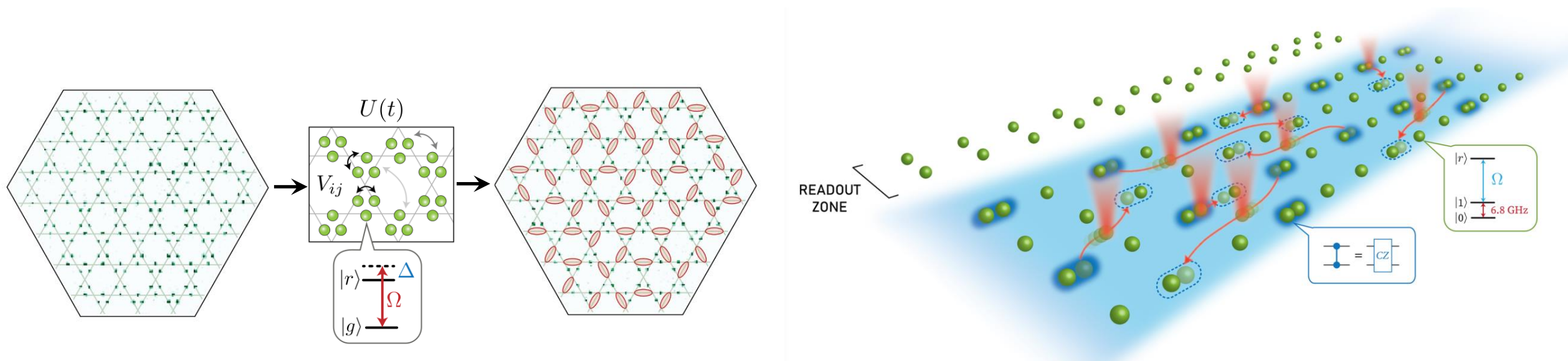
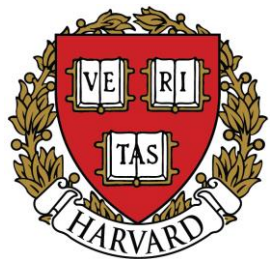


New frontiers in quantum simulation and computation with neutral atom arrays



Giulia Semeghini

Harvard University - Lukin Group



SPICE Workshop – Non-Equilibrium Emergence in Quantum Design
June 23rd, 2022

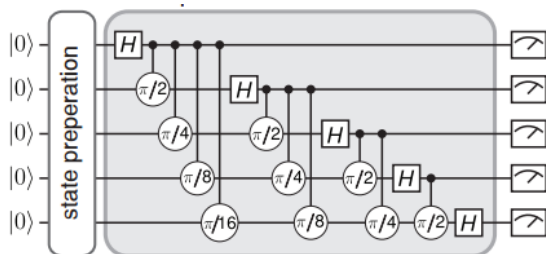


New frontiers in quantum science

Two key challenges:

- 1) building large scale **quantum machines**

Controllability:



Scalability:

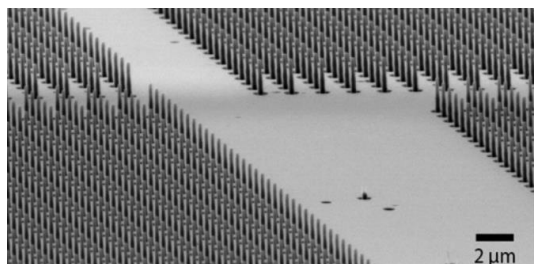
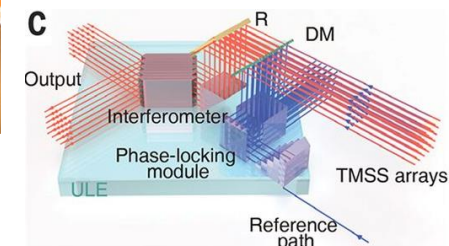
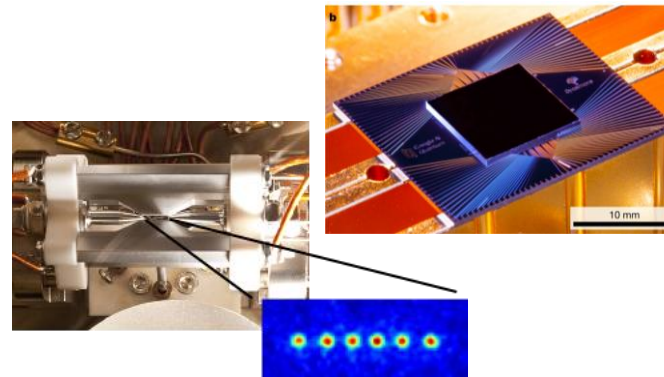
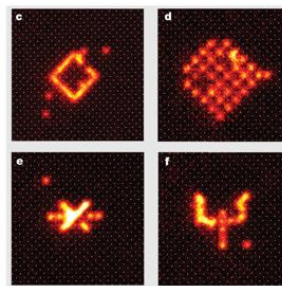


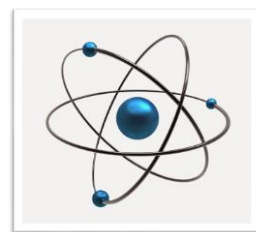
Image from Loncar's group

- 2) developing **new algorithms** and science **applications**

Several platforms: ions, atoms, superconductors, photons, defects...



Images from: Blatt, Bloch, Google, Jian-Wei Pan groups

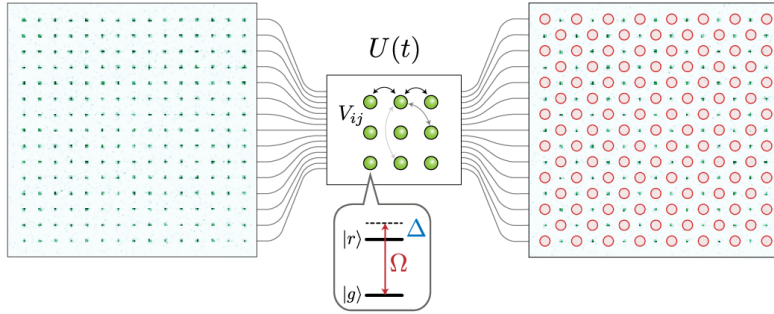


Our approach:
individual neutral atoms in optical tweezers



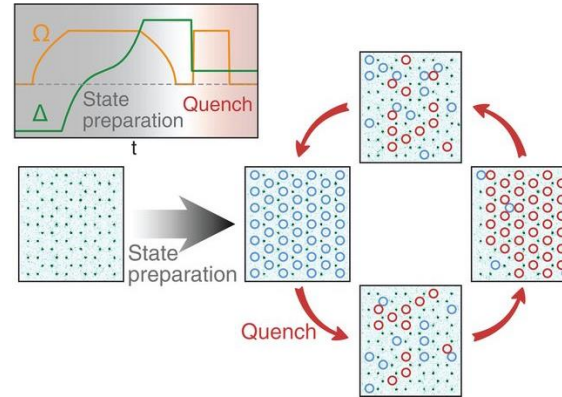
- Excellent isolation from the environment
- Highly scalable defect-free arrays (100s of **identical** atoms)
- Well-developed toolbox:
 - Initialization, readout
 - Strong, switchable Rydberg interactions
- **Scalability** \rightarrow **large-scale parallel photonic control**

Exploring quantum phases with 256 qubits



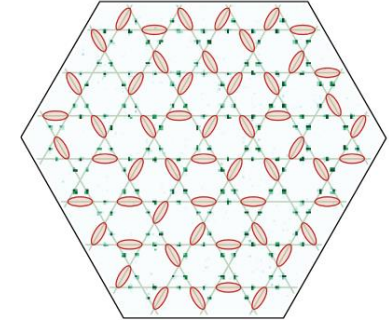
S. Ebadi et al, Nature (2021)

Non-equilibrium dynamics: many-body scars



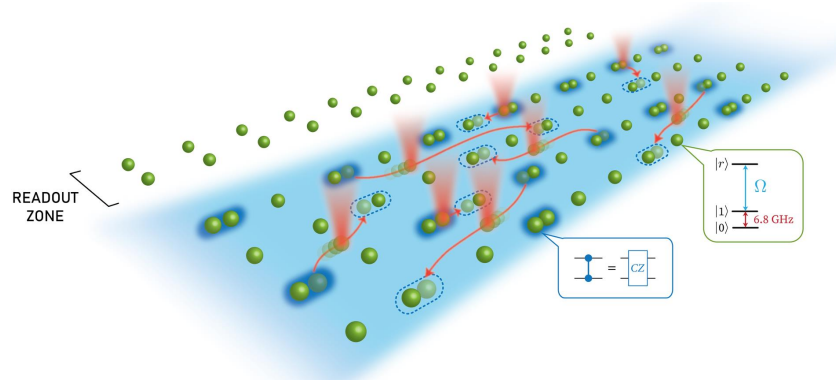
D. Bluvstein et al, Science (2021)

Topological spin liquids



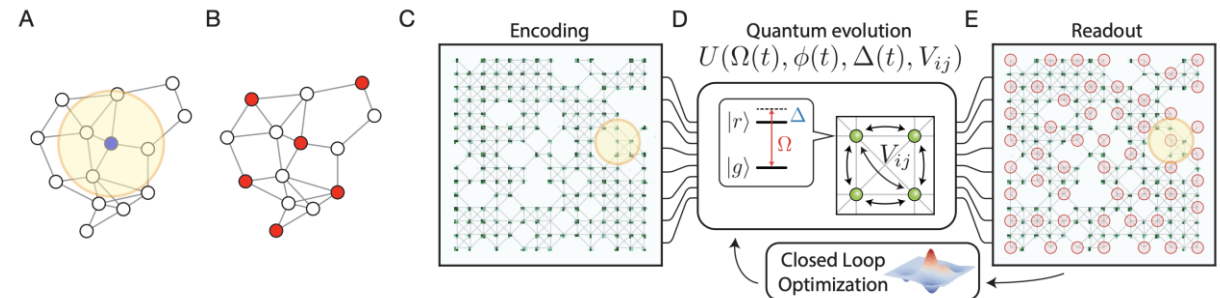
GS et al, Science 374, 1242 (2021)

Entanglement transport & universal programmability



D. Bluvstein et al, Nature (2022)

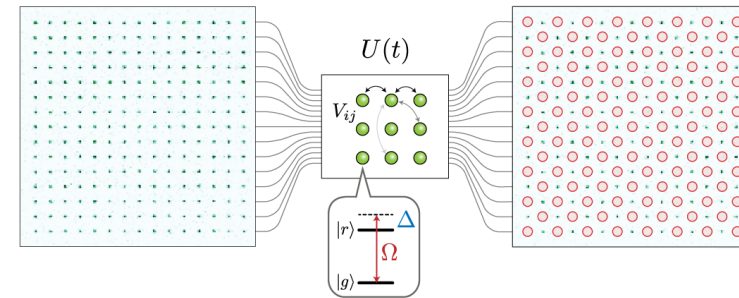
Accelerating combinatorial optimization



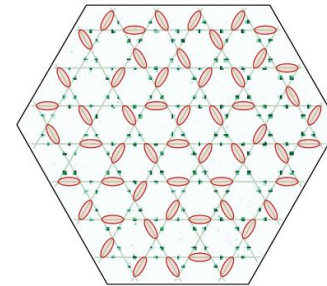
S. Ebadi et al, arXiv:2202.09372, Science (2022)

Outline

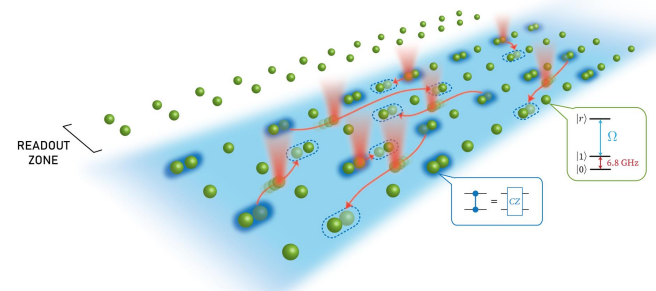
Programmable arrays
of Rydberg atoms



Experimental evidence of a
quantum spin liquid phase

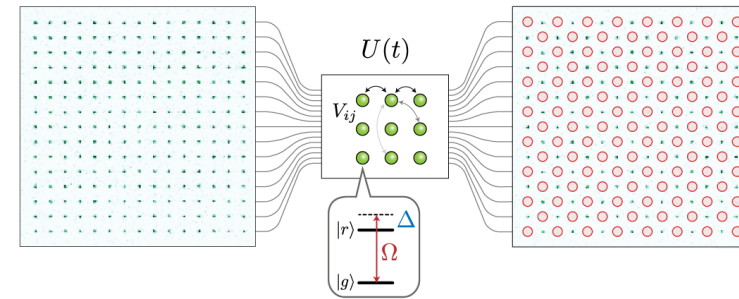


Dynamically reconfigurable
architecture for quantum
information processing

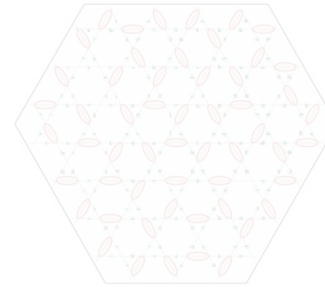


Outline

Programmable arrays
of Rydberg atoms



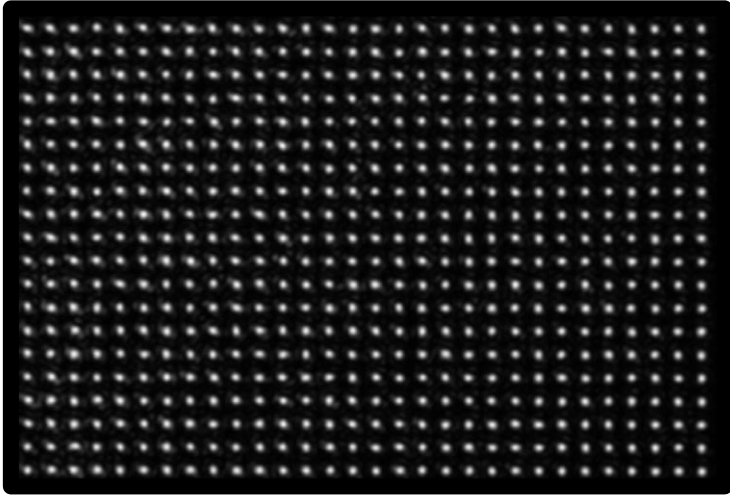
Experimental evidence of a
quantum spin liquid phase



Dynamically reconfigurable
architecture for quantum
information processing



2D array of optical tweezers



D. Kim, et al, Opt. Lett. 44, 3178 (2019)

Ebadi et al Nature 2021

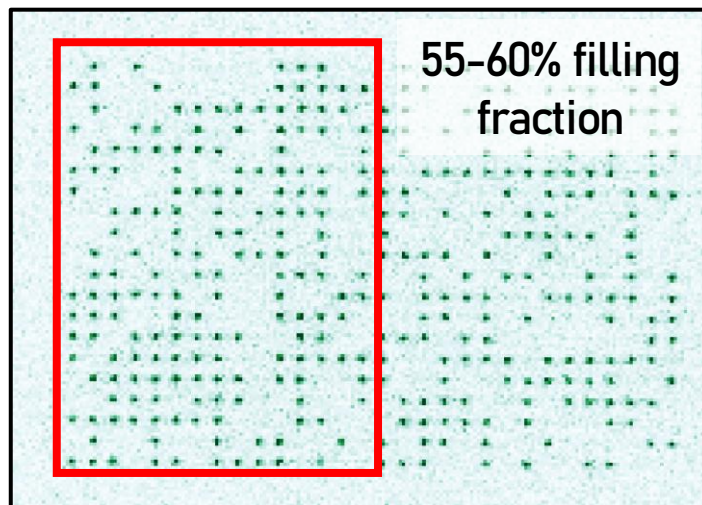
Related work: Browaeys' group;
Ahn, Regal, Endres, Kaufman, Saffman,
Thompson, Ni, Bakr, Bloch, Bernien, ...

