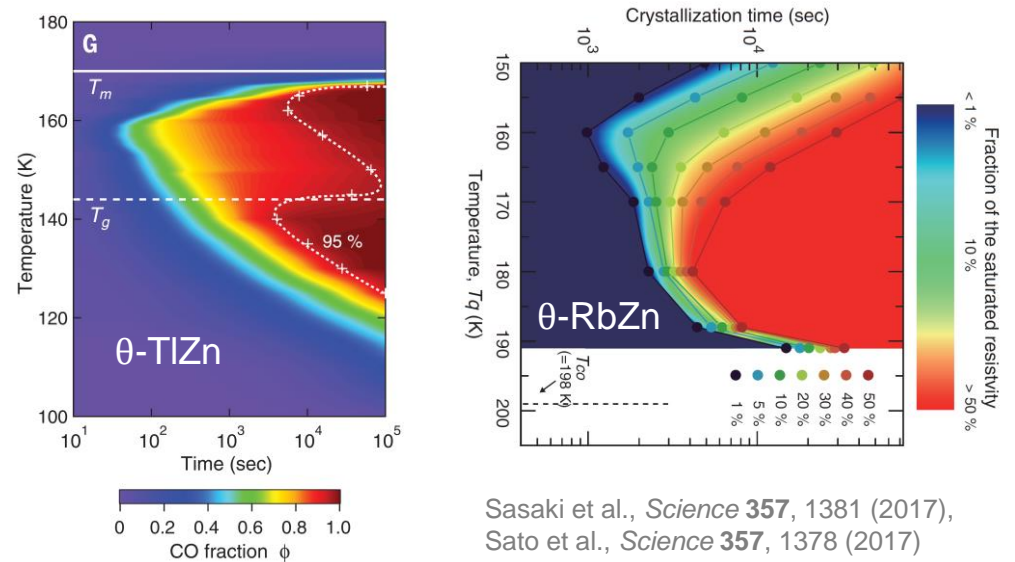
- Cooling-rate dependence
- Lack of long-range order
- Non-equilibrium dynamics

Slow and heterogeneous dynamics

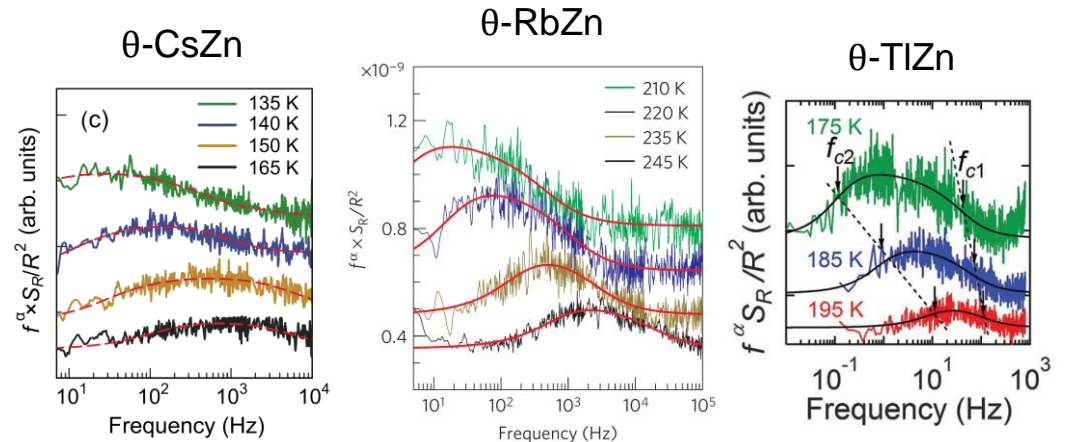
→ **Fluctuation spectroscopy**

Sato et al., *PRB* **89**, 121102(R) (2014),
 Kagawa et al., *Nat. Phys.* **9**, 419 (2013),
 Sasaki et al., *Science* **357**, 1381 (2017)

Electronic crystallization:



Sasaki et al., *Science* **357**, 1381 (2017),
 Sato et al., *Science* **357**, 1378 (2017)



Charge-Glass State

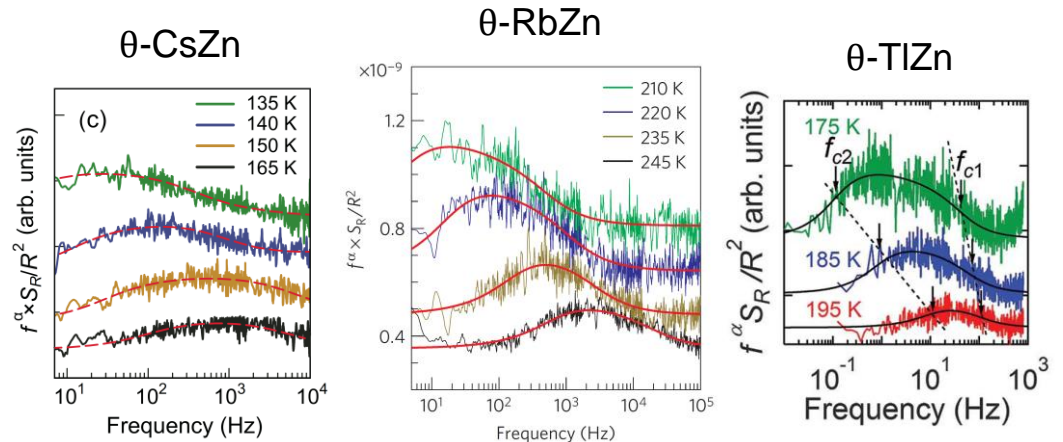
Open questions:

- Charge dynamics in the charge-crystal and charge-glass state?
- Involvement of structural degrees of freedom?

Slow and heterogeneous dynamics

→ **Fluctuation spectroscopy**

Sato et al., *PRB* **89**, 121102(R) (2014),
Kagawa et al., *Nat. Phys.* **9**, 419 (2013),
Sasaki et al., *Science* **357**, 1381 (2017)



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1) Quasi-2D Organic Charge-Transfer Salts

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3) Structural Transition in θ -(ET)₂CsM'(SCN)₄

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Fluctuation Spectroscopy

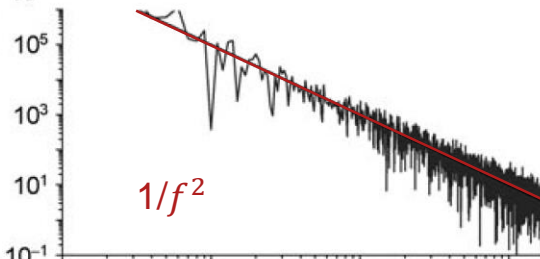
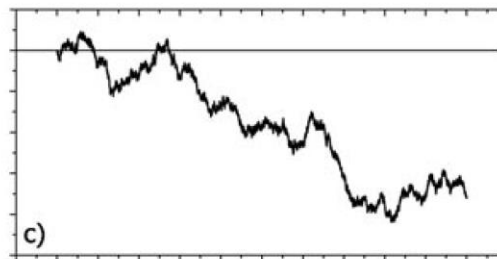
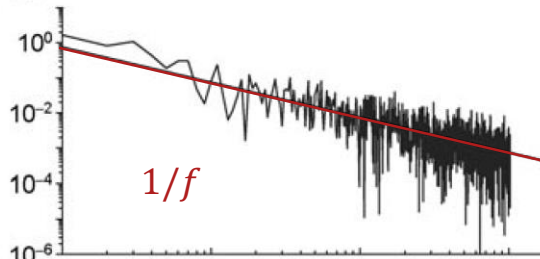
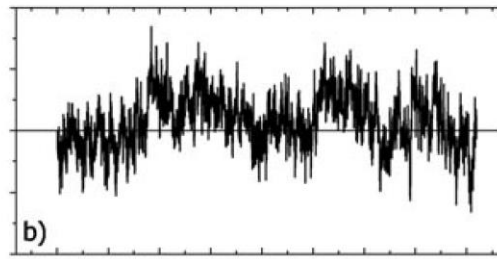
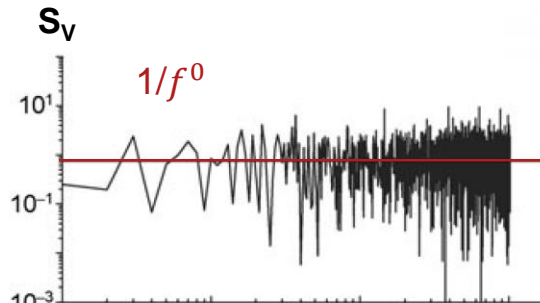
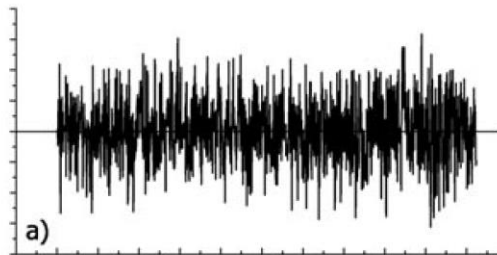
Different types of noise:

Müller, *ChemPhysChem* 12, 1222 (2011)

Power spectral density:

$$S_V(f) = 2 \lim_{T \rightarrow \infty} \frac{1}{T} \left| \int_{-T/2}^{T/2} \delta V(t) e^{-i2\pi f t} dt \right|^2$$

Voltage



time

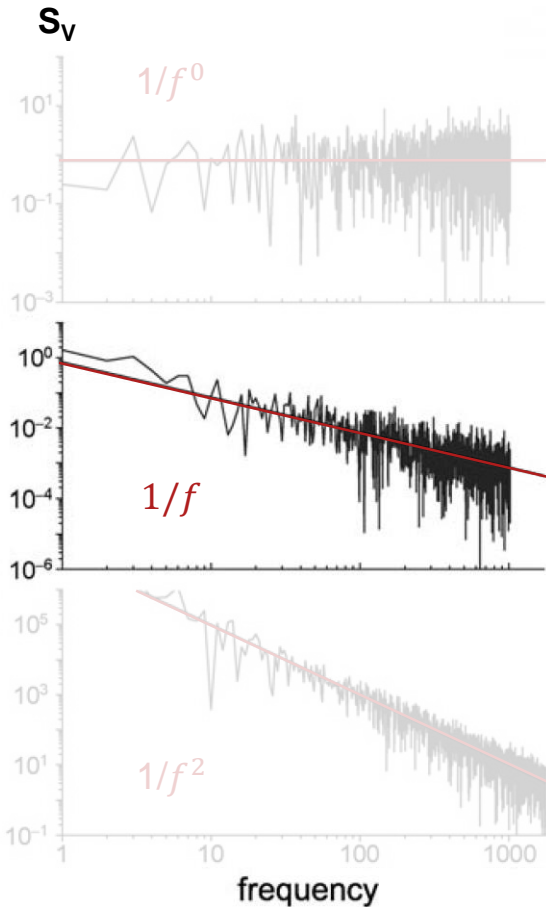
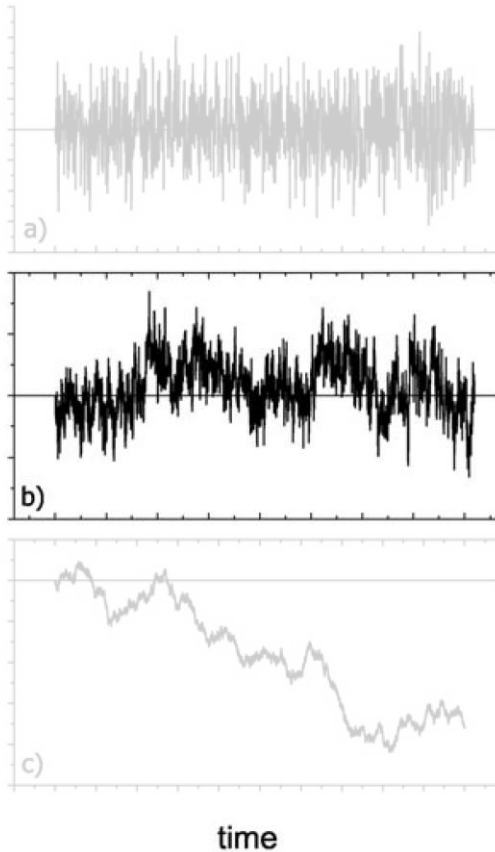
frequency

