

SPICE-Workshop on Altermagnetism

Spin crystalline group in magnetic materials

Qihang Liu (刘奇航)

Department of Physics and Shenzhen Institute for
Quantum Science and Technology
Southern University of Science and Technology (SUSTech)

2023. 5. 9

[Phys. Rev. X 12, 021016 \(2022\)](#)

[The Innovation 3, 100343 \(2022\)](#)

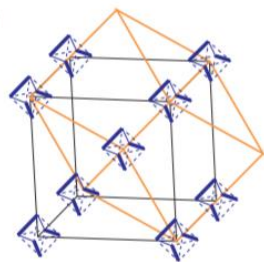
[arXiv:2303.04549 \(2023\)](#)

[arXiv:2301.12201 \(2023\)](#)

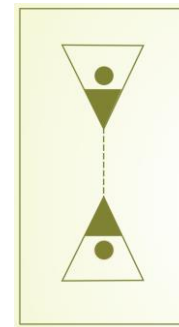
Previous connection to spintronics

- ❖ 2012-2013 **Northwestern University** (Postdoc with A. Freeman)
- ❖ 2013-2018 **CU Boulder** (Postdoc with A. Zunger)
- ❖ **2018.4 - SUSTech Associate Professor**

High global symmetry + Local symmetry breaking



Hidden Spin Polarization



Concept (LaOBiS₂):

QL* et al. *Nano Lett.* 13, 5264 (2013)

X. Zhang[†], QL†, J. Luo*, A. Zunger * et al.
Nat. Phys. 10, 387 (2014)

New materials:

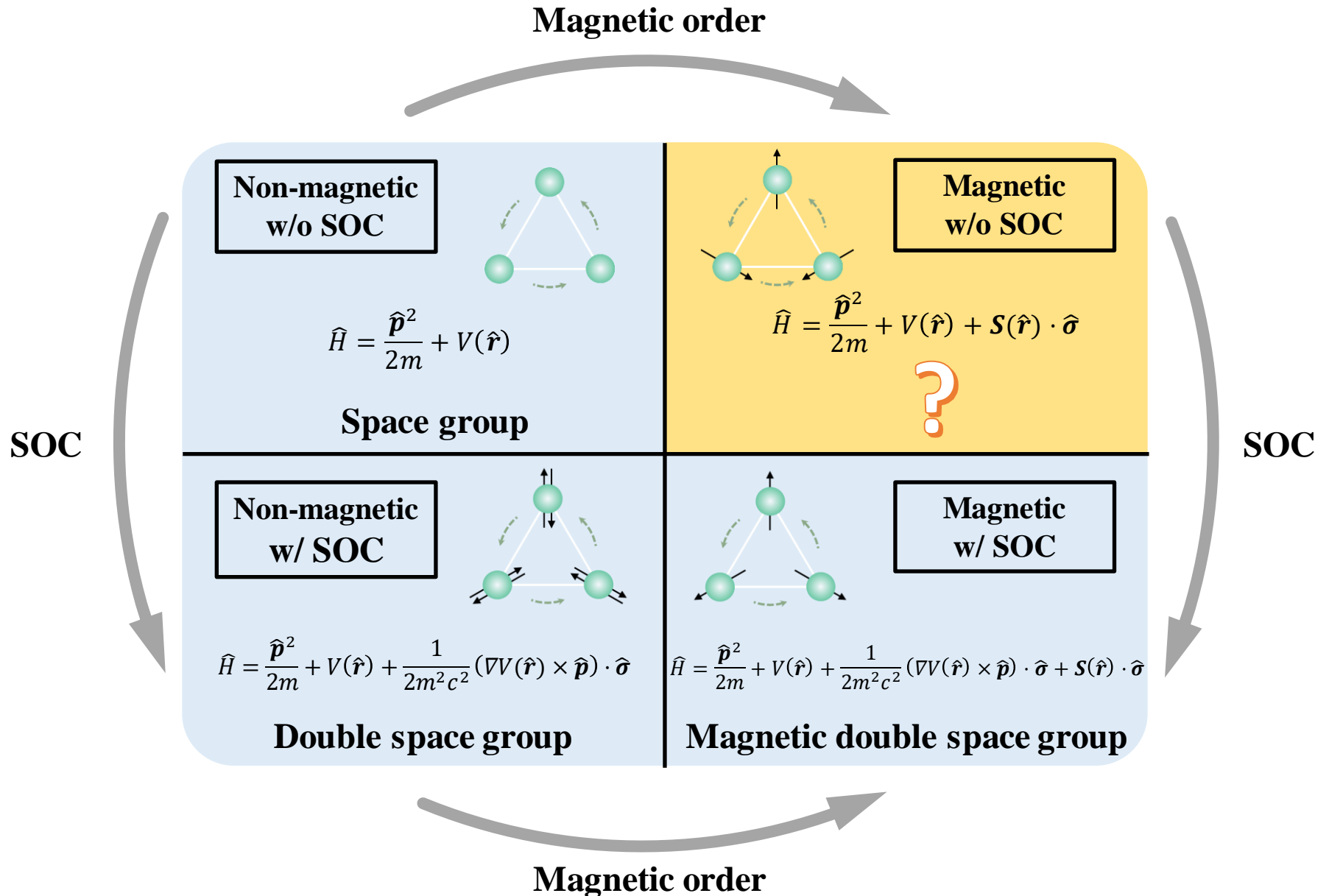
BiOI: C. Chen*, QL* et al. *PRL* 127, 126402 (2021)

CuMnAs: QL* et al. *PRL* 129, 276601 (2022)

Outline

- **Introduction of spin group**
- **Features of spin point/space group**
- **Topological phases protected by spin group symmetry**
 - **Kramers degeneracy and Z_2 topological phases**
 - **Chiral Dirac-like semimetal**

Symmetry considerations in crystals – Group Theory



Symmetry considerations in magnetic materials

w/ SOC

$$\hat{H} = \frac{\hat{\mathbf{p}}^2}{2m} + V(\hat{\mathbf{r}}) + \frac{1}{2m^2 c^2} (\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}) \cdot \hat{\boldsymbol{\sigma}} + \mathbf{S}(\hat{\mathbf{r}}) \cdot \hat{\boldsymbol{\sigma}}$$

$$\begin{aligned} & \widehat{C_n(\theta)} \left(\frac{1}{2m^2 c^2} (\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}) \cdot \hat{\boldsymbol{\sigma}} \right) \widehat{C_n(\theta)}^{-1} \\ &= \frac{1}{2m^2 c^2} (R_n(\theta) \nabla V(R_n(\theta)^{-1} \hat{\mathbf{r}}) \times R_n(\theta) \hat{\mathbf{p}}) \cdot \hat{\boldsymbol{\sigma}} \\ &= \frac{1}{2m^2 c^2} R_n(\theta) (\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}) \cdot \hat{\boldsymbol{\sigma}} \\ & \widehat{U_n(\theta)} \widehat{C_n(\theta)} \left(\frac{1}{2m^2 c^2} (\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}) \cdot \hat{\boldsymbol{\sigma}} \right) \widehat{C_n(\theta)}^{-1} \widehat{U_n(\theta)}^{-1} \\ &= \frac{1}{2m^2 c^2} R_n(\theta) (\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}) \cdot R_n(\theta) \hat{\boldsymbol{\sigma}} \\ &= \frac{1}{2m^2 c^2} (\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}) \cdot \hat{\boldsymbol{\sigma}} \end{aligned}$$

Magnetic double group with every spatial rotation operator attached to a spin rotation operator

