



Topological States in Thin Films of Cadmium Arsenide

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Acknowledgements



Binghao (Evan) Guo



Alex Lygo

Theory Collaborators:

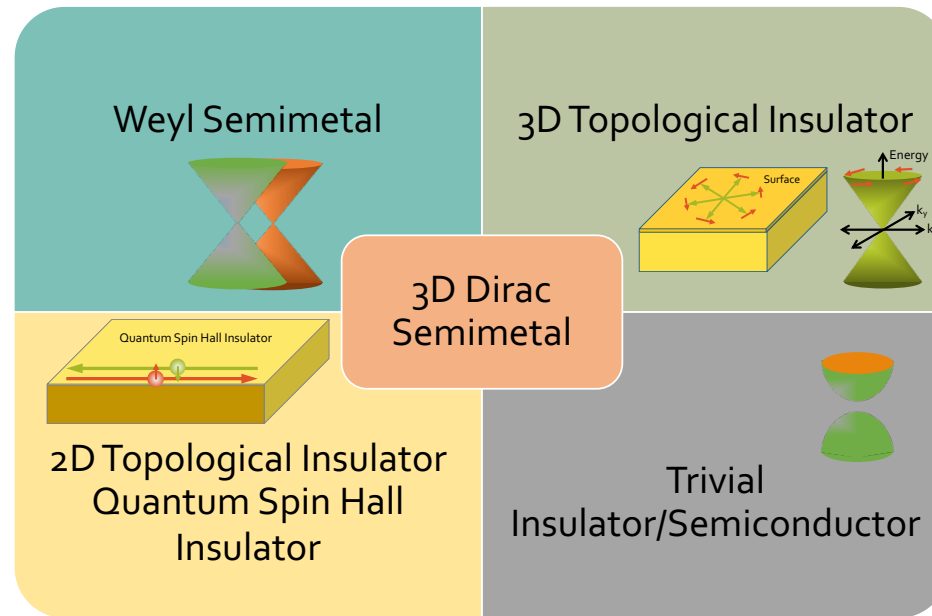
- Xi Dai (UCSB and HKUST)
- Wangqian Miao (UCSB)

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UCSB Quantum Foundry (NSF)
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Outline

- Cadmium arsenide films as a platform for creating different topological phases
- 2D TI state
- 2D Weyl semimetal
- Summary

Why topological semimetal thin films?



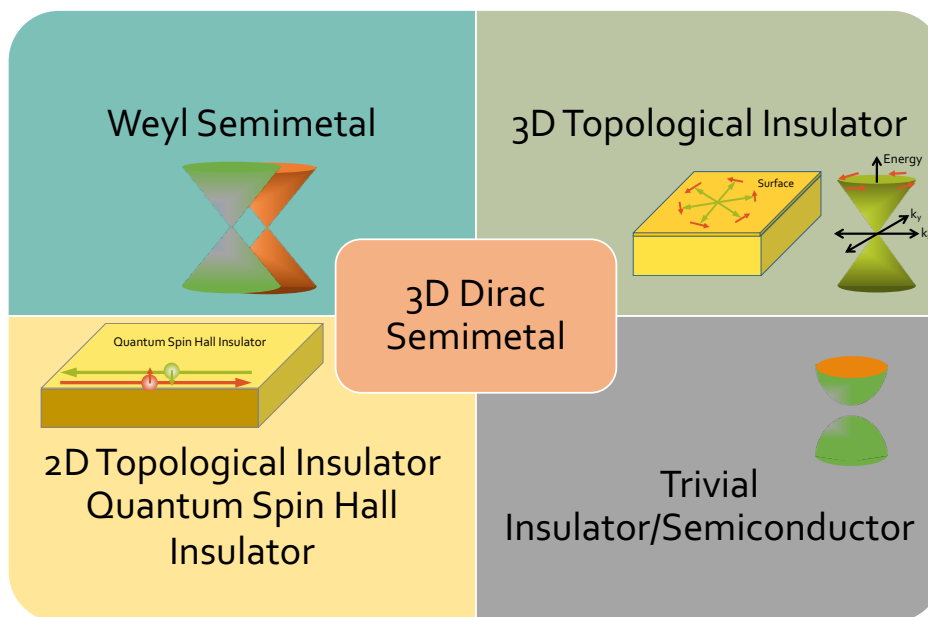
S. M. Young et al., Phys. Rev. Lett. 108, 140405 (2012).

3D Dirac semimetal thin films can be used to tune between topological phases, and manipulate and control their electronic states

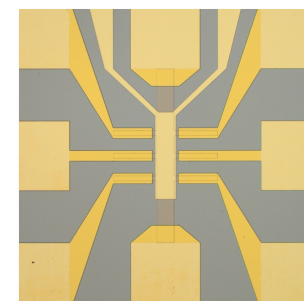
Why topological semimetal thin films?

No "ideal" WSMs:
many nodes, trivial
bands crossing Fermi
level obscure physics

Few good QSHI
candidates & lots of
open questions about
the existing QSHI



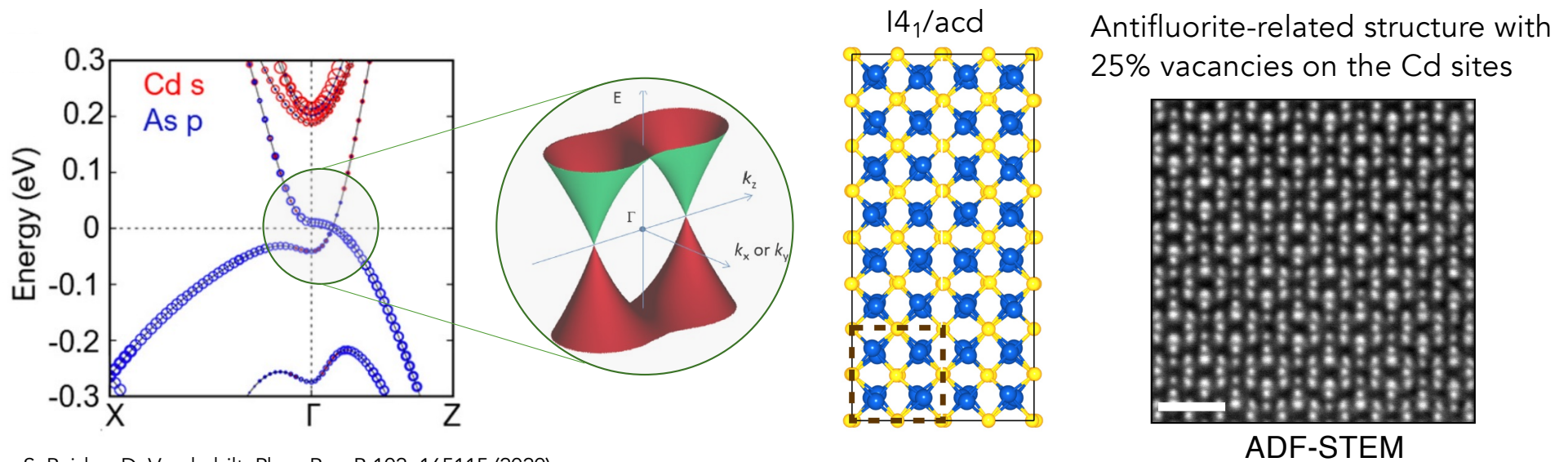
2D materials allow
control via electric field