Fluctuation Spectroscopy

Different types of noise:

Müller, *ChemPhysChem* 12, 1222 (2011)

Voltage



Power spectral density:

$$S_V(f) = 2 \lim_{T \rightarrow \infty} \frac{1}{T} \left| \int_{-T/2}^{T/2} \delta V(t) e^{-i2\pi f t} dt \right|^2$$

1/f noise:

$$S_V \propto 1/f^\alpha$$

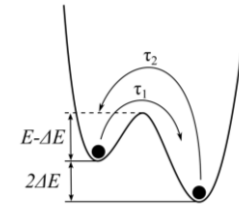
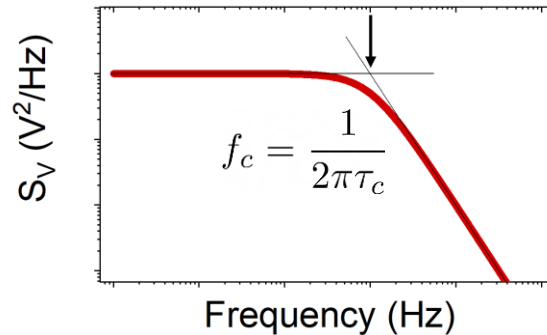
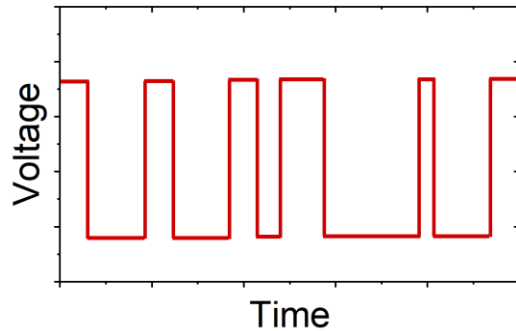
Frequency exponent:

$$\alpha = 0.8 - 1.4$$

Ubiquitous in nature:
Biology, geology, music,
condensed matter, ...

Fluctuation Spectroscopy

Two-level process:



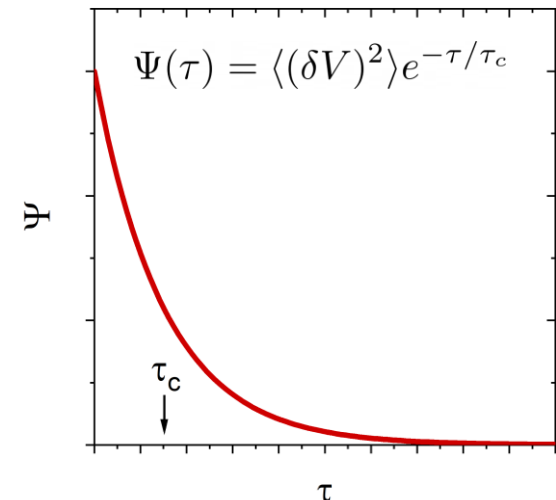
$$\tau_c = \tau_0 \exp\left(\frac{E_A}{k_B T}\right)$$

Wiener-Khinchine theorem:

$$S_V(f) = 4 \int_0^{\infty} \Psi(\tau) \cos(2\pi f\tau) d\tau \quad \propto \frac{\tau_c}{1 + 4\pi^2 f^2 \tau_c^2}$$

Autocorrelation function: $\Psi(\tau) = \langle \delta V(t) \delta V(t + \tau) \rangle$

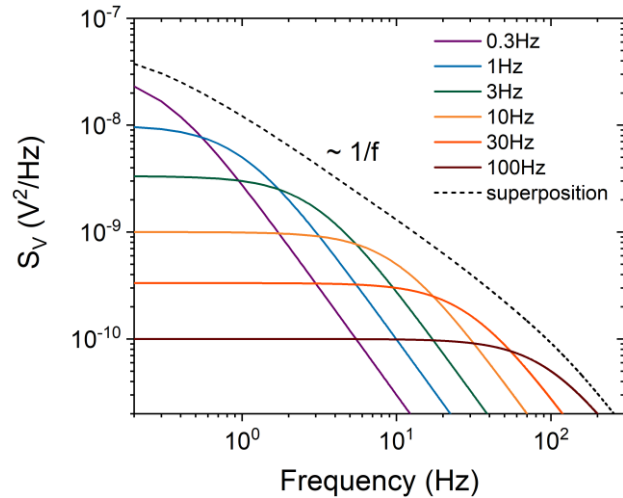
→ Microscopic kinetics of charge carriers



Fluctuation Spectroscopy

Simple model:

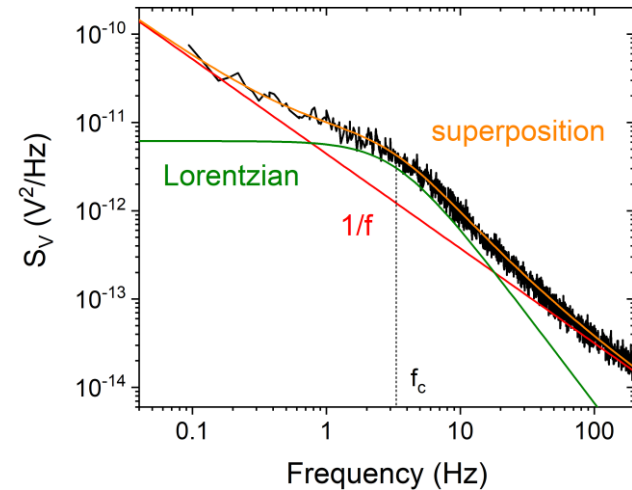
Superposition of many independent fluctuators



Dutta-Dimon-Horn (DDH) model Dutta et al., PRL 43, 9 (1979)

→ access to energy distribution of fluctuators

Dominating fluctuator:



T-dependence: $f_c = f_0 \exp\left(-\frac{E_A}{k_B T}\right)$

→ access to activation energy of fluctuator

What causes slow dynamics?

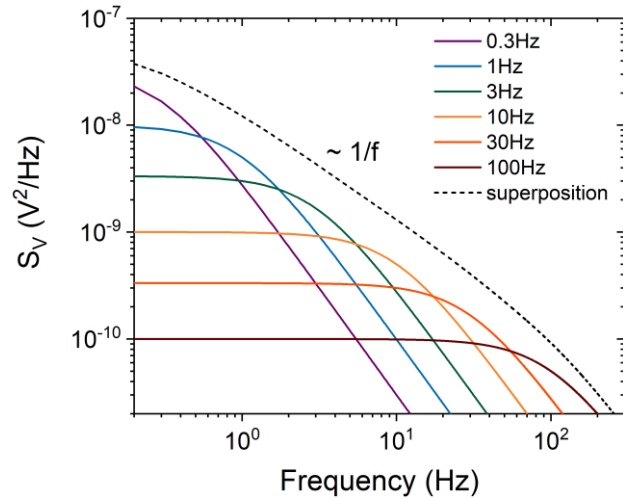
- large activation energy of localized process
 - collective effects

Parisi, Physik Journal 10, Nr. 8/9 (2011)

Fluctuation Spectroscopy

Simple model:

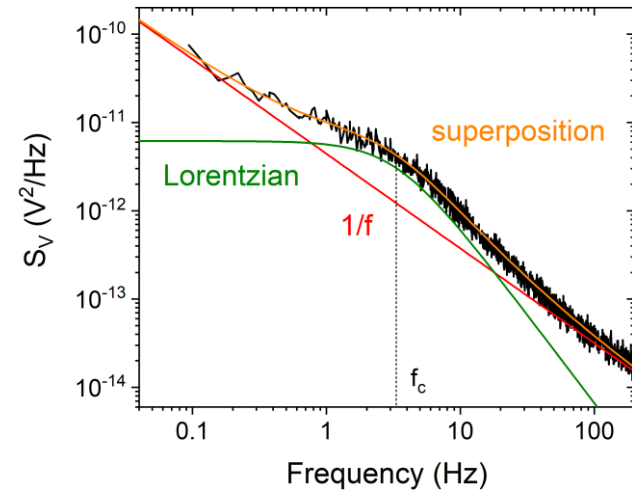
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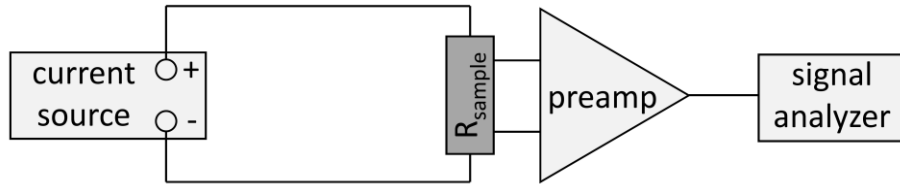
T-dependence: $f_c = f_0 \exp\left(-\frac{E_A}{k_B T}\right)$

→ access to activation energy of fluctuator

Useful for studying

- glassy transitions
- collective effects
- excitation of large clusters
- inhomogeneous current distribution
- percolation
- ...

