

School of Physics & Astronomy

FACULTY OF MATHS & PHYSICAL SCIENCES



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Skyrmions in chiral magnetic multilayers

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[@ChrisMarrows](https://twitter.com/ChrisMarrows)



Acknowledgements



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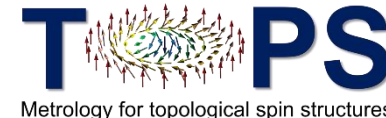
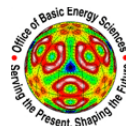


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– Paul Scherrer Inst., Switzerland

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Research Council



Chiral Magnetic Skyrmions

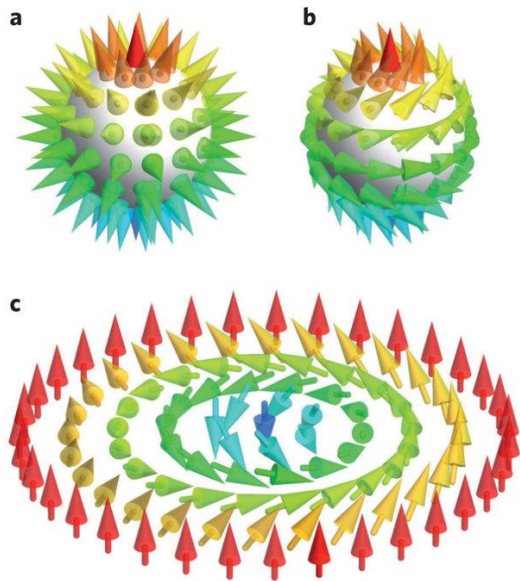


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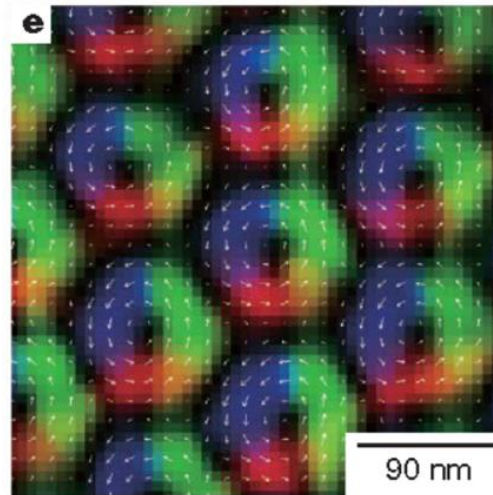
Topologically stable vector field object – integer winding numbers

“Combed hedgehog”

$$S = \frac{1}{4\pi} \int \hat{\mathbf{m}} \left(\frac{\partial \hat{\mathbf{m}}}{\partial x} \times \frac{\partial \hat{\mathbf{m}}}{\partial y} \right) dx dy$$



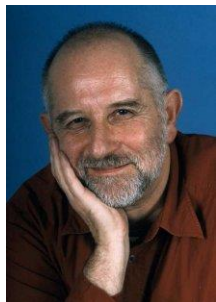
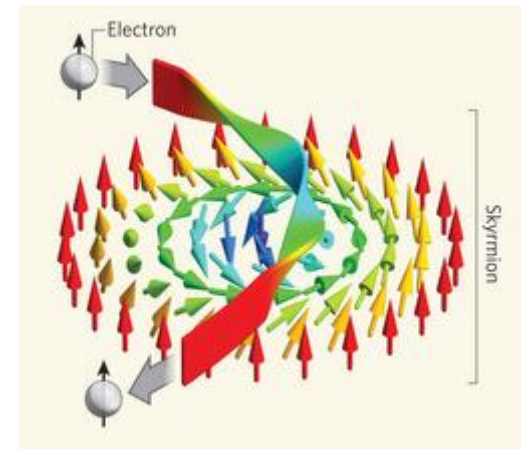
Skyrmion crystal



Fe_{0.5}Co_{0.5}Si - Yu Nature (2010)



Tony Skyrme FRS



Sir Michael Berry FRS

Emergent electrodynamics arising from Berry phase
Each skyrmion = φ_0 of fictitious magnetic flux
Moving skyrmions => effective electric field

Topology



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Topology is counting...

zero



one



two

...



Möbius strip



zero

one

Topology provides stability

“You cannot have half a hole in something”

Dzyaloshinskii-Moriya Interaction

requires structural inversion asymmetry + SOC

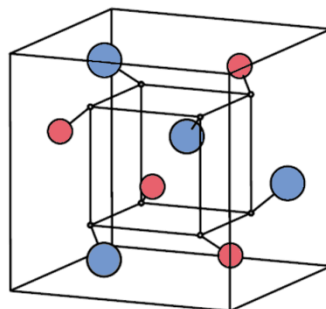


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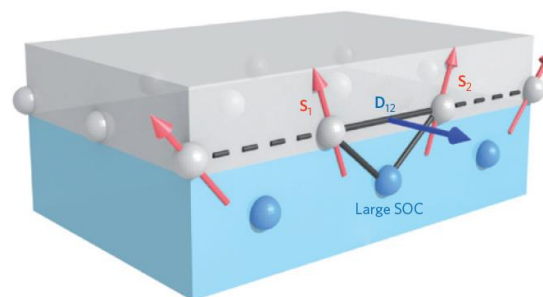
Chiral Interaction $E_{DM} = \mathbf{D} \cdot \mathbf{S}_1 \times \mathbf{S}_2$

bulk

crystal lacks inversion symmetry
e.g. B20 unit cell

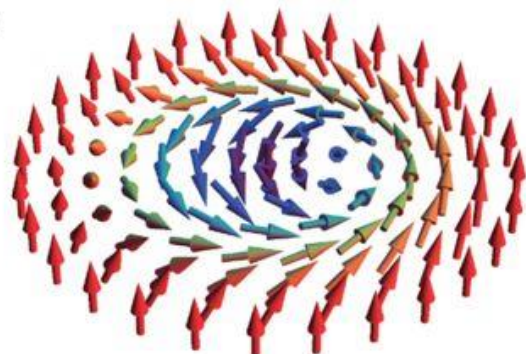


interface



asymmetric layers
(with different SOC)
around FM
break inversion symmetry

a



Bloch
skyrmion



Helicoid

b



Néel
skyrmion



Cycloid

Fert et al. Nature Nano. **8**, 152 (2013)

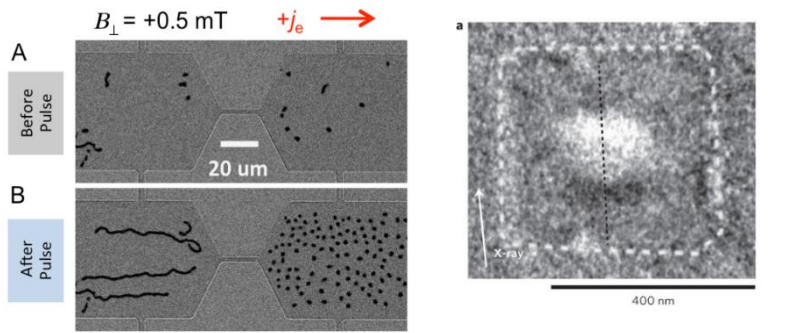
Al-Sharif, J. of Phys: Cond. Matter, **13**, 2807 (2001)

Room temperature magnetic skyrmions



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TRILAYER



$D_S = 700\text{-}2000 \text{ nm}$

Ta/CoFeB/TaO

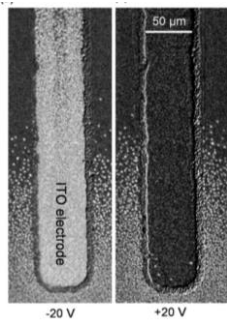
Jiang et al. Science (2015)

$D_S = 150 \text{ nm}$

Pt/Co/MgO

Boulle et al.

Nature Nanotech. (2016)

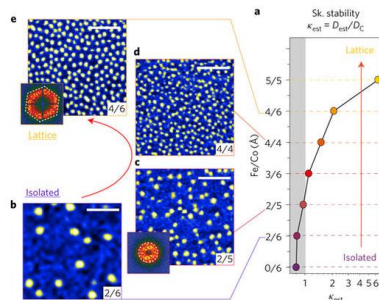


$D_S > 1000 \text{ nm}$

Pt/Co/AIO

Schott et al.

Nano Letters (2017)



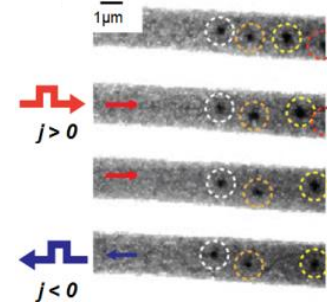
$D_S = 40\text{-}100 \text{ nm}$

Ir/Fe/Co/Pt

Soumyanarayanan et al.

Nature Materials (2017)

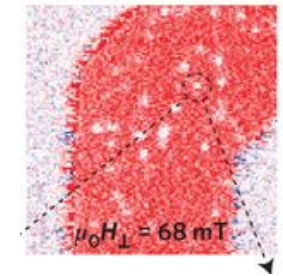
MULTILAYER



$D_S = 100\text{-}300 \text{ nm}$

[Pt/Co/Ta] $\times 15$

Woo et al. Nature Materials (2016)

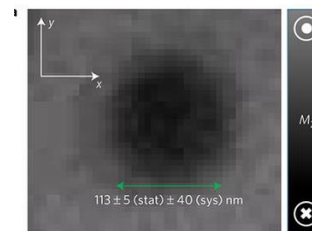


$D_S = 30\text{-}90 \text{ nm}$

[Pt/Co/Ir] $\times 10$

Moreau-Luchaire et al.

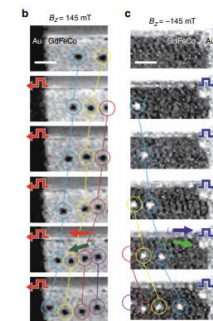
Nature Nanotech. (2016)



$D_S = 70 \text{ nm}\text{-}150 \text{ nm}$

[Pt/CoFeB/MgO] $\times 15$

Litzius et al. Nature Phys (2017)



$D_S = 180 \text{ nm}$

[Pt/GdFeCo/MgO] $\times 20$

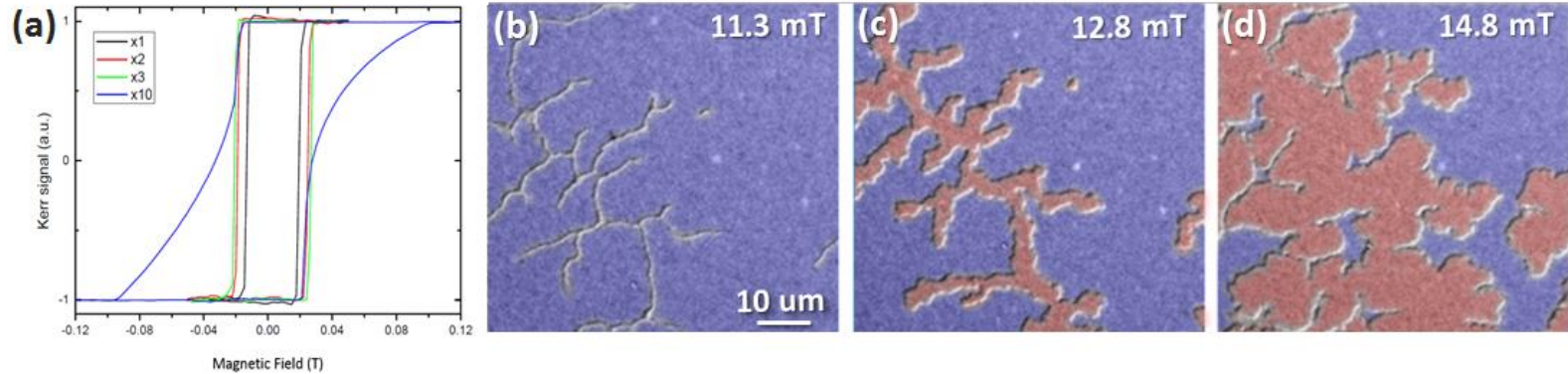
Woo et al. Nature Comm (2018)

LTEM under tilt: {Pt/Co/Ir} \times 2

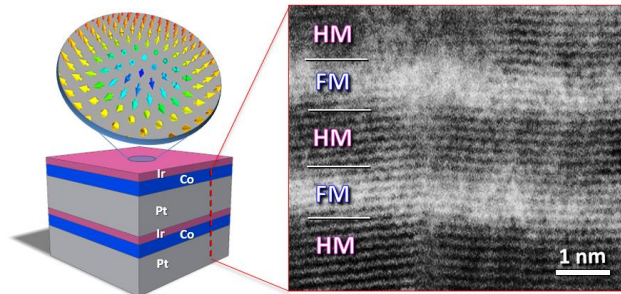


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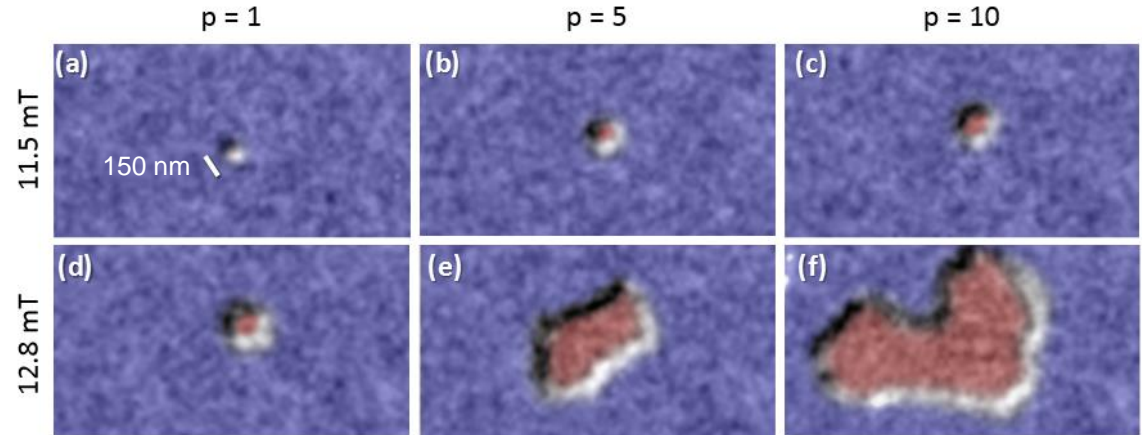
Chiral DWs and skyrmion bubbles