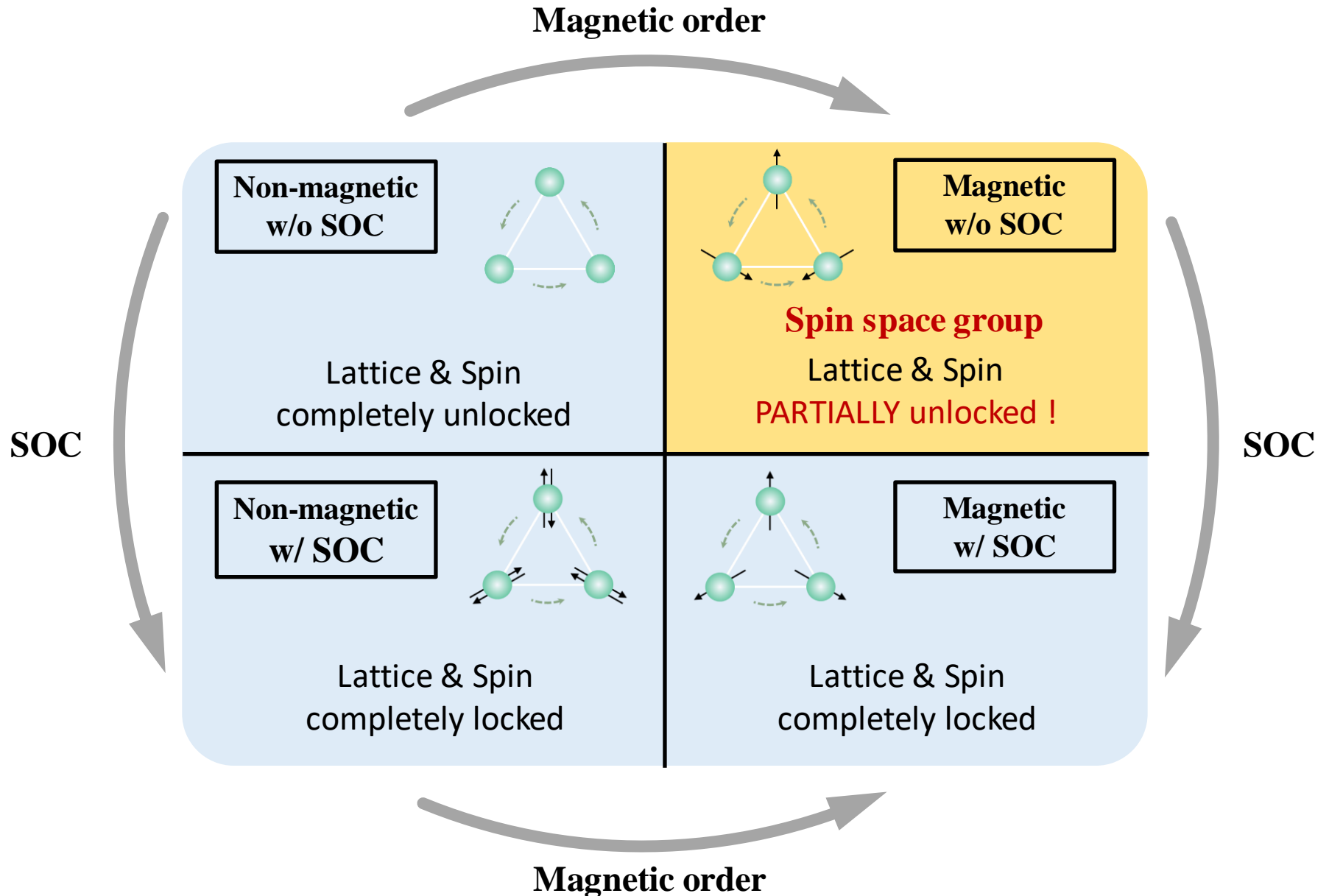
w/o SOC

$$\hat{H} = \frac{\hat{\mathbf{p}}^2}{2m} + V(\hat{\mathbf{r}}) + \mathbf{S}(\hat{\mathbf{r}}) \cdot \hat{\boldsymbol{\sigma}}$$

$$\begin{aligned} & \widehat{C_n(\theta)} \mathbf{S}(\hat{\mathbf{r}}) \cdot \hat{\boldsymbol{\sigma}} \widehat{C_n(\theta)}^{-1} \\ &= \widehat{C_n(\theta)} \mathbf{S}(\hat{\mathbf{r}}) \widehat{C_n(\theta)}^{-1} \cdot \hat{\boldsymbol{\sigma}} \\ &= \underline{\mathbf{S}(R_n(\theta)^{-1} \hat{\mathbf{r}})} \cdot \hat{\boldsymbol{\sigma}} \\ & \quad = \mathbf{S}(\hat{\mathbf{r}}) ? \\ & \widehat{U_m(\varphi)} \widehat{C_n(\theta)} \mathbf{S}(\hat{\mathbf{r}}) \cdot \hat{\boldsymbol{\sigma}} \widehat{C_n(\theta)}^{-1} \widehat{U_m(\varphi)}^{-1} \\ &= \mathbf{S}(R_n(\theta)^{-1} \hat{\mathbf{r}}) \cdot R_m(\varphi)^{-1} \hat{\boldsymbol{\sigma}} \\ &= \underline{R_m(\varphi) \mathbf{S}(R_n(\theta)^{-1} \hat{\mathbf{r}})} \cdot \hat{\boldsymbol{\sigma}} \\ & \quad = \mathbf{S}(\hat{\mathbf{r}}) \end{aligned}$$

Spin group with spin rotation operators detached from spatial rotation operators

Symmetry considerations in crystals – Group Theory



Original works about spin group symmetry

JOURNAL OF APPLIED PHYSICS

VOLUME 37, NUMBER 3

1 MARCH 1966

Space Group Theory for Spin Waves

W. BRINKMAN* AND R. J. ELLIOTT†

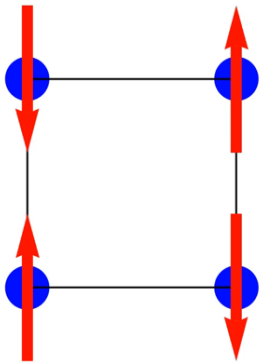
Department of Theoretical Physics, Oxford University, Oxford, England

Theory of spin-space groups

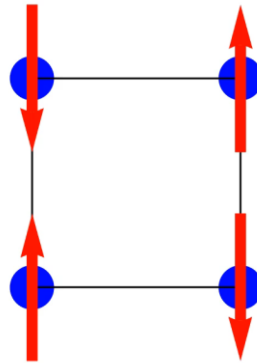
BY W. F. BRINKMAN† AND R. J. ELLIOTT

Department of Theoretical Physics, University of Oxford, England

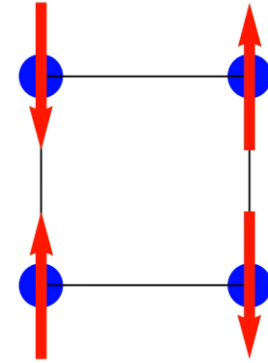
- Multiple-fold degeneracies of spin wave spectrum require an enhanced symmetry group called “spin group”, which consider spin rotation and spatial rotation separately.



Consider an AFM
spin arrangement

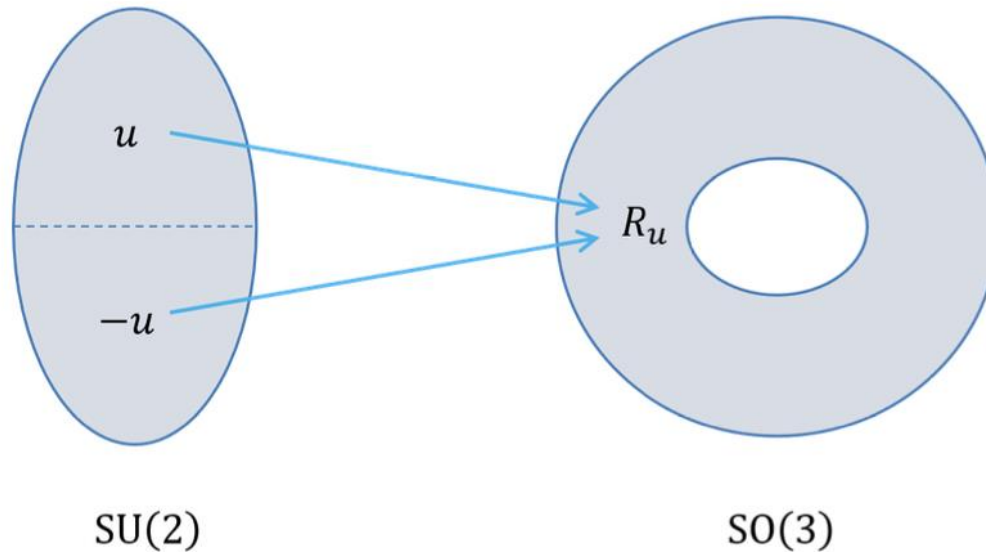


C_2 rotation in
magnetic group



C_2 rotation in
spin group

We are talking about spin “**crystalline**” group



In group theory, **spin group** is the double cover of $SO(N)$

$$\text{Spin}(1) = O(1) = \{+1, -1\}$$

$$\text{Spin}(2) = U(1) = SO(2)$$

$$\text{Spin}(3) = SU(2)$$

Outline

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- **Features of spin point/space group**
- Topological phases protected by spin group symmetry
 - Kramers degeneracy and Z_2 topological phases
 - Chiral Dirac-like semimetal
- Outlook

