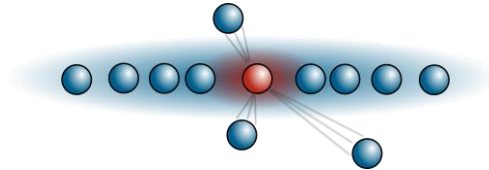


Equilibrium-, non-equilibrium & steady-state properties of quantum impurities in 1D Bose gases



Michael Fleischhauer

University of Kaiserslautern

Thanks to:



Martin Will



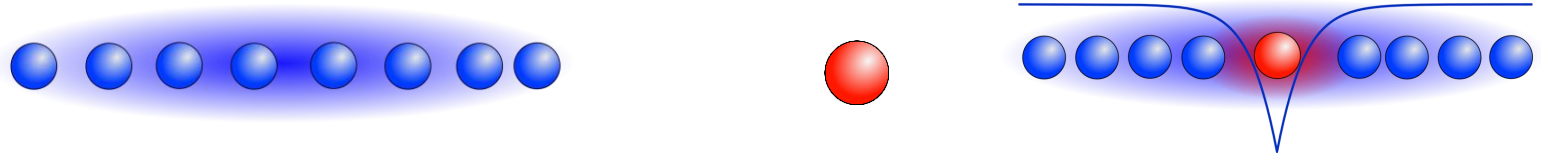
Jonas Jager



Gregori Astrakharchik



Jamir Marino



$$\hat{H} = \underbrace{\int dx \hat{\phi}^\dagger(x) \left[-\frac{\partial_x^2}{2m} + \frac{g_{\text{BB}}}{2} \hat{\phi}^\dagger(x) \hat{\phi}(x) \right] \hat{\phi}(x)}_{\text{homogeneous Bose gas}} + \underbrace{\frac{\hat{p}^2}{2M}}_{\text{single impurity}} + \underbrace{g_{\text{IB}} \int dx \hat{\phi}^\dagger(x) \delta(x - \hat{r}) \hat{\phi}(x)}_{\text{interaction}}$$

homogeneous Bose gas
2nd quantized

single impurity
1st quantized

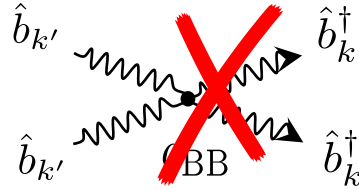
interaction

- weakly interacting Bose gas $\gamma \leq 1$ ($\xi n > 1$)
- heavy impurity $M \geq m$

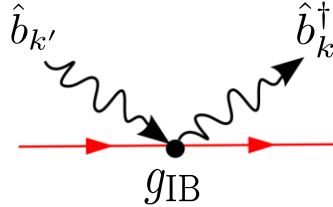
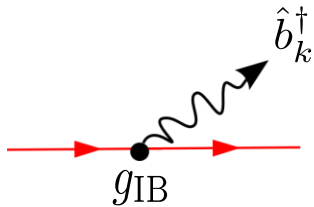
- **strong-coupling regime**
 $g_{\text{IB}} \gg g_{\text{BB}} \xi n$

1. Bogoliubov theory of Bose gas w/o impurity

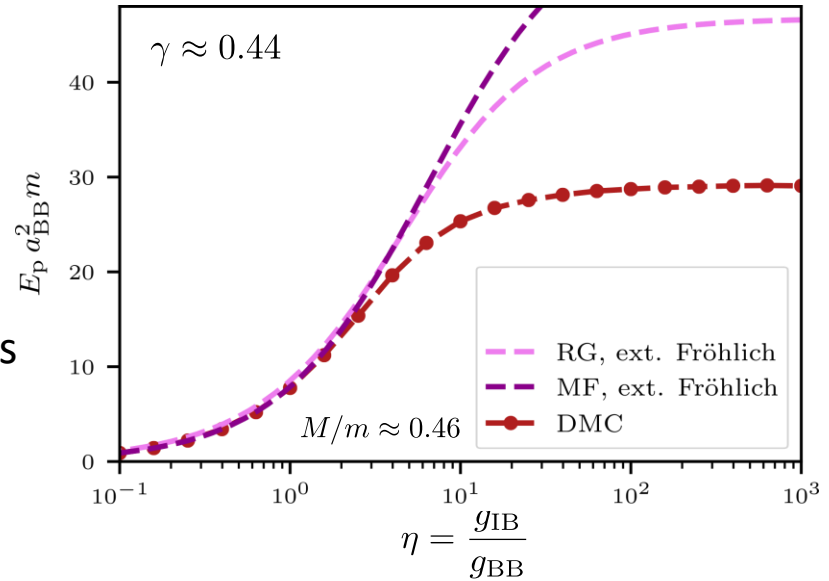
$$\hat{\phi}(x) = \sqrt{n} + \hat{\xi}(x)$$



2. impurity: scattering & generation of phonons



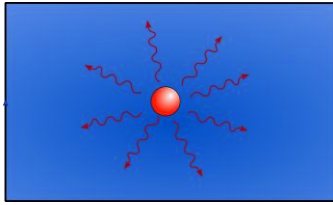
polaron energy:



Grusdt *et al.* *New J.Phys.*19, 103035 (2017)
(*experiment: Catani et al. Phys.Rev.A* 85, 023623 (2012))

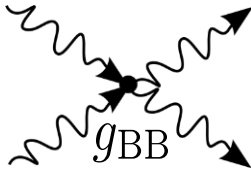
extended Fröhlich model

$$\hat{\phi}(x) = \sqrt{\bar{n}} + \hat{\xi}(x)$$



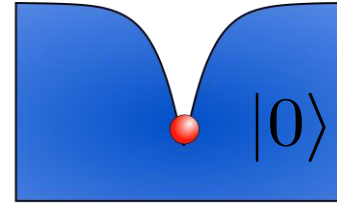
impurity generates phonons

$$\bar{n}_{\text{ph}} \sim g_{\text{IB}}^2 n$$



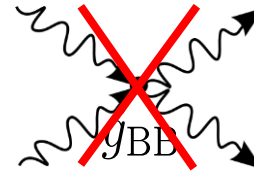
deformed condensate
quantum corrections

$$\hat{\phi}(x) = \phi(x) + \hat{\xi}(x)$$



phonons in deformed condensate

$$\bar{n}_{\text{ph}} = 0$$





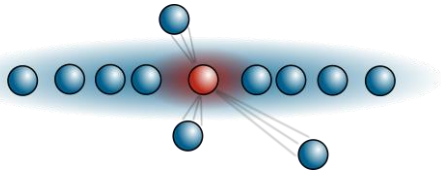
- Condensate deformation and modified phonons