Multilayer stack



Bursting skyrmion bubbles, p = number of field pulses



p = number of field pulses

Magnetic imaging using STXM

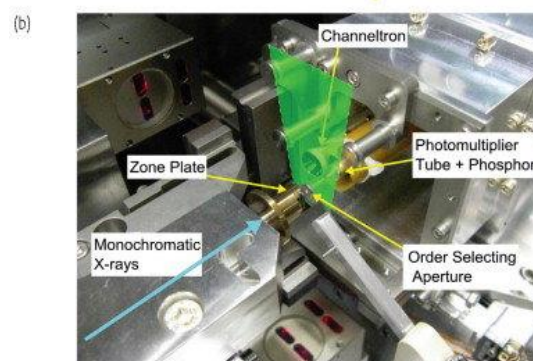
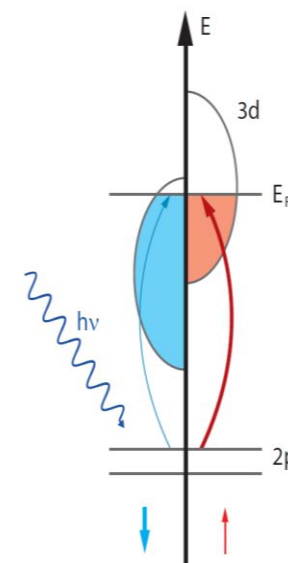
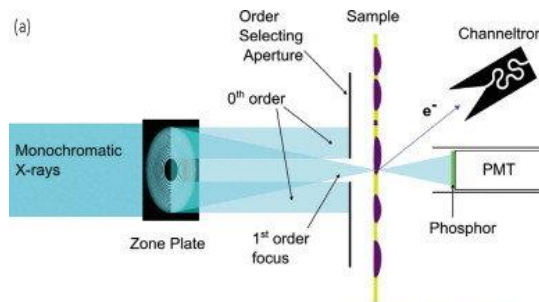


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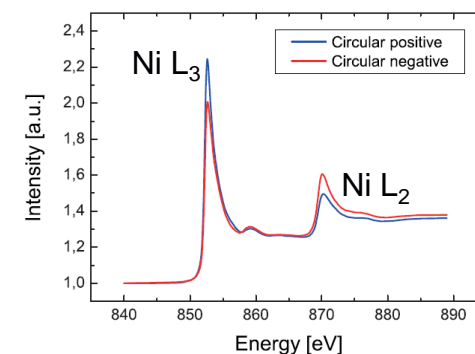
PolLux at Swiss Light Source, PSI



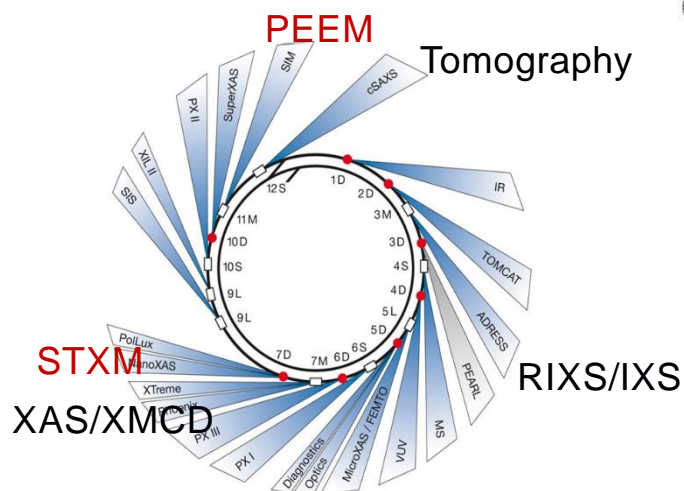
Zone plate optics



Watts and Ade, Materials Today (2012)



XMCD contrast

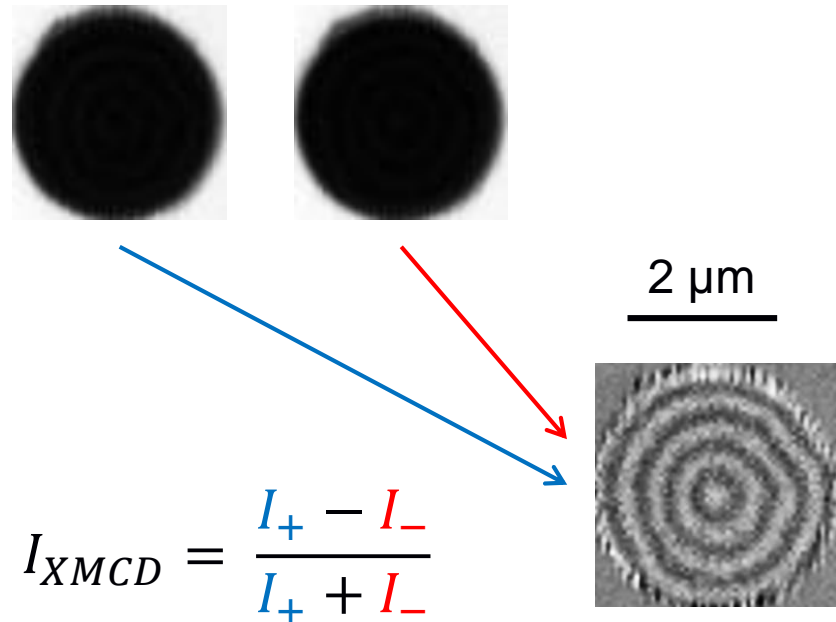
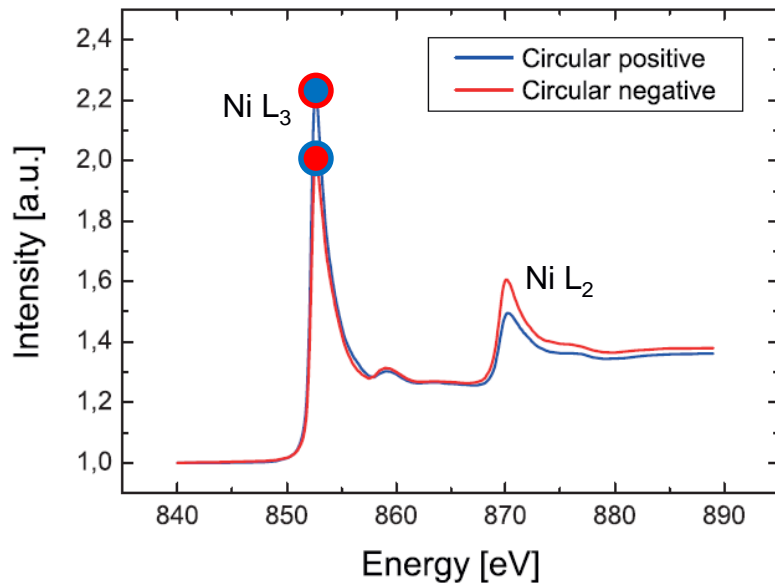


X-ray magnetic circular dichroism (XMCD)



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- Magnetic images can be acquired by employing the x-ray magnetic circular dichroism effect
- Different absorption of circularly-polarized x-rays depending on the local magnetization of the sample



Magnetic imaging of nanodiscs by STXM

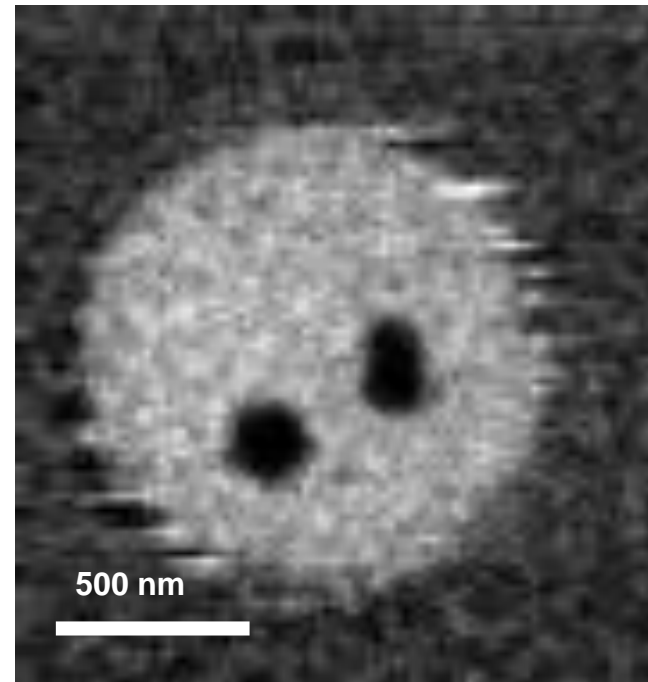


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Ta(4.6 nm)/Pt(7.5 nm)/[Co(0.7 nm)/Ir(0.5 nm)/Pt(2.3 nm)]_{x10}/Ir(1.3 nm)

STXM AT PoLux/PSI

- Co L₃ edge (779.5 eV)
- Room temperature
- Field applied out-of-plane
- X-ray absorption proportional to out-of-plane magnetization M_z
- Dark and bright contrast represents oppositely out-of-plane magnetized regions
- Spatial resolution ~25 nm
- Temporal resolution ~200 ps
- Skyrmion bubbles are observed

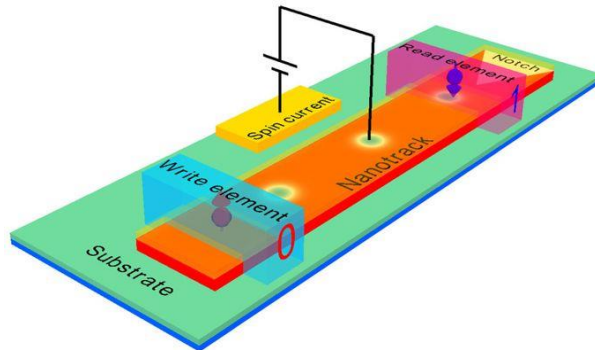


Skyrmions for future computing



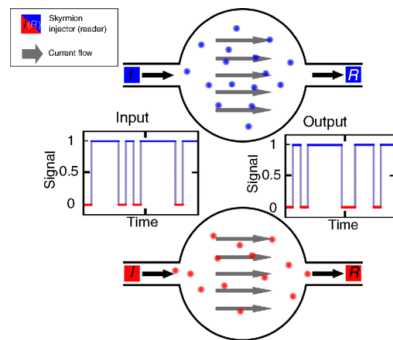
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Digital racetrack memory



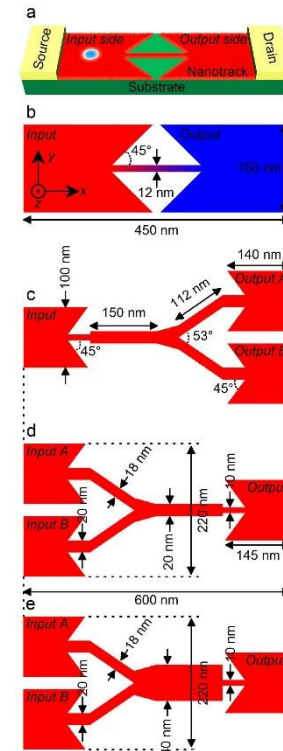
X. Zhang et al., *Sci. Rep.* **5**, 7643 (2015)

Neuromorphic computing



D. Pinna et al., *Phys. Rev. Applied* **9**, 064018 (2018)

Boolean logic



Fanout

OR

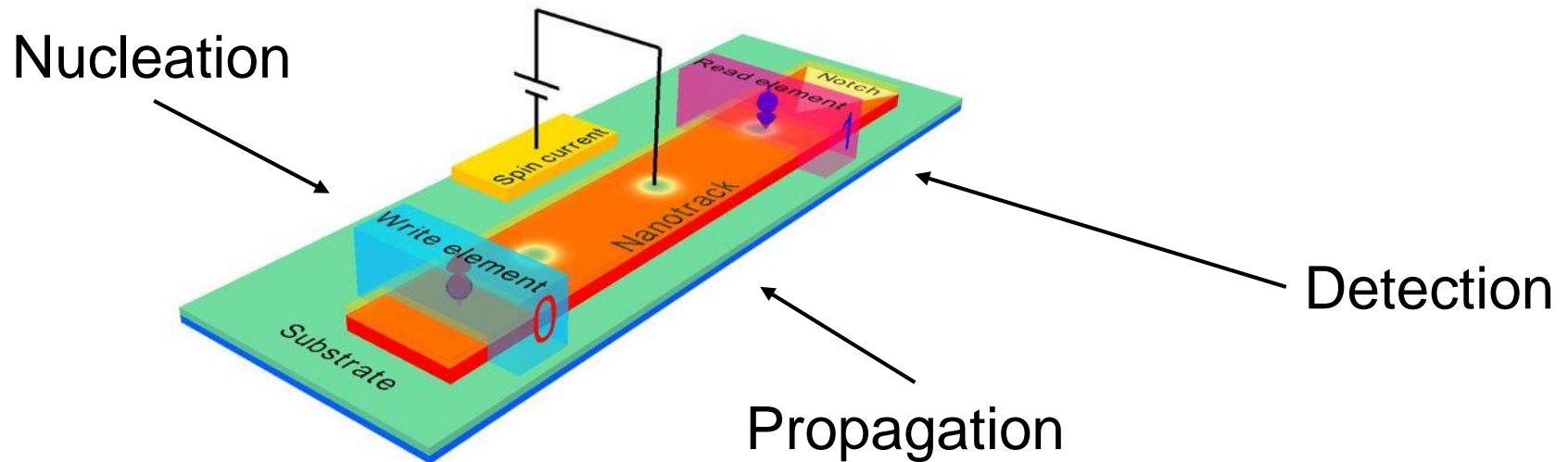
AND

X. Zhang et al., *Sci. Rep.* **5**, 9400 (2015)

Basic skyrmion operations



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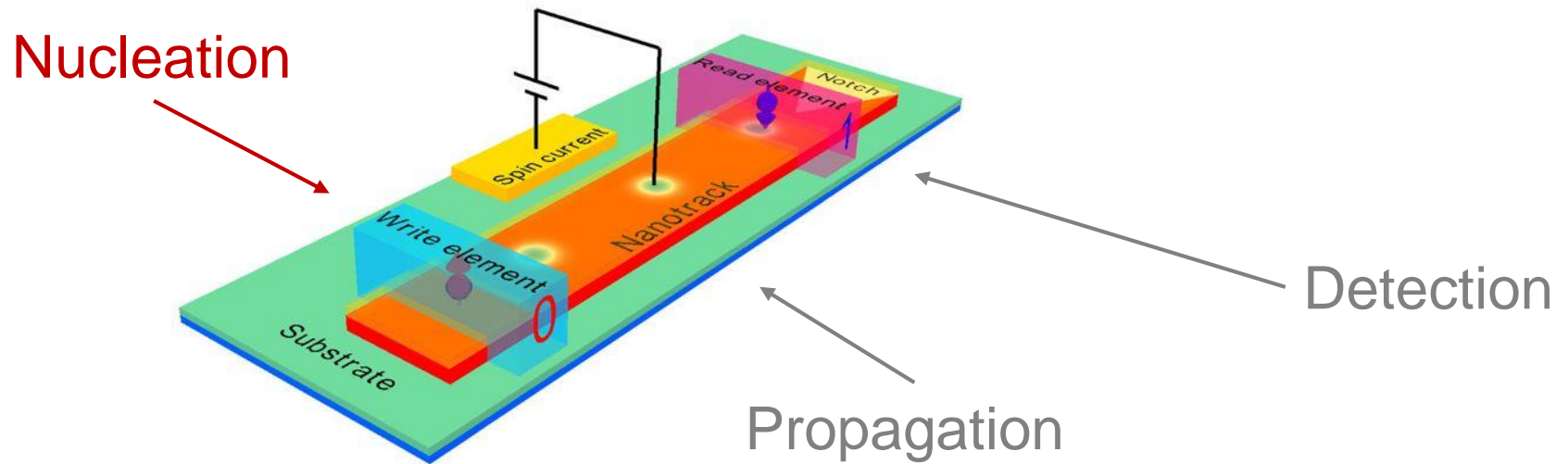