Separating spin & lattice operation — notations

Spin rotation

Lattice rotation

Spin point group: $\{U_m(\varphi), TU_m(\varphi) || C_n(\theta), IC_n(\theta)\}$

Spin space group: $\{U_m(\varphi), TU_m(\varphi) || C_n(\theta), IC_n(\theta) | \tau\}$

Lattice rotation & translation

There are 598 types of spin point group

D. B. Litvin and W. Opechowski, *Physica* **76**, 538 (1974); D. B. Litvin, *Acta Crystallogr. A* **33**, 279 (1977)

Theory to explore:

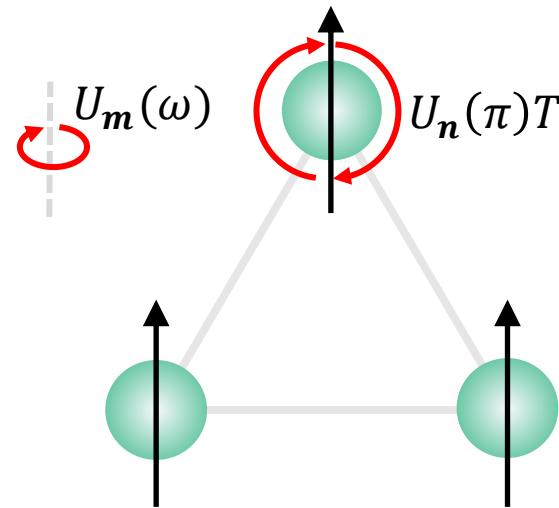
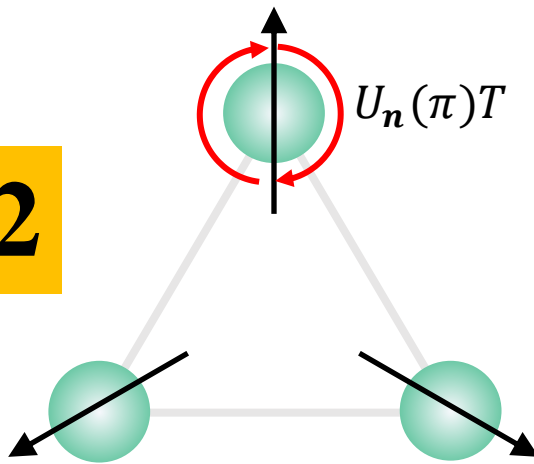
- Spin point groups for specific magnetic configuration?
- Classification of spin space groups?
- Representation theory of spin space groups?
- Application of spin group (Spintronics? Topology?)

Spin groups for **collinear** and **coplanar** spin arrangements

Coplanar: $G_{\text{SOP}} = Z_2^K$

Collinear: $G_{\text{SOP}} = SO(2) \times Z_2^K$

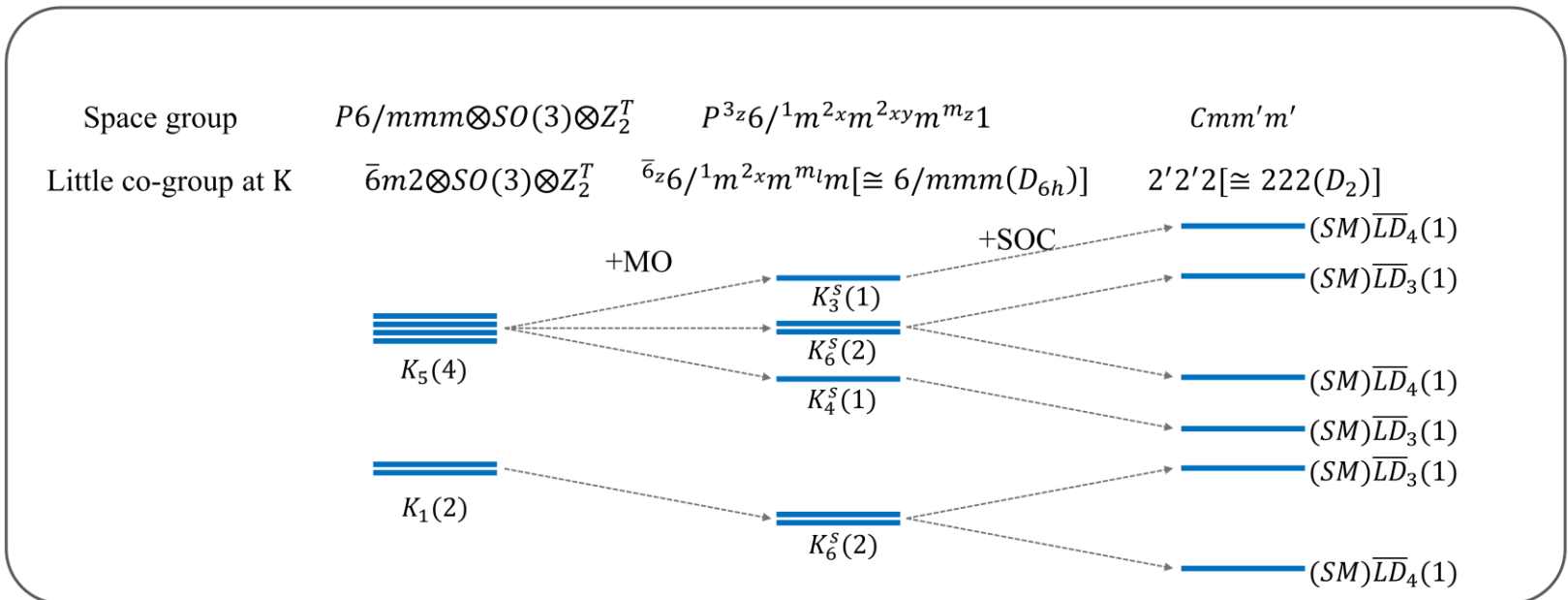
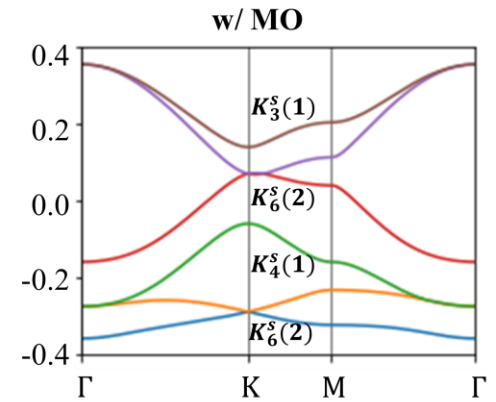
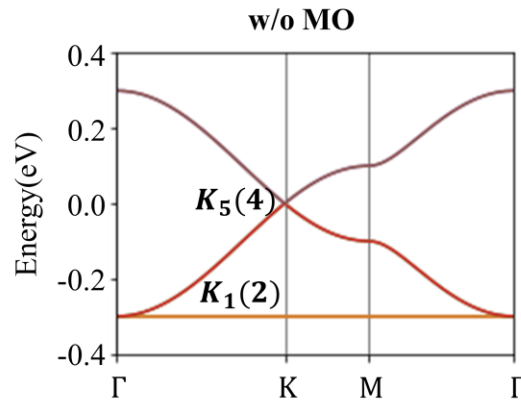
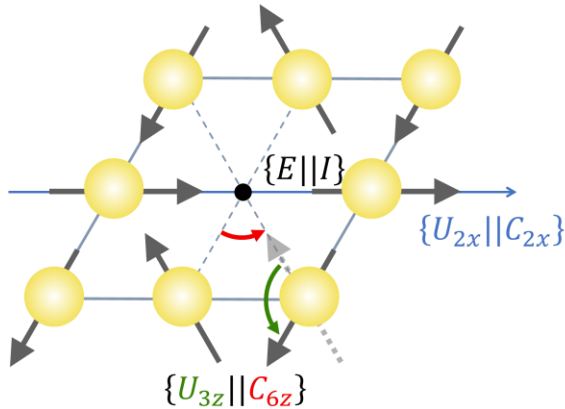
252



90

$$Z_2^K = \{\{E||E\}, \{TU_n(\pi)||E\}\}; SO(2) = \{\{U_m(\omega)T||E|0\}|\omega \in (0, 2\pi]\}$$

Degeneracies caused by spin space group



➤ 2-fold degeneracies occur at K are well explained by spin group symmetry

Applications of spin group symmetry (Before 2022)

Previous discussions

Landau theory of phase transition

Phys. Rev. B **26**, 6947 (1982);
I. A. Izyumov et al., *Phase transitions and crystal symmetry*;...

