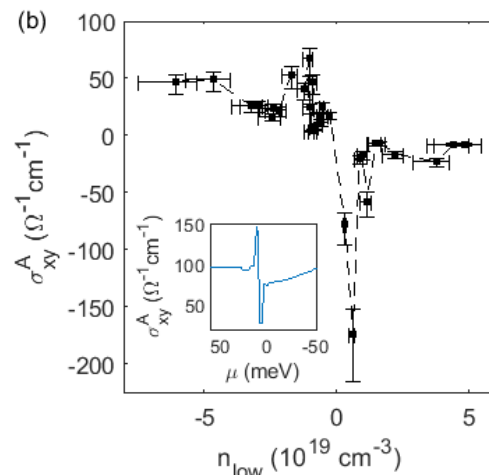
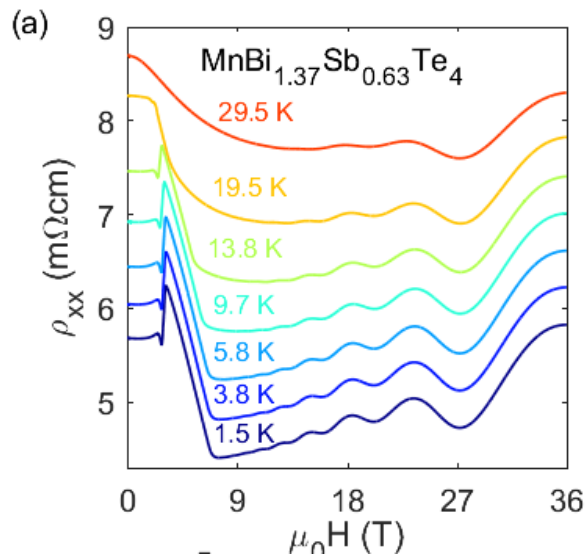


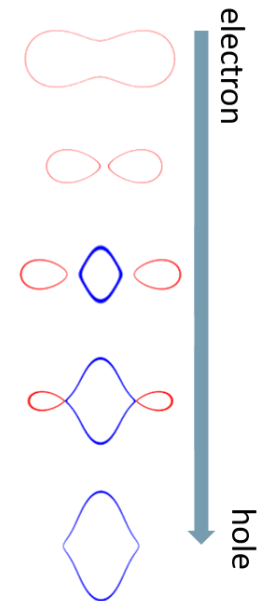
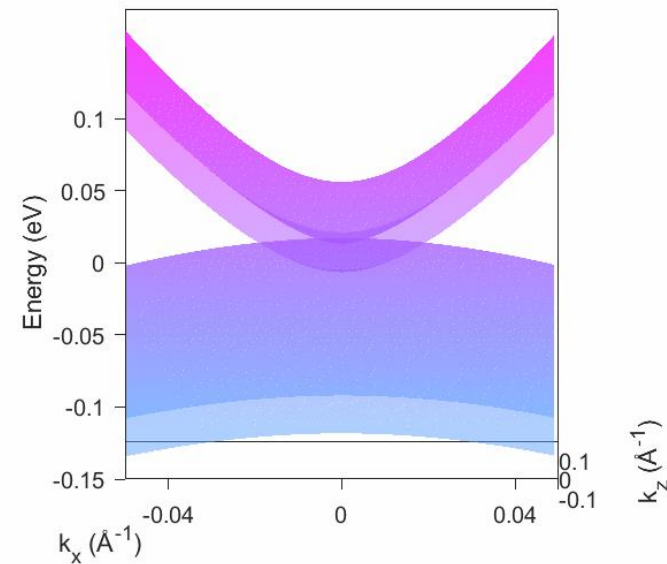
Realization of an **3D** ideal Weyl semimetal in $\text{MnBi}_{2-x}\text{Sb}_x\text{Te}_4$

Jiun-Haw Chu

Department of Physics, University of Washington



Type-II Weyl Semimetal



UW Quantum materials group

Qianni Jiang,, Yue Shi, Paul Malinowski, Aaron Wang,
Zhong Lin (joint with Xiaodong Xu)

UW Theory

Lingnan Shen, Chong Wang, Di Xiao

High Field Magnet Lab

Johannes Palmstrom, John Singleton
David Graf, Shalinee Chikara



Qianni Jiang
(Now GLAM Postdoc Fellow in Stanford)

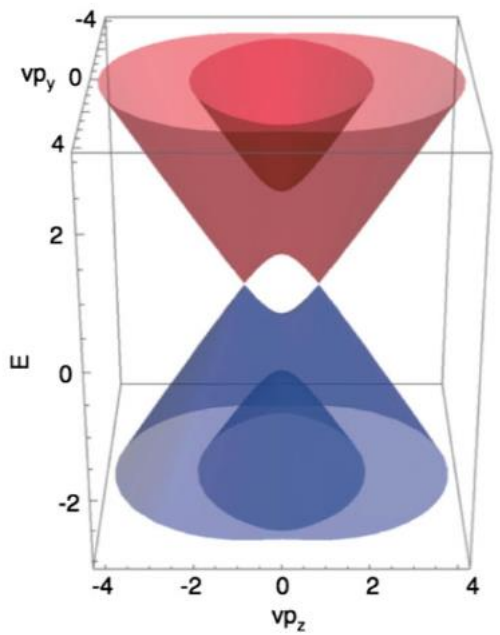


GORDON AND BETTY
MOORE
FOUNDATION

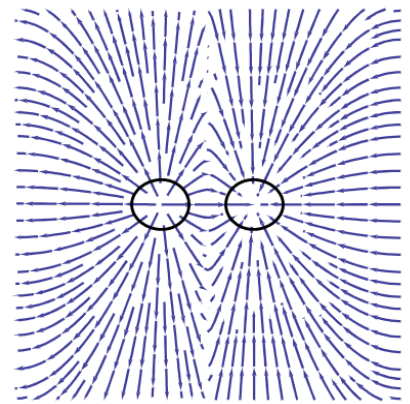
W Background of Weyl Semimetals

Weyl Equation

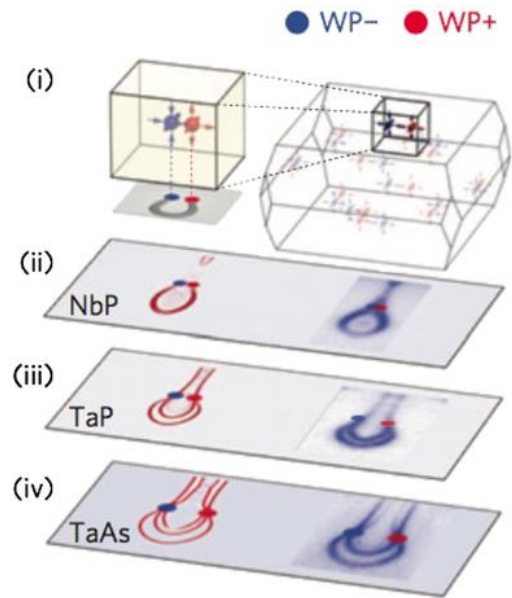
$$i \frac{\partial}{\partial t} \psi_{\pm} = \pm c \mathbf{p} \cdot \alpha \psi_{\pm}$$



Berry Flux

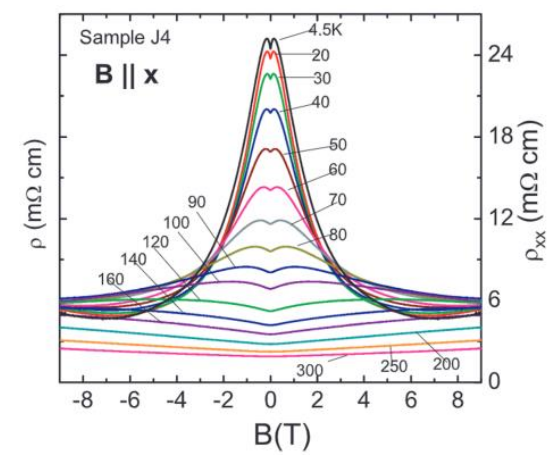


Fermi Arc



Low High Surface contribution
Low High Intensity (a.u.)

Chiral Anomaly



$$\frac{\partial}{\partial t} \rho_{xx} + \nabla \cdot \mathbf{J}_{xx} = -\chi \frac{e^3}{4\pi^2 \hbar^2} \mathbf{E} \cdot \mathbf{B}$$

Intrinsic anomalous Hall

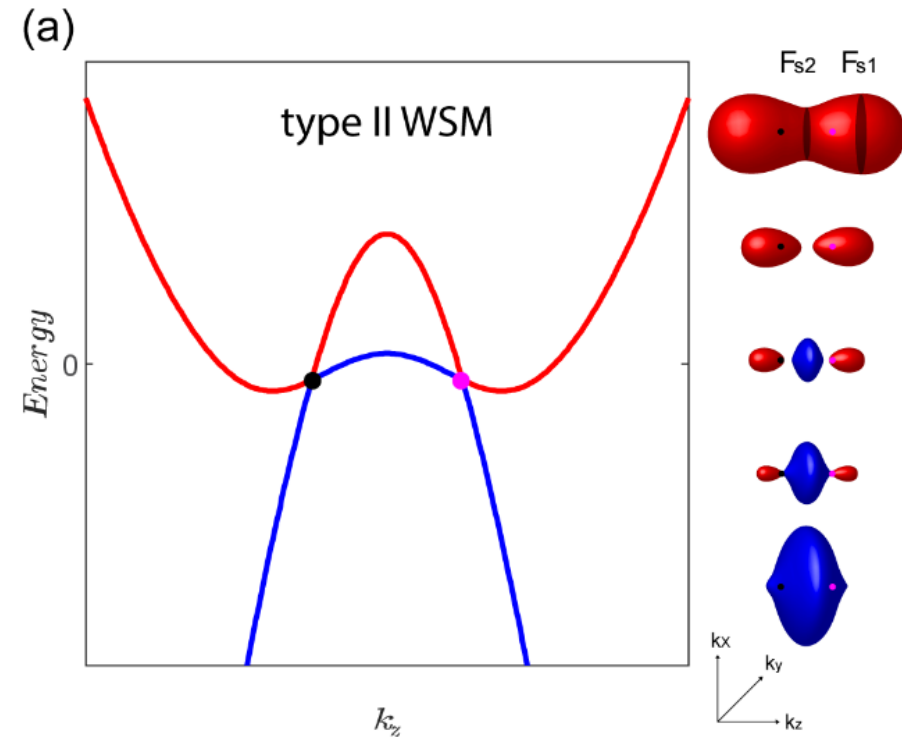
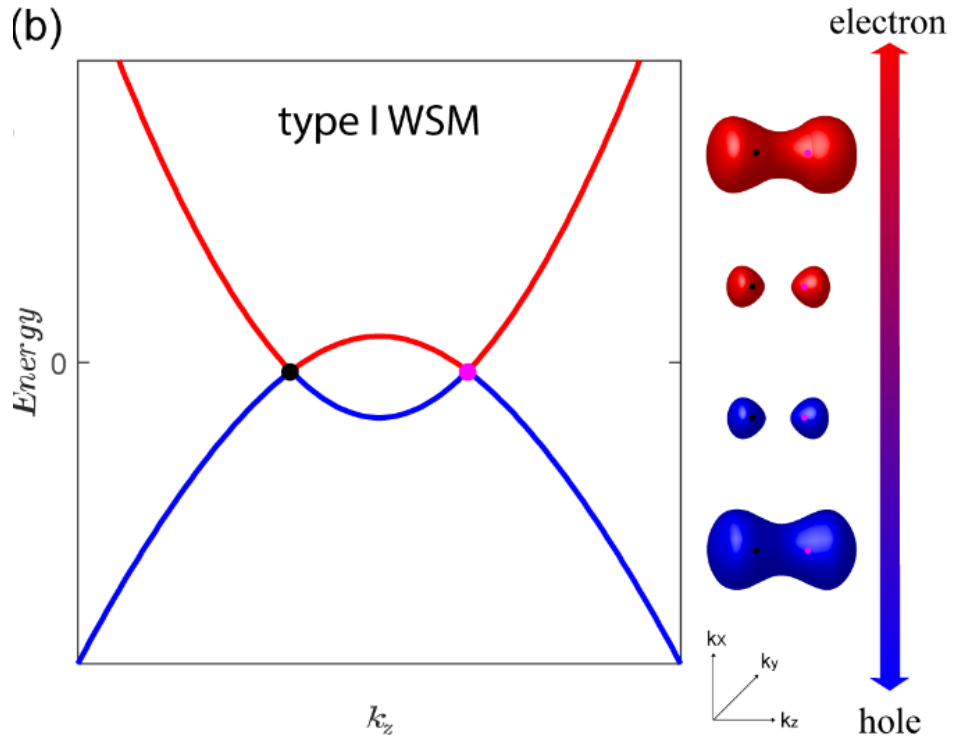
In an ideal Weyl semimetal:

$$\sigma_{xy}^A = \frac{e^2 Q}{2\pi^2 \hbar} \quad \text{Quantized}$$

Q : Weyl points separation

Z. Liu, et al. Nat. Mater. 15, 27 (2016)
 J. Xiong et al. Science 350,413-416 (2015)
 N.P. Armitage, et. al. Rev. Mod. Phys, 90(1), 015001 (2018)

W Two types of Weyl Semimetals



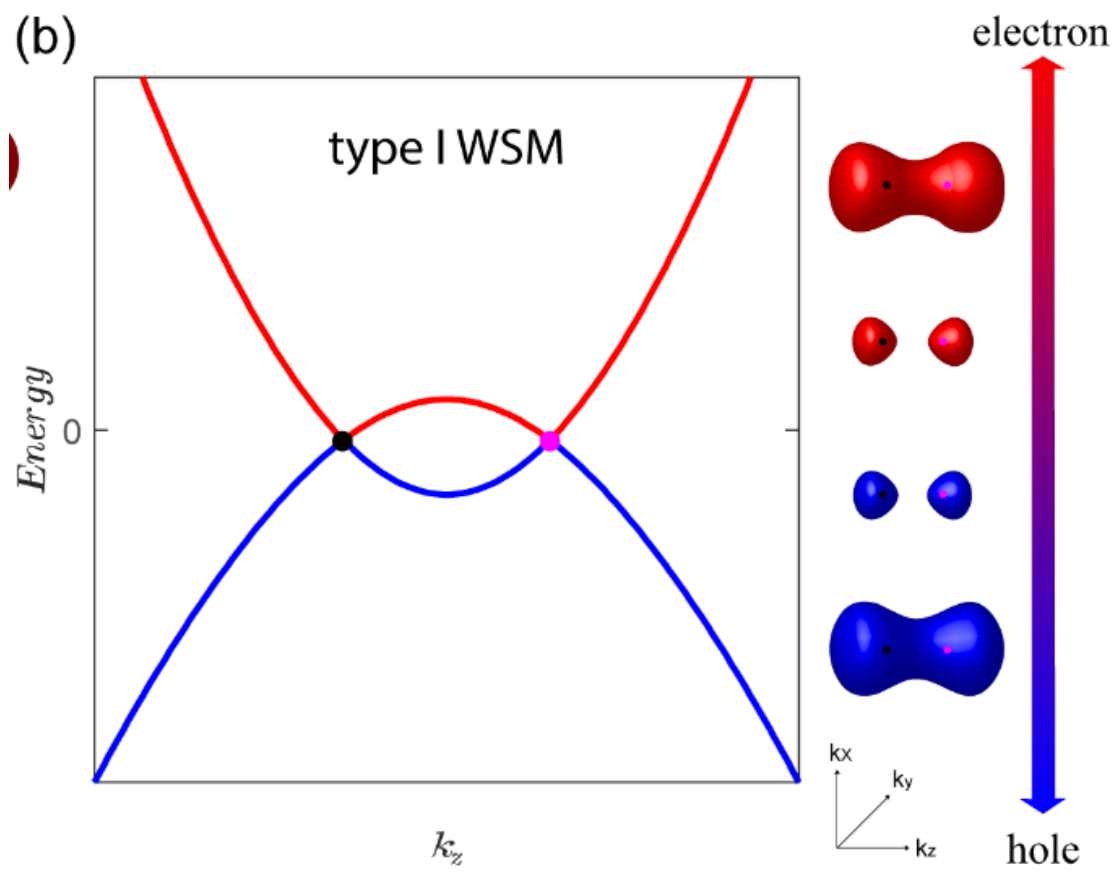
Z. Liu, et al. Nat. Mater. 15, 27 (2016)

