

Bosonic Quantum Error Correction

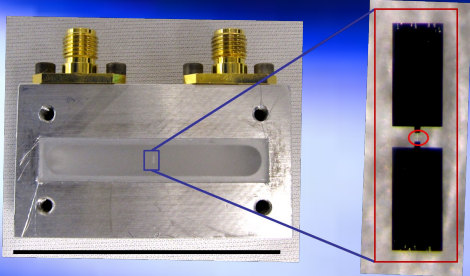
Liang Jiang
University of Chicago

SPICE Workshop on Non-Equilibrium Emergence in Quantum Design
June 22, 2022

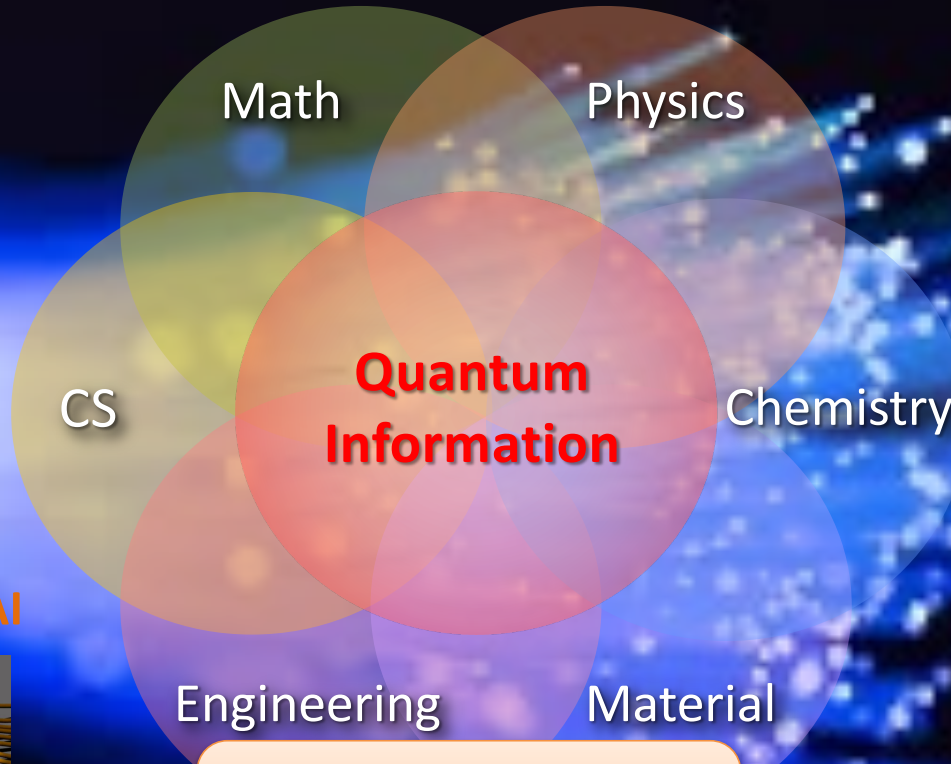
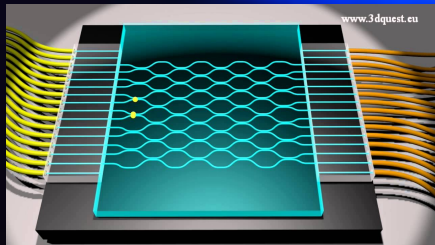


Various Applications of Bosonic Modes for Quantum Information

Quantum Computation



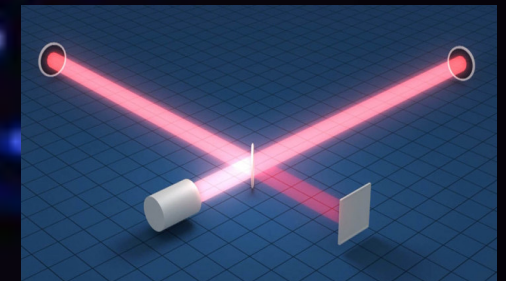
Quantum Simulation/Sampling/AI



Quantum Communication



Quantum Metrology



**Key Challenge:
Excitation loss error!**

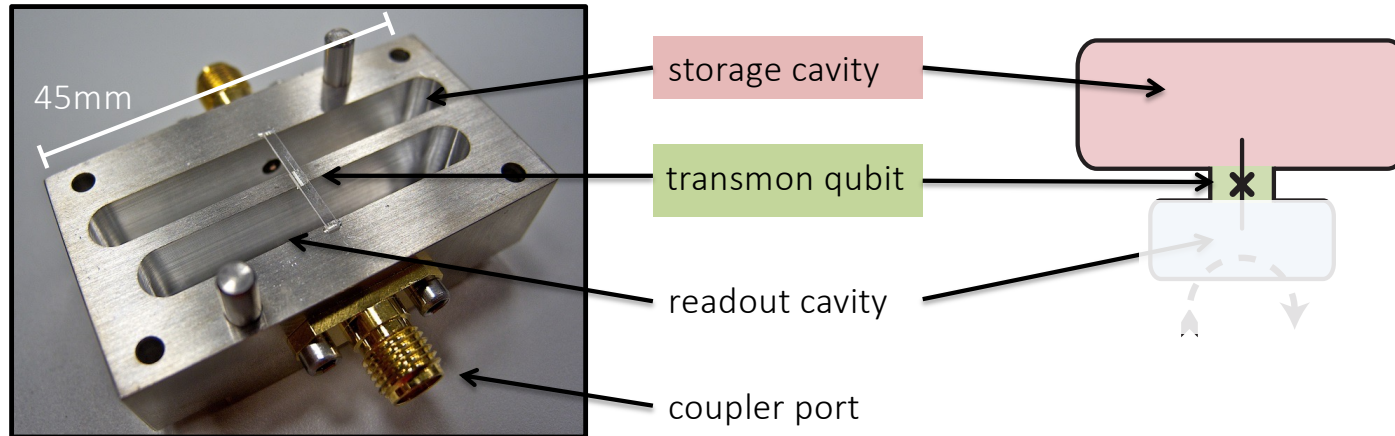
E.g., photons, phonons, magnons, atomic ensembles, ...

Quantum Control of Bosonic Modes

Key Challenges:

- How to achieve **universal control** of individual bosonic modes?
- How to **correct excitation loss errors** for quantum computing?
- How to implement quantum non-demolition (QND) parity measurement?

Superconducting Cavity-Qubit Module



Effective Hamiltonian:

$$H = \omega_s a^\dagger a + \omega_q |e\rangle\langle e| - \chi_s a^\dagger a |e\rangle\langle e|$$

- Long lived SC cavity ($T > 10$ ms $\gg T_{\text{qubit}} \sim 100$ us)
- Fast QND readout of the qubit (e.g., $F > 99.5\%$ in 300 ns)
- Adaptive control with FPGA

Strong Dispersive Regime:

$$\chi_s (\sim \text{MHz}) \gg \kappa_s, \gamma_q (\sim \text{kHz})$$

χ_s : dispersive coupling strength
 κ_s : photon decay rate
 γ_q : transmon decoherence rate

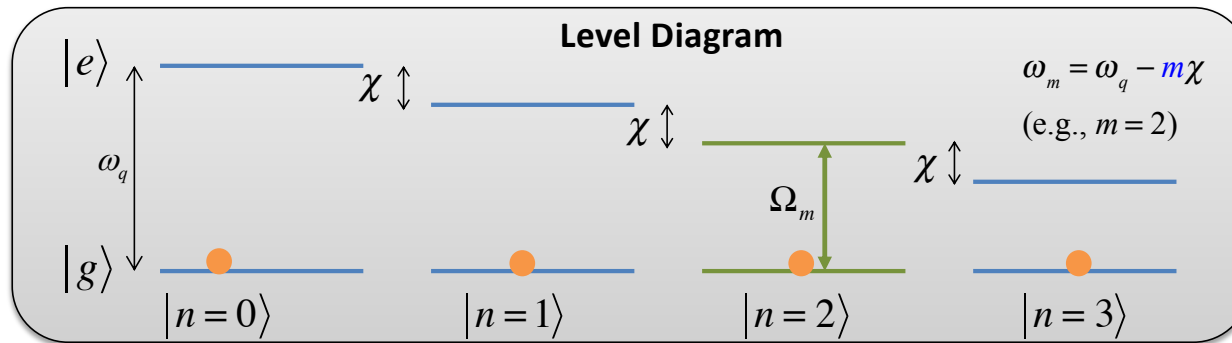
Control of Cavity-Qubit System

$$H = \omega_s a^\dagger a + \varepsilon(t) a^\dagger e^{i\omega_s t} + h.c.$$

Cavity drive

$$+ (\omega_q - \chi_s a^\dagger a) |e\rangle\langle e| + \sum_m \Omega_m(t) |e\rangle\langle g| e^{i(\omega_q - m\chi_s)t} + h.c.$$

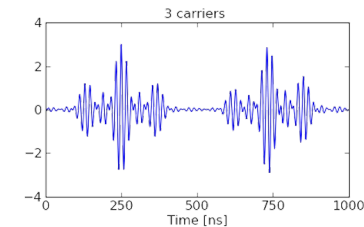
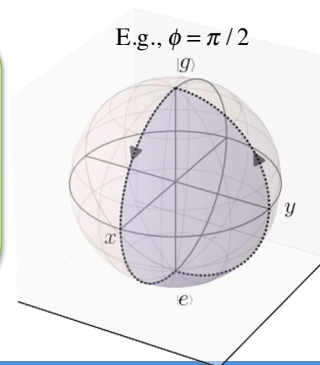
Qubit drive
(selective)



Geometric phase gate

Using two pi-pulses along different axes

$$|\psi'\rangle = \begin{cases} R_0^{(m)}(\pi) R_{\pi+\phi_m}^{(m)}(\pi) |n,g\rangle = e^{i\phi_m} |n,g\rangle & \text{for } n = m \\ |n,g\rangle & \text{for } n \neq m \end{cases}$$



