

Floquet time-crystals as sensors of AC fields

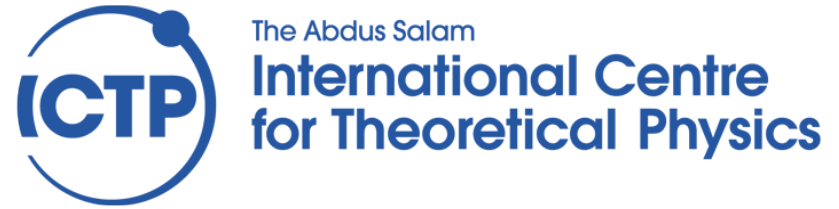
Fernando Iemini

Physics Institute, Universidade Federal Fluminense (UFF),
Niterói, Rio de Janeiro, Brazil*(on leave)

Johannes Gutenberg University, AvH Fellow in
NEUQUAM group– till 2024
(Henriette Herz programme)

F. Iemini, R. Fazio, A. Sanpera, **arXiv:2306.03927** (2023)

In collaboration with:



- Rosario Fazio
(ICTP / Università di Napoli)



UAB

Universitat Autònoma
de Barcelona

- Anna Sanpera
(ICREA/ Universitat
Autònoma de Barcelona)



Outline

- Introduction:
 - definition: what is a discrete time-crystal?
 - . from theoretical proposals to experimental observation.
- Sensing with DTC's

Introduction

- **Landau's symmetry breaking** is a cornerstone of modern physics:

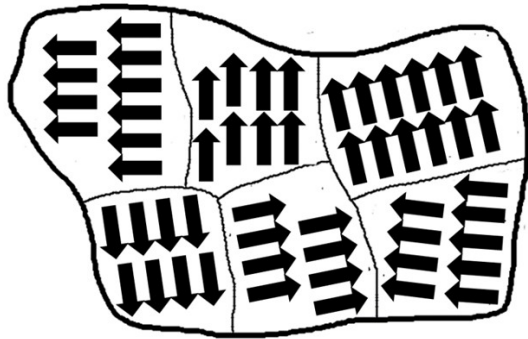
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Ferromagnetism
(spin rotation SB)

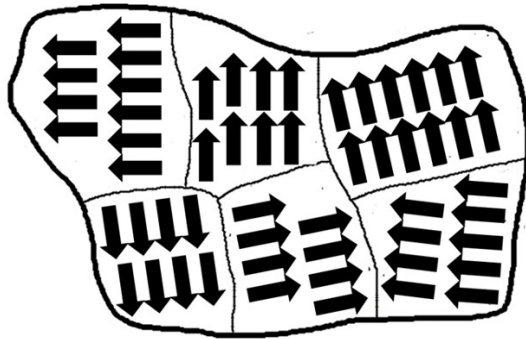


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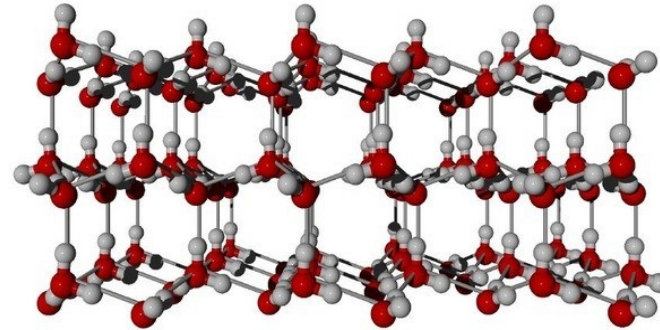
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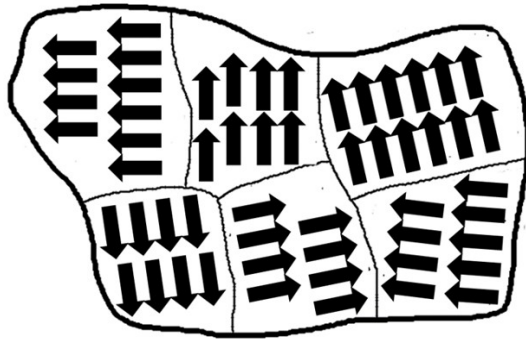


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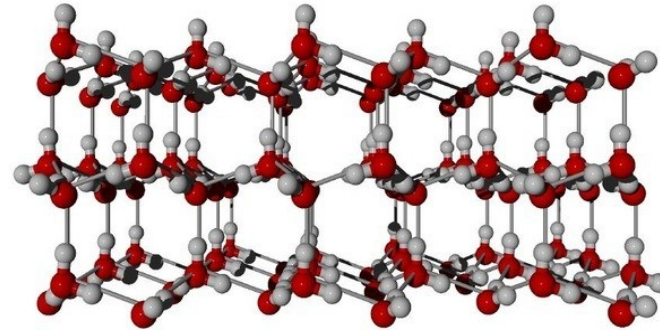
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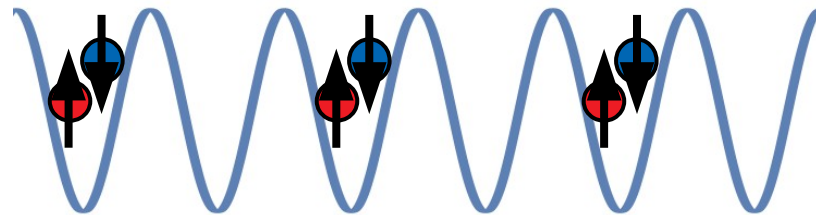
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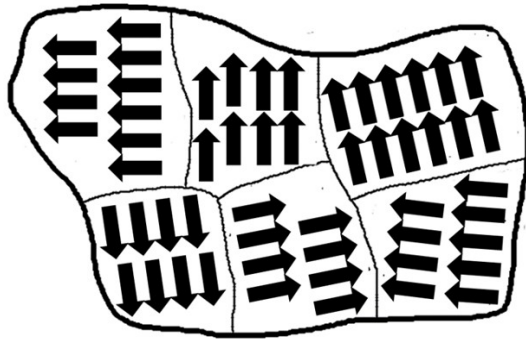
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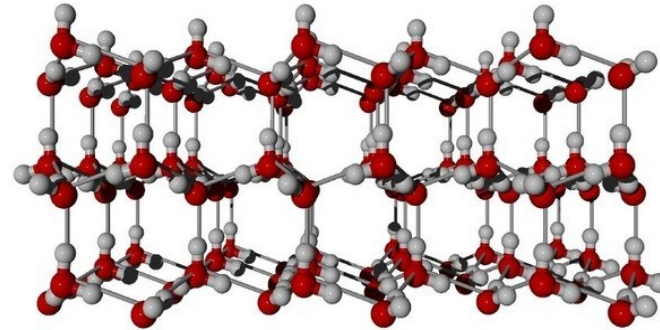
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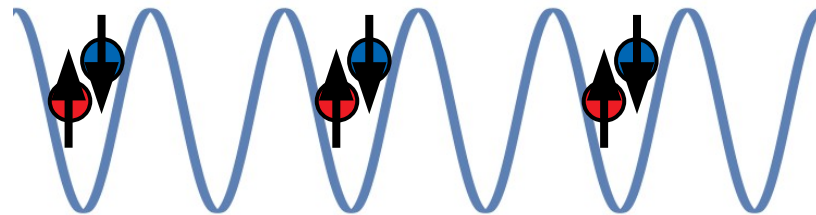
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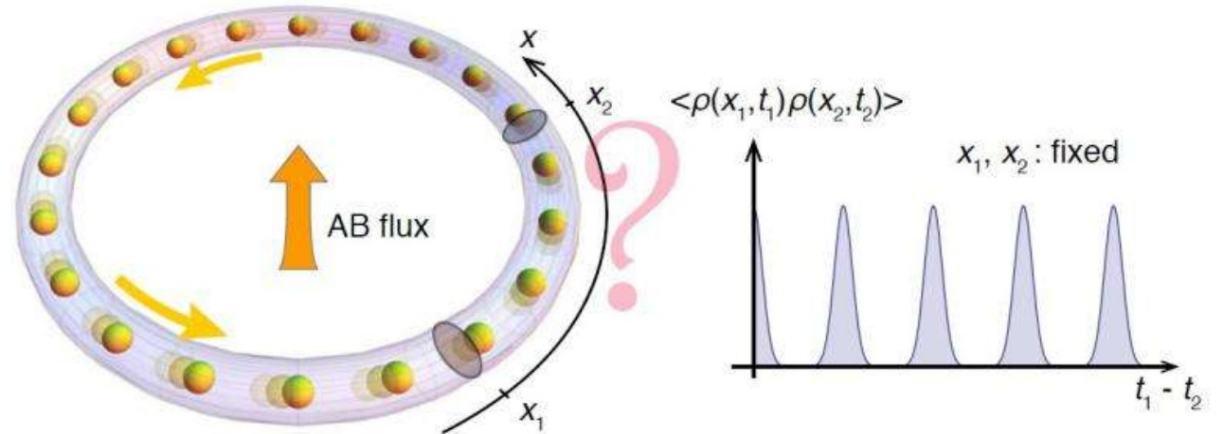
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- Can **time-translational invariance** be spontaneously broken?

(F. Wilczek, PRL. 109, 160401 (2012))

Time translation symmetry breaking

ring particle model: **ground states**
with moving “lumps”



$$\lim_{V \rightarrow \infty} \langle \rho(x, t) \rho(x', t') \rangle \longrightarrow f(t, t')$$
$$|x - x'| \rightarrow \infty$$

$\rho(x, t) =$ local order parameter

Time translation symmetry breaking

- intense discussion:

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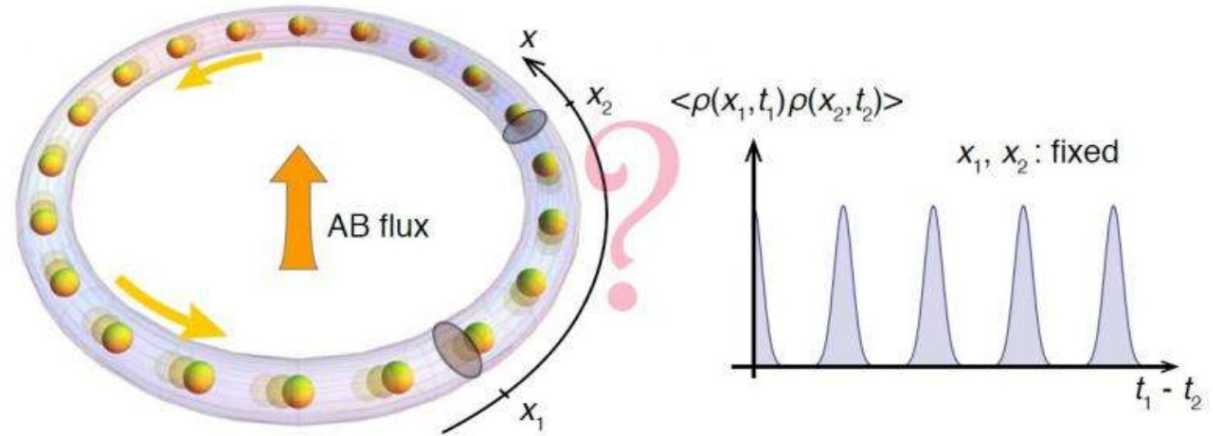
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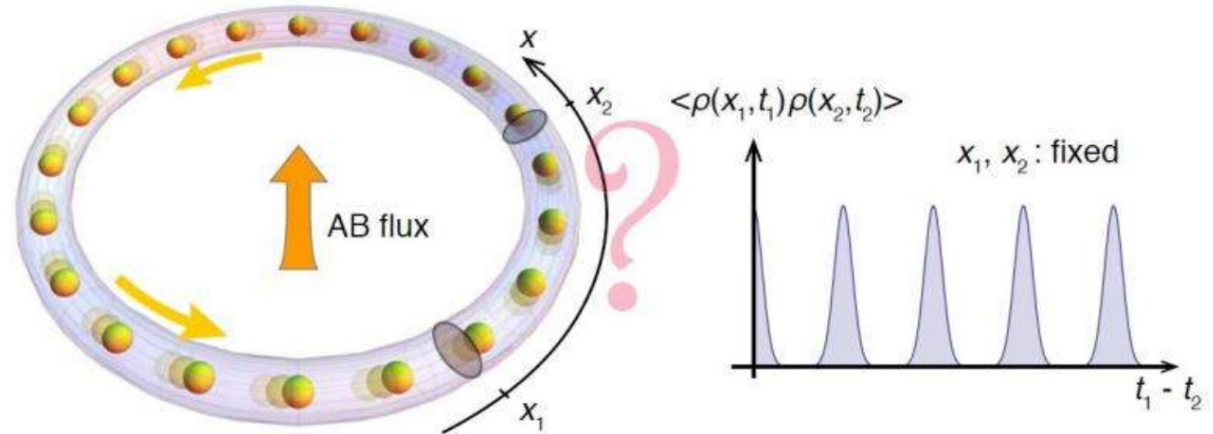
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- No-go theorem: systems in thermal equilibrium cannot manifest any time-crystalline behavior (*short-range Hamiltonians)

(H. Watanabe and M. Oshikawa, PRL 114, 251603 (2015))

- right context should be out of equilibrium: *e.g., preparing the system in an excited state, driven dynamics, ...*

- Important steps on periodically driven systems (**Floquet dynamics**):

$$\hat{H}(t + T) = \hat{H}(t) \quad T = \text{period}$$

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III) non-trivial oscillations persist for infinitely long times **only** the thermodynamic limit;

Floquet time-crystals

[D. V. Else, Bela Bauer, and Chetan Nayak, *PRL* 117, 090402 (2016)]