2D array of optical tweezers

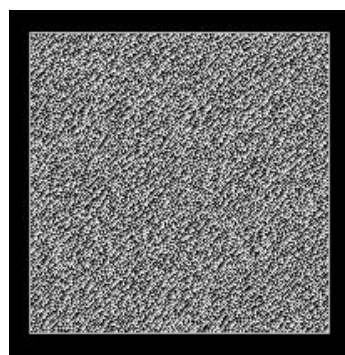


D. Kim, et al, Opt. Lett. 44, 3178 (2019)

Atoms:

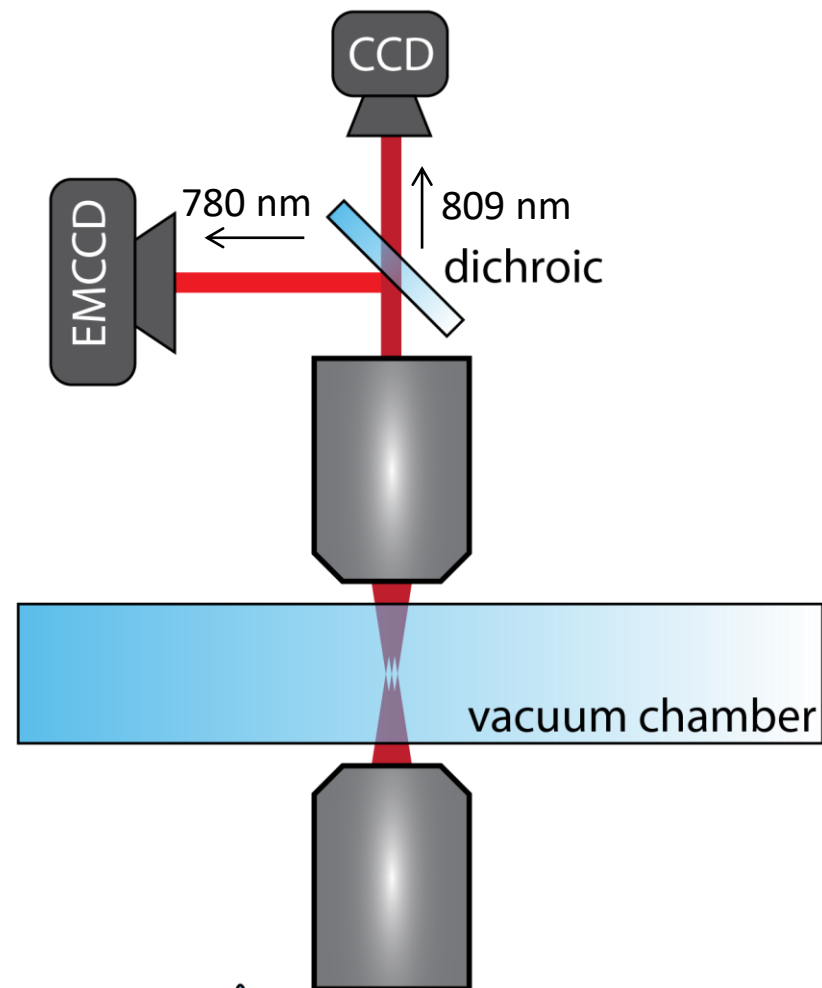


Deterministically fill
subset of traps?



phase
profile

Spatial Light
Modulator



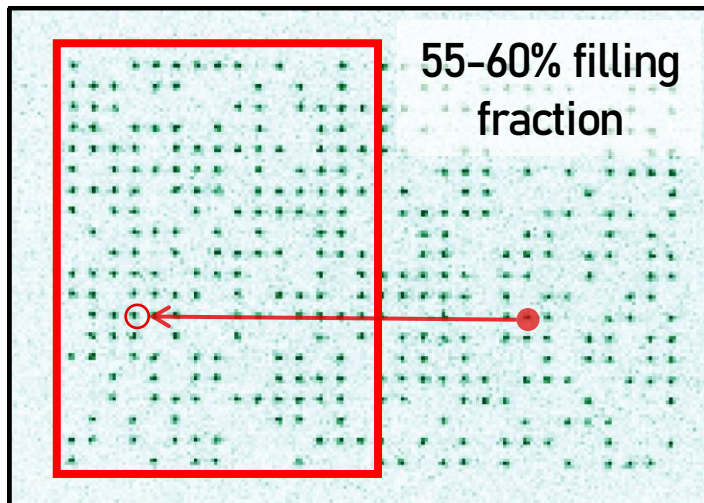
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2D array of optical tweezers

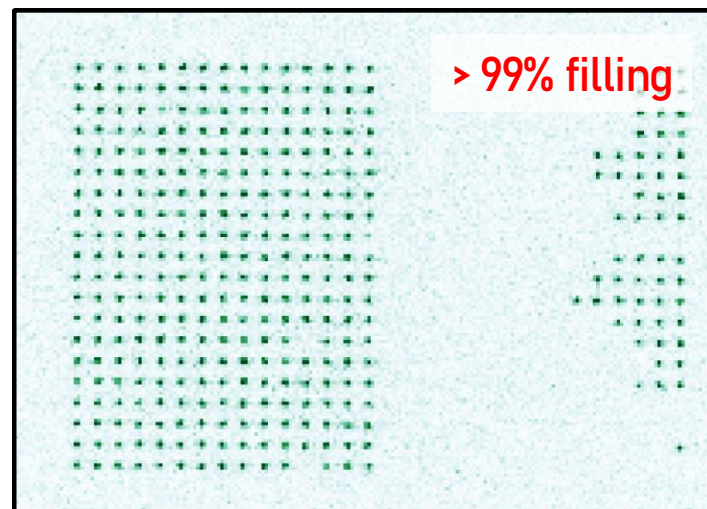


D. Kim, et al, Opt. Lett. 44, 3178 (2019)

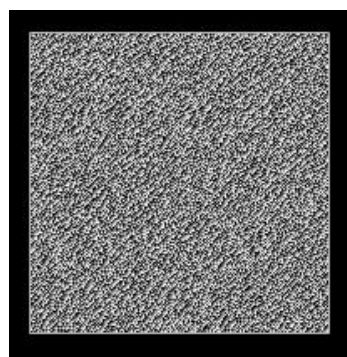
Atoms:



After sorting:



Deterministically fill
subset of traps?



phase
profile

Spatial Light
Modulator

SLM

movable
tweezers

AOD

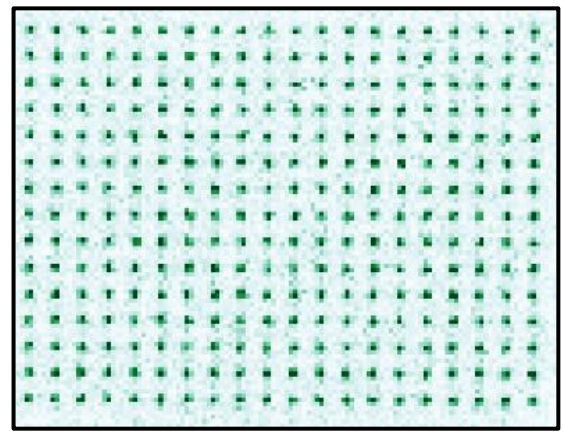
Acousto-Optical
Deflectors

vacuum chamber

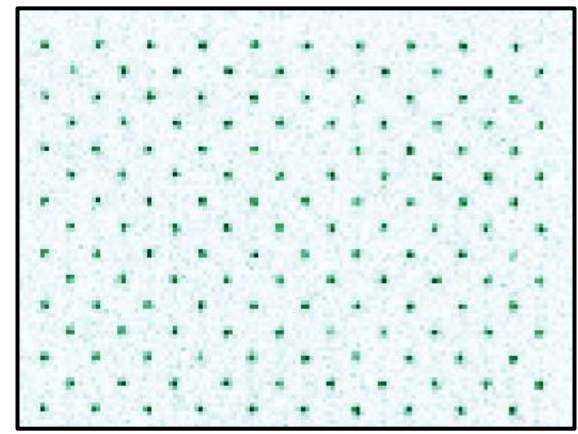
Ebadi et al Nature 2021
Related work: Browaeys' group;
Ahn, Regal, Endres, Kaufman, Saffman,
Thompson, Ni, Bakr, Bloch, Bernien, ...

Programmable 2D arrays

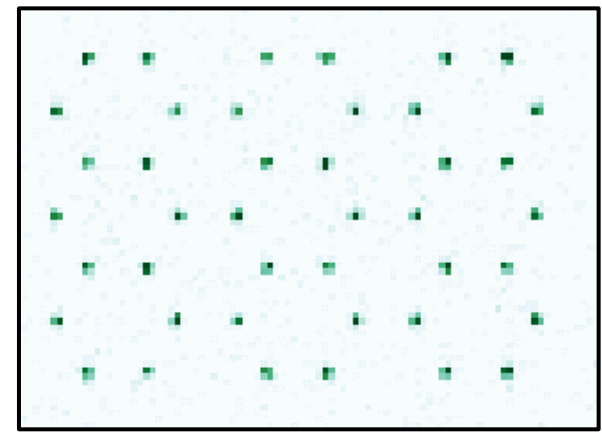
Square



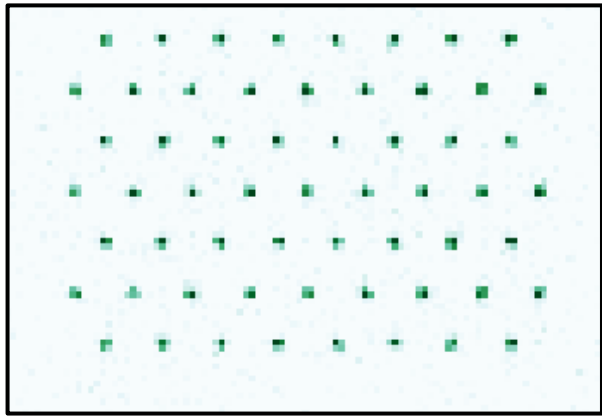
Tilted Square



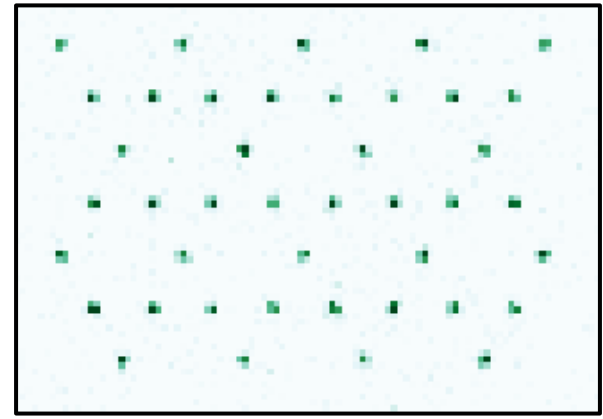
Honeycomb



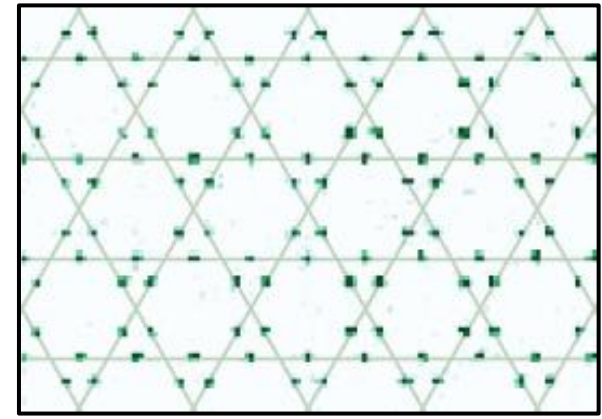
Triangular



Kagome



Link-kagome



Programmable 2D arrays

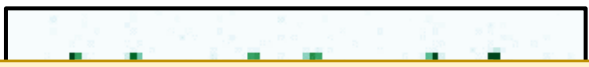
Square



Tilted Square

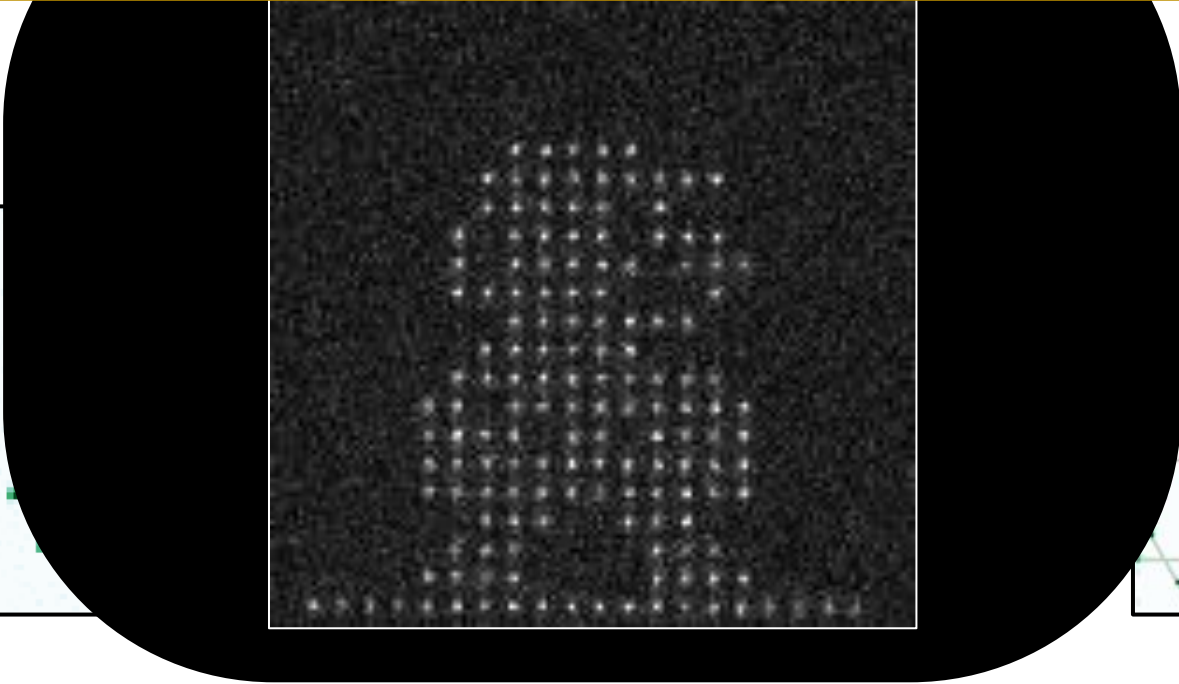
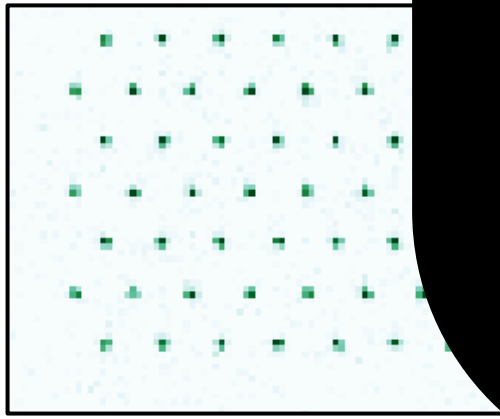


