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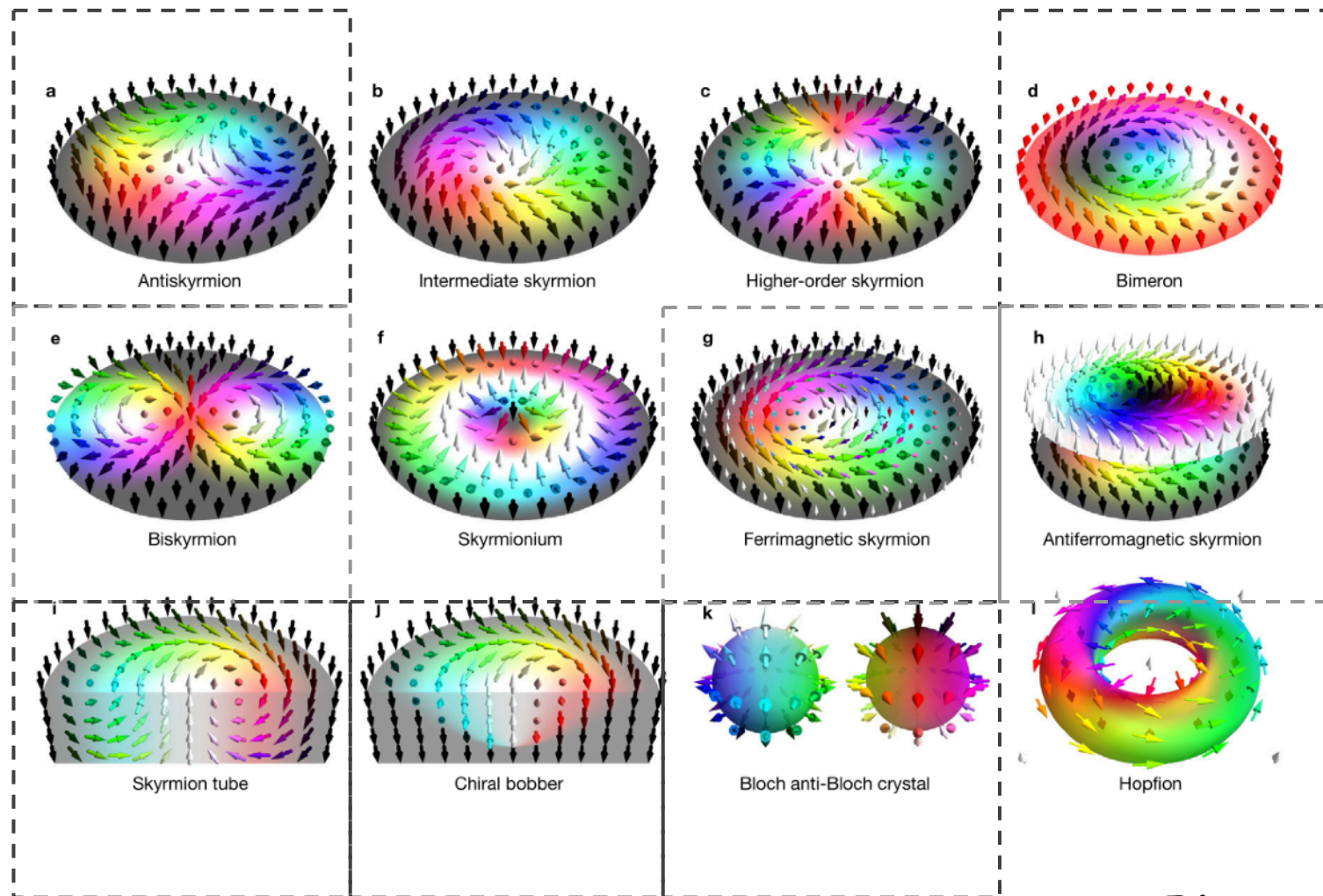
Long-periodic spin textures in vector magnetic fields as seen
by resonant elastic x-ray scattering

YRLG Workshop
Ingelheim
26-07-2023

Topological spin orders in magnetic materials

Particle-like spin objects hosting “topological charge”: a zoo of the topological spin textures

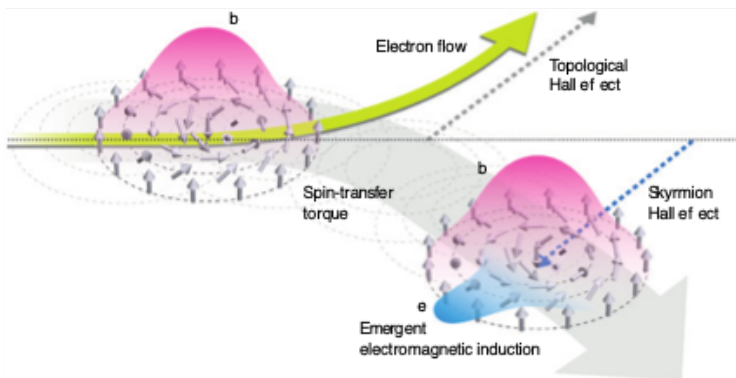
Theoretically proposed topological spin orders: $n = \frac{1}{4\pi} \int \mathbf{M} \cdot \left(\frac{\partial \mathbf{M}}{\partial x} \times \frac{\partial \mathbf{M}}{\partial y} \right) dx dy$



Discovered in bulk materials

Topological spin orders in magnetic materials

- Topological spin textures and emergent fields



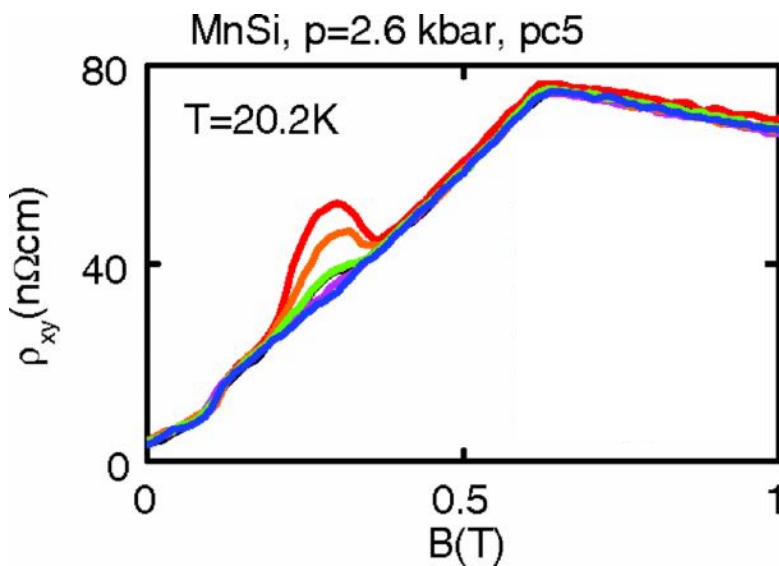
Electrons adiabatically traversing the topological spin texture adapt to the local \mathbf{M} and acquire a quantum-mechanical Berry phase