Combining **Floquet dynamics** + **MBL** (many-body localization) can support a **DTC**:

$$\hat{H}(t) = \hat{H}_{MBL} + \hat{X} \sum_{n=0}^{\infty} \delta(t - nT)$$

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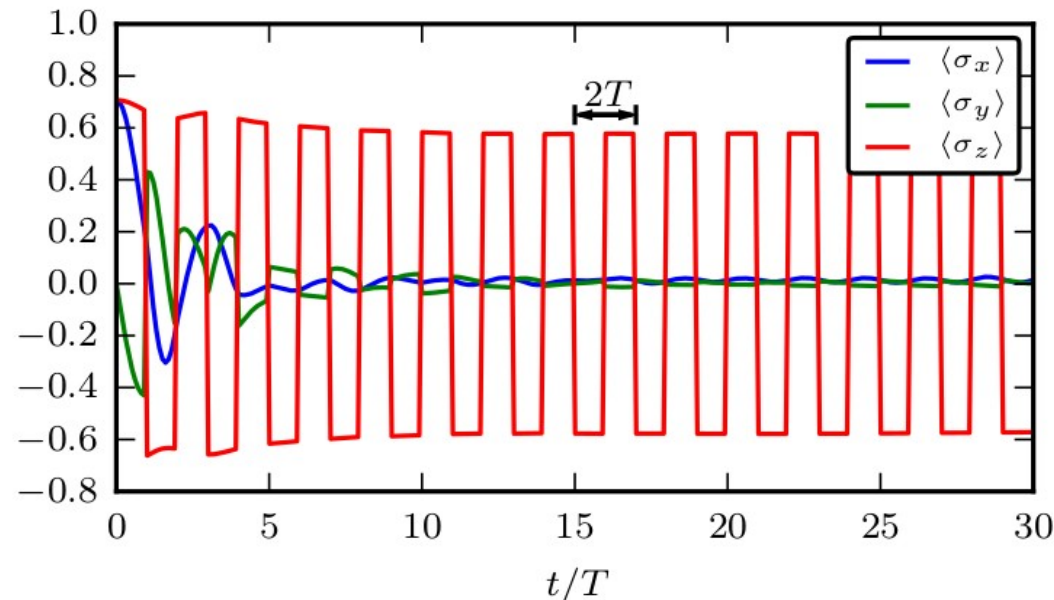
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period doubling for magnetization



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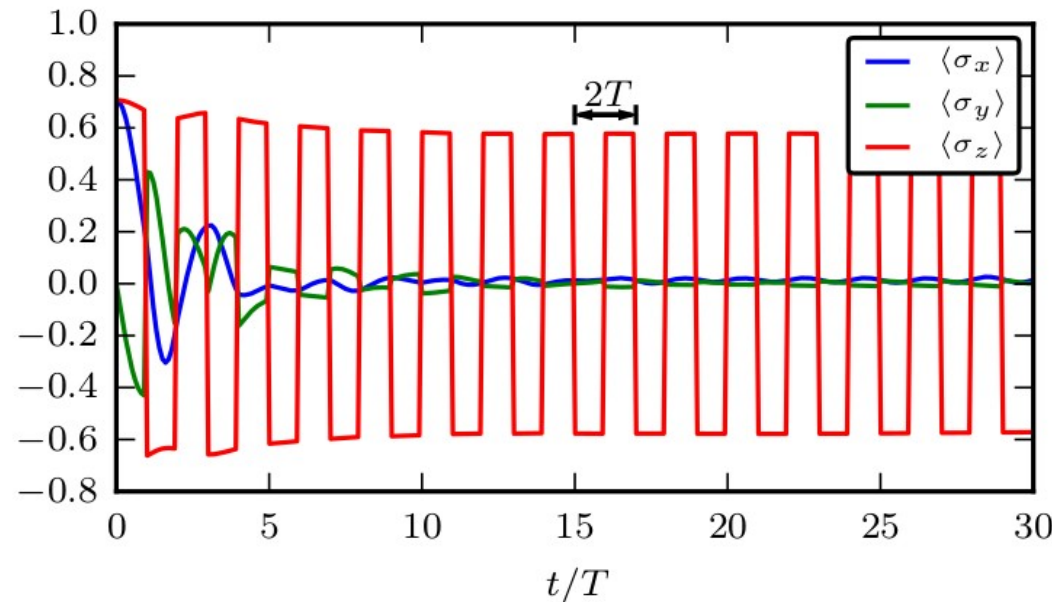
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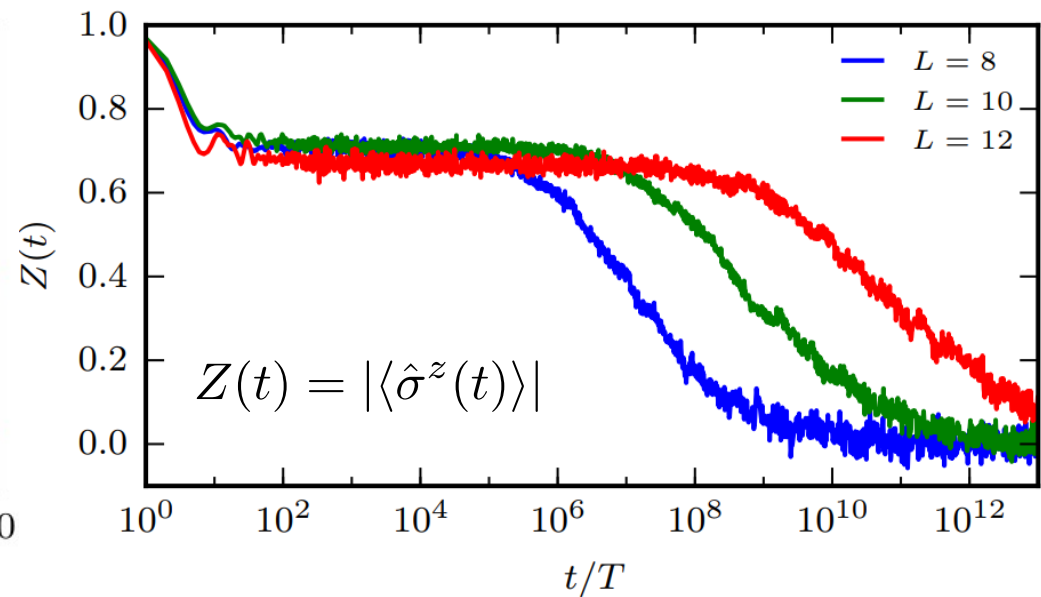
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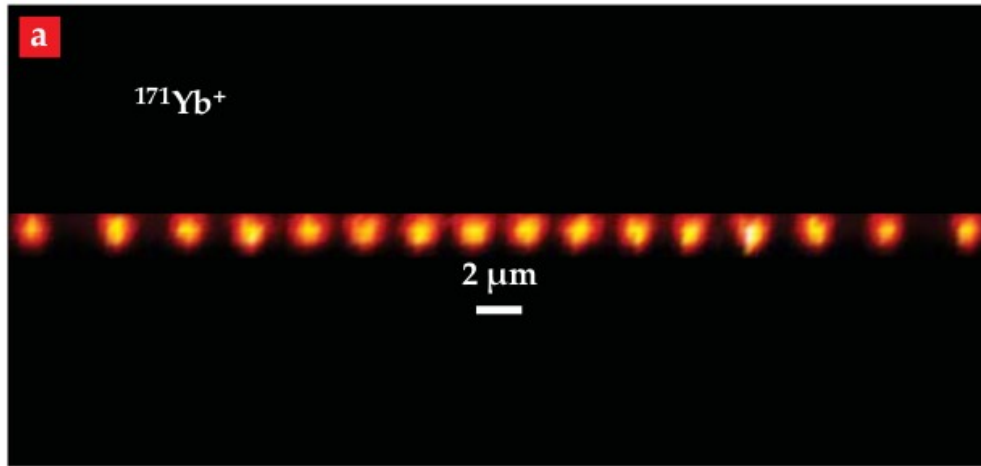