All need to be accomplished by electrical means for spintronics

Basic skyrmion operations



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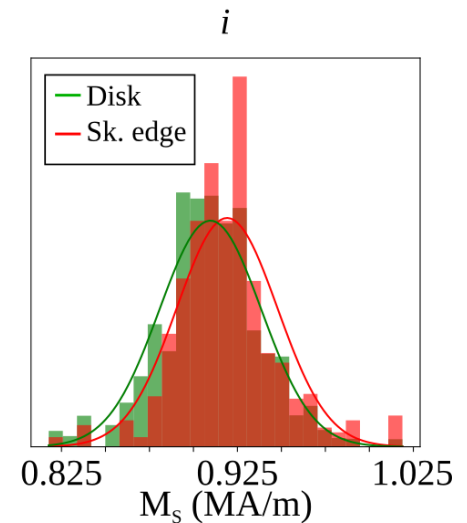
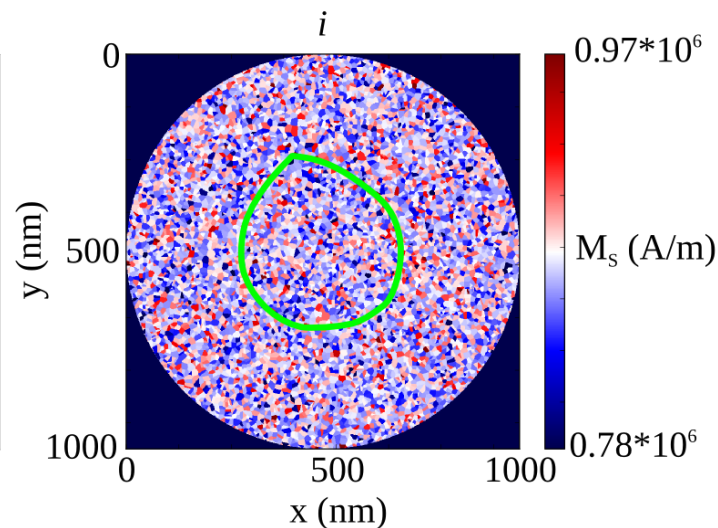
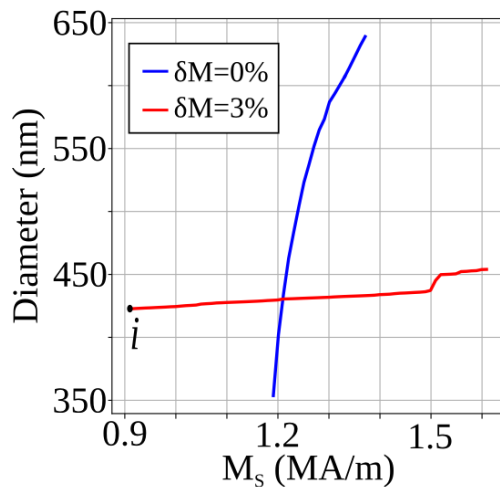
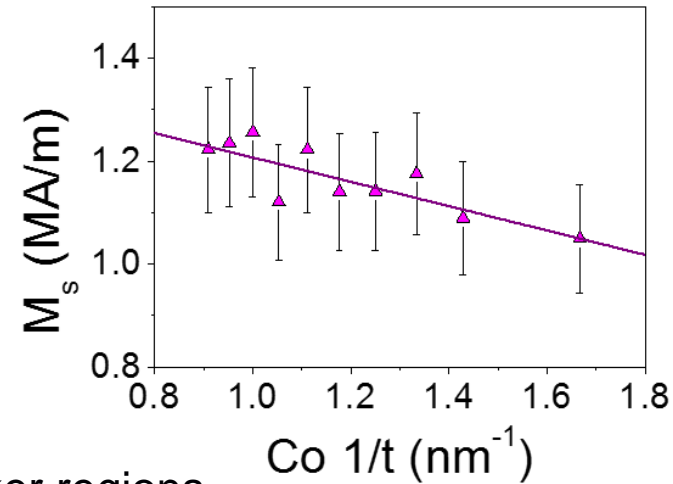
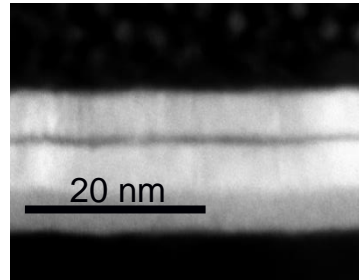
Disordered Simulation



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DISORDER

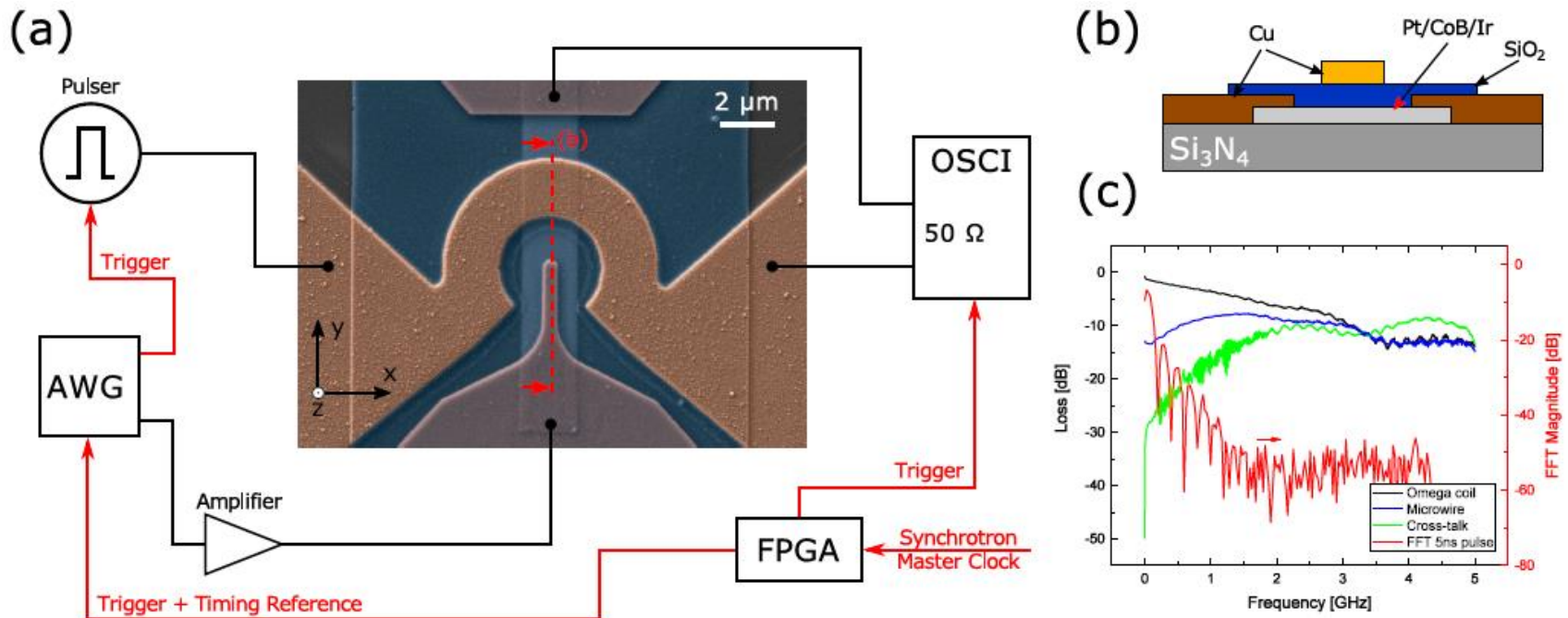
- Modulate thickness
- 10 nm grain size
- Average M_S with standard deviation of $\delta M/M$, 3%
- Stabilizes skyrmion bubble
- Domain boundary found in regions of high M_S i.e. thicker regions





Electrical nucleation in Pt/CoB/Ir

- Fabrication of point-contact geometries to facilitate the injection of magnetic skyrmions
- Design combined with Omega coils for investigation of gyration dynamics and for providing transient fields during nucleation/motion (5 step lithography process)

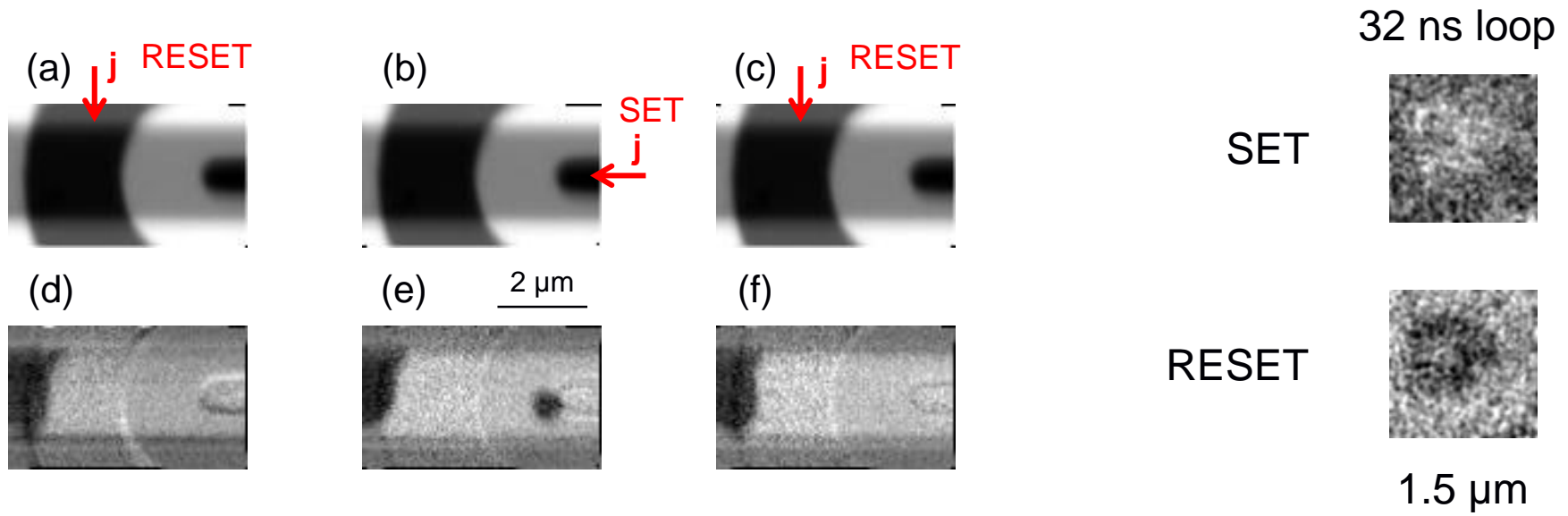


Skyrmion nucleation in Pt/CoB/Ir



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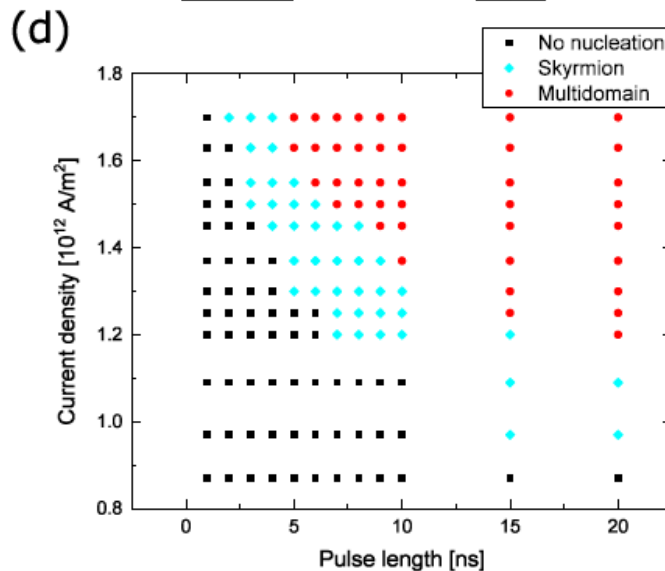
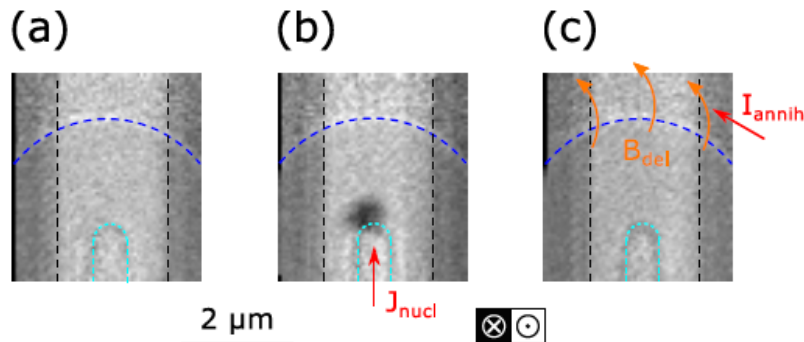
- Nucleation of magnetic skyrmions from the point contact (current densities $\sim 1-1.5 \times 10^{12}$ A/m²) - **SET**
- Annihilation of the skyrmion by pulsed magnetic field (generated by the omega coil) - **RESET**
- XMCD-STXM imaging at a time resolution of 2 ns
- All measurements were done at the remnant state (0 mT), normalised to average contrast throughout movie



Skyrmion nucleation in Pt/CoB/Ir



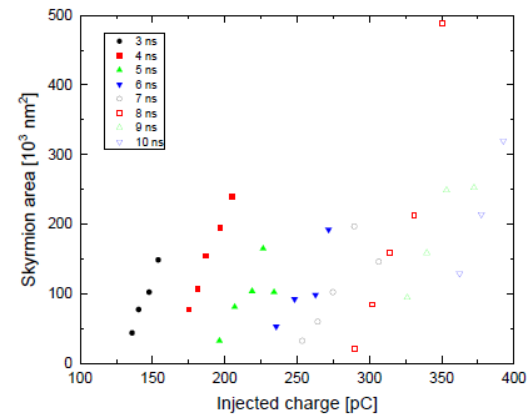
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Phase diagram at zero field

Broad range of values where single skyrmion is reliably nucleated by a single current pulse