Emergent magnetic field $\mathbf{B}_i^e = \frac{\hbar}{2} \epsilon_{ijk} \hat{n} \cdot (\partial_j \hat{n} \times \partial_k \hat{n})$

Emergent electric field $\mathbf{E}_i^e = \hbar \dot{\hat{n}} \cdot (\partial_i \hat{n} \times \partial_t \hat{n})$

Topological contribution to the Hall resistivity:

$$\rho_{yx} = \underbrace{R_0 B}_{\text{normal}} + \underbrace{R_S M}_{\text{anomalous}} + \underbrace{\rho_{yx}^T}_{\text{topological}}$$



The effective field in MnSi (skyrmion size of 15 nm)

$$B^{\text{eff}} = -\frac{h}{e} \left(\frac{\sqrt{3}}{2\lambda_S^2} \right) \approx -13.15 \text{ T.}$$

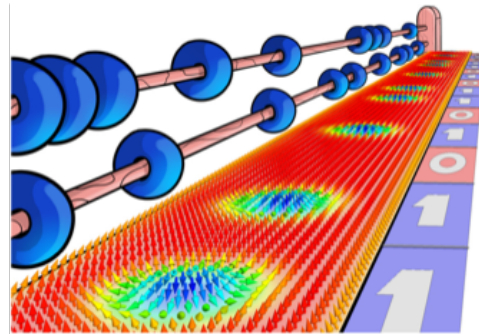
generated by the applied field of 0.3 T!

Topological spin orders in magnetic materials

INTRODUCTION

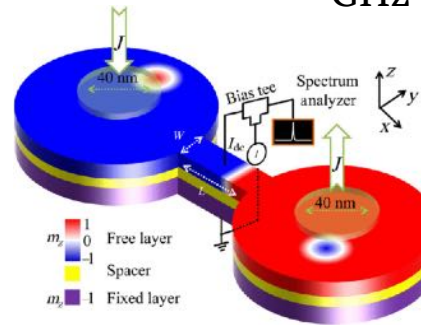
- Potential applications

Racetrack memories

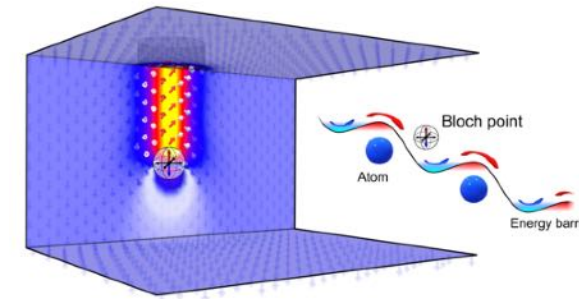


K. Everschor-Sitte, et al. J. Appl. Phys. 124, 240901 (2018)

GHz and THz nano-oscillators

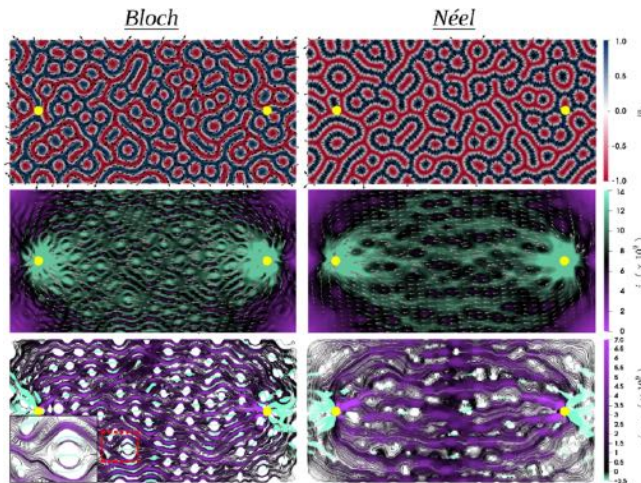


C. Jin, et al. Phys. Rev. App. 9, 044007 (2018)

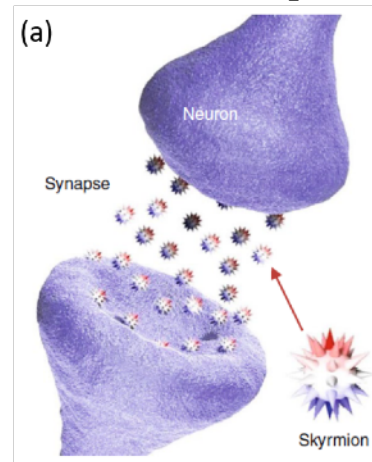


Y. Li, et al. Phys. Rev. Res., 2, 033006 (2020)

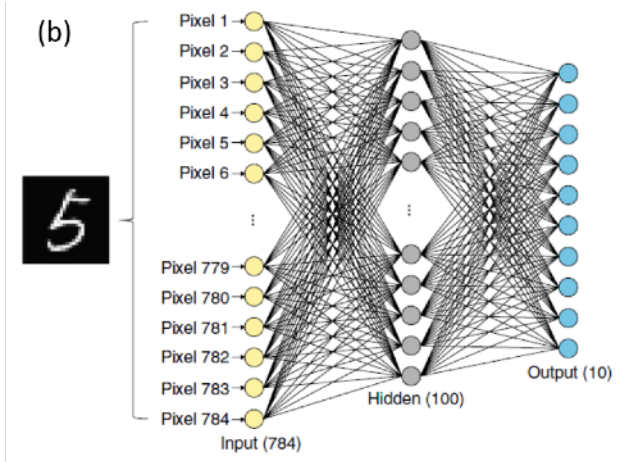
Unconventional computing



K. Everschor-Sitte, et al. AIP Advances 8, 055602 (2018)



K.M. Song, et al. Nature Electronics 3, 3 148-155 (2020)