S. M. Young et al., Phys. Rev. Lett. 108, 140405 (2012).

3D Dirac semimetal thin films can be used to tune between topological phases, and manipulate and control their electronic states

Three-dimensional Dirac Semimetal: Cadmium Arsenide (Cd_3As_2)



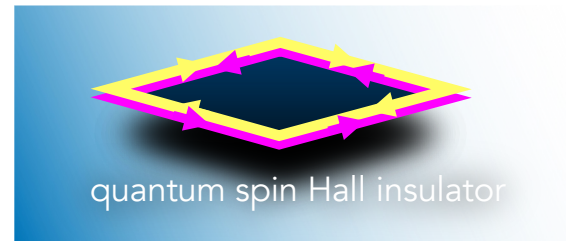
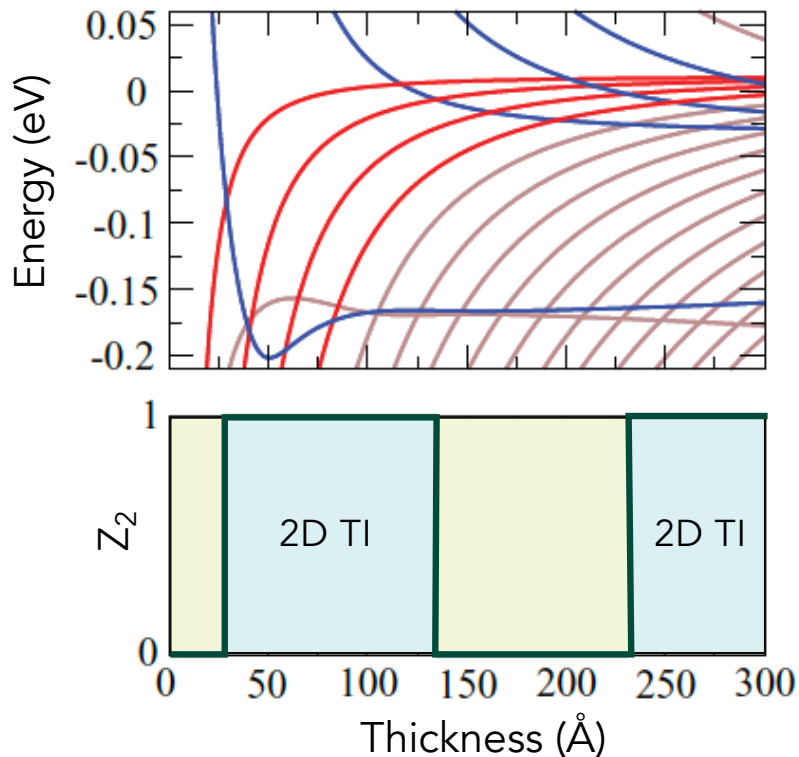
S. Baidya, D. Vanderbilt, Phys. Rev. B 102, 165115 (2020).

- Band inversion between Cd-5s and As-4p states at Γ
- Two crossings along Γ -Z protected by 4-fold symmetry
- High Fermi velocity
- High carrier mobilities

Wang et al., Phys. Rev. B 88, 125427 (2013).
Ali et al., Inorg. Chem. 53, 4062 (2014).
Borisenko, et al. Phys. Rev. Lett. 113, 165109 (2014).
Neupane et al., Nat. Commun. 5, 3786 (2014).
Liu et al., Nat. Mater. 13, 677 (2014).
Jeon et al., Nat. Mater. 13, 851 (2014).
Mosca Conte et al., Sci. Rep. 7, 45500 (2017).
A. J. Rosenberg, T. C. Harman, J. of Appl. Phys. 30, 1621 (1959).

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Topological States of Cadmium Arsenide: Quantum Spin Hall Insulator



“A 5 nm thick Cd₃As₂ film is a good QSH insulator with a gap of more than 100 meV”

PHYSICAL REVIEW B **88**, 125427 (2013)

Three-dimensional Dirac semimetal and quantum transport in Cd₃As₂

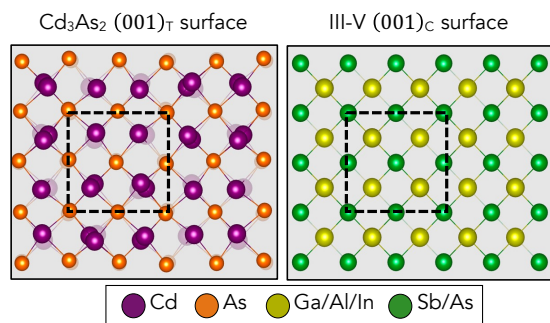
Zhijun Wang, Hongming Weng,^{*} Quansheng Wu, Xi Dai, and Zhong Fang[†]

Beijing National Laboratory for Condensed Matter Physics, and Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China

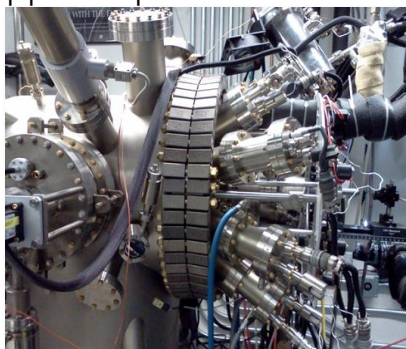
(Received 29 May 2013; revised manuscript received 5 September 2013; published 23 September 2013)

Based on the first-principles calculations, we recover the silent topological nature of Cd₃As₂, a well known semiconductor with high carrier mobility. We find that it is a symmetry-protected topological semimetal with a single pair of three-dimensional (3D) Dirac points in the bulk and nontrivial Fermi arcs on the surfaces. It can be driven into a topological insulator and a Weyl semimetal state by symmetry breaking, or into a quantum spin Hall insulator with a gap more than 100 meV by reducing dimensionality. We propose that the 3D Dirac cones in the bulk of Cd₃As₂ can support sizable linear quantum magnetoresistance even up to room temperature.

Molecular Beam Epitaxy of (001)Cd₃As₂



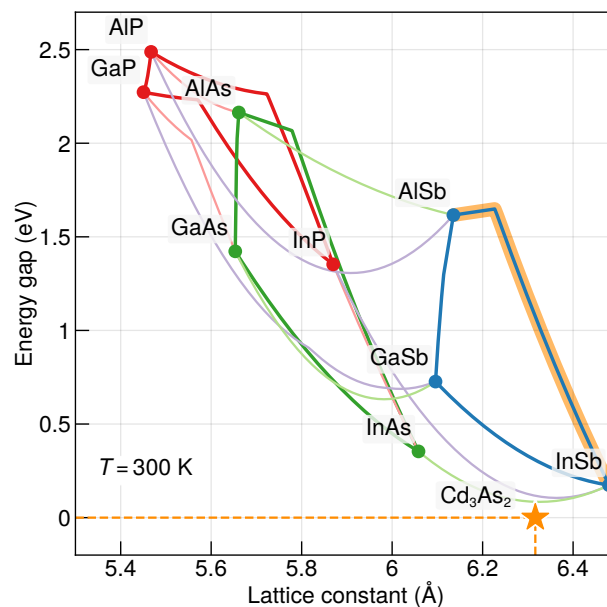
Applied Epi Modified GEN 2



- T. Schumann et al., APL Mater. 4, 126110 (2016).
- T. Schumann et al., Phys. Rev. B 95, 241113(R) (2017).
- D. A. Kealhofer et al. Phys. Rev. Mater. 3, 031201(R) (2019).
- M. Goyal et al., APL Mater., 8, 051106 (2020).