J. Xiong et al. Science 350,413-416 (2015)

N.P. Armitage, et. al. Rev. Mod. Phys, 90(1), 015001 (2018)

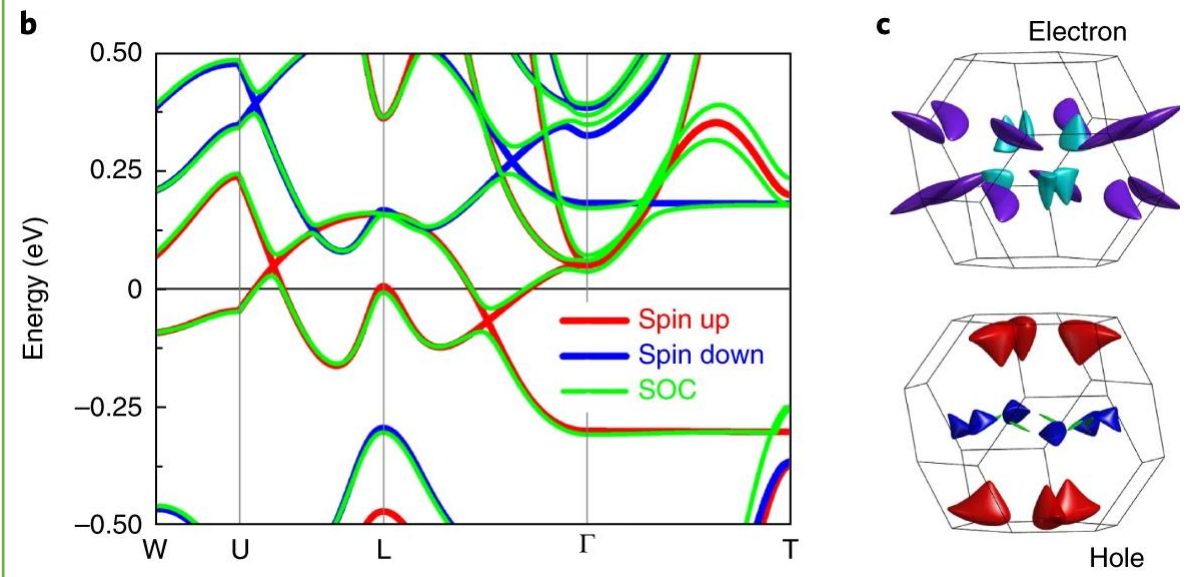
W The need for an ideal Weyl semimetal

WSM to the theorists



WSM in reality

$\text{Co}_3\text{Sn}_2\text{S}_2$ FM WSM



Liu, E., Sun, Y., Kumar, N. *et al.* Giant anomalous Hall effect in a ferromagnetic kagome-lattice semimetal. *Nature Phys* **14**, 1125–1131 (2018).

W The need for an ideal Weyl semimetal

PRL 107, 127205 (2011)

PHYSICAL REVIEW LETTERS

week ending
16 SEPTEMBER 2011

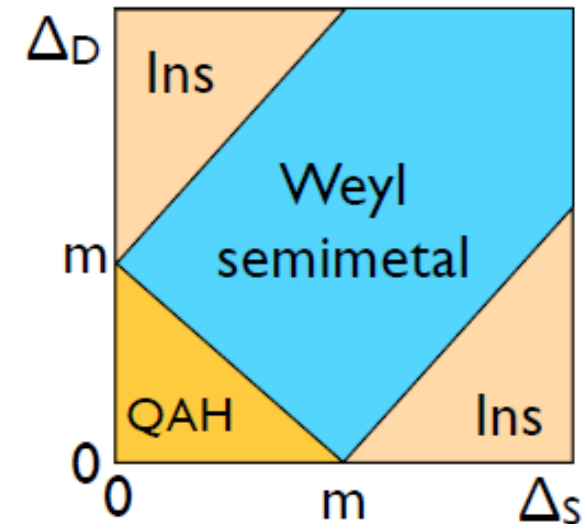
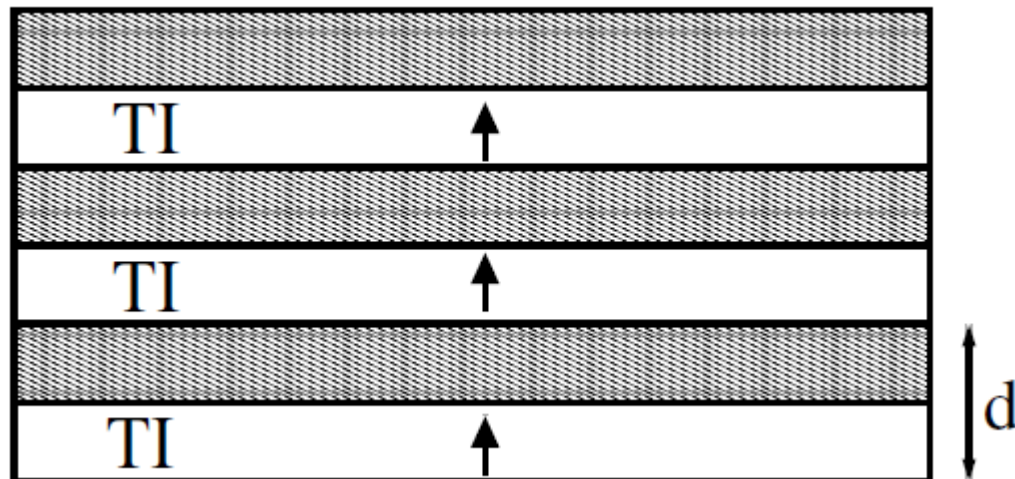
Weyl Semimetal in a Topological Insulator Multilayer

A. A. Burkov^{1,2} and Leon Balents²

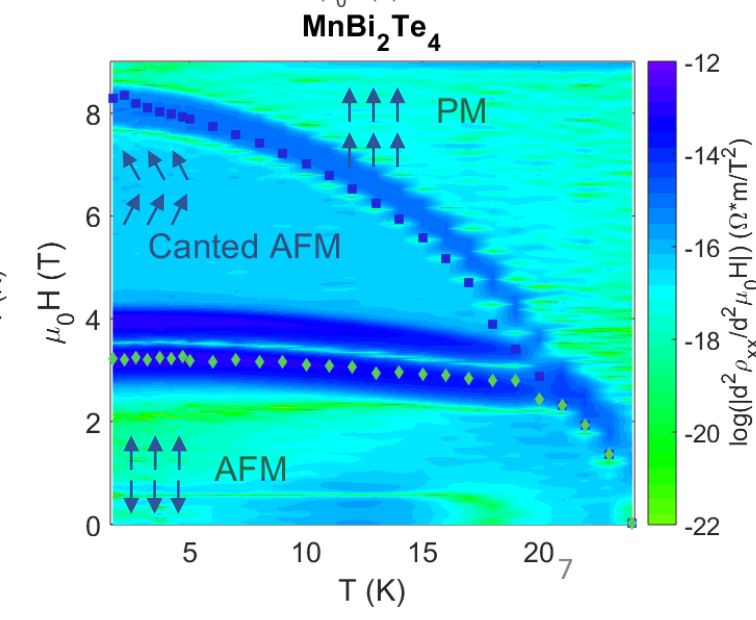
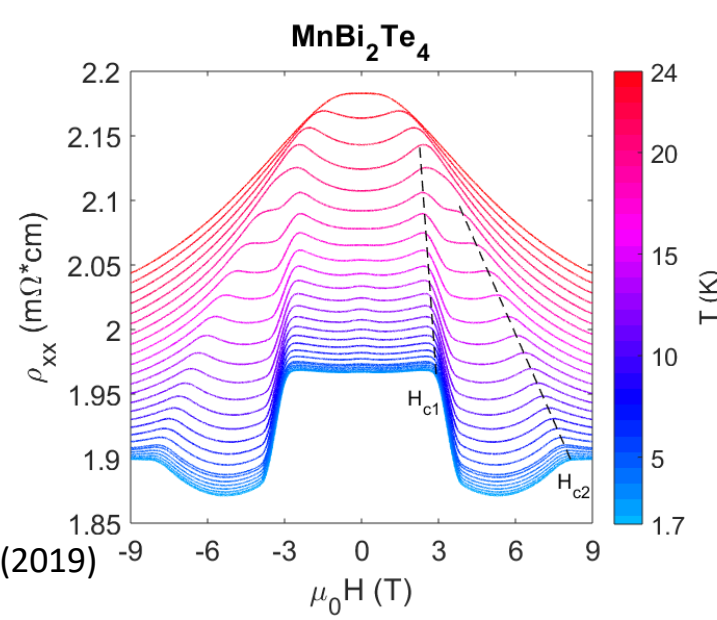
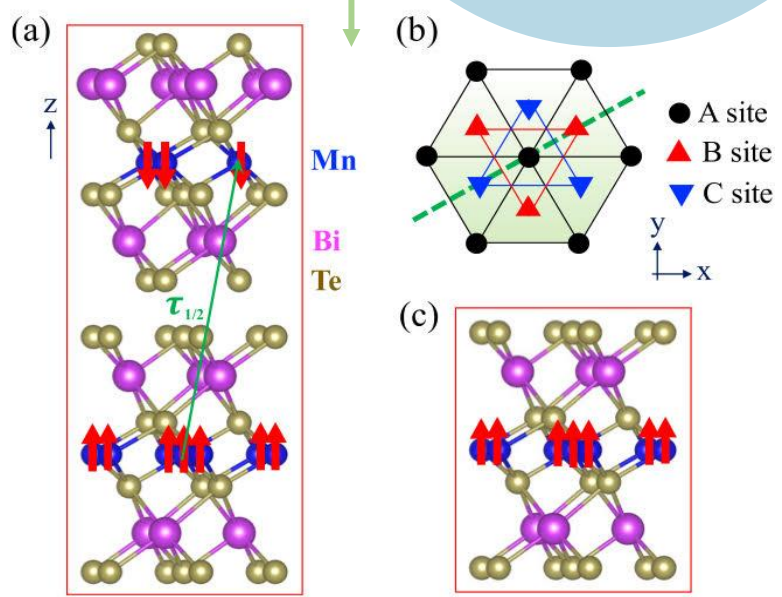
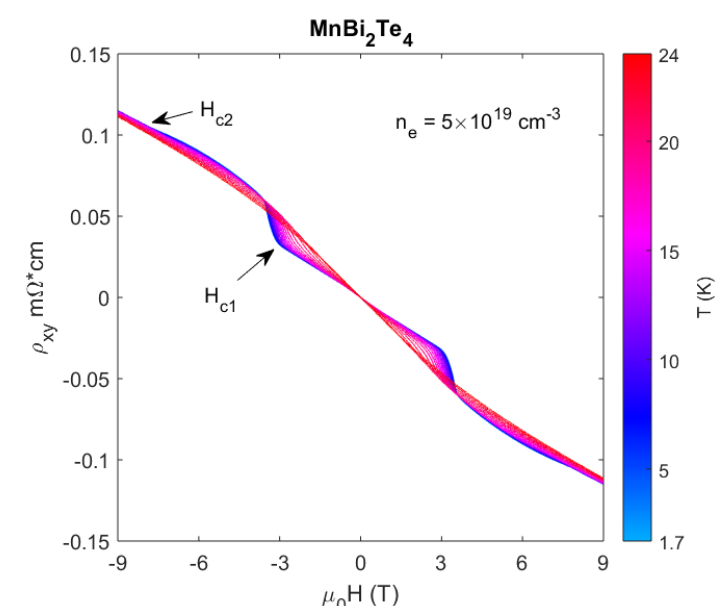
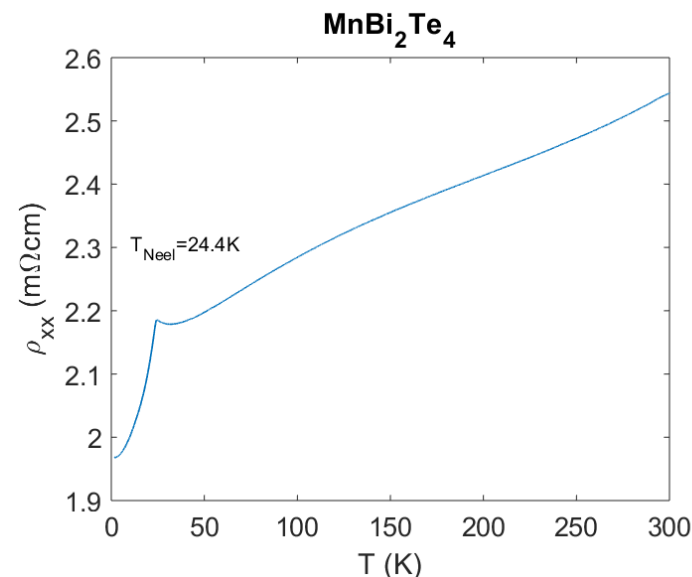
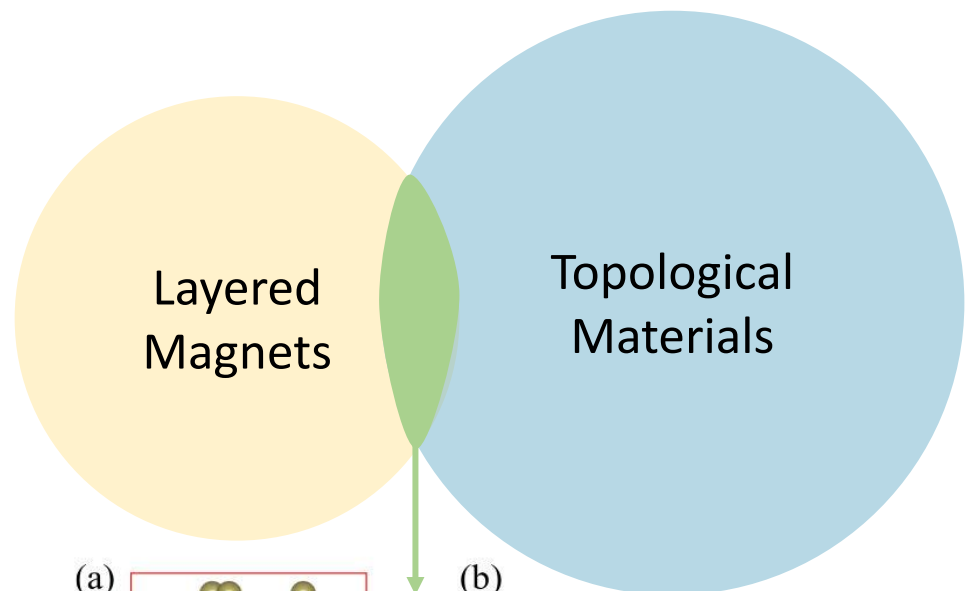
¹Department of Physics and Astronomy, University of Waterloo, Waterloo, Ontario N2L 3G1, Canada