Honeycomb

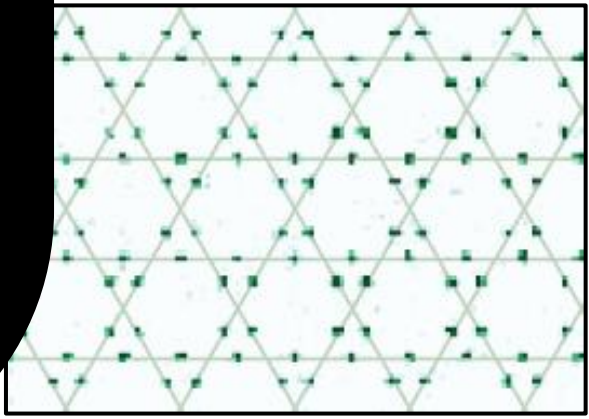


1) Programmable 2D arrays of neutral atoms

Triangular

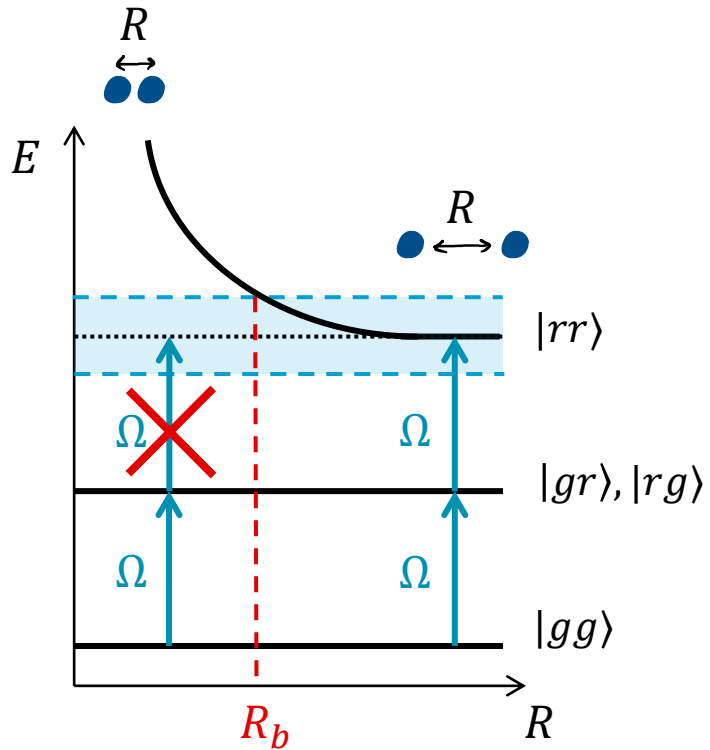


Link-kagome



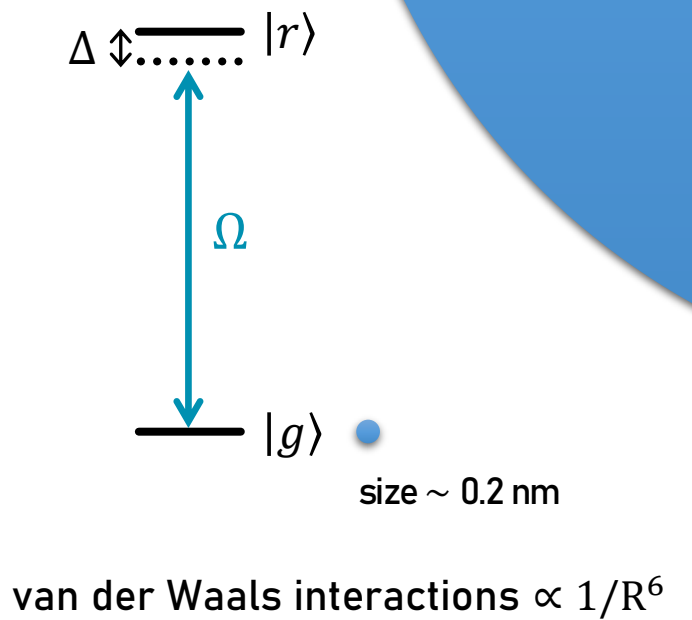
Rydberg states and long-range interactions

Rydberg blockade



blockade radius R_b :
 $V(R_b) = \Omega$

excitation to Rydberg state:



Rydberg states and long-range interactions

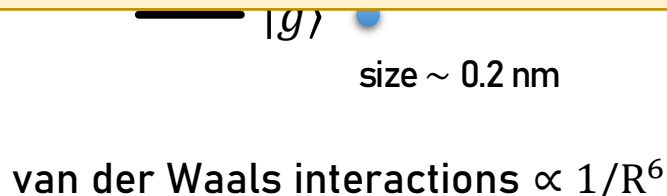
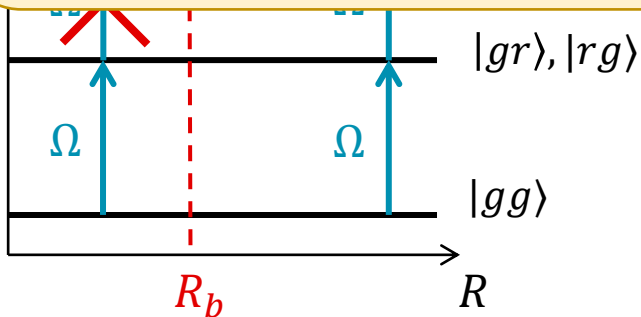
Rydberg blockade

excitation to Rydberg state:

size > 200nm



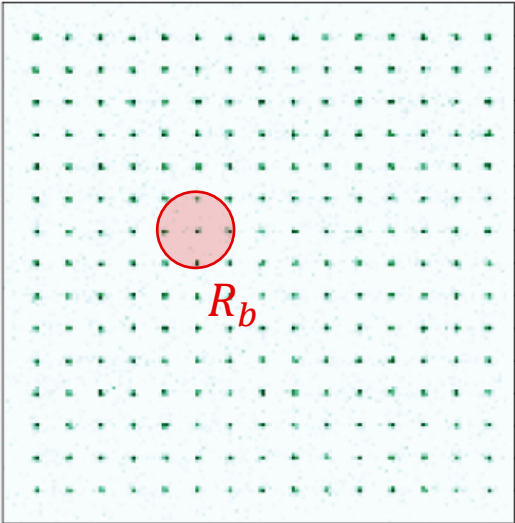
1) Programmable 2D arrays of neutral atoms
 2) with strong long-range interactions



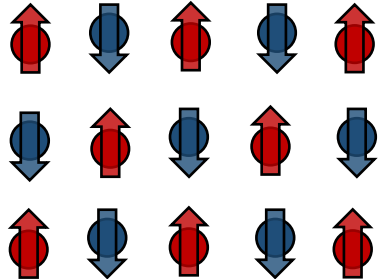
blockade radius R_b :
 $V(R_b) = \Omega$

Rydberg blockade and quantum many-body phases

Many-Body ground state
 $\Delta > 0$
 Maximize number of disconnected Rydberg atoms



antiferromagnetic state



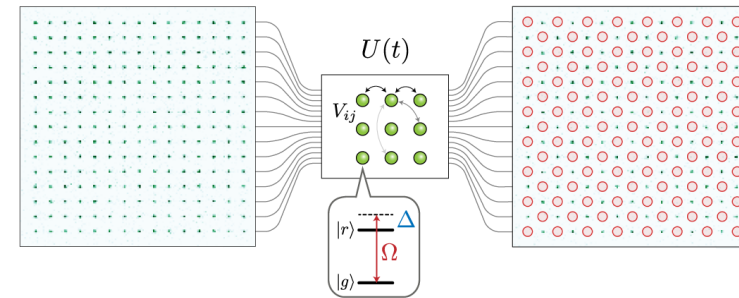
$$\mathcal{H} = \frac{1}{2} \Omega(t) \sum_i \sigma_x^{(i)} - \sum_i \Delta(t) n_i + \sum_{i < j} V_{ij} n_i n_j \quad n_i = |r_i\rangle\langle r_i|$$

drive term detuning interaction

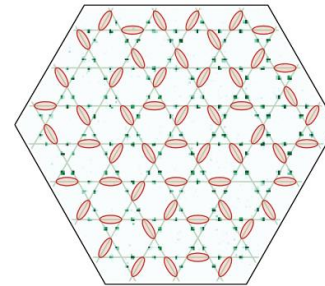
S. Ebadi et al., Quantum Phases of Matter on a 256-Atom Programmable Quantum Simulator, Nature 595, 227–232 (2021)
 P. Scholl et al., Quantum simulation of 2D antiferromagnets with hundreds of Rydberg atoms, Nature 595, 233–238 (2021)

Outline

Programmable arrays
of Rydberg atoms



Experimental evidence of a
quantum spin liquid phase



Dynamically reconfigurable
architecture for quantum
information processing



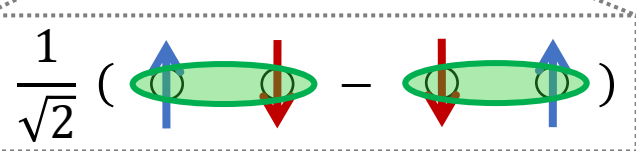
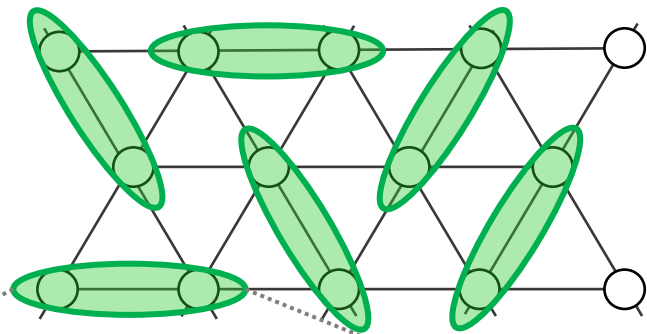
Quantum spin liquids

P. W. Anderson, *Science* 235 (1987)

The Resonating Valence Bond State in La_2CuO_4 and Superconductivity

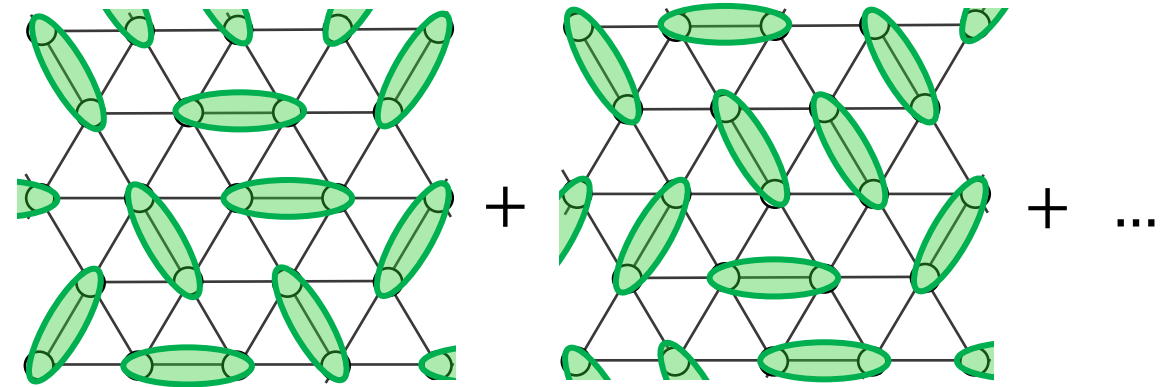
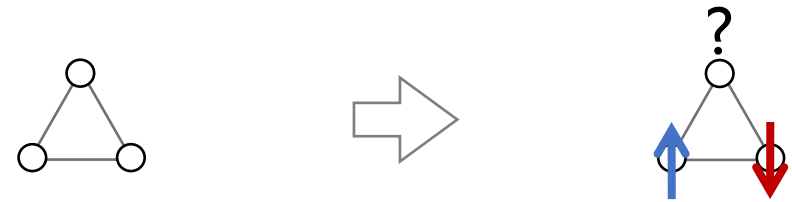
P. W. ANDERSON

The oxide superconductors, particularly those recently discovered that are based on La_2CuO_4 , have a set of peculiarities that suggest a common, unique mechanism: they tend in every case to occur near a metal-insulator transition into an odd-electron insulator with peculiar magnetic properties. This insulating phase is proposed to be the long-sought “resonating-valence-bond” state or “quantum spin liquid” hypothesized in 1973. This insulating magnetic phase is favored by low spin, low dimensionality, and magnetic frustration. The preexisting magnetic singlet pairs of the insulating state become charged superconducting pairs when the insulator is doped sufficiently strongly. The mechanism for superconductivity is hence predominantly electronic and magnetic, although weak phonon interactions may favor the state. Many unusual properties are predicted, especially of the insulating state.



entangled pair
(dimer)

spin 1/2 particles with AF interactions
on a frustrated lattice



resonating valence bond (RVB) state

Quantum spin liquids

P. W. Anderson, *Science* 235 (1987)

The Resonating Valence Bond State in La_2CuO_4 and Superconductivity

9000+ citations

P. W. ANDERSON

The oxide superconductors, particularly those recently discovered that are based on La_2CuO_4 , have a set of peculiarities that suggest a common, unique mechanism: they tend in every case to occur near a metal-insulator transition into an odd-electron insulator with peculiar magnetic properties. This insulating phase is proposed to be the long-sought “resonating-valence-bond” state or “quantum spin liquid” hypothesized in 1973. This insulating magnetic phase is favored by low spin, low dimensionality, and magnetic frustration. The preexisting magnetic singlet pairs of the insulating state become charged superconducting pairs when the insulator is doped sufficiently strongly. The mechanism for superconductivity is hence predominantly electronic and magnetic, although weak phonon interactions may favor the state. Many unusual properties are predicted, especially of the insulating state.

- no spatial order
- topological order
- long-range quantum entanglement
- anyonic non-local excitations
- emergent gauge fields
- robust ground state degeneracy
- link with high- T_c superconductivity
- application to fault-tolerant quantum computing → toric code

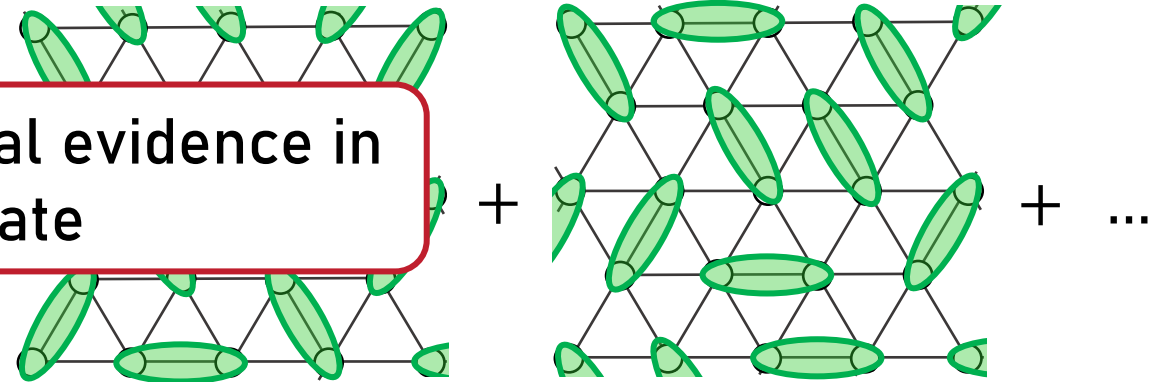
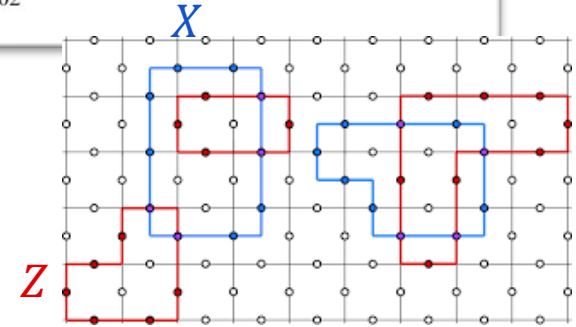
No conclusive experimental evidence in any system to date