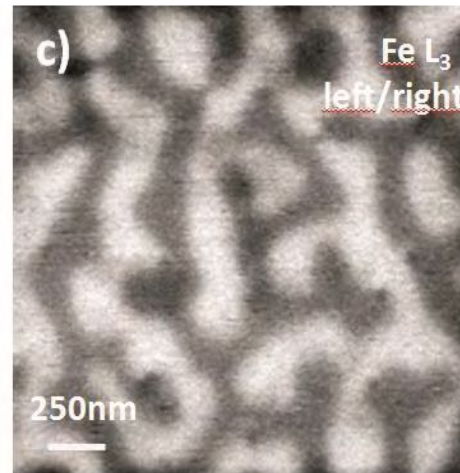
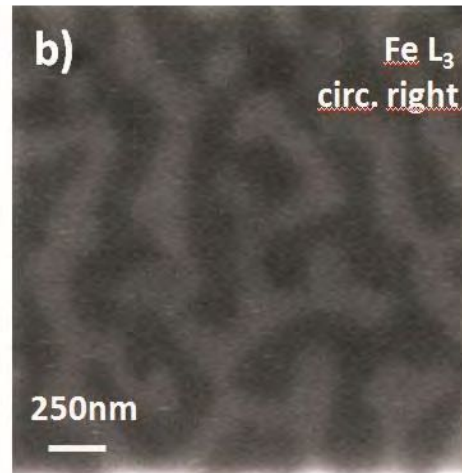
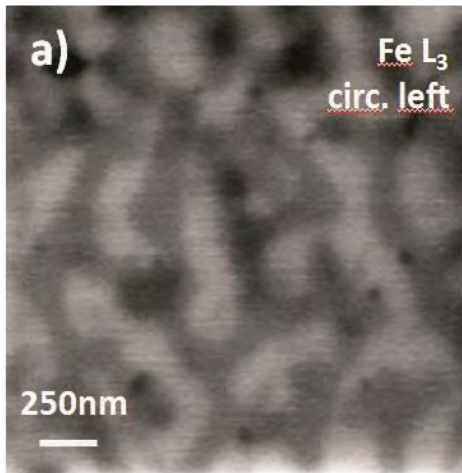
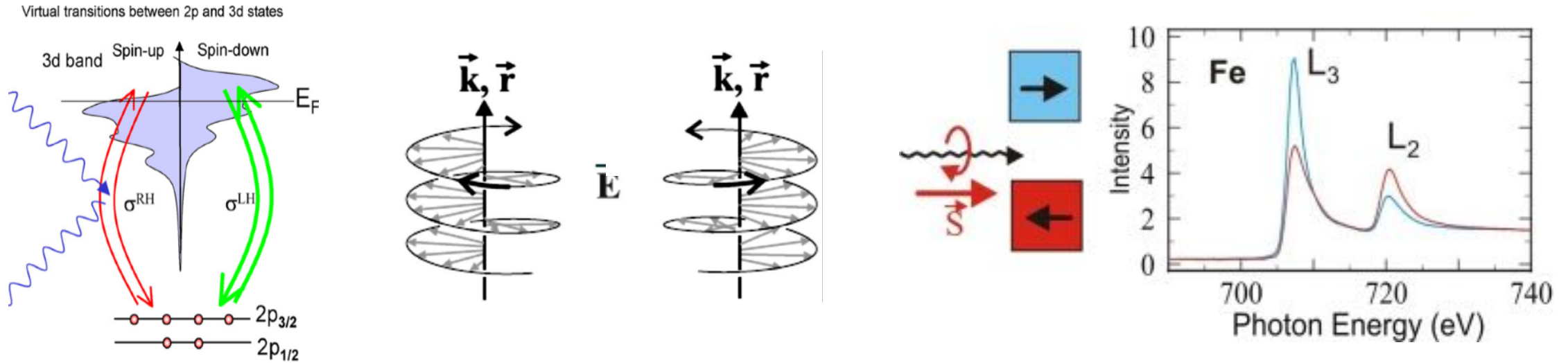


Resonant x-rays for magnetism

INTRODUCTION

Polarised x-rays - a tool to study static and dynamics magnetic orders

contrast mechanism for 'soft' x-rays (~100-2000 eV) - XMCD



$$I(z) = I_0 e^{-\mu_x z}$$

where μ is material-specific coefficient dependent on photon energy

XMCD microscopy, diffraction, scattering, etc.

Resonant x-rays for magnetism (scattering)

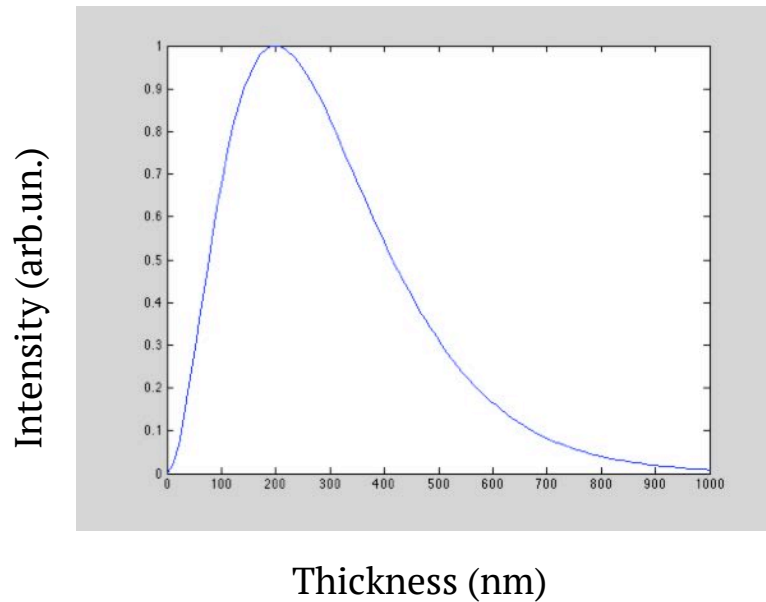
INTRODUCTION

- $3d-2p$ transition is used for resonant scattering and diffraction: $L_{2,3}$ edges of TM correspond to soft x-rays in the energy range of 460 (Ti L_3) – 950 eV (Cu L_2) and wavelength range $\lambda \sim 13-27 \text{ \AA}$
- Resonant x-rays allow both reciprocal (scattering) and real-space (imaging)