²Kavli Institute for Theoretical Physics, University of California, Santa Barbara, California 93106, USA

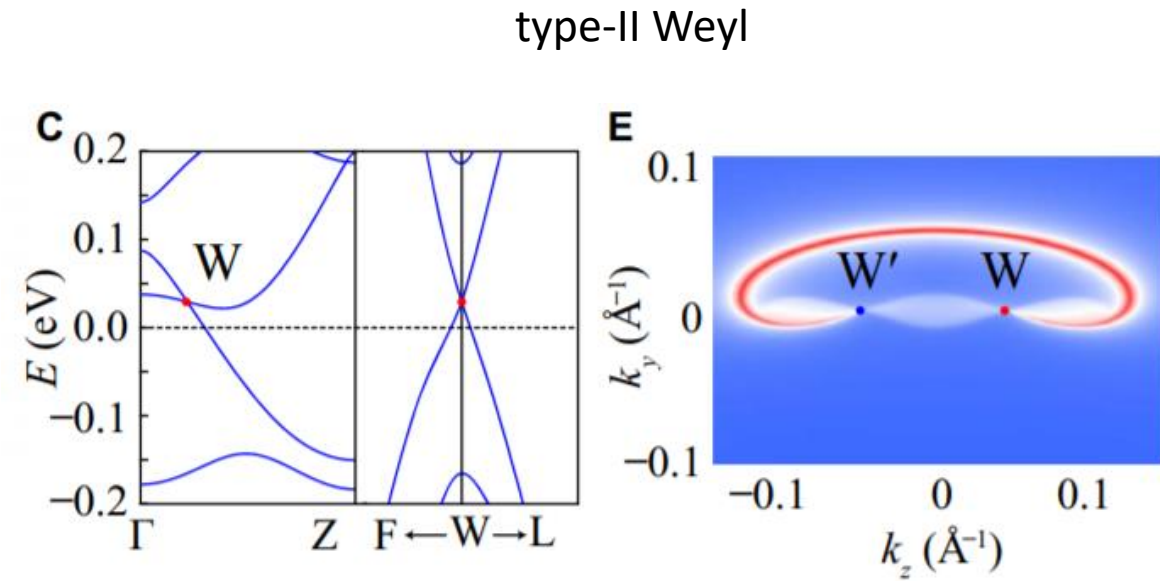
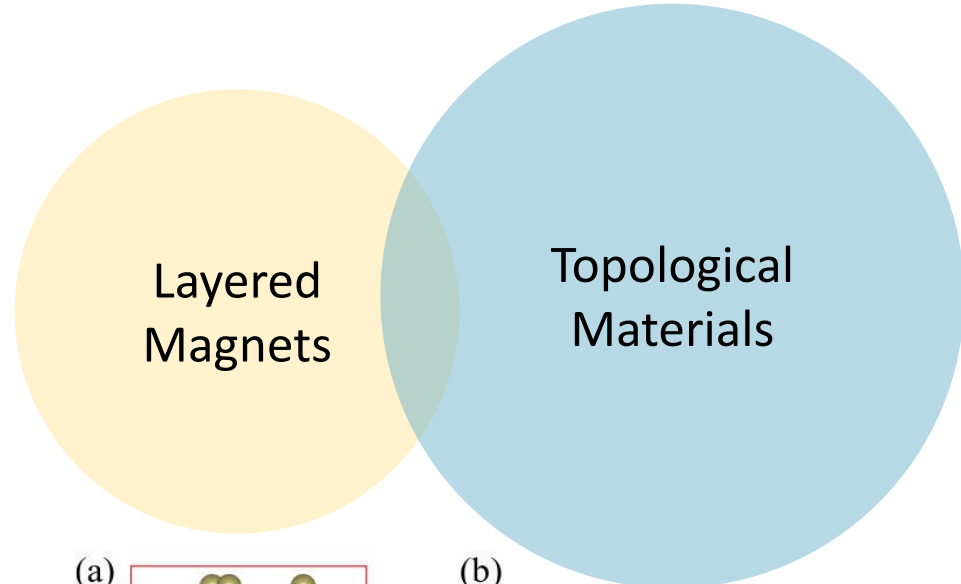
(Received 27 May 2011; published 16 September 2011)



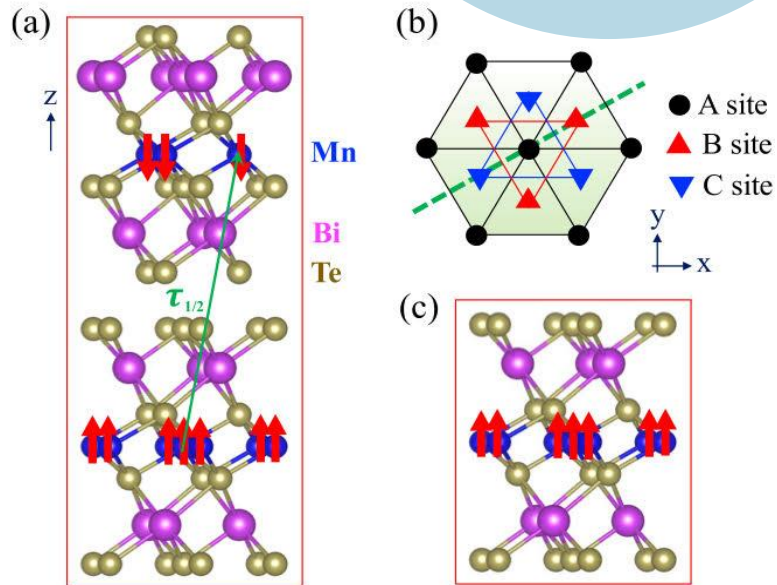
W Background of MnBi_2Te_4



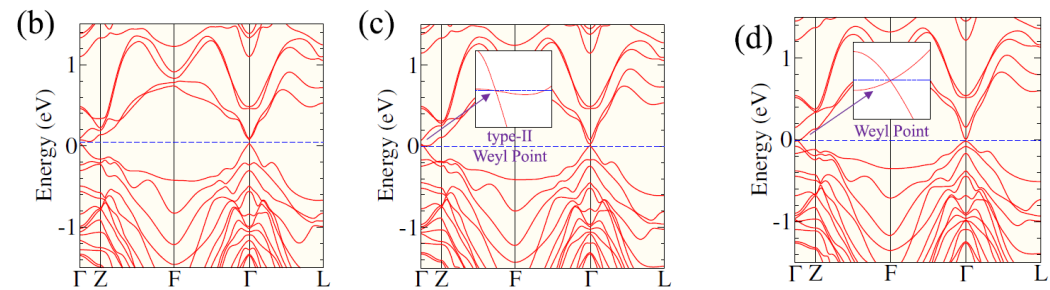
W MnBi₂Te₄ as a natural superlattice of FMs and TIs



(Li J., et al., Sci. Adv. 2019;5: eaaw5685)



Trivial semiconductor type-II Weyl type-I Weyl



(Zhang D., et al, PRL 122.20 (2019): 206401)

W MnBi₂Te₄ as a natural superlattice of FMIs and TIs

Layered
Magnets

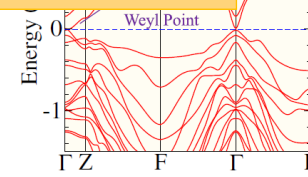
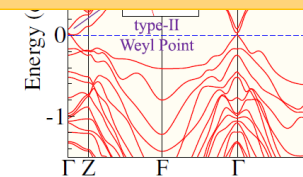
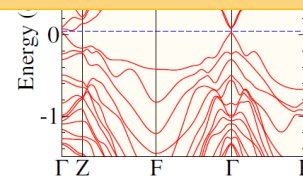
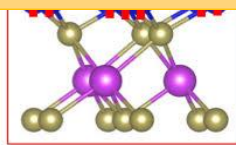
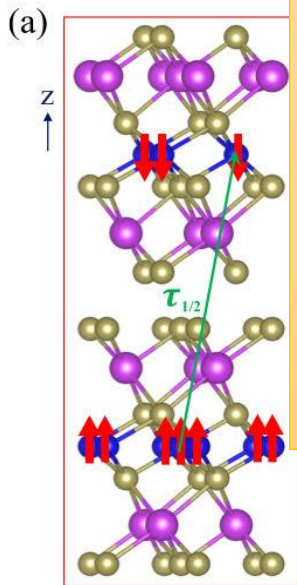
Topological

type-II Weyl

$C_{0.2}$ $E_{0.1}$

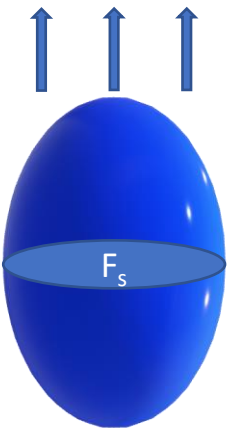
Challenge: the WSM requires an external field to induce the ferromagnetic state, incompatible with the ARPES.

Solution: use quantum oscillations to probe the electronic structure!

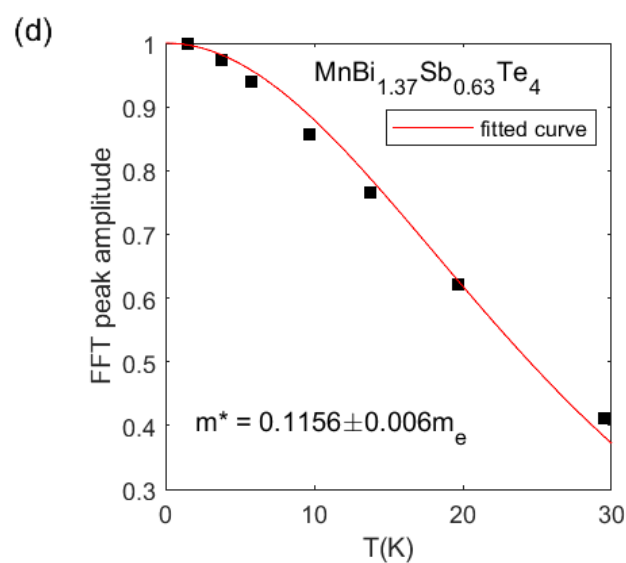
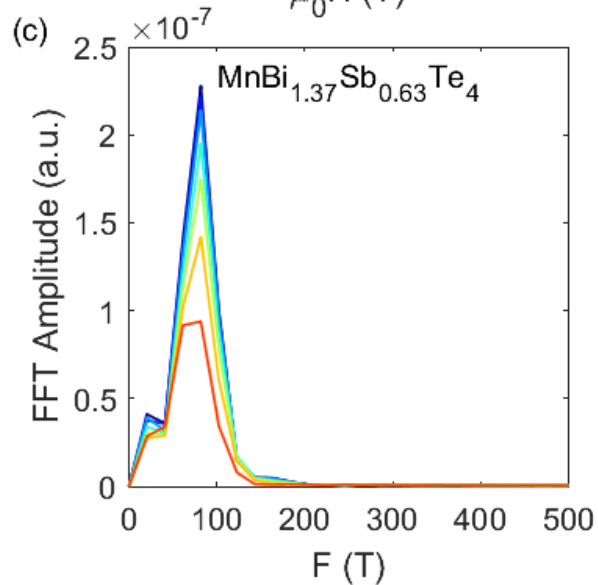
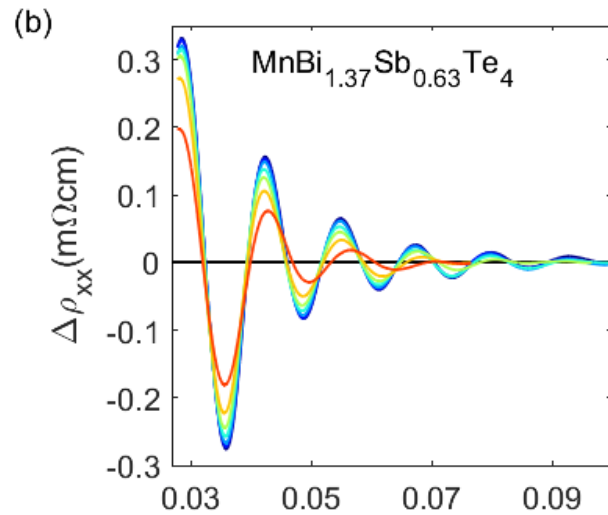
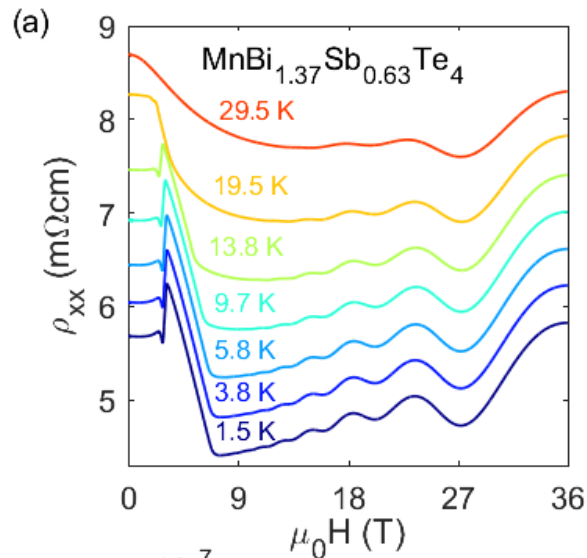