Neutron scattering

Thorpe, Proc. Phys. Soc. **91**, 903 (1967);
Acta Crystallogr. A **29**, 651 (1973);...

Electronic states of spiral magnet

J. Phys. Condens. Matter **3**, 8565 (1991)

Spin group symmetry

Recent applications

Spin splitting in AFM systems

J. Phys. Soc. Jpn. **88**, 123702 (2019);
Phys. Rev. B **102**, 014422 (2020);
PNAS **118**, e2108924118 (2021)...

Spin Hall effect in AFM systems

Phys. Rev. Lett. **119**, 187204 (2017);
Phys. Rev. Lett. **126**, 127701 (2021);...

Piezomagnetism

Nat. Commun. **12**, 2846 (2021)

Other possible applications ?

Topological phase of matter

- Introduction of spin group
- Features of spin point/space group
- **Topological phases protected by spin group symmetry**
 - **Kramers degeneracy and Z_2 topological phases**
 - Chiral Dirac-like semimetal

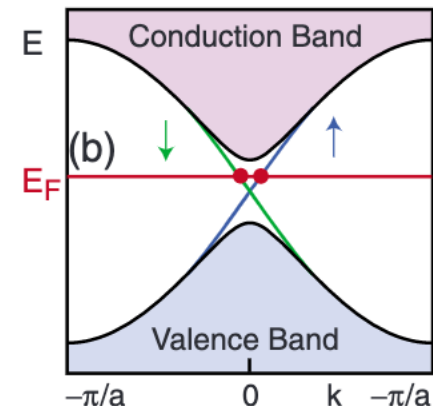
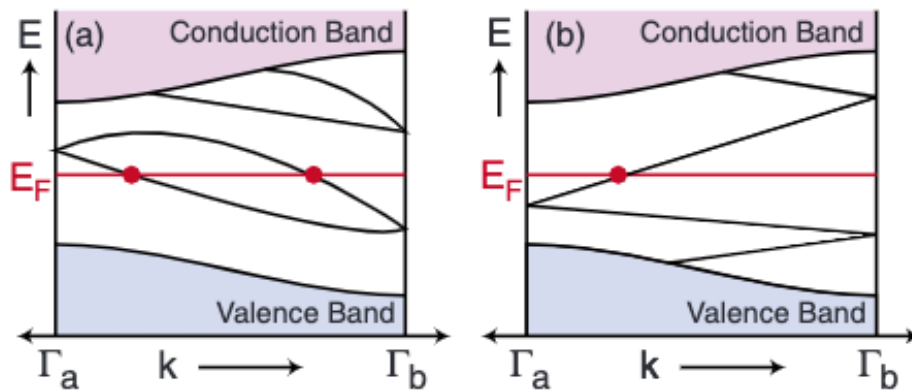
Kramers degeneracy and Z_2 topological classification

➤ Kramers degeneracy

1. An antiunitary symmetry operator Θ (i.e., $\langle \Theta\phi | \Theta\phi \rangle = \langle \phi | \phi \rangle$)
2. $\Theta^2 = -1$

➤ Z_2 topological classification in 2D

- Time reversal symmetry T of electronic system could protect Kramers degeneracy
- Degeneracy at time reversal invariant momenta (TRIM) for surface states
- Two types of surface states connection in edges of 2D systems corresponds to Z_2 topological trivial and nontrivial phase (Dirac surface states)



Rev. Mod. Phys. **82**, 3045 (2010)

Kramers degeneracy protected by spin group operations

- Find all symmetries that could protect Kramers degeneracy and are unbroken on certain surfaces

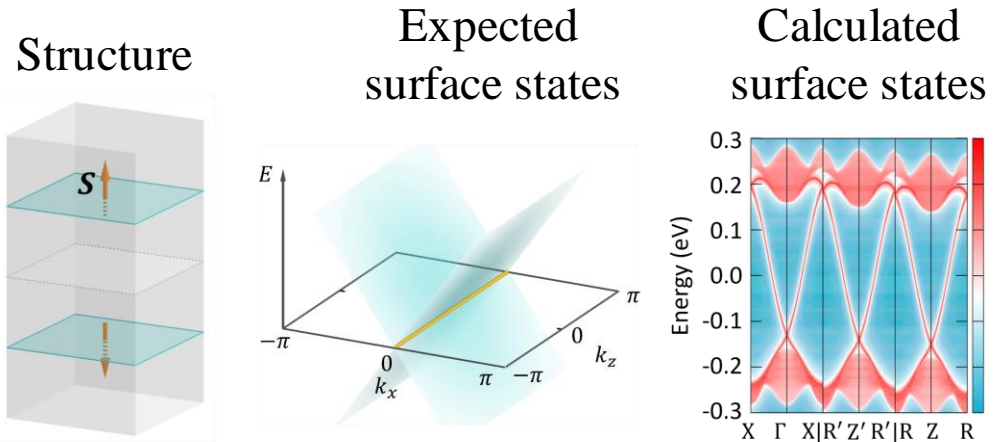
Spin group Symmetry	Momenta with protected 2- fold degeneracy	Surface with the symmetry	Possible surface states
$\{T E \tau_{z/2}\}$	TRIM within $k_z = 0$ plane	$(xy0)$	DP at $(0,0)$ or $(\pi, 0)$
$\{TU_z(\pi) m_{[001]} \tau_{x/2}\}$	$(\pi, 0, k_z)$ and (π, π, k_z) lines	(010)	DNL at $k_x = \pi$
$\{T C_z(\pi) 0\}$	$k_z = 0$ and $k_z = \pi$ planes	(001)	Possible double DP
$\{T m_{[001]} 0\}$	$(0,0, k_z), (0, \pi, k_z),$ $(\pi, 0, k_z)$ and (π, π, k_z) lines	$(xy0)$	DNL at $k_x = 0$ or $k_x = \pi$
$\{T m_{[001]} \tau_{x/2}\}$	$(0,0, k_z)$ and $(0, \pi, k_z)$ lines	(010)	DNL at $k_x = 0$
$\{TU_n(\pi) m_{[001]} \tau_{x/2}\}$	$(\pi, 0, k_z)$ and (π, π, k_z) lines	(010)	DNL at $k_x = \pi$
$\{TU_n(\pi) E \tau_{z/2}\}$	TRIM within $k_z = \pi$ plane	$(xy0)$	DP at $(0, \pi)$ or (π, π)

DP—Dirac point; DNL—Dirac nodal line

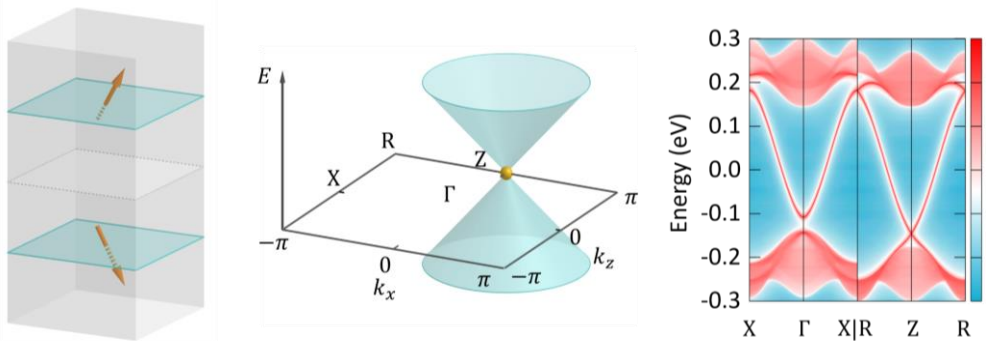
X. Wan*, QL* et al. PRX 12, 021016 (2022)