*persistence **only** in the thermodynamic limit*



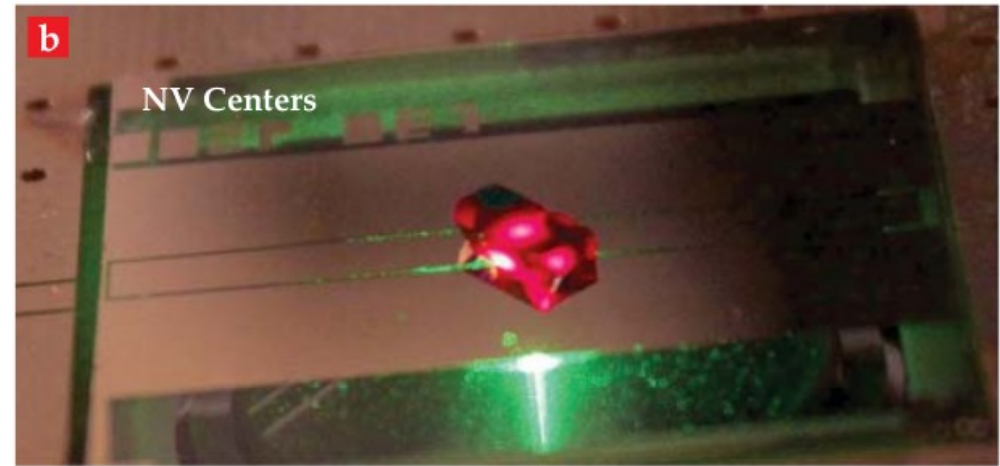
trapped ytterbium ions

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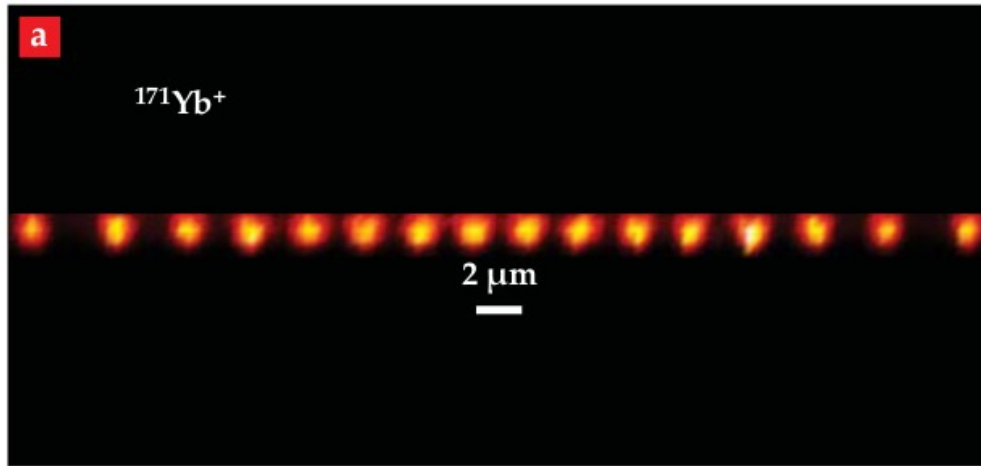
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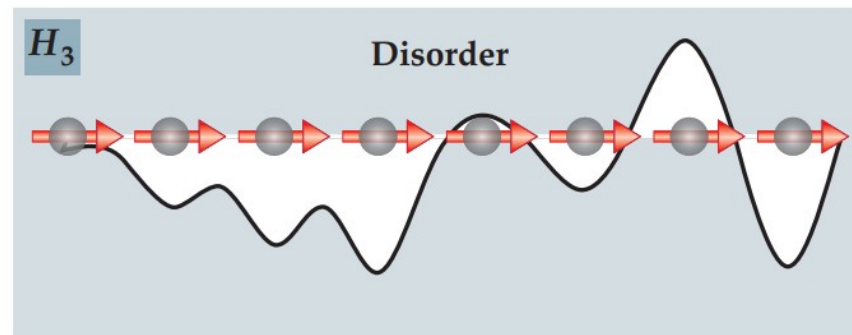
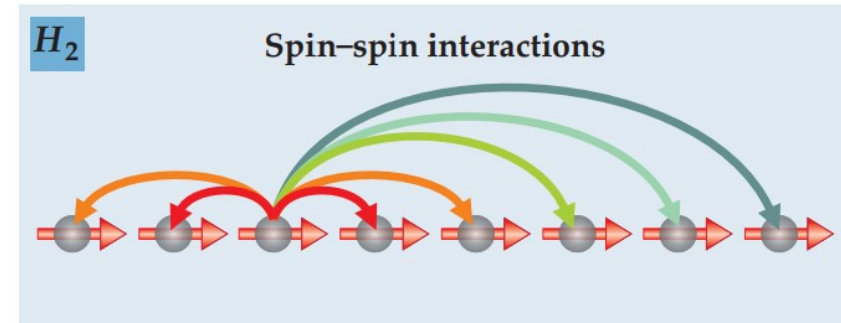
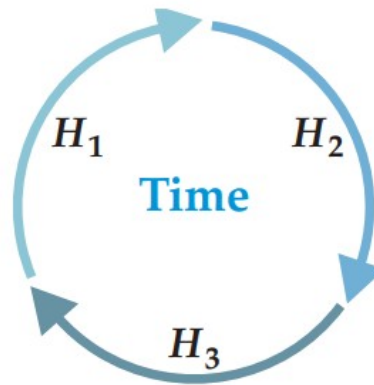
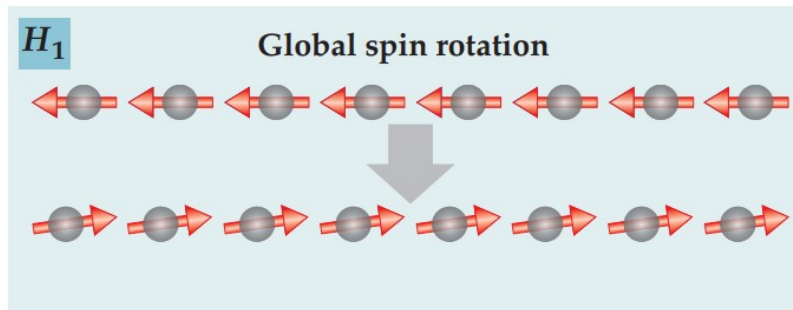
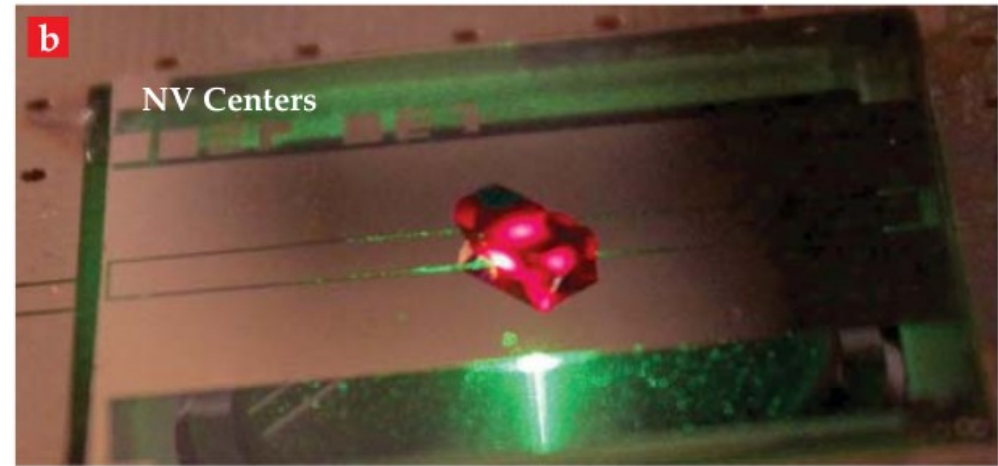
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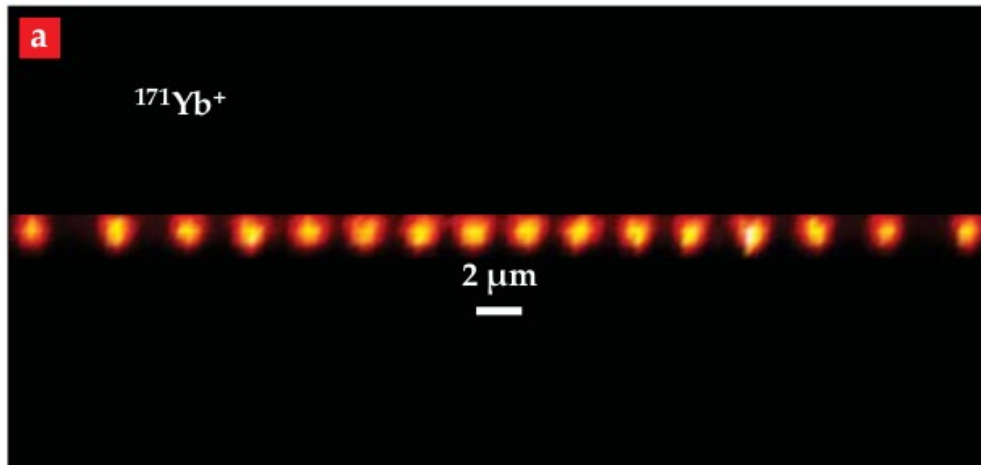
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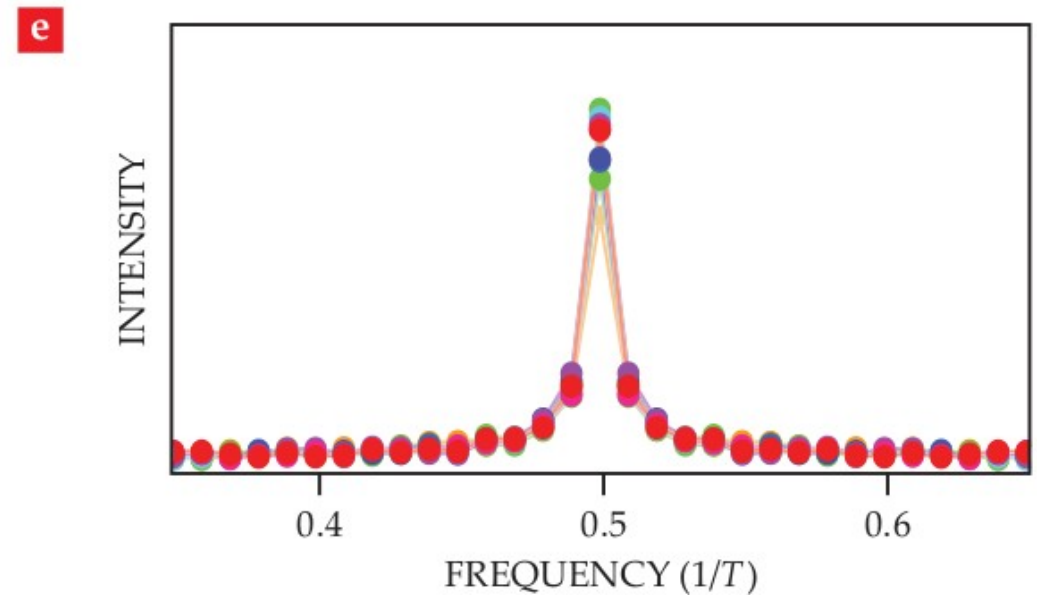
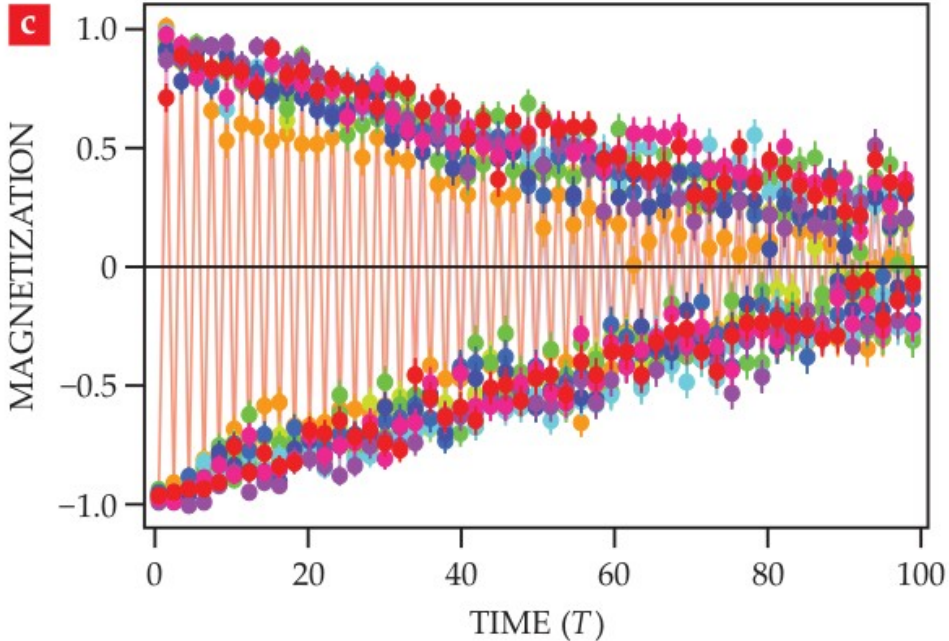
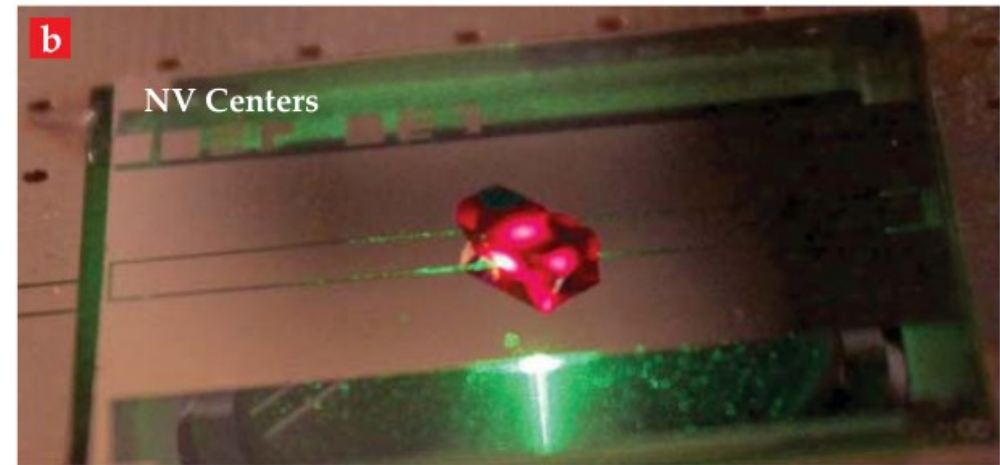
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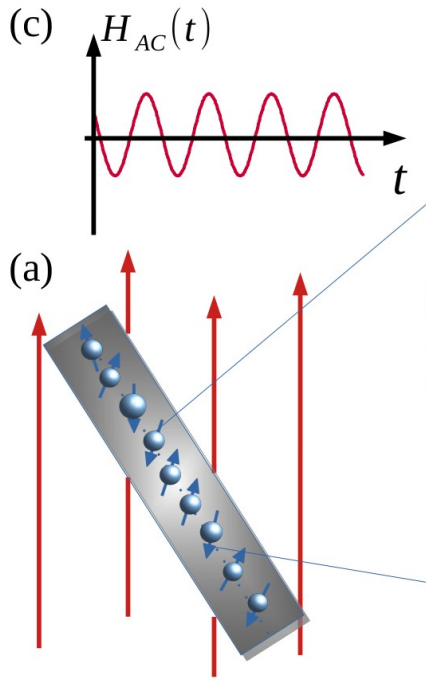
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Sensing AC fields

$$\hat{H}(t) = \hat{H}_0(t) + \hat{H}_{AC}(t)$$

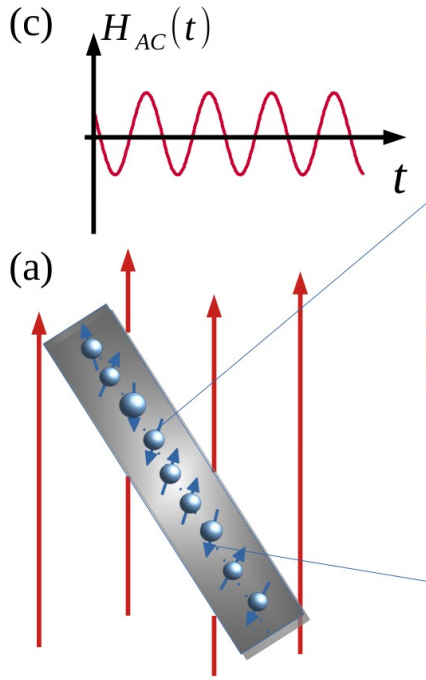
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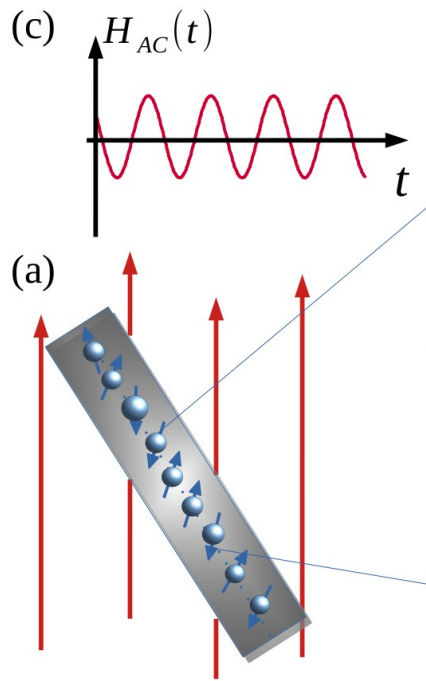
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improve performance

non-eq. interacting /
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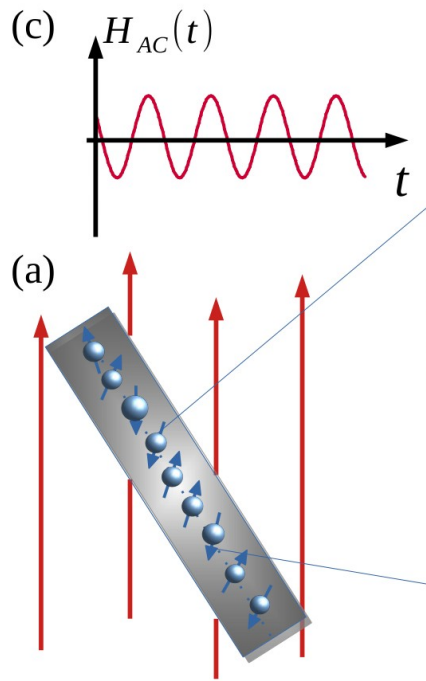
Multitude of approaches

- . Dieter Suter et al. Rev. Mod. Phys. 88 (2016).
- . Hengyun Zhou et al., PRX 10, 031003 (2020).
- . Soonwon Choi et al. ArXiv:1801.00042 (2017)
- . Utkarsh Mishra and Abolfazl Bayat, Scientific Reports 12, 14760 (2022)

- dynamical decoupling spins by a series of fast pulses – mitigate decoherence/interactions:
- strong interactions and fast drive: prepare and (prethermal) stabilize of GHZ-like states:
- “inspirational work”: non-equilibrium dynamics of an integrable (Ising chain) many-body system

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exploit long-range spatial and time ordering for enhanced metrological protocols.