W Shubnikov-de Haas Oscillations of $\text{MnBi}_{2-x}\text{Sb}_x\text{Te}_4$

$\mu_0 H \parallel [001]$



$$\Delta\rho_{xx} \propto R_T R_D \cos 2\pi \left(\frac{F}{B} + \phi \right)$$



Dingle damping factor:

$$R_D = \exp\left(-\frac{\alpha T_D m^*}{B}\right)$$

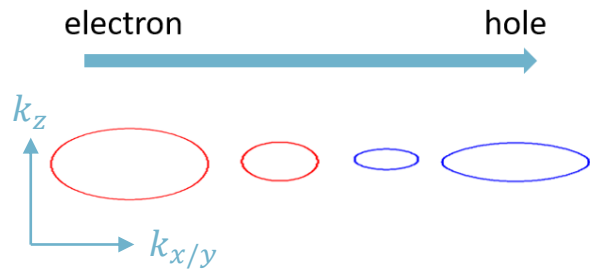
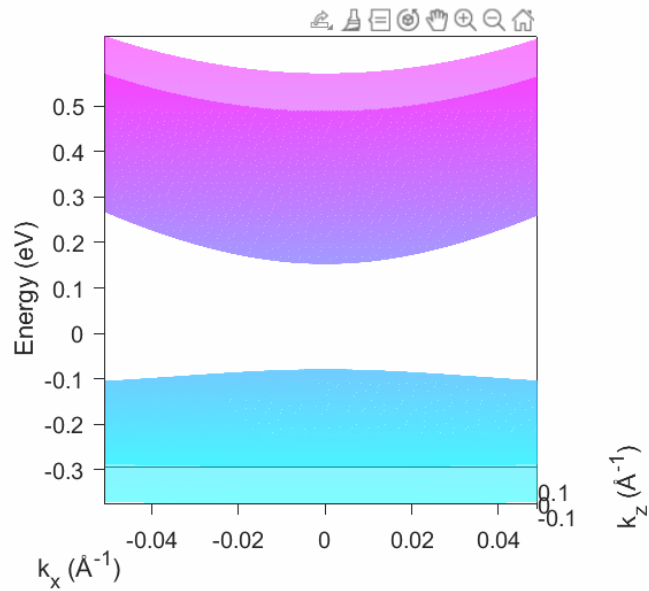
Thermal damping factor:

$$R_T = \frac{\alpha T m^*}{B \sinh(\alpha T m^* / B)}$$

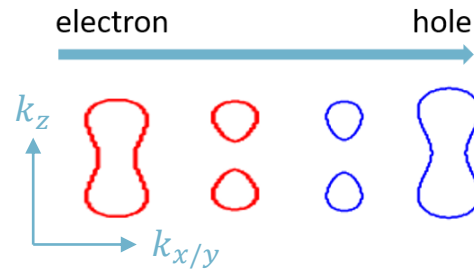
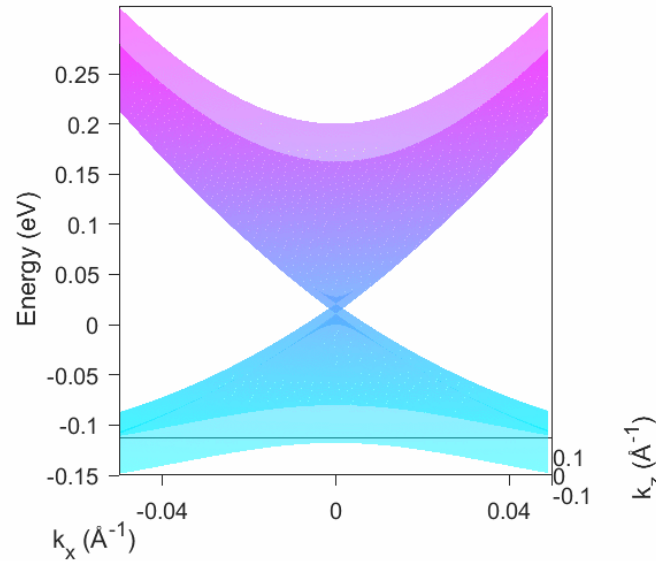
W Fermi surface evolution of Weyl Semimetals

Q: How to distinguish a normal semiconductor, an ideal type-I WSM and an ideal type-II WSM in this case?

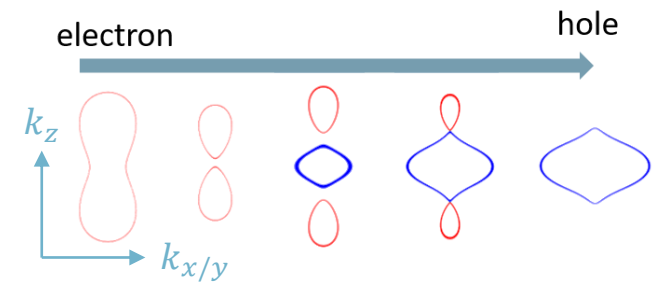
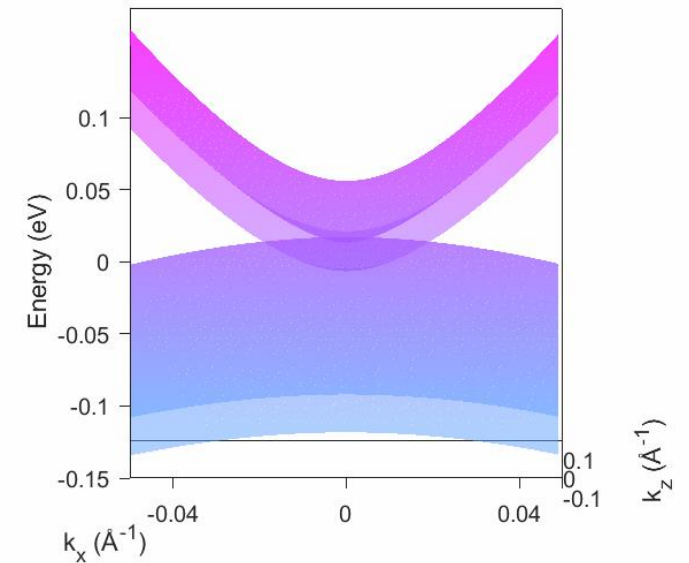
Normal Semiconductor

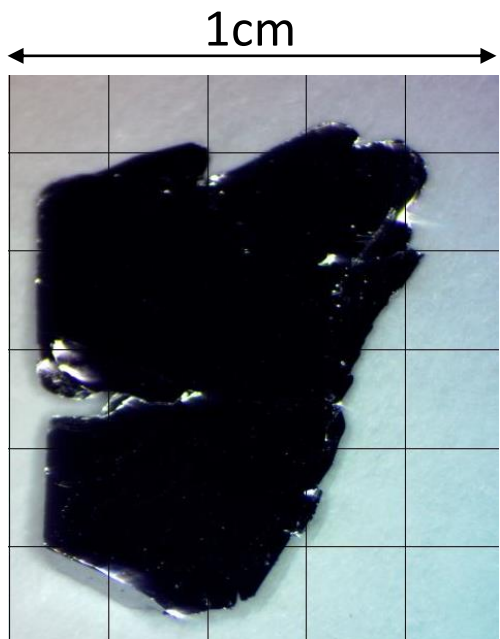


Type-I Weyl Semimetal

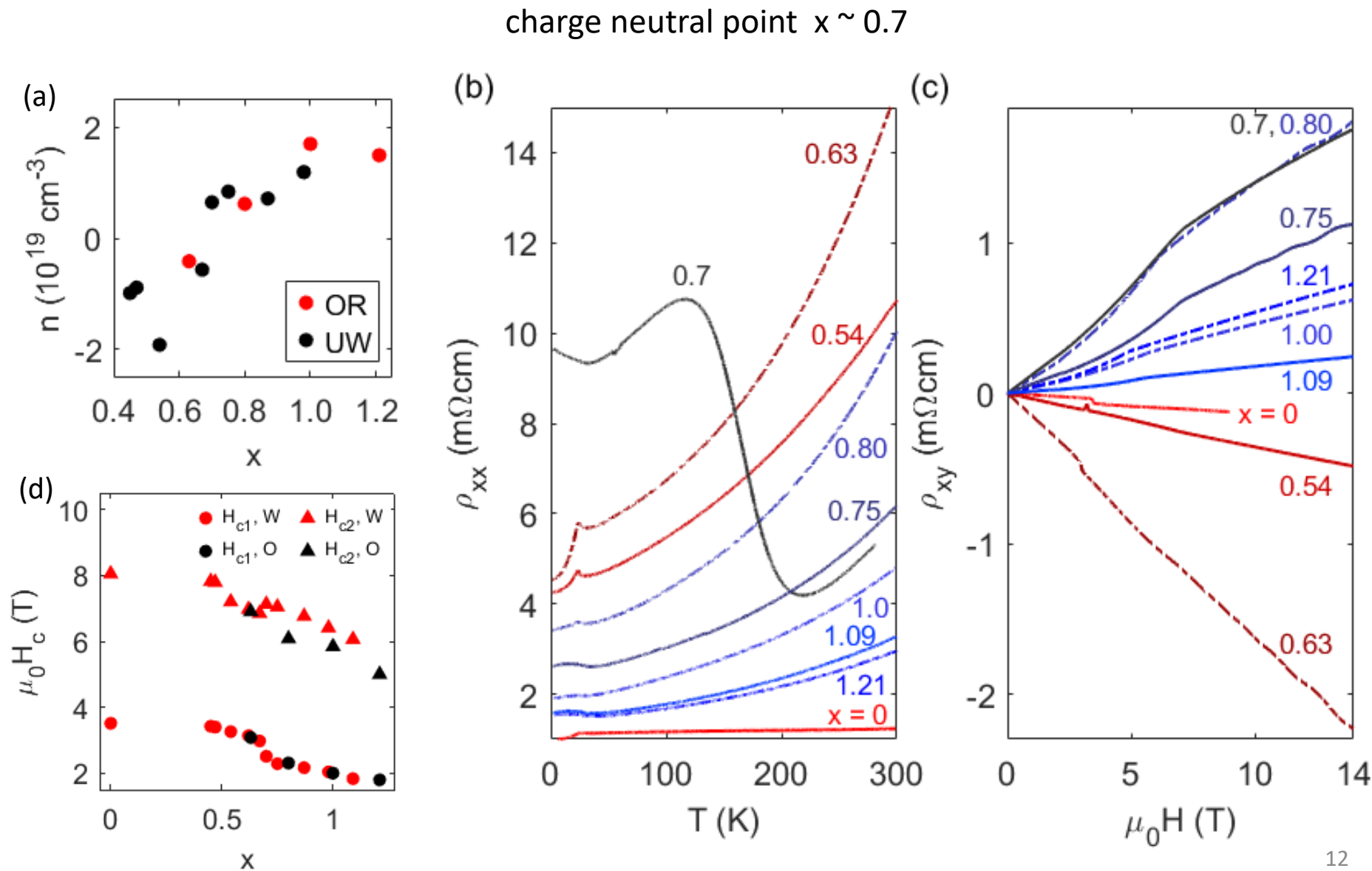


Type-II Weyl Semimetal

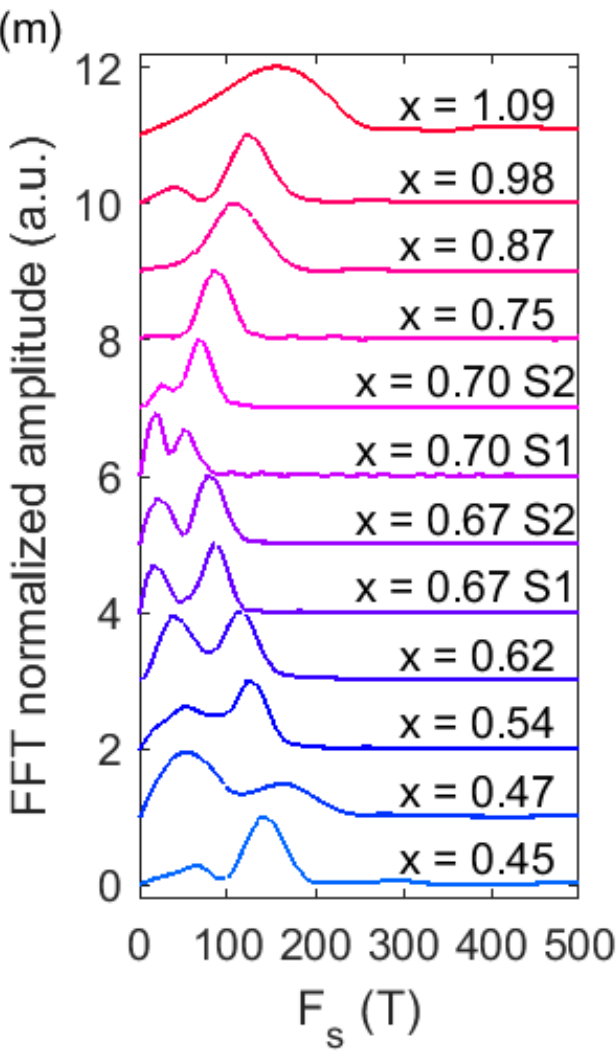
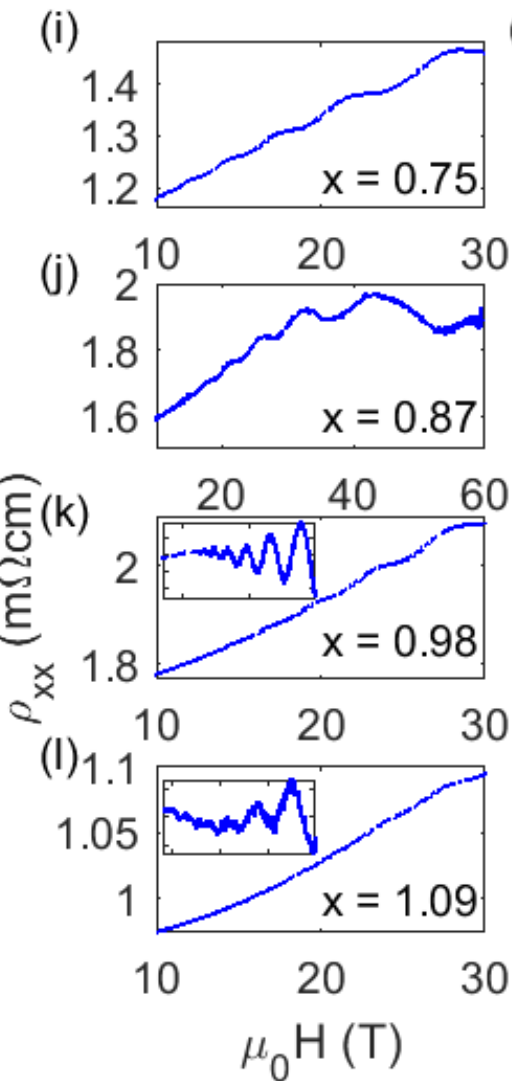
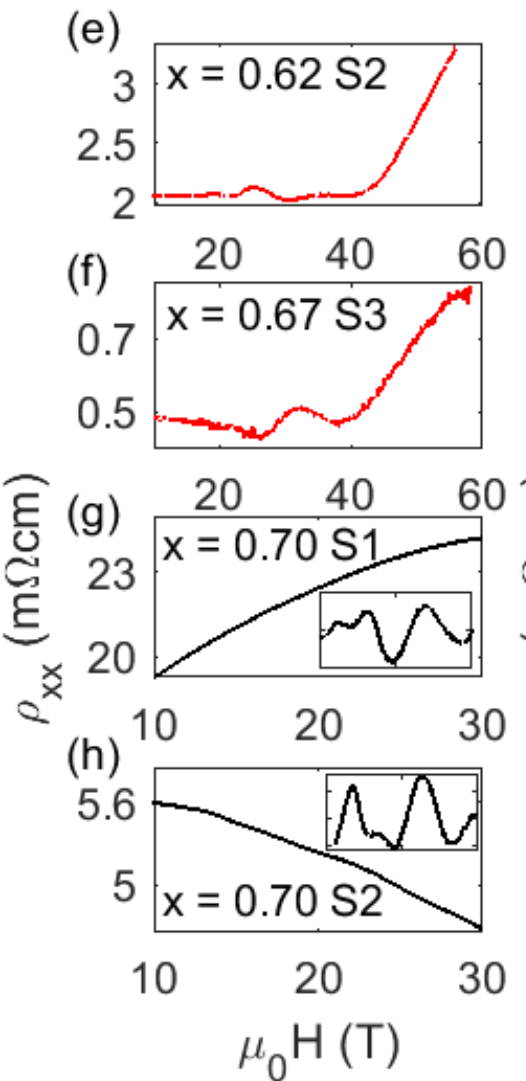
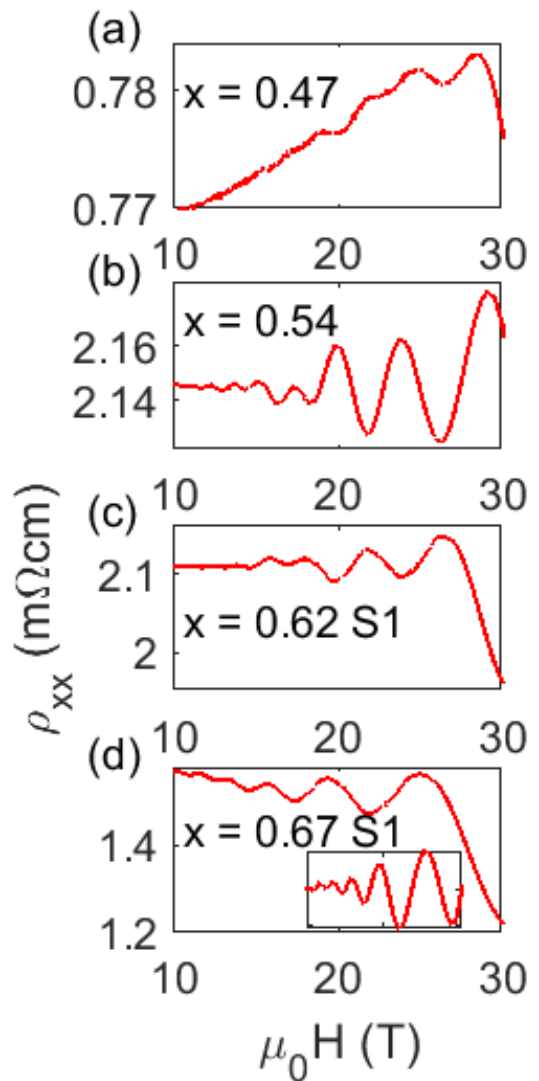