Skyrmion size determined by the charge contained in the pulse

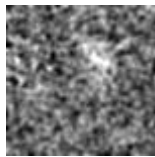


Skyrmion nucleation dynamics

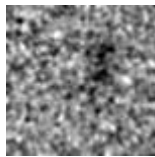


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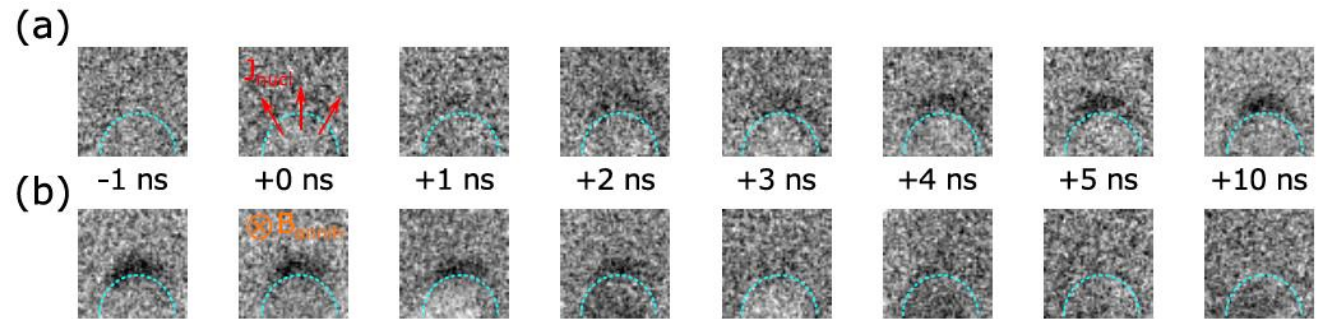
SET
30 ns loop



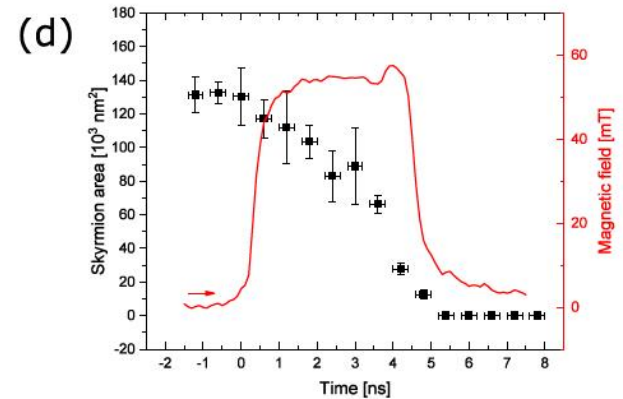
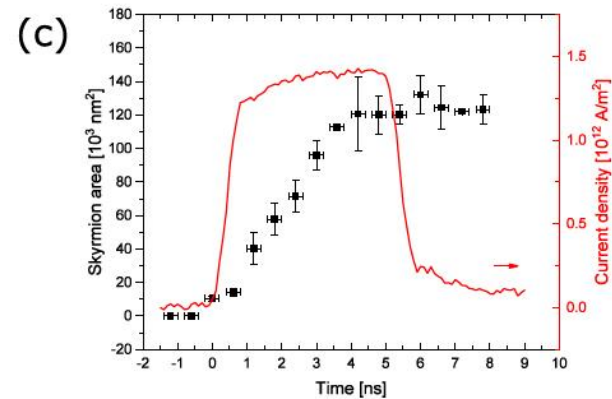
RESET
20 ns loop



2 μm



500 nm



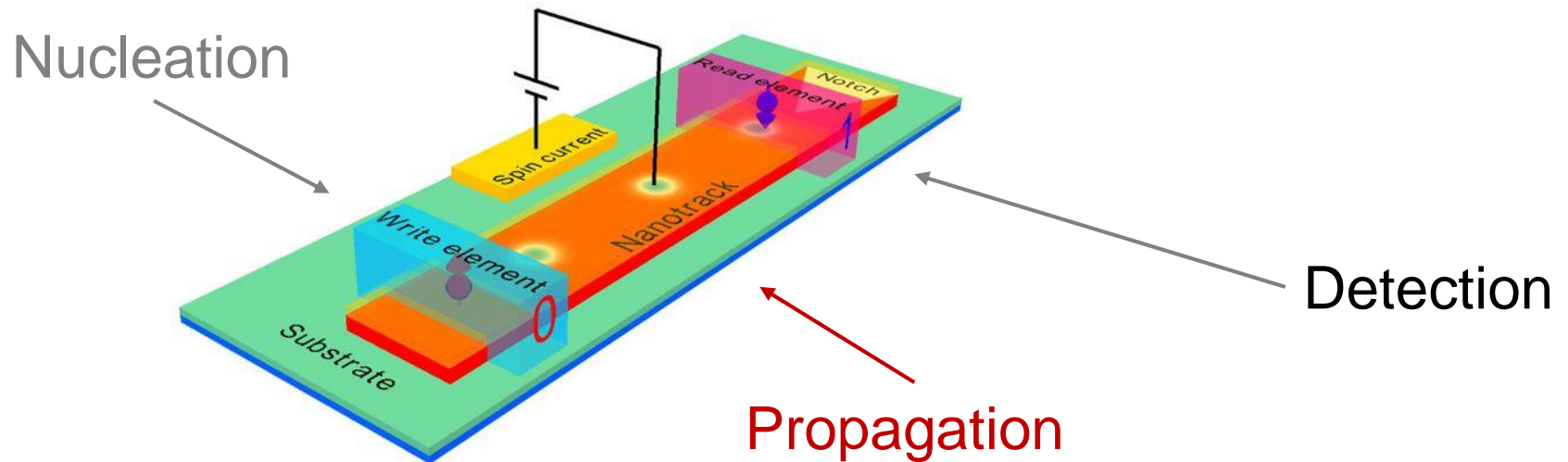
Improved temporal resolution to 200 ps

No incubation time for skyrmion nucleation or deletion

Basic skyrmion operations



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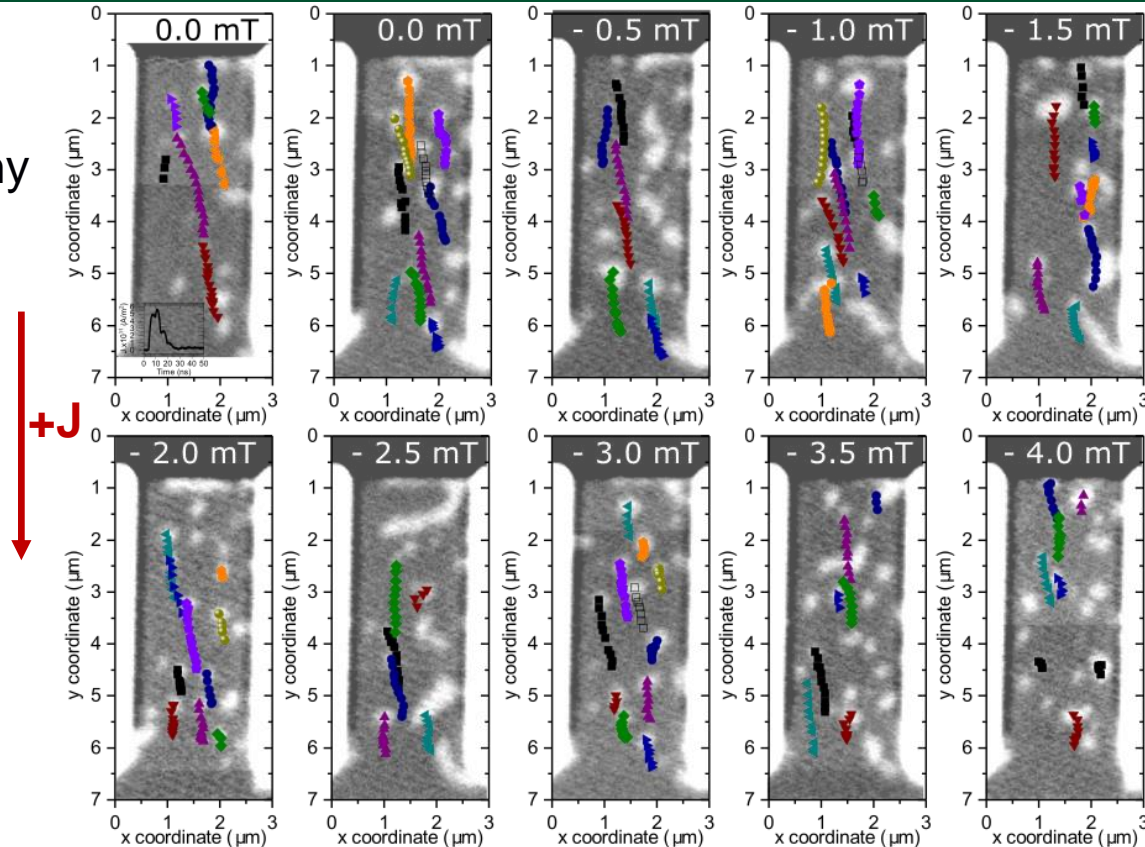


Skyrmion Motion



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- $[\text{Co}_{68}\text{B}_{32}$ (0.8 nm)/Ir (0.4 nm)/Pt (0.6 nm)]x5
- Lift off electron beam lithography
- 200 nm Silicon nitride membranes
- 3 μm wide wire
- Skyrmions nucleated with 9 ns current pulse in the mid peak 10^{12} A/m² range
- 9 ns pulse, peak J 5.6×10^{11} A/m²
- Two pulses between images
- Centre of skyrmion tracked with linear assignment problem algorithm, ImageJ Trackmate

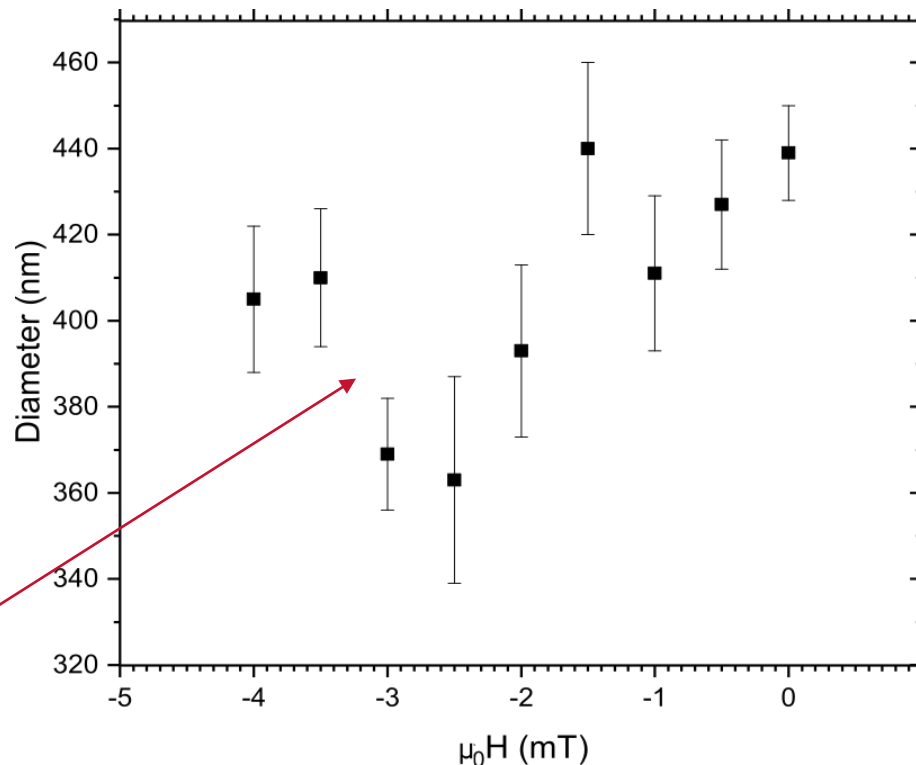


Skyrmion Motion



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- Average diameter at each field calculated by counting pixels of each skyrmion and converting area into a diameter assuming a cylindrical skyrmion
- Average diameter decreases with increasing field
- Access to a larger skyrmion size range



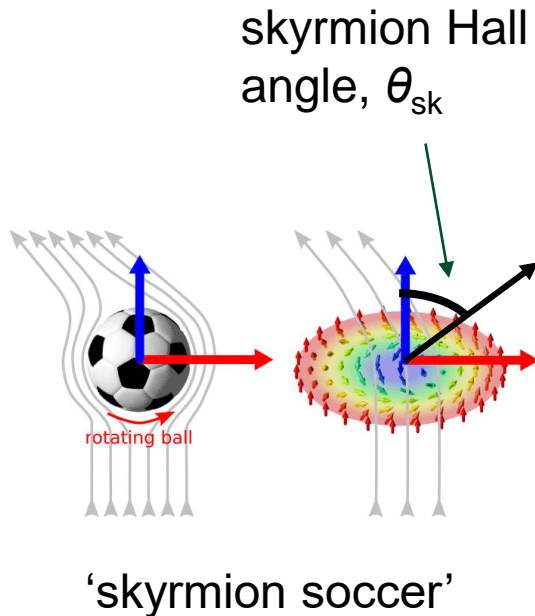
smaller
skyrmions
annihilated

Skyrmion Motion

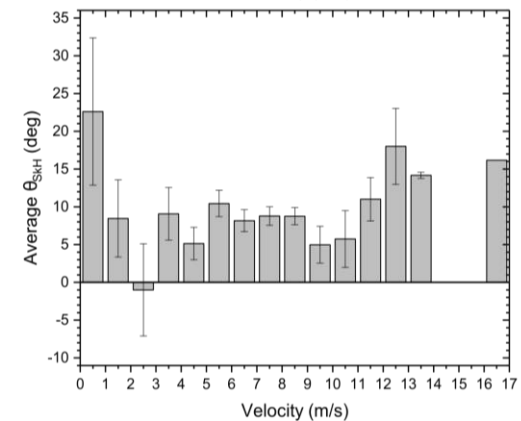
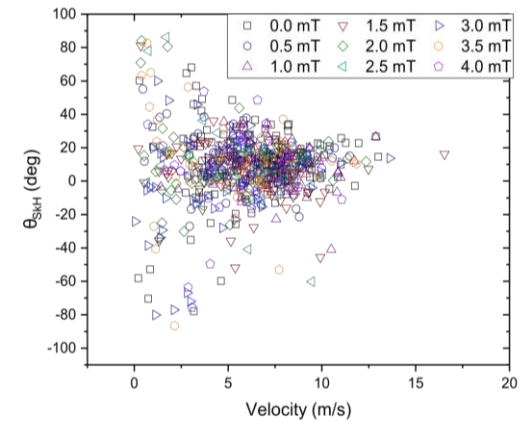
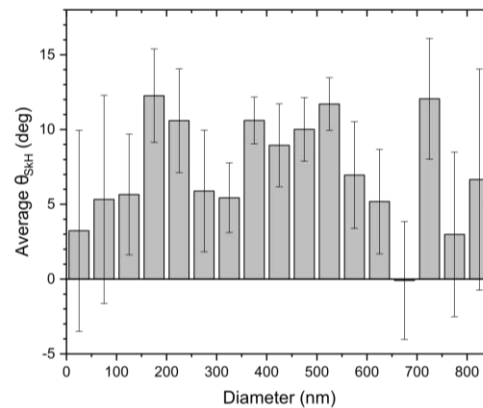
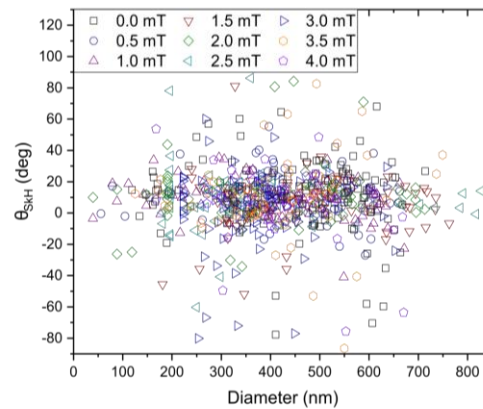