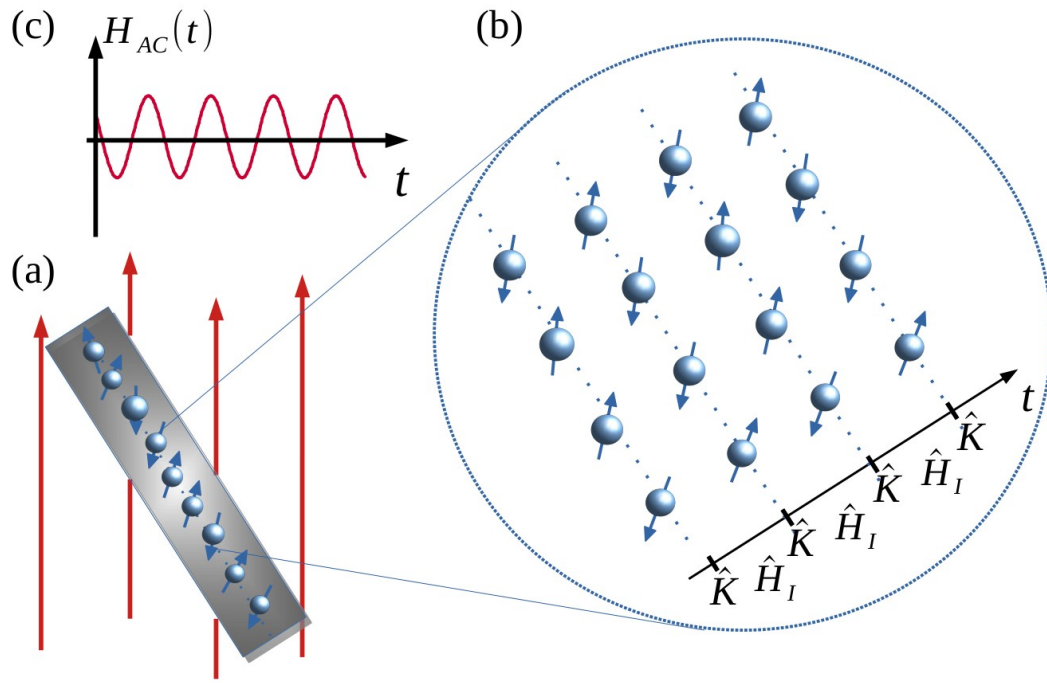
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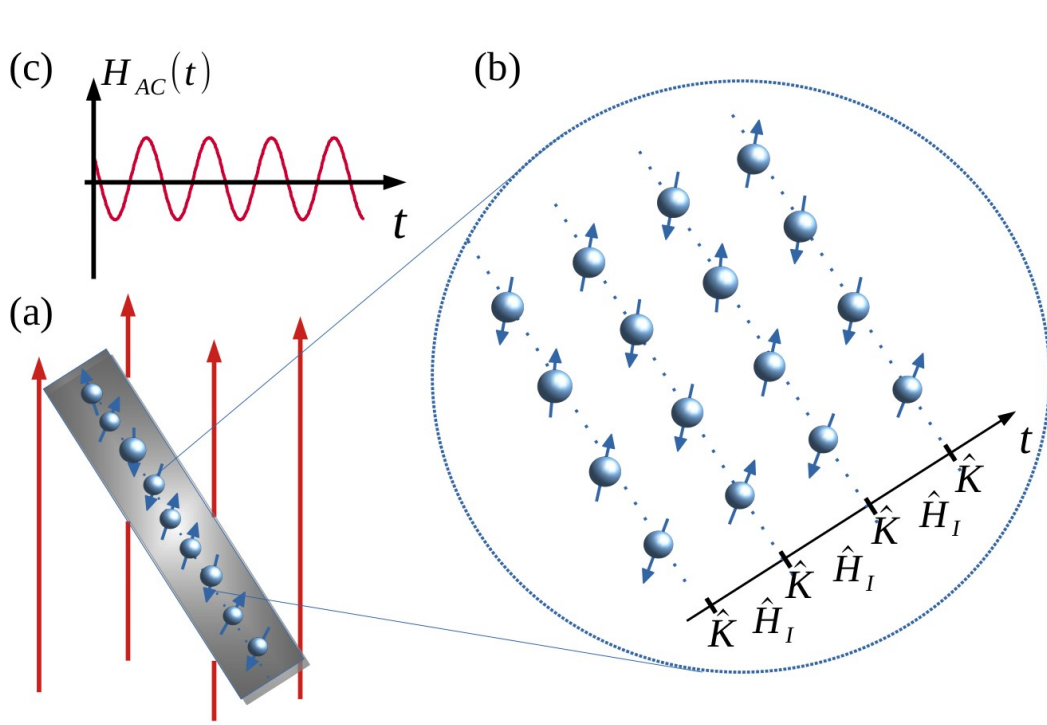
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Sensing steps:

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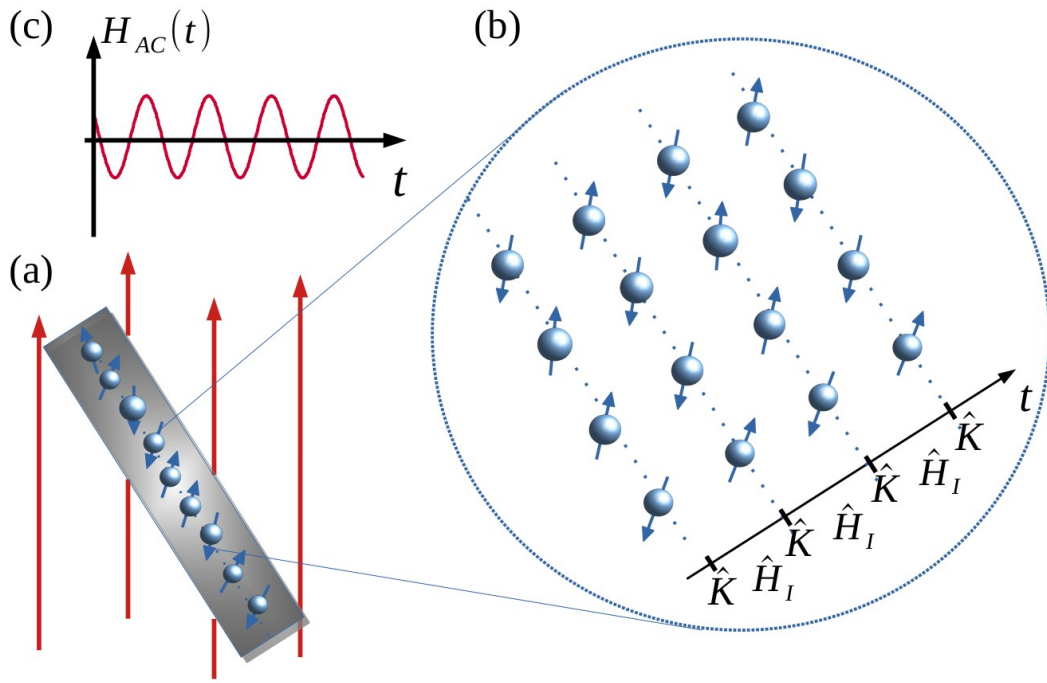
ii) interact with signal;

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iii) measurement on sensor state:

$$p_j(t) = \text{Tr}(\hat{E}_j \hat{\rho}(t, h_{ac}))$$

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Least uncertainty/maximum precision:

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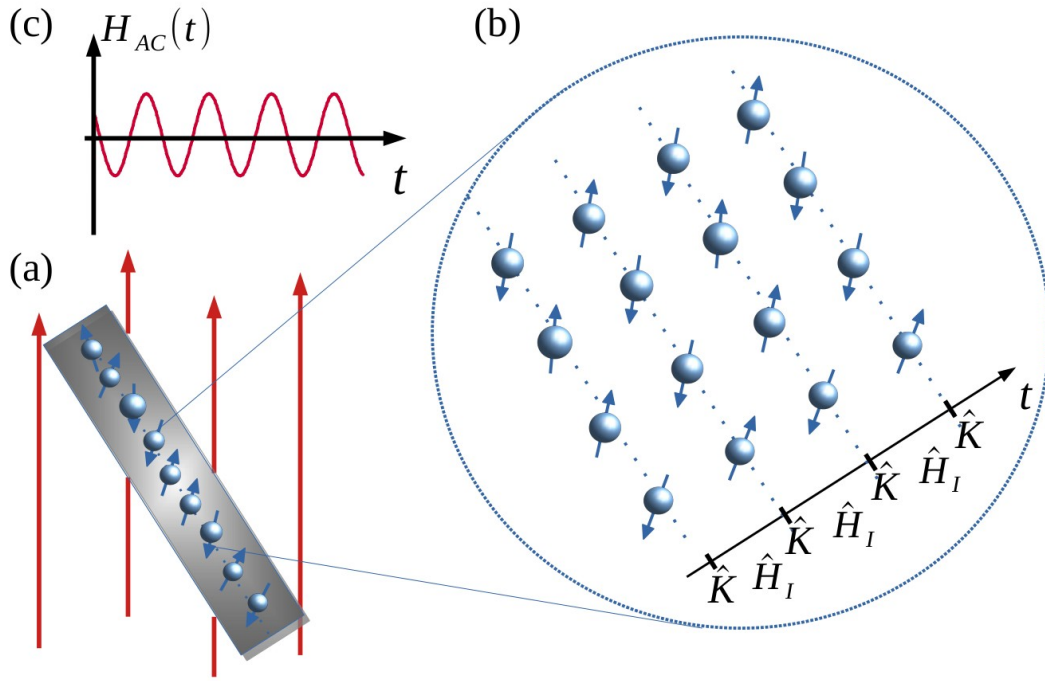
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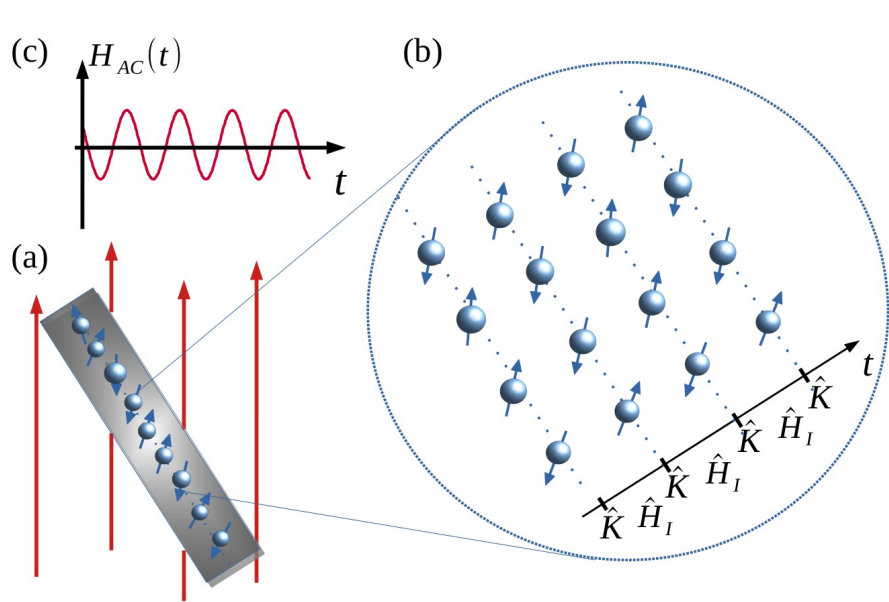
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$$F_{h_{ac}}(t)/4 = \langle \partial_{h_{ac}} \psi(t, h_{ac}) | \partial_{h_{ac}} \psi(t, h_{ac}) \rangle - |\langle \psi(t, h_{ac}) | \partial_{h_{ac}} \psi(t, h_{ac}) \rangle|^2$$

(quantum Fisher information)

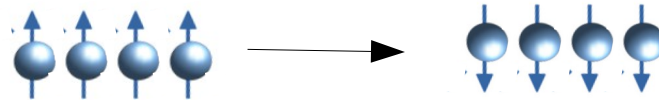
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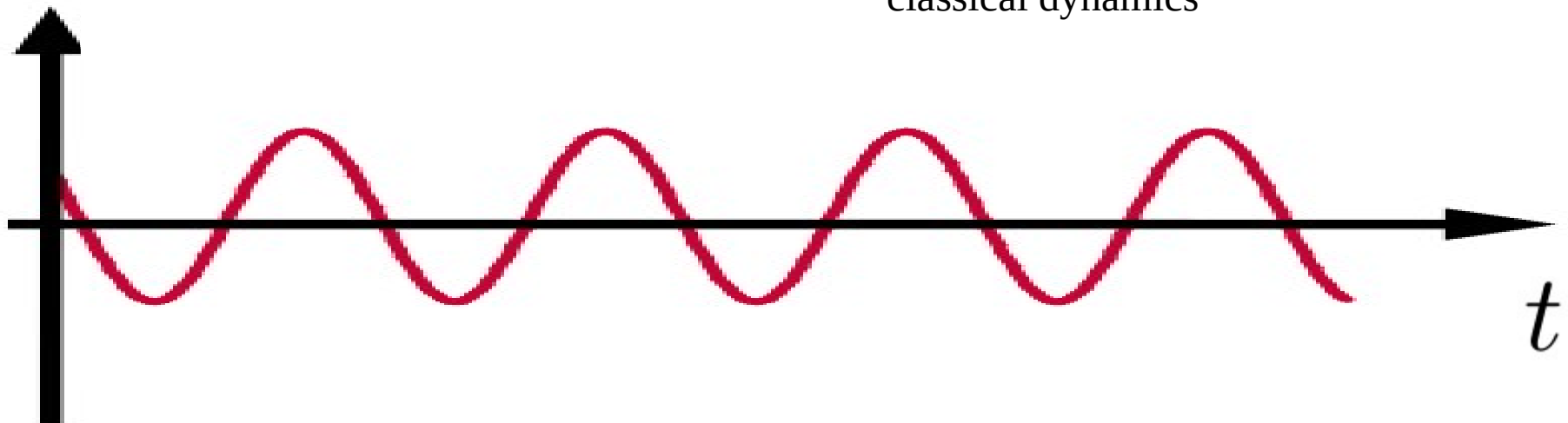
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Case with $\phi = \pi$, $h_i^x = 0$, $h_{ac} \rightarrow 0$

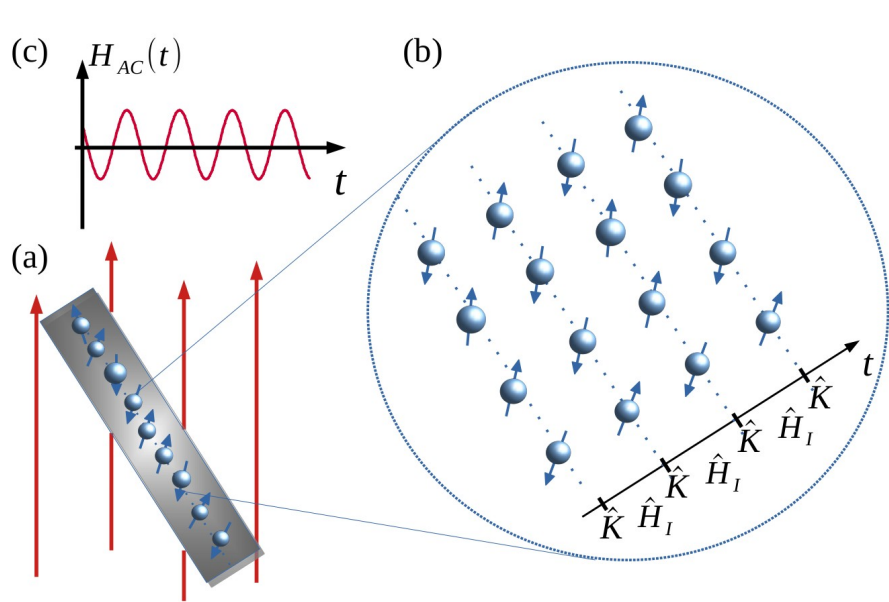


'classical dynamics'

