

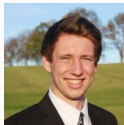
# Efficient method for quantum impurity problems out of equilibrium

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July 22, 2023



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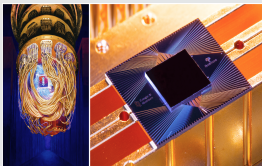
# Quantum systems out of equilibrium

## Thermalization



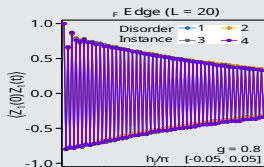
ETH or MBL?

## Quantum computing



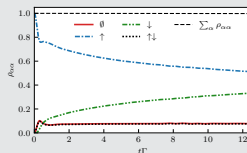
Quantum advantage?

## Dynamical phases



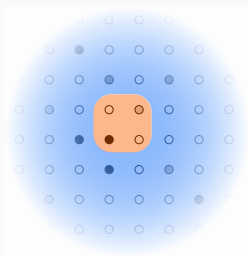
Time crystals, ...

## Impurity models



Kondo physics, DMFT,  
NMR

# Real time evolution of local observables



$$\langle O(t) \rangle_{\rho(t=0)=\rho_0} \quad \text{or} \quad \langle O(0)O(t) \rangle_{\rho_0}$$

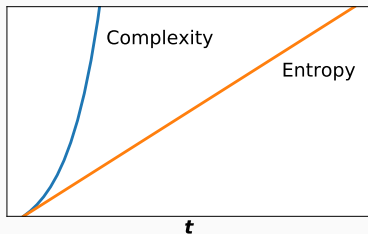
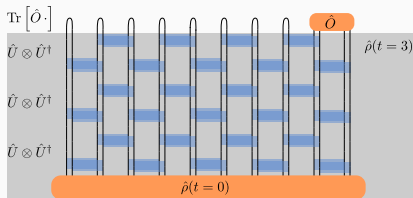
- Measurable quantity
- Transport properties: (Sub)diffusion vs ballistic
- Localization, Edge modes

# Wavefunction based approaches

- Represent the density matrix at different time  $\rho(t)$
- Obtain observables by computing  $\text{Tr}[\rho(t)O]$

## Problem:

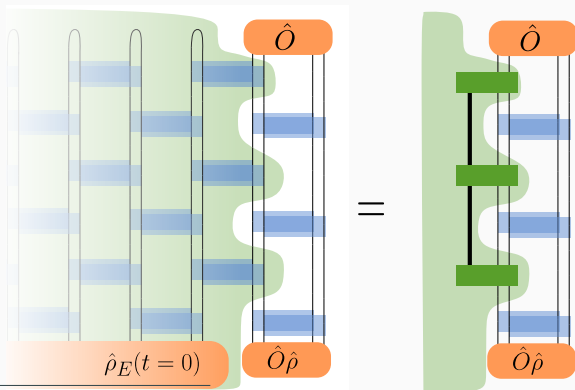
Complexity grows exponentially with time!



# Influence functional approach<sup>1</sup>

## Idea:

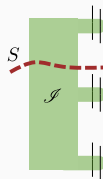
Describe the local subsystem as a quantum system coupled to a bath.



<sup>1</sup>Feynman and Vernon Jr 1963.

# Temporal Entanglement

The numerical power of the Influence matrix approach is limited by **temporal entanglement** instead of spatial entanglement



- Represents the "Memory" of the quantum bath
- Low temporal entanglement  $\Rightarrow$  efficient MPS representation
- Examples for low temporal entanglement:
  - Very chaotic systems  $U \approx e^{i\frac{\pi}{4}Z_i Z_{i+1} + hZ_i} e^{i\frac{\pi}{4}X_i}$
  - Close to integrability  $U \approx e^{ij_x X_i X_{i+1} + ij_y Y_i Y_{i+1}}$
  - Dissipative system<sup>3</sup>
  - Many body localized systems - Exact disorder averaging!<sup>4</sup>

<sup>1</sup>Lerose, Sonner, and Abanin 2021a.

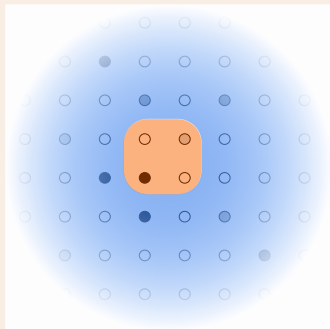
<sup>2</sup>Lerose, Sonner, and Abanin 2021b; Giudice et al. 2021; Thoenniss, Lerose, and Abanin 2022.

<sup>3</sup>Sonner, Lerose, and Abanin 2021; Mi et al. 2022.

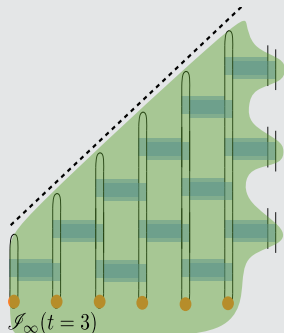
<sup>4</sup>Sonner, Lerose, and Abanin 2022.