Experiment

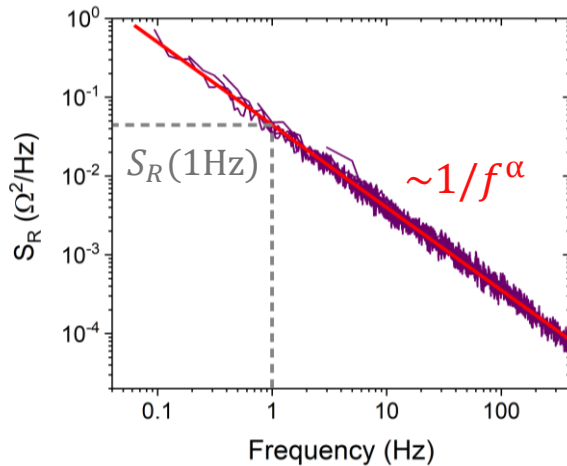
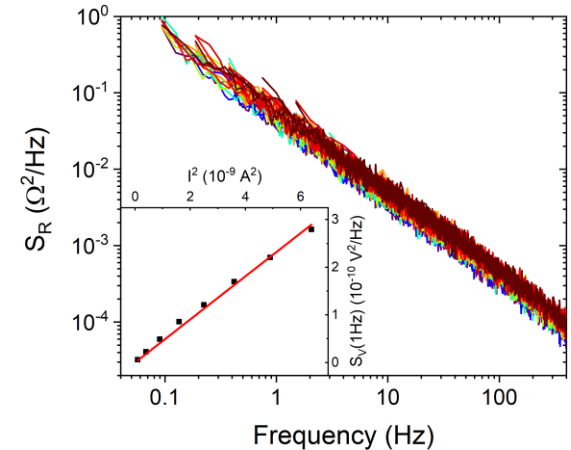
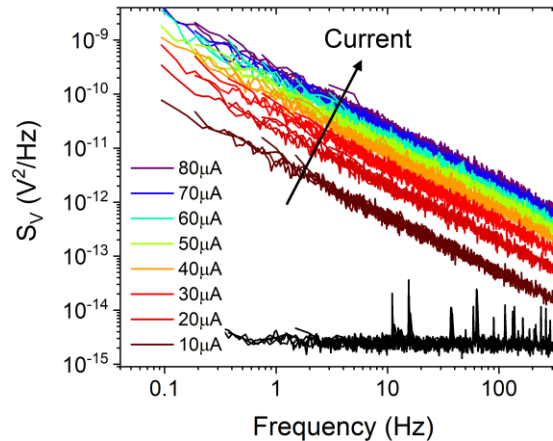


Frequencies: mHz-kHz
Impedances: mΩ-GΩ

Current/voltage scaling:

$$S_V = \frac{\gamma_H V^2}{n_c \Omega} \frac{1}{f^\alpha} \propto V^2 \propto I^2$$

Hooge, *Phys. Lett. A* **29**, 139 (1969)



- Normalized noise magnitude $S_R(1\text{Hz})/R^2$
 - Frequency exponent α
- in dependence of T, B, E , etc.

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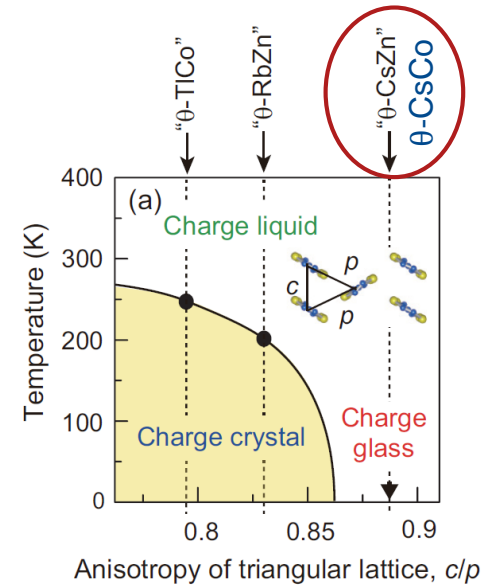
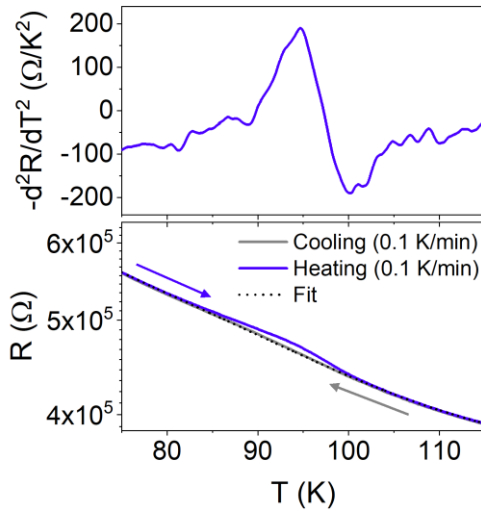
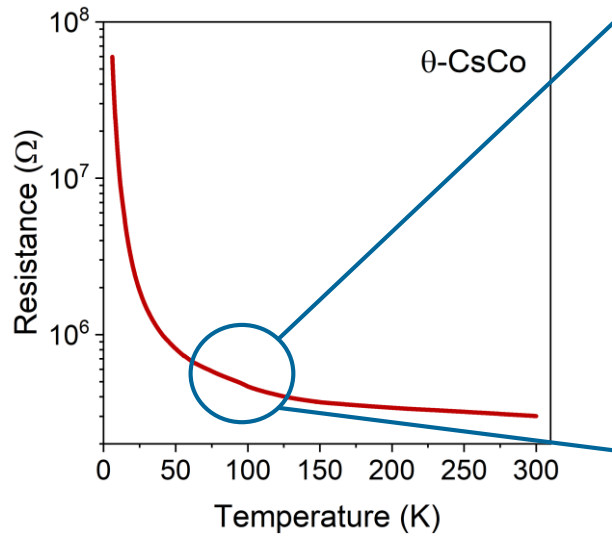
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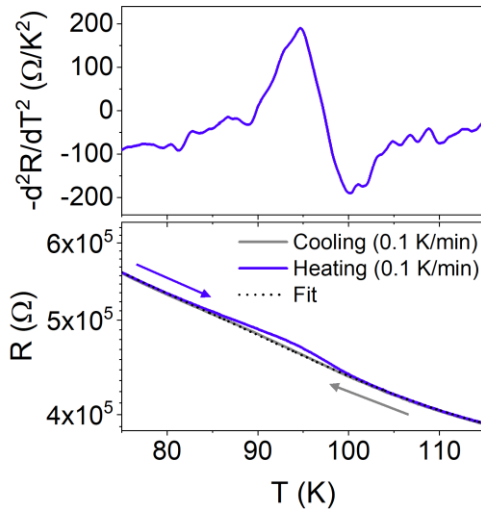
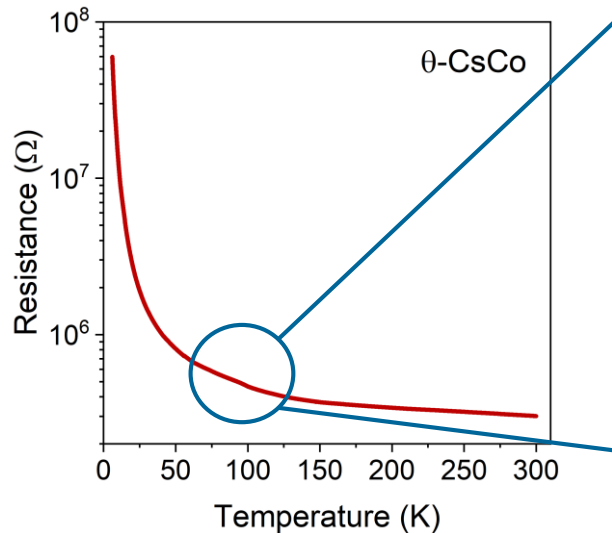
Resistance of θ -CsCo

Resistance anomaly at $T_g \sim 95\text{K}$

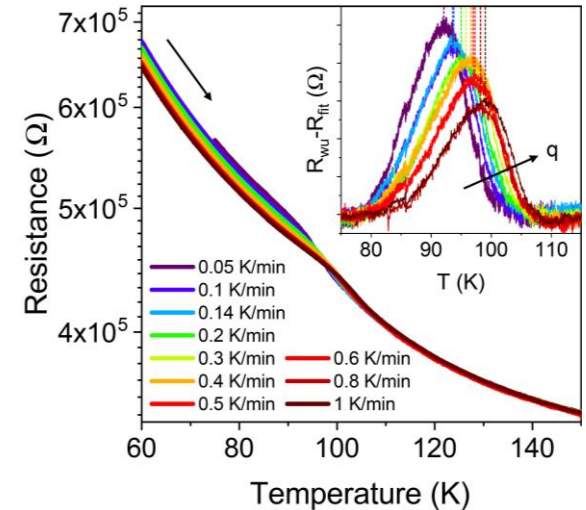


Resistance of θ -CsCo

Resistance anomaly at $T_g \sim 95\text{K}$



Cooling-rate dependence:

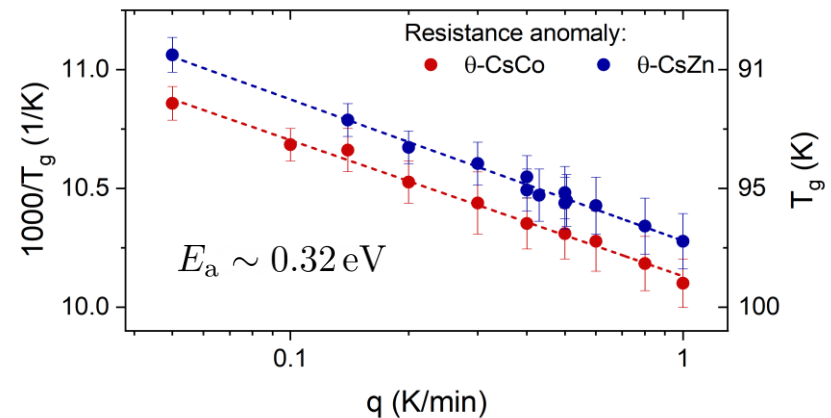


- Hysteresis and q -dependence
- Consistent with measurements on θ -CsZn
- Previously observed Sato et al., PRB 89, 121102(R) (2014)

Criterion for T_g :

$$-|q| \cdot \frac{d\tau}{dT} \Big|_{T_g} \approx 1$$

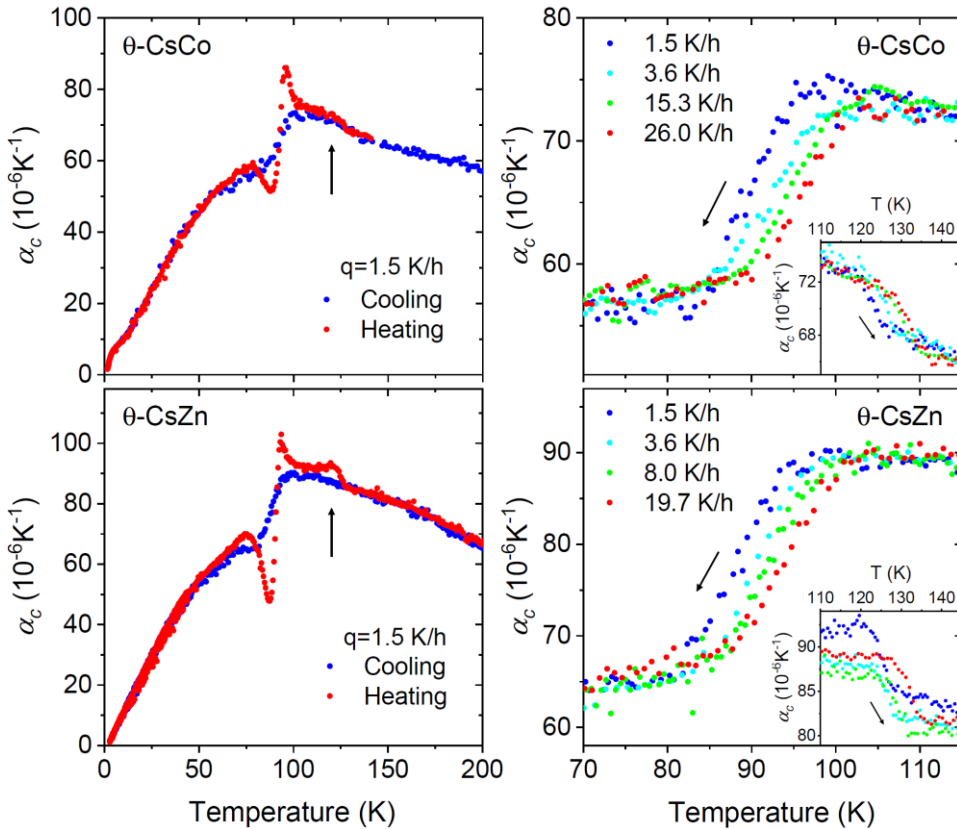
$$\Rightarrow \ln |q| \sim -\frac{E_a}{k_B} \cdot T_g^{-1}$$