$\hat{H}_{AC}(t)$



Floquet time-crystals as sensors of AC fields

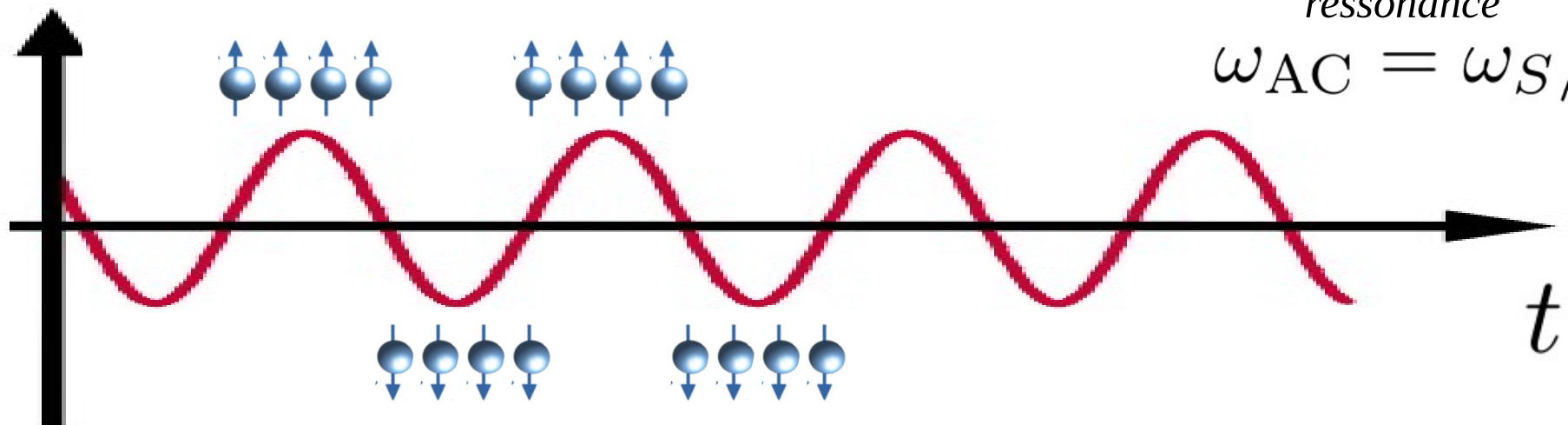


$$\hat{H}_0(t) = \sum_i \left(J_i \hat{\sigma}_i^z \hat{\sigma}_{i+1}^z + \sum_{\alpha=x,z} h_i^\alpha \hat{\sigma}_i^\alpha \right) + \sum_{n=1}^{\infty} \delta(t - nT_S) \phi \hat{S}^x,$$

$$\hat{H}_{AC}(t) = h_{AC} \sin(\omega_{AC} t + \theta_{AC}) \hat{S}_z$$

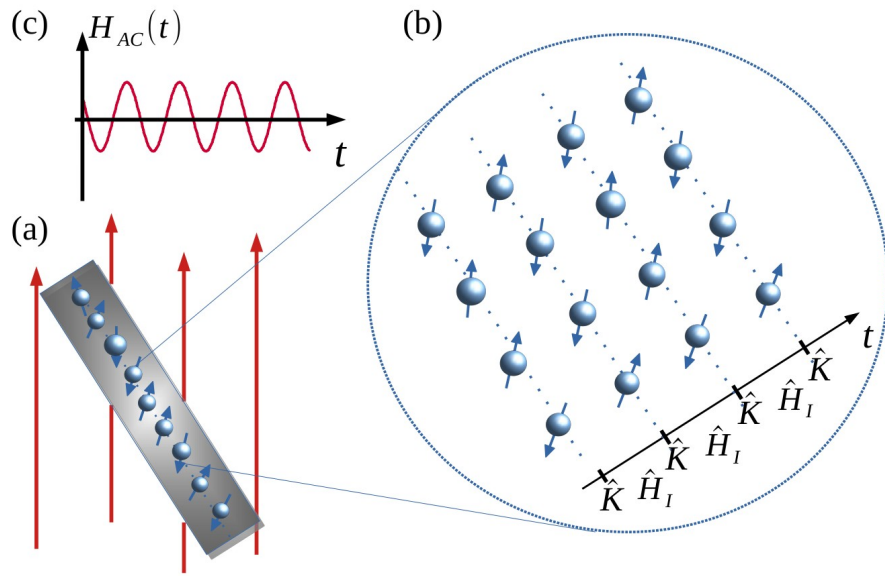
Case with $\phi = \pi$, $h_i^x = 0$, $h_{ac} \rightarrow 0$

$\hat{H}_{AC}(t)$



“period doubling
ressonance”
 $\omega_{AC} = \omega_S/2$

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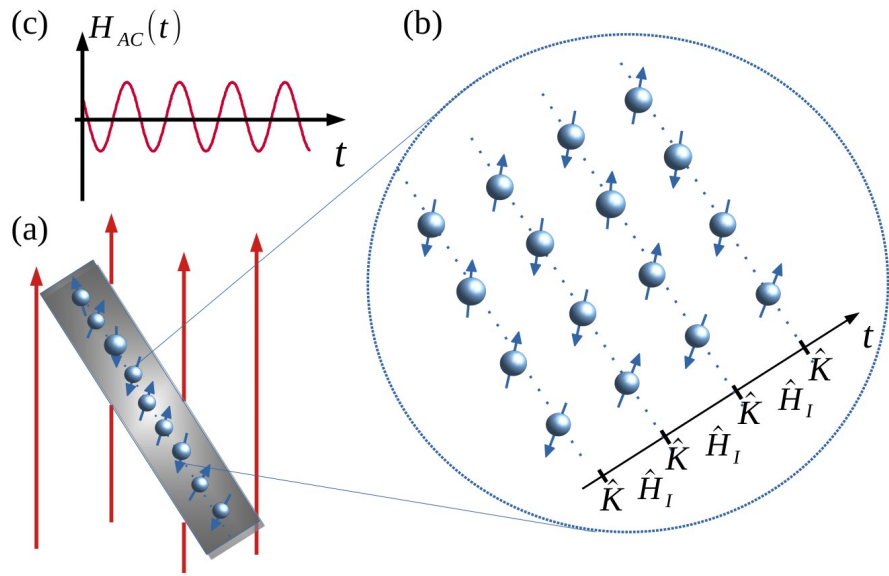
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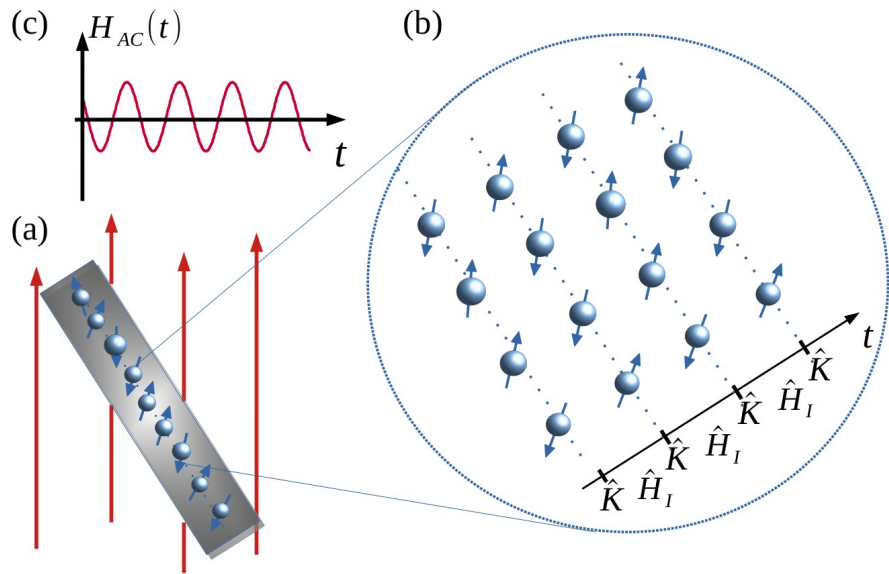
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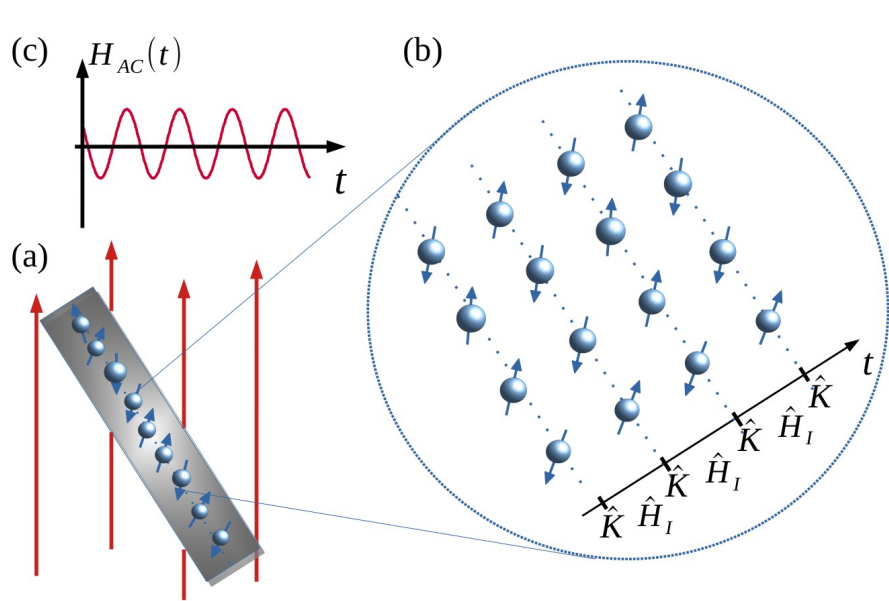
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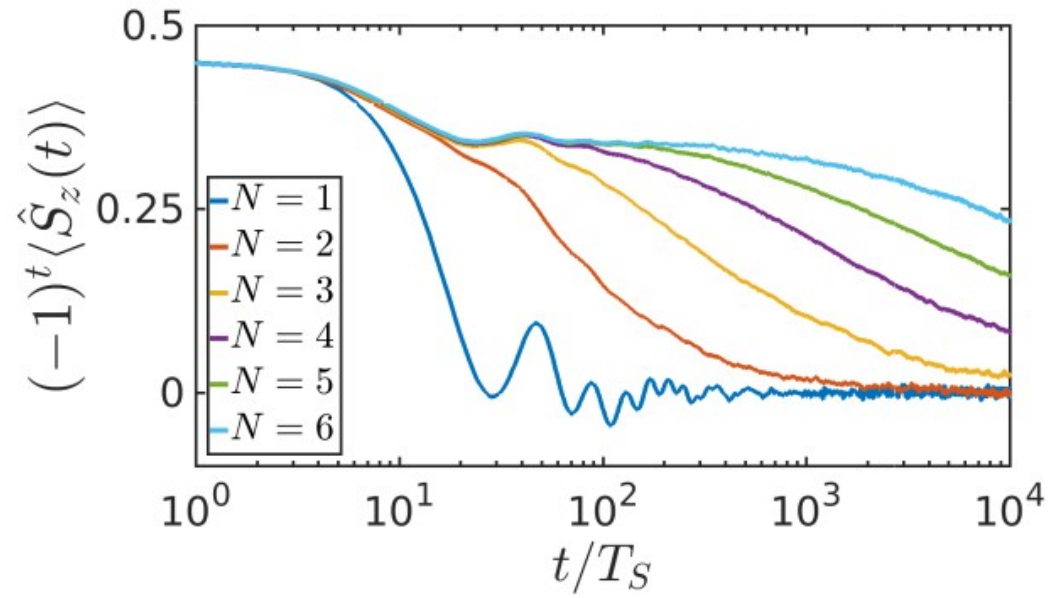
$$\varphi|_{\omega_S=2\omega_{AC}+\epsilon} : t^* \sim \theta_{AC}/\epsilon$$

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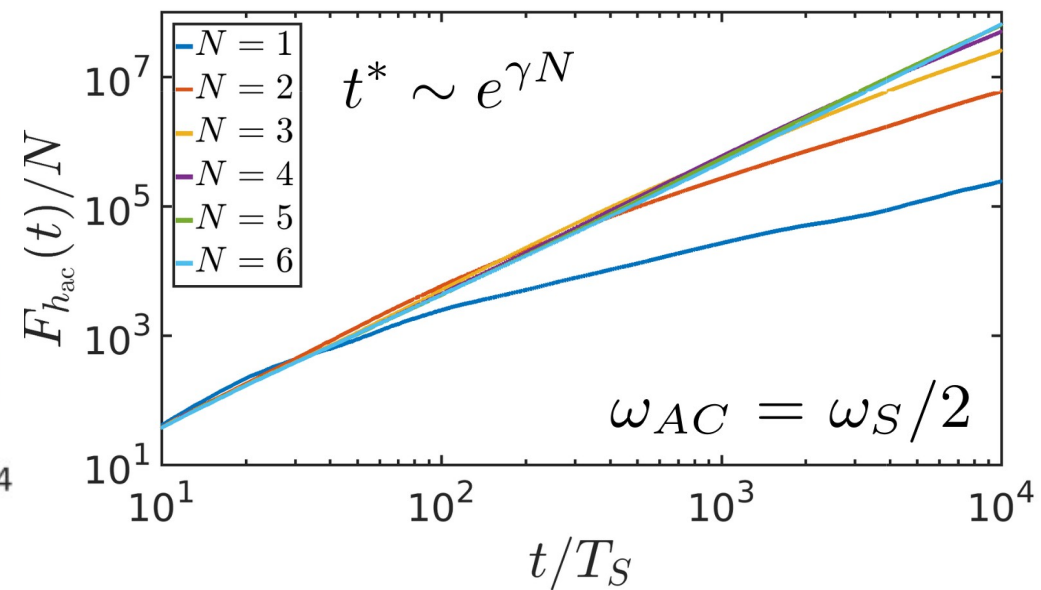
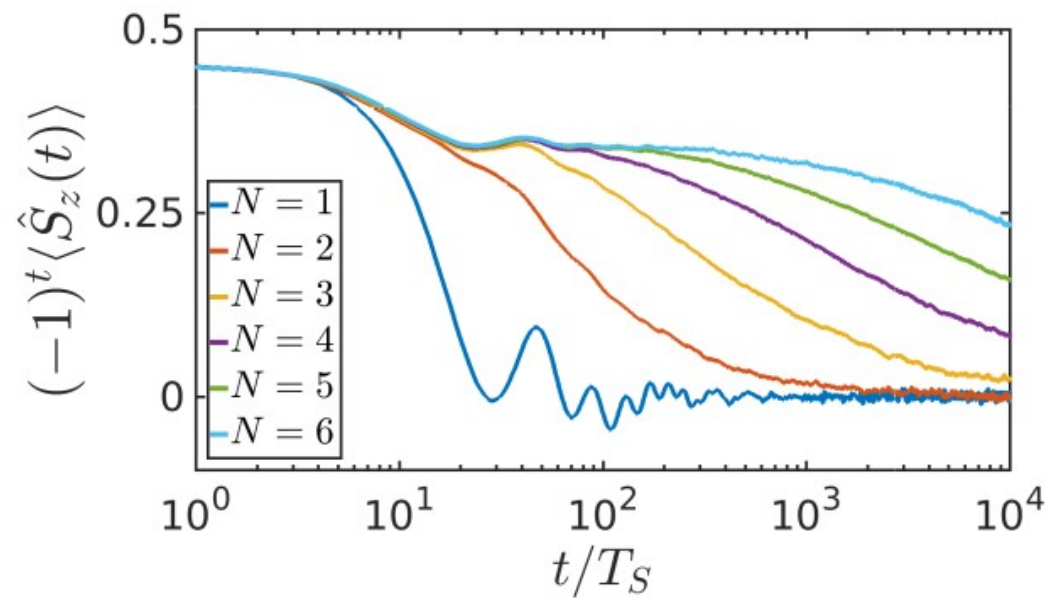
(iii) off-resonance time scale for the optimal growth of Fisher information.