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- Each individual motion event
- Skyrmion Hall angle scatter decreases with increasing velocity
- Skyrmion Hall angle scattered around $9 \pm 2^\circ$



Everschor-Sitte, JAP **115**,
172602 (2014) (2013)



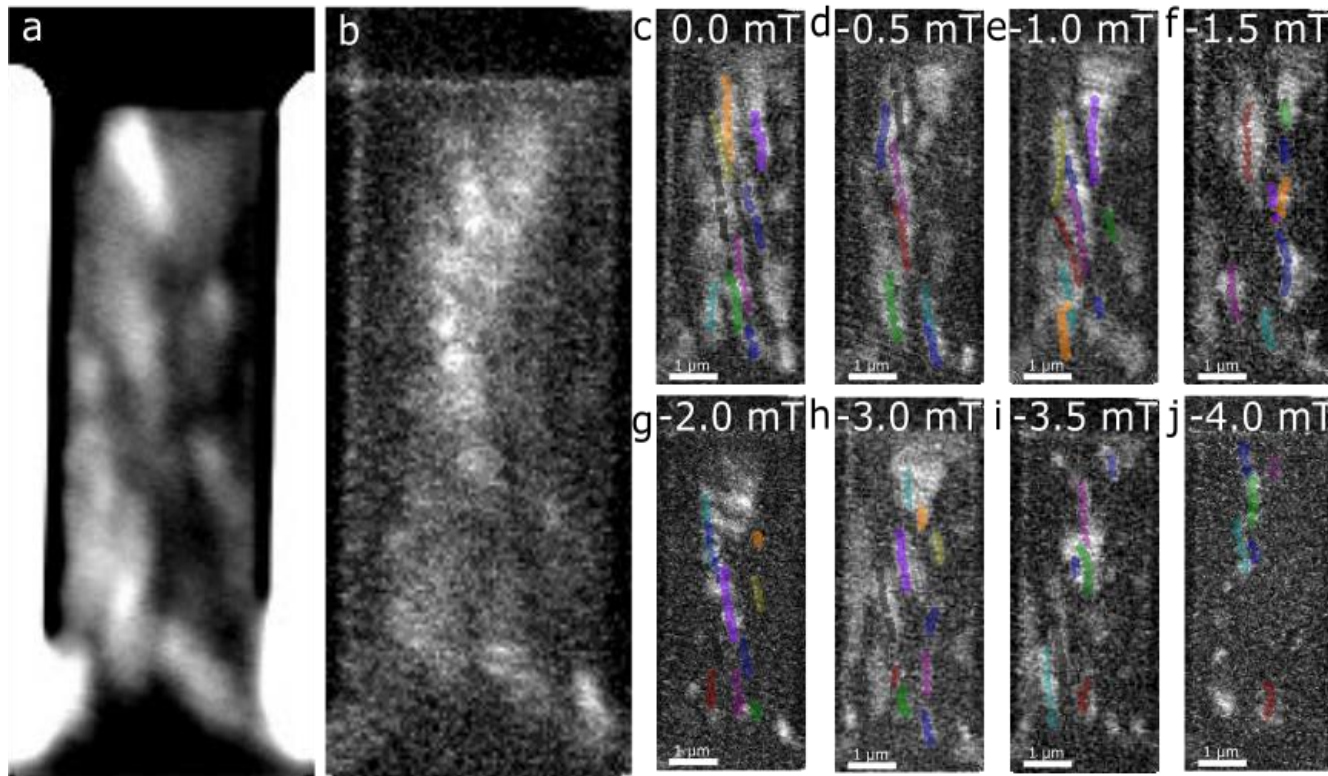
Zeissler et al., Nature Commun. **11**, 428 (2020)

Skyrmion Motion



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- Average of all images (a) skyrmion positions (b) skyrmion motion.
- Shows skyrmion prefer certain sites and to move along certain tracks – local environment dominates skyrmion motion

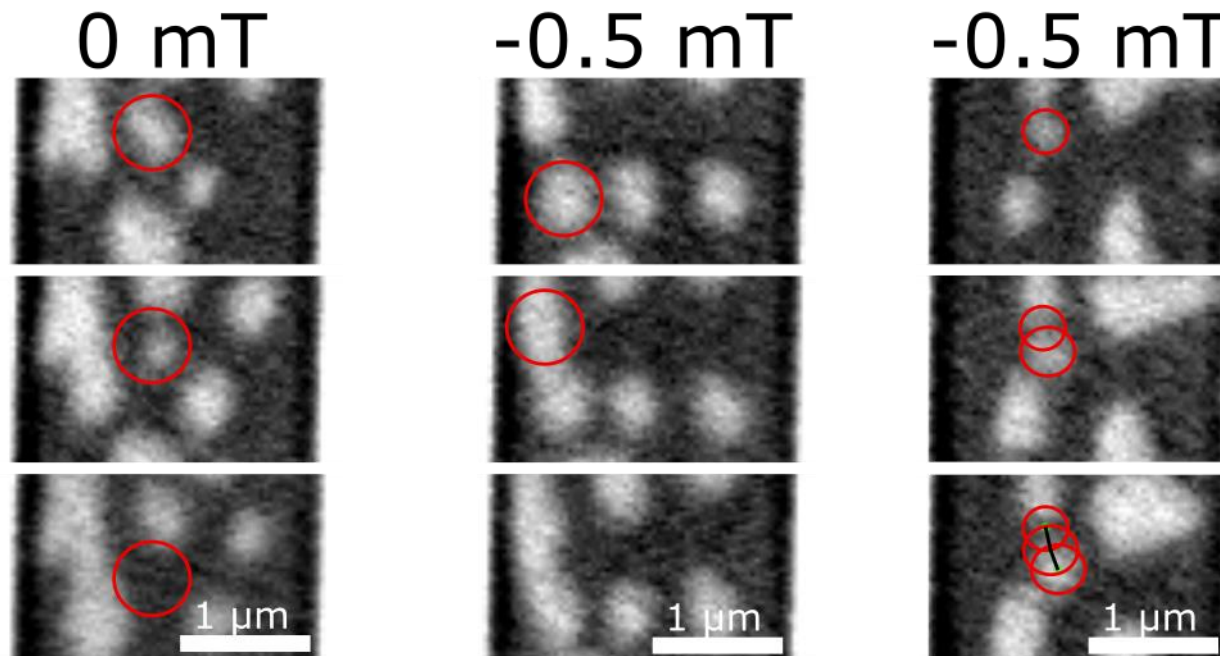


Other skyrmions matter



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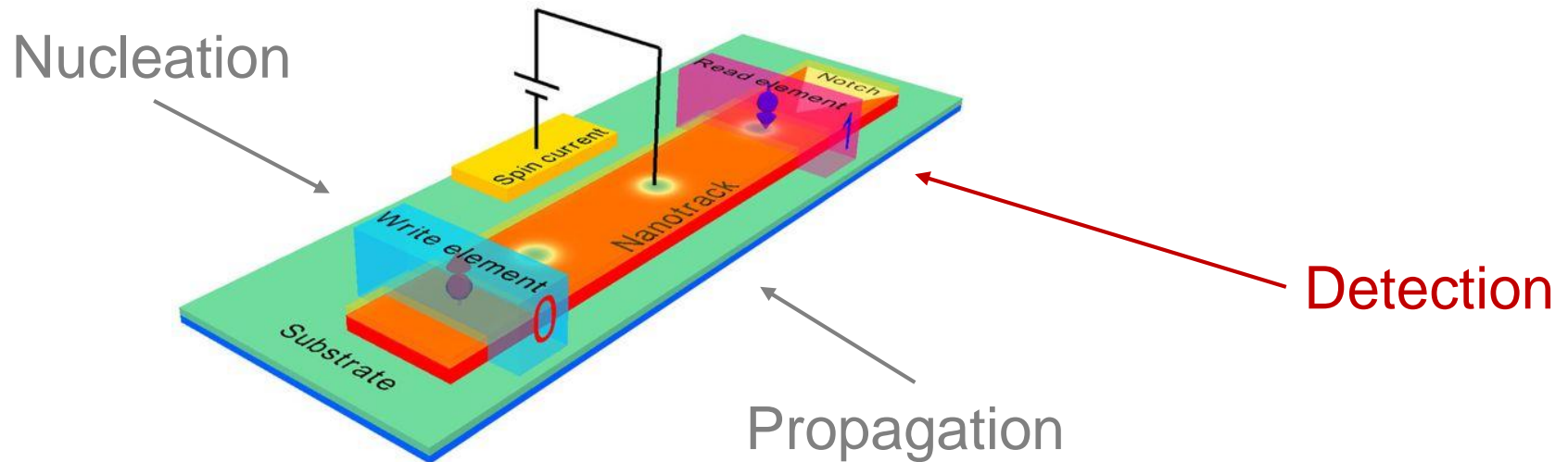
- Skyrmion annihilation due to pinning
- Skyrmion fusion
- Skyrmion-skyrmion deflection – leads to a change in skyrmion Hall angle



Basic skyrmion operations



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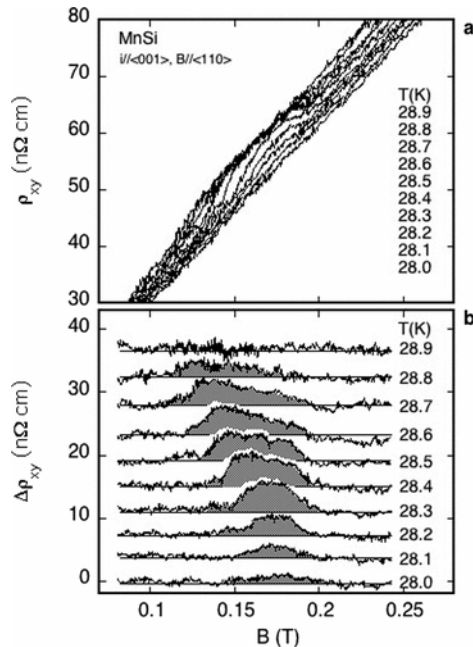


Topological Hall effect in B20 materials

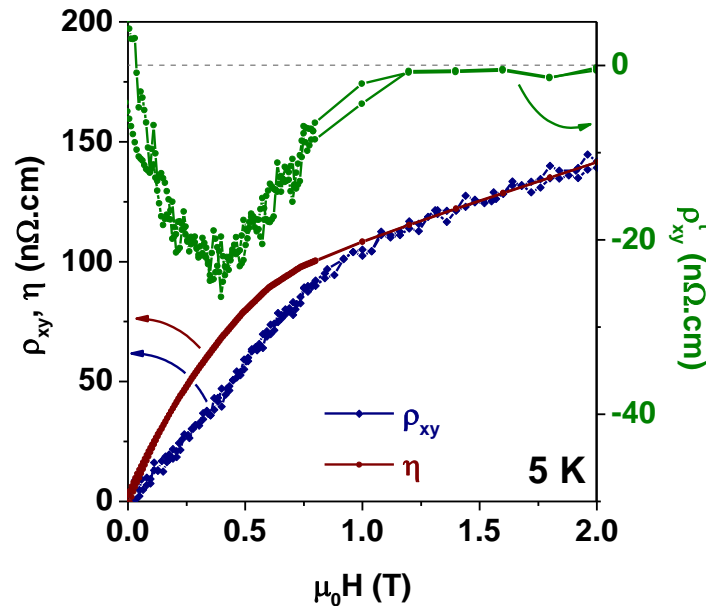


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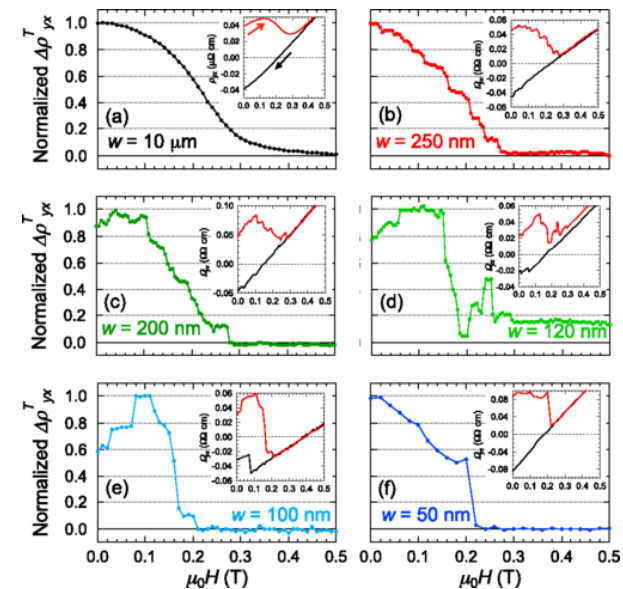
MnSi crystal



FeGe thin film



FeGe nanoscale Hall bar



Neubauer et al., PRL (2013)

Porter et al. PRB (2014)

Kanazawa et al. PRB (2015)

Widely accepted as a signature for the presence of Bloch skyrmions in B20s

$$\rho_{xy}(H) = \rho_{xy}^o(H) + \rho_{xy}^a(H) + \rho_{xy}^t(H)$$

$$\rho_{xy}^t = PR_0 B_{\text{eff}}$$

Anomalous and topological Hall effect measurements in thin films

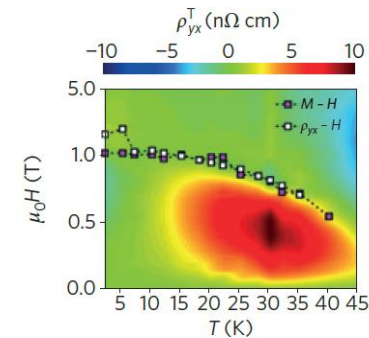
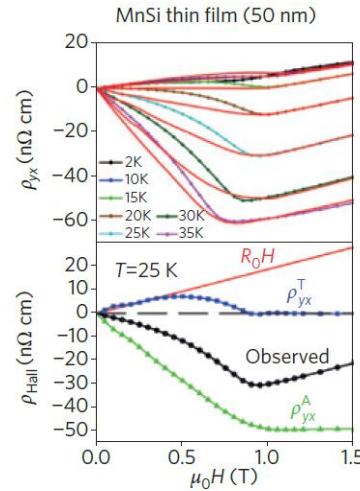