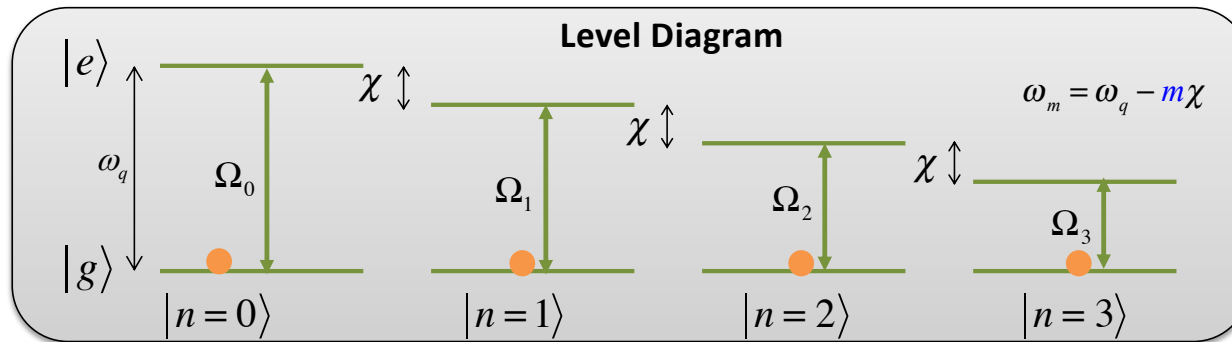
SNAP (Selective Number-dependent Arbitrary Phase) Gate

$$H = \omega_s a^\dagger a + \varepsilon(t) a^\dagger e^{i\omega_s t} + h.c.$$

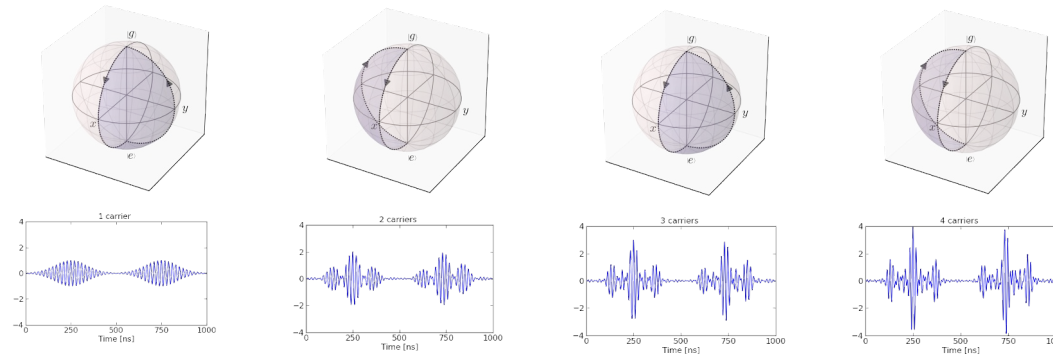
Cavity drive

$$+ (\omega_q - \chi_s a^\dagger a) |e\rangle\langle e| + \sum_m \Omega_m(t) |e\rangle\langle g| e^{i(\omega_q - m\chi_s)t} + h.c.$$

Qubit drive
(selective)



Simultaneous
Geometric
phase gates
for all $|n\rangle$



Theory: Krastanov, L.J., et al., PRA 92, 040303; Experiment: Heeres, L.J., Schoelkopf, et al. PRL 115, 137002 (2015)

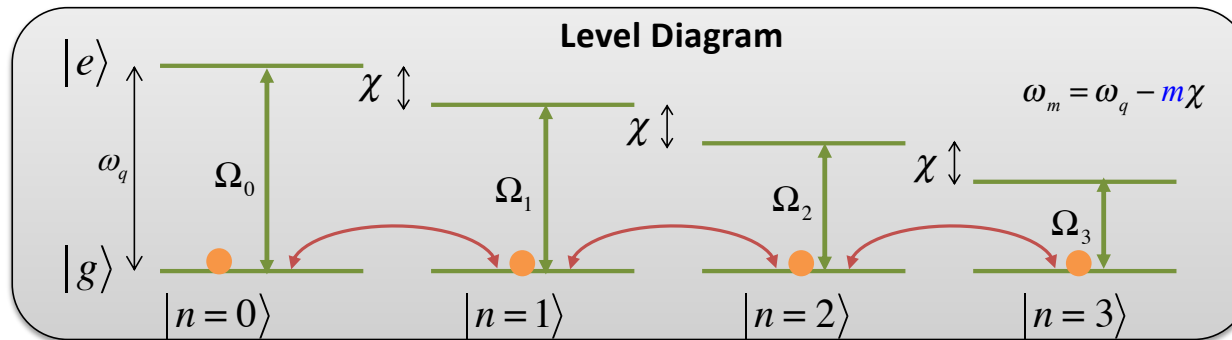
SNAP (Selective Number-dependent Arbitrary Phase) Gate

$$H = \omega_s a^\dagger a + \varepsilon(t) a^\dagger e^{i\omega_s t} + h.c.$$

Cavity drive

$$+ (\omega_q - \chi_s a^\dagger a) |e\rangle\langle e| + \sum_m \Omega_m(t) |e\rangle\langle g| e^{i(\omega_q - m\chi_s)t} + h.c.$$

Qubit drive
(selective)



Qubit Drive $|g\rangle \rightarrow |g\rangle$

Simultaneous Geometric phase gate

$$|\psi\rangle = \sum_n c_n |n, g\rangle \longrightarrow |\psi'\rangle = \sum_n e^{i\phi_n} c_n |n, g\rangle$$

SNAP Gate:

$$S_{\vec{\phi}} = \sum_n e^{i\phi_n} |n\rangle\langle n|$$

Cavity Drive $|g\rangle\langle g|$

Displace the cavity field with qubit in $|g\rangle$

$$|\psi\rangle = \sum_n c_n |n, g\rangle \longrightarrow |\psi'\rangle = D(\alpha) \sum_n c_n |n, g\rangle$$

Displacement:

$$D_\alpha = \exp[\alpha a^\dagger - \alpha^* a]$$

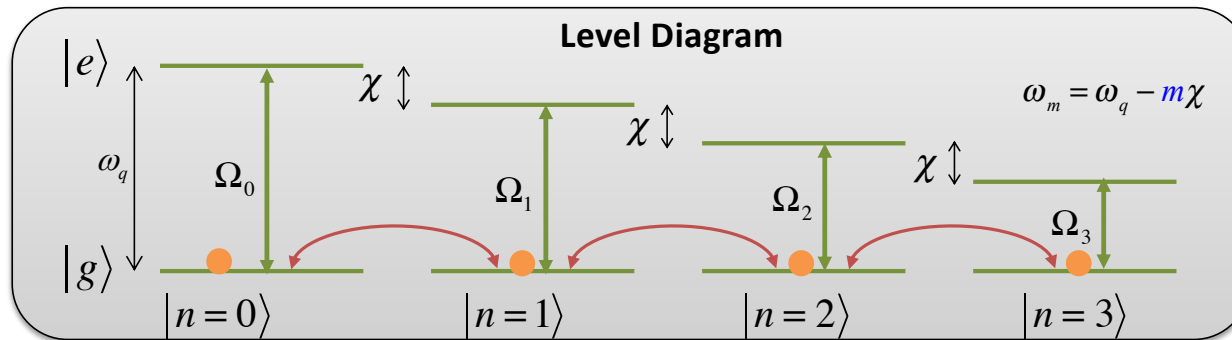
SNAP (Selective Number-dependent Arbitrary Phase) Gate

$$H = \omega_s a^\dagger a + \varepsilon(t) a^\dagger e^{i\omega_s t} + h.c.$$

Cavity drive

$$+ (\omega_q - \chi_s a^\dagger a) |e\rangle\langle e| + \sum_m \Omega_m(t) |e\rangle\langle g| e^{i(\omega_q - m\chi_s)t} + h.c.$$

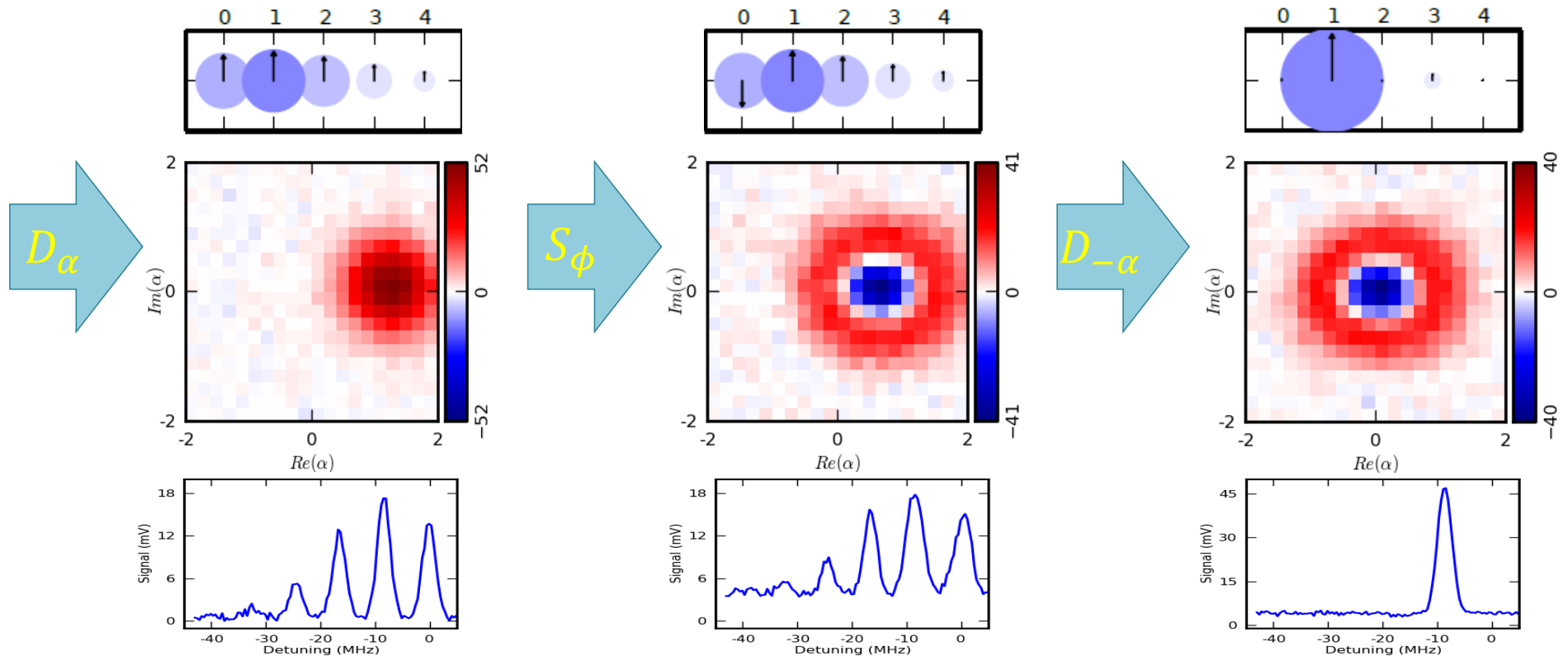
Qubit drive
(selective)