Thickness: transmission vs magnetic scattering intensity

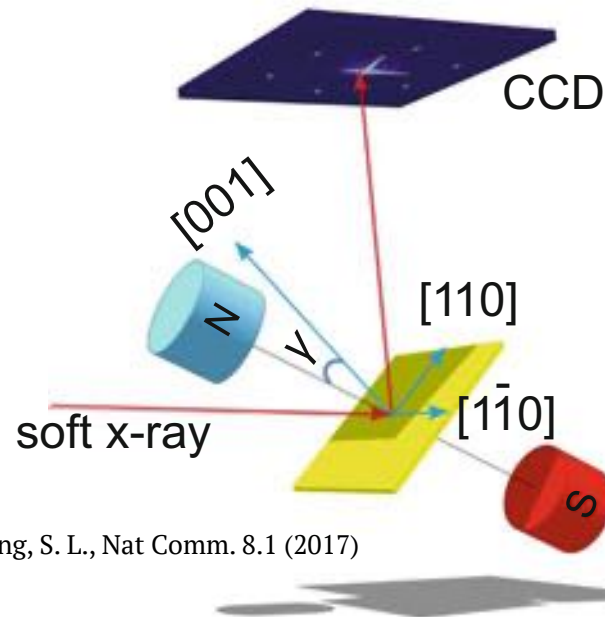
$$I_{mag} \sim M^2 = (\mu A z)^2$$

$$I_{trans} = I_0 \exp(-l_{att} z)$$



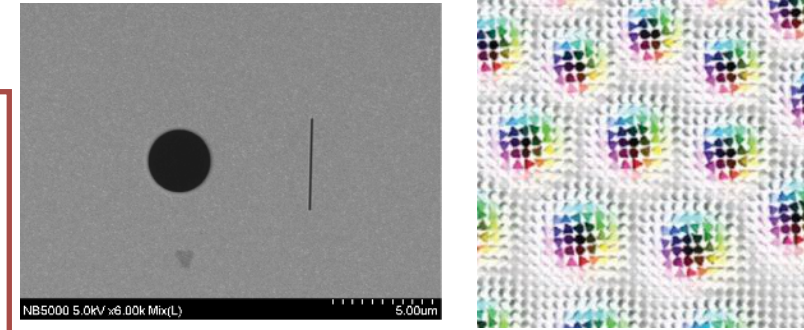
Limited by a few 100s of nm

$$f^n = \underbrace{(\mathbf{e}_n \cdot \mathbf{e}'_n) f_c^n}_{charge} + \underbrace{i(\mathbf{e}_n \times \mathbf{e}'_n) \cdot \mathbf{M}^n f_{m1}^n}_{magnetic} + \underbrace{(\mathbf{e}_n \cdot \mathbf{M}^n)(\mathbf{e}'_n \cdot \mathbf{M}^n) f_{m2}^n}_{magnetic}$$



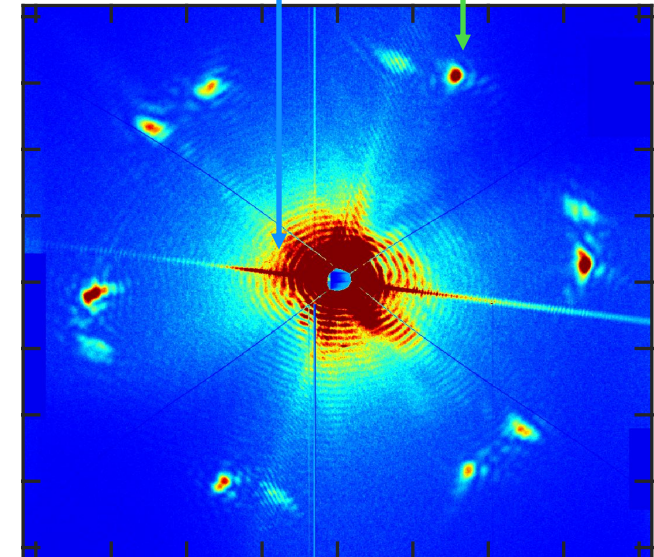
Zhang, S. L., Nat Comm. 8.1 (2017)

Rather weak effect (need a bright x-ray source)



$$\frac{\sigma_m}{\sigma_c} \sim 10^{-2}$$

charge magnetic



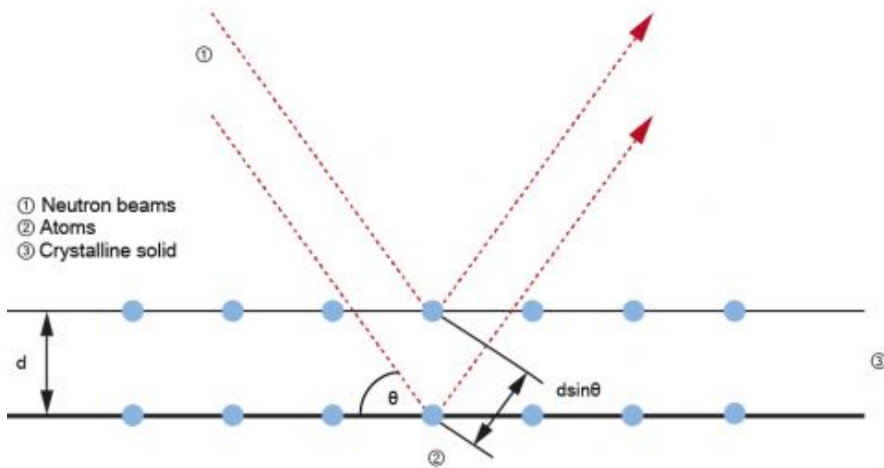
1 μm -thick Cu_2OSeO_3

V.U., unpublished

Small-angle scattering: a tool to study nanometric orders

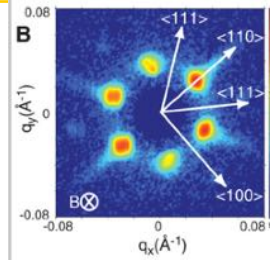
SANS (neutrons)

Bragg diffraction



typical SANS setting

$\lambda \sim 5 \text{ \AA}$
 $d \sim 800 \text{ \AA}$
 $\theta \sim 0.2^\circ$ - small!

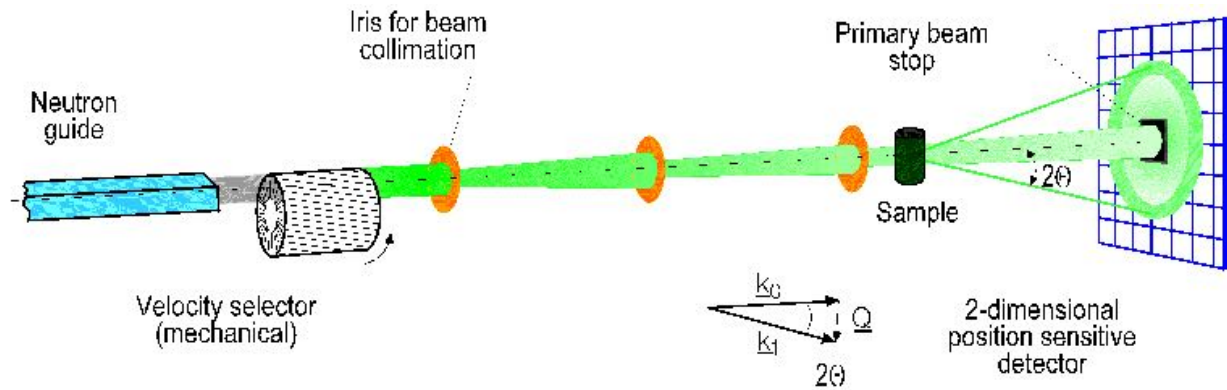


pixel size $\sim 1 \times 1 \text{ cm}^2$

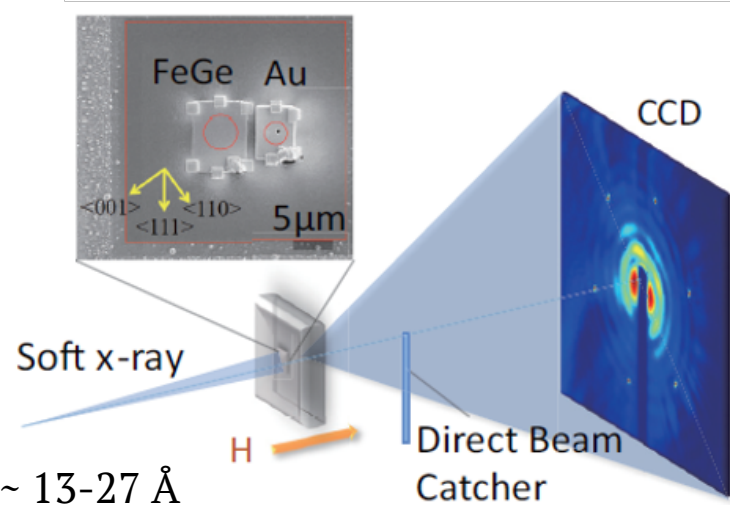
SAXS (resonant x-rays)