General case: $\phi \neq \pi$ $h_i^x \neq 0$
(linear response $h_{AC} \rightarrow 0$)



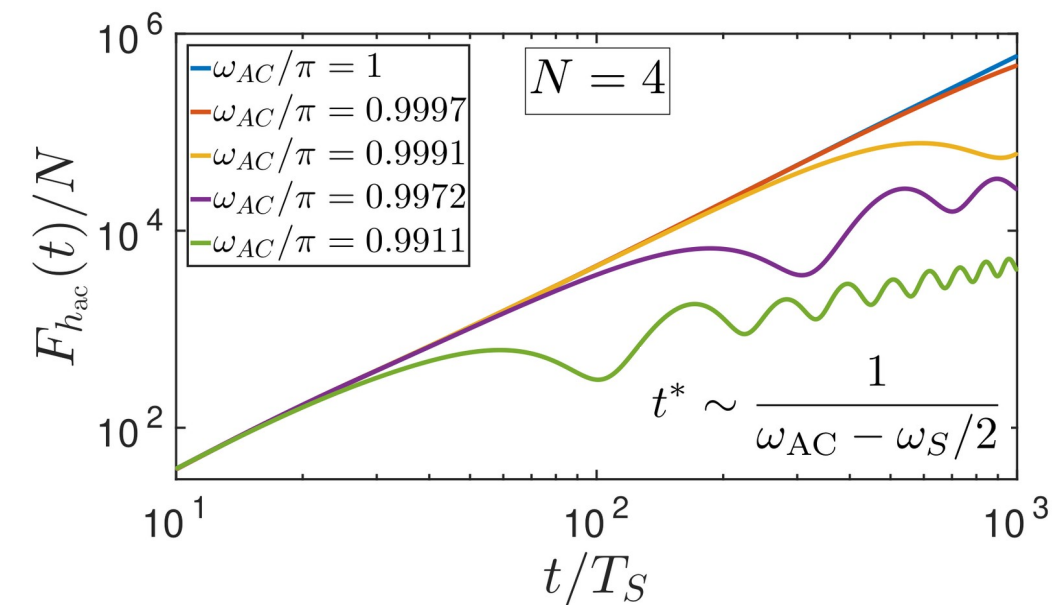
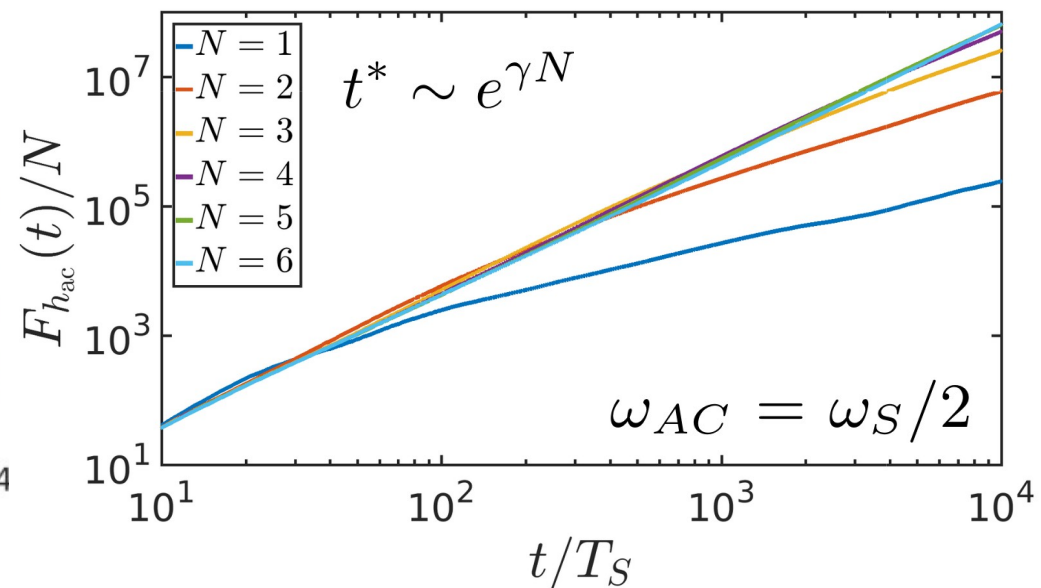
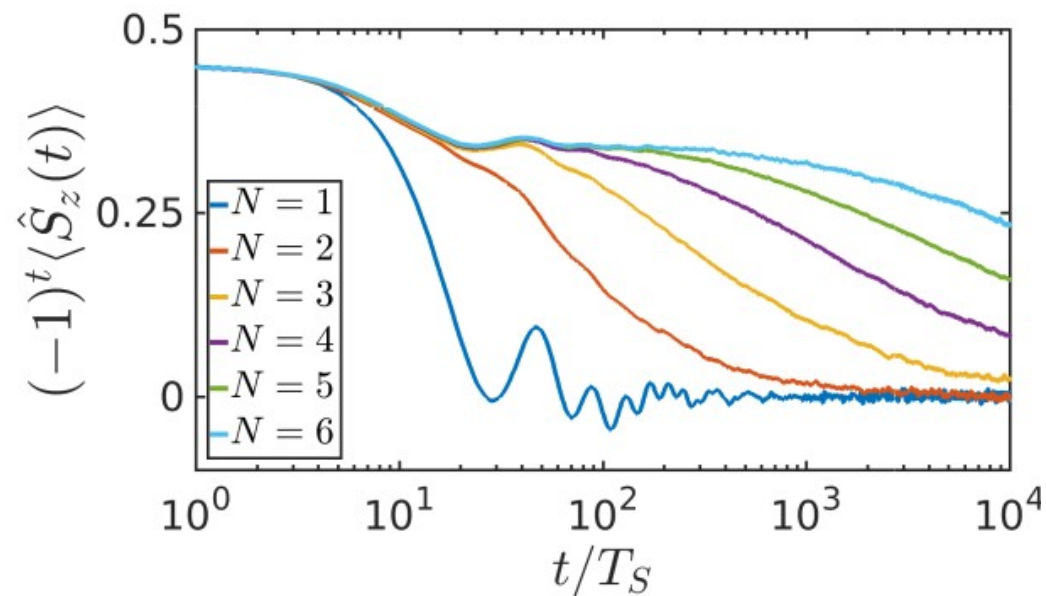
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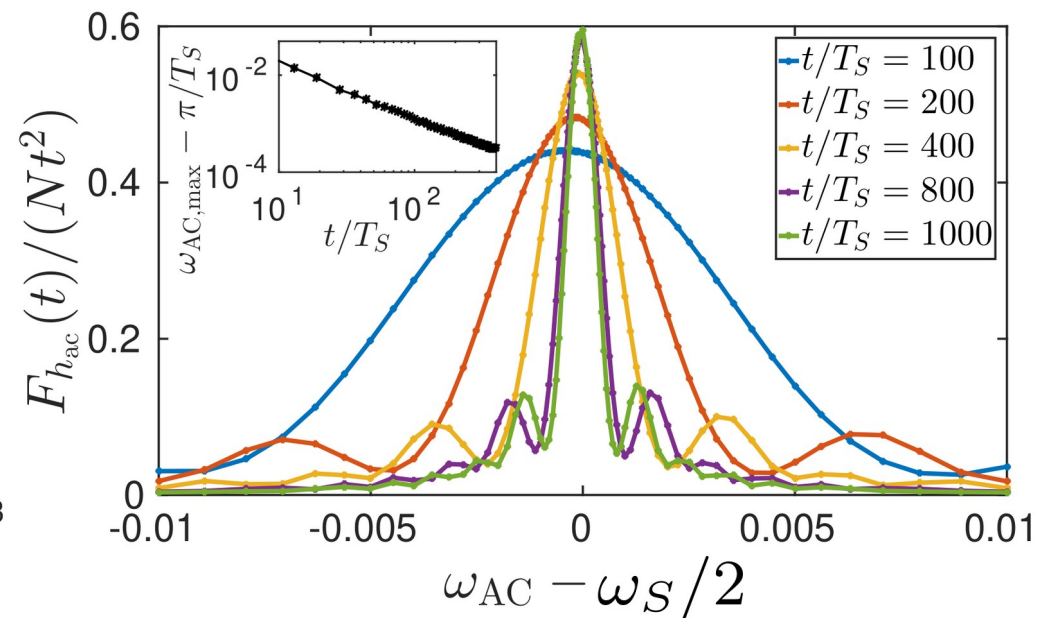
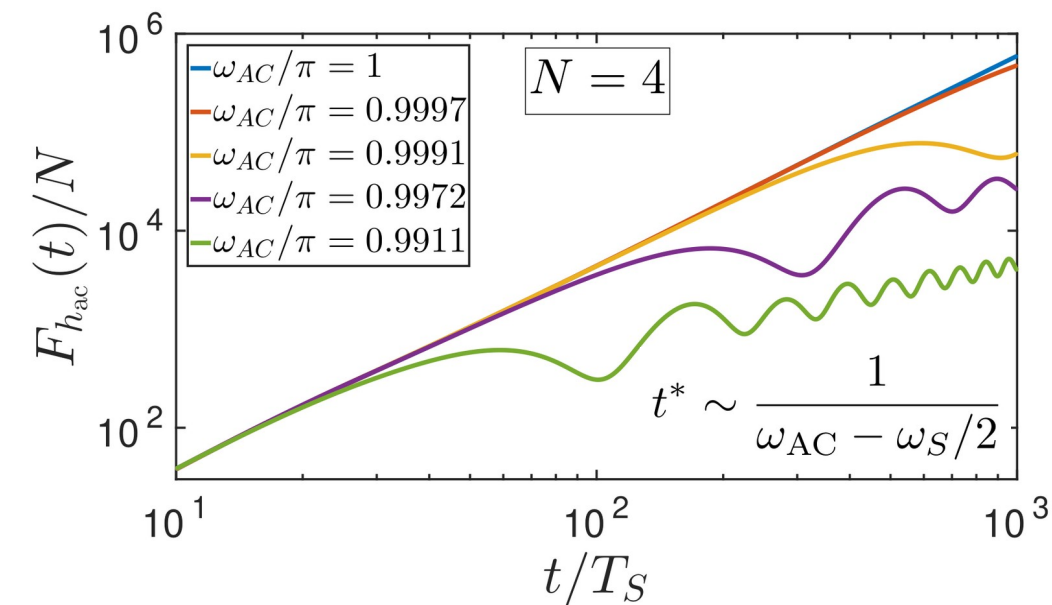
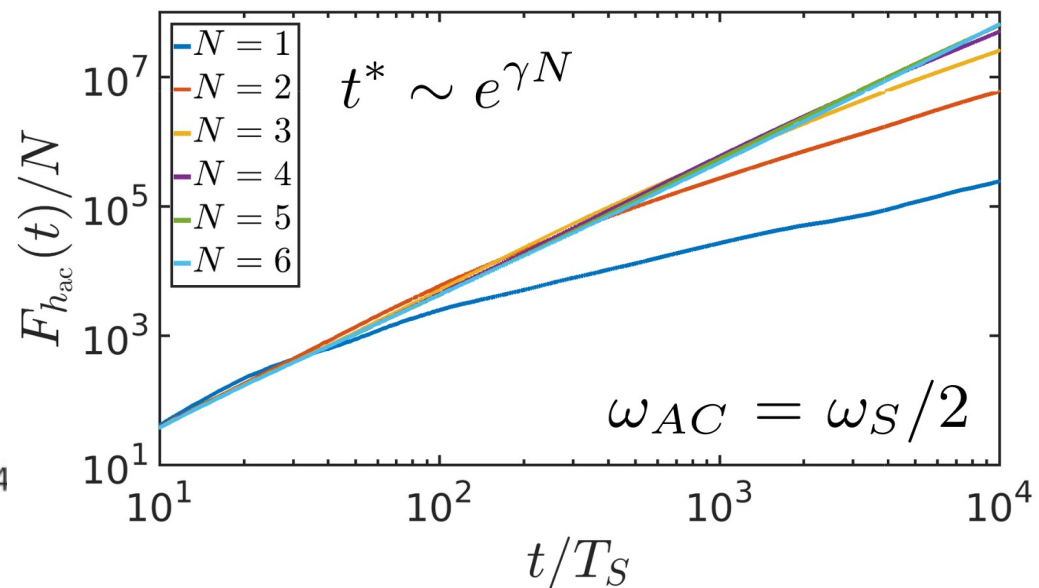
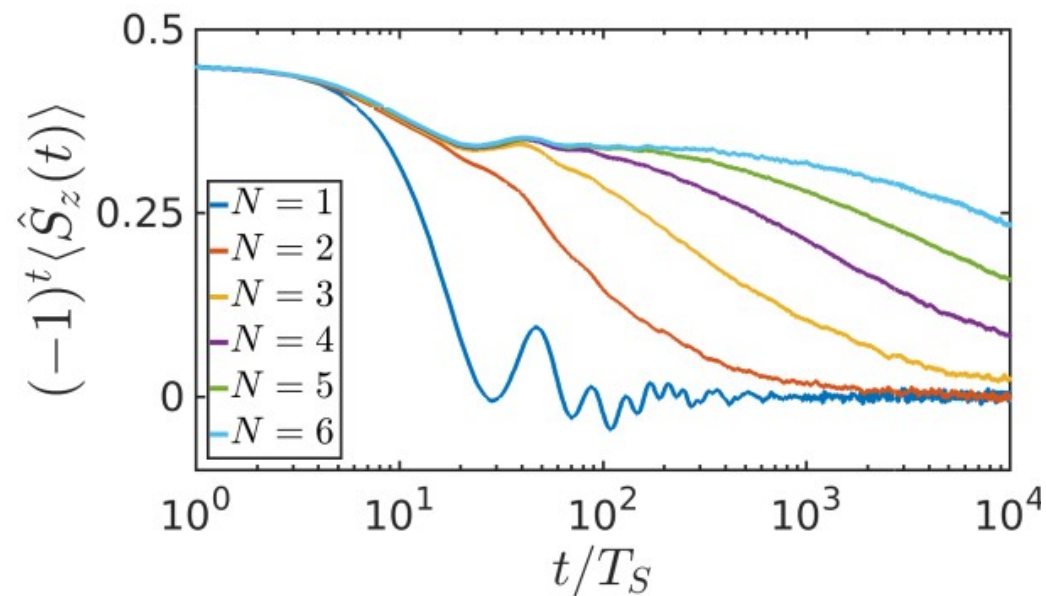
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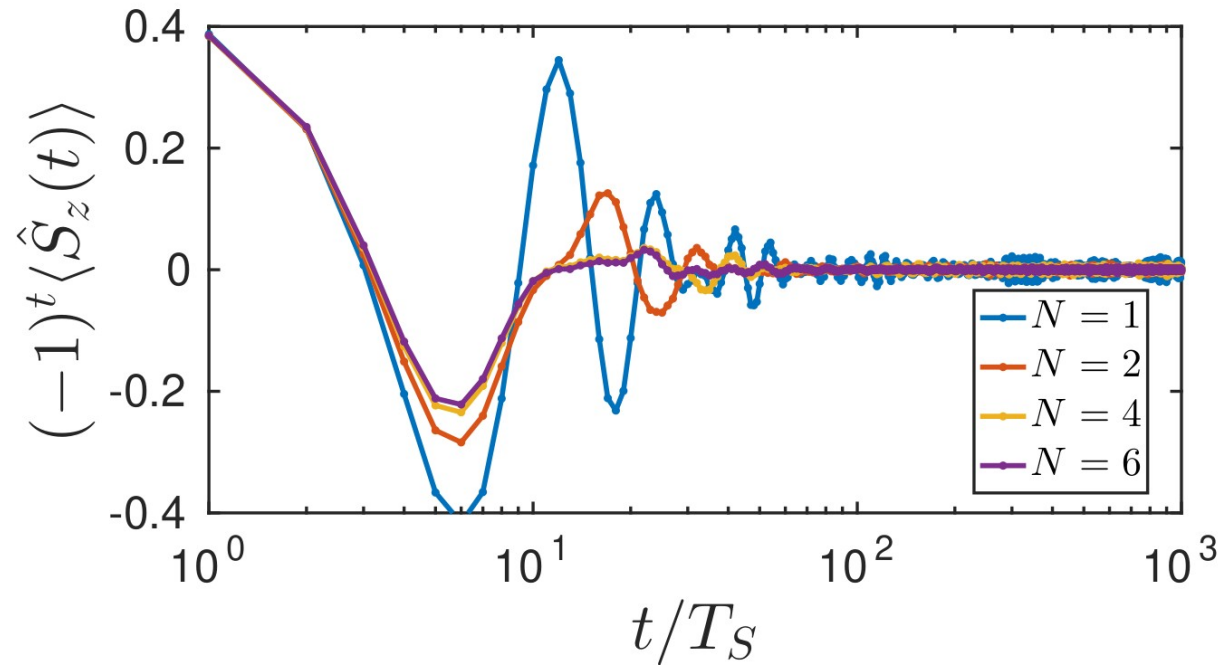
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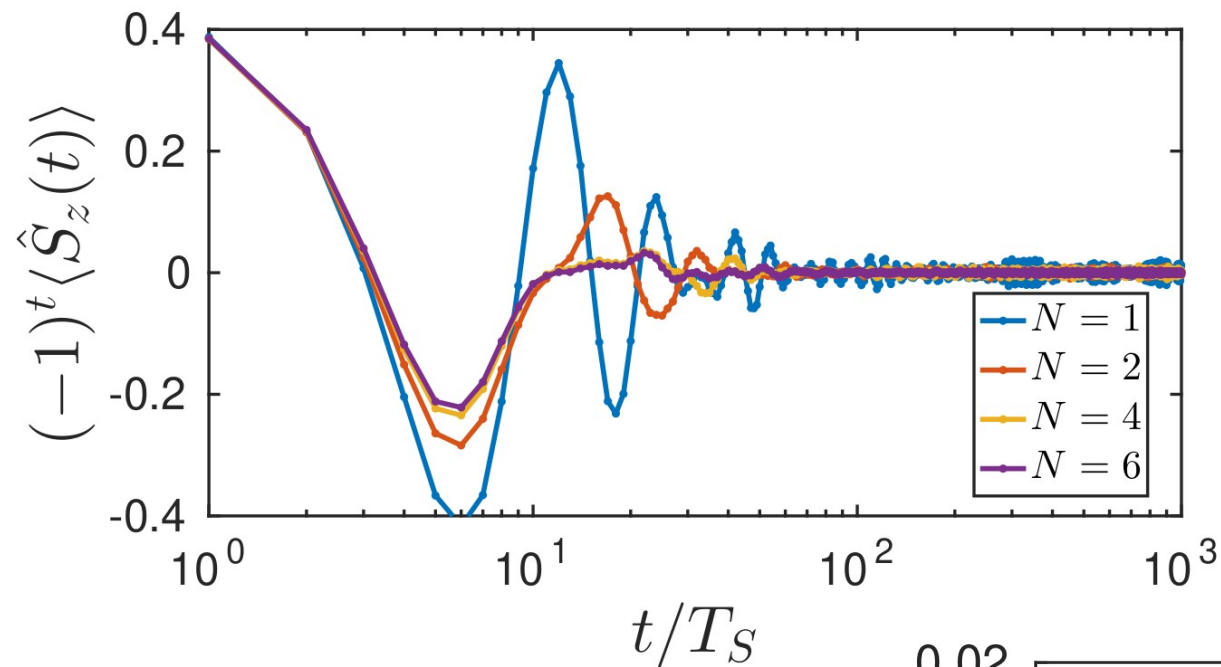
$$\max_{\omega} F_{h_{ac}}(t) \sim t^{\beta}, \beta > 2$$