- Polaron interactions and bi-polarons

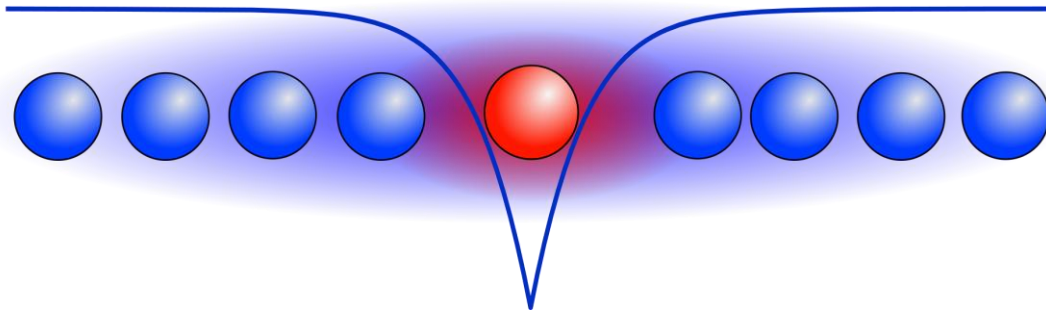


- Adiabatic & quench dynamics of a quantum impurity



- Control of superfluid flow by noisy impurities

Single Bose-polaron



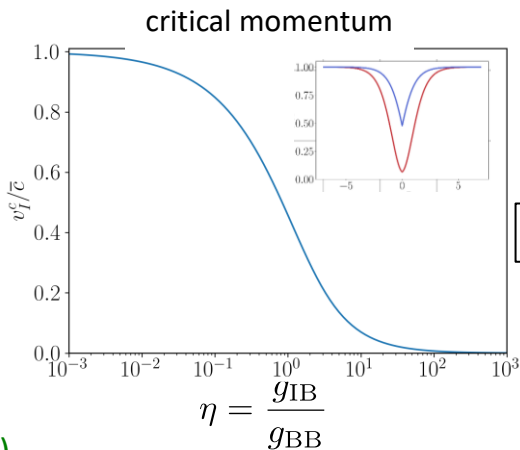
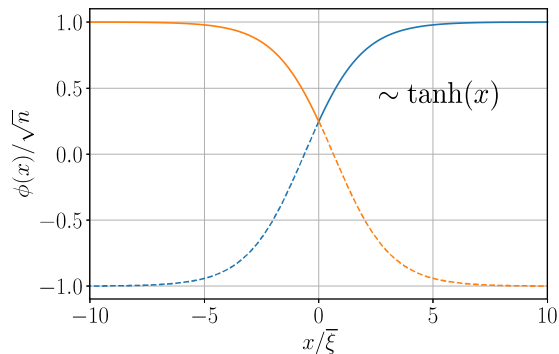
Jonas Jager, Ryan Barnett, Martin Will and Michael Fleischhauer
Phys. Rev. Res. **2**, 033142 (2020)

Bogoliubov with deformed phonons

- **Mean-field:** $\hat{\phi}(x) = \phi(x) \quad \tilde{m} = \frac{mM}{m+M}$

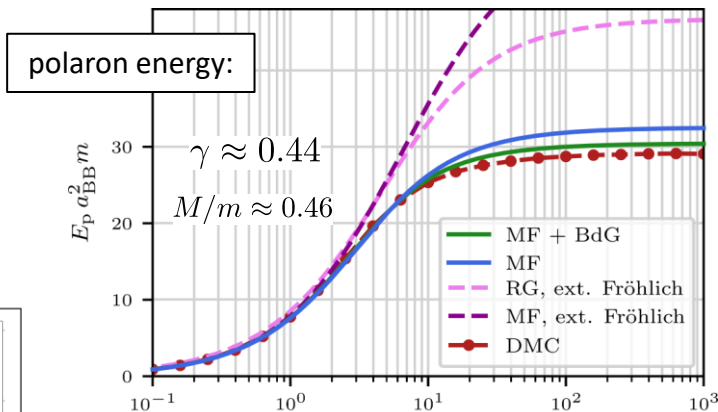
$$\left[-\frac{\partial_x^2}{2\tilde{m}} + \frac{i}{M} (p - P_{\text{BEC}}) \partial_x + g|\phi(x)|^2 + g_{\text{IB}}\delta(x) - \mu \right] \phi(x) = 0$$

⇒ solved by two cut grey solitons (2 solutions)

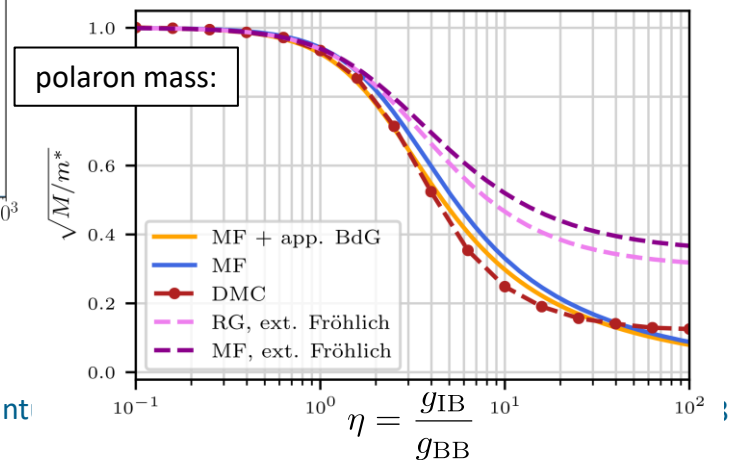


Hakim Phys.Rev.E **55**, 2835 (1996)

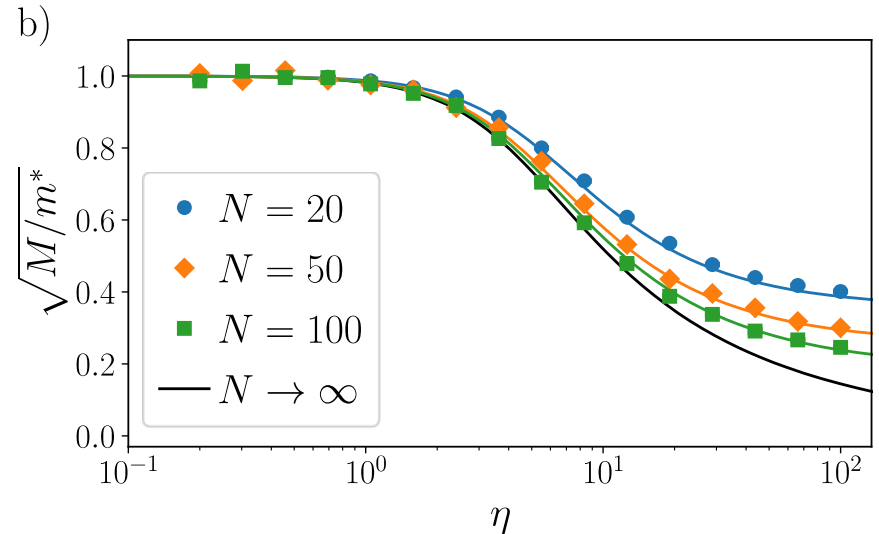
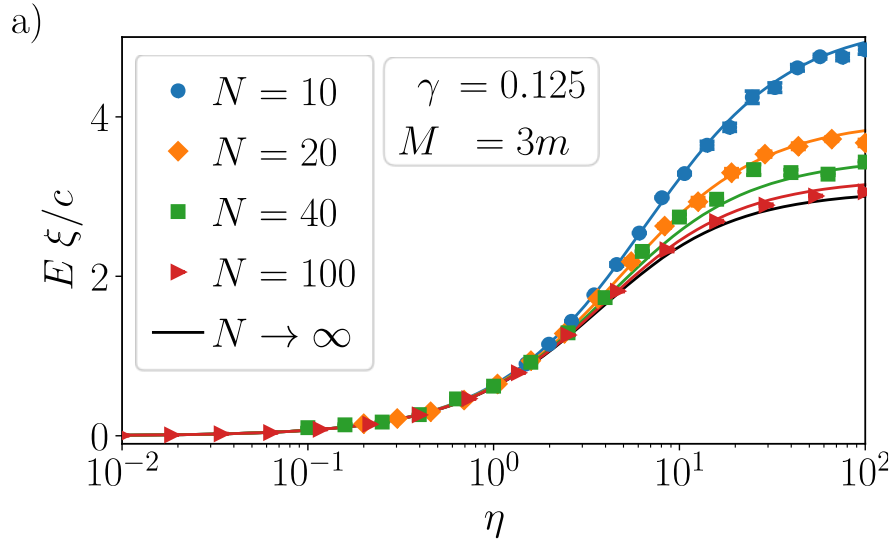
- **BdG:** $\hat{\phi}(x) = \phi(x) + \hat{\xi}(x)$