Single crystal of $\text{MnBi}_{2-x}\text{Sb}_x\text{Te}_4$



W SdH Oscillations under a 31T DC Field



Heavily hole doped

↑

charge neutrality

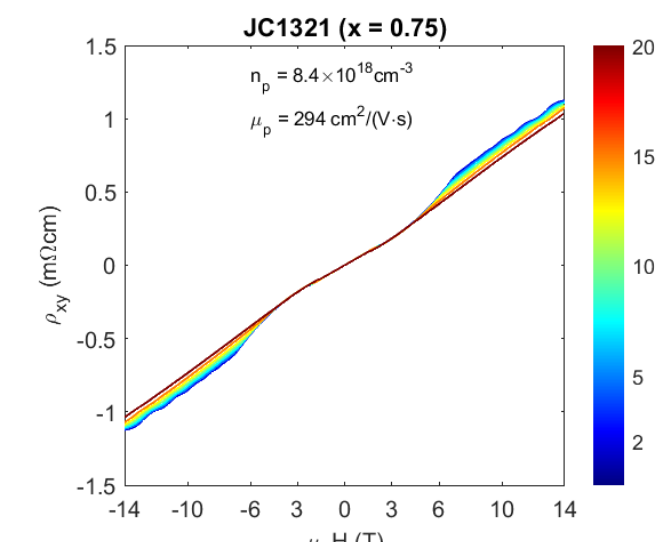
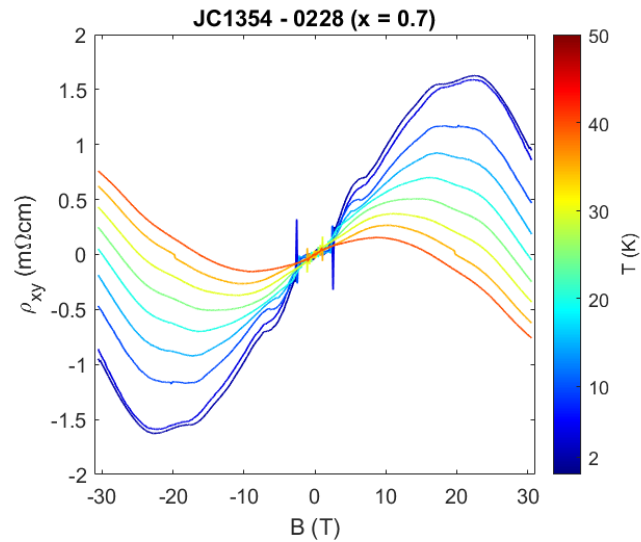
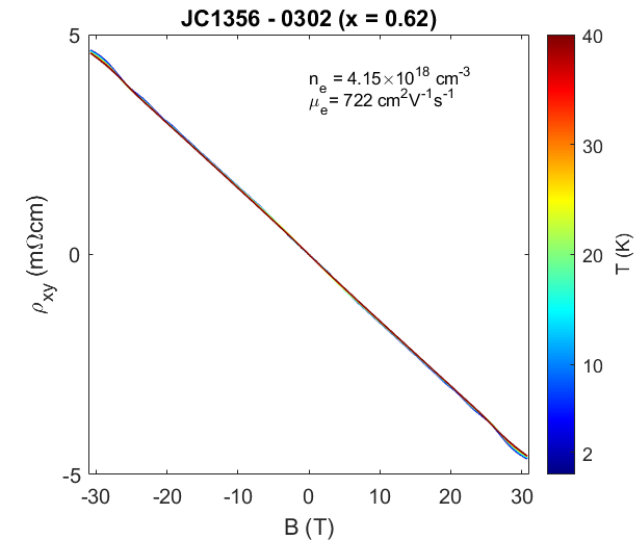
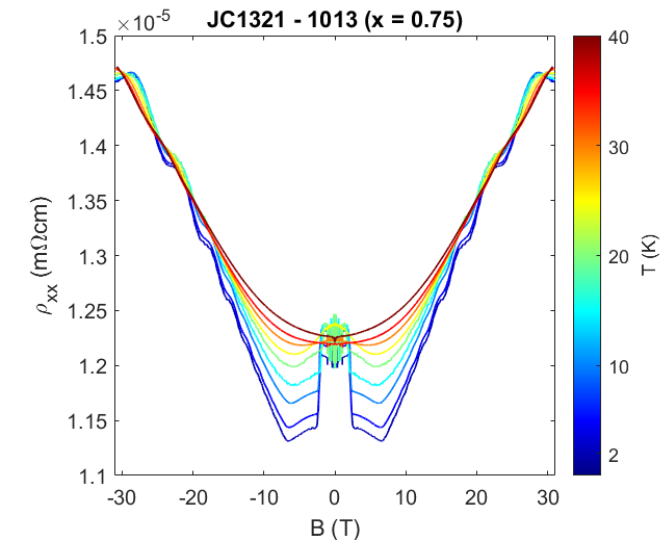
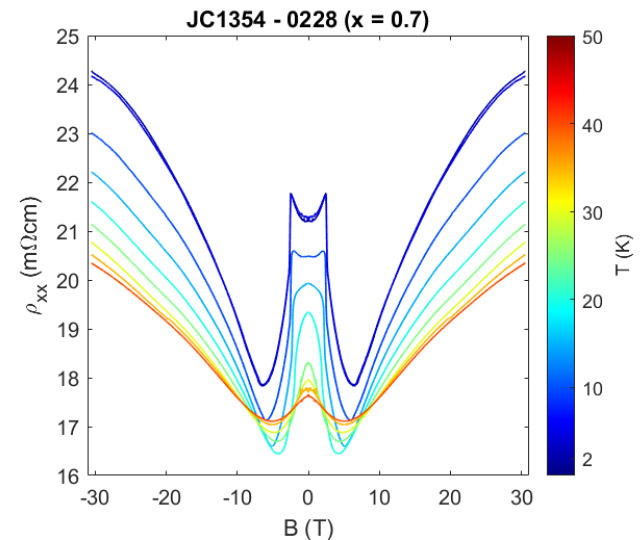
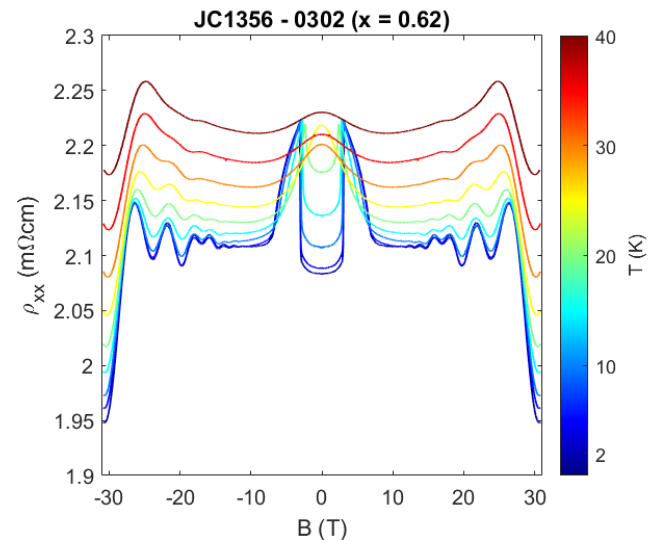
↑

heavily Electron doped



SdH Oscillations under a 31T DC Field

$\mu_0 H \parallel [001]$

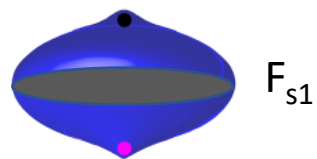
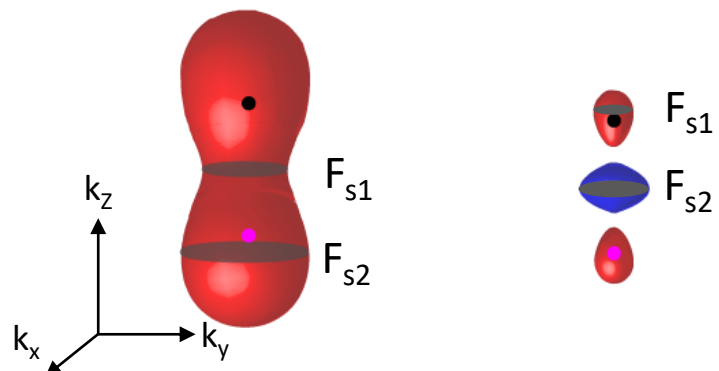
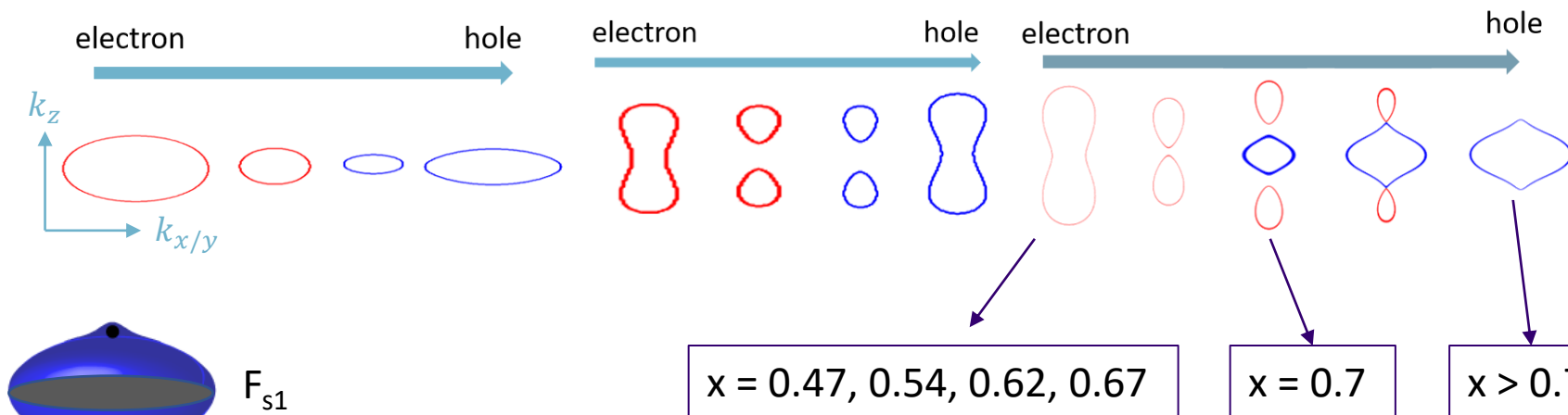
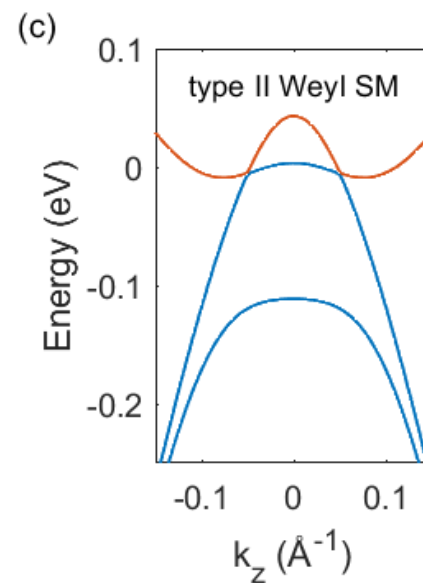
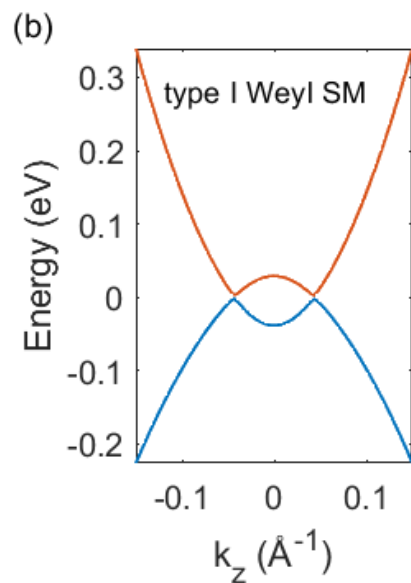
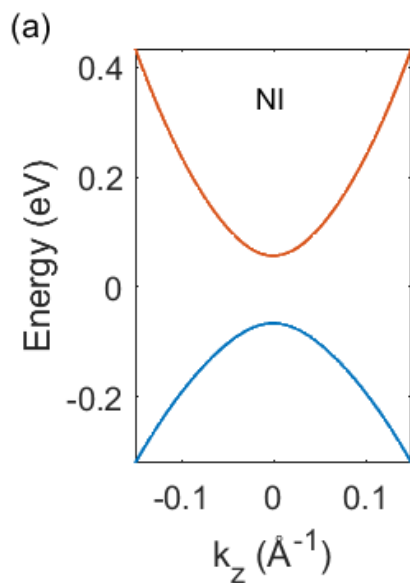
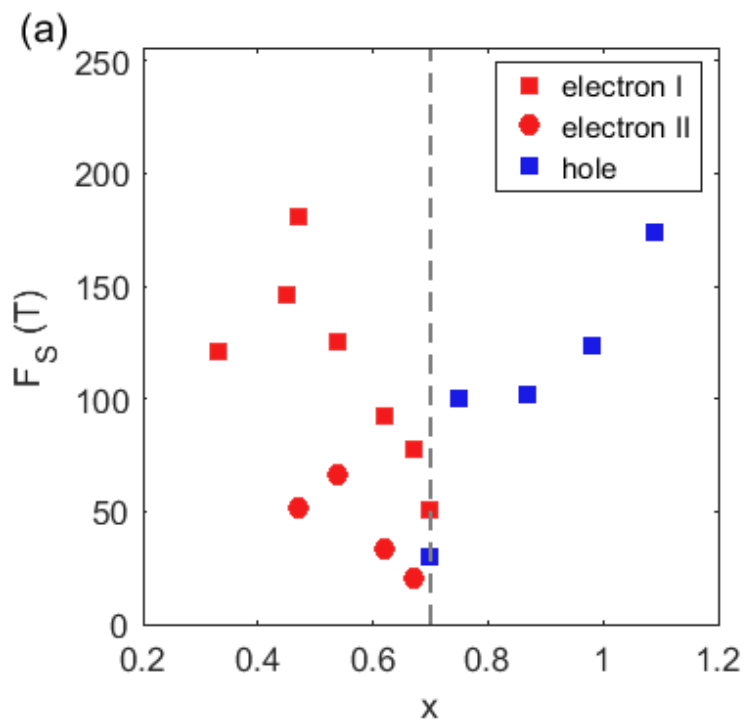