Thermal Expansion of θ -CsCo

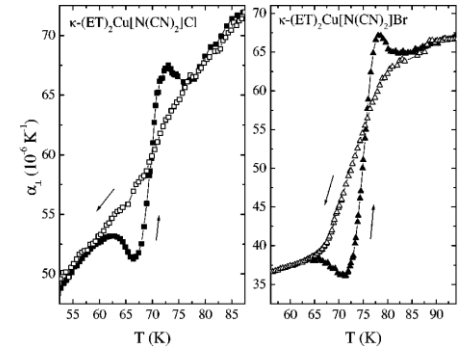
by Dr. Y. Saito, Prof. M. Lang

Cooling-rate dependence:

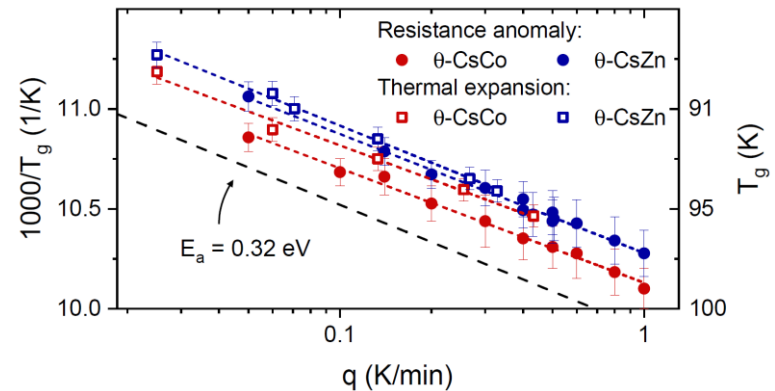
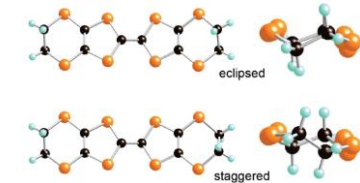


- Structural glass transition at $T_g \sim 95\text{K}$
- Connected to resistance anomaly
- Second glassy transition at $T_g^\dagger \sim 120\text{K}$

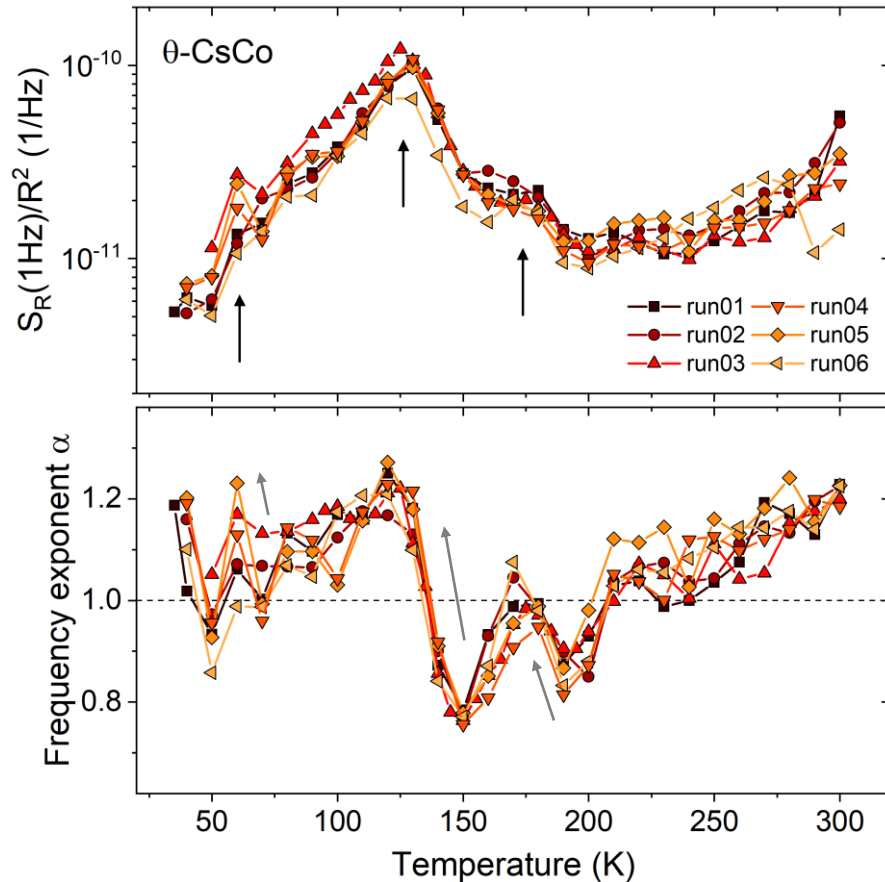
cf. κ -(ET)₂X:



Müller et al., *PRB* **65**, 144521 (2002)
Müller et al., *New J. Phys.* **17**, 083057 (2015)

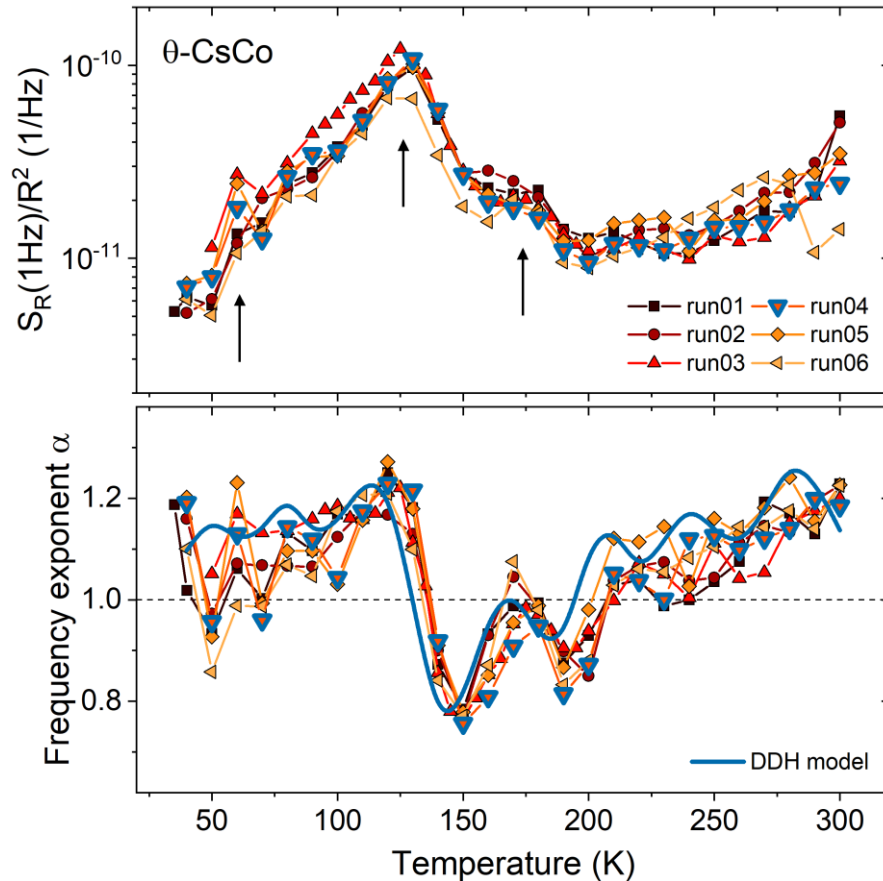


Resistance Noise in θ -CsCo



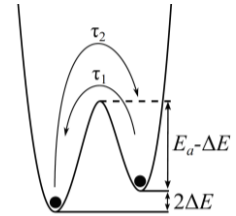
- Noise maximum at $T \sim 130\text{K}$
 - Spectral weight shifts to lower frequencies
- Charge dynamics slows down when approaching T_g
- Cf. θ -CsZn
Sato et al., *JPSJ* 85, 123702 (2016)
 - Small features at 175K and 60K

Resistance Noise in θ -CsCo



Dutta-Dimon-Horn model

Dutta et al., *PRL* **43**, 9 (1979),
Raquet et al., *PRB* **59**, 12435 (1999)

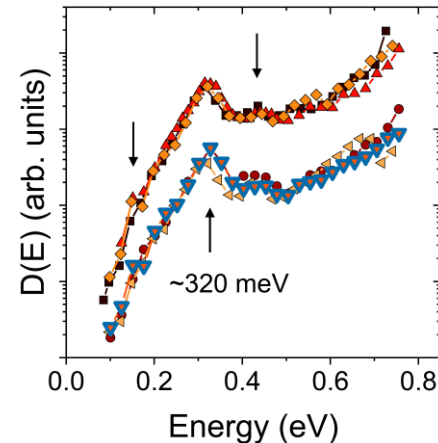


- Independent fluctuators
- Thermally activated two-level process

$$\tau = \tau_0 \exp\left(\frac{E_A}{k_B T}\right)$$

- Distribution of activation energies

$$\frac{S_R(f)}{R^2}(T) = \frac{1}{\pi f} \int_0^\infty g(T) \frac{D(E)}{\cosh[(E - E_f)/k_B T]} dE$$



Frequency-Dependent Noise Maximum

- cf. $\kappa\text{-(ET)}_2\text{X}$
 - similar glassy freezing of EEG in $\theta\text{-(ET)}_2\text{X}$
- role of ethylene endgroups for CO:
 - **order-disorder structural transition of EEG**

Alemaný *et al.*, *J. Phys.: Condens. Matter* **27**, 465702 (2015)