Fault-tolerant quantum computation by anyons

A.Yu. Kitaev*

L.D. Landau Institute for Theoretical Physics, 117940, Kosygina St. 2, Germany

Received 20 May 2002



resonating valence bond (RVB) state

Quantum spin liquids in Rydberg atom arrays

R. Samajdar



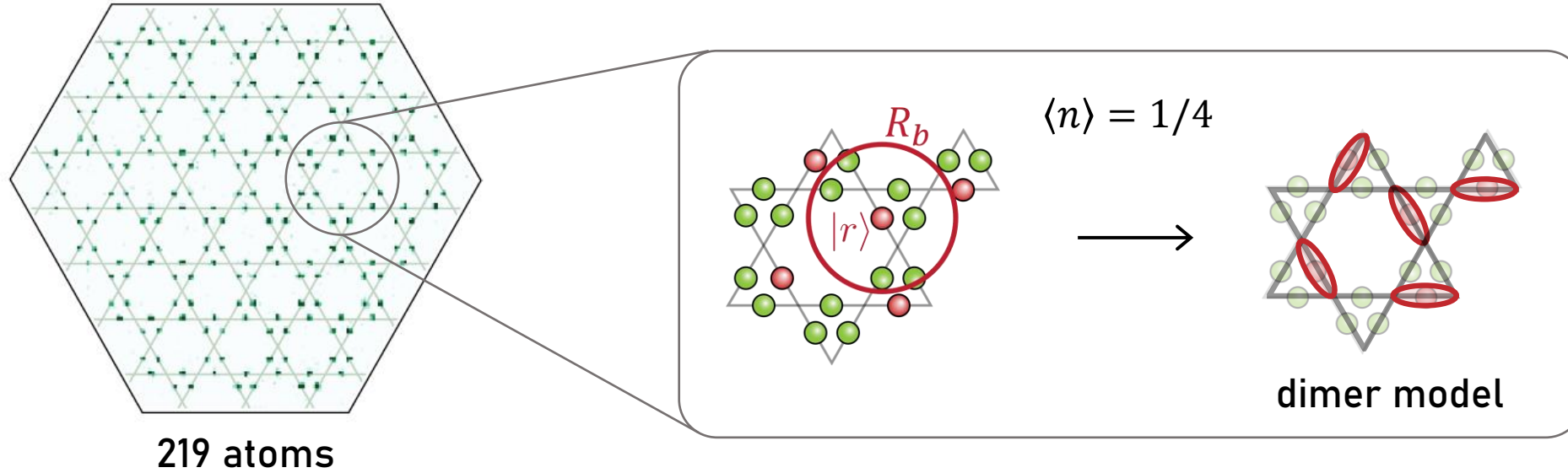
S. Sachdev



R. Verresen



A. Vishwanath



quantum spin liquid state: superposition of all dimer coverings

$$|\psi_{QSL}\rangle = \left| \begin{array}{c} \text{dimer covering 1} \\ \text{dimer covering 2} \\ \text{dimer covering 3} \\ \text{dimer covering 4} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 5} \\ \text{dimer covering 6} \\ \text{dimer covering 7} \\ \text{dimer covering 8} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 9} \\ \text{dimer covering 10} \\ \text{dimer covering 11} \\ \text{dimer covering 12} \end{array} \right\rangle + \left| \begin{array}{c} \text{dimer covering 13} \\ \text{dimer covering 14} \\ \text{dimer covering 15} \\ \text{dimer covering 16} \end{array} \right\rangle + \dots \longrightarrow \mathbb{Z}_2 \text{ topological order}$$

Quantum spin liquids in Rydberg atom arrays

R. Samajdar



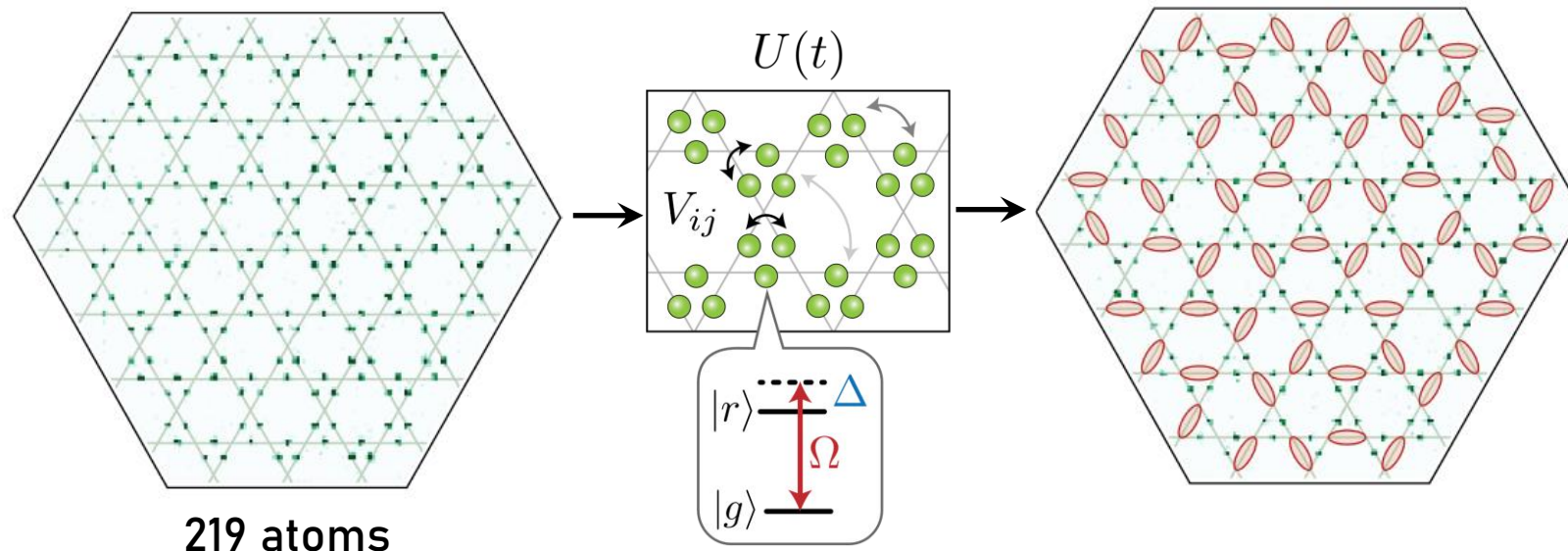
S. Sachdev



R. Verresen



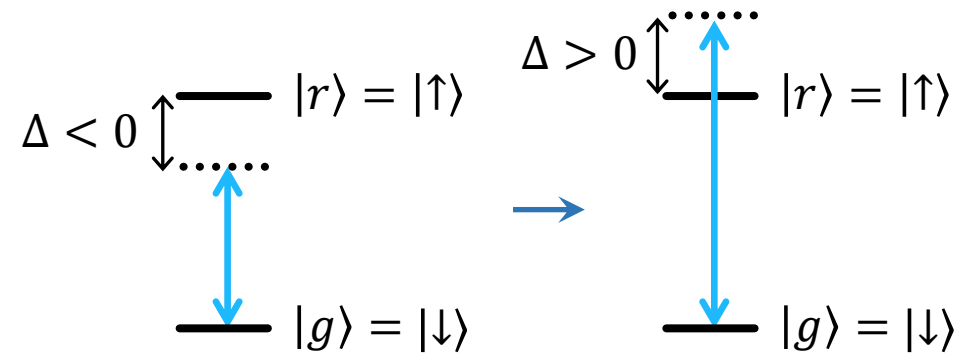
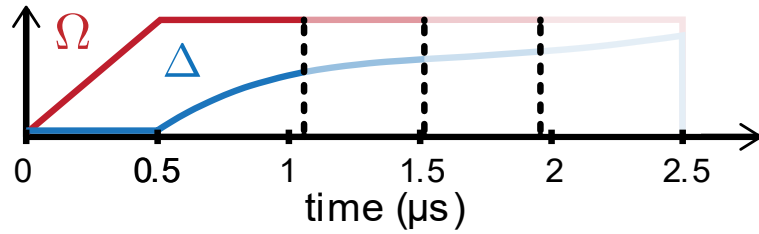
A. Vishwanath



quantum spin liquid state: superposition of all dimer coverings

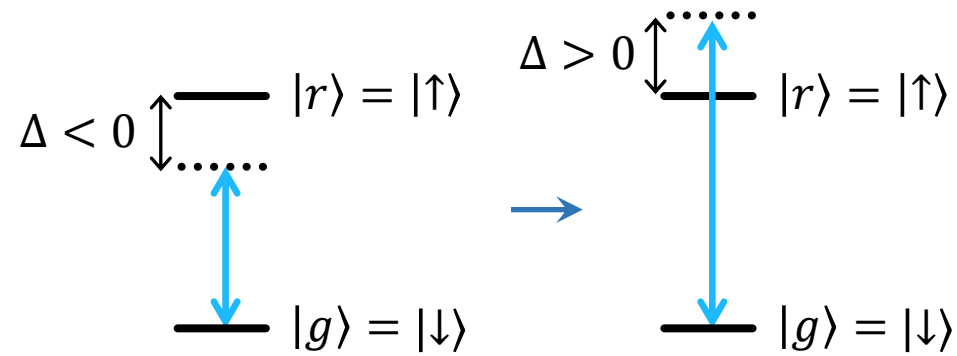
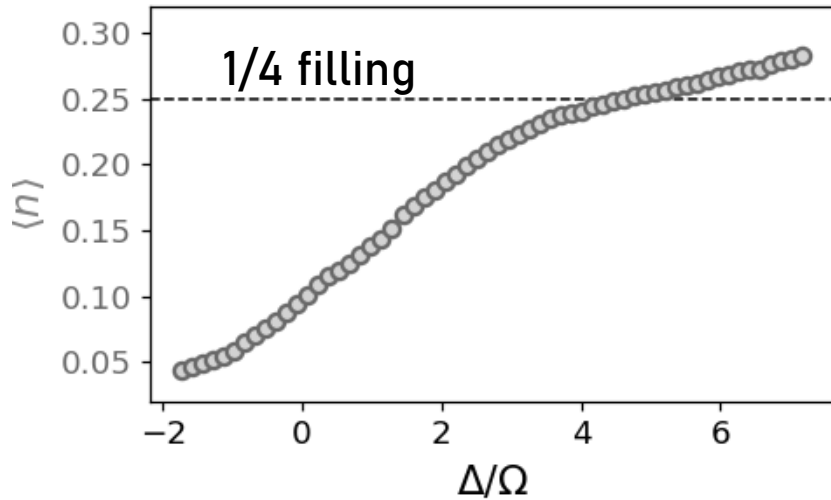
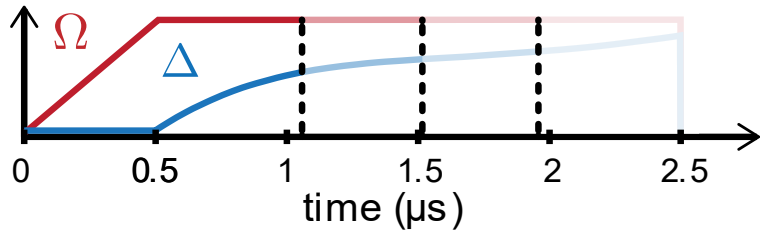
$$|\psi_{QSL}\rangle = \left| \begin{array}{c} \text{dimer covering 1} \\ \text{dimer covering 2} \\ \text{dimer covering 3} \\ \text{dimer covering 4} \end{array} \right\rangle + \dots \longrightarrow \mathbb{Z}_2 \text{ topological order}$$

Quasi-adiabatic preparation of a **dimer phase**



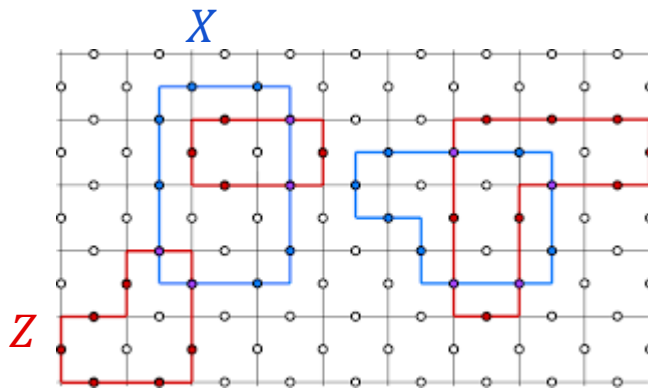
$$\mathcal{H} = \frac{1}{2}\Omega(t) \sum_i \sigma_x^{(i)} - \sum_i \Delta(t) n_i + \sum_{i<j} V_{ij} n_i n_j$$

Quasi-adiabatic preparation of a **dimer phase**



$$\mathcal{H} = \frac{1}{2} \Omega(t) \sum_i \sigma_x^{(i)} - \sum_i \Delta(t) n_i + \sum_{i < j} V_{ij} n_i n_j$$

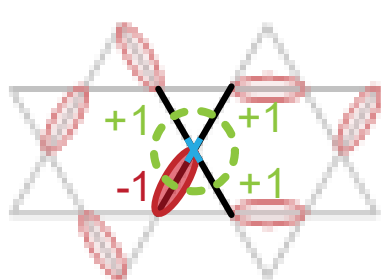
topological string operators associated with a \mathbb{Z}_2 quantum spin liquid (toric code)



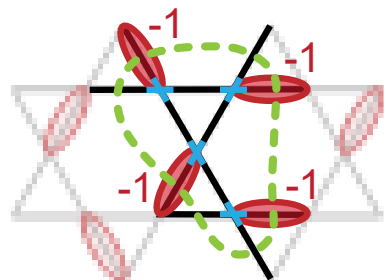
Probing topological correlations

1 Identify dimer phase

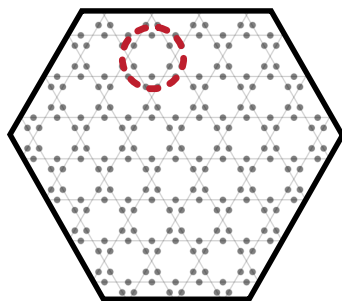
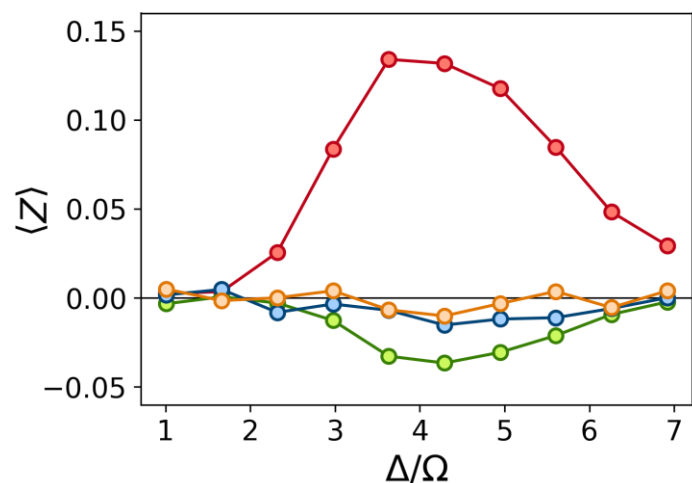
diagonal string operator Z :
parity of dimers along a string



$$\langle Z \rangle = -1$$



$$\langle Z \rangle = (-1)^{\# \text{ enclosed vertices}}$$

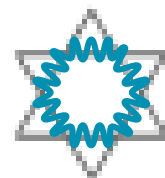


$$|\psi_{QSL}\rangle = \left| \begin{array}{c} \text{dimer configuration 1} \\ \text{dimer configuration 2} \\ \text{dimer configuration 3} \\ \text{dimer configuration 4} \end{array} \right\rangle + \dots$$