slightly electron-doped

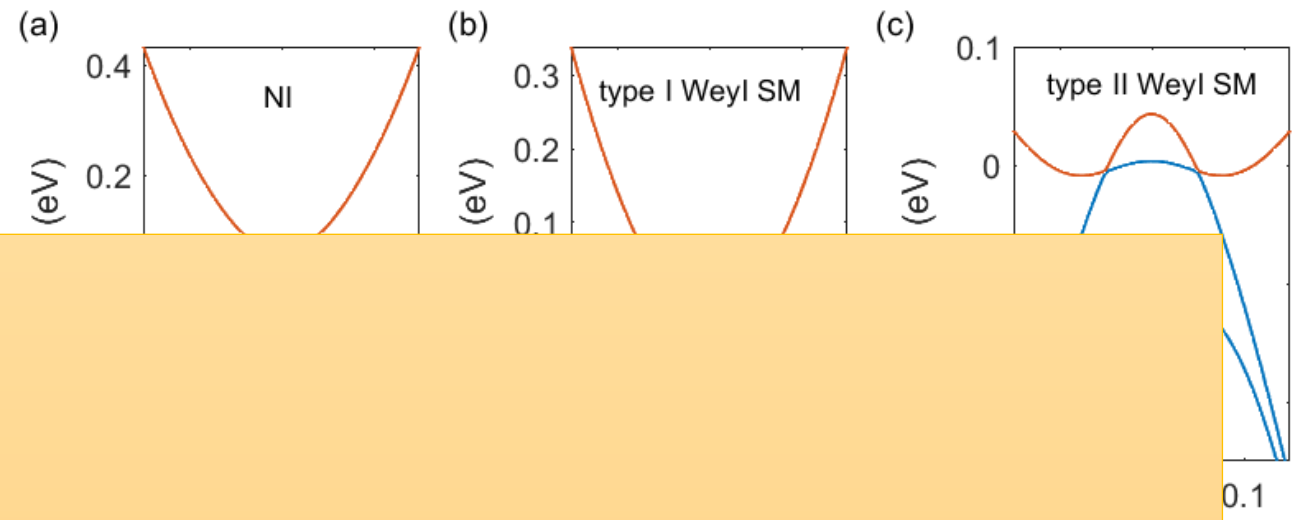
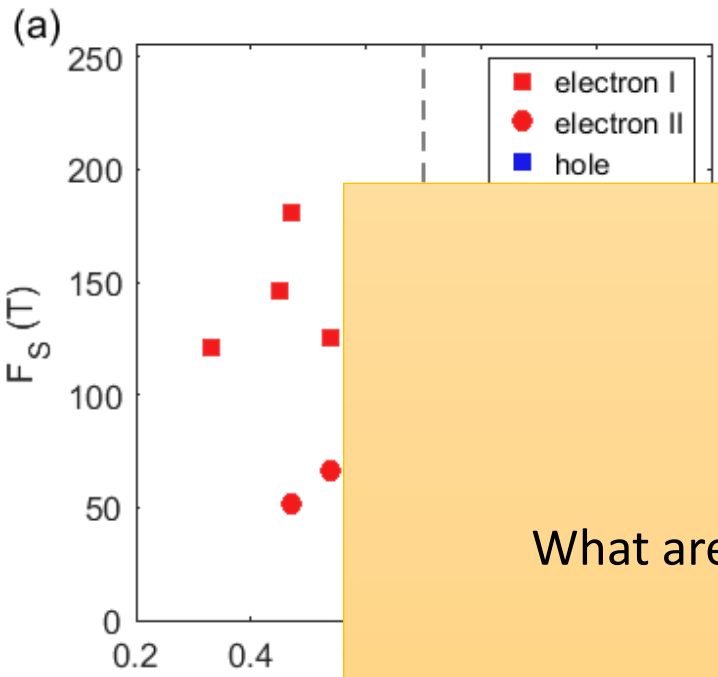
near charge neutral point

slightly hole-doped

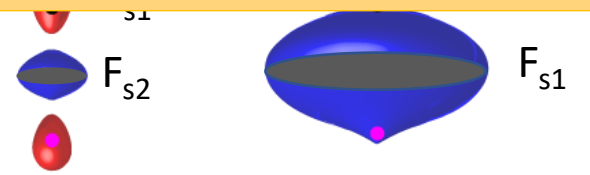
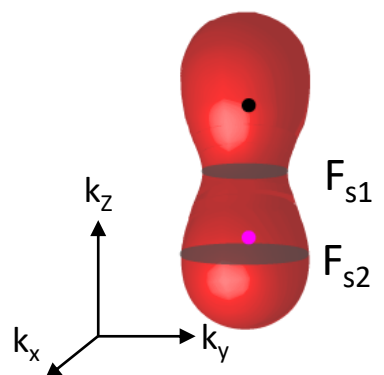
W Evolution of Fermi Pockets



W Evolution of Fermi Pockets



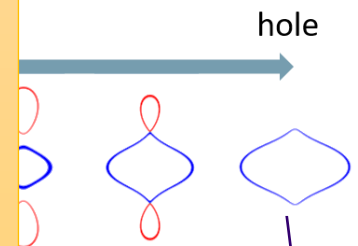
What are the observable consequences of an ideal type-II WSM?



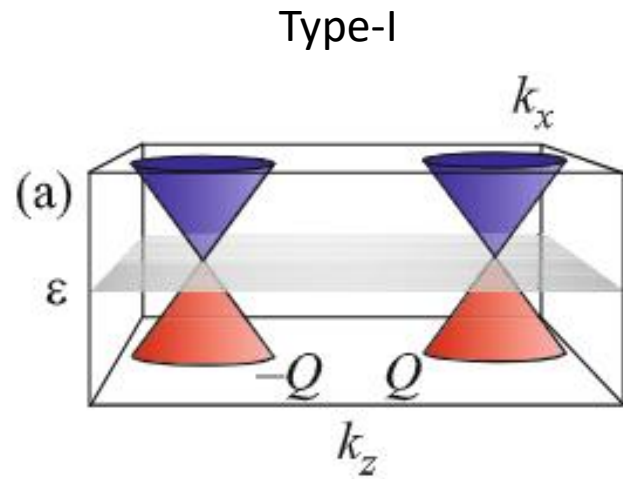
$x = 0.47, 0.54, 0.62, 0.67$

$x = 0.7$

$x > 0.7$

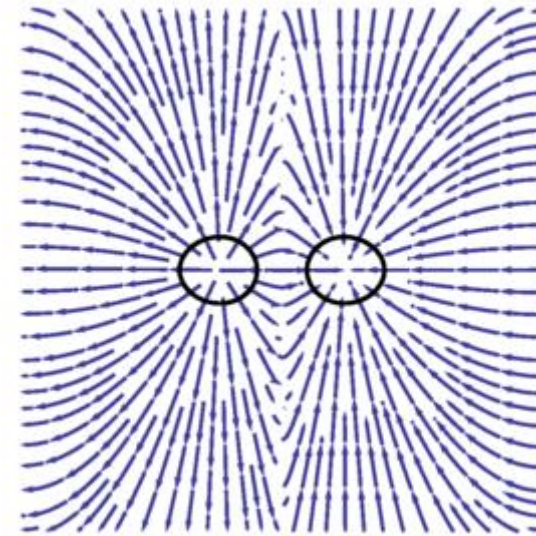


W Anomalous Hall Conductivity of Ideal Weyl Semimetal



$$\sigma_{xy} = \frac{e^2 Q}{2\pi^2 \hbar}$$

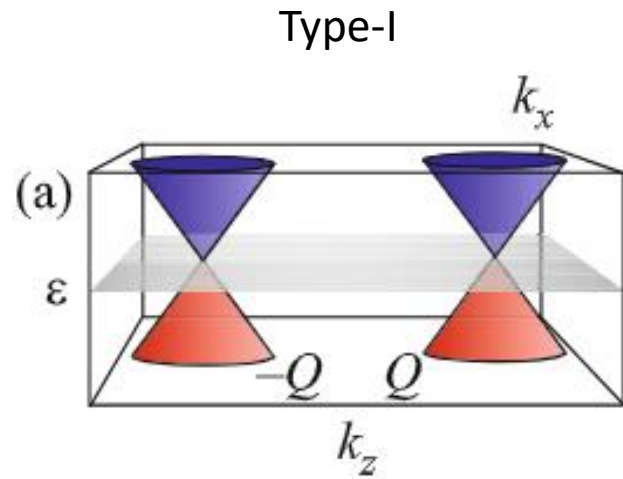
Q : Weyl points separation



Free carrier contribution:

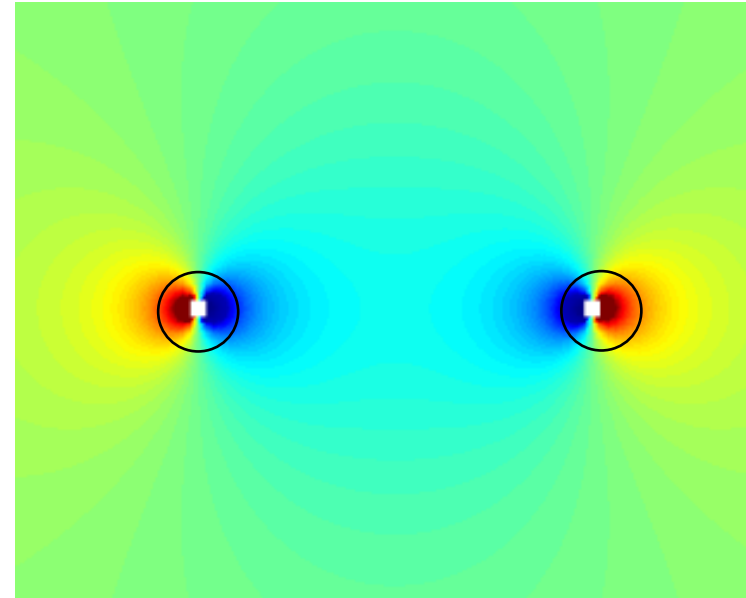
Integrate the z-component of Berry curvature over the occupied states

W Anomalous Hall Conductivity of Ideal Weyl Semimetal



$$\sigma_{xy} = \frac{e^2 Q}{2\pi^2 \hbar}$$

Q : Weyl points separation

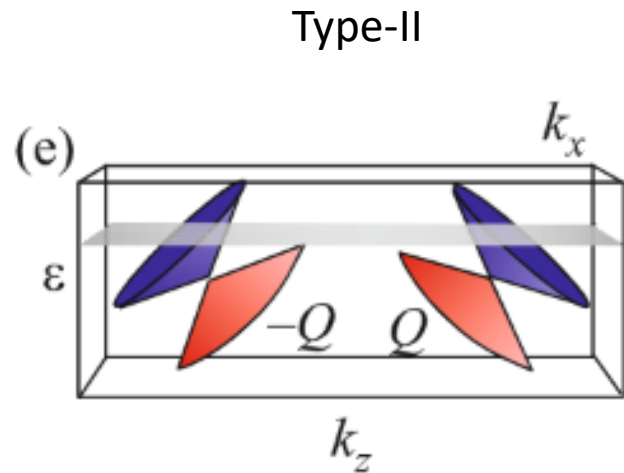


Fermi Surface contribution:

Integrate the z-component of Berry curvature over the occupied states

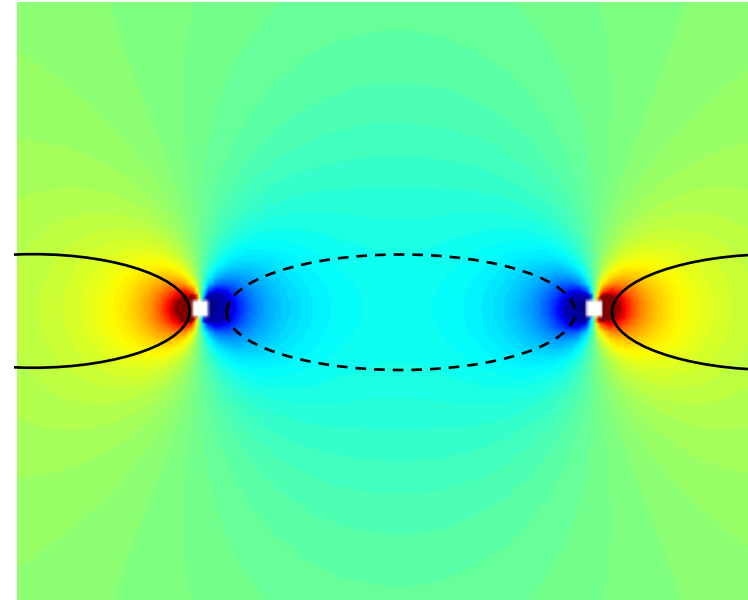
In type-I WSM the Fermi Surface contribution cancel!