High quality epitaxy requires:

- Symmetry matched + chemically similar substrate
- GaSb substrate + buffer layer: Al_{0.45}In_{0.55}Sb ternary alloy
- GaSb cap

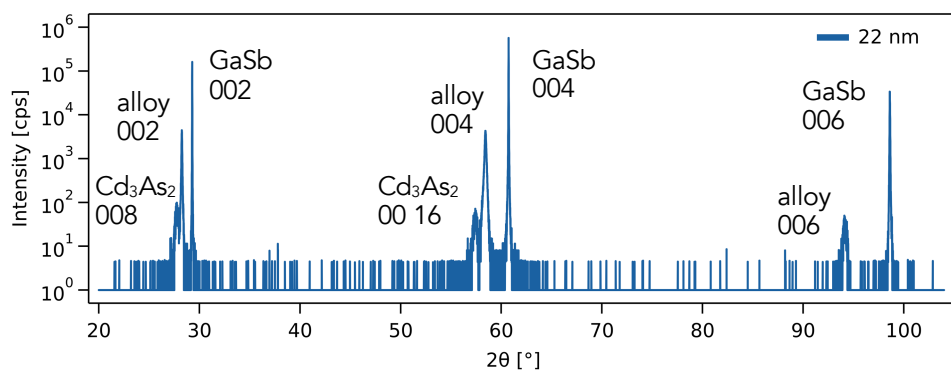


AlO _x	25 nm
GaSb	3 nm
Cd ₃ As ₂ QW	18 nm to 22 nm
Al _{0.45} In _{0.55} Sb	1500 nm
GaSb	50 nm
GaSb (001) substrate	

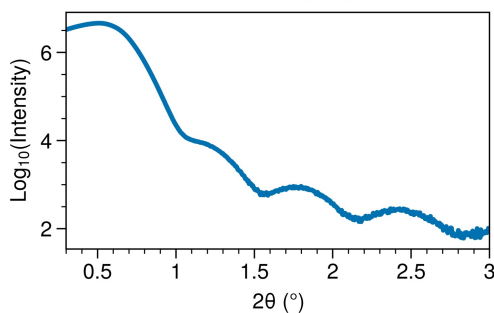
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Molecular Beam Epitaxy of (001)Cd₃As₂

Wide-angle, out-of-plane 2θ-ω scan

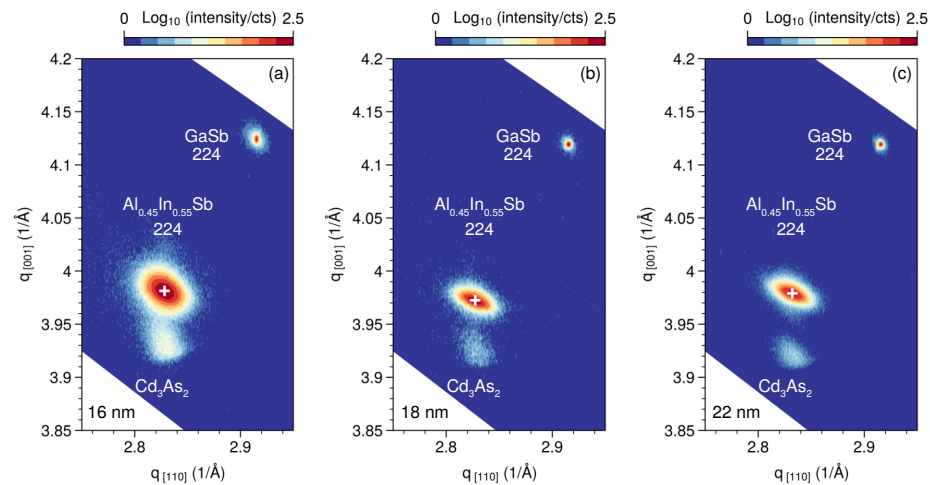


10 nm Cd₃As₂

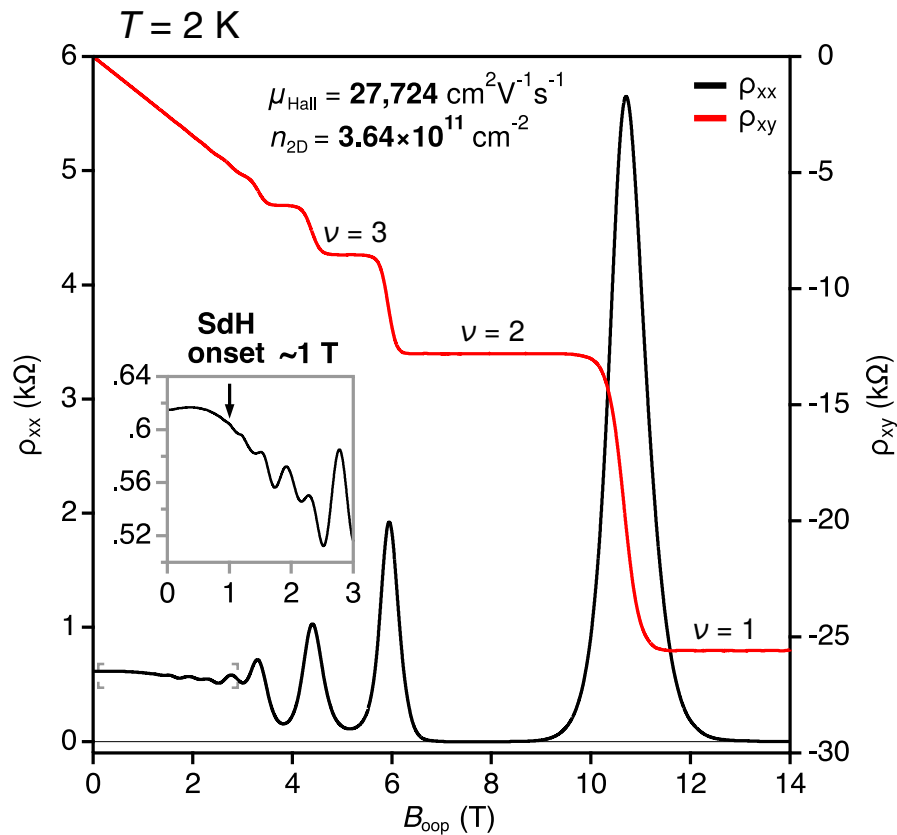


X-ray reflectivity:
smooth Cd₃As₂ layers
even at 10 nm

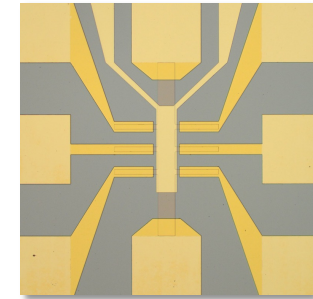
Reciprocal space maps show that Cd₃As₂ is coherently strained to buffer alloy