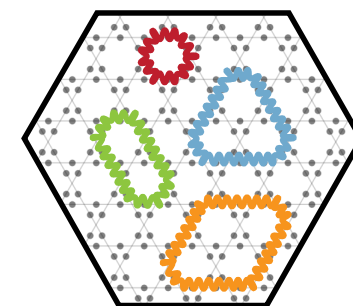
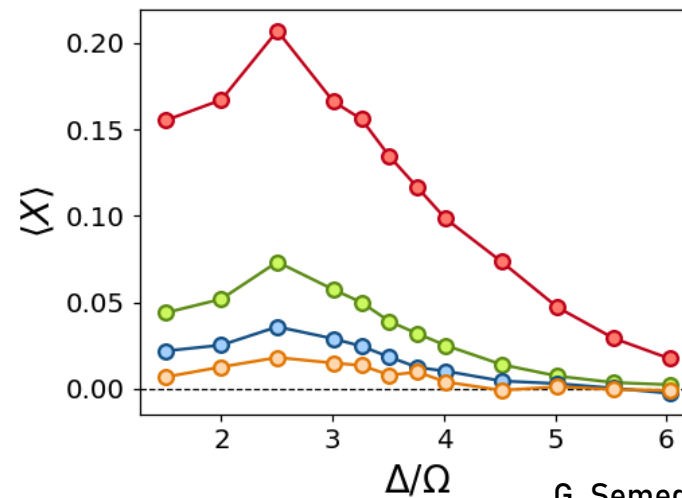
2 Probe coherence between dimer states

off-diagonal string operator X :

$$\langle X \rangle = \langle \text{dimer state 1} | \text{dimer state 2} \rangle + \dots$$

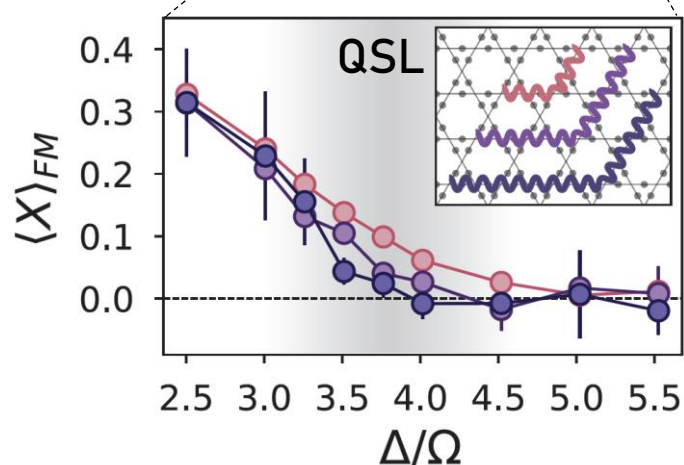
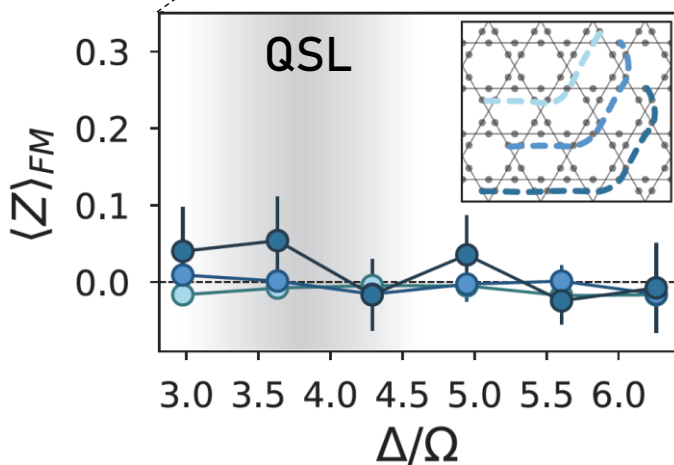
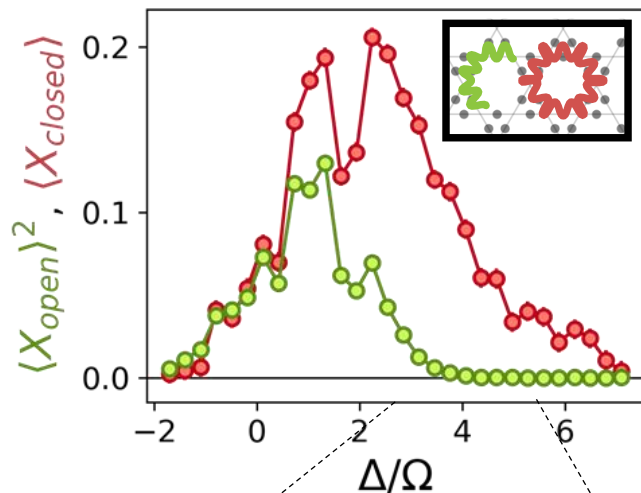
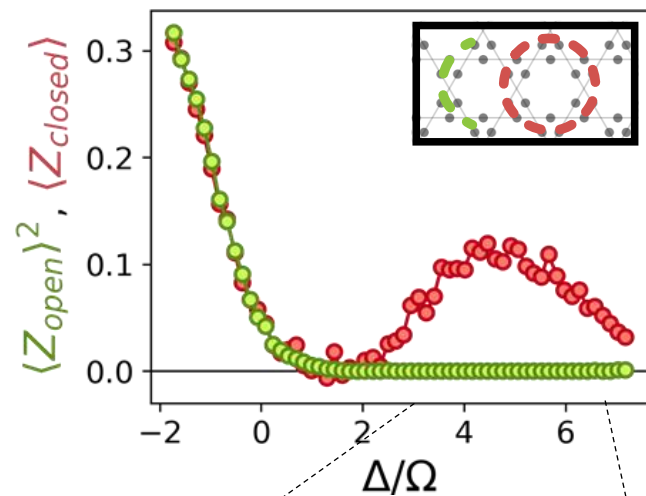


$\langle X \rangle > 0 \rightarrow$ coherence between dimer coverings



Order parameters for QSL phase

closed loops vs open strings



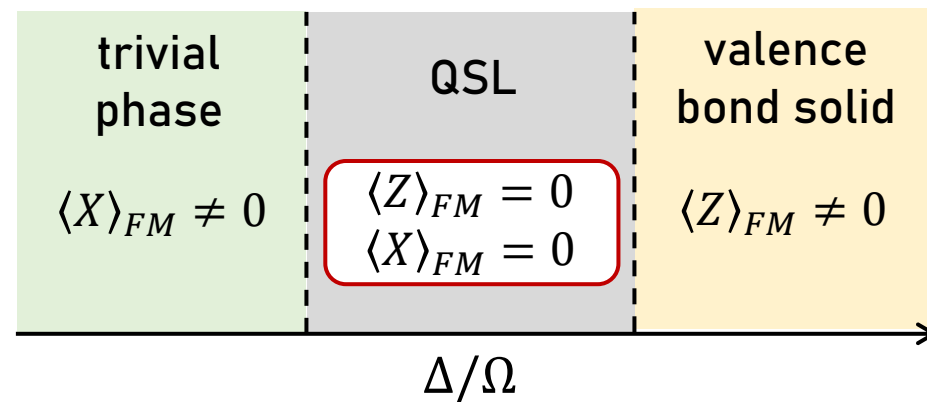
closed loops: detect non-trivial topological correlations

open strings: distinguish QSL from nearby phases

FM string order parameters:

$$\langle Z \rangle_{FM} = \frac{\text{[Lattice with blue open strings]}}{\sqrt{\text{[Lattice with blue closed loops]}}}$$

$$\langle X \rangle_{FM} = \frac{\text{[Lattice with purple open strings]}}{\sqrt{\text{[Lattice with purple closed loops]}}}$$

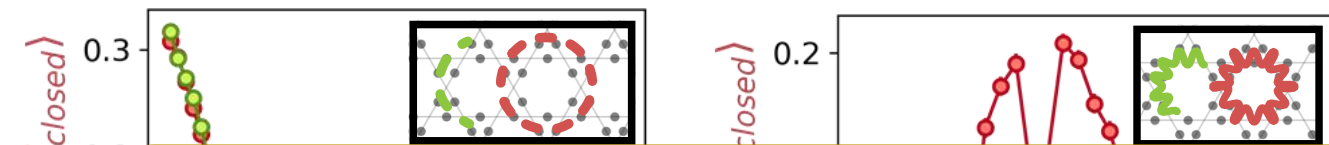


Order parameters for QSL phase

closed loops vs open strings

closed loops: detect non-trivial topological correlations

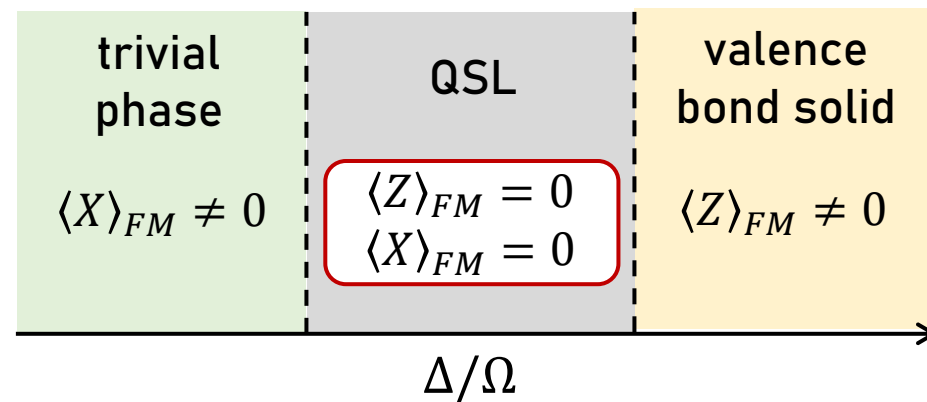
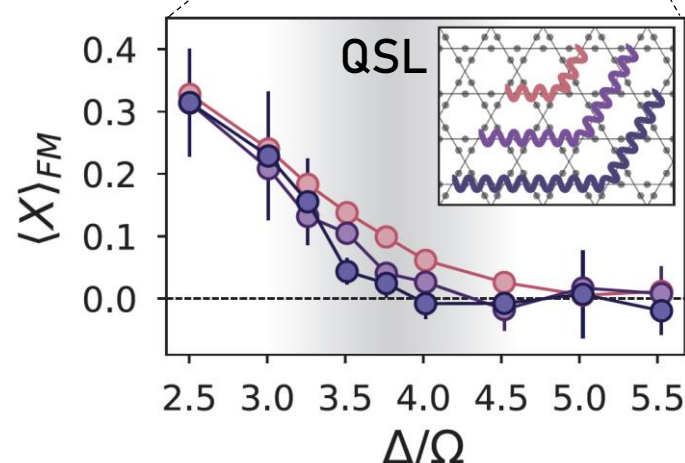
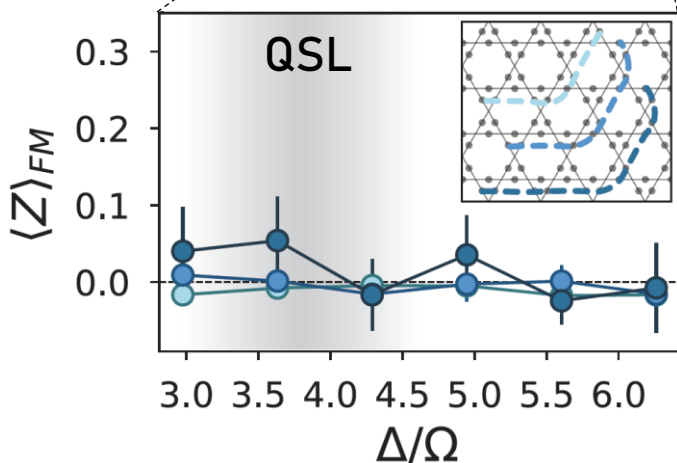
open strings: distinguish QSL from nearby phases



onset of a quantum spin liquid phase!

$$|\psi_{QSL}\rangle = \left| \begin{array}{c} \text{Lattice with red ovals} \end{array} \right\rangle + \left| \begin{array}{c} \text{Lattice with red ovals} \end{array} \right\rangle + \left| \begin{array}{c} \text{Lattice with red ovals} \end{array} \right\rangle + \left| \begin{array}{c} \text{Lattice with red ovals} \end{array} \right\rangle + \dots$$

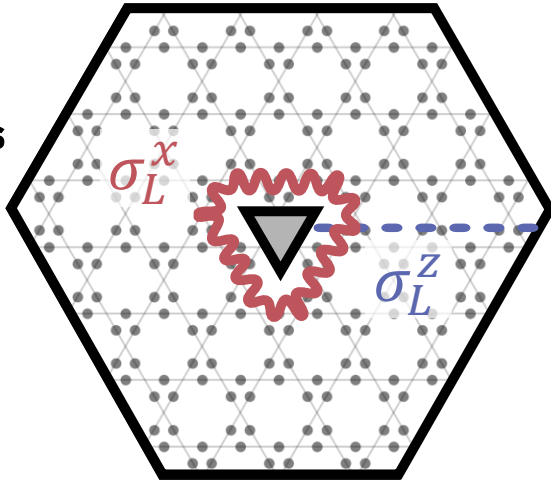
$$\langle X \rangle_{FM} = \left| \begin{array}{c} \text{Lattice with purple wavy lines} \end{array} \right\rangle + \left| \begin{array}{c} \text{Lattice with purple wavy lines} \end{array} \right\rangle + \dots$$