W Anomalous Hall Conductivity of Ideal Weyl Semimetal



$$\sigma_{xy} = \frac{e^2 Q}{2\pi^2 \hbar}$$

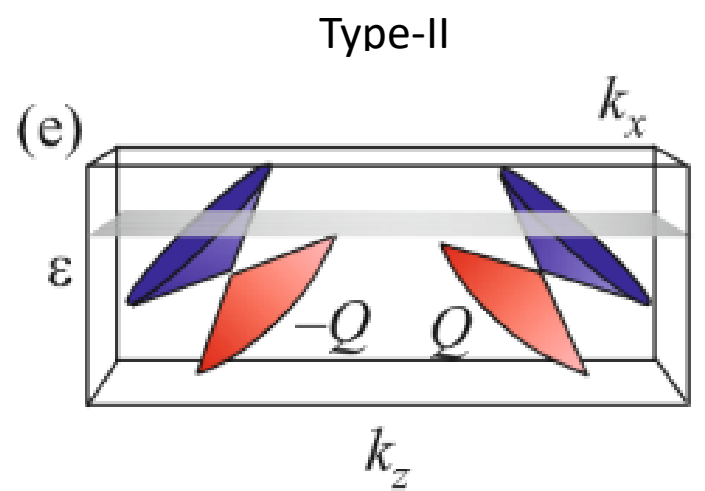
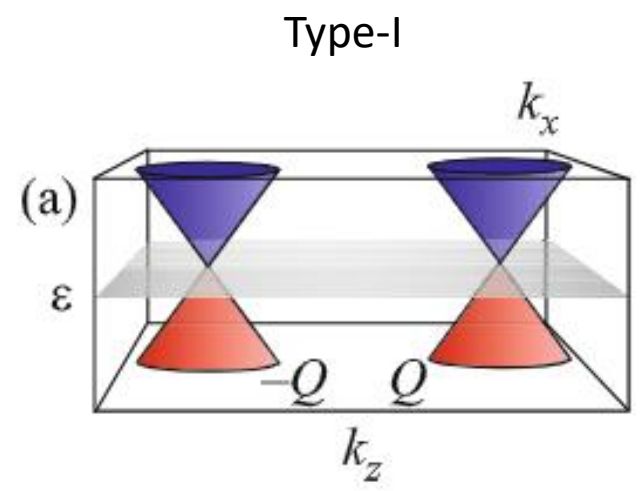
Q : Weyl points separation



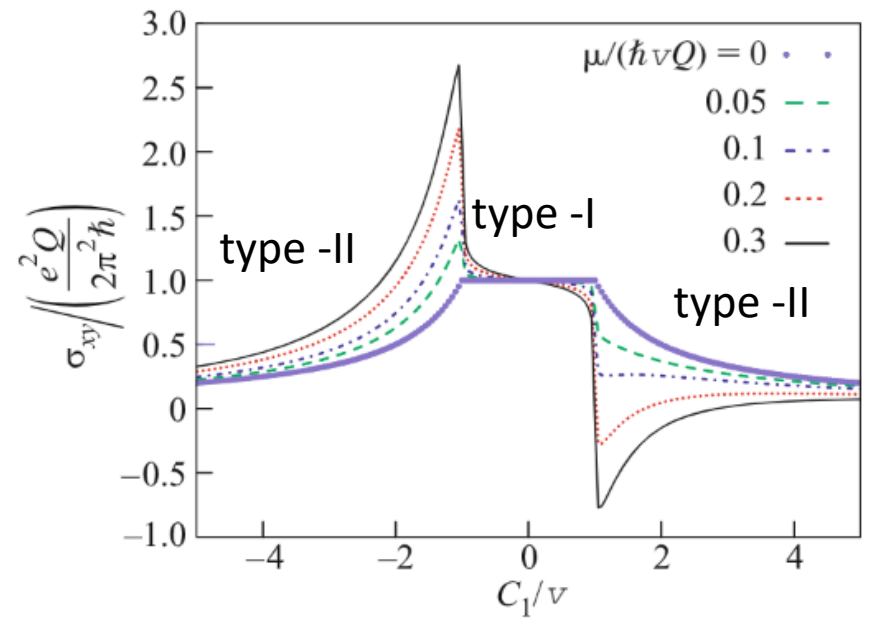
Fermi Surface contribution:

Integrate the z-component of Berry curvature over the occupied states

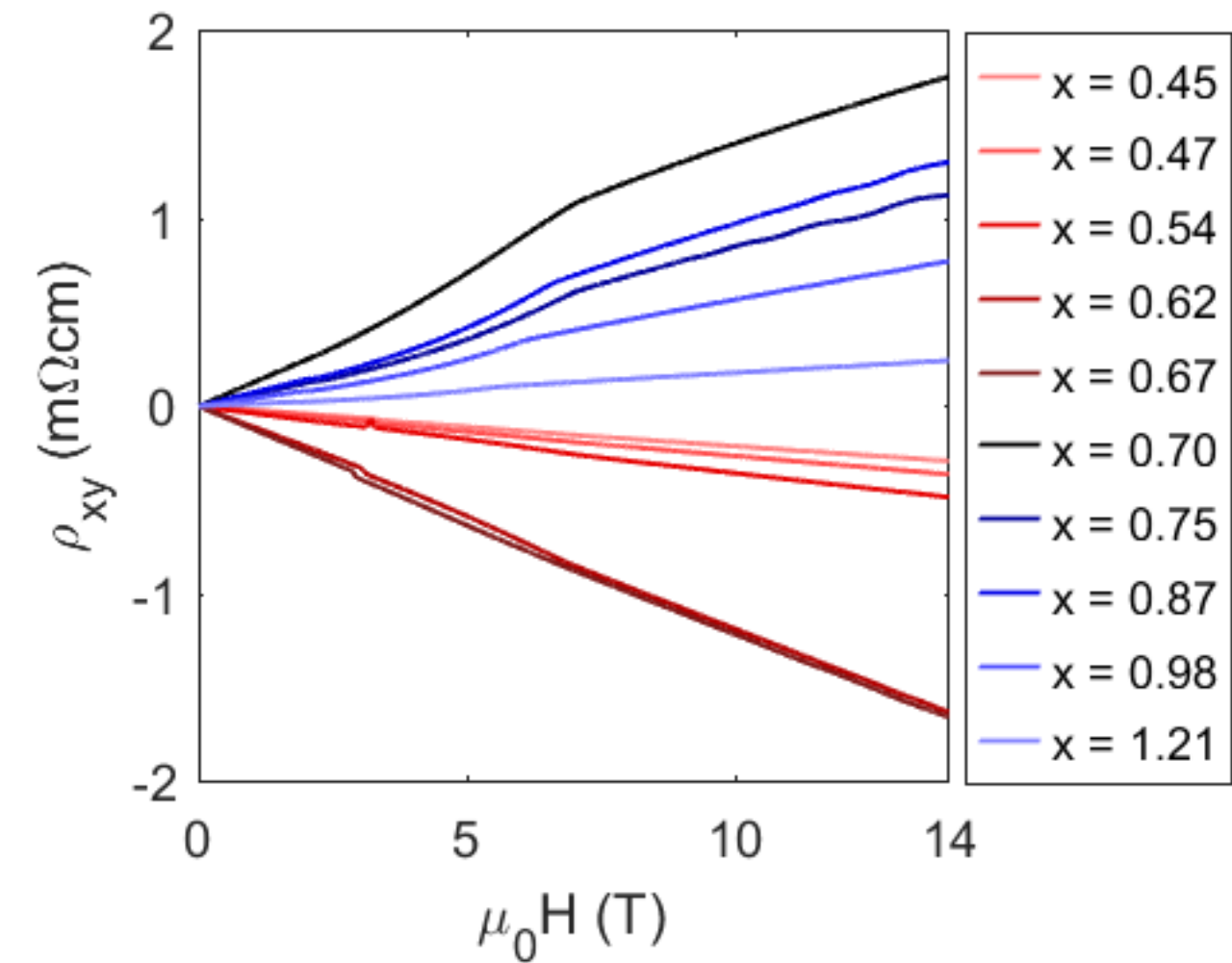
In type-II WSM the Fermi Surface contributions diverge!



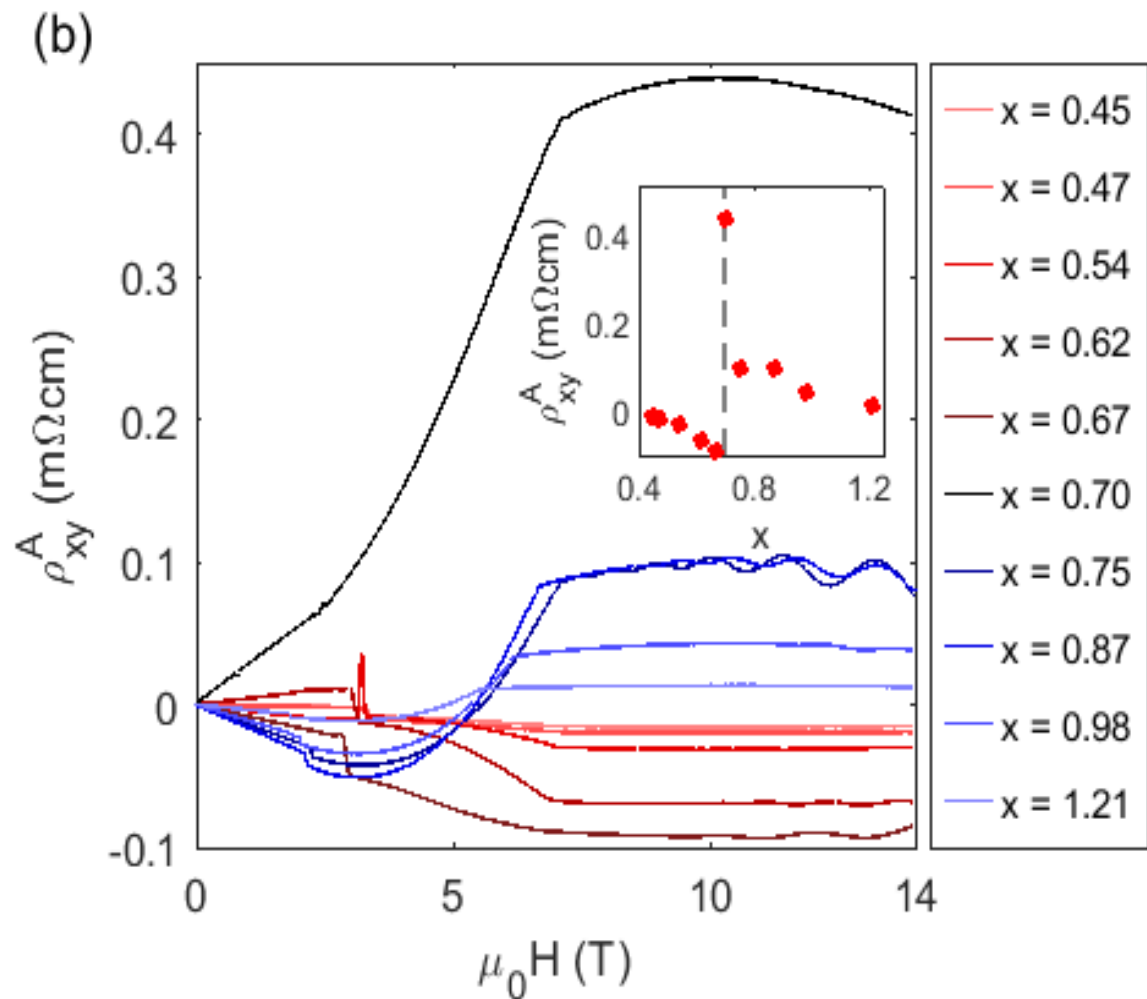
$$\sigma_{xy} = \frac{e^2 Q}{2\pi^2 \hbar}$$

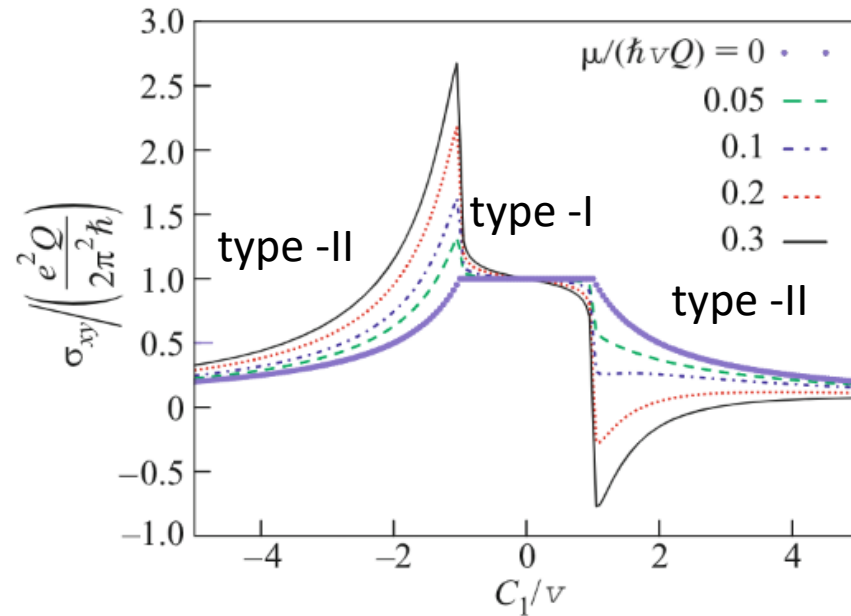
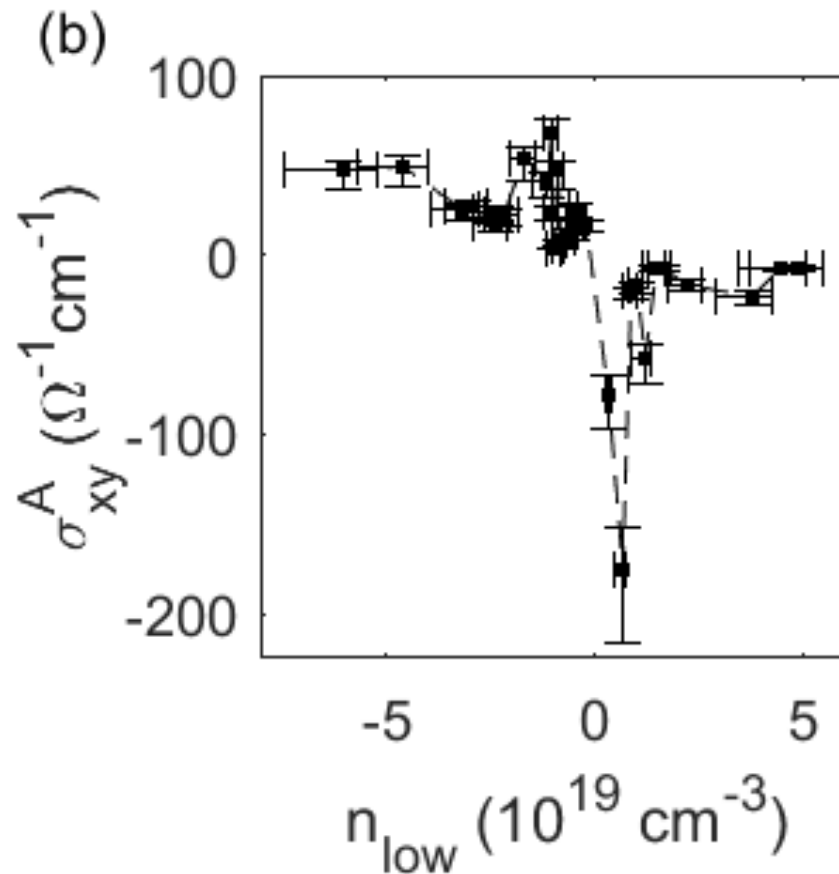


$$\sigma_{xy} = \frac{e^2}{2\pi^2 \hbar} \left[Q + \frac{\mu}{4\hbar} \left(\frac{1}{C_2} - \frac{1}{C_1} \right) \left(\ln \left| \frac{2\hbar v \Lambda}{\mu} \right| - 1 \right) \right]$$