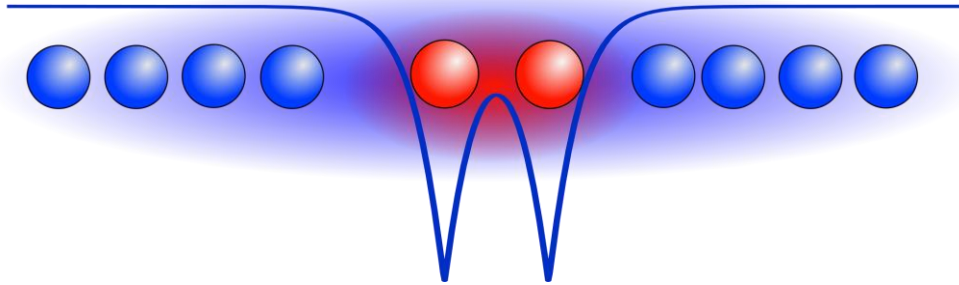


$p = 0$



- Finite-size effects:



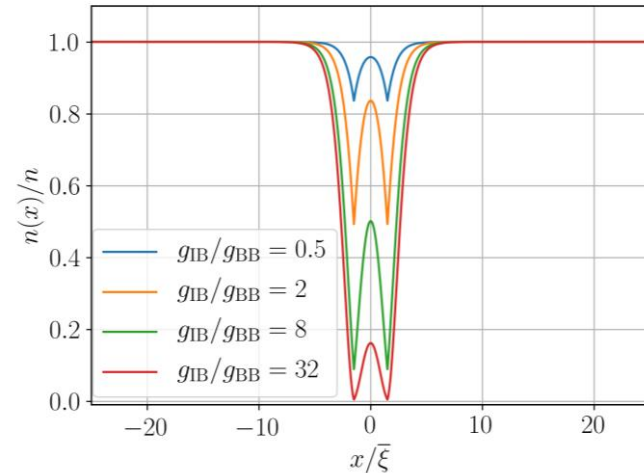
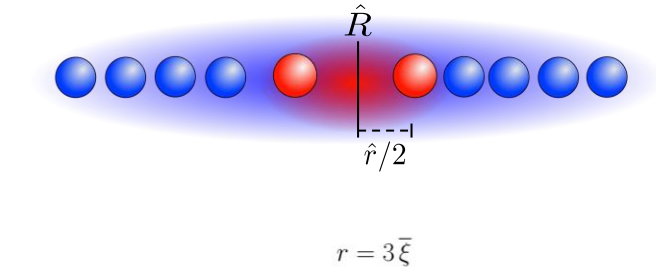
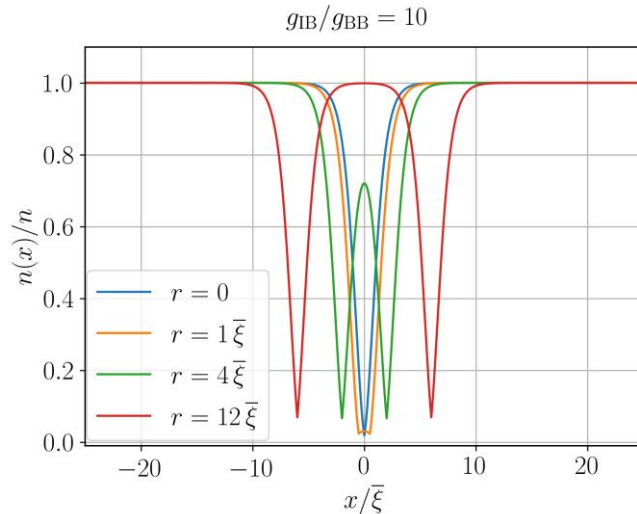


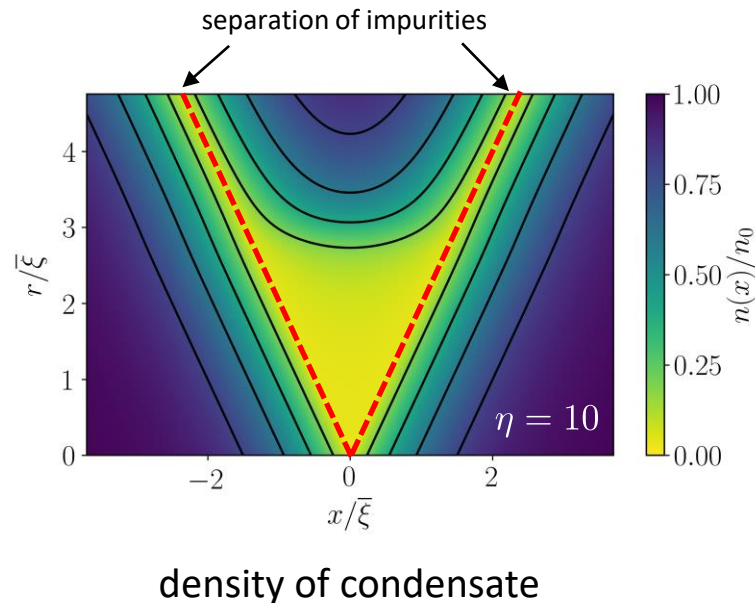
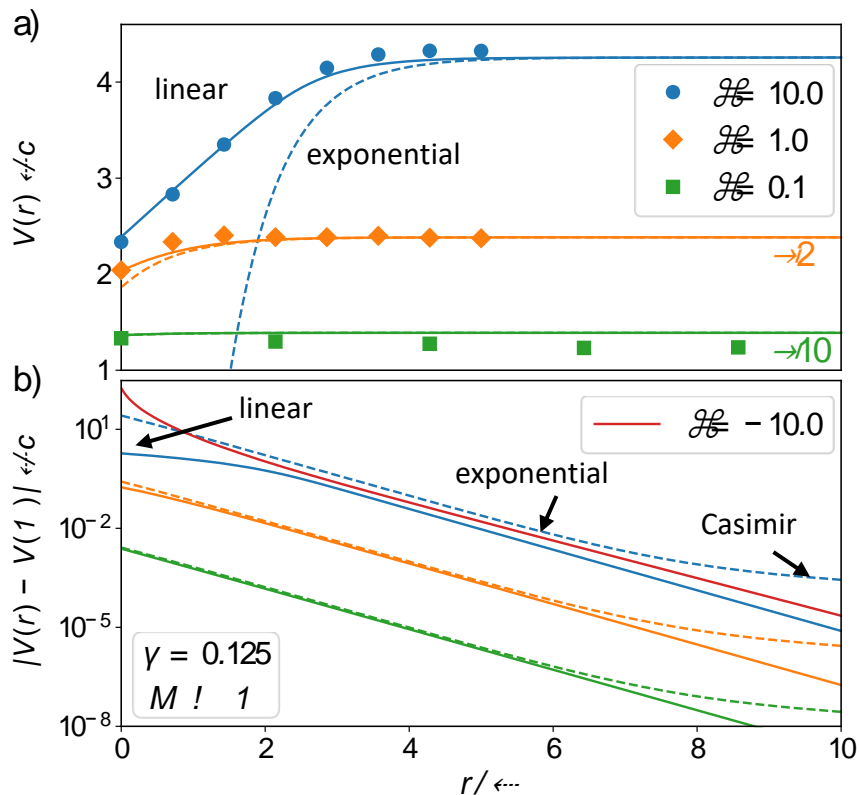
Martin Will, Grigori Astrakharchik and Michael Fleischhauer
Phys. Rev. Lett. **127**, 103401 (2021)

relative impurity coordinates: \hat{p}, \hat{r}
center of mass: \hat{P}, \hat{R}

- Lee-Low-Pines removes \hat{P}
- **Born-Oppenheimer:** $M \gg m$

Mean-field for $P = 0$:





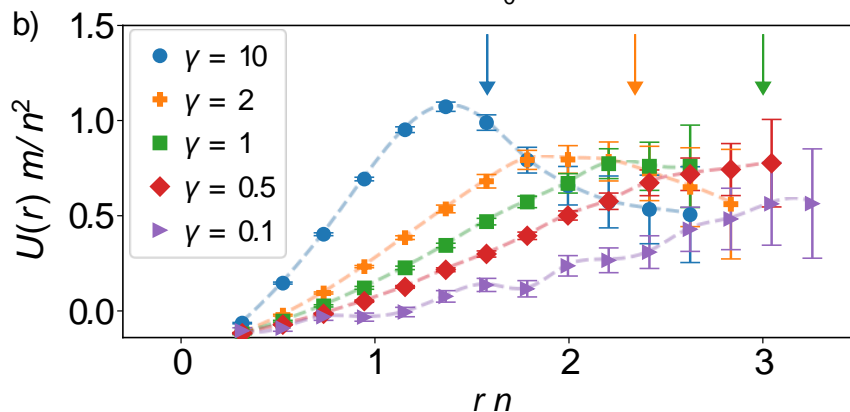
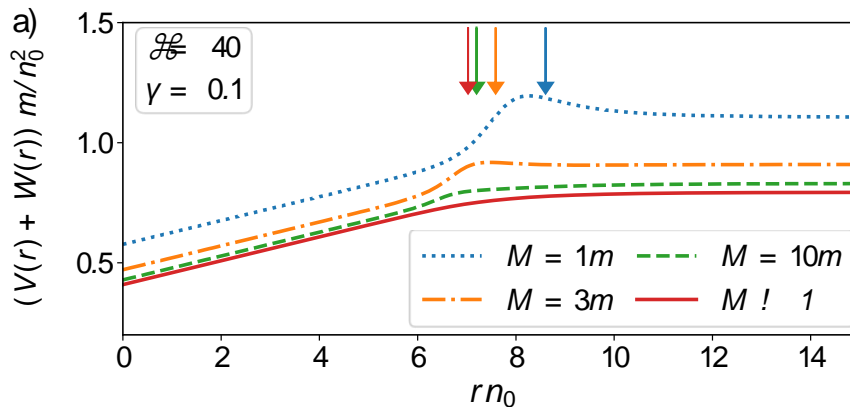
$$\left[-\frac{1}{M} \partial_r^2 + V(r) + W(r) - E \right] \Psi(r) = 0$$

- additional potential

$$W(r) = \frac{1}{M} \int dx |\partial_r \phi(x, r)|^2$$

potential maximum

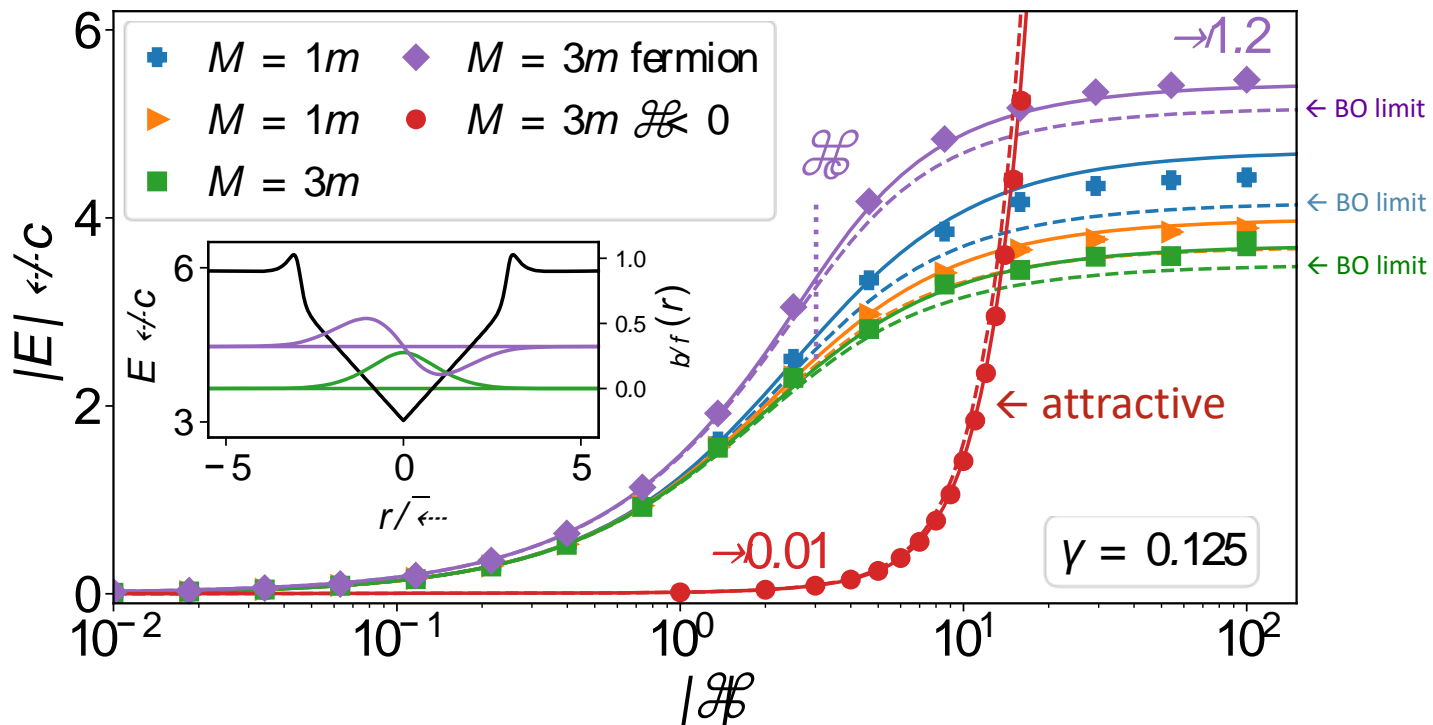
$$r_{\max} = \frac{\pi}{\sqrt{2m_r \mu}}$$