# Obtaining the Influence matrix

## Free fermions



## One dimensional models<sup>1</sup>



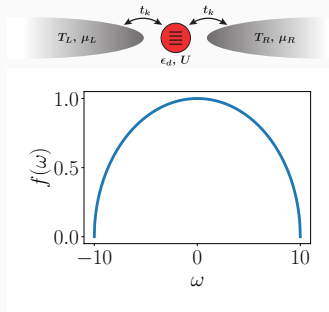
<sup>1</sup>Lerose, Sonner, and Abanin 2022; Frías-Pérez and Bañuls 2022; Banuls et al. 2009

# Model

## Single Impurity Anderson model

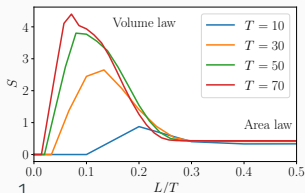
$$H = \sum_{k,\sigma} \left( \epsilon_{k,\sigma} c_{k,\sigma}^\dagger c_{k,\sigma} + V_{k,\sigma} d_\sigma^\dagger c_{k,\sigma} + h.c. \right) + U \left( n_\uparrow - \frac{1}{2} \right) \left( n_\downarrow - \frac{1}{2} \right)$$

- Impurity with on-site interaction coupled to free fermions
- Bath defined by spectral density  $f(\omega) = 2\pi \sum_k |V_k|^2 \delta(\epsilon_k - \omega)$
- Kondo effect: Effective impurity spin forms singlet with the fermions in the bath
- Quantum embedding approaches

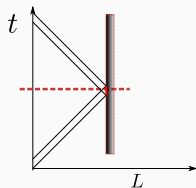
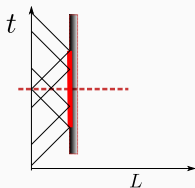
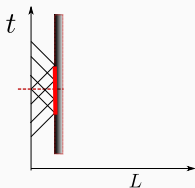




# Entanglement barrier<sup>2</sup>



Quasiparticle picture<sup>1</sup>:



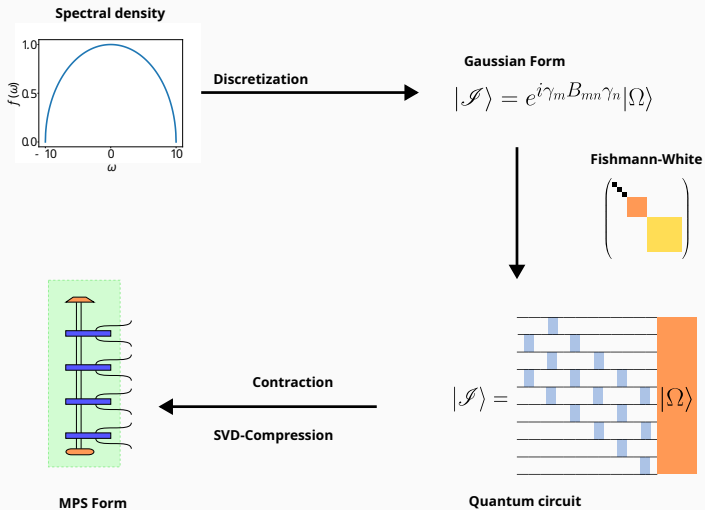
Reflection at the boundary cause high temporal entanglement in finite system IM!

**But:** thermodynamic IM has low temporal entanglement

<sup>1</sup>Calabrese and Cardy 2005.

<sup>2</sup>Lerose, Sonner, and Abanin 2022.

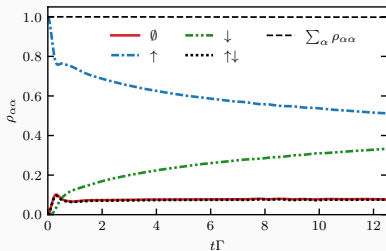
# Algorithm<sup>1</sup>



<sup>1</sup>Thoenniss, Lerose, and Abanin 2022; Thoenniss et al. 2023.

# Spin relaxation<sup>2</sup>

- Initialize impurity in the  $|\uparrow\rangle$  state
- Slow spin relaxation in the Kondo regime
- Competitive with state of the art QMC<sup>1</sup>



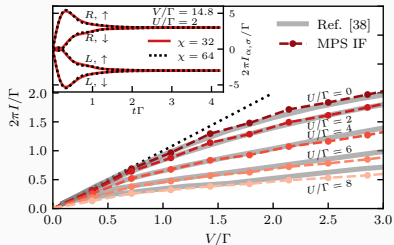
<sup>1</sup>Cohen et al. 2015

<sup>2</sup>Thoenniss et al. 2023.

# Transport<sup>2</sup>



- Left and right reservoir at different chemical potential
- Small bond dimension  $\chi = 32, 64$  (!)
- Steady state reachable
- Comparable with QMC <sup>1</sup>



<sup>1</sup>Bertrand et al. 2019

<sup>2</sup>Thoenniss et al. 2023.

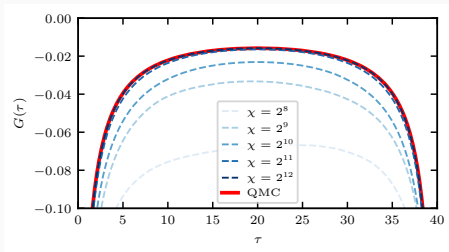


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## Imaginary time Green's function

$$G(\tau) = \text{Tr}(e^{-(\beta-\tau)H} c^\dagger e^{-\tau H} c)$$

- Central step for Dynamic mean field theory (DMFT)
- Requires relatively high bond dimension
- Multi-orbital possible, no sign problem!



<sup>1</sup>Kloss et al. 2023.

# Summary and Outlook

Compute local observables by finding a compressed representation of the large quantum system as a bath (IM)

- Complexity of the IM can be low even in traditionally difficult regimes
- For free fermions there is an efficient algorithm to compute the IM directly
- Natural for impurity problems, DMFT!
- For one dimensional interactive models, we can be efficiently obtain the IM through sequential contraction<sup>1</sup>

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<sup>1</sup>Lerose, Sonner, and Abanin 2022; Frías-Pérez and Bañuls 2022.