Molecular Beam Epitaxy of (001)Cd₃As₂



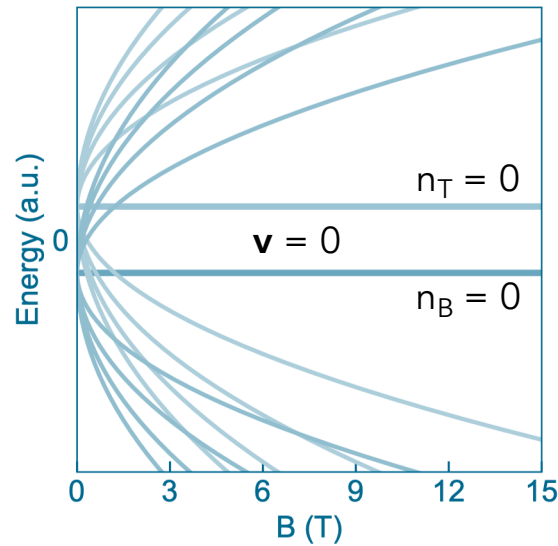
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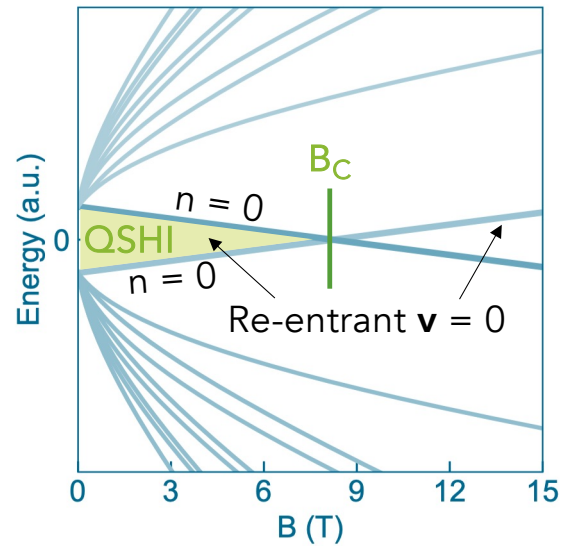
- Lithographically defined Hall bar
- Around zero field: Hall mobility, area density
- Low field: estimate quantum mobility from onset of Shubnikov-de Haas oscillations at 1 T
- High field: formation of Landau levels + quantized plateaus in Hall resistivity, with simultaneously vanishing longitudinal resistivity

Topological Insulators in a **Perpendicular** Magnetic Field

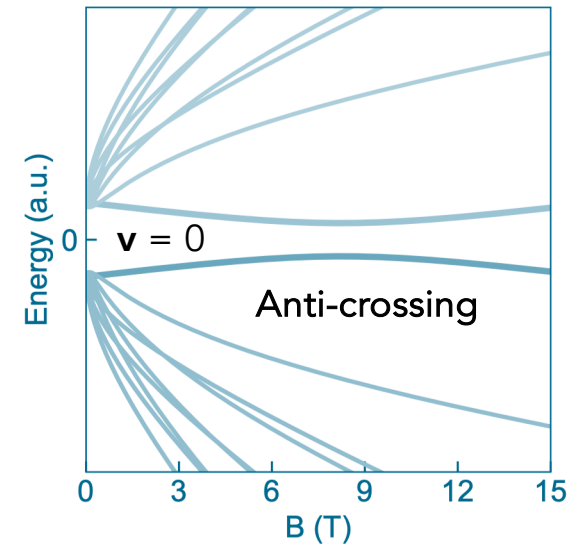
3D Topological Insulator
with energy offset
between the top and
bottom surfaces



2D Topological Insulator
no SIA

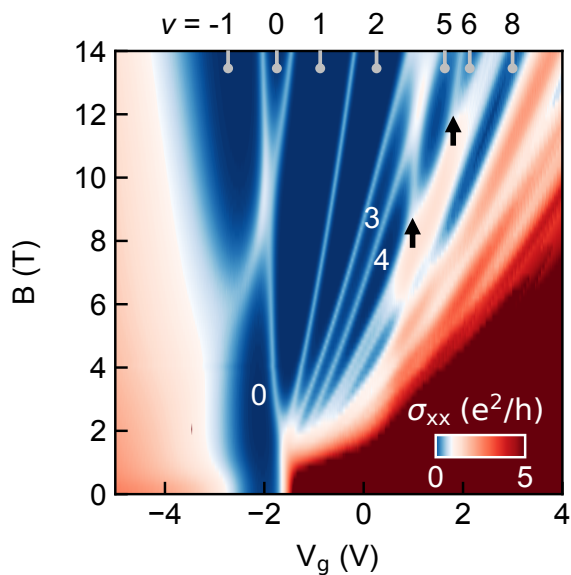


2D Insulator
(un-inverted due to SIA)

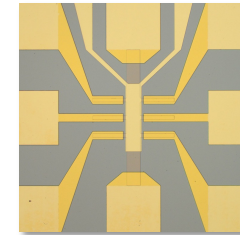
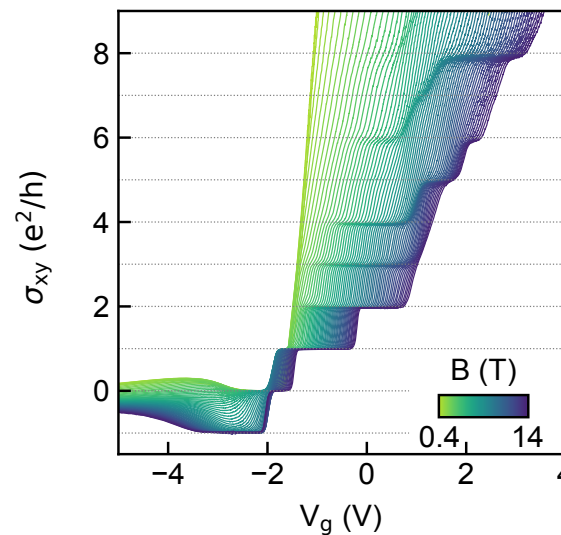


Landau levels and quantum Hall effect in a 20 nm film

Color map of σ_{xx} shows the evolution of Landau levels with field and gate voltage