**Universal Control of
cavity-qubit system with
Cavity & Qubit Drives.**

Application 1: Arbitrary State Preparation

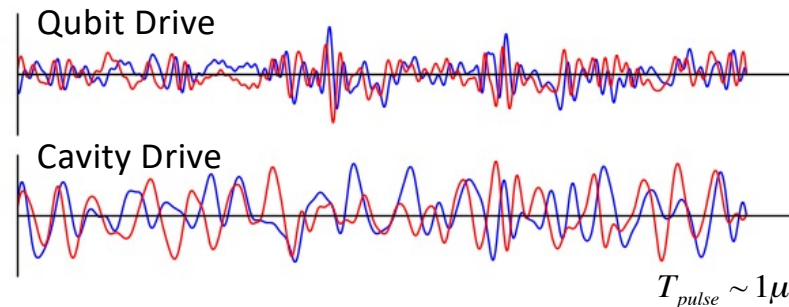
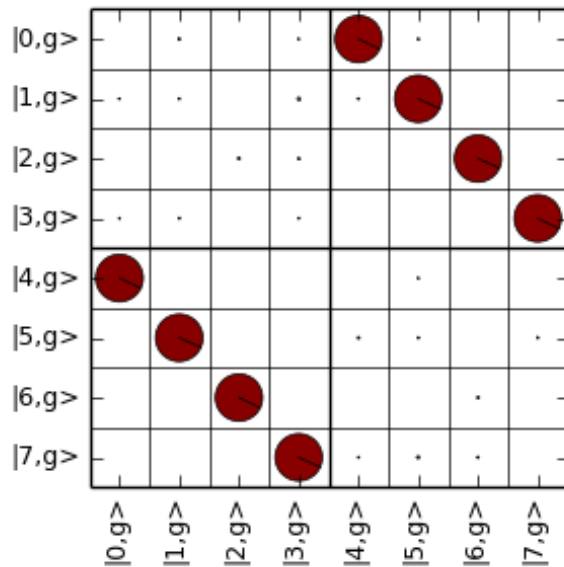
Three-pulse scheme: $D_{-\alpha} S_{\phi} D_{\alpha} |0\rangle \approx |1\rangle$ prepares single photon state deterministically.



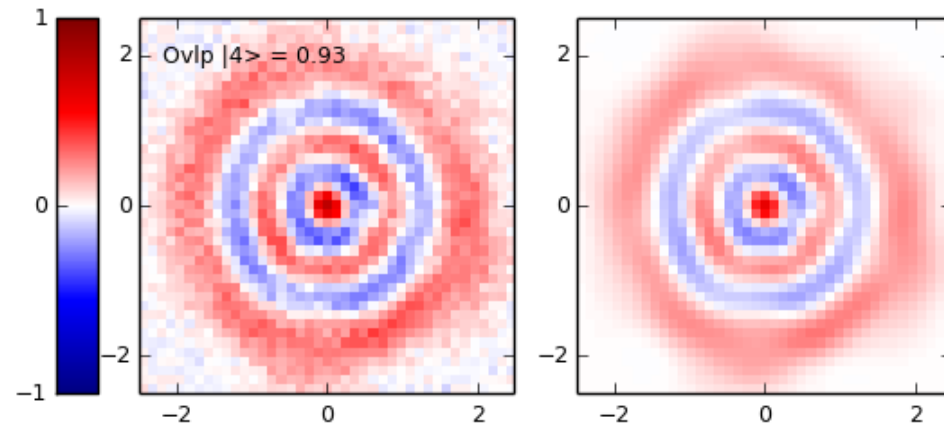
Application 2: Arbitrary Unitary Control of Bosonic Mode

GRAPE Optimized Pulses (simultaneous qubit & cavity drives)

Target Unitary



Tomography

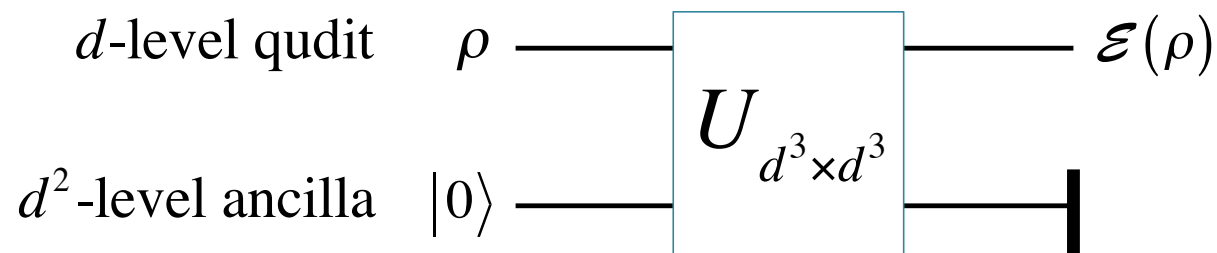


Application 3: Arbitrary Quantum Process

Most general CPTP (*Complete Positive Trace Preserving*) Map:

Kraus representation of CPTP channel

$$\rho \rightarrow \mathcal{E}(\rho) = \sum_{i=1}^N K_i \rho K_i^\dagger \quad \text{with Kraus rank } N \leq d^2 \quad \text{and} \quad \sum_{i=1}^N K_i^\dagger K_i = I.$$



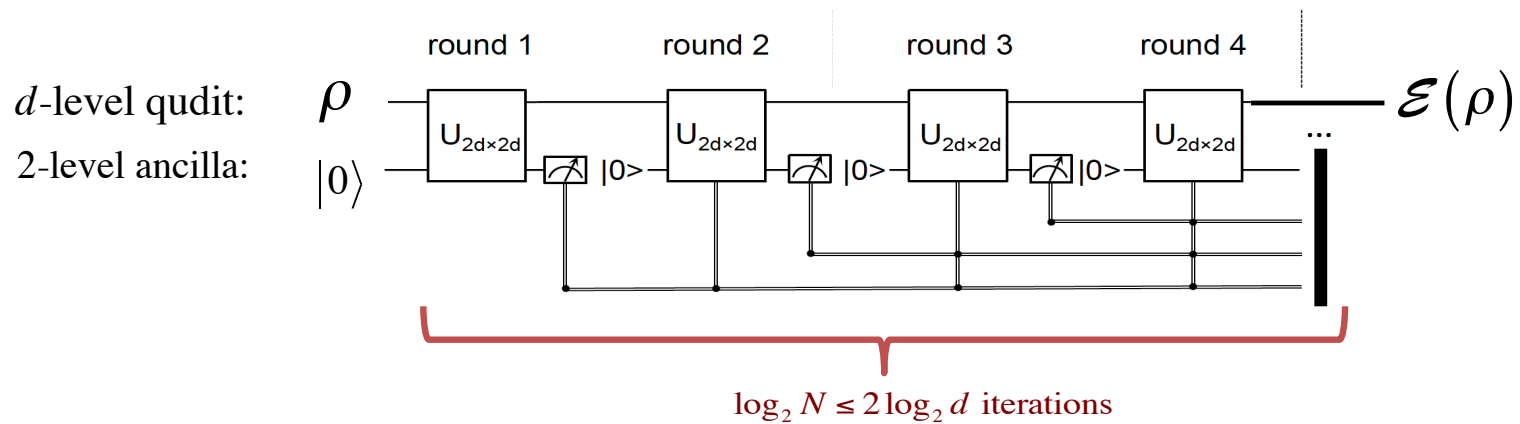
Challenging to use **d^2 -level ancilla** and implement $d^3 \times d^3$ large unitary operation.