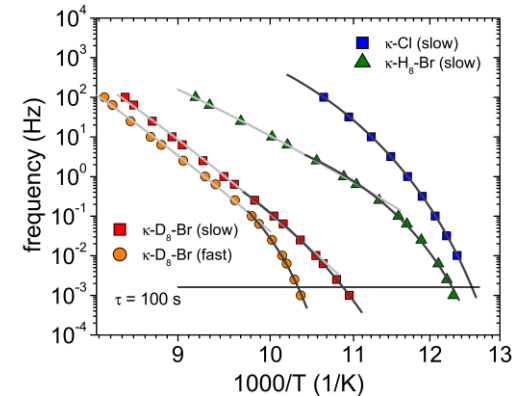
Vogel-Fulcher-Tammann law:

$$f = f_0 \exp\left(\frac{-DT_{\text{VF}}}{T - T_{\text{VF}}}\right)$$

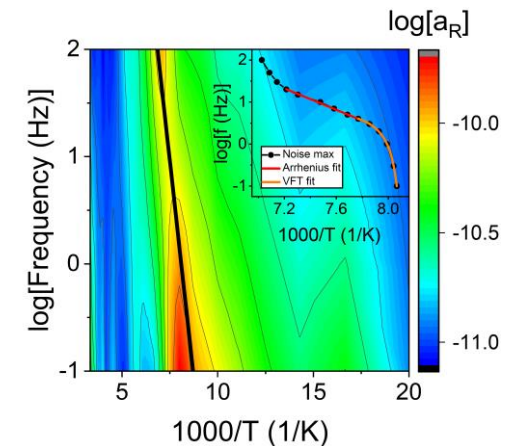
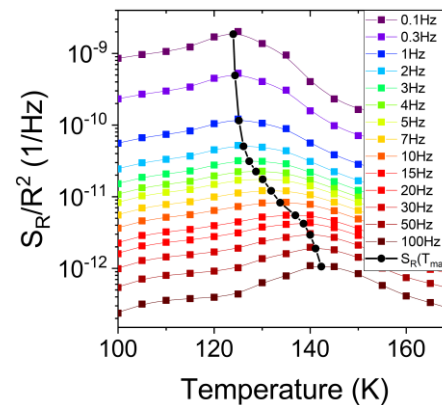
$\theta\text{-CsCo}$:

- unusual frequency dependence
 - **due to coupling of structural and charge dynamics**

$\kappa\text{-(ET)}_2\text{X}$:
fragile glass former



Müller *et al.*, *New J. Phys.* **17**, 083057 (2015)



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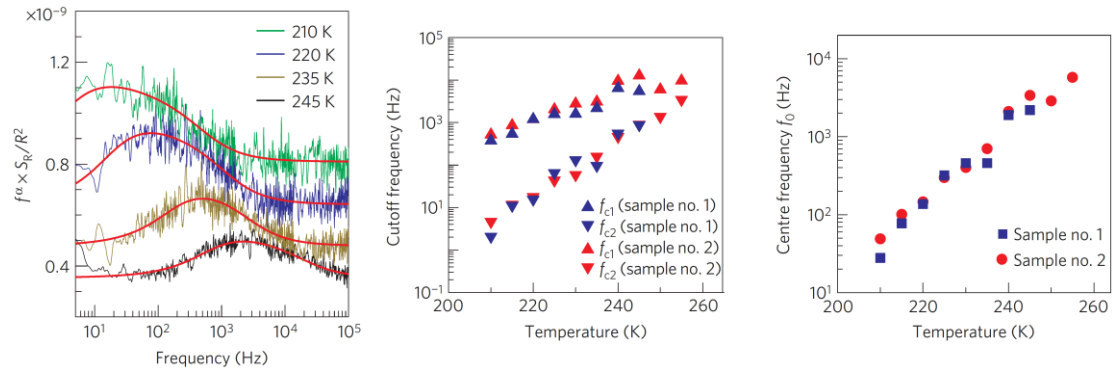
- Resistance noise in θ -(ET)₂RbZn(SCN)₄
- Resistance noise in θ_m -(ET)₂TIZn(SCN)₄

Previous Results in the CL State

Fluctuation spectroscopy:

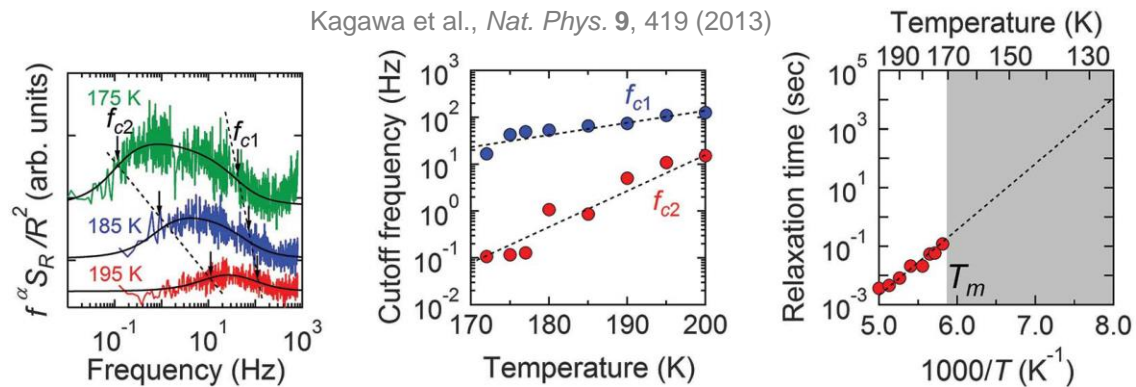
Slow and heterogeneous dynamics above T_{MI}

θ -RbZn



Kagawa et al., *Nat. Phys.* **9**, 419 (2013)

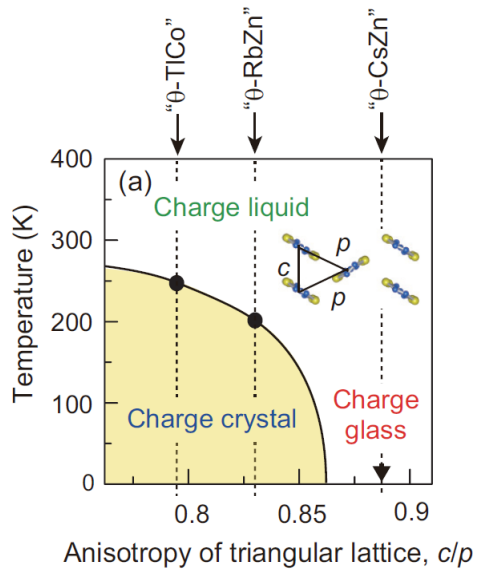
θ -TiZn



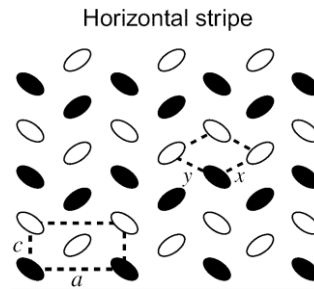
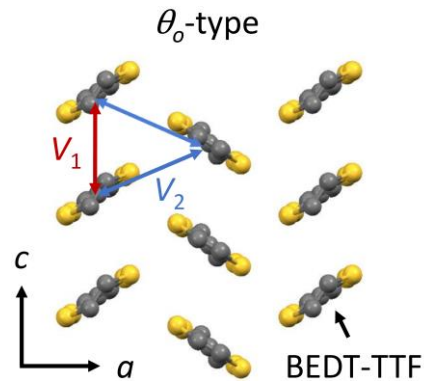
Sasaki et al., *Science* **357**, 1381 (2017)

→ extend noise measurements to CC and CG state

Comparison of the CC and CG State



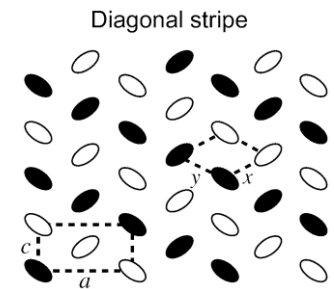
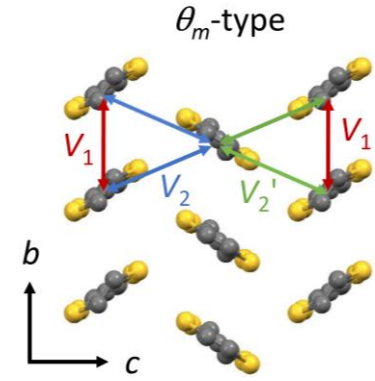
Orthorhombic θ_o -RbZn



$$T_{MI} \sim 200\text{K}$$

$$q_c \sim 5\text{K/min}$$

Monoclinic θ_m -TiZn



$$T_{MI} \sim 175\text{K}$$

$$q_c > 50\text{K/min}$$

Kagawa et al., *Nat. Phys.* **9**, 419 (2013)

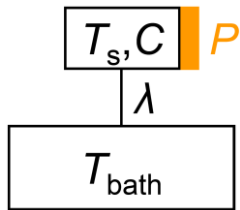
Sasaki et al., *Science* **357**, 1381 (2017)

Varying

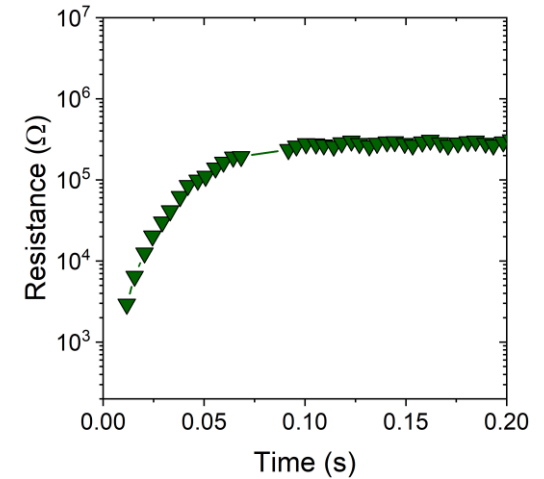
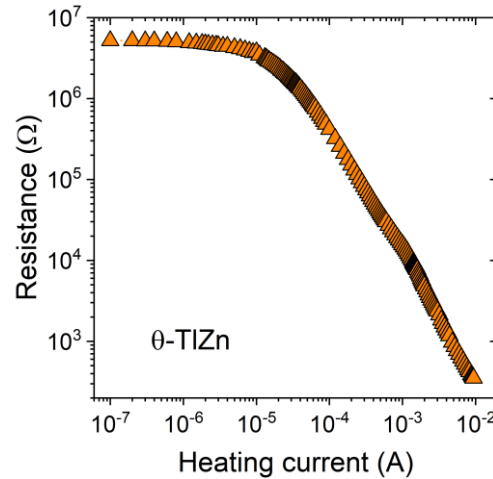
- CO pattern
- degree of frustration
- strength of el.-ph. coupling
- critical cooling rates