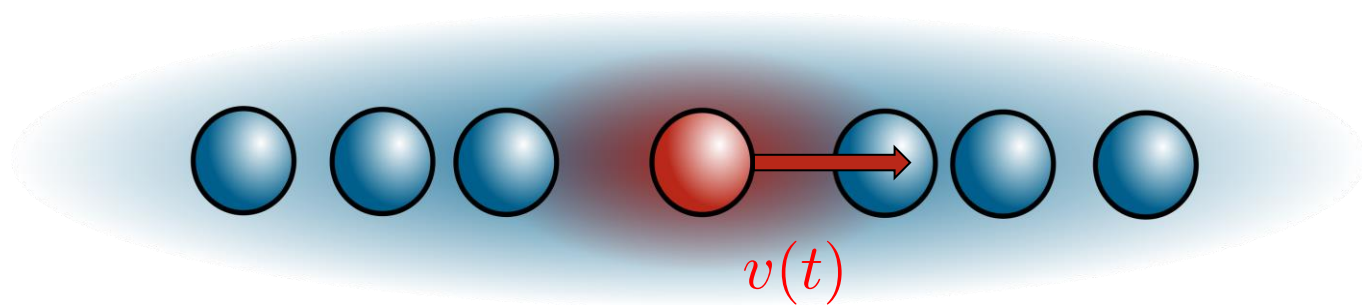


QMC



Bi-polaron bound states

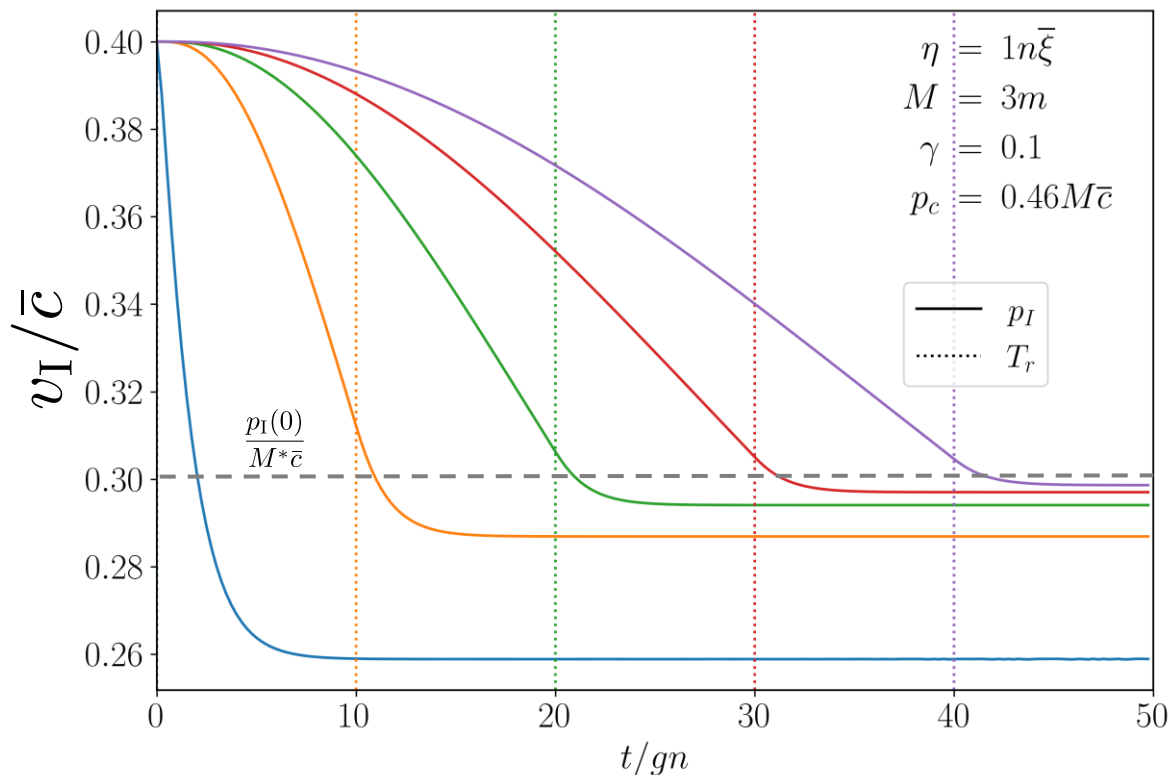




Martin Will and Michael Fleischhauer
in preparation

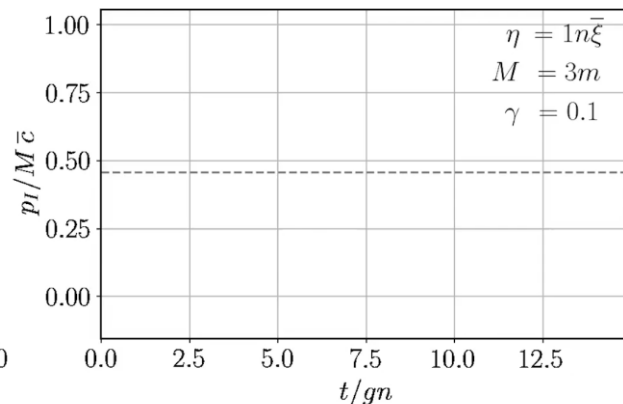
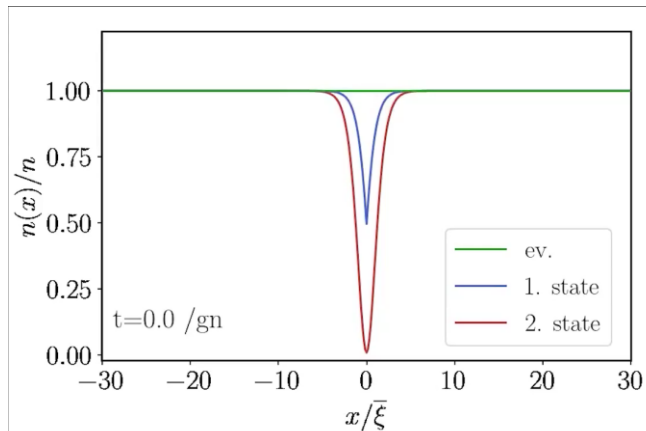
$$v_I < c$$

linear ramp of $g_{IB}(t)$
in time T_c

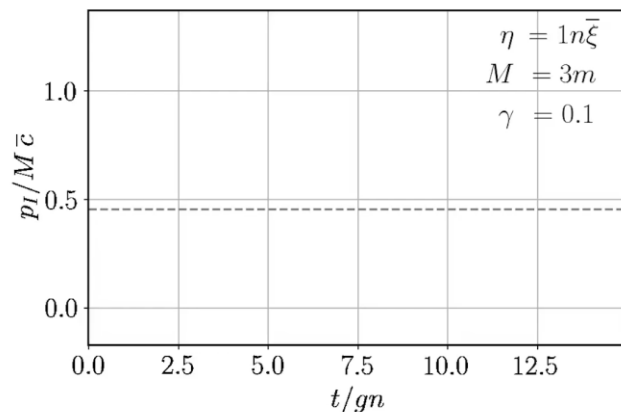
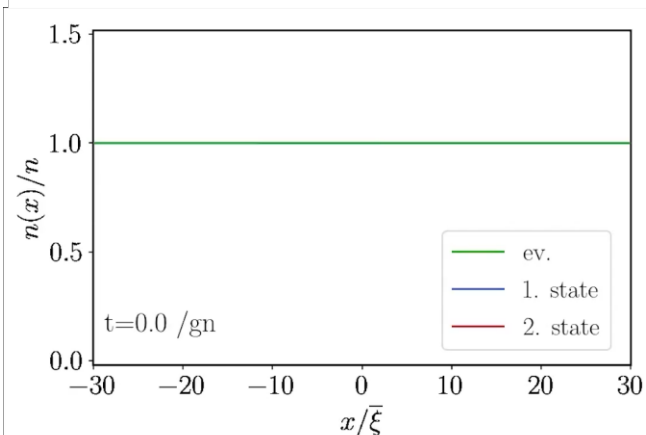