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BULK DMI SYSTEMS

- B20 e.g. MnSi

$$\rho_{xy}^T \sim 4 - 10 \text{ n}\Omega \text{ cm}$$



YuFan Li et al., PRL 110, 117202 (2013)

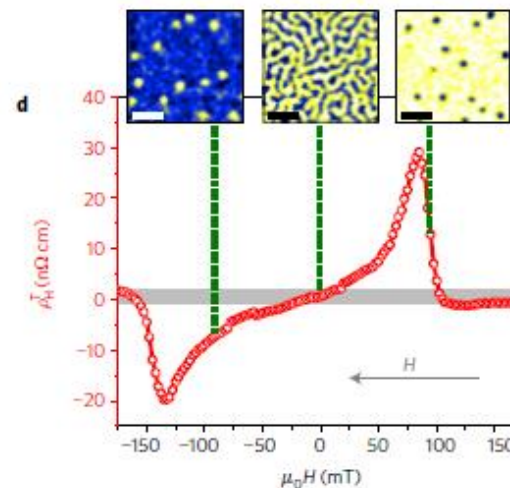
INTERFACIAL DMI SYSTEMS

- Ir/Fe/Co/Pt

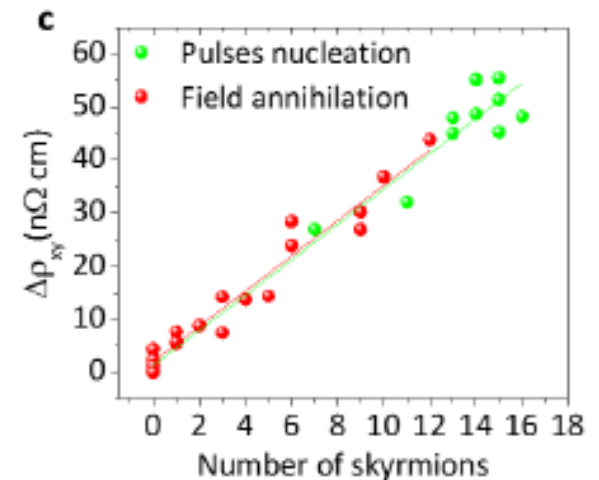
$$\rho_{xy}^T = 30 \text{ n}\Omega \text{ cm}$$

- $[\text{Pt}/\text{Co}/\text{Ir}]_{\text{x}20}$

$$\rho_{xy} = 3.5 \pm 0.5 \text{ n}\Omega \text{ cm}$$



A. Soumyanarayanan
Nature Materials 16 898 (2017)



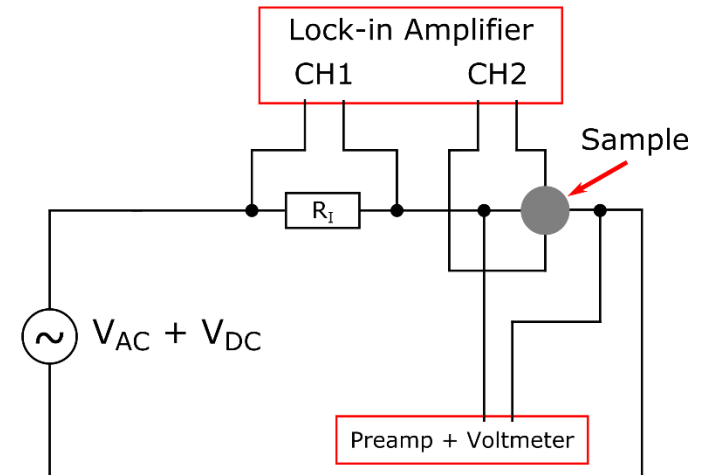
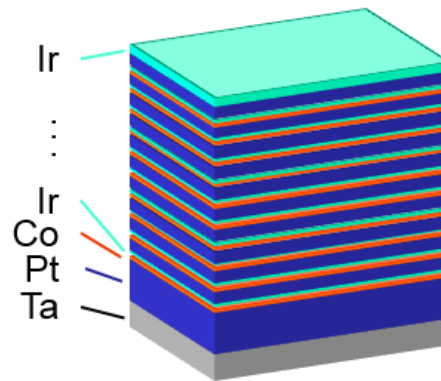
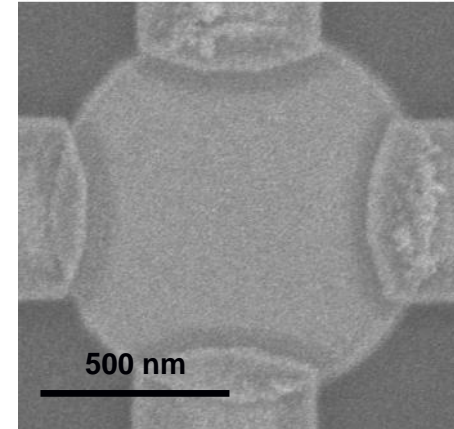
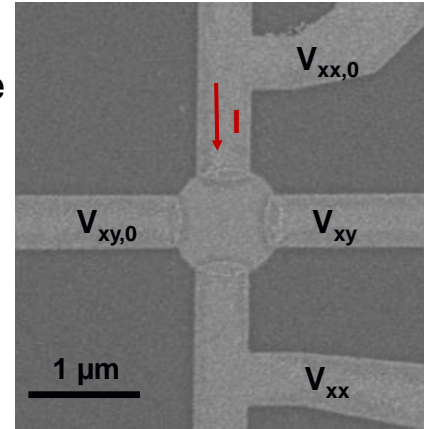
D. Maccariello et al. Nature Nanotech. (2018)

In situ electrical measurement set-up



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- Ta(3.5 nm)/Pt(3.8 nm)/[Co(0.5 nm)/Ir(0.5 nm)/Pt(1.0 nm)]_{x10} Pt(3.2 nm)
- DC sputtered on 200 nm silicon nitride membrane
- Top down e-beam lithography
- 1000 nm diameter
- Quasi 4 point measurements
- 10 k Ω series resistor
- 17 Hz AC voltage
- DC offset voltage



Field-driven reversal & skyrmion nucleation

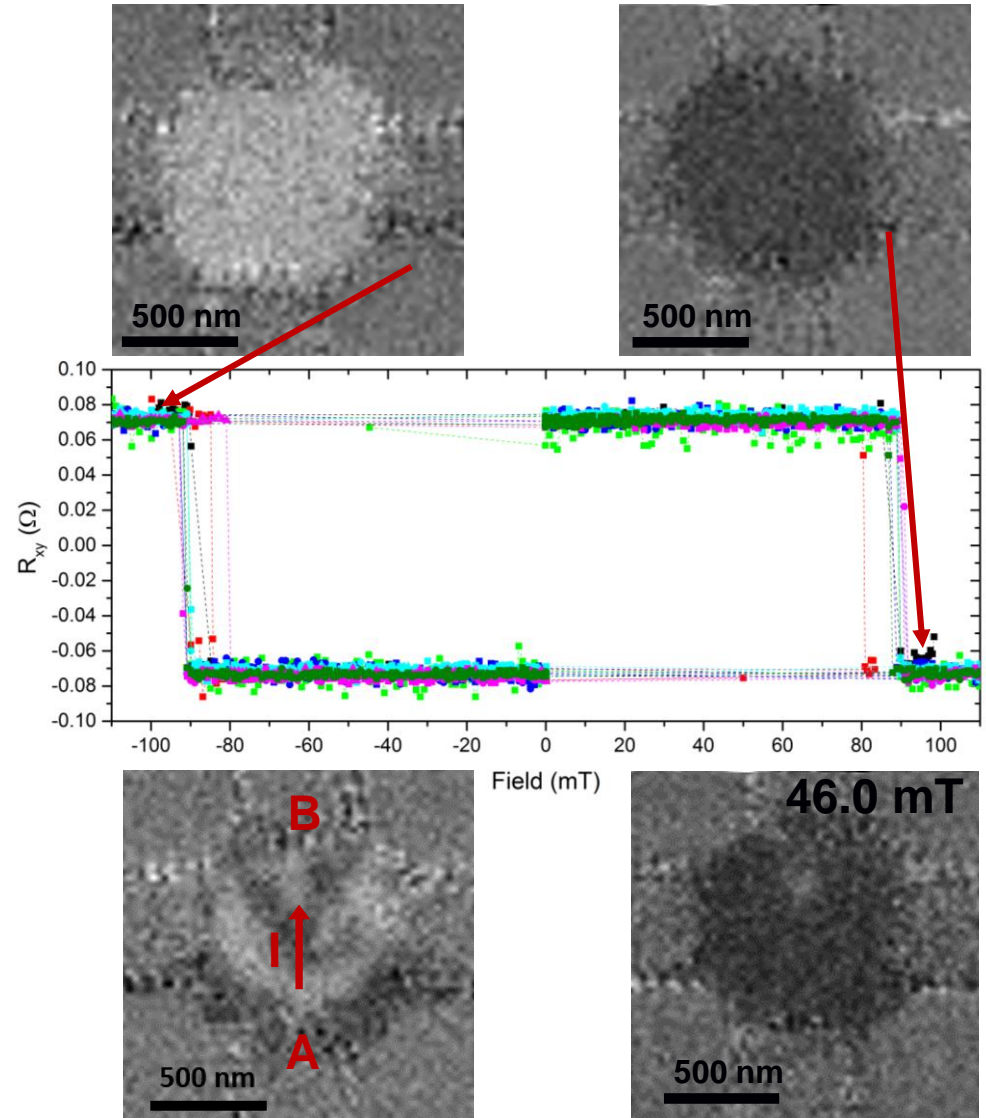


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- Out-of-plane field applied
- XMCD contrast
BLACK = +90 mT
WHITE = -90 mT
- Clear Hall signal

Skyrmion nucleation protocol

- +100 mT saturation
- Return to 0 mT
- 5x 5 ns 0.5 V pulses separated by 200 ns
 $J \sim 7 \times 10^{11} \text{ A/m}^2$
- Apply positive field to collapse domains
- Last object left before saturation is a skyrmion

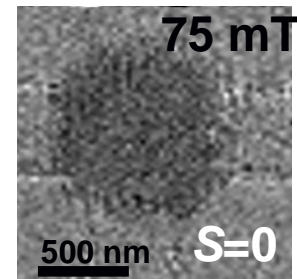
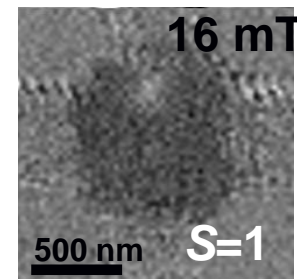
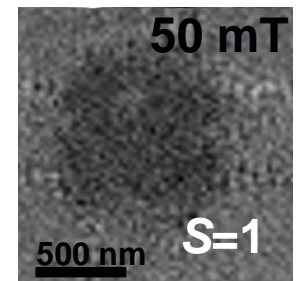
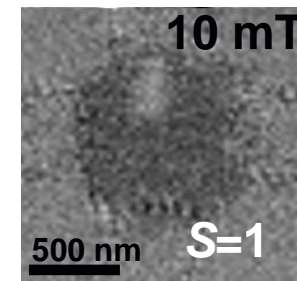
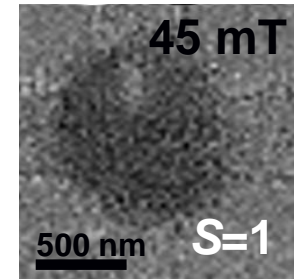
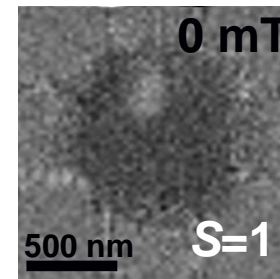
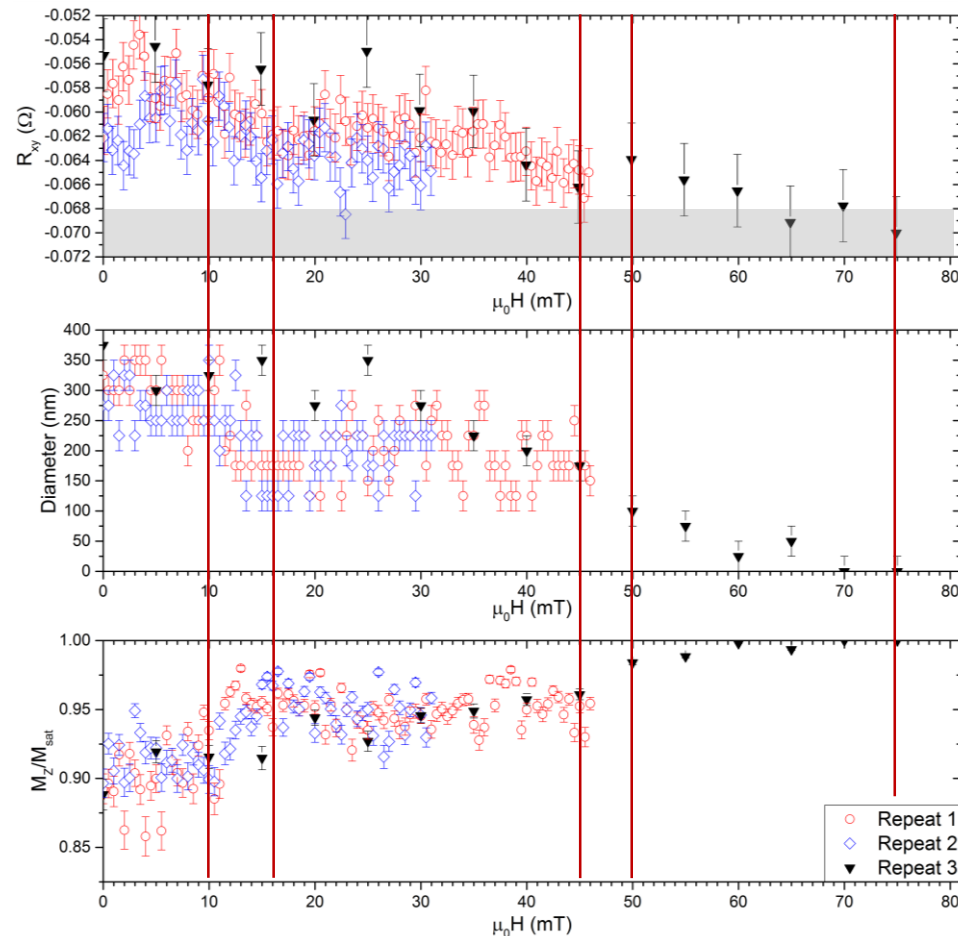