Comparison of the CC and CG State

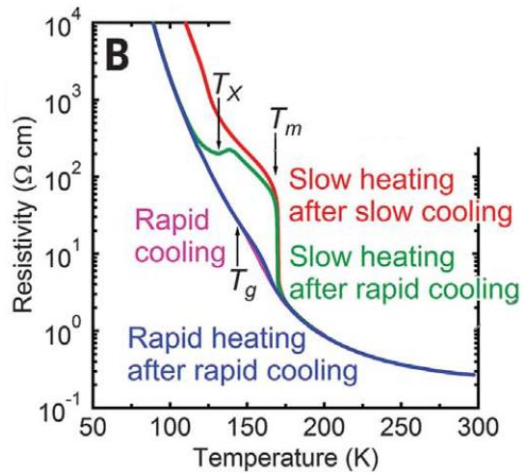
Heat pulse method



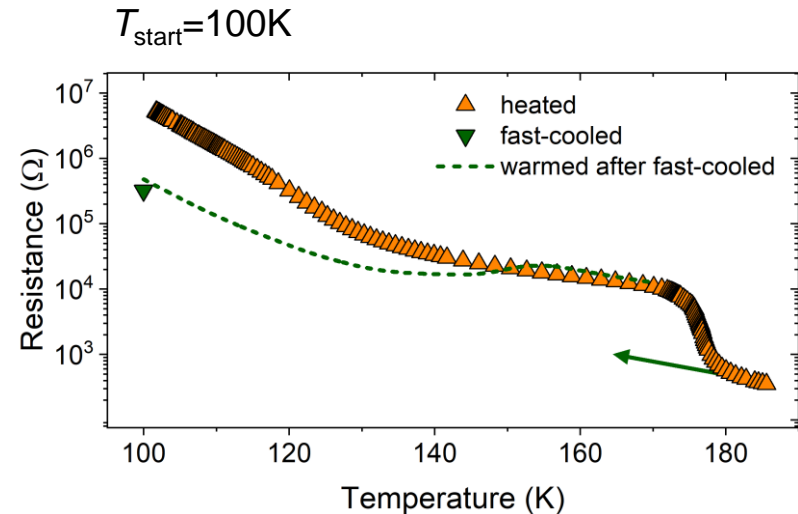
Hartmann et. al., *PRB* **90**, 195150 (2014)



→ Cooling rates of $q \sim 700\text{K/s}$

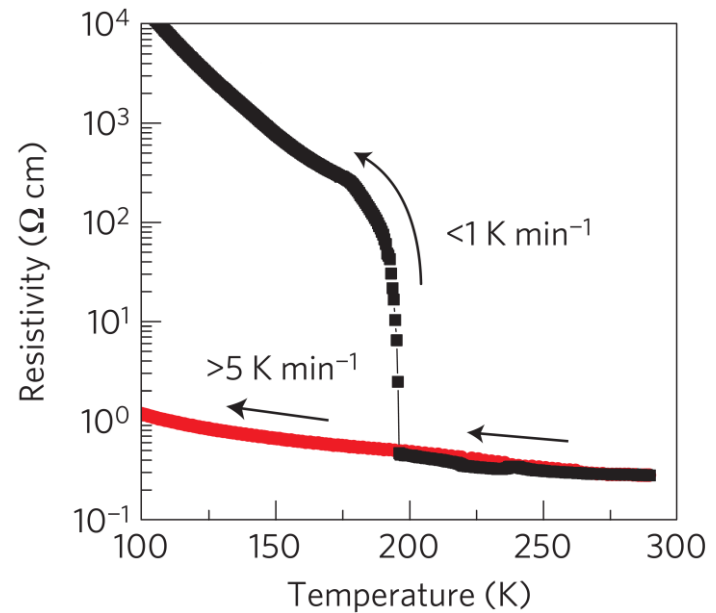
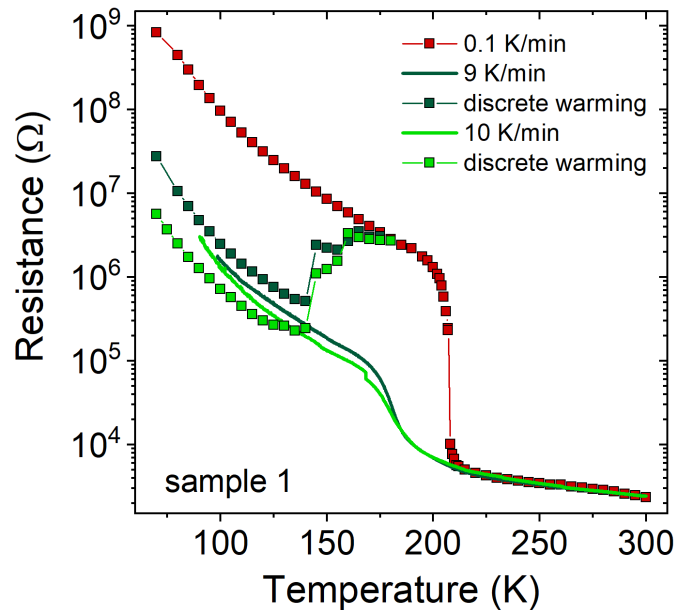


Sasaki et al., *Science* **357**, 1381 (2017)



Resistance of θ -RbZn

Cooling-rate dependence



Kagawa et al., *Nat. Phys.* **9**, 419 (2013)

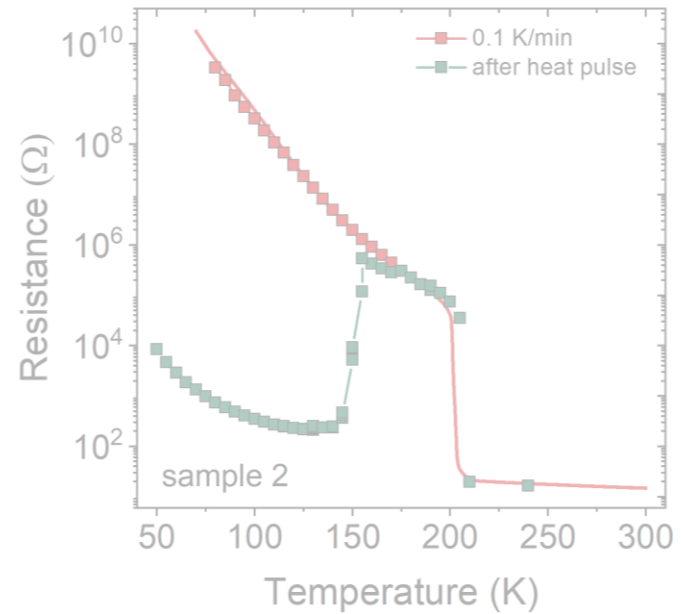
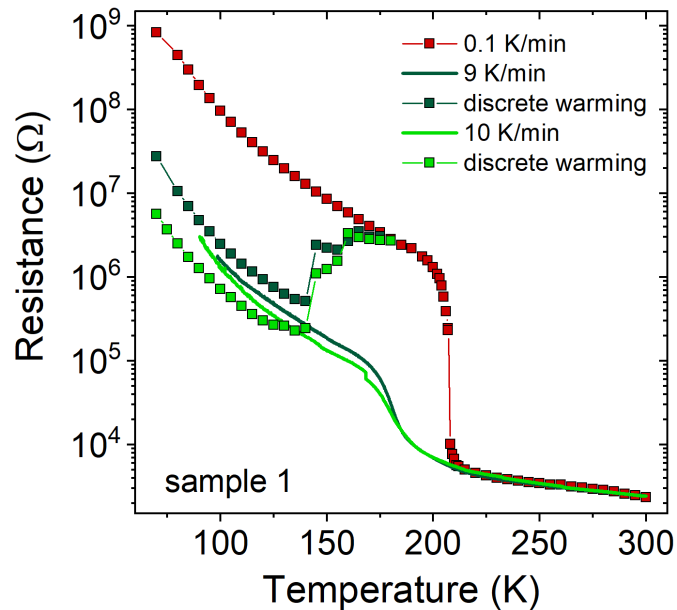
Critical cooling rate $q_c \sim 5 \text{ K/min}$?!

→ dependence on sample quality?

(samples from Prof. H. M. Yamamoto)

Resistance of θ -RbZn

Cooling-rate dependence



Critical cooling rate $q_c \sim 5 \text{ K/min}$?!

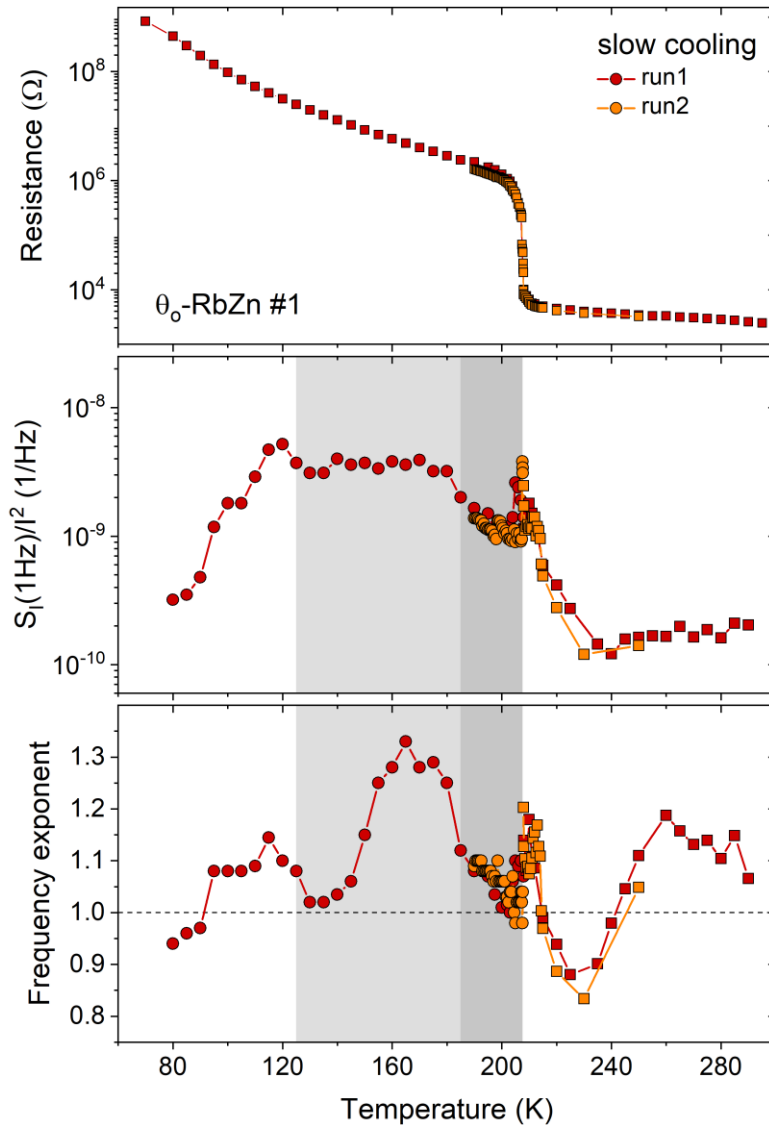
→ dependence on sample quality?