$$M = 3m$$

$$p_I(0) = 0.1M\bar{c}$$



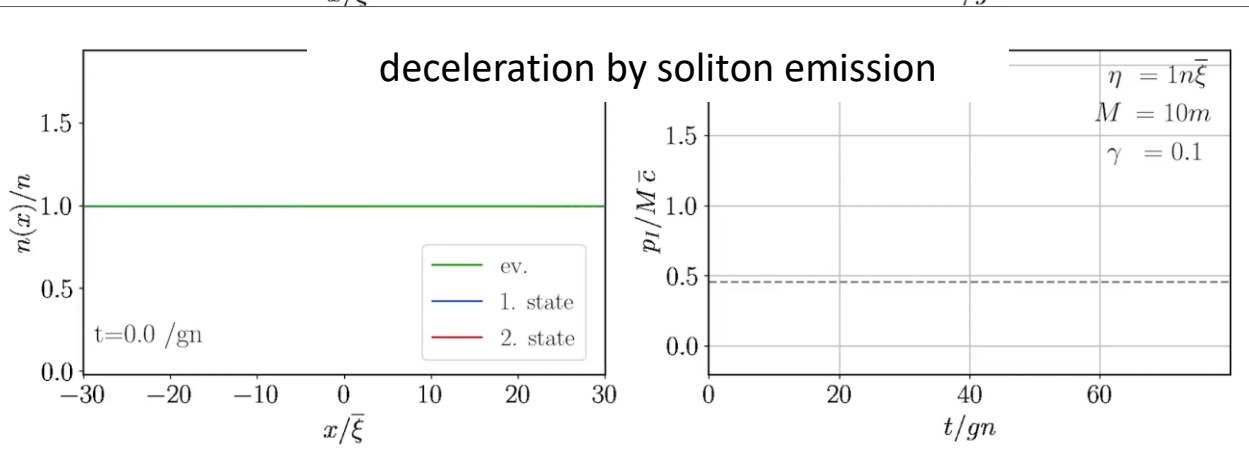
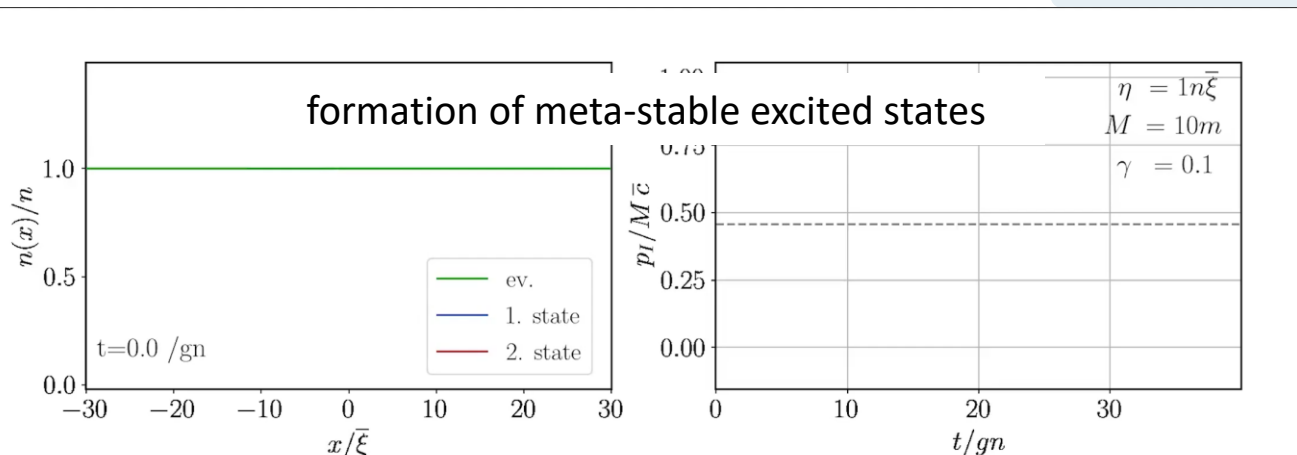
$$p_I(0) = 1.3M\bar{c}$$