Application 3: Arbitrary Quantum Channel of d-Level System

Most general CPTP map for d-level system:

$$\rho \rightarrow \mathcal{E}(\rho) = \sum_{i=1}^N K_i \rho K_i^\dagger \quad \text{with Kraus rank } N \leq d^2 \quad \text{and} \quad \sum_{i=1}^N K_i^\dagger K_i = I.$$

Key Idea: Repeatedly use the 2-level ancilla & adaptive control

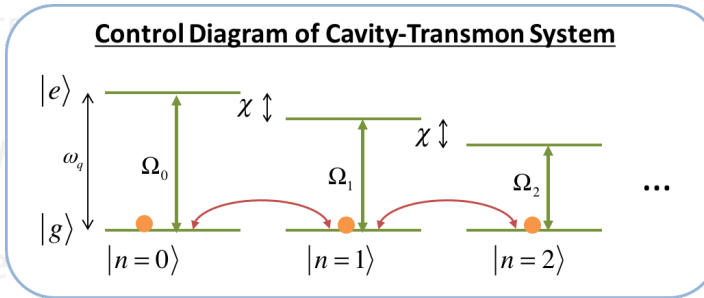


Application 3: Arbitrary Quantum Channel of d-Level System

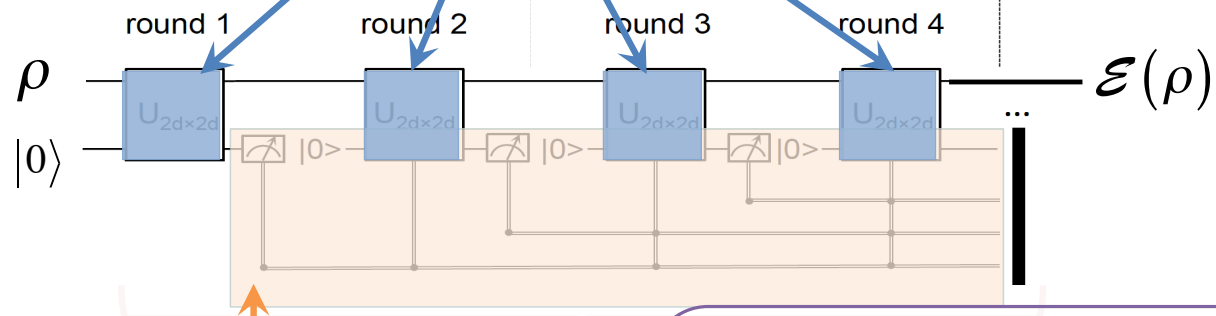
Most general CPTP

$$\rho \rightarrow \mathcal{E}(\rho) = \sum_{i=1}^N K_i \rho K_i^\dagger$$

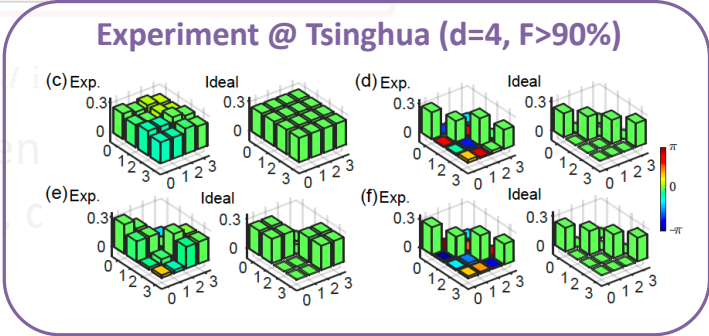
Key Idea: Repeated



d-level qudit:
2-level ancilla:



QND Measurement of Transmon & Adaptive Control with FPGA



Summary of Quantum Control over Bosonic Modes

State
Preparation

$$|0\rangle \rightarrow |\phi_{\text{target}}\rangle = U|0\rangle = \sum_{n=0}^{d-1} c_n |n\rangle$$

Unitary
Gates

$$\mathcal{U}(\rho) = U\rho U^\dagger$$

Quantum
Channel

$$\rho \rightarrow \mathcal{E}(\rho) = \sum_{i=1}^N K_i \rho K_i^\dagger$$

Quantum Error Correction of Bosonic Loss Errors

Key Challenges:

- How to achieve **universal control** of individual bosonic modes?
- How to **correct excitation loss errors** for quantum computing?
- How to implement quantum non-demolition (QND) parity measurement?

Multi-Qubit Systems vs Oscillator Systems

Qubits

- Hilbert space consists of tensor products of two-level systems.



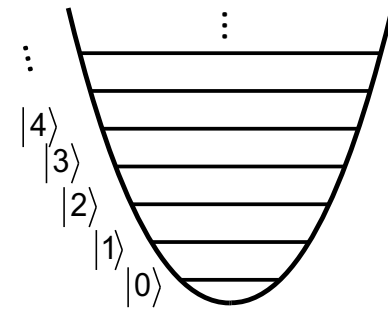
- Dominant errors expressible in terms of the Pauli operators

$$\{\sigma_X, \sigma_Y, \sigma_Z\}$$

- Distance of code: # of tolerable qubit loss
- Hardware inefficient (large overhead)

Oscillators

- Hilbert space is a ladder of Fock states $|n\rangle$ for $n \geq 0$.



- Dominant excitation loss error (and/or dephasing error)

$$\{a^\ell\} \text{ with } a|n\rangle = \sqrt{n}|n-1\rangle$$

- Distance: # of tolerable excitation loss
- Hardware efficient (e.g., single oscillator)

Zoo of Bosonic Codes

Fock states $|0\rangle, |1\rangle$ only

- **Dual-rail:** Chuang Yamamoto 1995
- **Linear optics:** Knill Laflamme Milburn 2001
- **Parity check:** Ralph Hayes Gilchrist 2005
- Any codes correcting against qubit amplitude damping: [37,64-73] in arxiv:1708.05010v2



Higher Fock state superpositions

- Chuang Leung Yamamoto 1997
- Wasilewski Banaszek 2007
- **Noon:** Boto Kok Abrams Braunstein Williams Dowling 2000
- **Bin:** Michael Silveri Brierley Albert Salmilehto L.J. Girvin 2016
- χ^2 : Niu Chuang Shapiro 2017



Position/momentum eigenstates

- Lloyd Slotine 1998
- Braunstein 1998



Coherent

- **Can we use single bosonic mode to correct loss errors?**

- **Cat (no loss protection):** Niset Andersen Cerf 2008
- **Cat:** Leghtas Kirchmair Vlastakis Schoelkopf Devoret Mirrahimi 2013
- **Coherent state constellation (scheme):** Lacerda Renes Scholz 2016
- **Multimode cat:** Albert Mundhada Grimm Touzard Devoret L.J. 2018

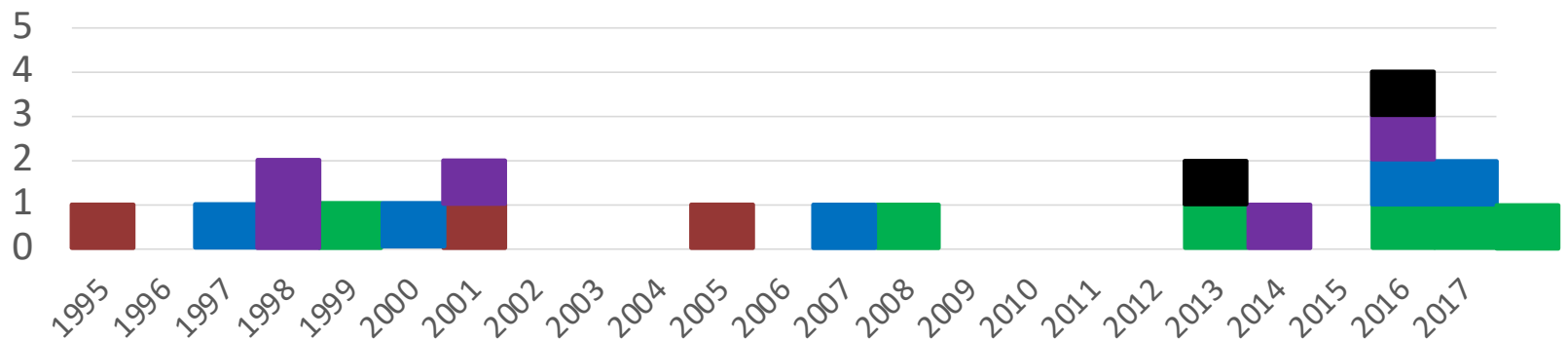
Hybrid – multi-qubit codes

- **Hybrid optical qubits:** Lee Jeong 2013
- **Very small logical qubit:** Kapit 2016



Bosonic Codes by year

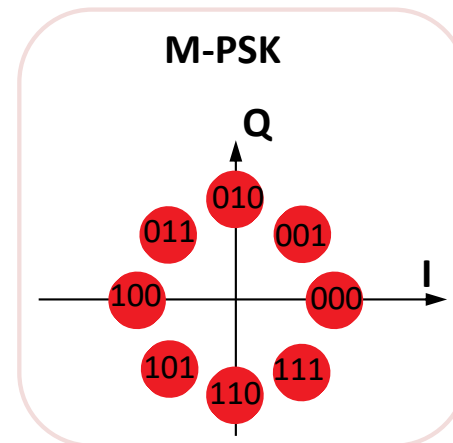
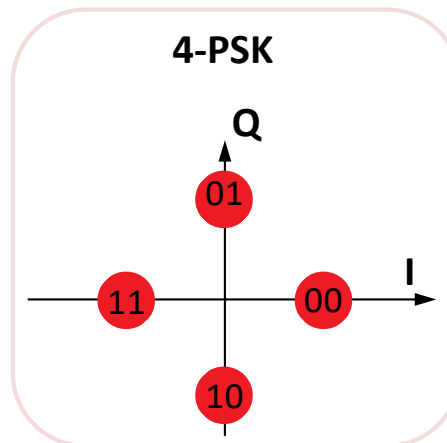
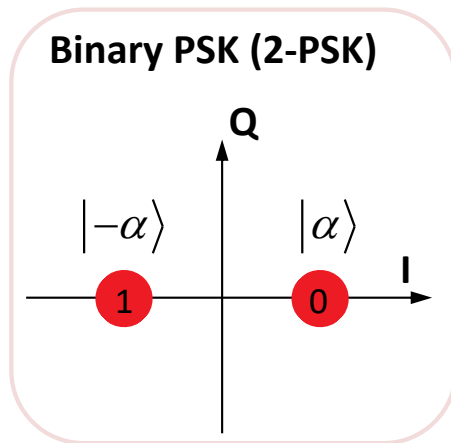
Prepared by:
Victor Albert
(Zookeeper)



See more at Error Correction Zoom: <https://errorcorrectionzoo.org/>

Classical Encoding – Phase-Shift Keying (PSK)

- **Idea:** Encode information in the phase of a coherent state, $\alpha = |\alpha|e^{i\phi}$



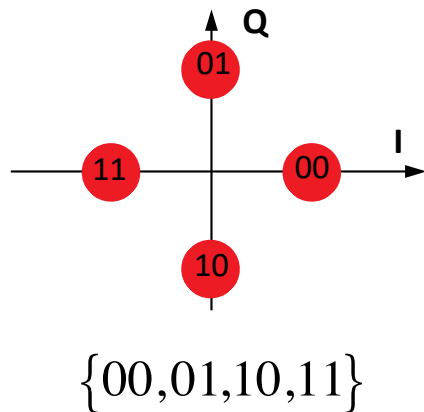
- Encode more information by increasing the number of phase angles
- Widely used in Wifi, Bluetooth, etc.