Quantum Hall plateaus

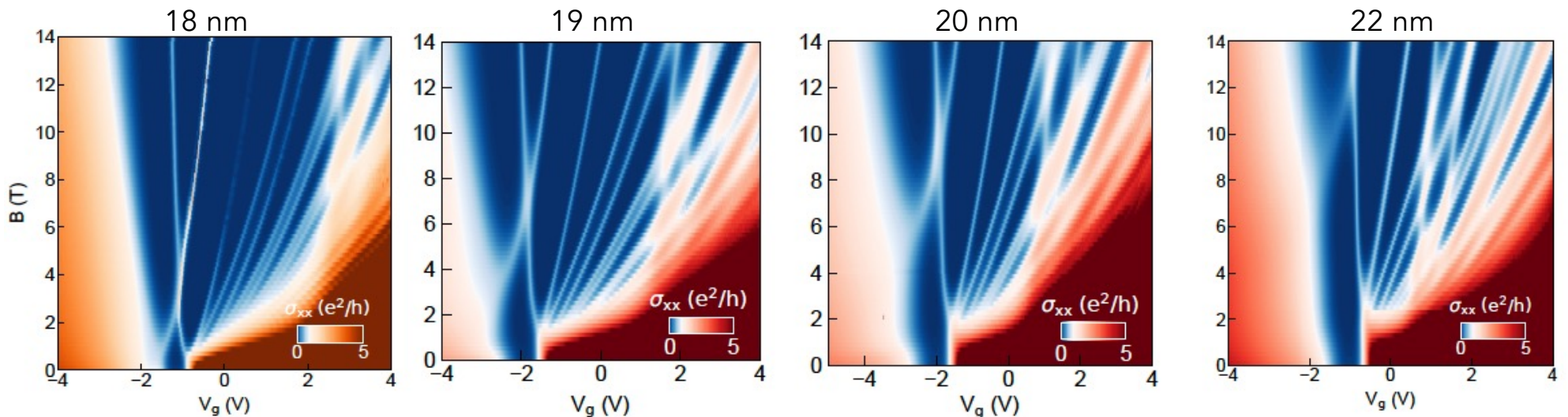


Gated Hall bar

- Signature of inverted spectrum: crossing of $n = 0$ Landau levels at a critical field
- Re-entrant $\nu = 0$ plateau due to crossing of zeroth Landau levels
- Additional higher energy subband

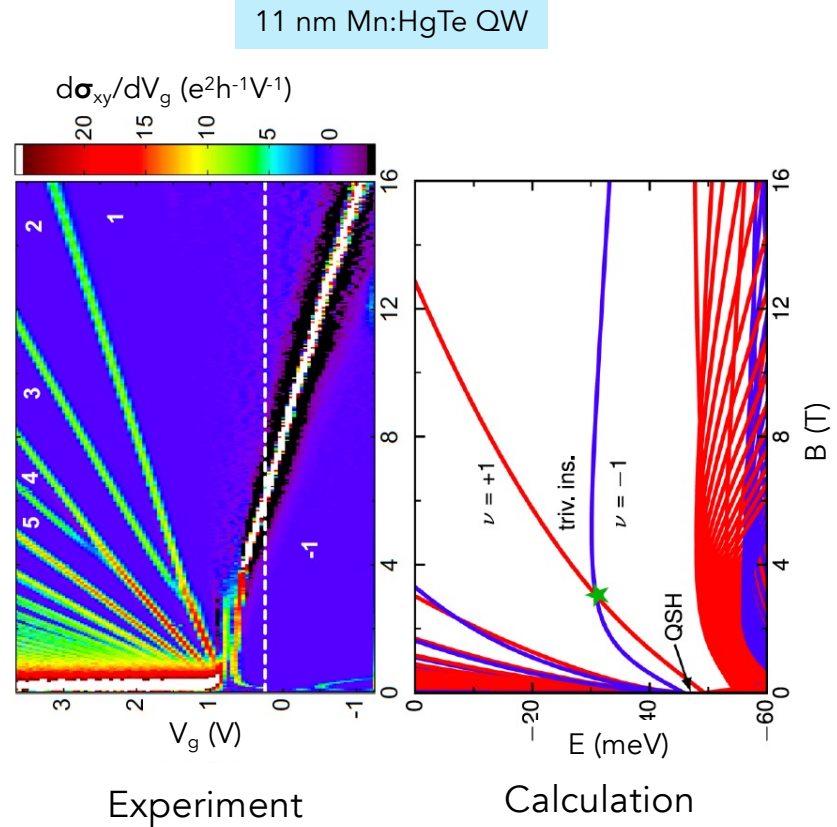
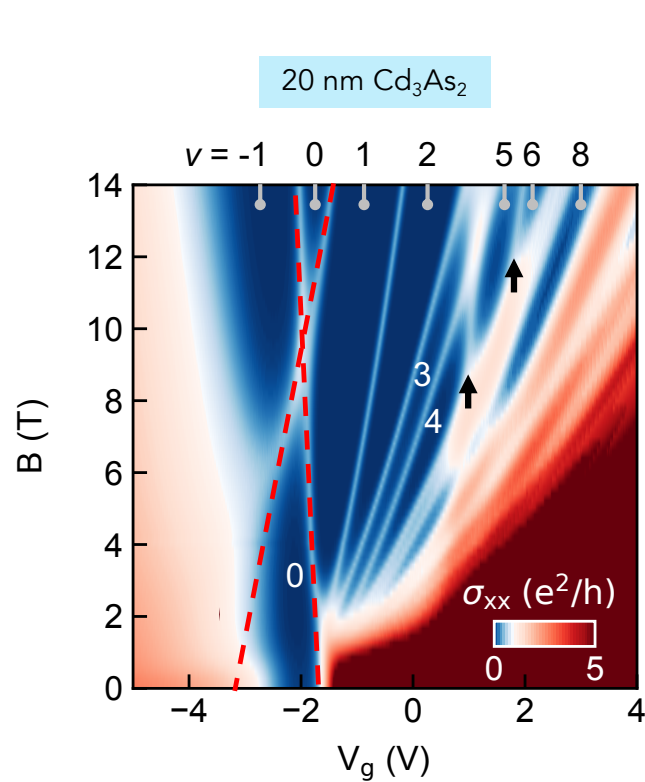
Landau Level Spectrum: Thickness Dependence