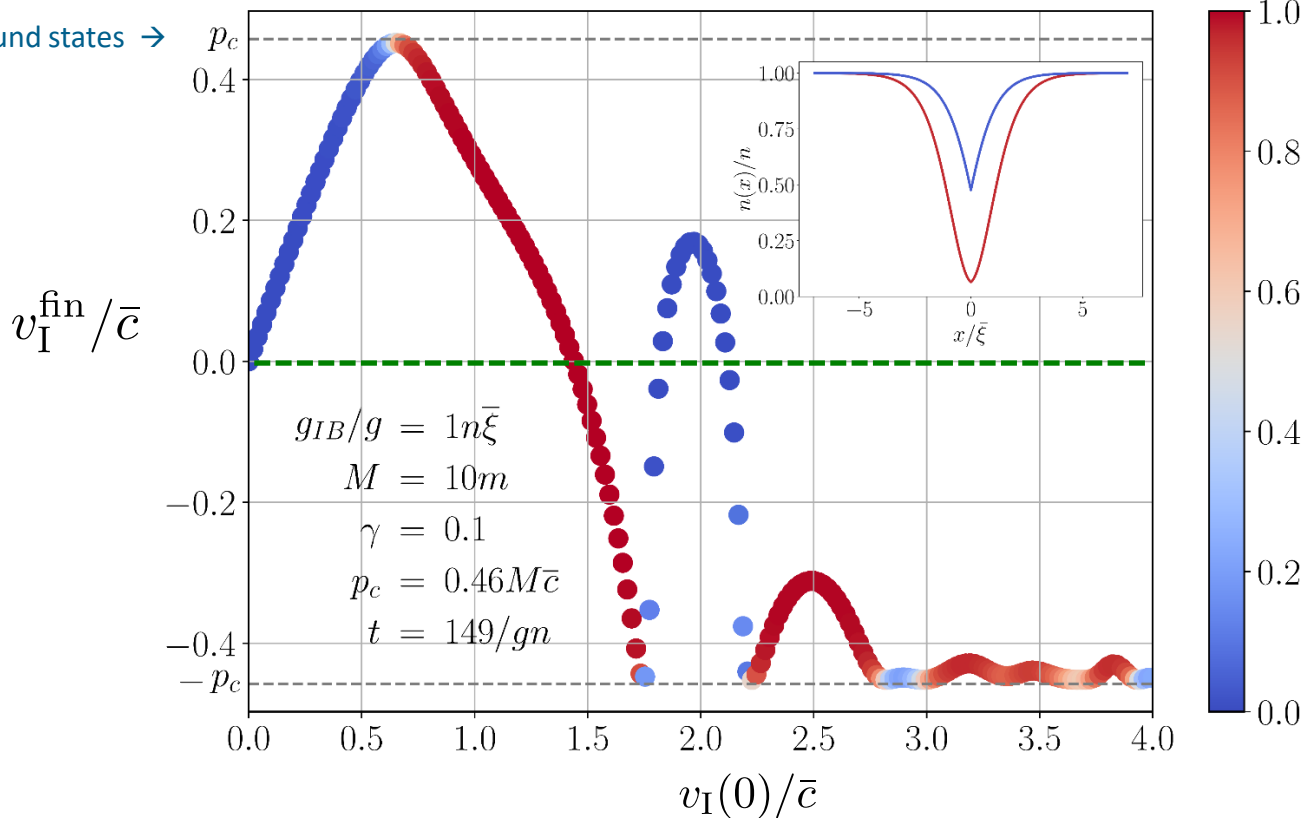


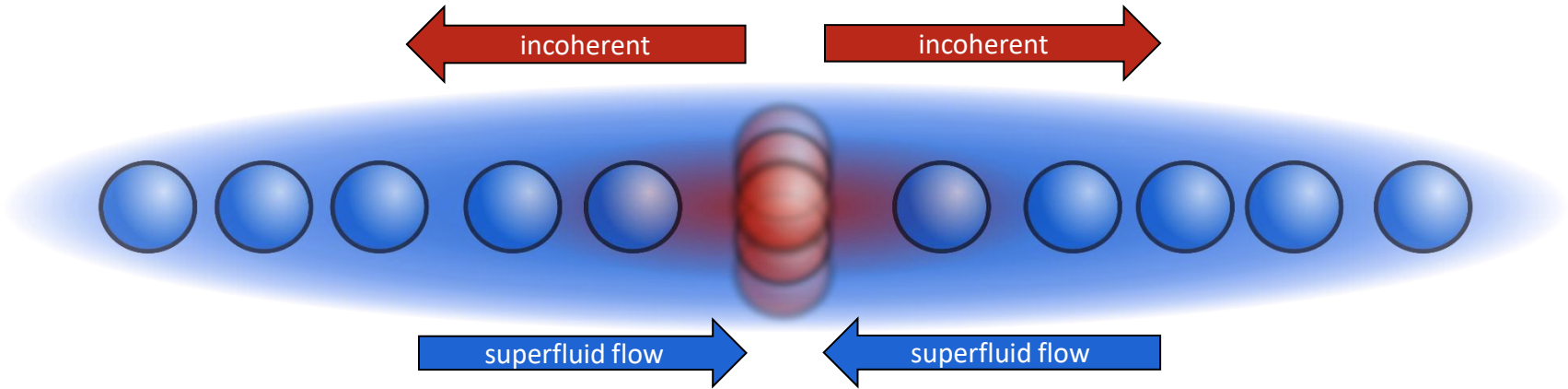
$$M = 10m$$

$$p_I(0) = 1 M \bar{c}$$

$$p_I(0) = 2 M \bar{c}$$



existence of bound states \rightarrow 



Martin Will, Jamir Marino and Michael Fleischhauer
in preparation

D. A. Zezyulin, V. V. Konotop, G. Barontini and H. Ott, *Phys. Rev. Lett.* **109**, 020405

D. Sels and E. Demler, *Ann. Phys.* **412**, 168021 (2020)

$$i\partial_t\Phi(x,t) = \left[-\frac{1}{2m}\partial_x^2 + iv\partial_x + g|\Phi(x,t)|^2 + \eta(t)V(x) \right] \Phi(x,t)$$

Random variable: $\overline{\eta(t)} = 0$ $\overline{\eta(t)\eta(t')} = \delta(t-t')$

- **Stratonovich** stochastic differential equation:

$$d\Phi(x,t) = -i \left[-\frac{1}{2m}\partial_x^2 + iv\partial_x + g|\Phi(x,t)|^2 \right] \Phi(x,t)dt - i\Phi(x,t)V(x)dW$$

Stochastic increment: $dW = \eta(t)dt$

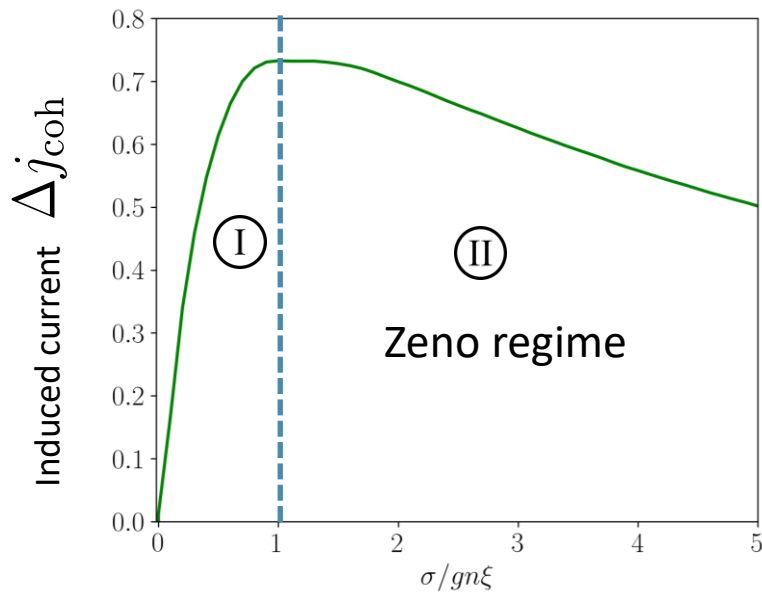
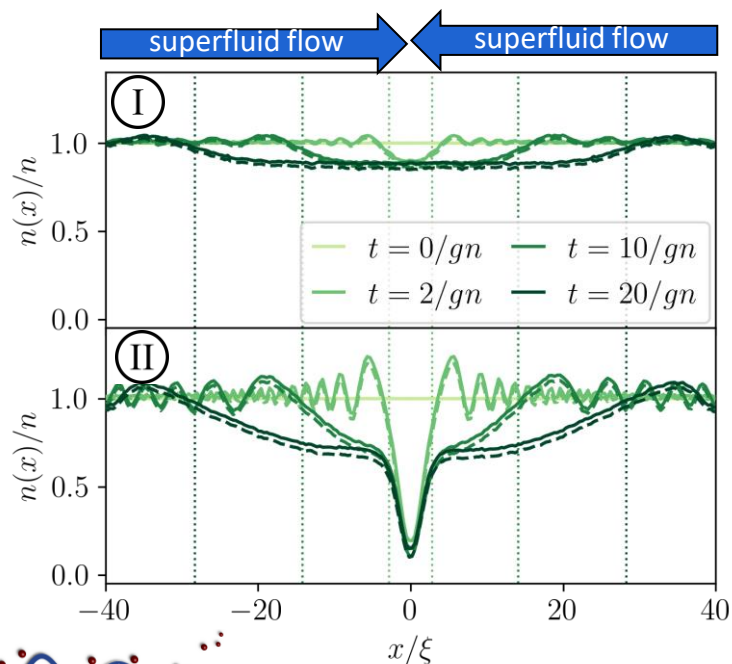
Noise correlated with field: $\overline{\Phi(x,t)dW} \neq 0$

Stationary impurity $v = 0$

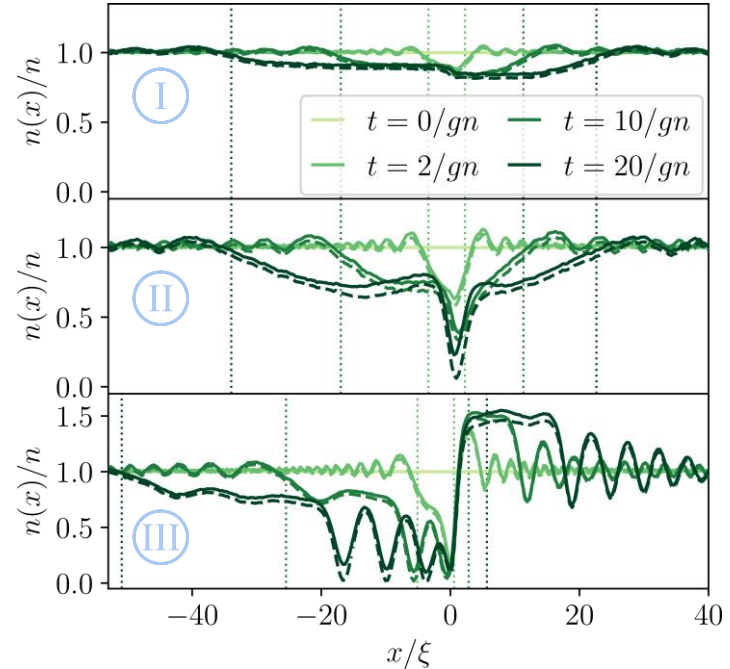
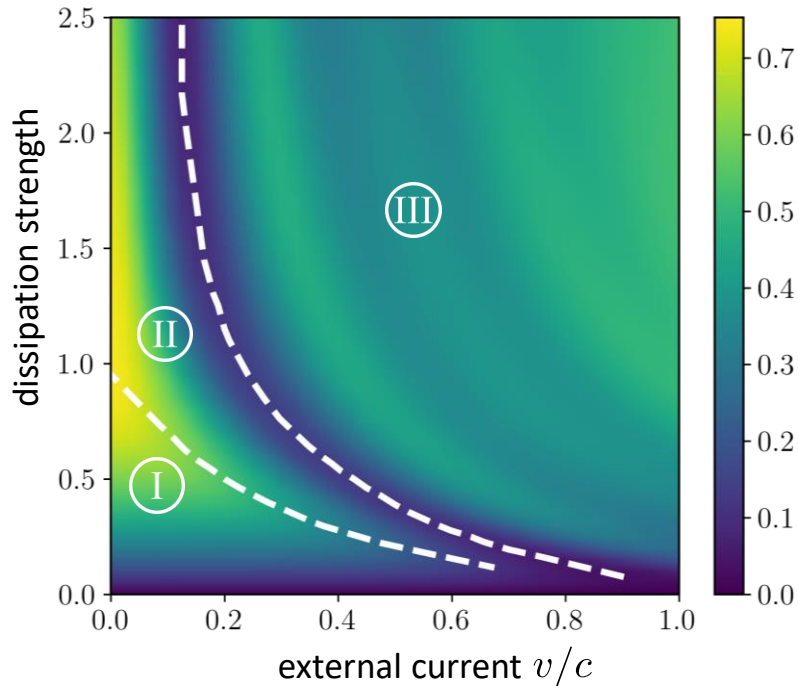
- average field (superfluid amplitude)