(Classical) PSK v.s. (Quantum) Cat Codes

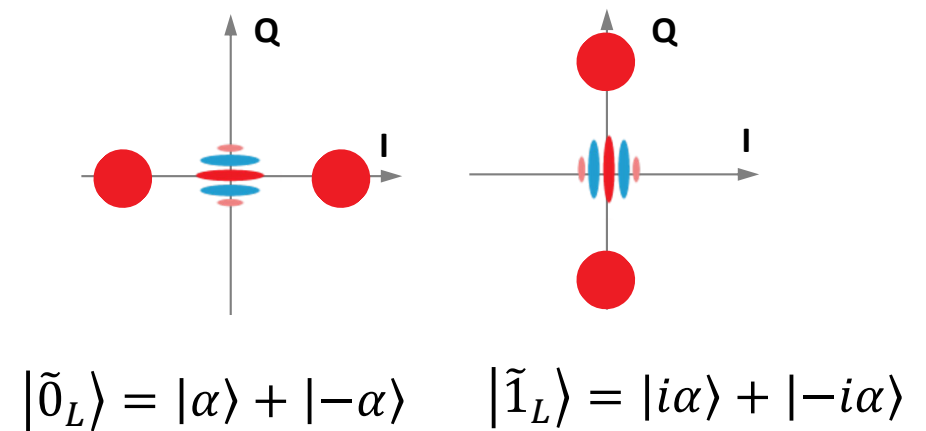
Classical Phase-Shift Keying (PSK):

- Convey data by modulating phase of a pulse signal



(Quantum) Cat Codes:

- **Superposition** of coherent states
- Fixed (even) **parity** logical subspace



Cat Codes

Logical States: $|\tilde{0}_L\rangle = |\alpha\rangle + |-\alpha\rangle$ $|\tilde{1}_L\rangle = |i\alpha\rangle + |-i\alpha\rangle$
 within **even** number subspace

Circle
of
Life

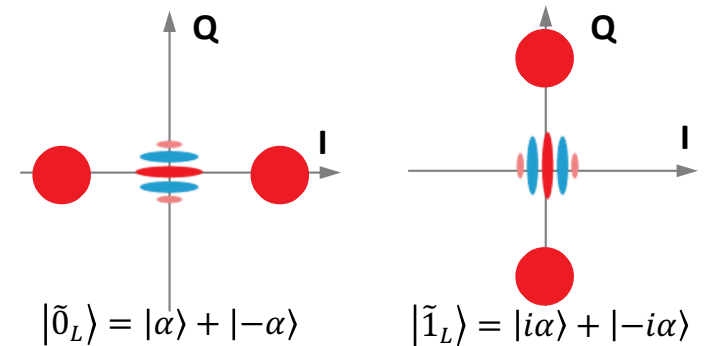
No excitation loss: $|\Psi^{[0]}\rangle = c_g(|\alpha\rangle + |-\alpha\rangle) + c_e(|i\alpha\rangle + |-i\alpha\rangle)$ **even**

1 excitation loss: $|\Psi^{[1]}\rangle = c_g(|\alpha\rangle - |-\alpha\rangle) + ic_e(|i\alpha\rangle - |-i\alpha\rangle)$ **odd**

2 excitation loss: $|\Psi^{[2]}\rangle = c_g(|\alpha\rangle + |-\alpha\rangle) - c_e(|i\alpha\rangle + |-i\alpha\rangle)$ **even**

3 excitation loss: $|\Psi^{[3]}\rangle = c_g(|\alpha\rangle - |-\alpha\rangle) - ic_e(|i\alpha\rangle - |-i\alpha\rangle)$ **odd**

4 excitation loss: $|\Psi^{[0]}\rangle = c_g(|\alpha\rangle + |-\alpha\rangle) + c_e(|i\alpha\rangle + |-i\alpha\rangle)$ **even**



Properties:

1. **Detect single excitation loss by measuring parity**
2. Only need to track single photon losses (mod 4)

Correct single loss error!

Other choices of codes - Kitten/Binomial Codes

Kitten Code:

$$|0_L\rangle = (|0\rangle + |4\rangle)/\sqrt{2}$$

$$|1_L\rangle = |2\rangle$$

$$E = \{I, a\}$$

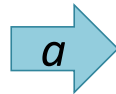
Binomial Code:

$$|0_L/1_L\rangle = \frac{1}{\sqrt{2^N}} \sum_{p \text{ even/odd}} \sqrt{\binom{N+1}{p}} |p \cdot (S+1)\rangle$$

$$E = \{I, a, a^2, \dots, a^S\}$$

$$|0_L\rangle = (|0\rangle + |4\rangle)/\sqrt{2}$$

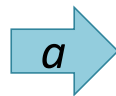
$$|1_L\rangle = |2\rangle$$



$$a|0_L\rangle = \sqrt{2}|3\rangle$$

$$a|1_L\rangle = \sqrt{2}|1\rangle$$

$$c_0 \frac{(|0\rangle + |4\rangle)}{\sqrt{2}} + c_1 |2\rangle$$



$$\sqrt{2}(c_0|3\rangle + c_1|1\rangle)$$

Quantum Error Correction using Cat Code

Key Challenges:

- How to **correct excitation loss errors** for quantum computing?
- How to achieve **universal control** of individual bosonic modes?
- How to implement **quantum non-demolition (QND)** parity measurement?

Recall: Cat Codes

Logical States: $|\tilde{0}_L\rangle = |\alpha\rangle + |-\alpha\rangle$ $|\tilde{1}_L\rangle = |i\alpha\rangle + |-i\alpha\rangle$
 within **even** number subspace

Circle
of
Life

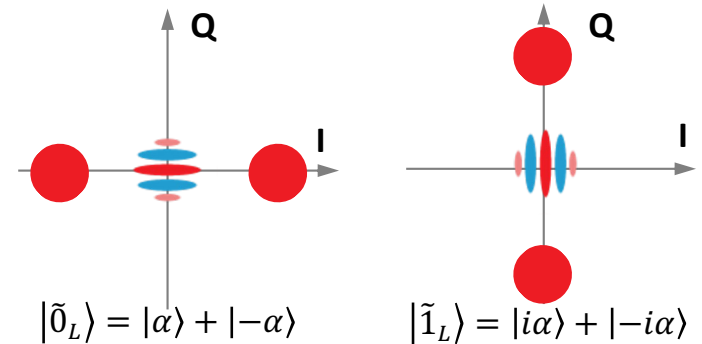
No excitation loss: $|\Psi^{[0]}\rangle = c_g(|\alpha\rangle + |-\alpha\rangle) + c_e(|i\alpha\rangle + |-i\alpha\rangle)$ **even**

1 excitation loss: $|\Psi^{[1]}\rangle = c_g(|\alpha\rangle - |-\alpha\rangle) + ic_e(|i\alpha\rangle - |-i\alpha\rangle)$ **odd**

2 excitation loss: $|\Psi^{[2]}\rangle = c_g(|\alpha\rangle + |-\alpha\rangle) - c_e(|i\alpha\rangle + |-i\alpha\rangle)$ **even**

3 excitation loss: $|\Psi^{[3]}\rangle = c_g(|\alpha\rangle - |-\alpha\rangle) - ic_e(|i\alpha\rangle - |-i\alpha\rangle)$ **odd**

4 excitation loss: $|\Psi^{[0]}\rangle = c_g(|\alpha\rangle + |-\alpha\rangle) + c_e(|i\alpha\rangle + |-i\alpha\rangle)$ **even**



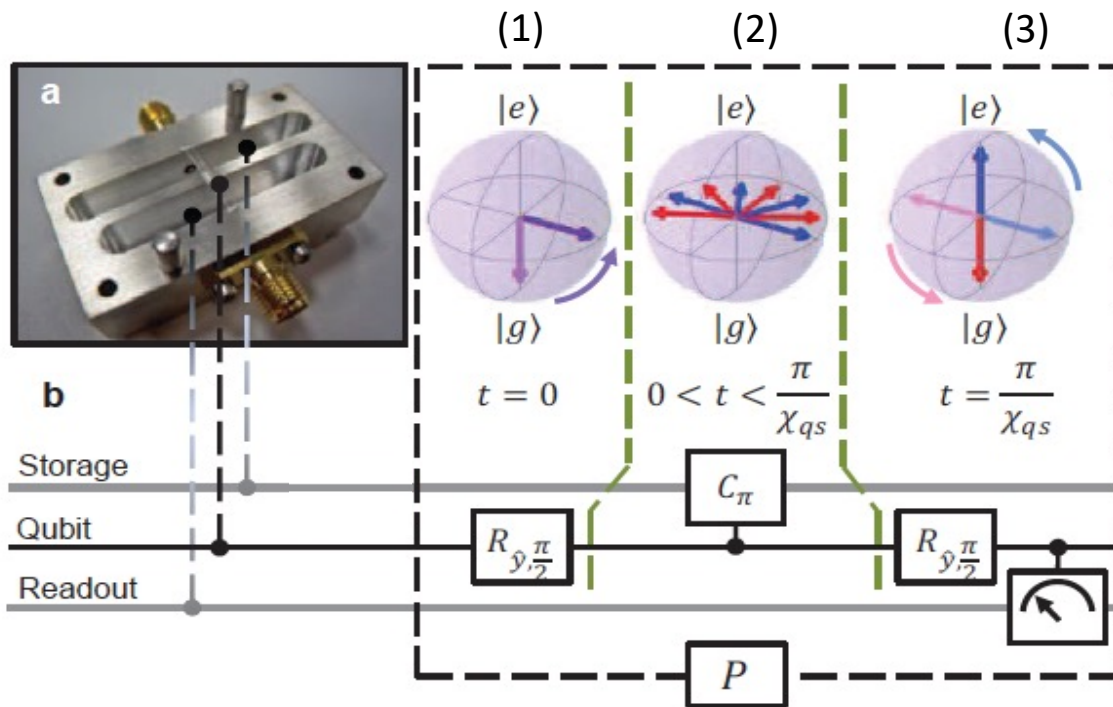
Properties:

1. **Detect single excitation loss by measuring parity**
2. Only need to track single photon losses (mod 4)

Can correct single loss error!