Color maps of σ_{xx} show the evolution of Landau levels



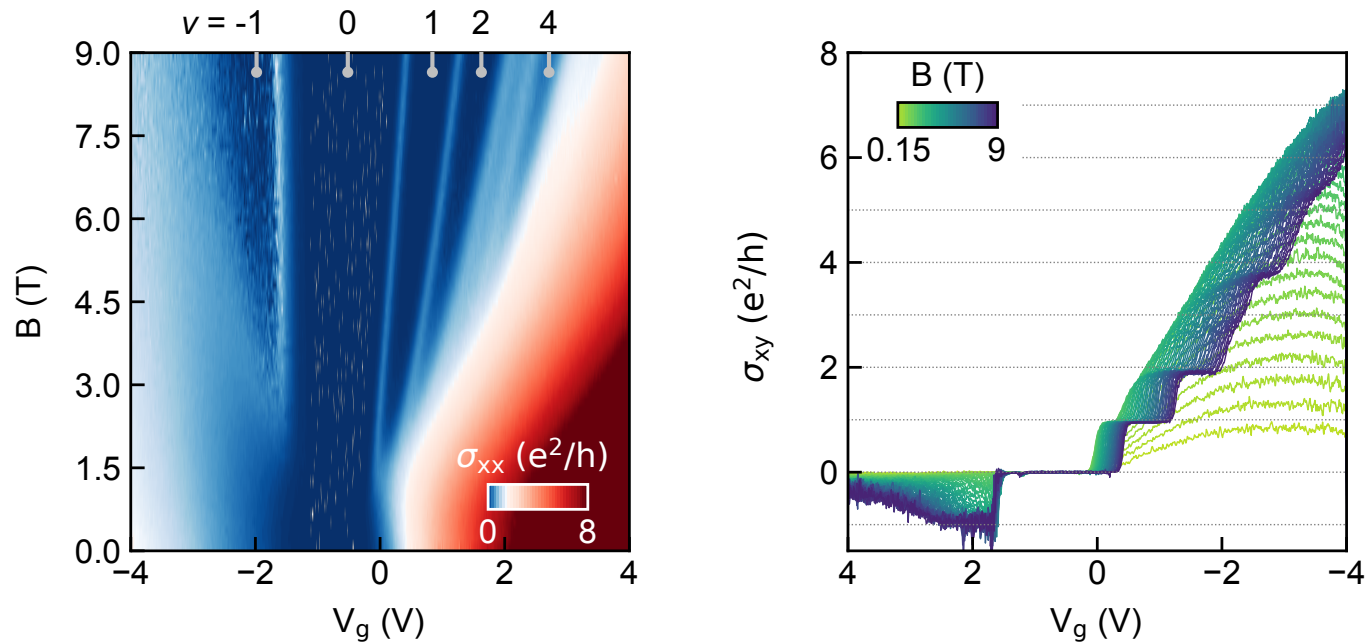
- Signature of inverted spectrum: crossing of $n = 0$ Landau levels at a critical field
- Re-entrant $\nu = 0$ plateau due to crossing of zeroth Landau levels
- Systematic behavior with thickness
- Second conduction subband moves out to higher energies with decreasing thickness

Comparison with HgTe Quantum Wells



S. Shamim, ..., L. W. Molenkamp, Nat. Comm. 13, 2682 (2022)

Trivial 2D Insulator: 14 nm film



Electron and hole-like zeroth Landau levels are “un-inverted” and do not cross
Film thickness can be used to tune between topological and trivial states, just as predicted

Cadmium arsenide as a novel two-dimensional topological insulator

- “Textbook” 2D TI seen in LL spectroscopy
- Should host helical edge states
- Mobilities are high for a near surface electronic system
- Facilitates hybrid structures with superconductors
- Tunability by strain and film thickness
- Good platform for novel quantum devices

Two-dimensional topological insulator in an **in-plane magnetic field**

- Out-of-plane field is responsible for the orbital effect and Landau quantization
- In-plane magnetic field: main effect on energy bands is the Zeeman energy:

$$\Delta E_Z \sim g\mu_B B_{ip}$$

- The g-factor is sample/material specific
- In narrow-gap semiconductors, like a 2D TI, the Zeeman field may be sufficient to close the band gap

S. Sun, ... X. Dai, Phys. Rev. B 106, L241105 (2022).

F. Dominguez, ... E. M. Hankiewicz, SciPost Phys. Core 5, 024 (2022).

O. E. Raichev, Phys. Rev. B 85, 045310 (2012).

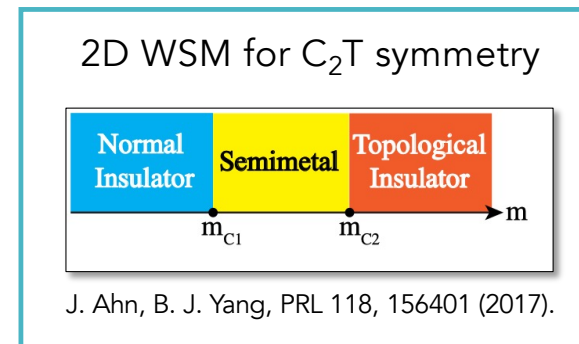
M. Kubisa, K. Ryczko, Phys. Rev. B 104, L161406 (2021).

A. A. Zyuzin, ... A. A. Burkov, Phys. Rev. B 83, 245428 (2011).

Symmetry considerations for a gapless phase in Cd_3As_2 :

- Zero-field crystallographic point group: $4/mmm$
- With in-plane field, magnetic point group: $2'/m'$

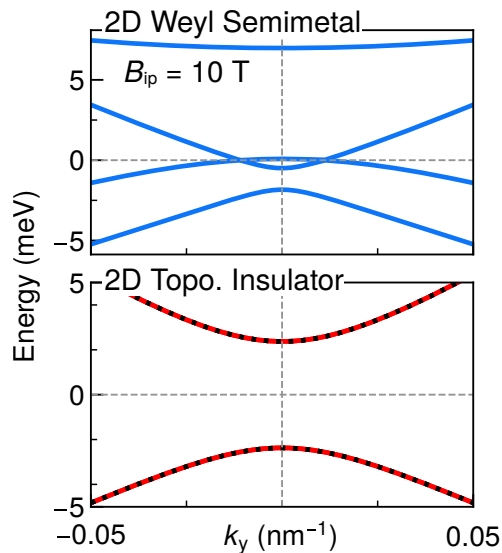
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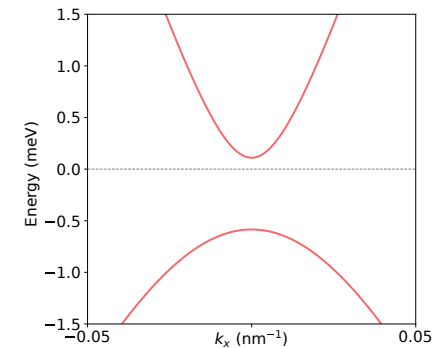
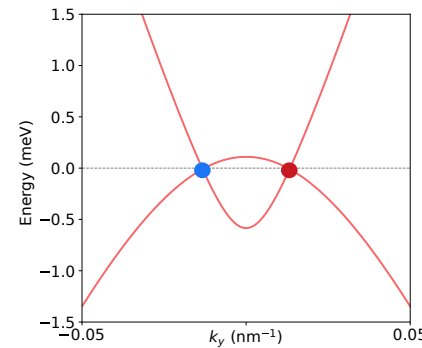
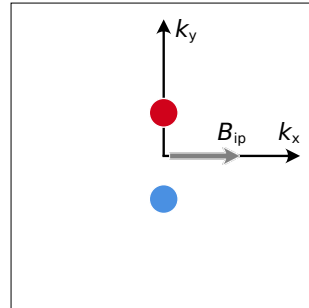
C_2T symmetry guarantees that the Berry curvature to be zero everywhere except at the Weyl points.

Two-dimensional topological insulator in an **in-plane magnetic field**

Ideal 2D Weyl semimetal, with two nodes at the charge neutrality level



B. Guo et al., submitted



Calculations by Wangqian Miao and Xi Dai for an 18 nm (001) Cd₃As₂ film
Requires accurate g -factor calculations that consider the combined effects of remote bands and quantum confinement*

* Z. Song,... X. Dai, arXiv:1512.05084 (2019).