Also known as:

- Resonant Elastic X-ray Scattering (REXS)
- X-ray Magnetic Resonant Scattering (XRMS)
- Resonant X-ray Magnetic Scattering (RXMS)
- Soft X-ray Magnetic Resonant Scattering (SXRMS)
- Resonant Magnetic Small-angle X-ray Scattering (RMSAXS)
- Small-angle Resonant Elastic X-ray Scattering (SAREXS)
- Transmission REXS (T-REXS)
- Anomalous Small-angle X-ray Scattering (ASAXS)
- Resonant Soft X-ray Scattering (RSoXS), etc.



Large instrument!



$\lambda \sim 13-27 \text{ \AA}$

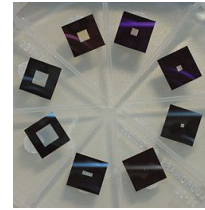
pixel size $\sim 15 \times 15 \mu\text{m}^2$



Sample preparation for soft x-ray experiment

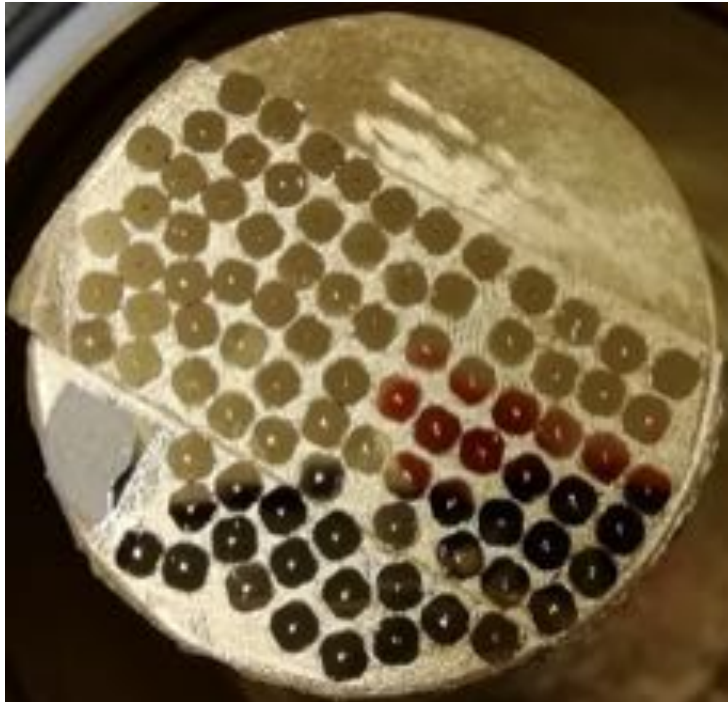
- Sample size and shape are adjusted according to the attenuation length of soft x-ray beam (~50 - 1000 nm thick for TM) and desire to use the coherence (aperture size of 1-20 μm)
- Method: thin plates prepared by Focused Ion Beam (FIB) milling

Commercial membranes (Silson Ltd., UK)
50 - 500 nm



Silicon nitride is almost transparent for soft x-rays

1 μm Au-coated
 Si_3N_4 membranes



Zeiss NVision
at EMF PSI

Sample preparation for soft x-ray experiment

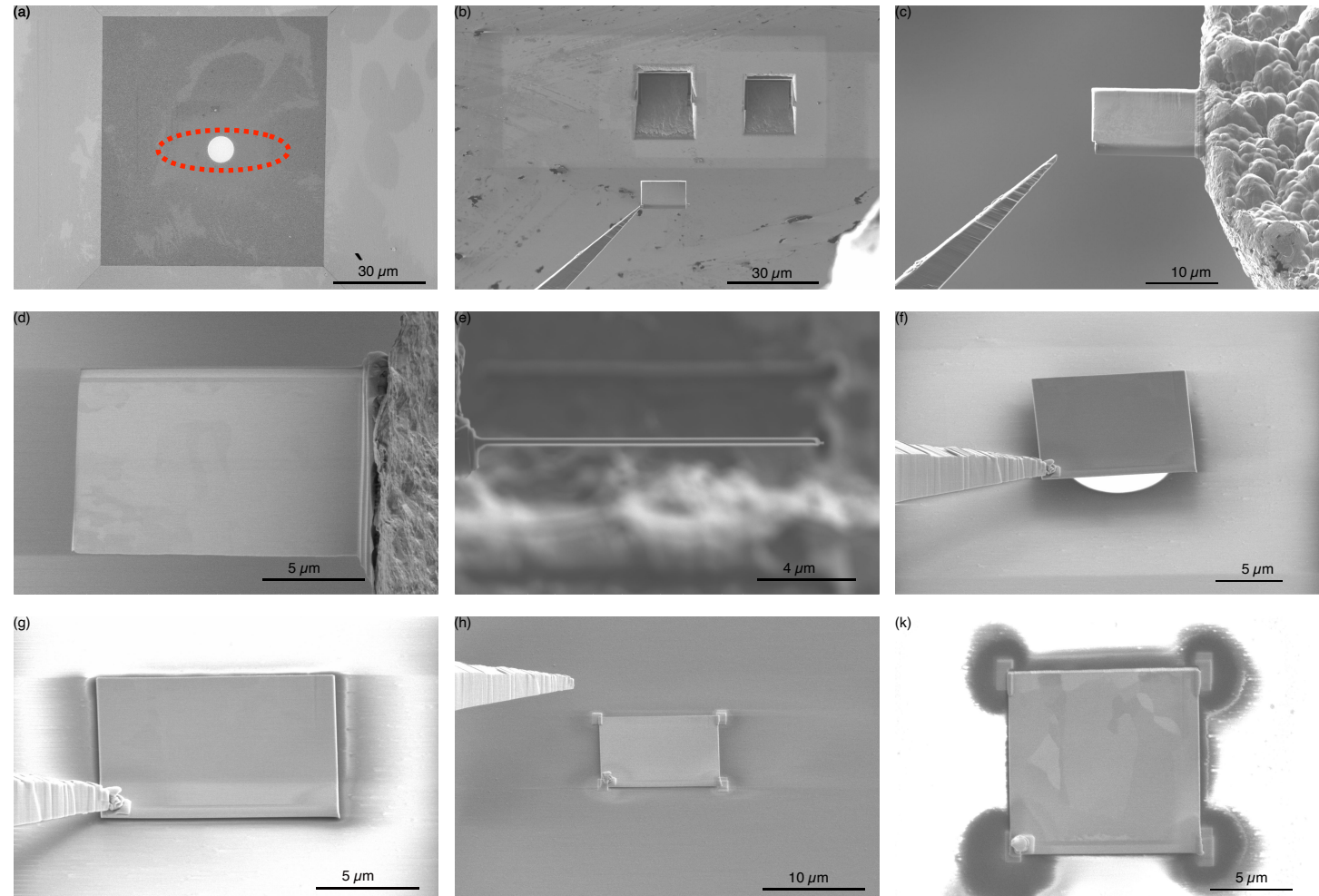
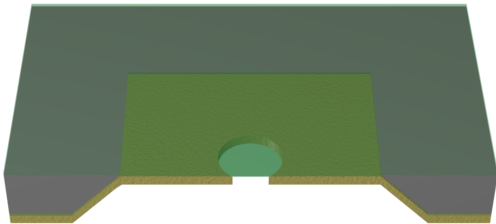
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- Method: thin plates prepared by Focused Ion Beam (FIB) milling