Towards a topological qubit

non-trivial topology:
array with a hole

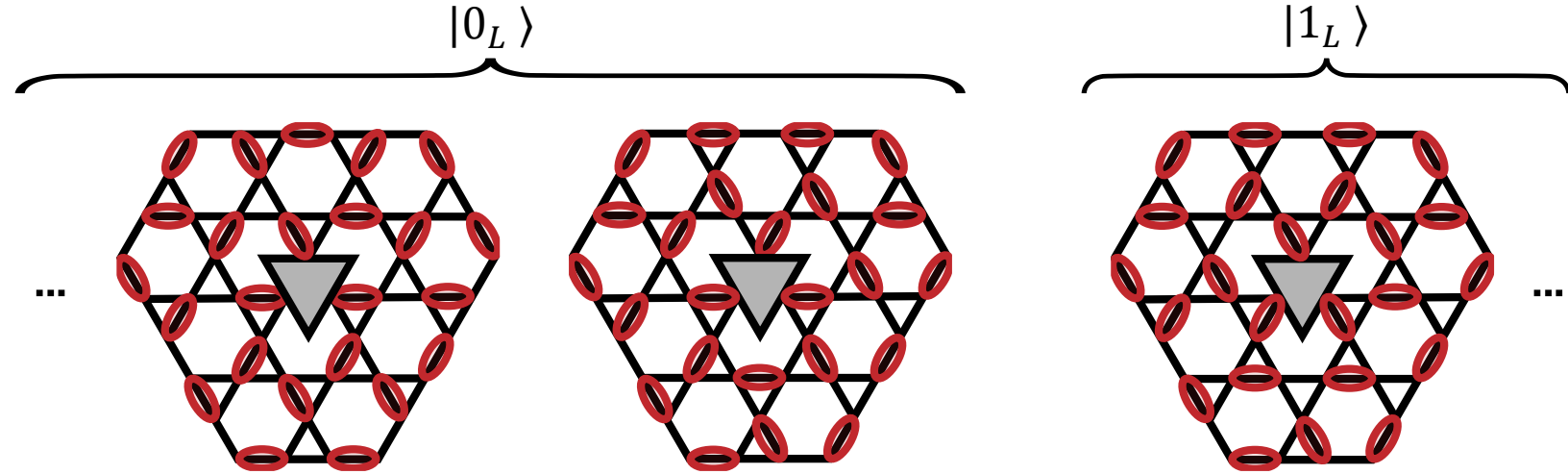
logical operators



two distinct topological sectors



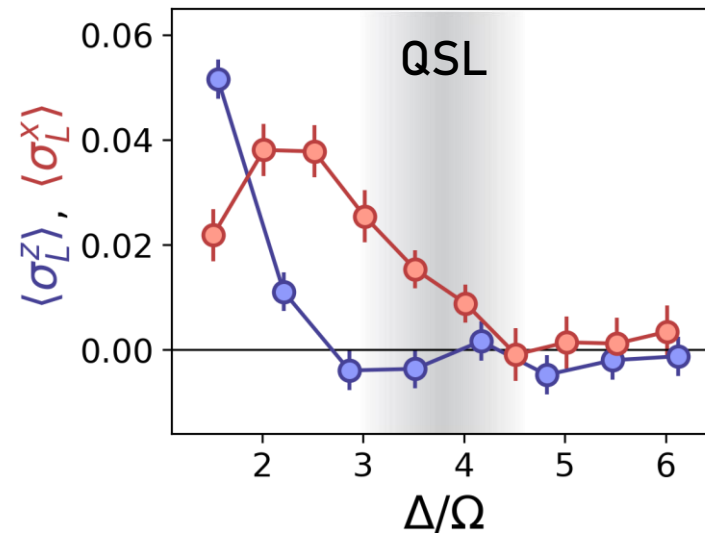
logical states of the topological qubit



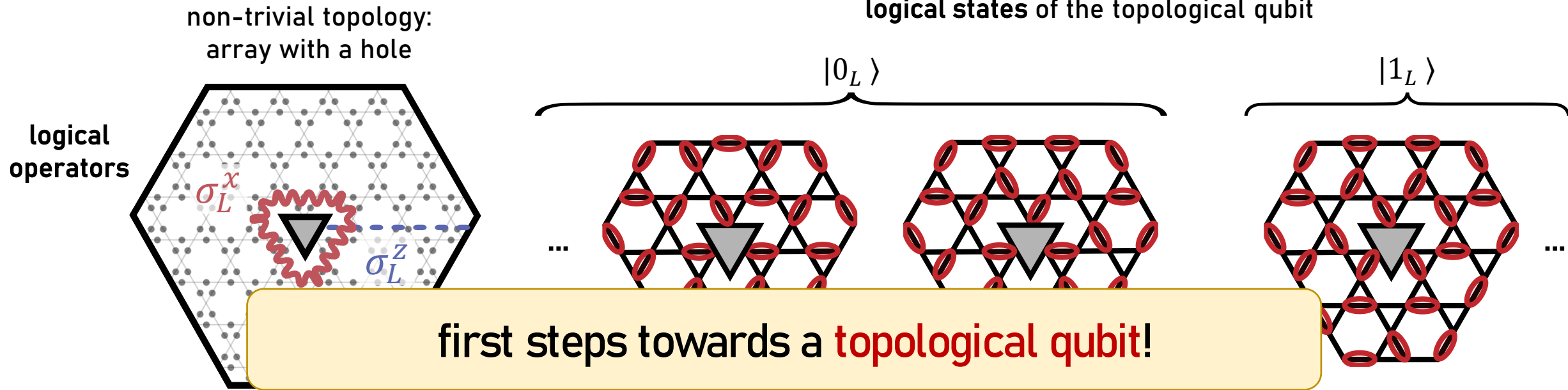
Two-fold degenerate ground state:

$$|+\rangle \sim \frac{|0_L\rangle + |1_L\rangle}{\sqrt{2}} \rightarrow \langle \sigma_L^x \rangle = +1, \langle \sigma_L^z \rangle = 0$$

$$|-\rangle \sim \frac{|0_L\rangle - |1_L\rangle}{\sqrt{2}} \rightarrow \langle \sigma_L^x \rangle = -1, \langle \sigma_L^z \rangle = 0$$



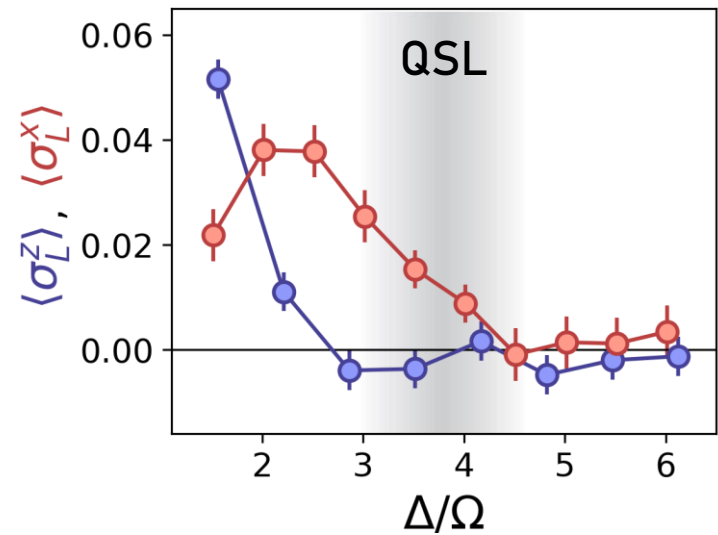
Towards a topological qubit



Two-fold degenerate ground state:

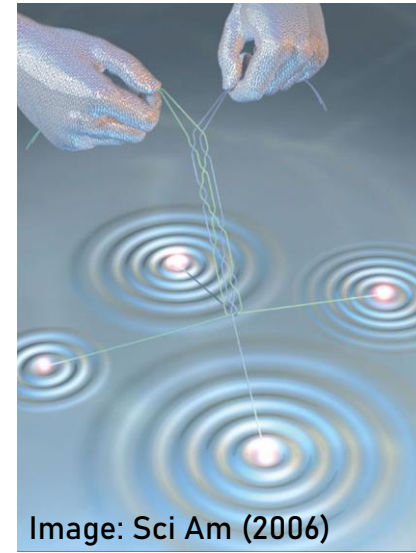
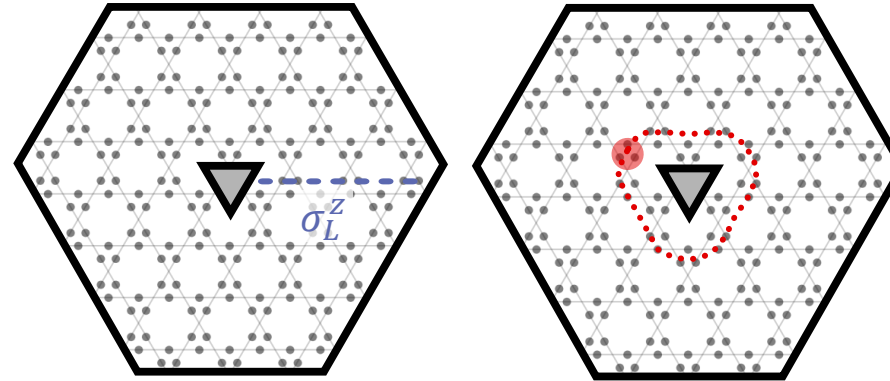
$$|+\rangle \sim \frac{|0_L\rangle + |1_L\rangle}{\sqrt{2}} \rightarrow \langle \sigma_L^x \rangle = +1, \langle \sigma_L^z \rangle = 0$$

$$|-\rangle \sim \frac{|0_L\rangle - |1_L\rangle}{\sqrt{2}} \rightarrow \langle \sigma_L^x \rangle = -1, \langle \sigma_L^z \rangle = 0$$

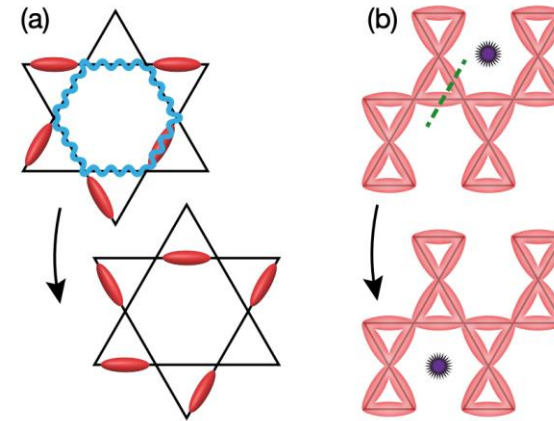


Outlook

- topological qubit encoding and manipulation requires arrays of local addressing laser beams (dynamically tunable on fast time scales)



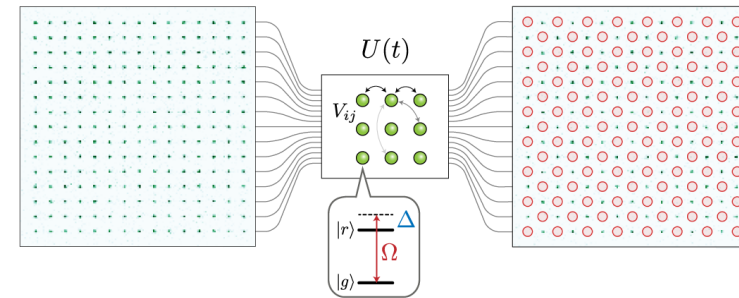
- error correction on the topological phase detect and annihilate anyons using ancilla atoms



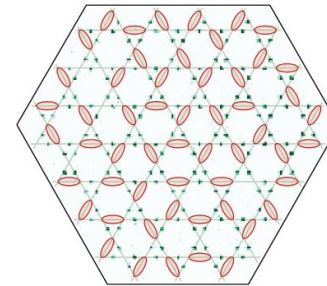
- platform for exploring fault-tolerant quantum information processing

Outline

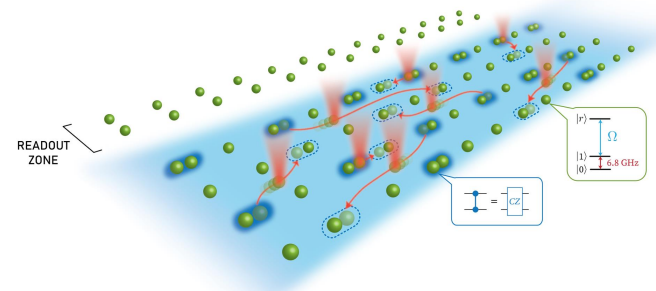
Programmable arrays
of Rydberg atoms



Experimental evidence of a
quantum spin liquid phase



Dynamically reconfigurable
architecture for quantum
information processing

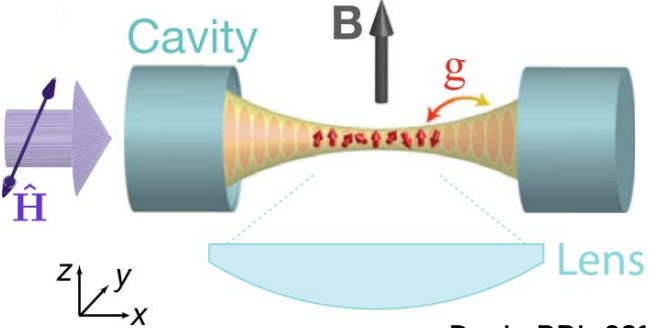


Dynamically reconfigurable architecture

Nonlocal connectivity and entanglement: approaches

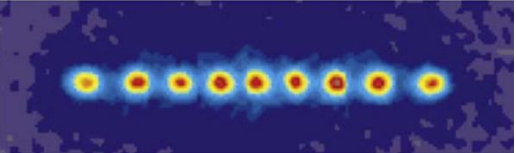
Quantum data bus

Cavity-mediated interactions



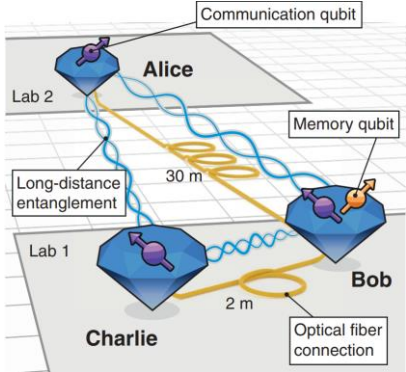
Davis PRL 2019

Shared motional modes

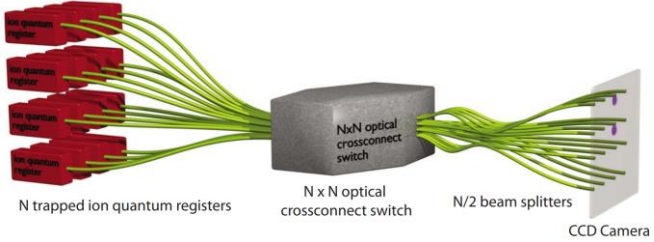


Innsbruck

Quantum networks



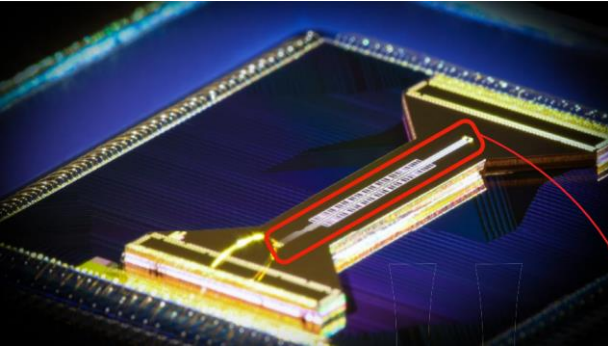
Pompili ... Hanson (Science 2021)



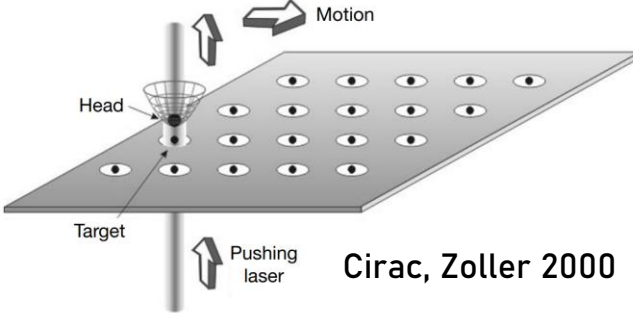
Monroe ... Duan, Kim

See also SC qubit transduction to optical photons

Quantum ion "CCD"