Ergodic Sensor



- fast thermalization time $\sim O(T)$;

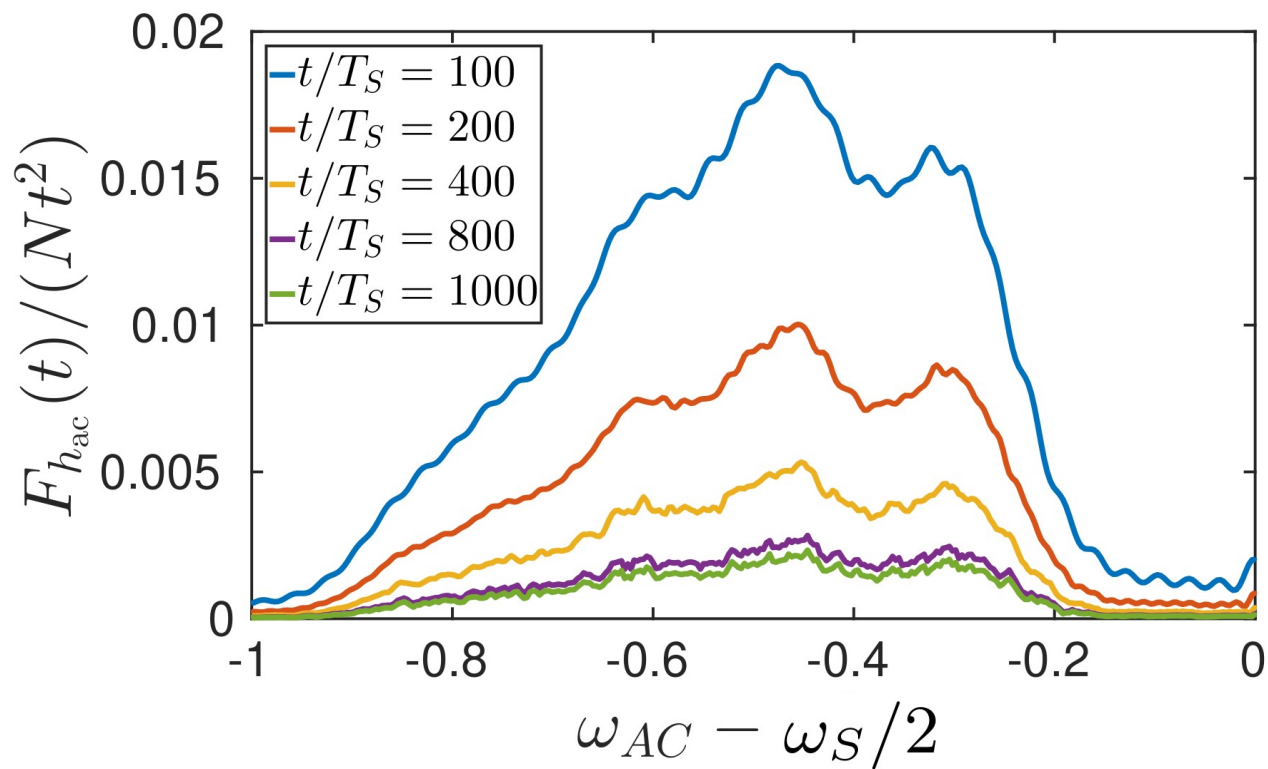
Ergodic Sensor



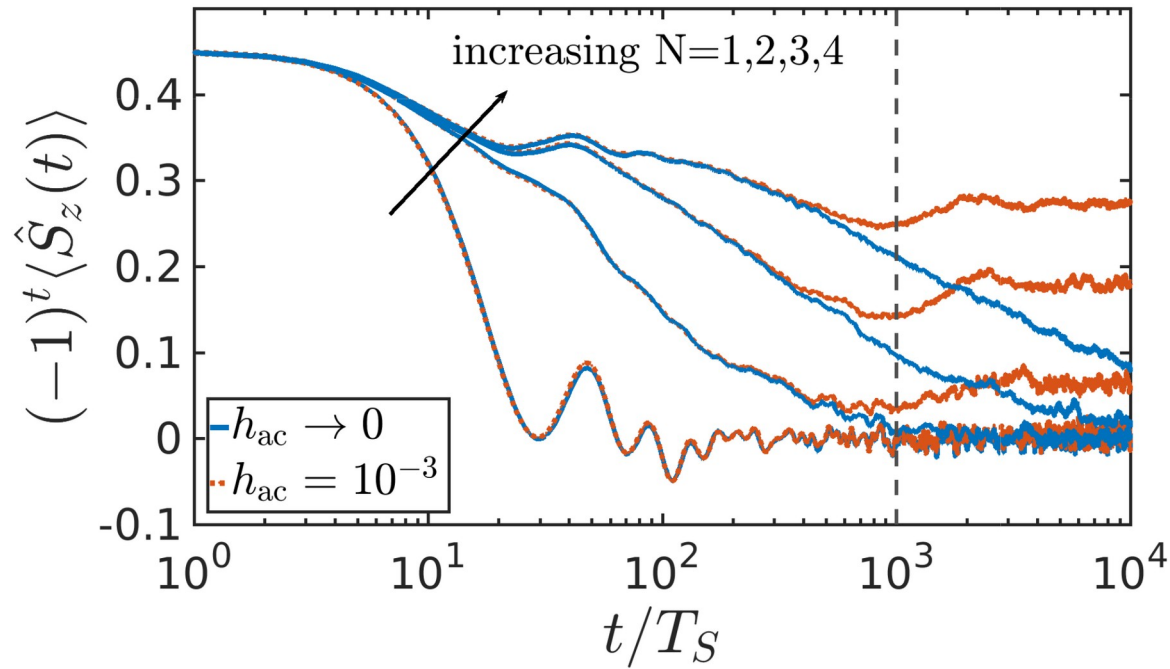
- fast thermalization time $\sim O(T)$;

- noisy response;