Z_2 topological phases realized by a tight binding model

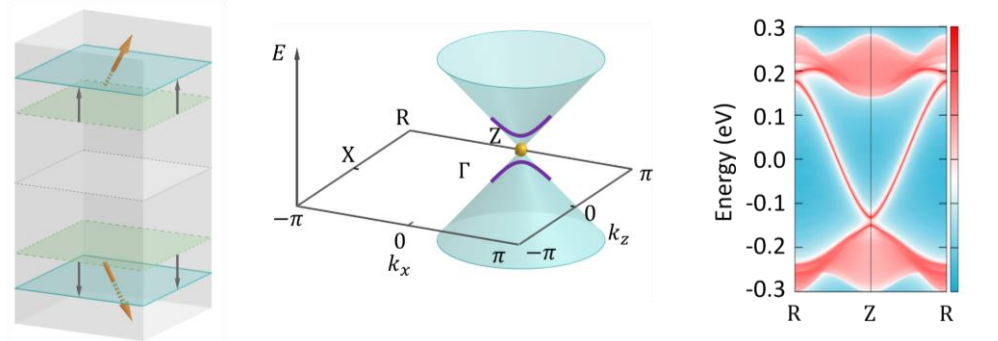
- Z_2 Topological insulator protected by $\{T || m_{[001]} | 0\} \& \{TU_n(\pi) || E | \tau_z/2\}$



- Z_2 Topological insulator protected by $\{TU_n(\pi) || E | \tau_z/2\}$

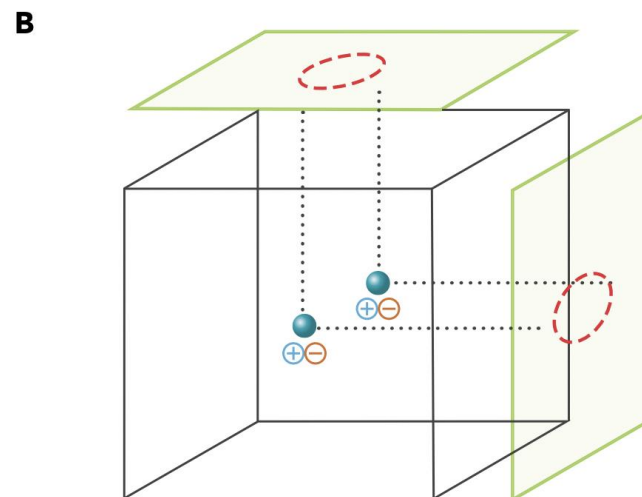
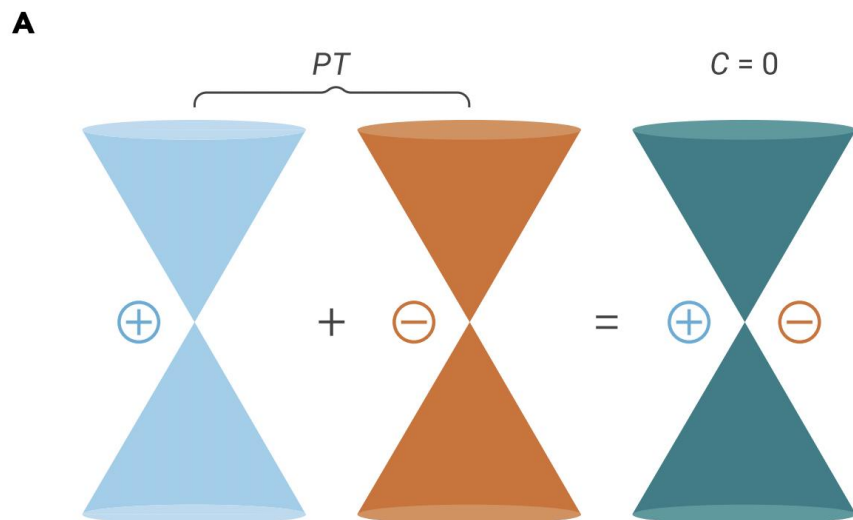


- Trivial insulator

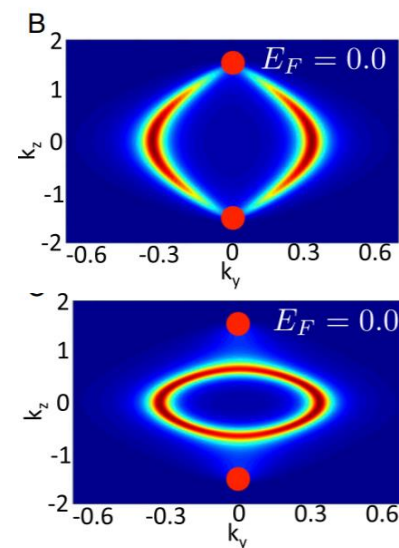


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- **Topological phases protected by spin group symmetry**
 - Kramers degeneracy and Z_2 topological phases
 - **Chiral Dirac-like semimetal**

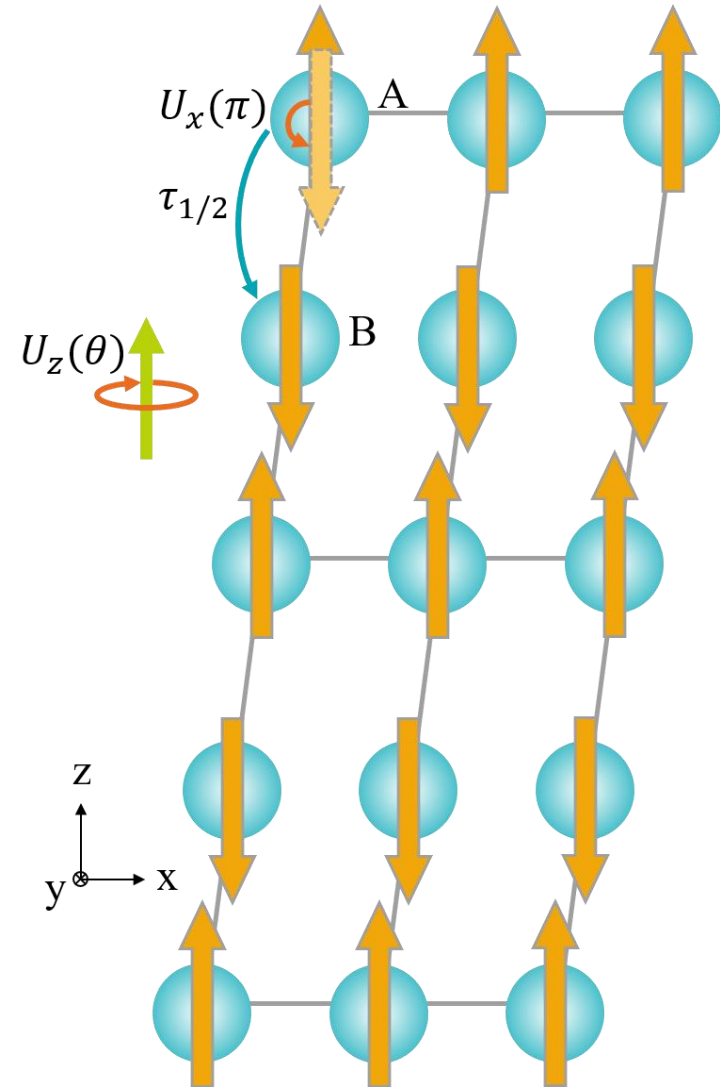
Conventional Dirac semimetal



- A Dirac cone is composed of two Weyl cones with opposite chirality
- Appear at high-symmetry line/point
- The surface states are not topologically protected