Skyrmion: shrinking and annihilation



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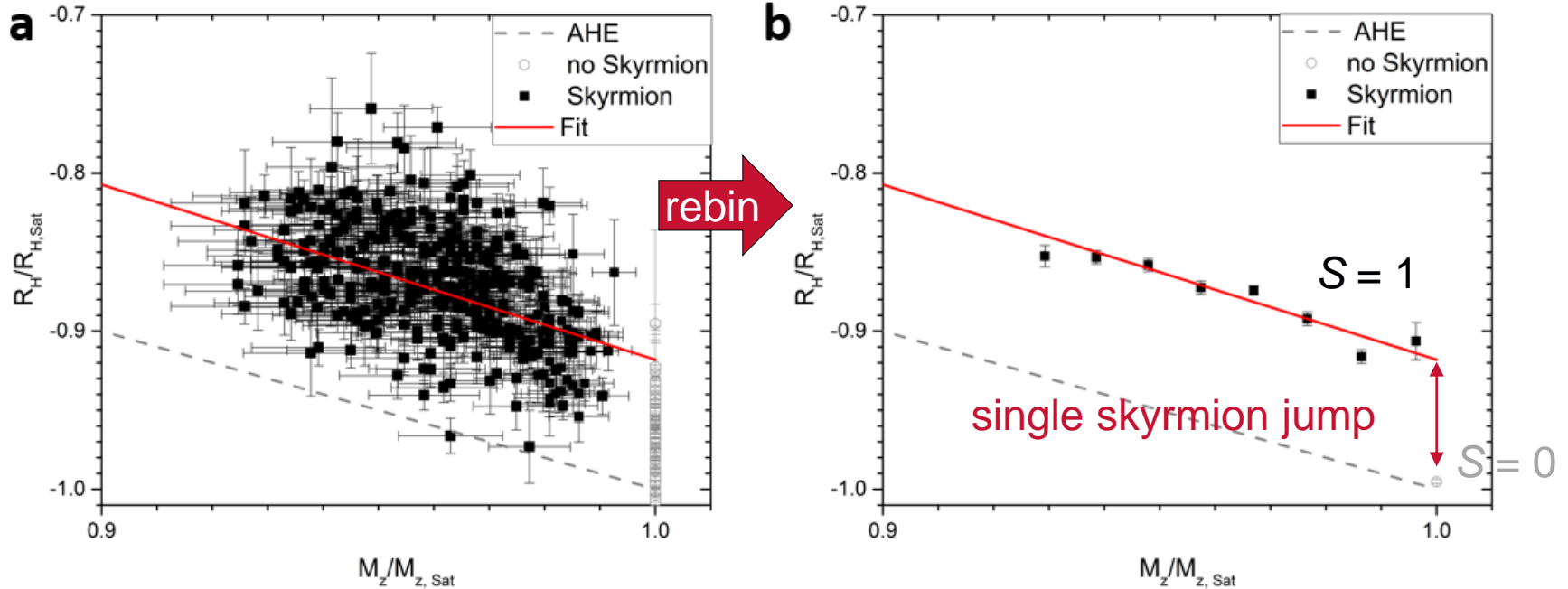
- Return to 0 mT - 300 nm Skyrmions
- Increase field – shrinks skyrmion ; Resistance follows diameter
- 24 % change in the Hall resistance when skyrmion is present



Single skyrmion Hall signal



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Fitted slope = 1.1 ± 0.1

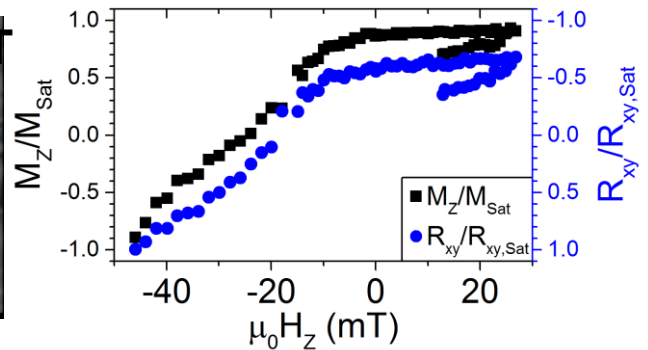
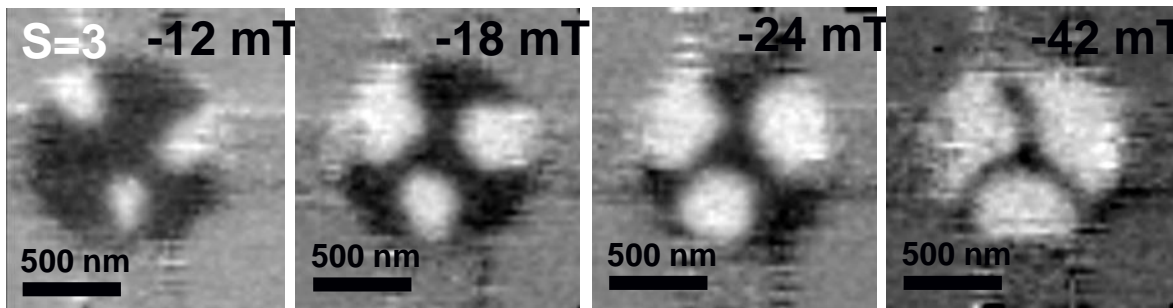
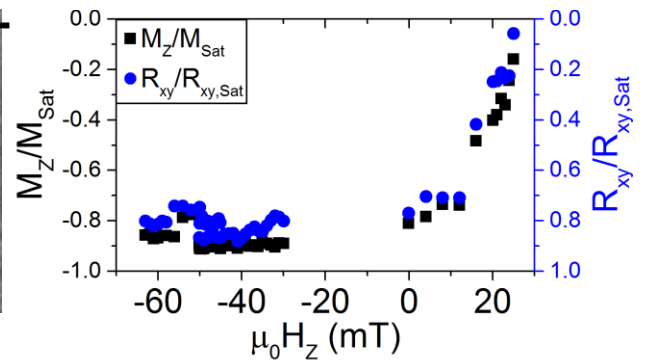
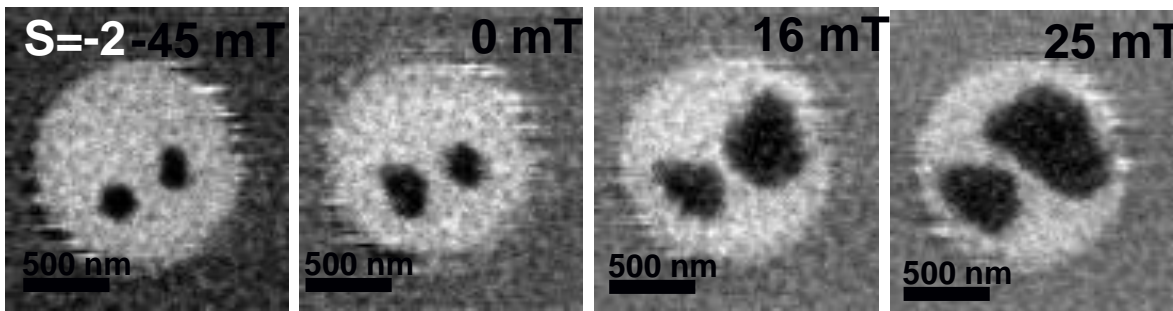
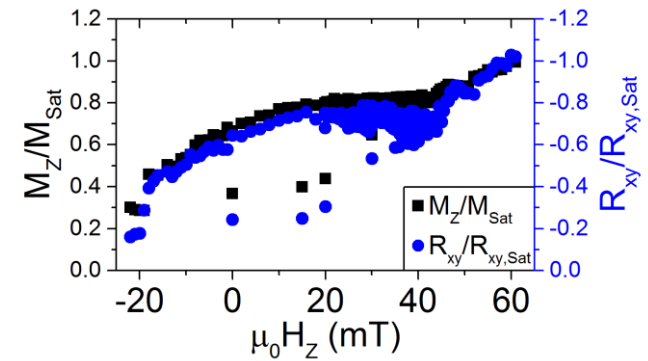
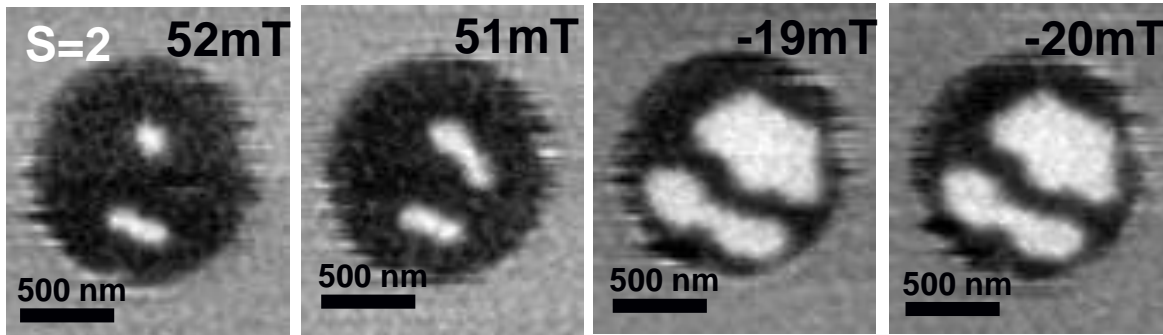
AHE ✓

Fitted jump = 8.2 ± 0.5 % of $R_{H,Sat} \sim 11 \pm 1$ n Ω cm

Multiple skyrmion states



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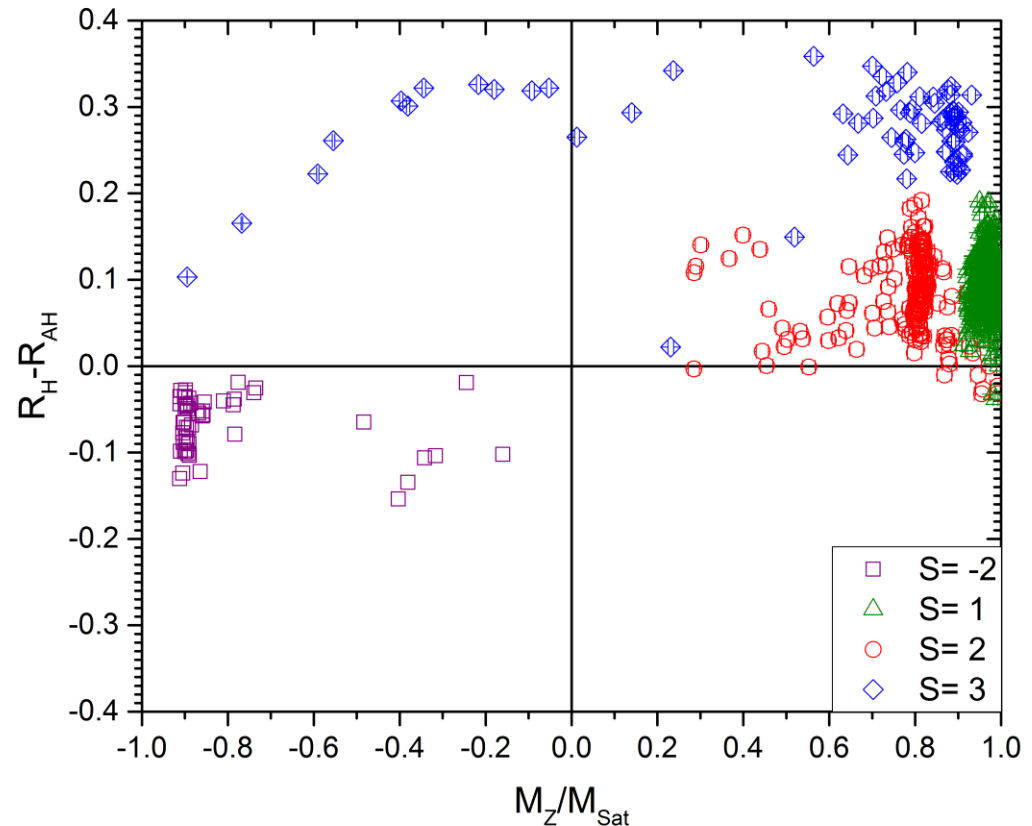


Hall resistance contributions



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- Normalise Hall resistance and magnetisation to saturated values
- $R_{xy} = R_0 B + \mu_0 R_S M + R_{int}$
- Measured $R_0 = -(1.9 \pm 0.2) \times 10^{-11} \Omega \text{m/T}$
- $R_H = R_{xy} - R_0 B = \mu_0 R_S M + R_{int}$
- Plot R_H vs M_Z
- Intercept at R_{int}
- Plot $R_H - R_{AH}$ vs M_Z to see R_{int}
- This is the signal that remains once ordinary and anomalous effects are accounted for.

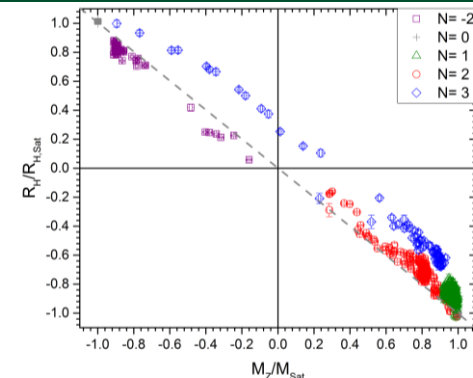


Topological Hall resistivity?