Question: How to perform QND measurement of parity?

QND Measurement of Photon Number Parity



Procedure:

(1). Transmon initialization: $|g\rangle + |e\rangle$

(2). Phase accumulation:

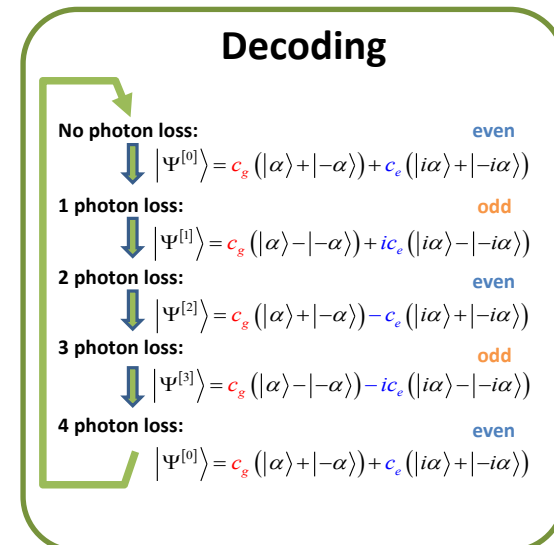
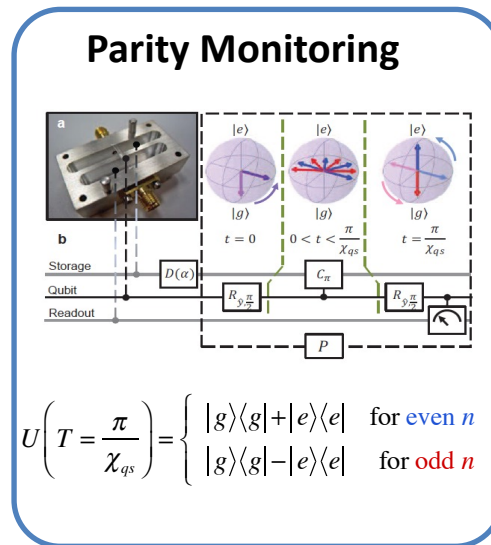
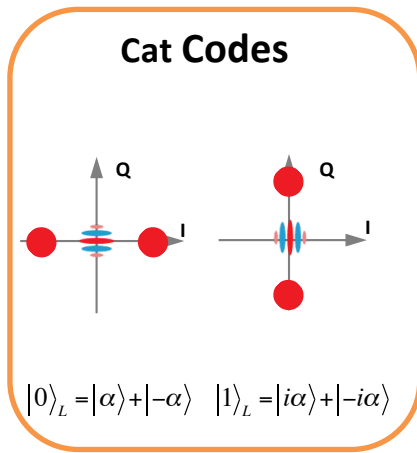
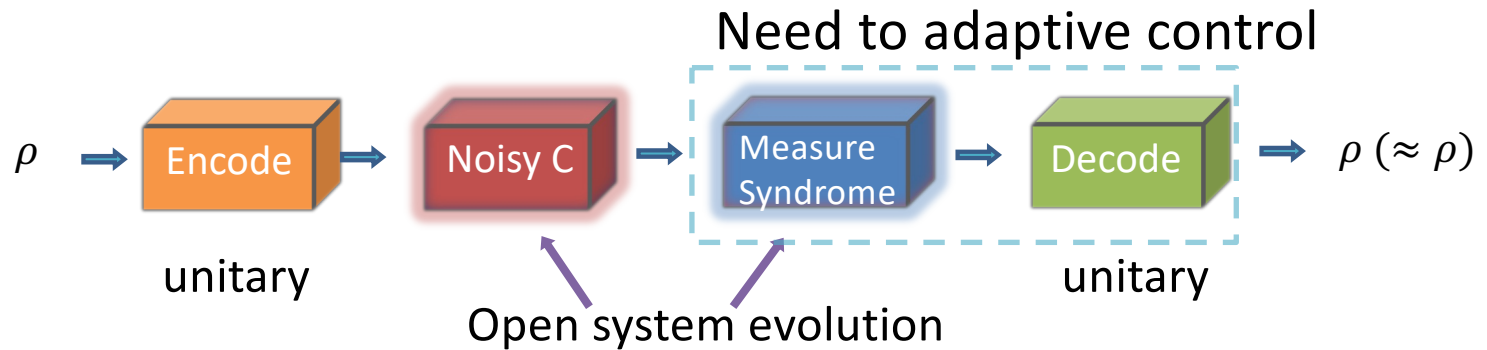
$$H = -\chi a^\dagger a |e\rangle\langle e|$$

$$U\left(T = \frac{\pi}{\chi}\right) = \begin{cases} |g\rangle\langle g| + |e\rangle\langle e| & \text{for even } n \\ |g\rangle\langle g| - |e\rangle\langle e| & \text{for odd } n \end{cases}$$

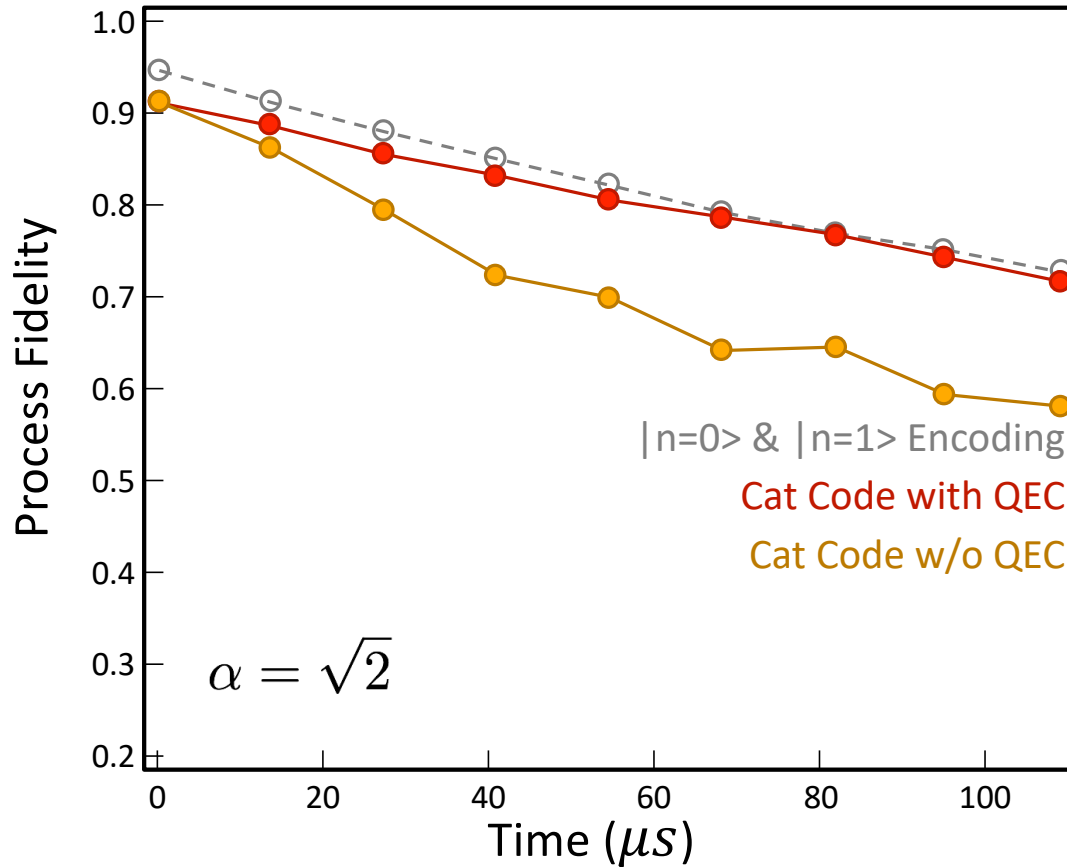
(3). Transmon measurement in

$|g\rangle \pm |e\rangle$ basis

Quantum Error Correction (QEC)



QEC Experiment – Reaching Break Even



$\tau \approx 290 \mu s$

$\tau \approx 320 \mu s$

$\tau \approx 130 \mu s$

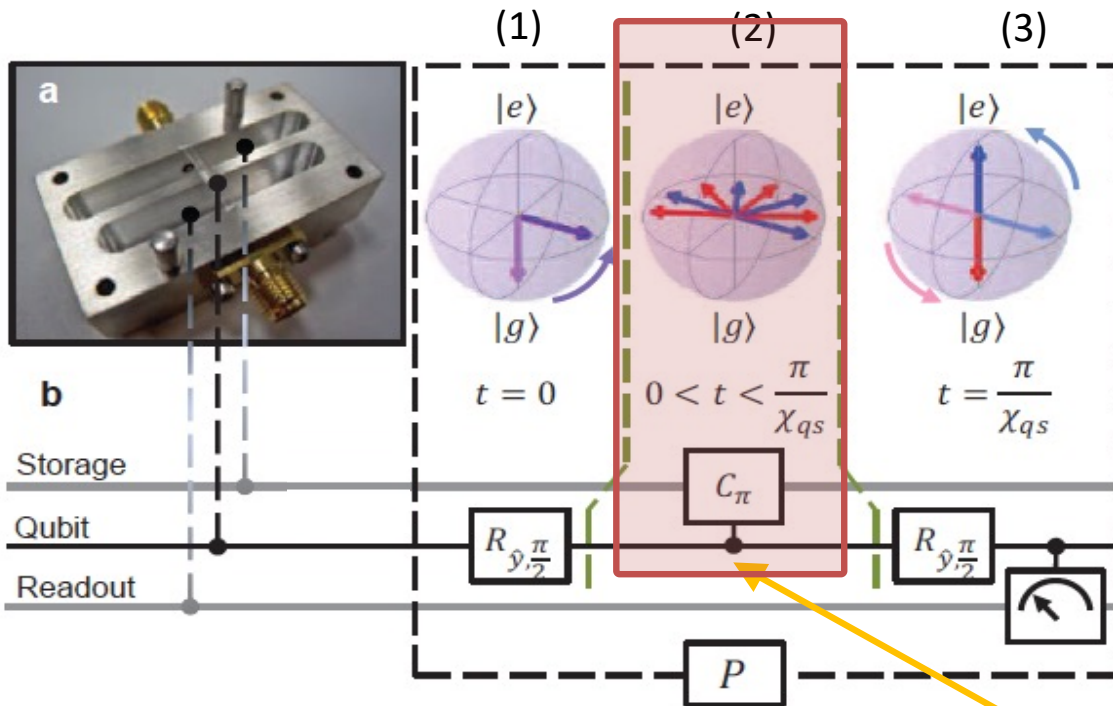
Limited by Transmon decay (prob $\sim 10^{-2}$) during measurement, causing cavity dephasing