Hidden SU(2) symmetry in certain AFM systems without SOC



- Consider a collinear AFM structure with type-IV magnetic space group.

$$u_x^{1/2} \equiv \{U_x(\pi) || E | \tau_{1/2}\} = -ie^{-ik \cdot \tau_{1/2}} \tau_x \otimes \sigma_x$$

$$u_y^{1/2} \equiv \{U_y(\pi) || E | \tau_{1/2}\} = -ie^{-ik \cdot \tau_{1/2}} \tau_x \otimes \sigma_y$$

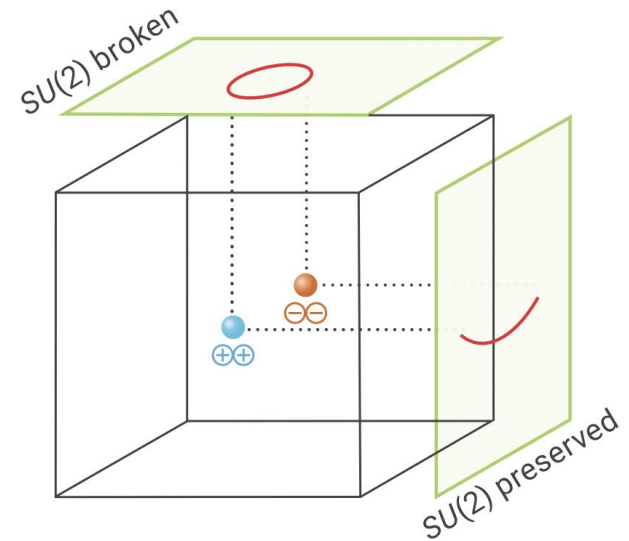
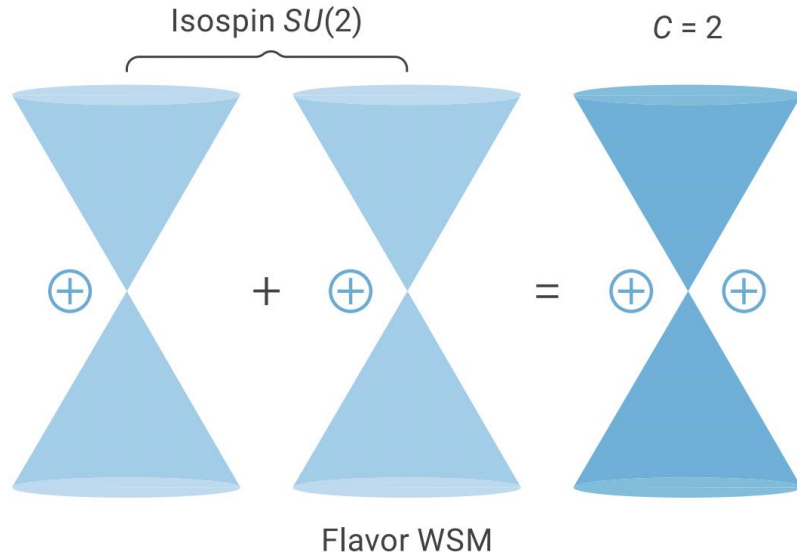
$$u_z \equiv \{U_z(\pi) || E | 0\} = -i\tau_0 \otimes \sigma_z$$

Generators of $su(2)$ Lie algebra for an arbitrary k

$$\Rightarrow SU(2) = \{\exp(-i\theta \cdot \rho)\},$$

$$\rho = \left(\frac{1}{2} \tau_x \otimes \sigma_x, \frac{1}{2} \tau_x \otimes \sigma_y, \frac{1}{2} \tau_0 \otimes \sigma_z \right)$$

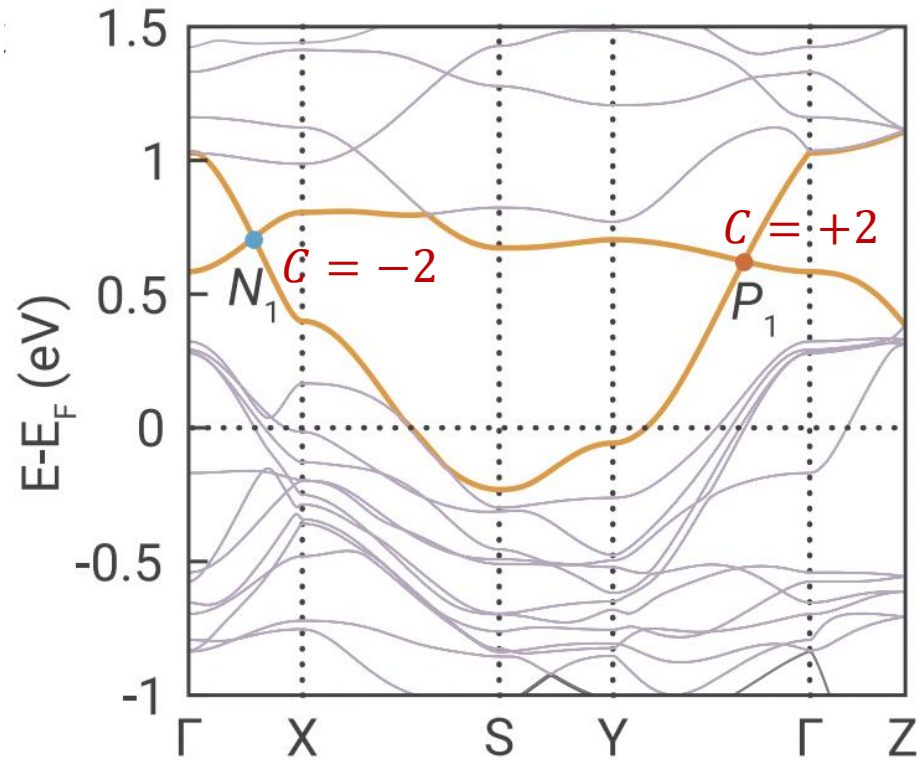
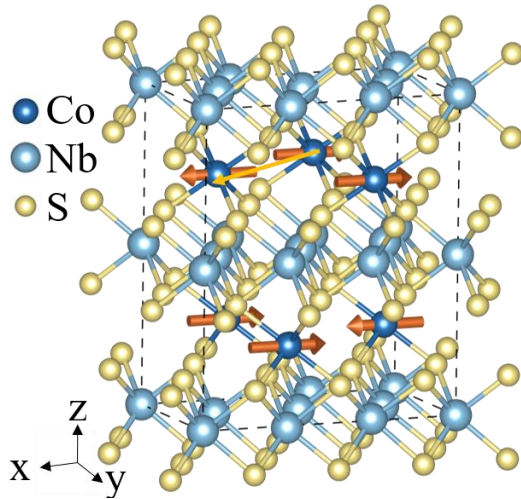
A Dirac-like fermion but with chirality



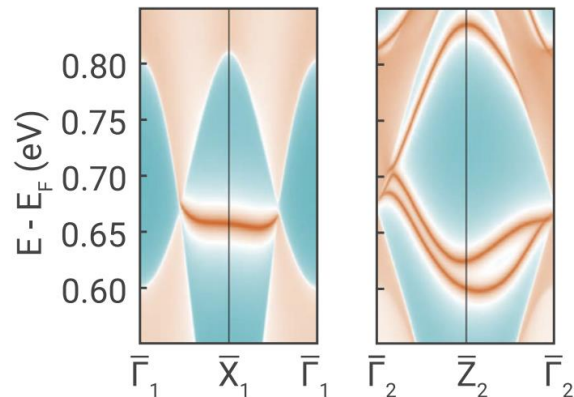
- Two Weyl quasiparticles with same chirality
- Robust Fermi arcs
- Dirac-like on certain surfaces while Weyl-like on other surfaces