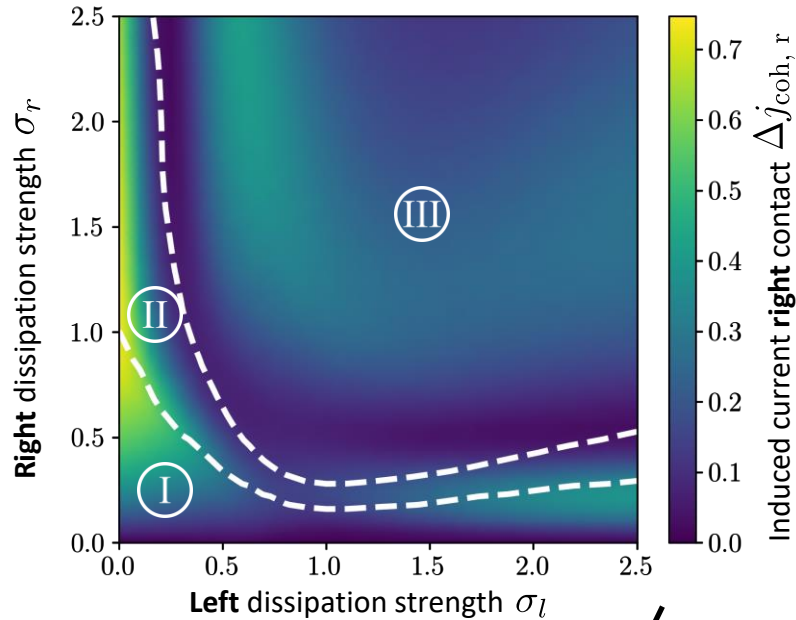
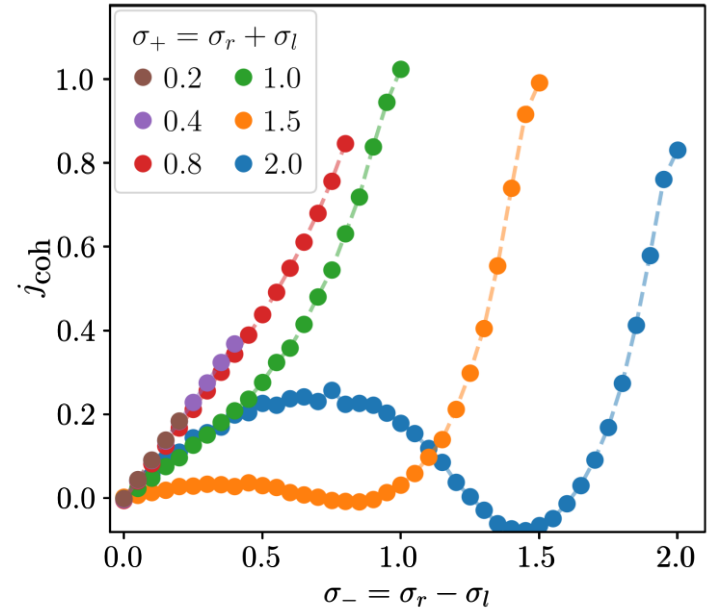
$$i\partial_t \overline{\Phi(x,t)} = -\frac{\partial_x^2}{2m} \overline{\Phi(x,t)} + g \overline{|\Phi(x,t)|^2 \Phi(x,t)} - \frac{i}{2} \overline{V(x)^2 \Phi(x,t)}$$

$$V(x)^2 = 2\sigma\delta(x)$$

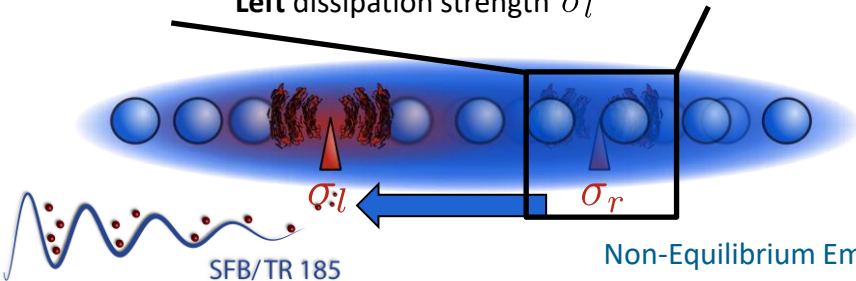


- New soliton regime III



Current **between** contacts

Black line:

Single contact at dissipation strength σ_- 

Bose polaron

- Backaction of impurity to condensate: deformed condensate and modified Bogoliubov phonons

Bi-polaron

- BO & beyond BO: linear binding potential on short distances

Dynamics of impurity injected into BEC

- slow down also for $v < c$ even in adiabatic regime
- slow down by (i) density wave emission (ii) soliton emission
- momentum reversal possible $|p_{\text{fin}}| \leq p_c$

Noisy impurity

- Control of superfluid and incoherent flow
- Different dynamical regimes: (i) linear response (ii) Zeno / negative differential conductivity (iii) soliton regime



$$\begin{aligned}
 V(r, \eta) = & gn^2 r \left(\frac{1}{2} - \frac{4 + 2\nu}{3(\nu + 1)^2} \right) \\
 & + \frac{4}{3} gn^2 \bar{\xi} \frac{1}{\sqrt{1 + \nu}} \left\{ 2 E(\operatorname{am}(u, \nu), \nu) - \frac{3\sqrt{\nu}}{2 + 2\nu} \operatorname{cd}(u, \nu) \left[1 + \nu + \frac{2}{3} \nu \operatorname{cd}(u, \nu)^2 \right] \right. \\
 & \left. + \sqrt{2\nu + 2} - \frac{\nu}{1 + \nu} \operatorname{cd}(u, \nu) \operatorname{sn}(u, \nu) \left[\nu \operatorname{cd}(u, \nu)^2 + 2\nu + 1 \right] \right\}
 \end{aligned}$$