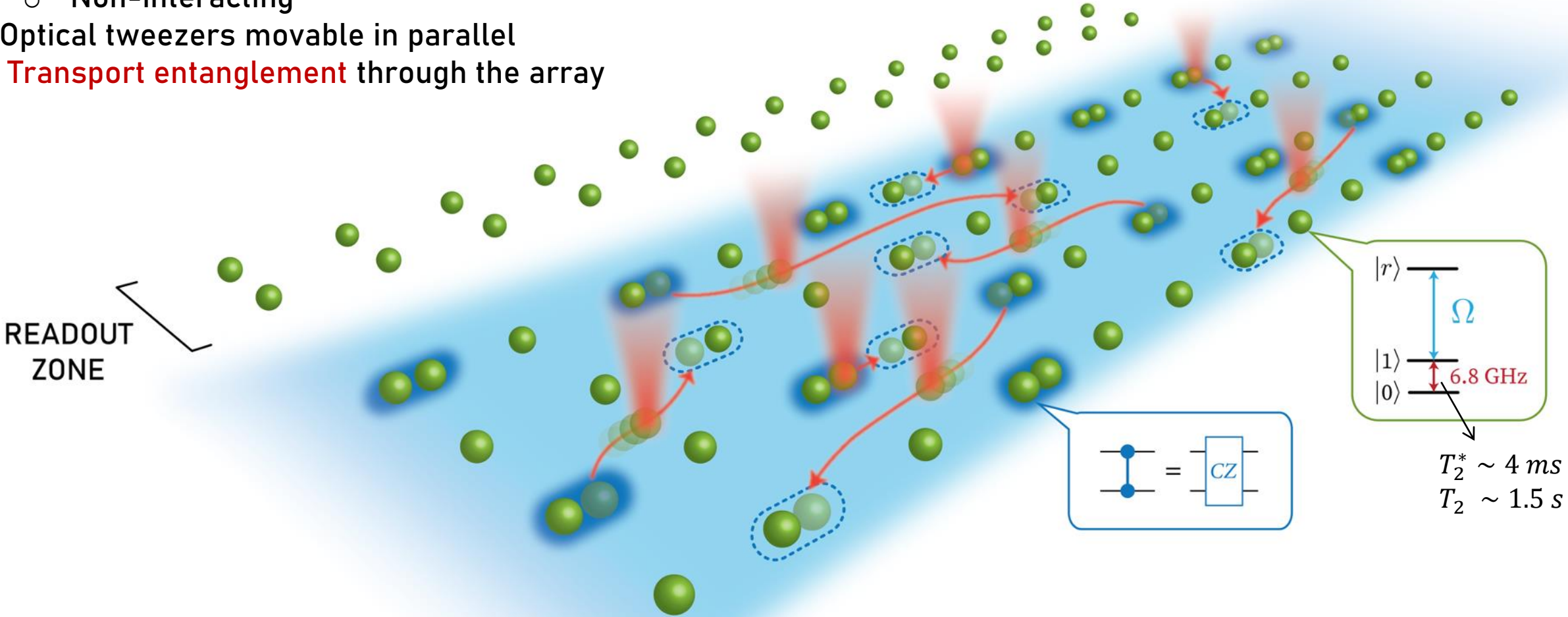
Honeywell 2021



Cirac, Zoller 2000

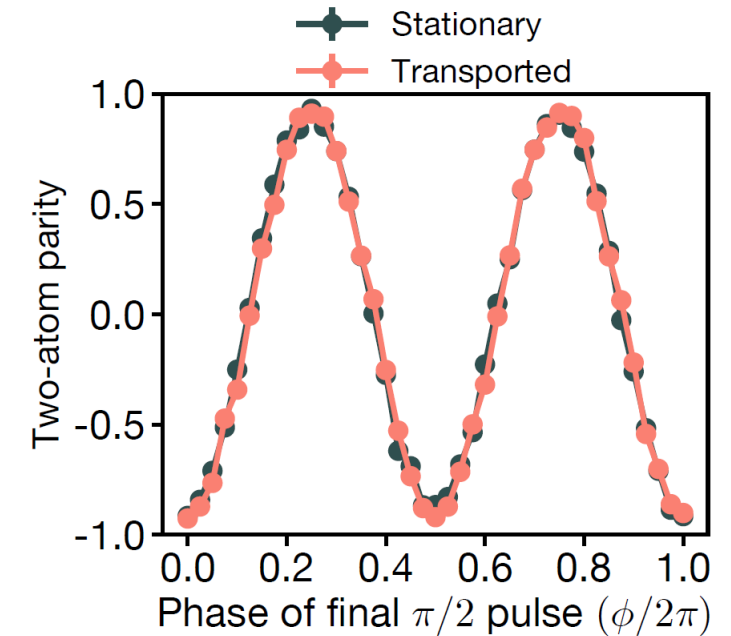
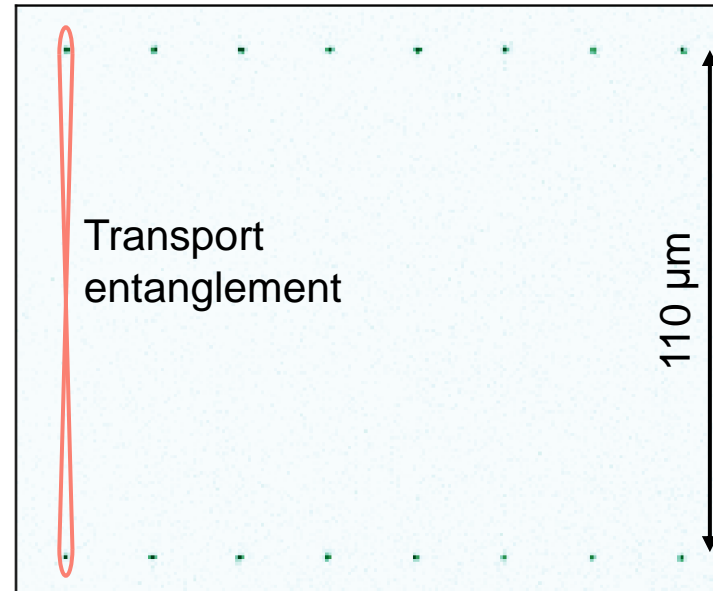
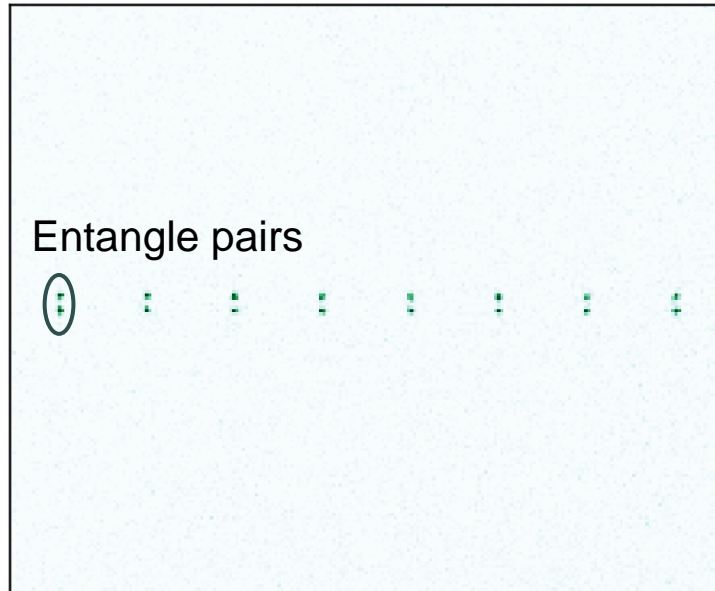
Dynamically reconfigurable architecture

- Use Rydberg for generating **entanglement**
- Store information in hyperfine qubit
 - Long **coherence times**
 - Non-interacting
- Optical tweezers movable in parallel
- **Transport entanglement** through the array

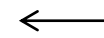


Dynamically reconfigurable architecture

Transporting entanglement across the array



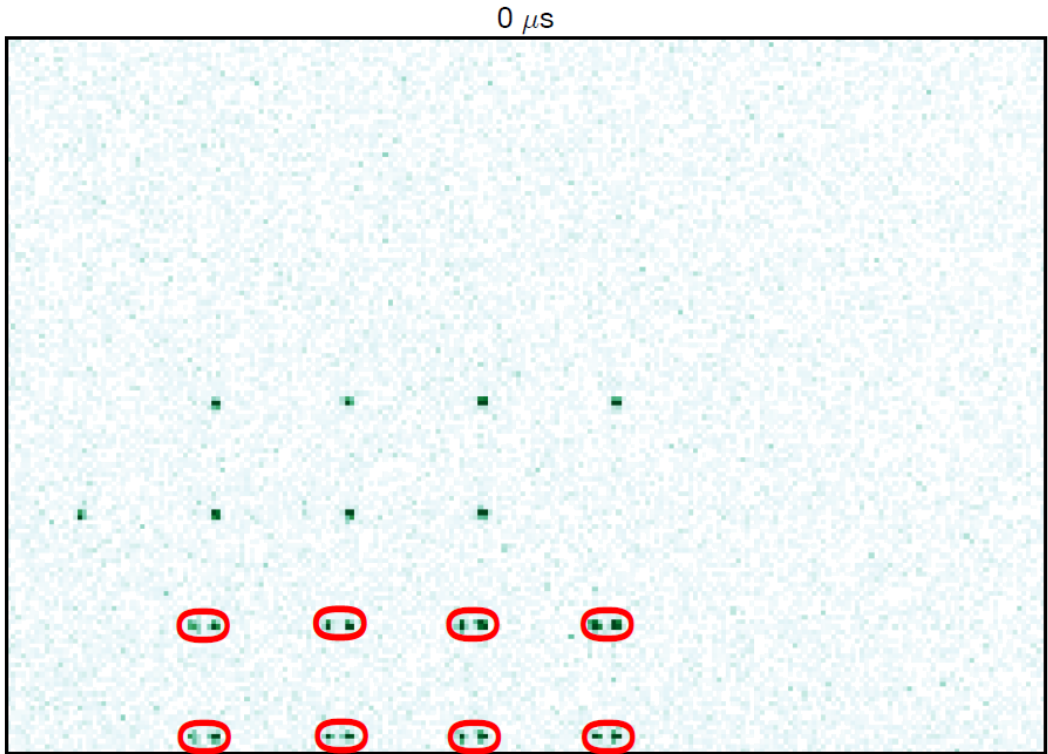
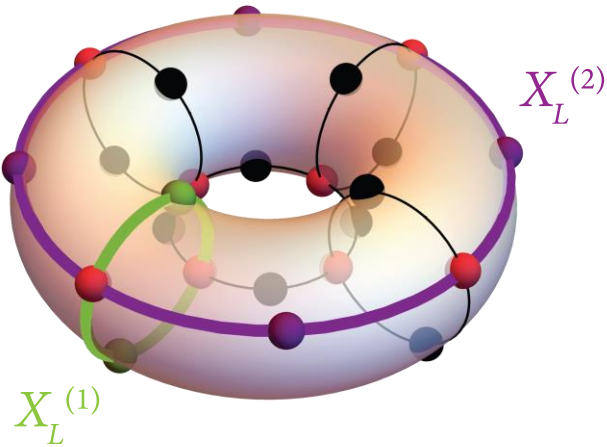
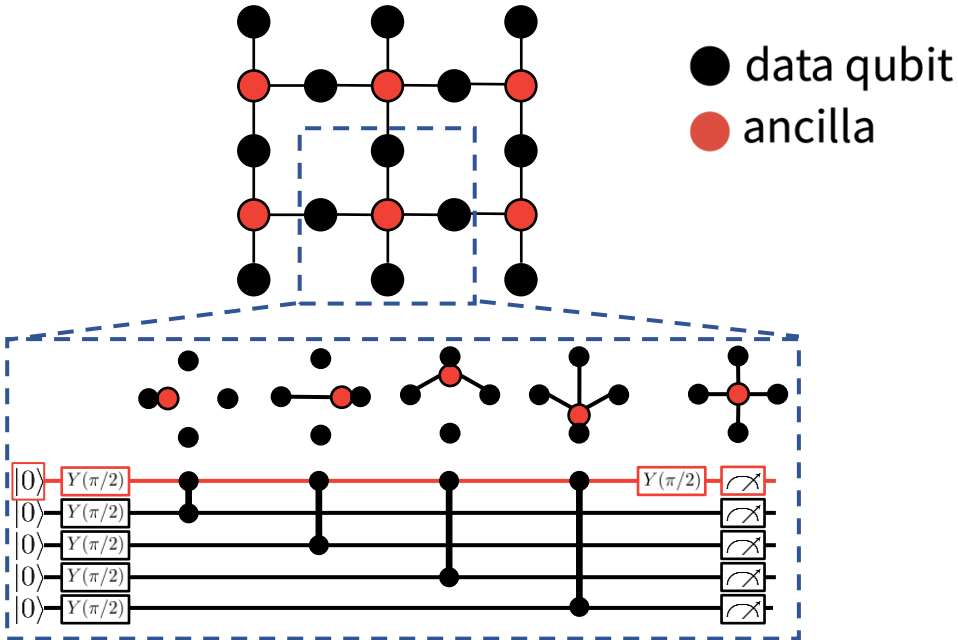
Many potential applications:
complex quantum computing architectures
& new probes for many-body phases



Coherence is preserved when
transporting the atoms over a hundred μm
in a few hundred μs ($\sim 10^{-4} T_2$)

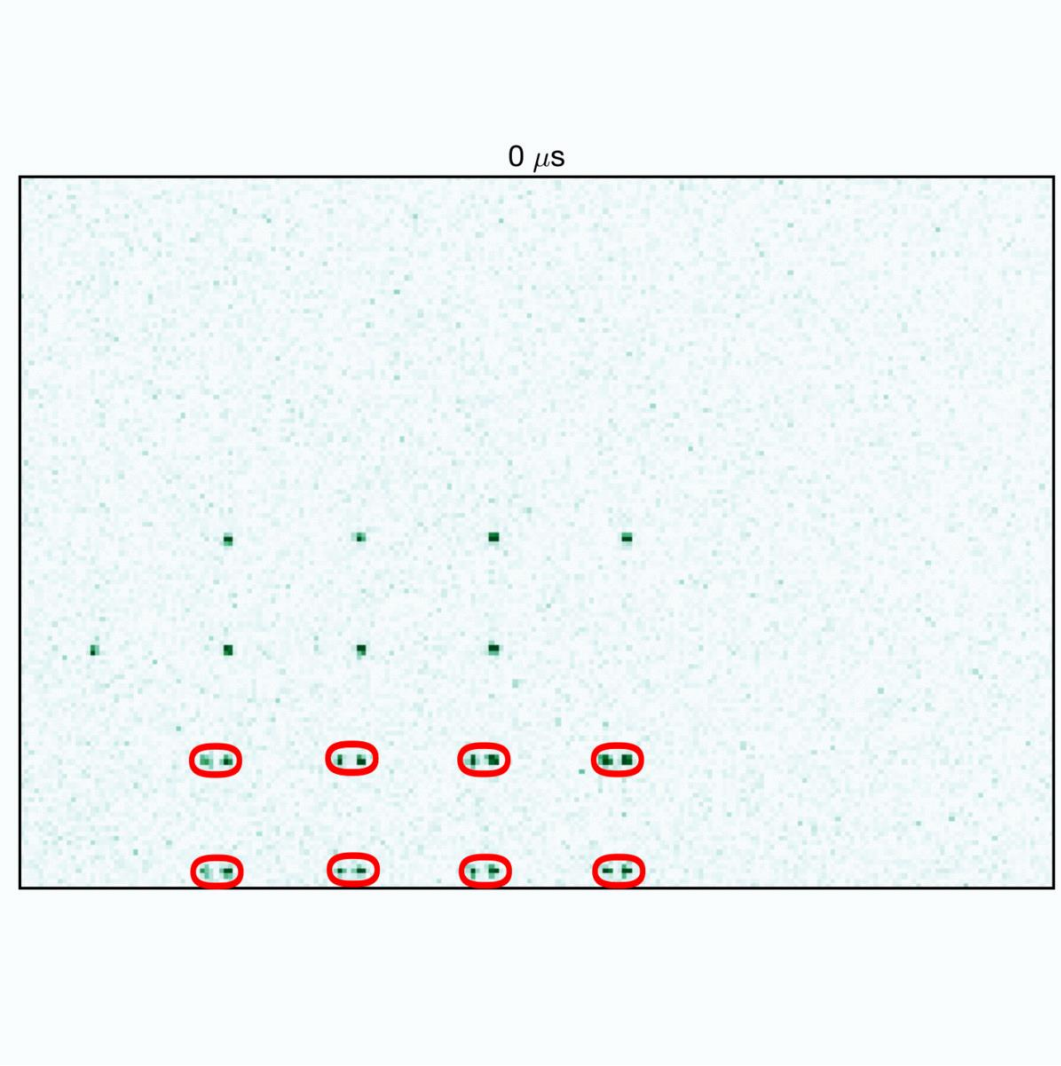
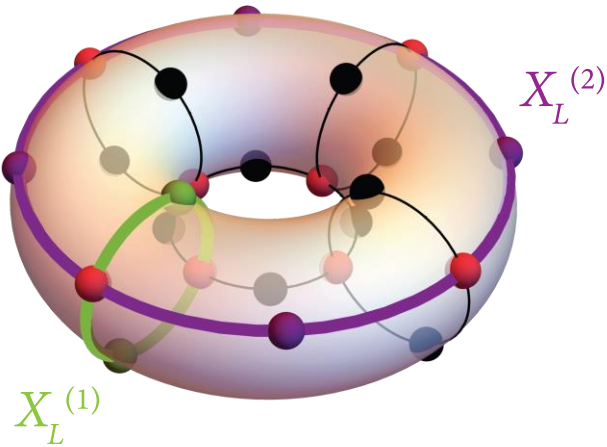
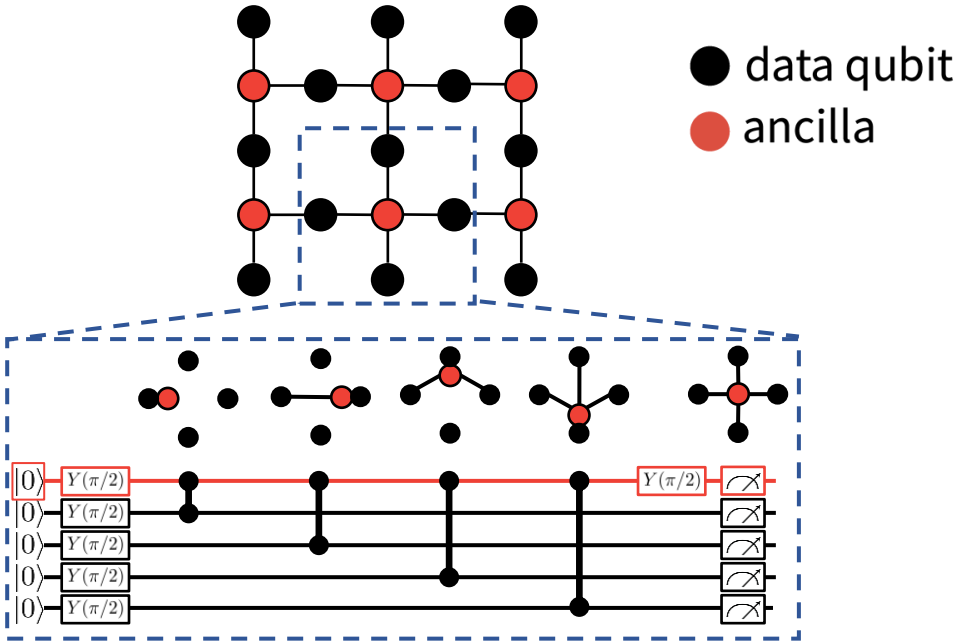
Dynamically reconfigurable architecture