(samples from Prof. H. M. Yamamoto)

Heat pulse method

→ full avoidance of CO transition

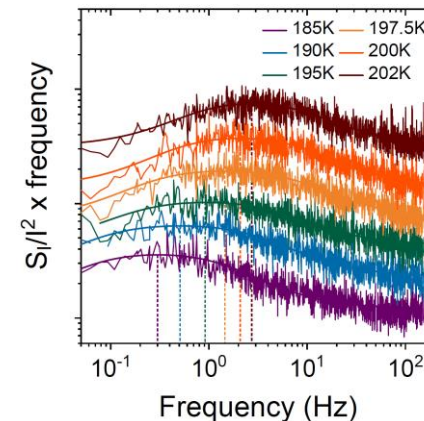
θ -RbZn: Charge-Crystal State



- $T_{CO} \sim 200\text{K}$
 - Large peaks at T_{CO}
 - Noise increases already above T_{CO}
- Slowing down of charge cluster and/or structural dynamics

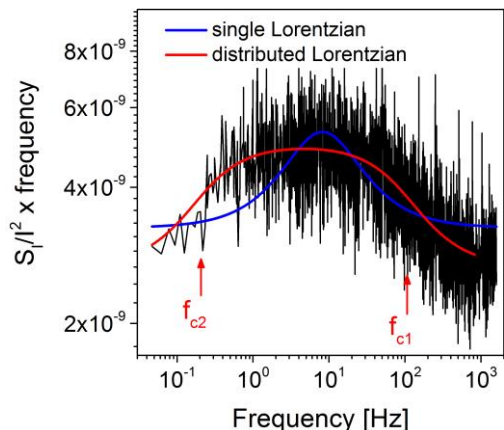
$T \sim 125\text{-}205\text{K}$:

- superimposed Lorentzian spectra



θ -RbZn: Charge-Crystal State

Dominating two-level fluctuators ($T \sim 125$ - 205 K)

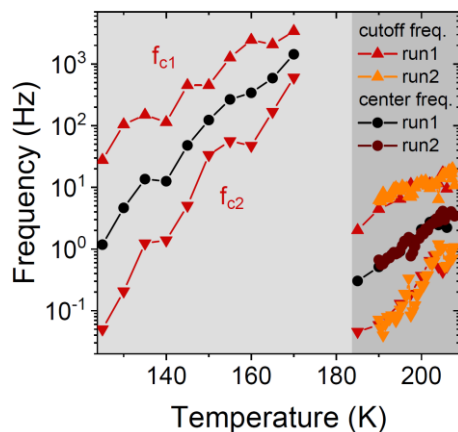


- Distributed Lorentzian model

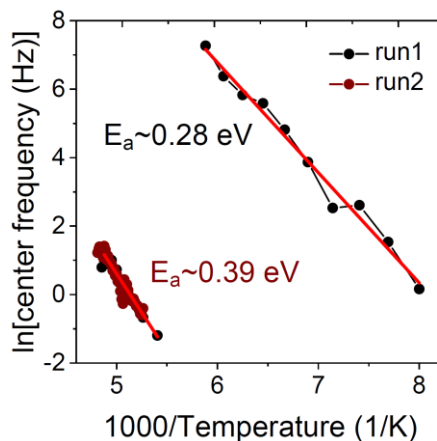
Kagawa et al., *Nat. Phys.* **9**, 419 (2013)

- Characteristic energies
- Slow and heterogeneous dynamics (cf. CL phase)

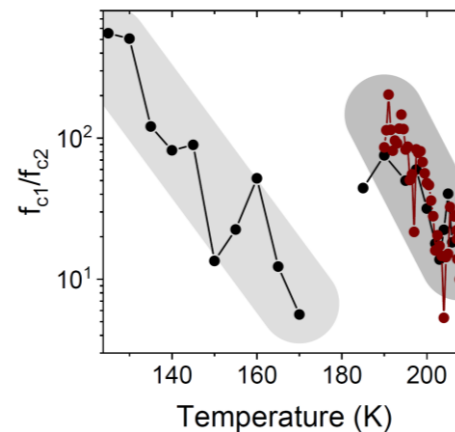
Frequencies:



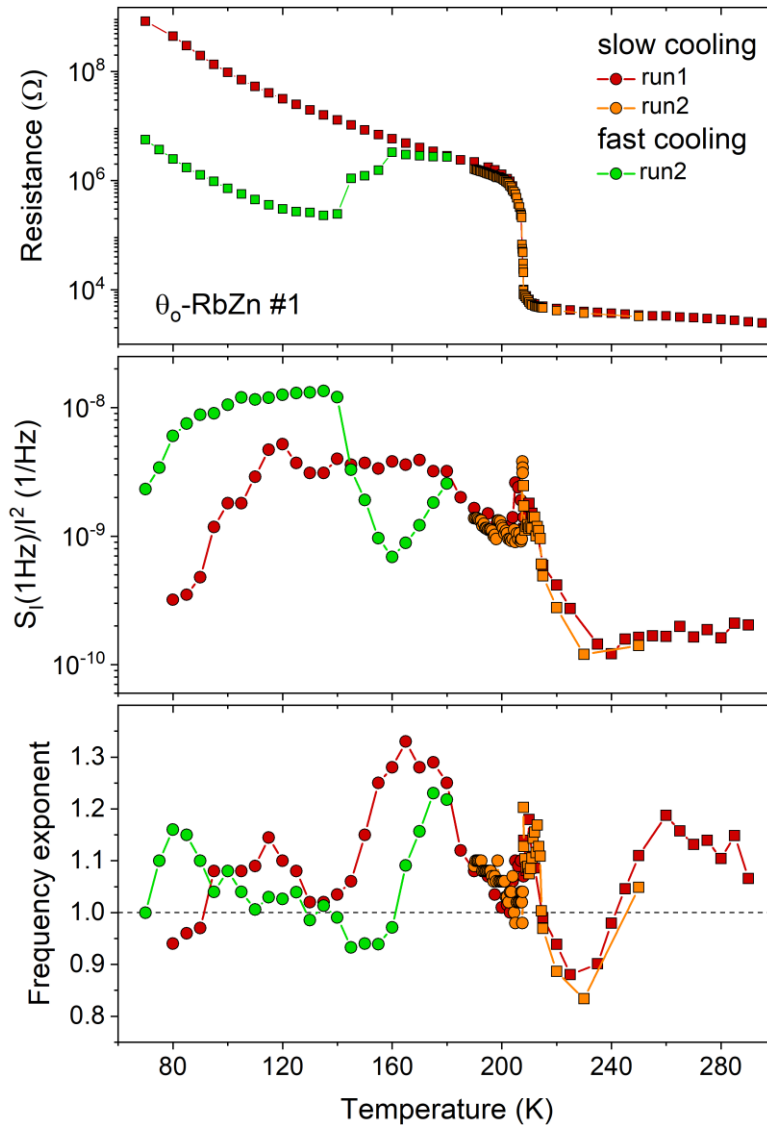
Energies:



Linewidth:



θ -RbZn: Charge-Glass State



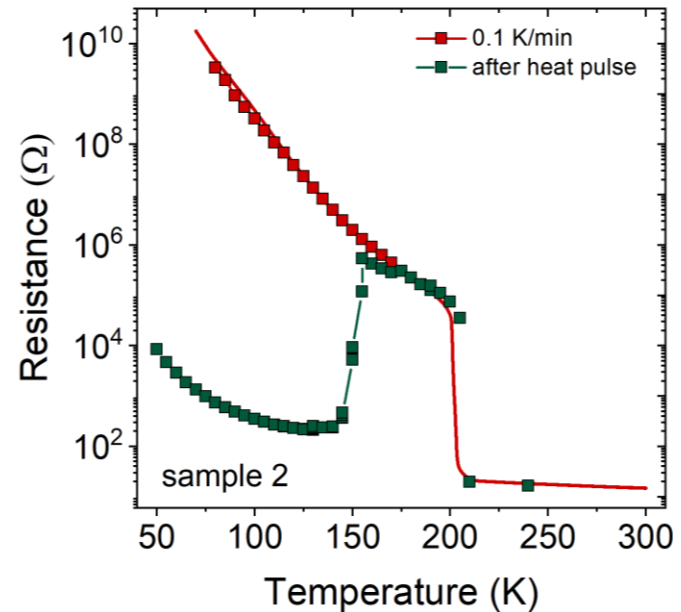
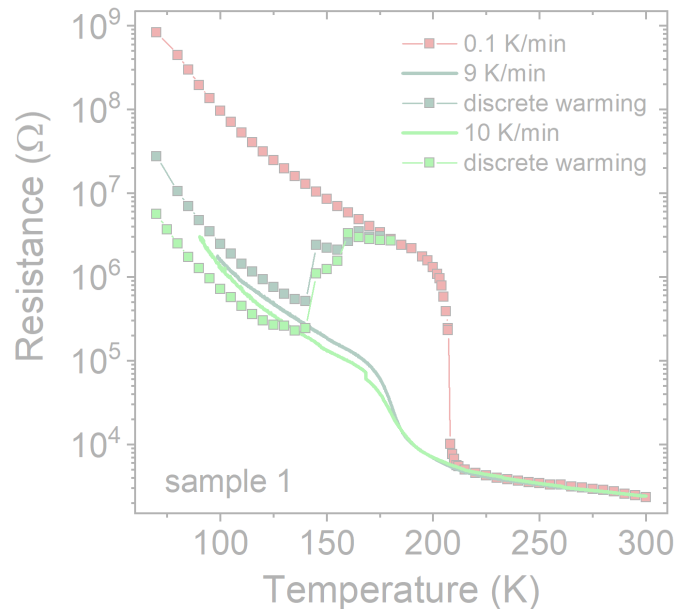
Partly quenched state

- Relaxation at $T \sim 145\text{K}$
- Complex noise behavior:
Under-/overshoot of noise magnitude

→ Due to mixed/inhomogeneous state?

θ -RbZn: Charge-Glass State

Fully quenched state



Critical cooling rate $q_c \sim 5\text{K/min}$?!

→ dependence on sample quality?