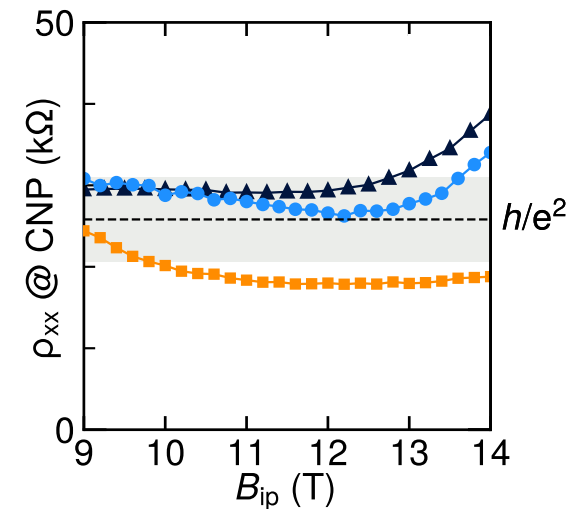
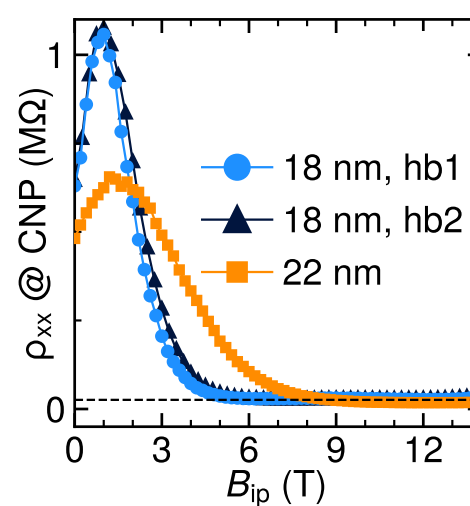
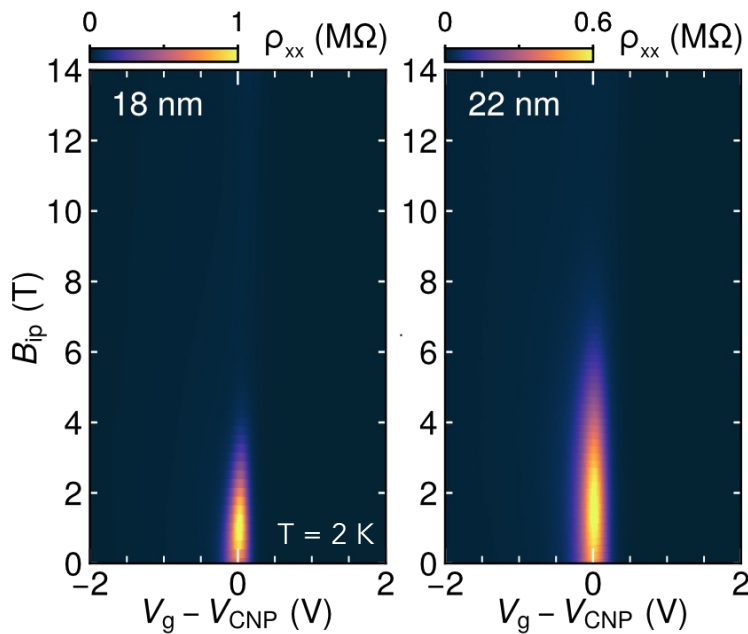
S. Sun,... X. Dai, Phys. Rev. B 101, 125118 (2020).

Zeeman field-induced 2D Weyl Semimetal: Experiment

Gap closing and metallic phase

- Non-monotonic evolution of peak resistivity with in-plane field
- Band gap closing: sharp drop past 1 T

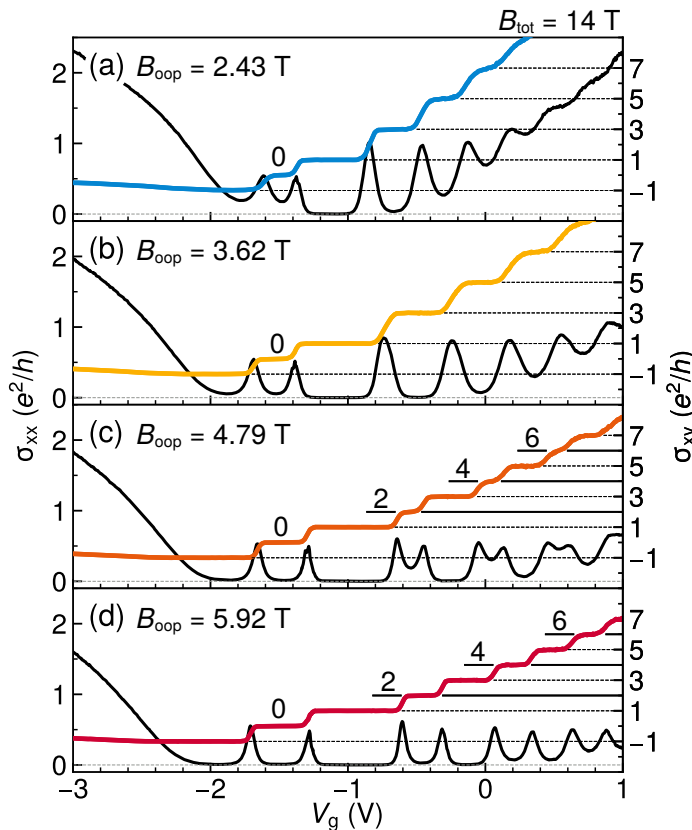
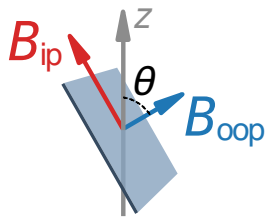
- Initial increase and peak possibly related to current-carrying edge states
- Resistivities on the order of $\sim h/e^2$ in a wide range of B-fields
- Evidence for the symmetry-protection of the metallic phase



Peak resistivities at the charge neutrality point

Zeeman field-induced 2D Weyl Semimetal: Experiment

Odd-integer quantum Hall effect



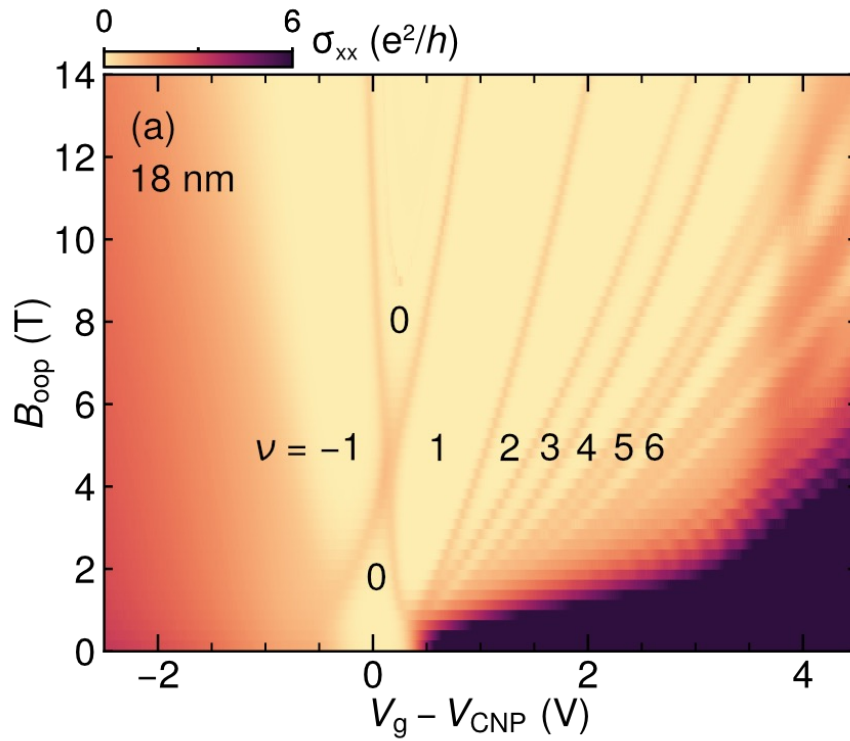
Sample:

- 18 nm QW
- T = 2 K

- The two $n = 0$ Landau levels are spin resolved
- Connect the $\nu = 0$ and $\nu = \pm 1$ plateaus
- $|n| > 1$ Landau levels are **doubly degenerate** and produce an **odd integer quantum Hall effect**
- Odd-integer QHE implies the existence of a pair of massive Weyl fermions:
 $\nu = (n + 1/2)$ per node
- Valley degeneracy is lifted at higher out-of-plane fields; requires further investigation

Comparison with oop field experiments

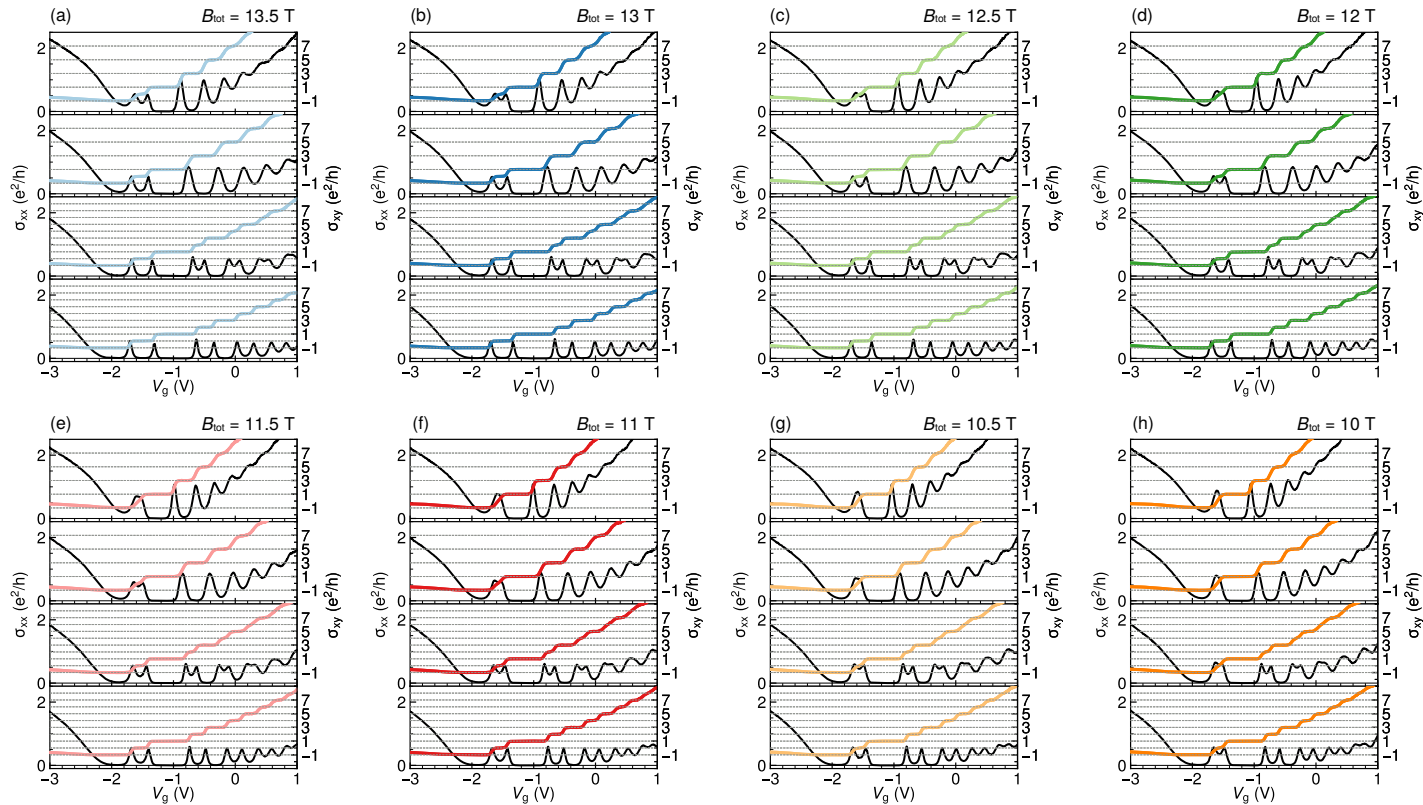
Out-of-plane magnetic field, no in-plane field: 2D TI phase
QHE with all degeneracies are lifted



- The filling factor $\nu = 0$ state appears twice because of crossing Landau levels
- All degeneracies are lifted: even and odd integer filling factors
- Spin degeneracy is immediately lifted

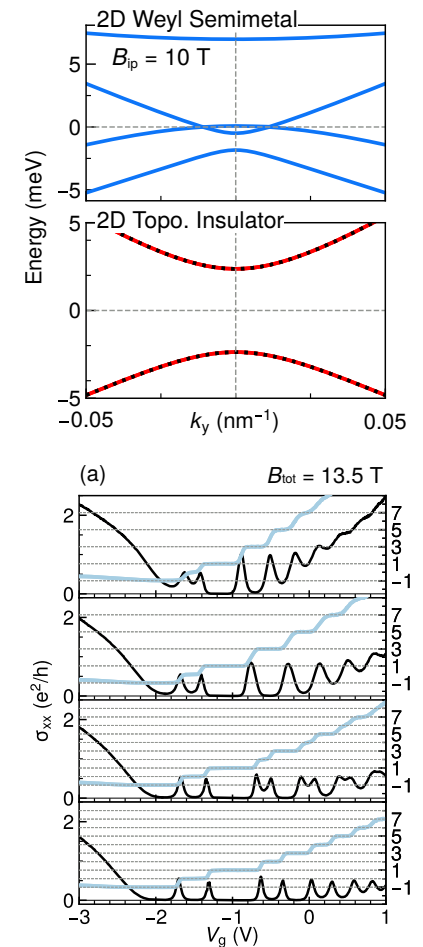
Zeeman field-induced 2D Weyl Semimetal: Experiment

Odd-integer quantum Hall effect



Two-dimensional Weyl semimetal phase

- Strong in-plane Zeeman field converts the 2D TI into a 2D WSM
- Ideal WSM
 - Only two Weyl nodes
 - Nodes at the charge neutrality point
 - No other bands
- Odd-integer quantum Hall effect
- Symmetry protected
- Host for other topological states, e.g. combination with superconductors



Summary

- Cd_3As_2 is a platform for creating, and controlling different topological phases
- “Ideal” 2D TI
- “Ideal” 2D Weyl semimetal
- Tunability by strain, film thickness, magnetic fields
- Excellent transport properties
- Promising platform for novel quantum devices

Thank you!