Materials realization

- CoNb_3S_6 is a representative material of chiral Dirac-like semimetal



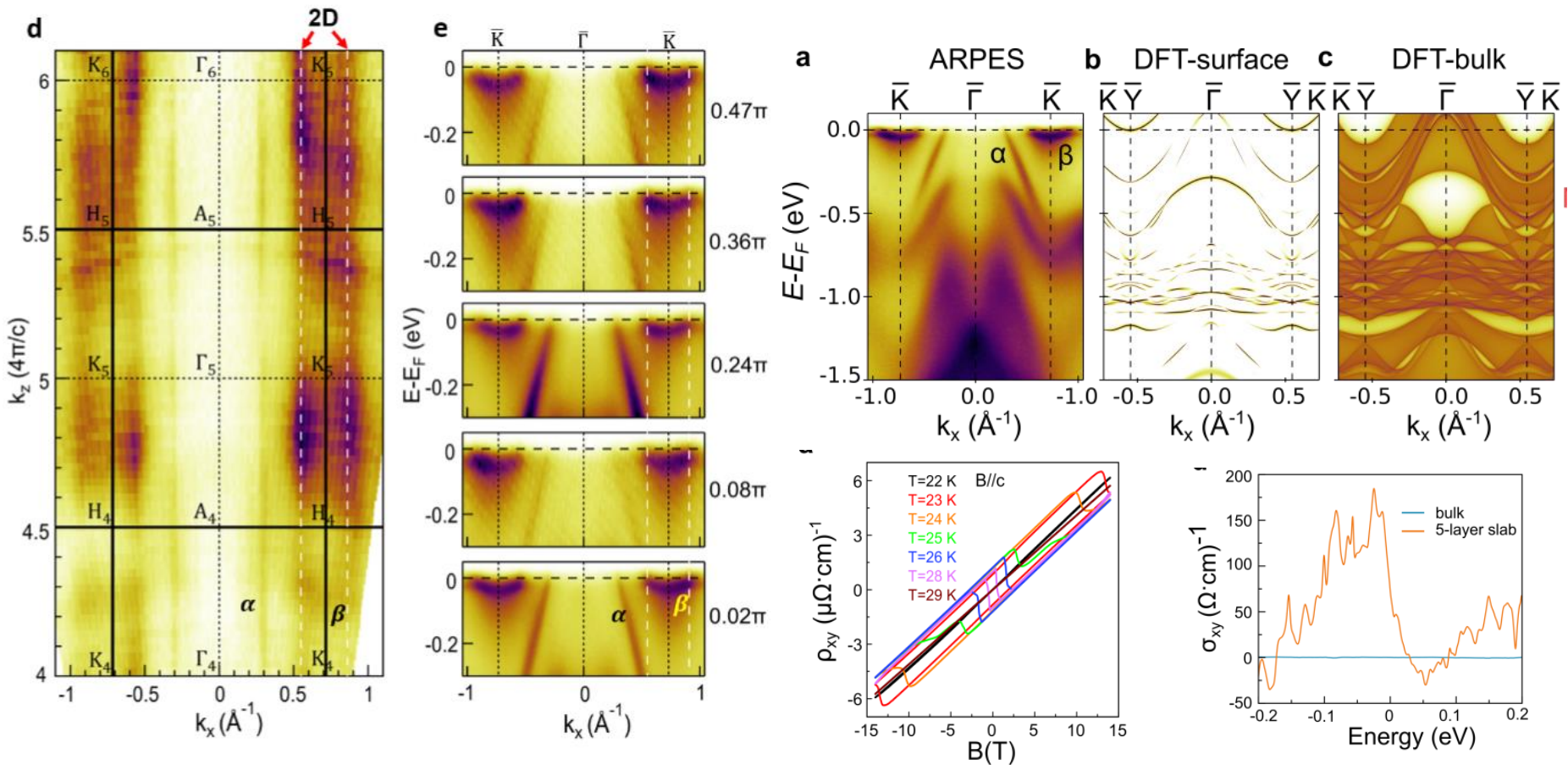
D $SU(2)$ preserved $SU(2)$ broken



- P_1 and P_2 have chirality $+2$
- N_1 and N_2 have chirality -2 .

Experimental verification by ARPES and neutron diffraction

➤ The comparison between ARPES and DFT reveal the Fermi arc surface states of CoNb_3S_6



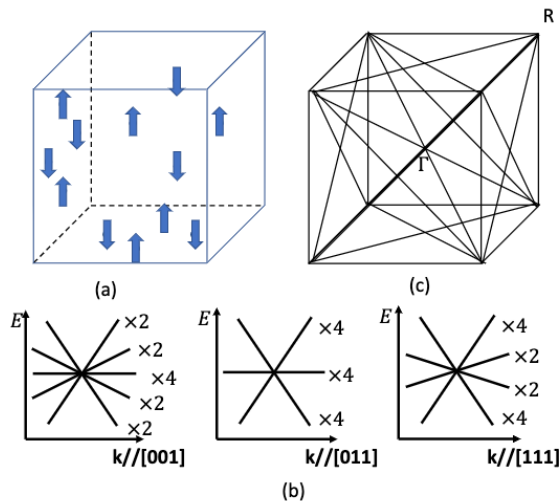
➤ k_z dispersion shows 2D nature

➤ Mysterious Anomalous Hall effect explained by chiral Dirac fermion

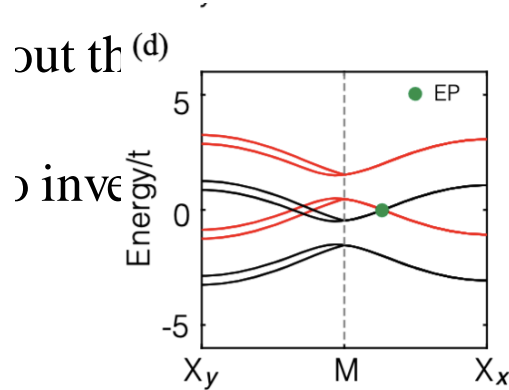
Summary and What's next ?

Spin crystalline group — symmetry description of magnetic materials in SOC-free limit:

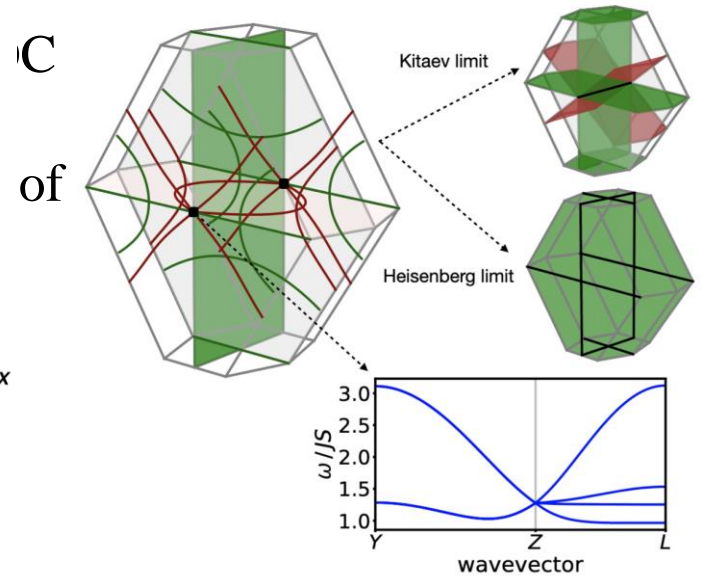
- 1) New symmetry operations, new degeneracies, and new quasiparticles
- 2) New topological phases and new topological classifications