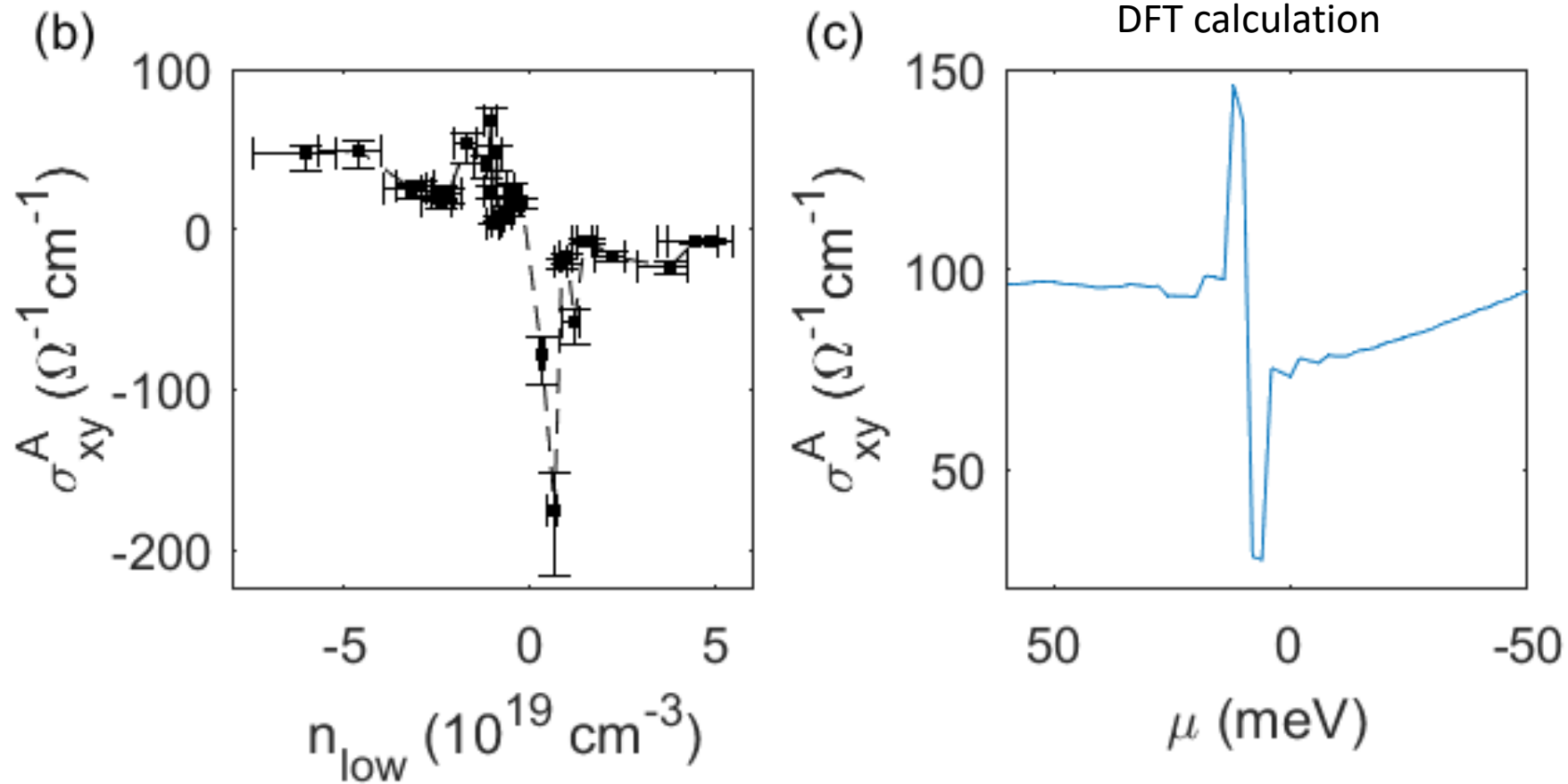
$$\rho_{xy}^A = \rho_{xy} - R_H \mu_0 H$$



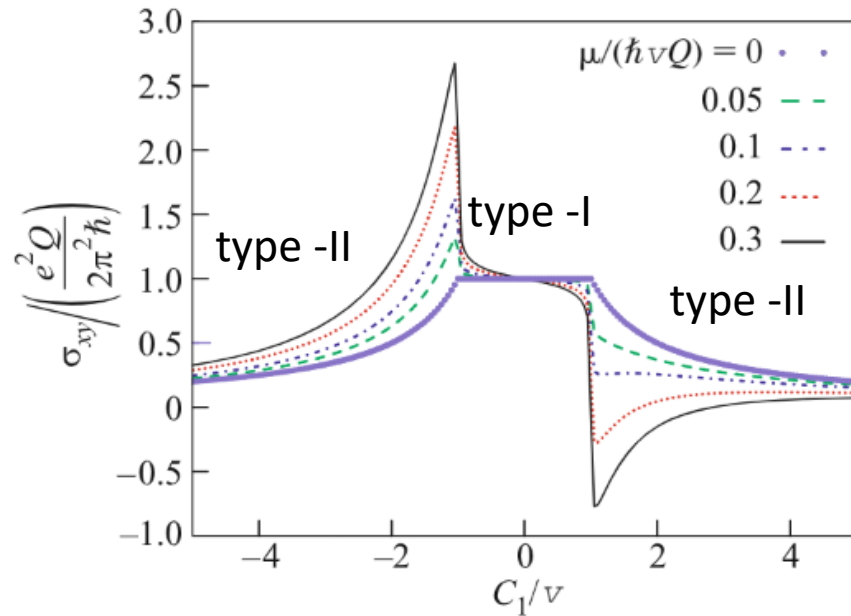
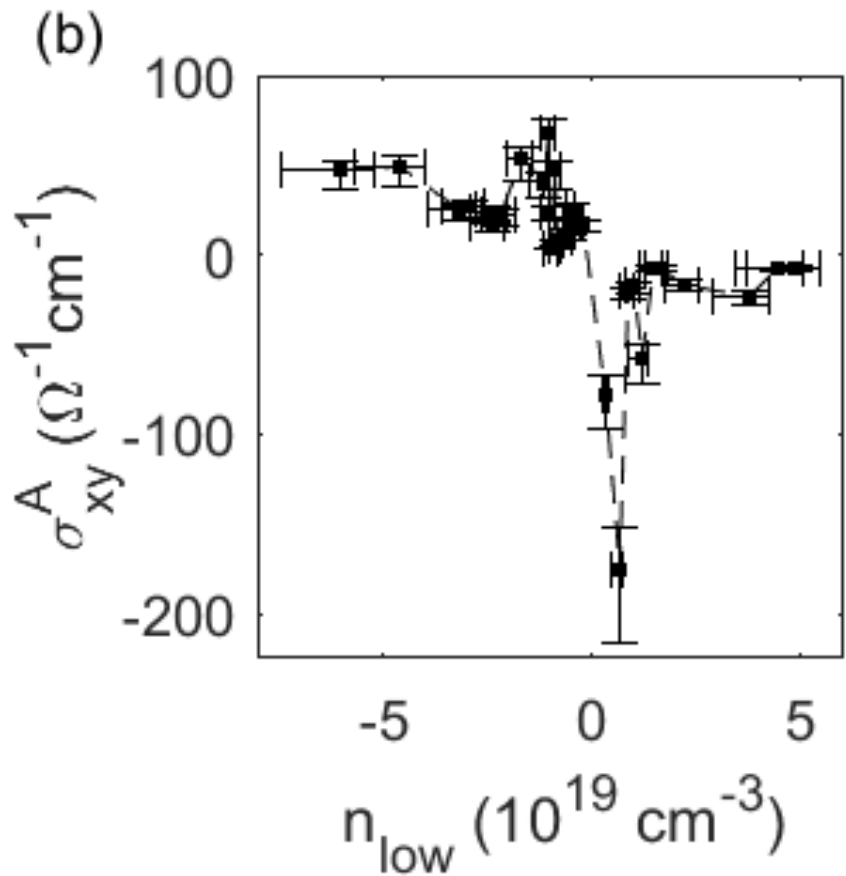


A. A. Zyuzin and R. P. Tiwari, JETP Letters, 2016, 103, 11, 717–722 (2016)

The evolution of anomalous Hall conductivity matches with the ideal type-II WSM case.



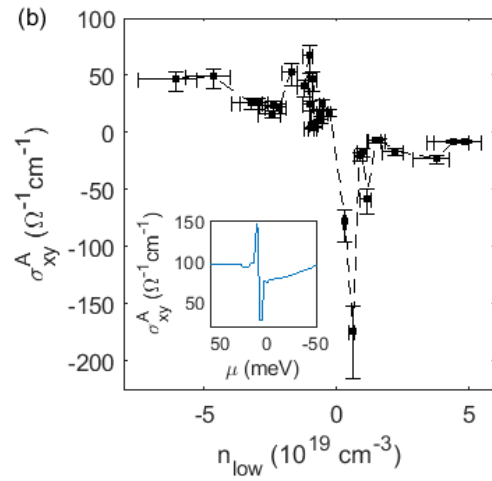
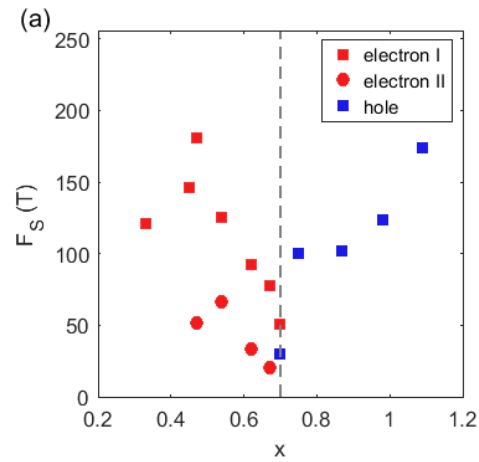
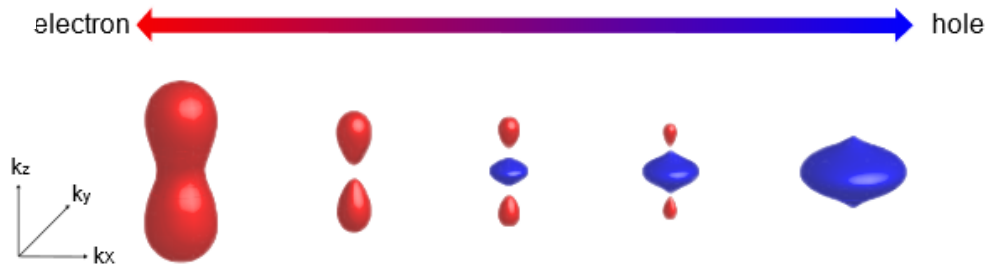
The evolution of anomalous Hall conductivity matches with the ideal type-II WSM case.



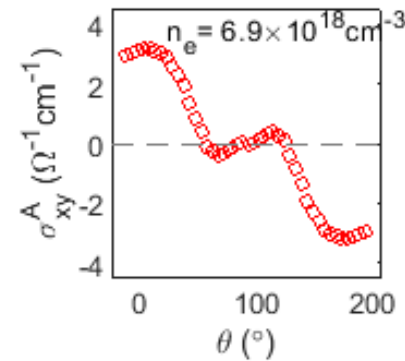
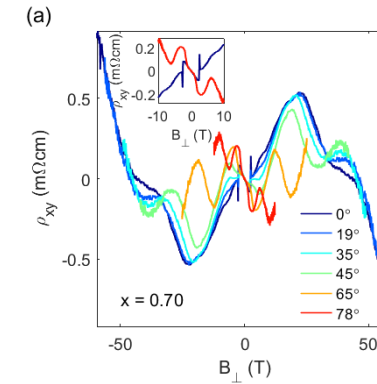
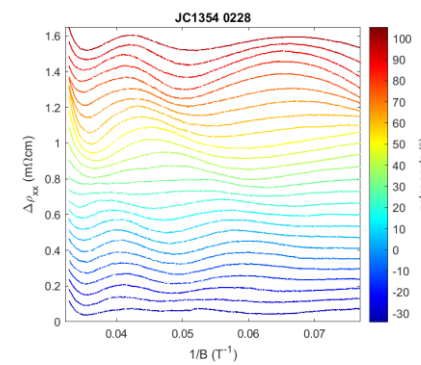
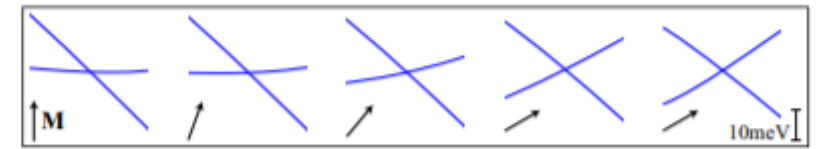
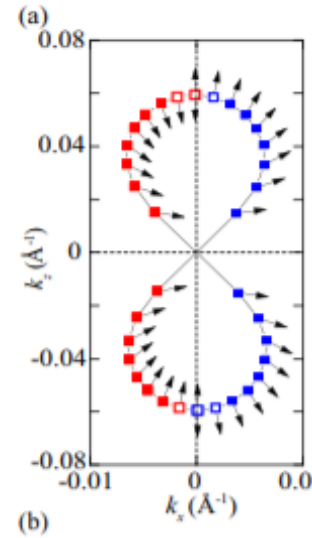
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The evolution of anomalous Hall conductivity matches with the ideal type-II WSM case.

- An ideal type-II Weyl semimetal phase in Field-induced FM $\text{MnBi}_{2-x}\text{Sb}_x\text{Te}_4$



- Evidence of a type-II to type-I WSM transition by rotating the magnetic field



Thank you!