Toric code on a 3D torus



Dynamically reconfigurable architecture

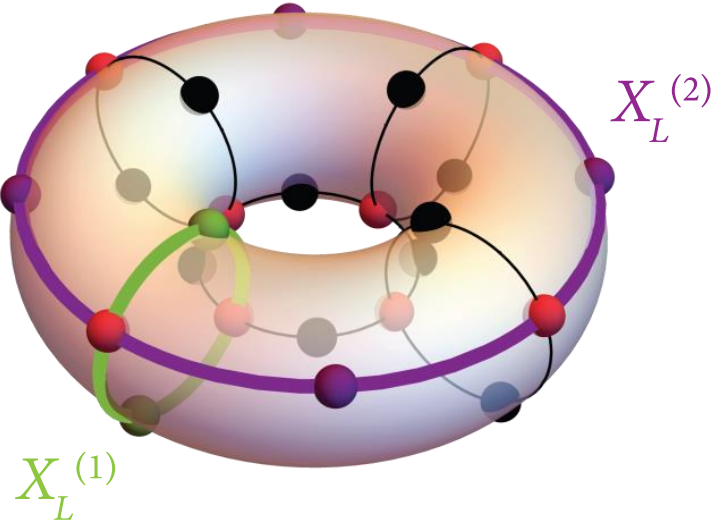
Toric code on a 3D torus



Dynamically reconfigurable architecture

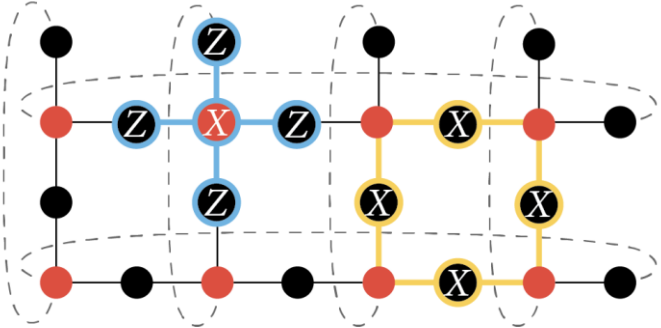
Toric code on a 3D torus

● data qubit
● ancilla

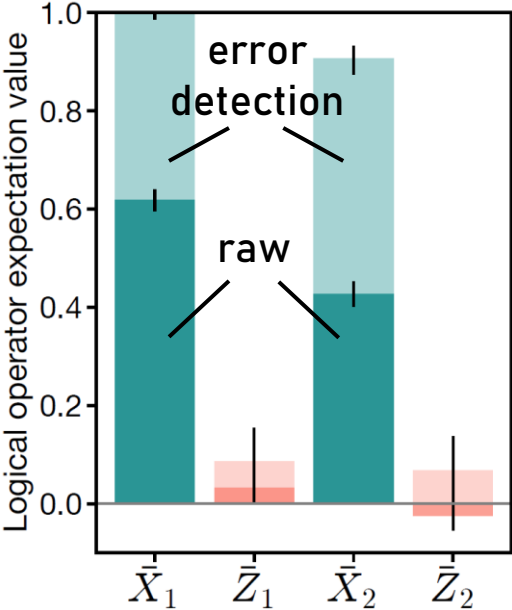
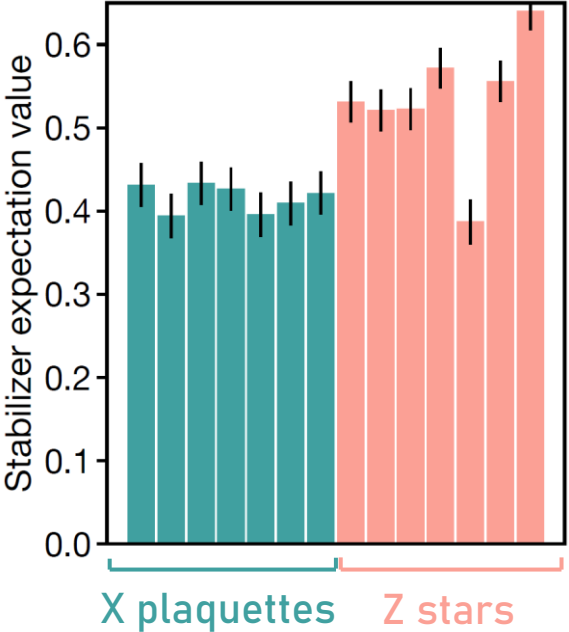
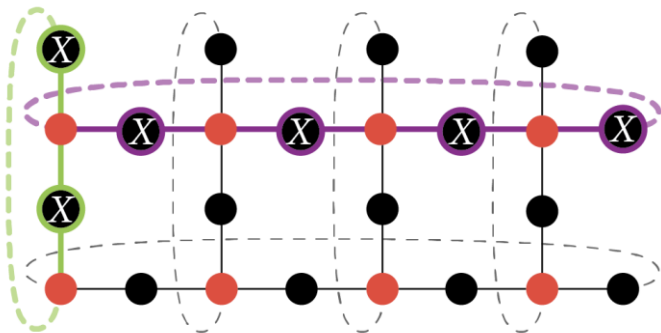


realization of **QEC** via highly parallel atom control

Stabilizers



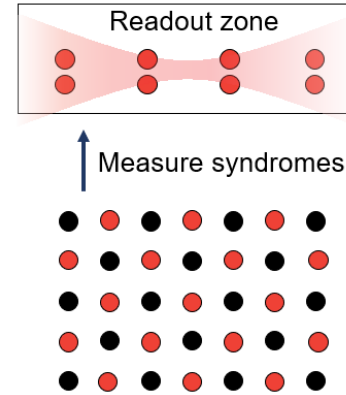
Two logical qubits!



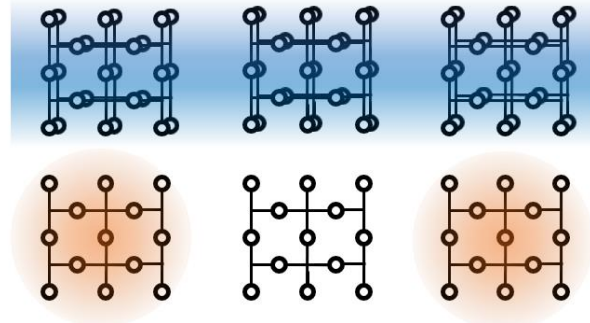
Outlook

- Logical state preservation

Repetitive, mid-circuit measurement and correction (repetition code, surface code)



- Logical qubit gates



- Fault-tolerant logical algorithms (~10s logical qubit) **empowered by highly-parallel optical control**

Fiber Amplifier

2020-08-04 20:36:20

Swann

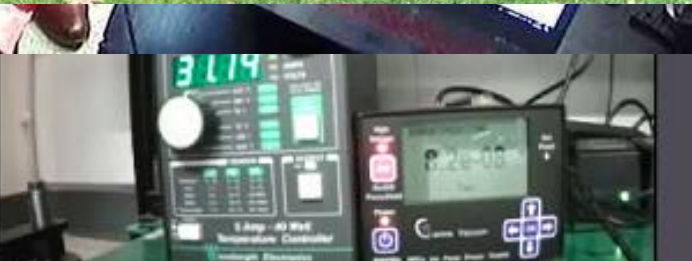
Rydberg lasers

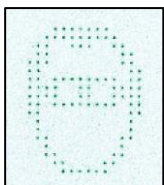
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Swann

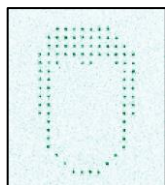
Lab life during the pandemic

March 17th, 2020

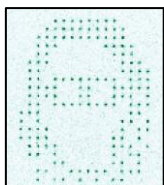




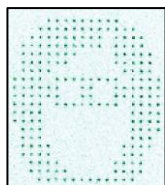
Harry Levine



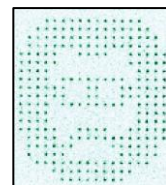
Tout Wang



Sepehr Ebadi



GS



Alex Keesling



Dolev Bluvstein



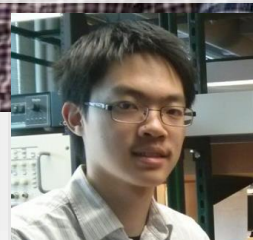
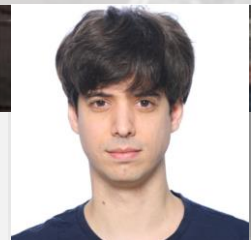
Ahmed Omran



Tom Manovitz

Simon Evered

Harry Zhou



Pls: M. Lukin




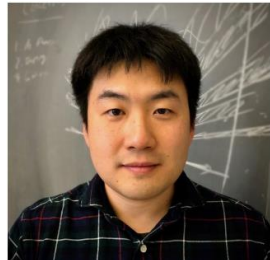
V. Vuletic



M. Greiner



THEORY COLLABORATORS

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 A. Vishwanath	 S. Sachdev	 H. Pichler	 S. Choi

Funding: NSF, CUA, MPHQ, Vannevar Bush, AFOSR MURI, DARPA, DOE, NDSEG, QuEra, AWS

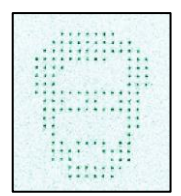
Interested in joining the newborn Semeghini's group at Harvard as a PhD student or PostDoc?



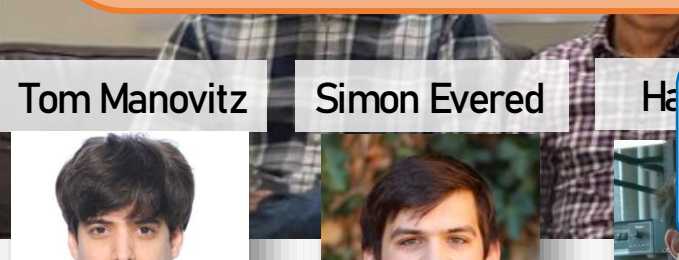
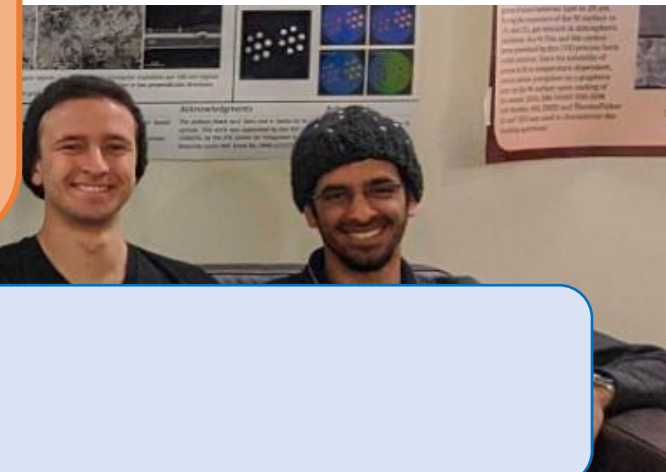
Contact: semeghini@fas.harvard.edu



Bluvstein



Ahmed Omran



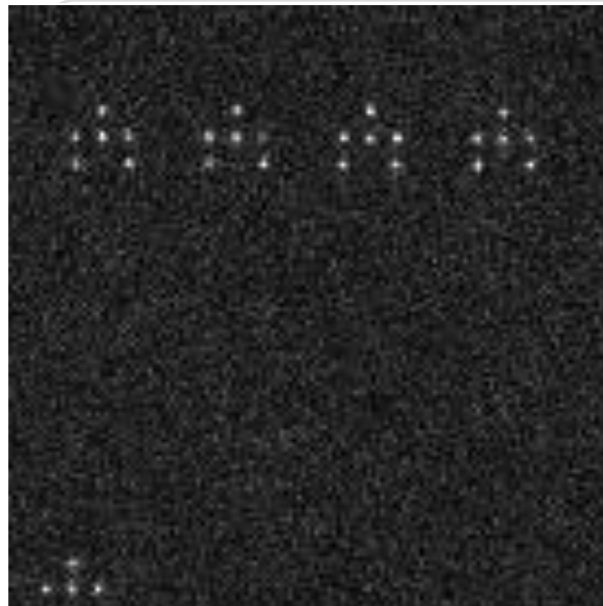
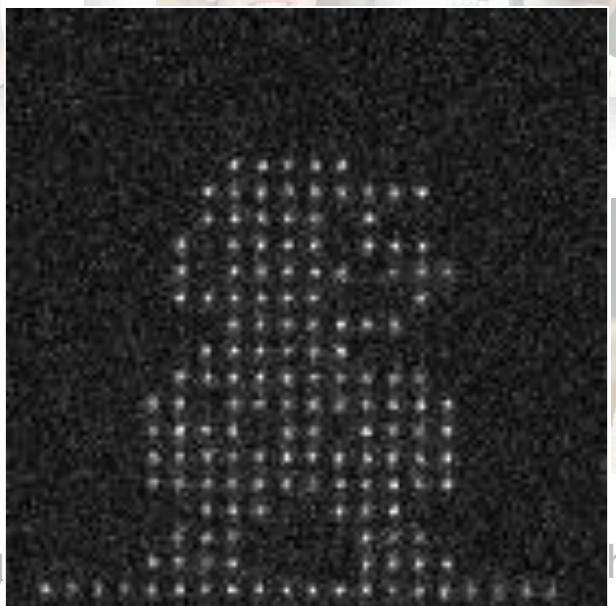
Tom Manovitz

Simon Evered

Ha

Thank you!

Pls:



Funding: MURI, DARPA, DOE, NUSEG, QuEra, AWS