Membrane preparation

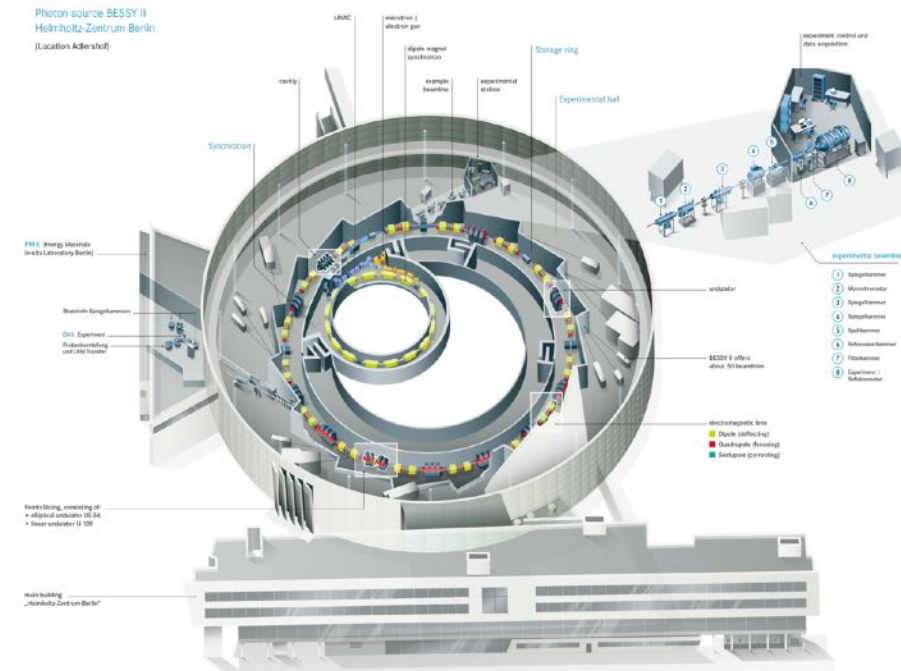
Si_3N_4 membrane: thickness of 200 nm (transparent for soft X-rays)



Coating membrane with a few micron of Au to block the beam



- **Location:** Adlershof, Berlin, Germany
- **Operator:** Helmholtz-Zentrum Berlin für Materialien und Energie (Financing: 90% federal government, 10% State of Berlin)
- **Commissioned:** September 1998
- **Upgrades:** October 2012: Top-up mode and fast-orbit feedback
- **Circumference:** 240 metres
- **Bending magnets:** 32
- **Beam tubes:** approx. 45
- **Electron energy:** 1.7 gigaelectronvolts (GeV)
- **Nominal beam current:** 300 milliamperes (mA)
- **Energy of the synchrotron radiation:** 1 to 150 kiloelectronvolts (keV)
- **Duration of the light pulse:** 20 picoseconds (after the upgrade to BESSY VSR: 1.5 and 15 ps)
- **Measurement time per year:** 40 weeks p.a., approx. 6,600 hours or 800 shifts p.a. (2017)
- **User visits from guest researchers:** 3,000 p.a.
- **Strategic partners:** Max Planck Society, PTB, BAM, Berliner universities



PM-2 beamline:

- **Source:** Dipole
- **Energy range:** 20-1600 eV
- **Flux:** $\sim 5 \cdot 10^9$ photons/sec
- **Polarization:** 77% circular, linear horizontal
- **Focus:** from $60 \times 50 \mu\text{m}$ to $800 \times 800 \mu\text{m}$
- **Beam availability:** 24 h/d, 6 d/w, 40 w/y
- **User beam:** $\sim 70\%$

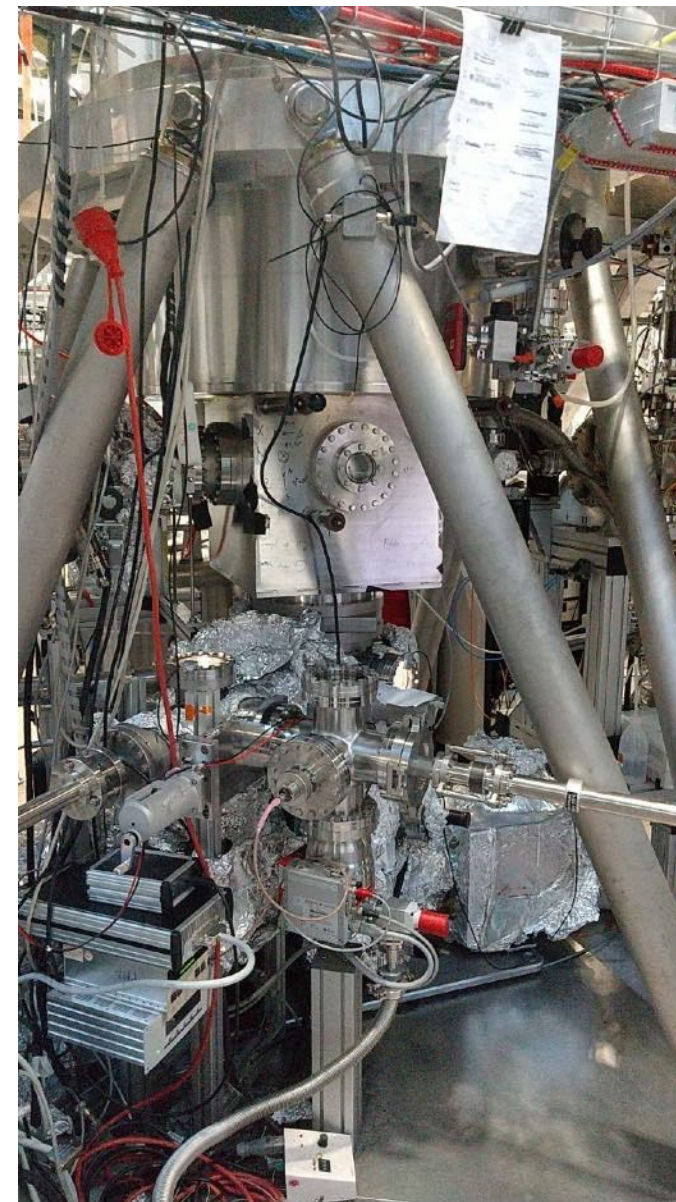
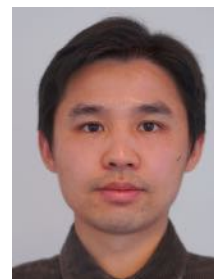
VEKMAG instrument:

- **Pressure:** $\sim 10^{-9}$ mbar
- **Temperature:** 350 mK - 500 K
- **Magnetic field:** 9 T (X), 2 T (Y), 1 T (Z), up to 1 T in any 3D direction
- **Sample manipulator:** x, y, z, θ , azimuthal (ext. holder)
- **Detectors:** drain current: TEY, photodiodes: FY, transmission, reflection ($z, 2\theta$), 2k-CCD, 4k-CCD (scattering in transmission).
- **Techniques:** XAS, XMCD, XLMD, XFMR, XRMS, SAXS
- **Unique capabilities:** downstream deposition chamber for in-situ experiments and vacuum transfer; in-situ cleaving; electric transport capability; XFMR; SAXS; laser (to be commissioned in 2023-2024)

Instrument scientist: Dr. Florin Radu



Dr. Victor Ukleev Dr. Chen Luo



EXAMPLE I

High-resolution SAXS for capturing weak effects in cubic helimagnets

Bak-Jensen model of B20s:

$$H_{\text{EX}} = -\frac{1}{2} \sum J_{\mathbf{R}-\mathbf{R}'} \mathbf{S}_{\mathbf{R}} \cdot \mathbf{S}_{\mathbf{R}'};$$

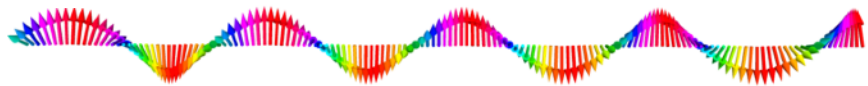
$$H_{\text{DM}} = \frac{1}{2} \sum D_{\mathbf{R}-\mathbf{R}'} (\nabla - \nabla') [\mathbf{S}_{\mathbf{R}} \times \mathbf{S}_{\mathbf{R}'}];$$

$$H_{\text{AEI}} = \frac{1}{2} \sum F_{\mathbf{R}-\mathbf{R}'} \{ (\nabla_x S_{\mathbf{R}}^x) (\nabla'_x S_{\mathbf{R}'}^x) + (\nabla_y S_{\mathbf{R}}^y) (\nabla'_y S_{\mathbf{R}'}^y) + (\nabla_z S_{\mathbf{R}}^z) (\nabla'_z S_{\mathbf{R}'}^z) \};$$

$$H_Z = -\mathbf{H} \sum \mathbf{S}_{\mathbf{R}},$$

$$H_{\text{CA}} = K \sum \{ (S_{\mathbf{R}}^x)^4 + (S_{\mathbf{R}}^y)^4 + (S_{\mathbf{R}}^z)^4 \}.$$

Helical state



Conical state



spiral wvector $k = 2\pi/l$

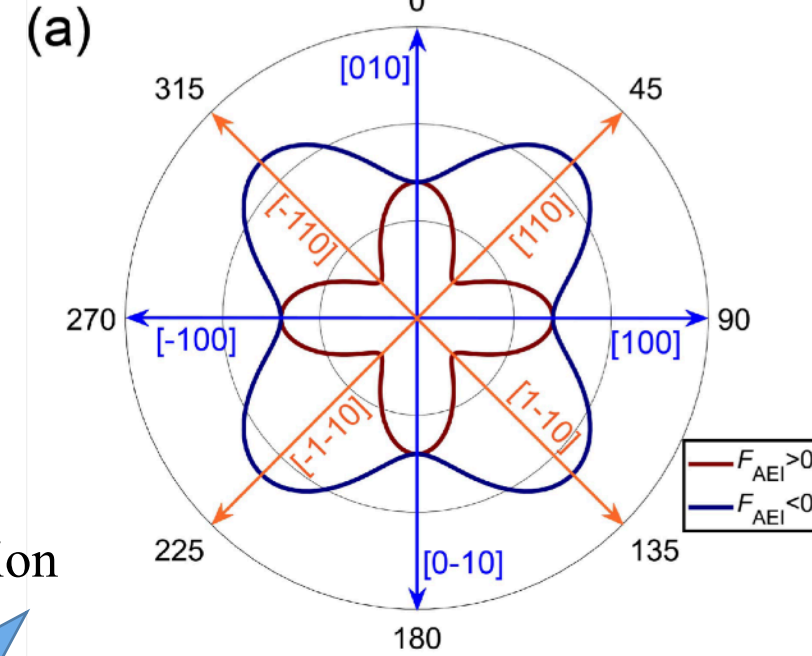
$$k = S \frac{D}{A} \left(1 - \frac{SF}{2A} L(\hat{k}) \right) \approx S \frac{D}{A},$$

standard "D/J" "weak" AEI contribution (a few % to D/J)

L - Cubic invariant:

$$\sin^2 \vartheta (\sin^2 \varphi \cos^2 \varphi + \cos^2 \vartheta);$$

2/3 for [111], 0 for [100]



Anisotropic k -vector in the (001) plane

The k -vector anisotropy is the only way to isolate the AEI

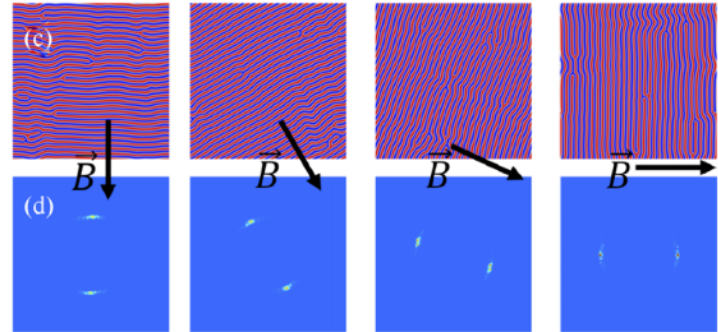
High-resolution SAXS for capturing weak effects in cubic helimagnets