Related Experiments:

Binomial Code: Hu, et al., Nat Phys 15, 503 (2019)

GKP Code: Campagne-Ibarcq, et al., Nature 584, 368 (2020)

QND Measurement of Photon Number Parity



Procedure:

(1). Transmon initialization: $|g\rangle + |e\rangle$

(2). Phase accumulation:

$$H = -\chi a^\dagger a |e\rangle\langle e|$$

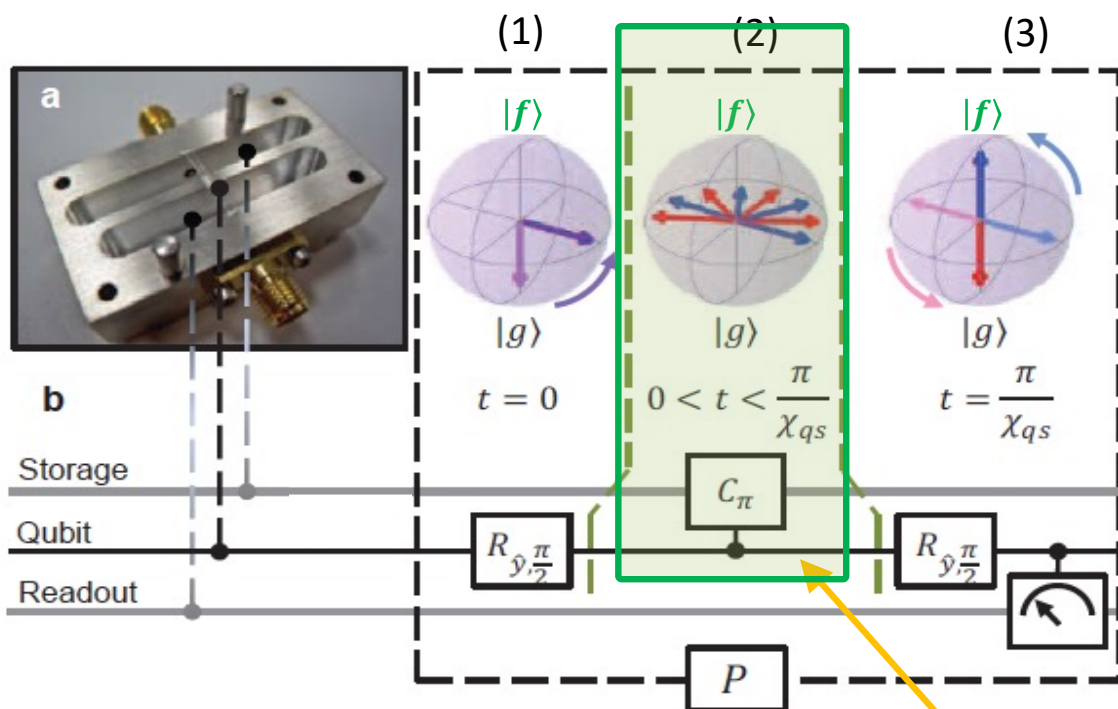
$$U\left(T = \frac{\pi}{\chi}\right) = \begin{cases} |g\rangle\langle g| + |e\rangle\langle e| & \text{for even } n \\ |g\rangle\langle g| - |e\rangle\langle e| & \text{for odd } n \end{cases}$$

(3). Transmon measurement in $|g\rangle \pm |e\rangle$ basis

Not Fault Tolerant!
Vulnerable to Transmon Decay ($|e\rangle \rightarrow |g\rangle$)

How to Go beyond Break Even?

-- Use **Fault-tolerant** Parity Measurement



Procedure:

(Transmon decay: $\hat{E} = |e\rangle\langle f|$)

1. Transmon initialization: $|g\rangle + |f\rangle$
2. Phase accumulation:

$$H = -\chi a^\dagger a (|e\rangle\langle e| + |f\rangle\langle f|)$$

$$U\left(T = \frac{\pi}{\chi}\right) = \begin{cases} |g\rangle\langle g| + |f\rangle\langle f| & \text{for even } n \\ |g\rangle\langle g| - |f\rangle\langle f| & \text{for odd } n \end{cases}$$

3. Transmon measurement in basis:

- $|g\rangle \pm |f\rangle$ for even/odd parity
- $|e\rangle$ for transmon T1 decay, with unitary $U = e^{i\pi a^\dagger a}$ of the cavity.

Fault Tolerant!

Robust against Transmon Decay ($|f\rangle \rightarrow |e\rangle$)

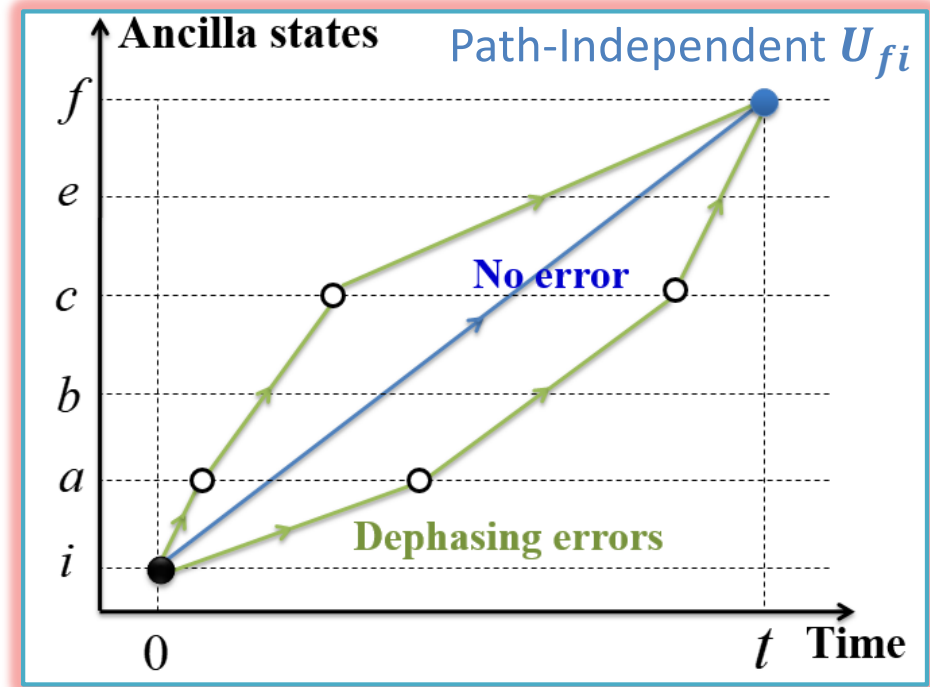
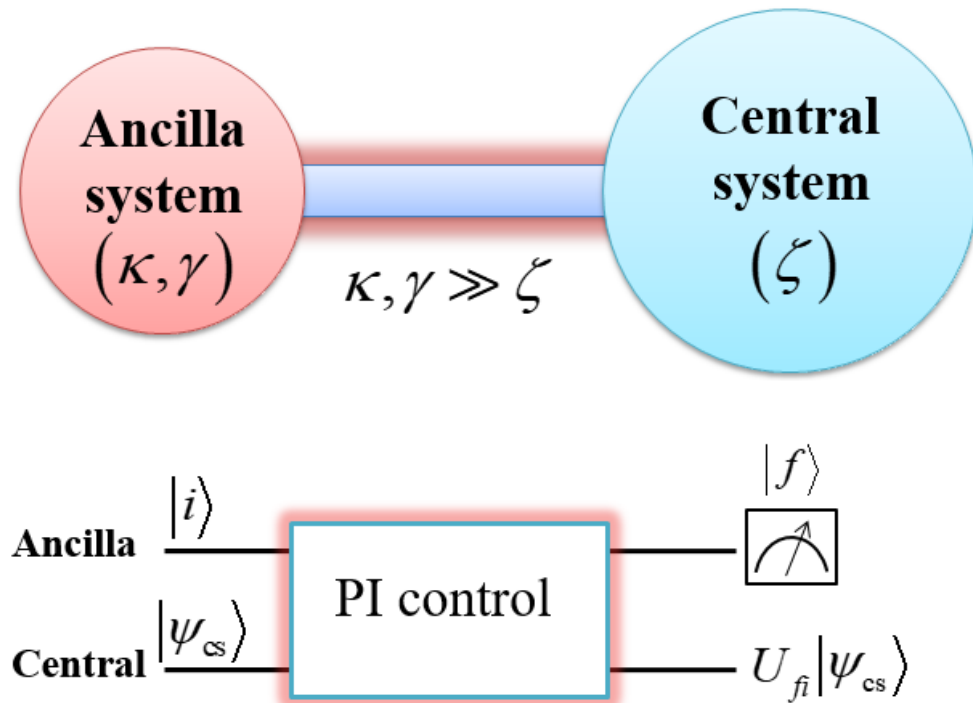
Key Idea: **Error Transparency (Path-Independency):**

$$H \hat{E} |\phi_{enc}\rangle = \hat{E} H |\phi_{enc}\rangle$$

for $\hat{E} = |e\rangle\langle f|$ & $|\phi_{enc}\rangle \in \text{span}\{|g\rangle, |f\rangle\}$