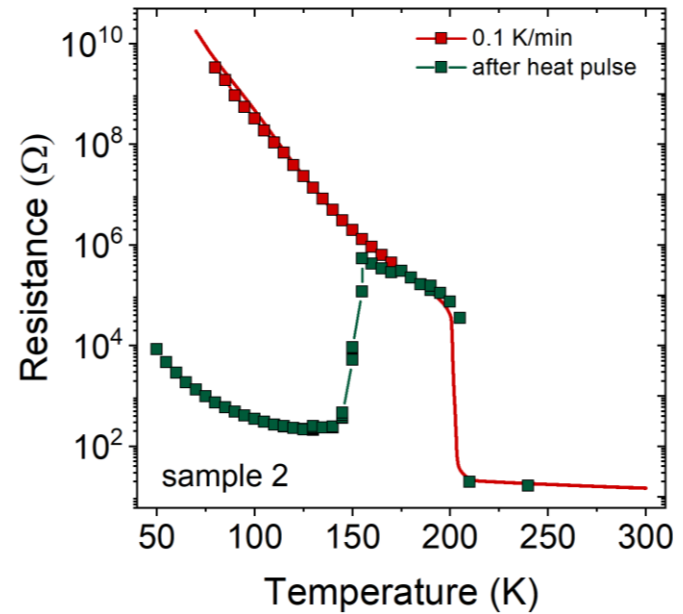
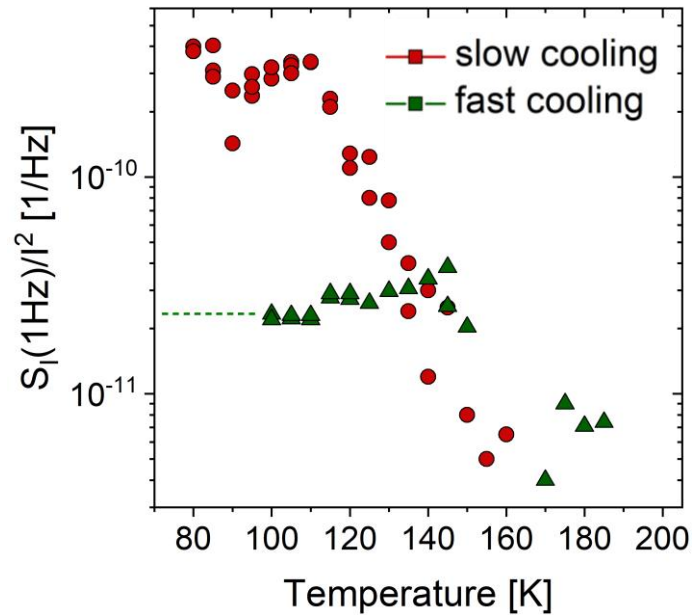
(samples from Prof. H. M. Yamamoto)

Heat pulse method

→ full avoidance of CO transition

θ -RbZn: Charge-Glass State

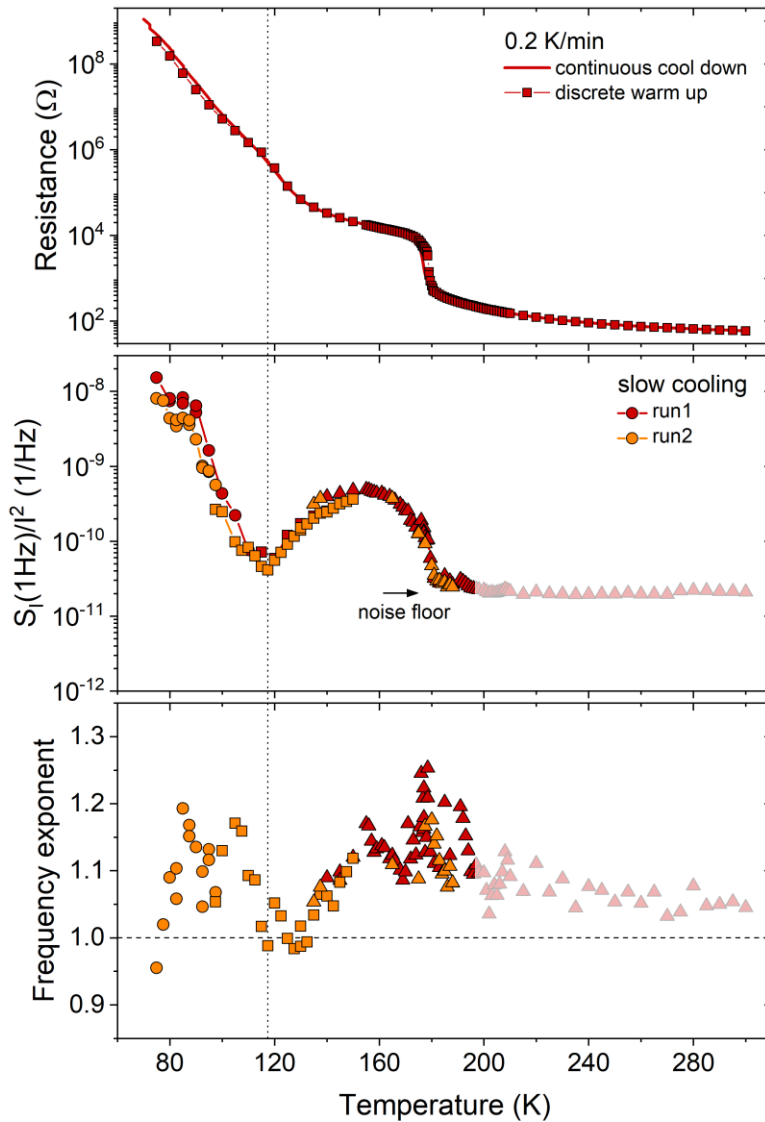
Fully quenched state



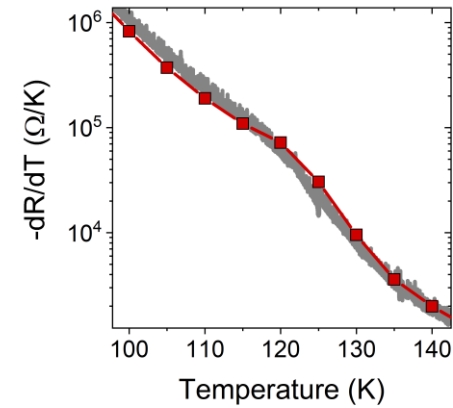
→ Noise in CG state much lower than in CC state

cf. θ_m -TiZn

θ_m -TlZn: Charge-Crystal State

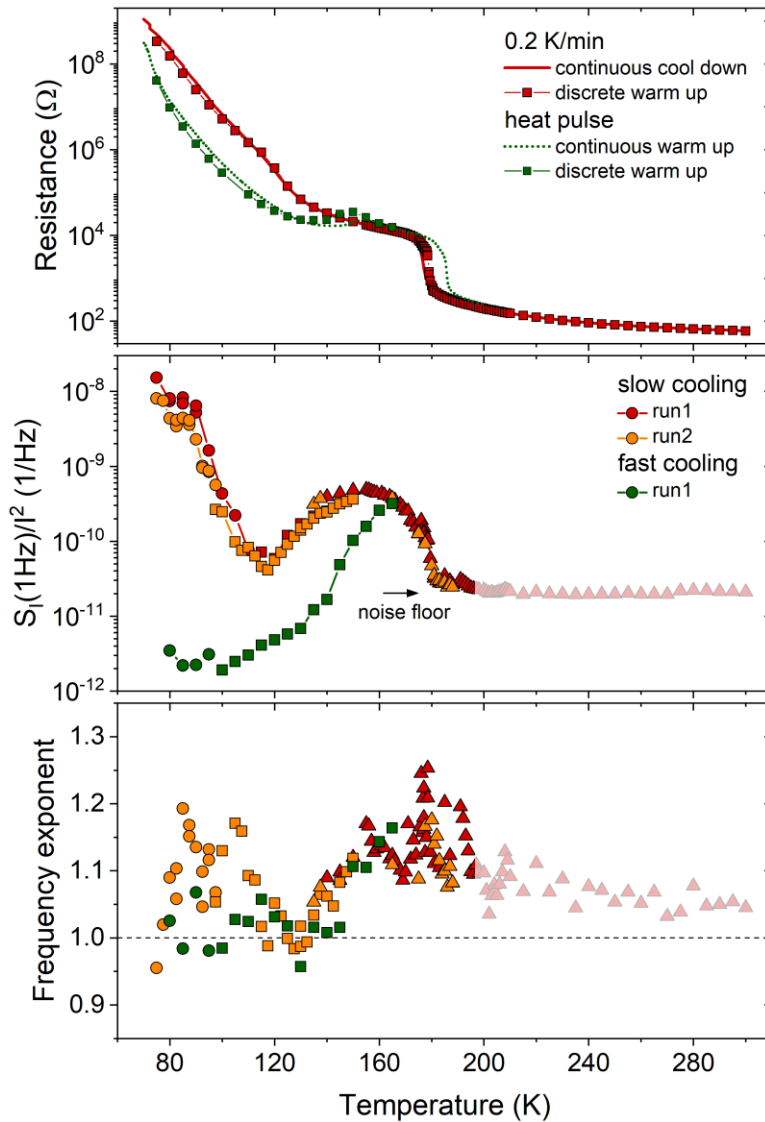


- $T_{CO} \sim 175\text{K}$
- Resistance anomaly at $T^* \sim 120\text{K}$



- Broad noise maximum around $T \sim 150\text{K}$
- Large peaks at transition

θ_m -TiZn: Charge-Glass State



- Heat pulse suppresses CO transition

- Relaxation at $T \sim 110\text{K}$

- Much lower noise magnitude in CG

→ Inhomogeneous current distribution in CC

→ Larger cluster size in CC

→ Lower energy of fluctuators in CG

Summary

Investigation of the charge dynamics in θ -(ET)₂MM'(SCN)₄ reveals

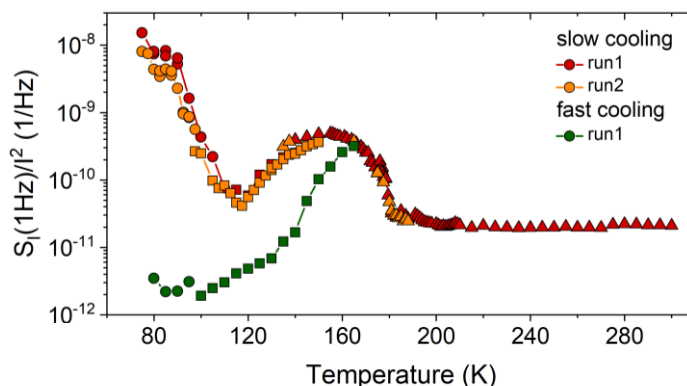
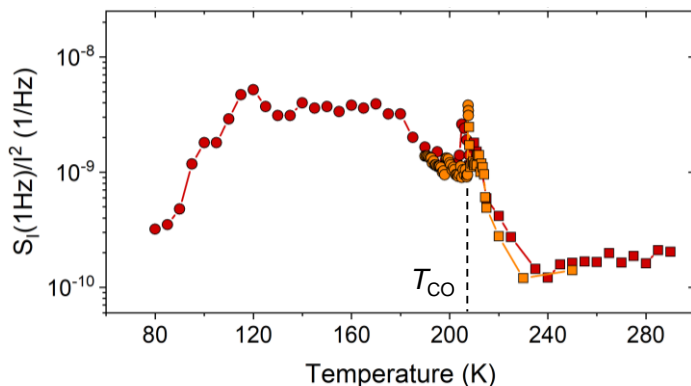
→ structural glass transition in θ -(ET)₂CsM'(SCN)₄

Thomas & Saito et al., *PRB* **105**, L041114 (2022)

→ strong fluctuations at phase transition

→ large differences in CC and CG state

Thomas et al., *PRB* **105**, 205111 (2022)



Open questions:

- Relation between lattice and electron dynamics?
- Influence of disorder by x-ray irradiation?
- Role of ferroelectricity? Current-dependent noise

cf. β' -(ET)₂ICl₂, κ -(ET)₂Cu[N(CN)₂]Cl, κ -(ET)₂Hg(SCN)₂Cl, κ -(BETS)₂Mn[N(CN)₂]₃

Müller et al., *PRB* **102**, 100103(R) (2020) Thomas et al., *Phys. Status Solidi B*, 1800746 (2019)