- below standard (classical) limit for longer times;



beyond linear response, $h_{AC} \neq 0$, $\omega_{AC} = \omega_S/2$

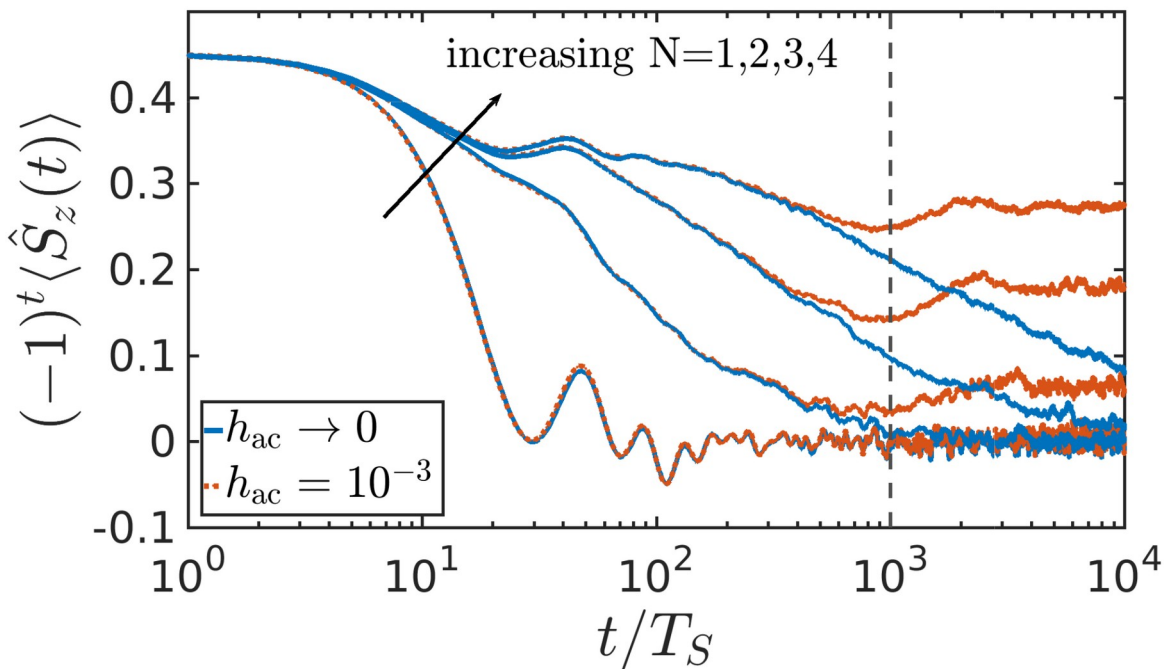


- sets a characteristic time scale,

$$t_{AC} \sim h_{ac}^{-1}$$

(dominant AC fields over dynamics)

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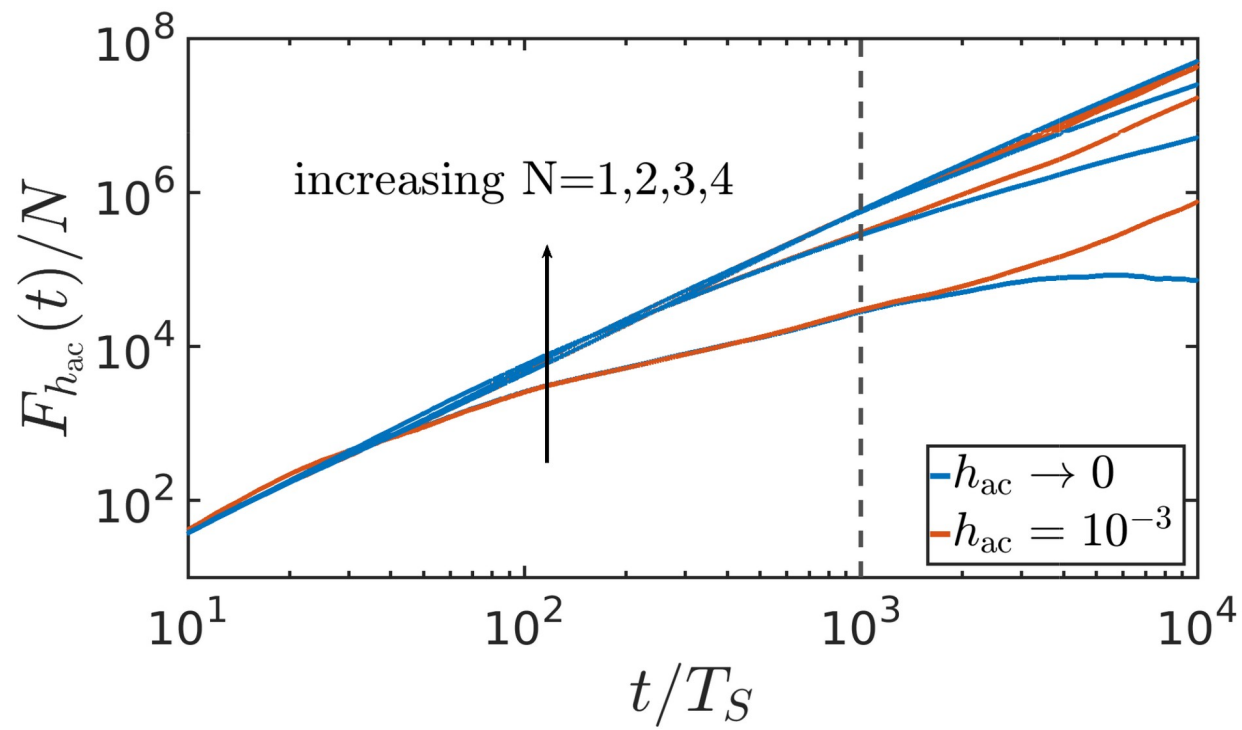
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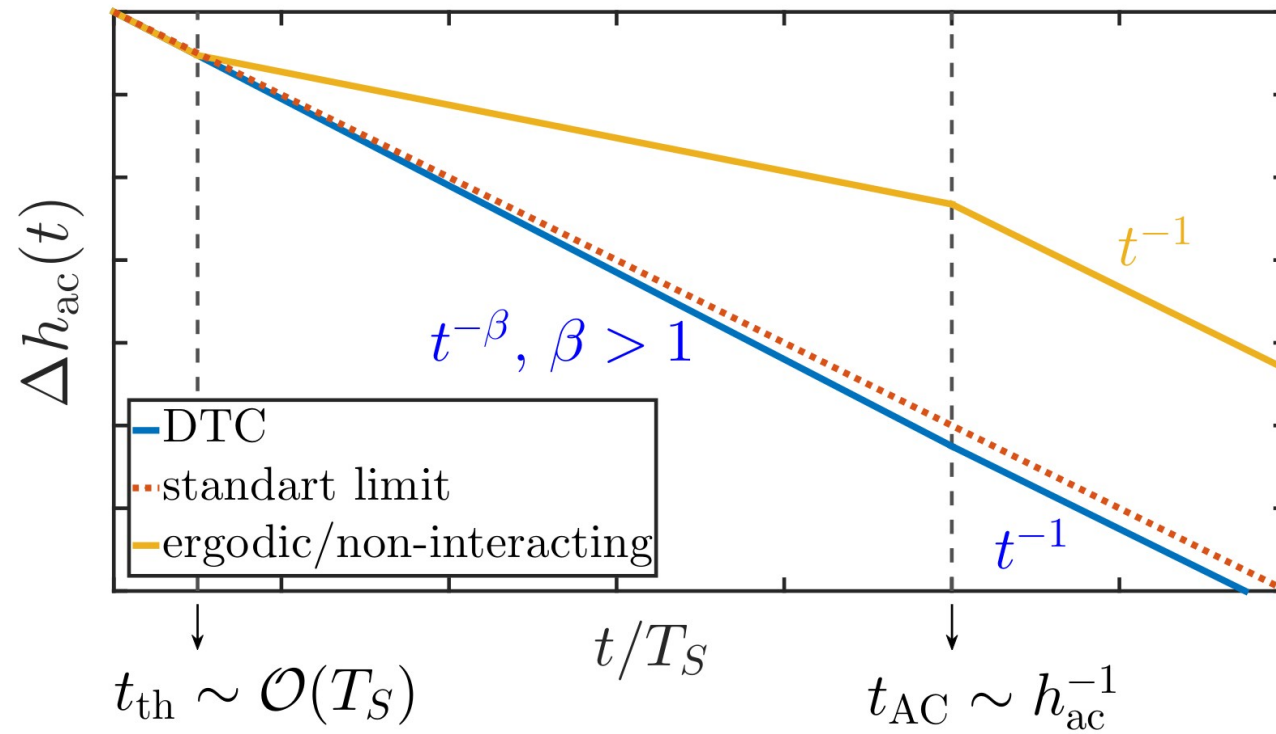
(dominant AC fields over dynamics)

- recovers Fisher quadratic growth,

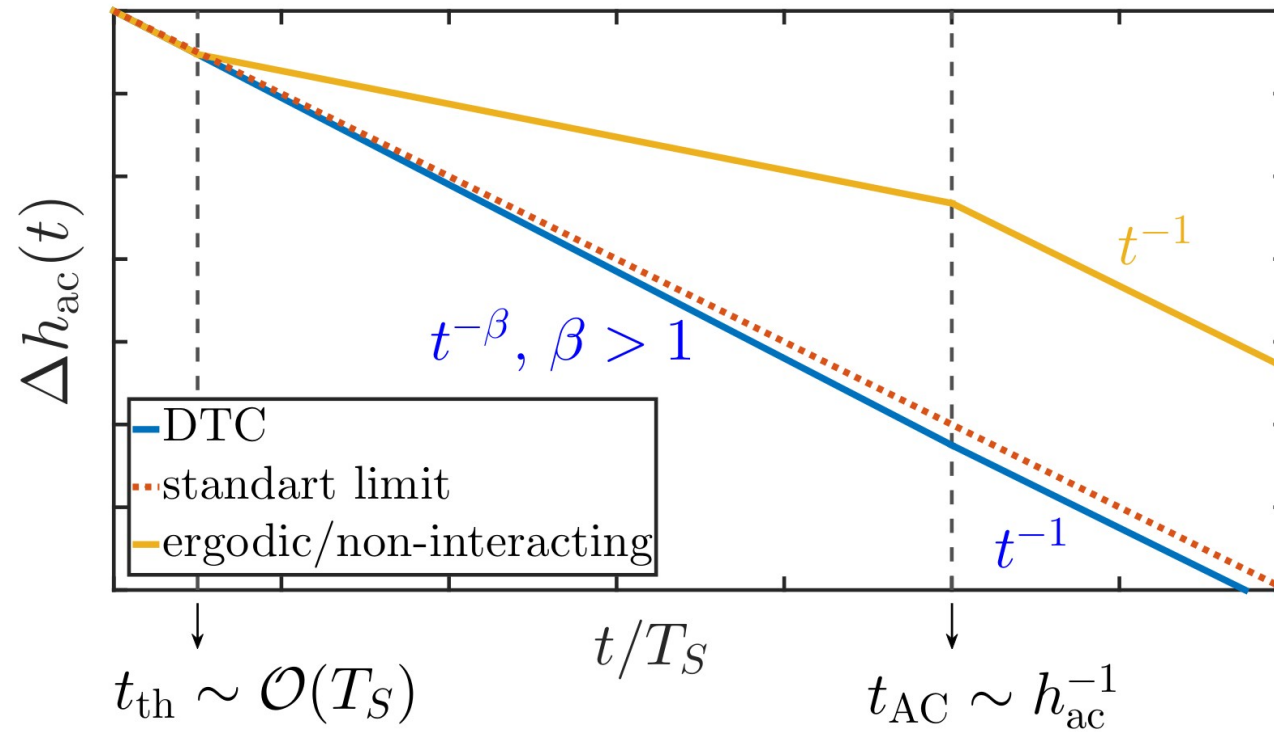
$$F_{h_{ac}}(t) \sim t^2$$



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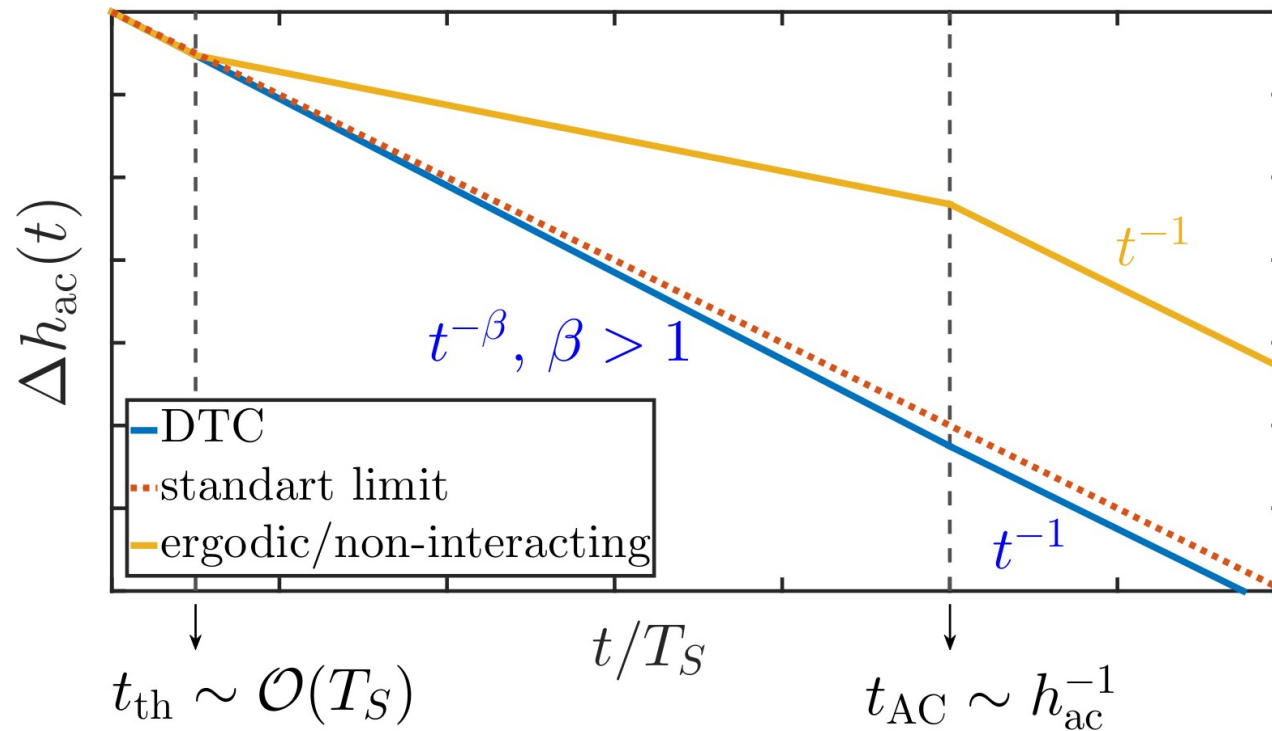
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In summary, Discrete Time Crystal sensors offer several advantages:

- **slow heating** – exponentially slow with system size N
- useful quantum correlations - **overcome classical limit;**
- **robust** protocol.

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