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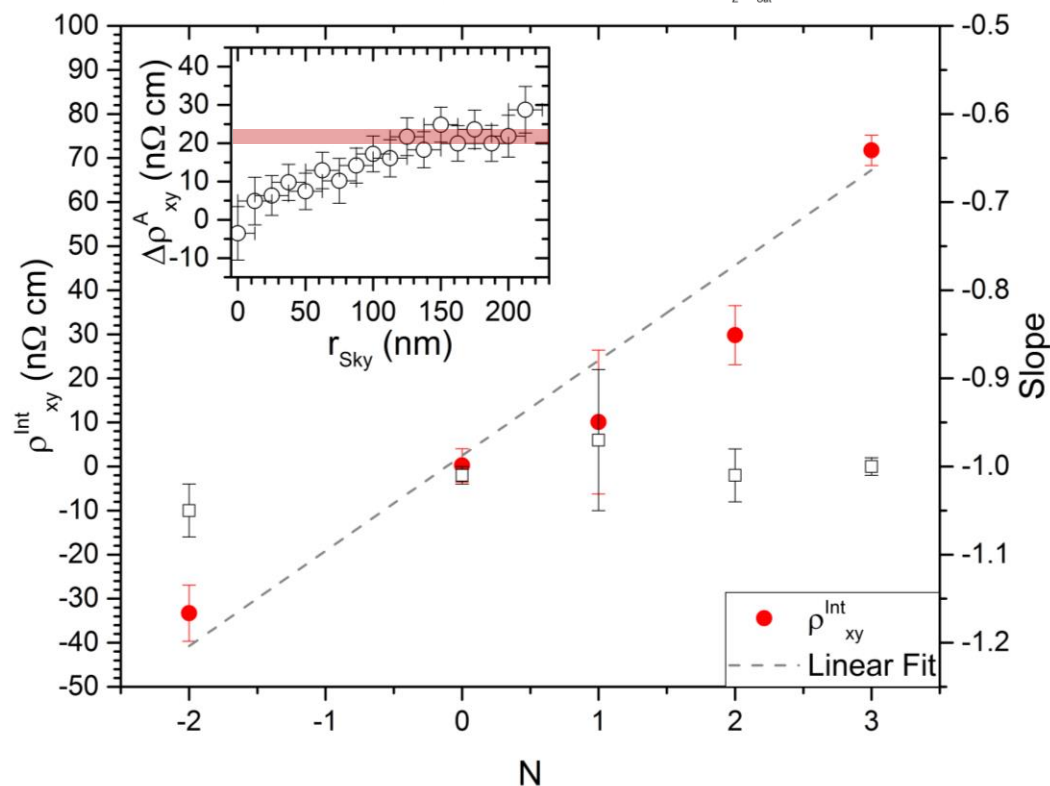
- Linear fit of normalised R_H vs M_z for different N
- Anomalous Hall contribution:
Slope of -1 independent of N
- Intercept ρ_{xy}^{Int} proportional to N
- $\rho_{xy}^{Int} = 22 \pm 2$ n Ω cm per N
- Comparable to anomalous Hall signal for $r_{sky} > 125$ nm



TOPOLOGICAL HALL RESISTIVITY
expected from Berry phase theory

- $\rho_{xy}^T = (P R_0 \phi_0 / A) S \approx 0.003$ n Ω cm

Signal of similar magnitude observed
by Raju et al. arXiv:1708.04084



ARTICLE

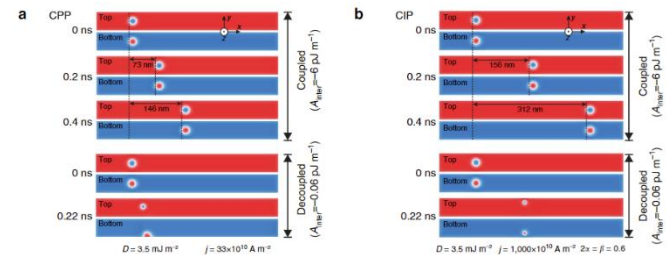
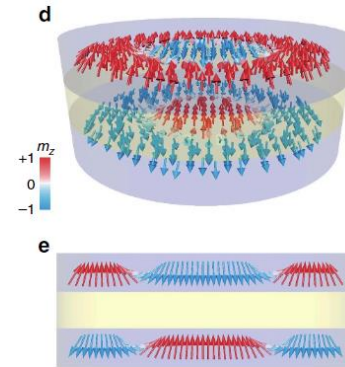
Received 21 Apr 2015 | Accepted 26 Nov 2015 | Published 19 Jan 2016

DOI: 10.1038/ncomms10293 OPEN

Magnetic bilayer-skyrmions without skyrmion Hall effect

Xichao Zhang^{1,2}, Yan Zhou^{1,2} & Motohiko Ezawa³

Magnetic skyrmions might be used as information carriers in future advanced memories, logic gates and computing devices. However, there exists an obstacle known as the skyrmion Hall effect (SkHE), that is, the skyrmion trajectories bend away from the driving current direction due to the Magnus force. Consequently, the skyrmions in constricted geometries may be destroyed by touching the sample edges. Here we theoretically propose that the SkHE can be suppressed in the antiferromagnetically exchange-coupled bilayer system, since the Magnus forces in the top and bottom layers are exactly cancelled. We show that such a pair of SkHE-free magnetic skyrmions can be nucleated and be driven by the current-induced torque. Our proposal provides a promising means to move magnetic skyrmions in a perfectly straight trajectory in ultra-dense devices with ultra-fast processing speed.

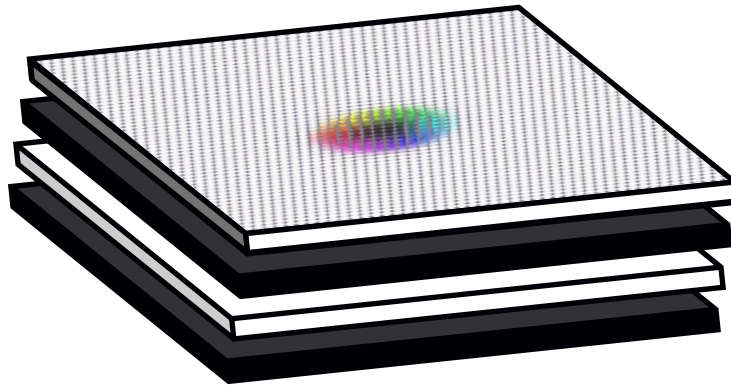


- Antiferromagnetically coupled layers
- No skyrmion Hall effect – no net topological charge

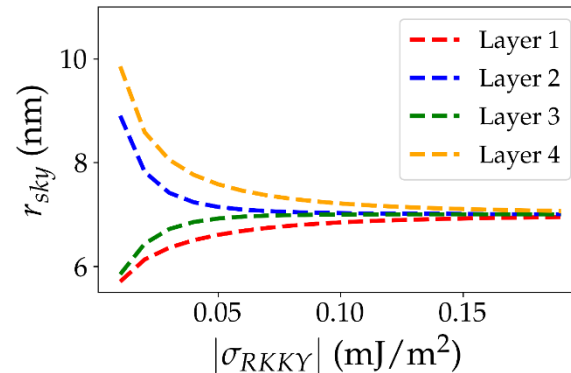
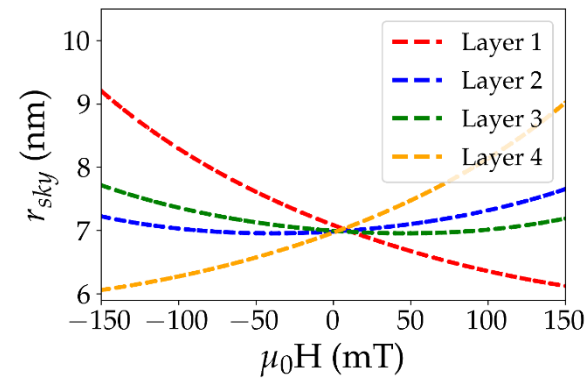


SAF skyrmion breathing modes: our model

MuMax³

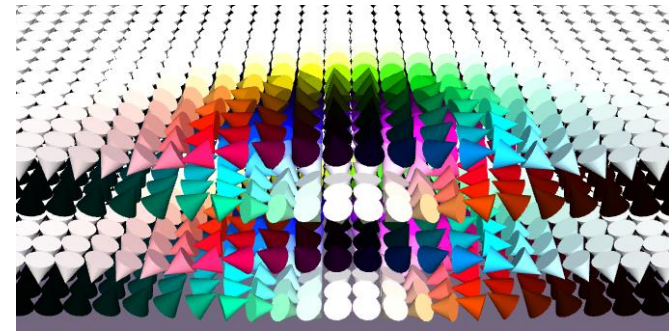
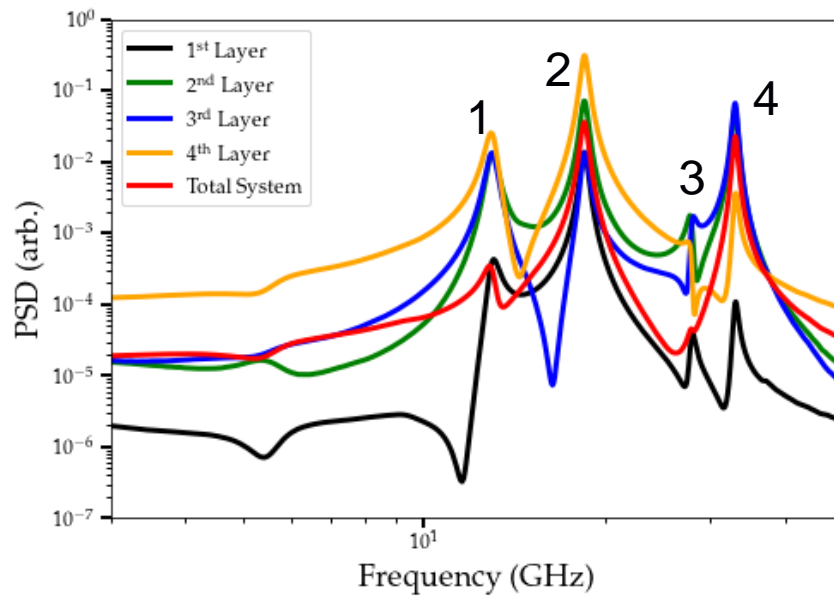


- 100 nm square
- Periodic boundary conditions
- $M_s = 1 \times 10^6$ A/m, $A = 15 \times 10^{-12}$ J/m, $K = 1 \times 10^6$ J/m³, $D = 2.5 \times 10^{-3}$ J/m²





SAF skyrmion breathing modes



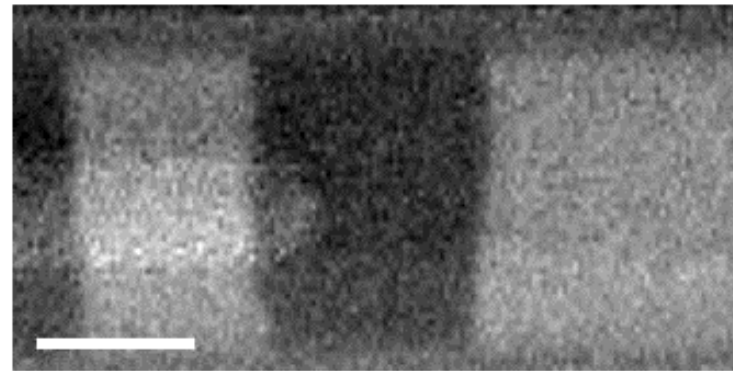
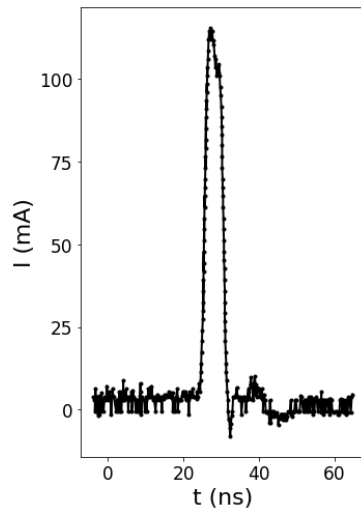
4 layers => 4 modes

Domain wall motion – current pulse response



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Nominally 5 ns pulses



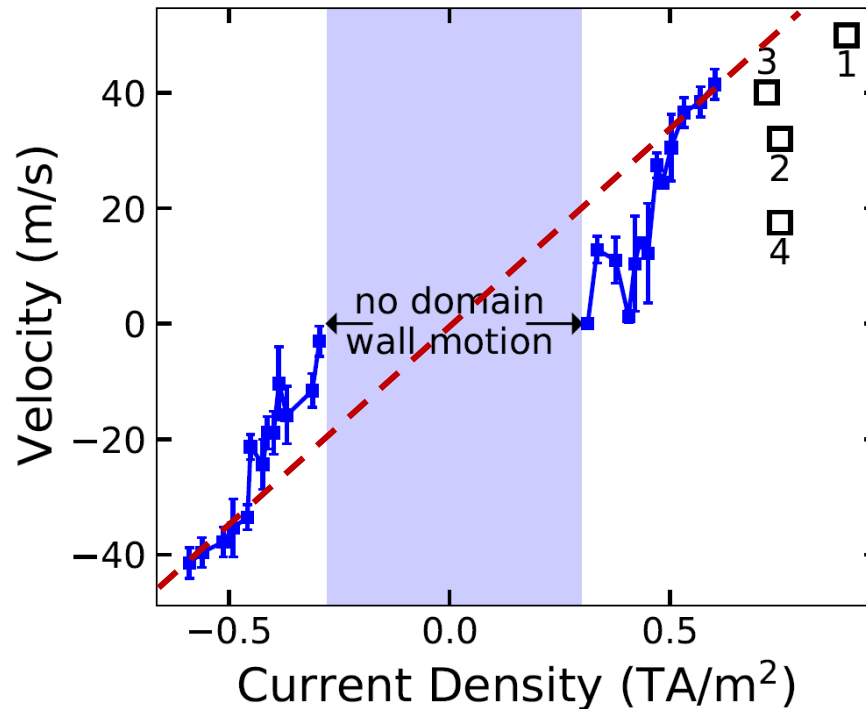
1 μm

Sequences of 10 positive,
then 10 negative pulses
 $3.6 - 5.0 \times 10^{11} \text{ A/m}^2$
XMCD images

Domain wall motion – current pulse response



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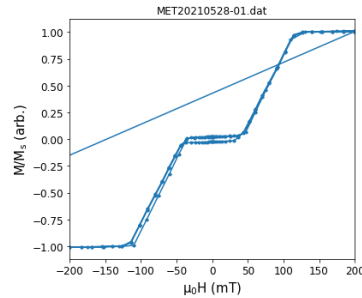
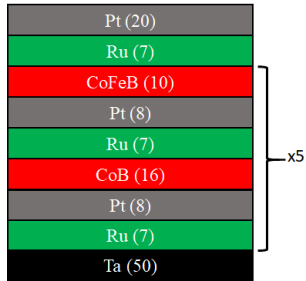
[1] S. Finizio *et al.*, Nano Lett. 19, 375–380 (2018). [2,4] 19S. Finizio, S. *et al.*, Appl. Phys. Lett. 117, 212404 (2020). [3] 20S. Finizio, S. *et al.*, Appl. Phys. Lett. 116, 182404 (2020).



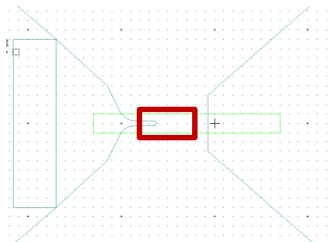
Electrical SAF skyrmion injection

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10 ns pulses



As-grown domain pattern



Wire 2 μm wide

Injector 250 nm radius of curvature



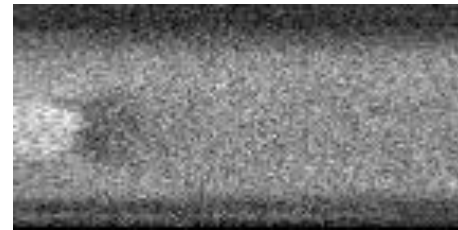
Co edge



Fe edge



$0.8 \times 10^{12} \text{ A/m}^2$



$0.9 \times 10^{12} \text{ A/m}^2$



$1.1 \times 10^{12} \text{ A/m}^2$

Co edge +50 Oe

- Deterministic **nucleation** of single skyrmions observed at 200 ps resolution
Finizio et al., Nano Letters **19**, 7246 (2019)
- Observed **propagation** quasistatically through complex energy landscape
Zeissler et al., Nature Commun. **11**, 428 (2020)
- Hall signal for **detection** has two parts
 - size-dependent part \longrightarrow AHE \checkmark
 - size-independent part, 22 ± 2 n Ω cm per skyrmion
 \longrightarrow revised Berry-phase THE or new mechanism?
Zeissler et al., Nature Nanotech. **13**, 1161 (2018)