Use Noisy Ancilla for High Fidelity Operations

-- Path-Independent (PI) Quantum Gates

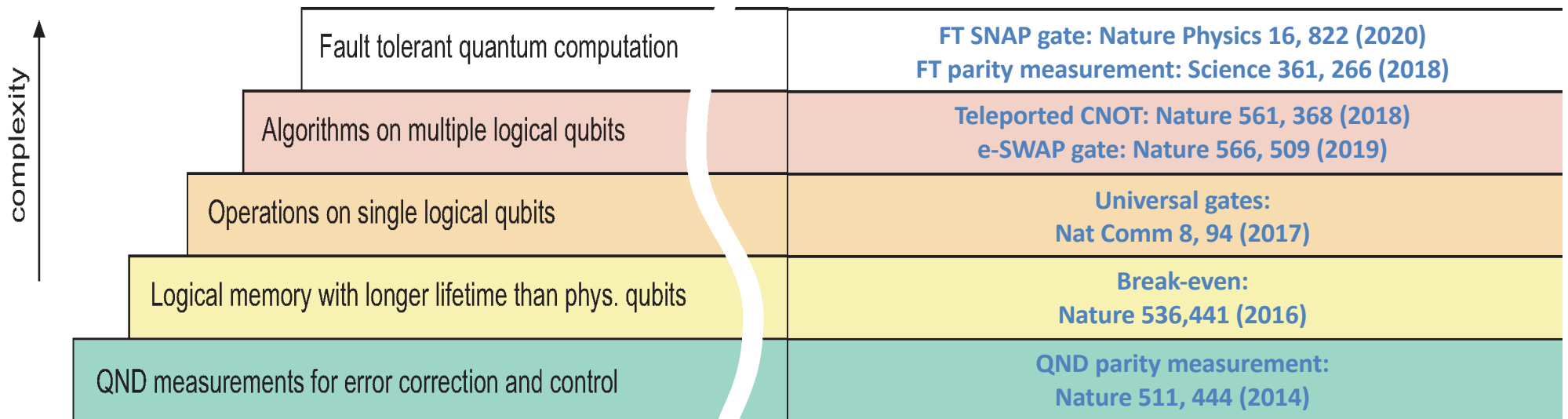


Note: {Error Transparent Gates} \subset {Path Independent Gates}

Ordinary matrix algebra & PI matrix algebra (Isomorphism)

	Ordinary matrix algebra \mathcal{A} (for only the ancilla system):	PI matrix algebra $\mathcal{A}_{PI} \cong \mathcal{A}$ (for the ancilla and central systems):
Basis	$\mathcal{B}_A = \{ m\rangle\langle n \}_{m,n=0}^{d-1}$	$\mathcal{B} = \{ m\rangle\langle n \otimes U_{mn}\}_{m,n=0}^{d-1}$
Jump operator	$(a\rangle\langle b) \cdot (c\rangle\langle d) = \delta_{bc} a\rangle\langle d $	$(a\rangle\langle b \otimes U_{ab}) \cdot (c\rangle\langle d \otimes U_{cd}) = \delta_{bc} a\rangle\langle d \otimes U_{ad}$
Hamiltonian	$H_A(t) = \sum_{m,n} \varepsilon_{mn}(t) m\rangle\langle n $	$H(t) = \sum_{m,n} \varepsilon_{mn}(t) m\rangle\langle n \otimes U_{mn}$
Propagator	$W_A(t_2, t_1) = \sum_{m,n} \xi_{mn}(t_2, t_1) m\rangle\langle n $	$W(t_2, t_1) = \sum_{m,n} \xi_{mn}(t_2, t_1) m\rangle\langle n \otimes U_{mn}$
Path independence	$ r\rangle\langle a a\rangle\langle b b\rangle\langle c \cdots e\rangle\langle i = r\rangle\langle i $	$(r\rangle\langle a \otimes U_{ra}) \cdot (a\rangle\langle b \otimes U_{ab}) \cdot (b\rangle\langle c \otimes U_{bc}) \cdots (e\rangle\langle i \otimes U_{ei}) = r\rangle\langle i \otimes U_{ri}$

Quantum Computation with Bosonic QEC



Devoret & Schoelkopf, Science 339, 1169 (2013)

Key Challenges:

- How to achieve control of bosonic modes?
- How to implement QND parity msmt & fault-tolerance?
- *How to couple remote bosonic modes?*
- *How to achieve code-independent coupling gate?*

Ideas:

- ✓ Transmon SNAP gate & optimal control
- ✓ Dispersive Ramsey msmt & Error transparent design
- ✓ *Teleported CNOT gate*
- ✓ *E-SWAP gate*

Various Approaches for Bosonic QEC

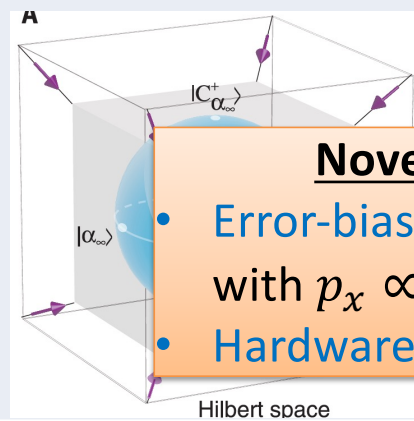
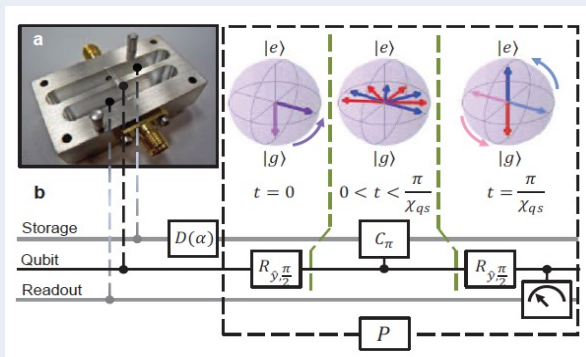
Discrete QEC

Continuous QEC

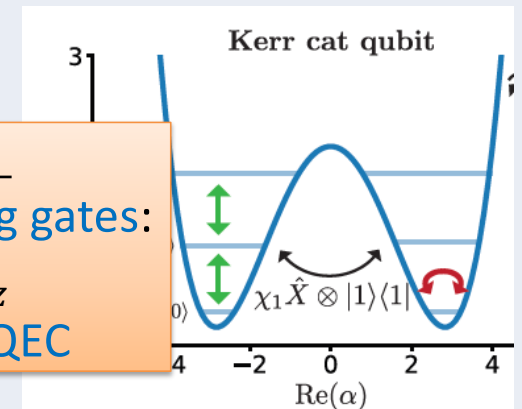
Syndrome Measurement

Engineered Dissipation

Hamiltonian Protection



- Novel Feature:**
- Error-bias preserving gates: with $p_x \propto e^{-n} \ll p_z$
 - Hardware-Efficient QEC



E.g., Cat, Binomial, GKP codes

E.g., Dissipative-cat qubit

E.g., Kerr-cat qubit

Experiments on cat/binomial/GKP codes

- Ofek, et al., Nature 536, 441 (2016)
- Hu, et al., Nat Phys 15, 503 (2019)
- Campagne-Ibarcq, et al., Nature 584, 368 (2020)

Experiments on Dissipative-cat qubit:

- Leghtas, et al., Science 347, 853 (2015)
- Lescanne, et al., Nature Physics 16, 509 (2020)
- Leghtas, MM 22', A28.00004

Experiment on Kerr-cat qubit:

- Grimm, et al., Nature 584, 205 (2020)

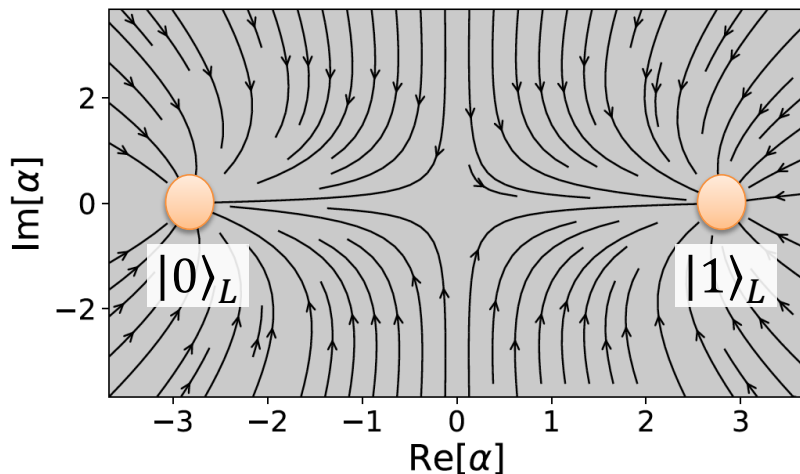
Real-time adaptive control

Engineer dissipation to remove entropy

Energy gap to suppress noise, no cooling

Outlook: *Hardware Efficient* Bias-Preserving QEC

Exponentially suppressed bit-flip errors



Key feature:

- Apply continuous QEC to exponentially suppress bit-flip errors:

$$p_x, p_y \propto e^{-2\alpha^2} \ll p_z$$

- Design universal gates preserving error bias (no mixing p_z into $p_{x,y}$)
- Use repetition-like code to achieve fault-tolerance

Hardware-Efficient Proposals for Fault-Tolerant QC with Bias-Preserving Gates

- Dissipative-cat qubit: Guillaud, Mirrahimi, PRX 9, 0411053 (2019); Chamberland, L.J., et al., PRXQ 3, 010329 (2022)
- Kerr-cat qubit: Puri, L.J., Girvin, et al., Science Advances 6, eaay5901 (2020); Qian, L.J. et al., arXiv:2105.13908
- Colored Kerr-cat qubit: Putterman, Iverson, Xu, L.J., Painter, Brandao, Noh, PRL (accepted), arXiv:2107.09198.
- Code optimization for biased noise: Xu, Mannucci, Seif, Kubica, Flammia, L.J., arXiv:2203.16486 (2022).