Symmetry invariants
arxiv:2105.12738

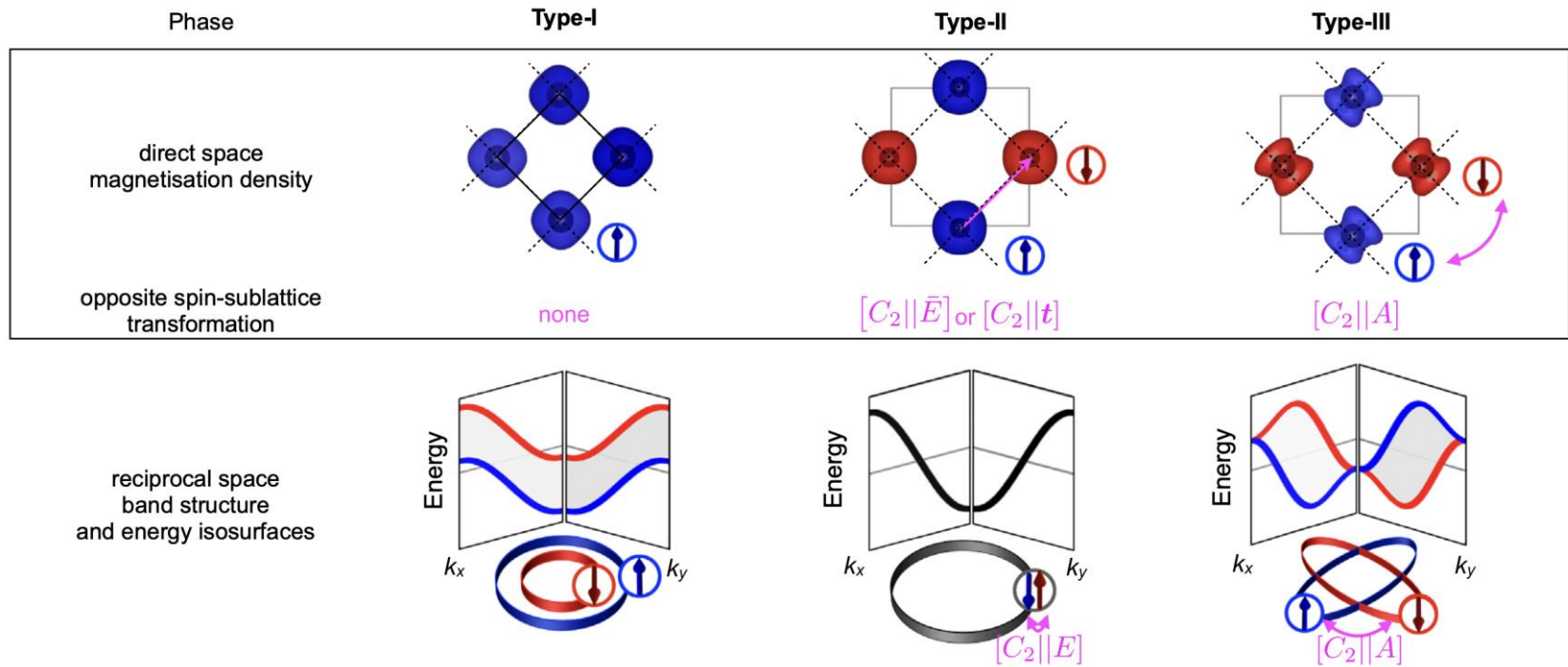


New fermions
PRL 127, 176401 (2021)



Magnon topology
PRB 105, 064430 (2022)

Altermagnetism and spintronics



RuO₂, FeSb₂, MnF₂, CrSb, MnTe, VNb₃S₆...

PRX 12, 031042 (2022); PRX 12, 040501 (2022)

Spin splitting torque

PRL 128, 197202 (2022)

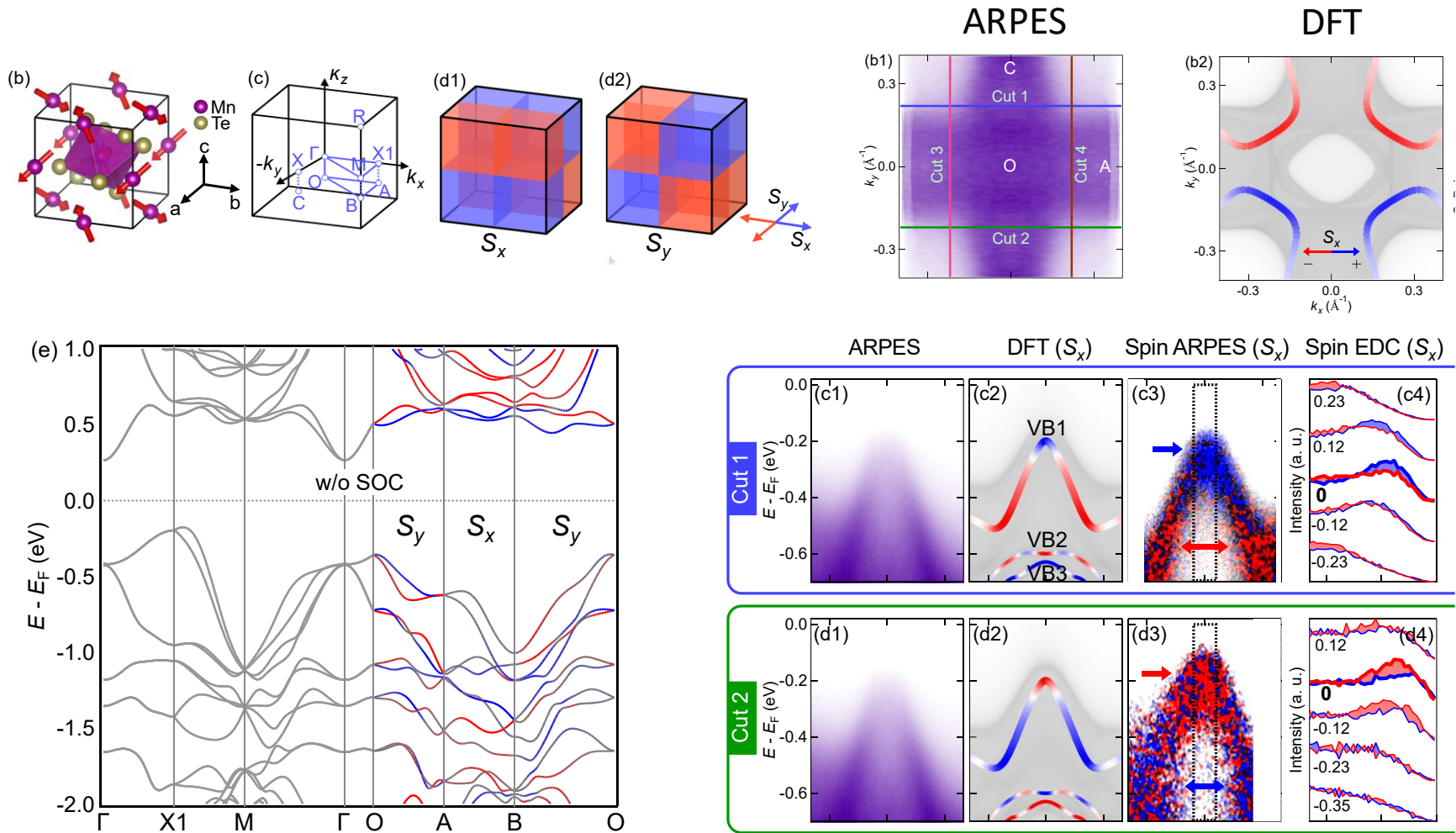
Spin Hall effect

PRL 126, 127701 (2021)

Tunneling Magnetoresistance

PRX 12, 011028 (2022)

Plaid-like spin splitting in a noncoplanar antiferromagnet MnTe_2



Plaid-like spin texture
in the $k_z = -0.4 \pi/c$ (O-A-B-C) plane

Spin-ARPES shows symmetric/antisymmetric
spin polarization along k_x/k_y axes

Acknowledgements

Collaborators

Prof. Xiangang Wan (NJU) Theory

Prof. Chaoyu Chen (SUSTech) ARPES

Prof. Chang Liu (SUSTech) ARPES

Prof. Liusuo Wu (SUSTech) Neutron diffraction

Pengfei Liu



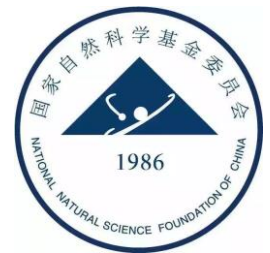
Jiayu Li



Ao Zhang



Xiaobing Chen



Please also check our posters:

5149: Antiferromagnetic Chern insulators

5151: Observation of plaid-like spin splitting in a noncoplanar
antiferromagnet

Thank you for your attention !

[Phys. Rev. X 12, 021016 \(2022\)](#)

[The Innovation 3, 100343 \(2022\)](#)

[arXiv:2303.04549 \(2023\)](#)

[arXiv:2301.12201 \(2023\)](#)



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SOUTHERN UNIVERSITY OF SCIENCE AND TECHNOLOGY