- Recently, we proposed a SAXS-based method to unambiguously quantify the AEI constant via tuning the spiral orientation by the vector field
- Anisotropic exchange interaction is called to explain the new tilted conical and low-T skyrmion phases in cubic chiral insulator Cu_2OSeO_3

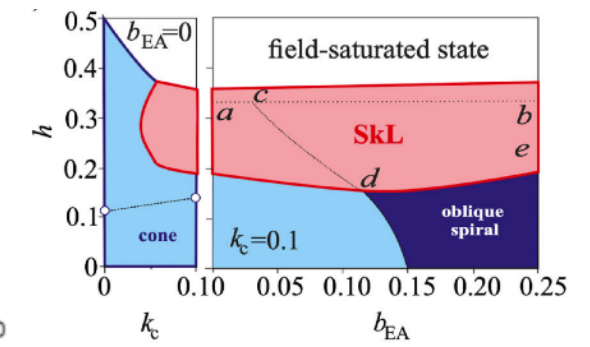
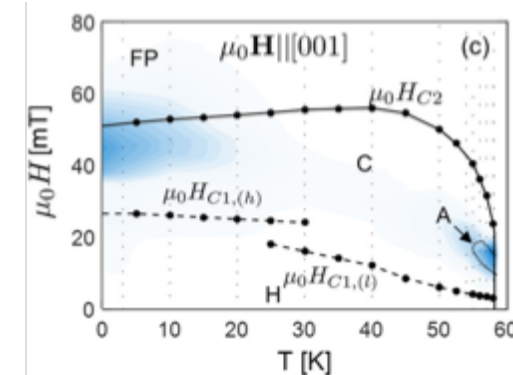
Recent data on **FeGe**

$$\mathbf{Q} = \frac{SD\hat{c}}{A_{ex}} \left(1 - \frac{SFI(\hat{c})}{2A_{ex}} \right),$$

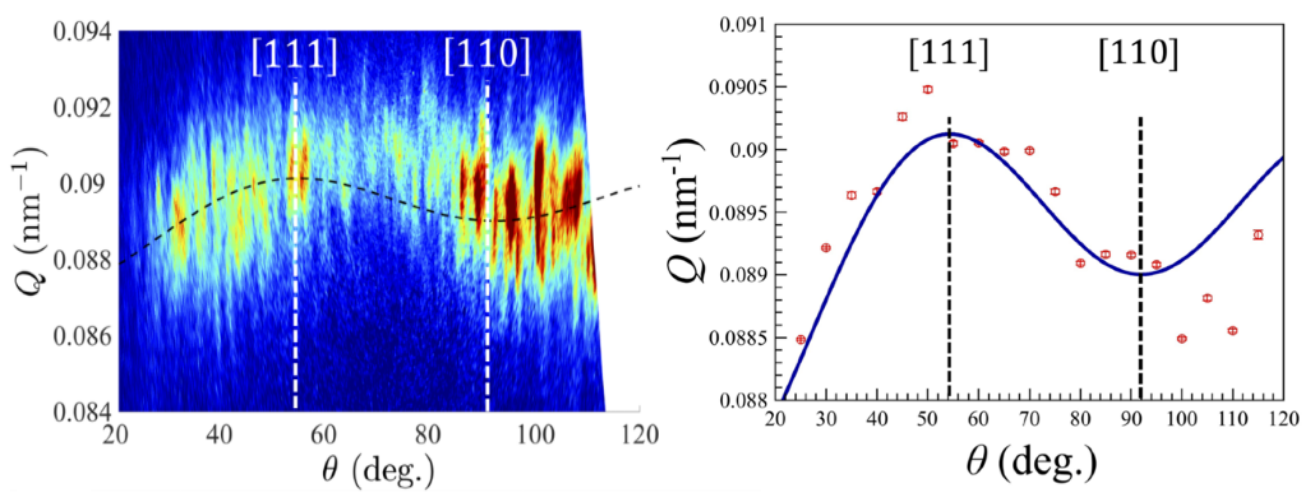
$$I(\hat{c}) = I(\theta) = \frac{\sin^2 \theta (1 + 3 \cos^2 \theta)}{2}.$$



What about **Cu_2OSeO_3** ?



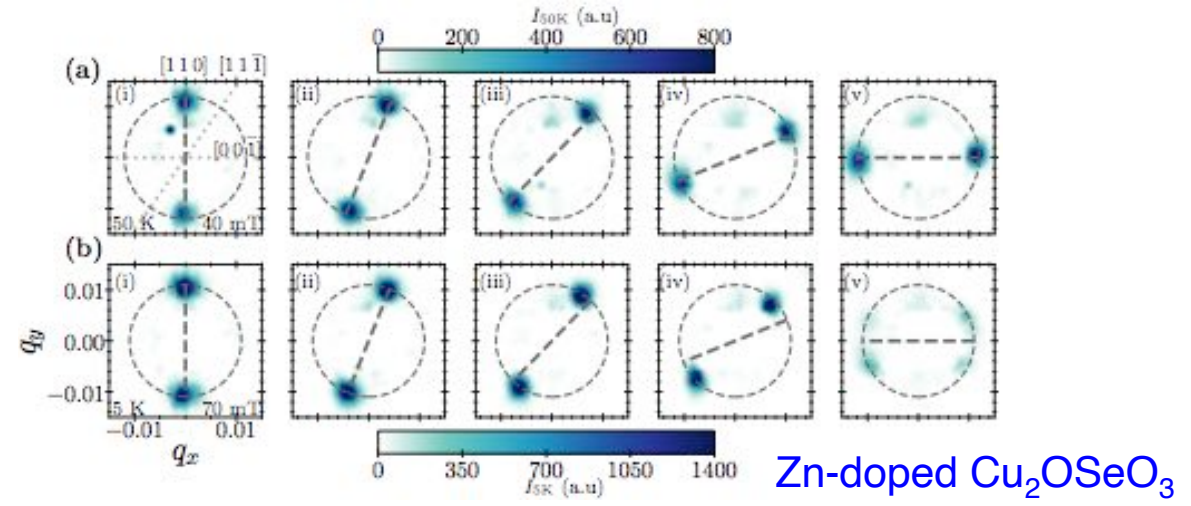
F. Qian, et al., Science Advances 4.9 (2018): eeat7323.



PHYSICAL REVIEW RESEARCH 3, 013094 (2021)

Signature of anisotropic exchange interaction revealed by vector-field control of the helical order in a FeGe thin plate

Victor Ukleev^{1,*} Oleg Utesov^{2,3,4} Le Yu,^{1,5,6} Chen Luo⁷ Kai Chen,⁷ Florin Radu⁷ Yuichi Yamasaki,^{8,9} Naoya Kanazawa,¹⁰ Yoshinori Tokura,^{10,11,12} Taka-hisa Arima,^{11,13} and Jonathan S. White¹⁴

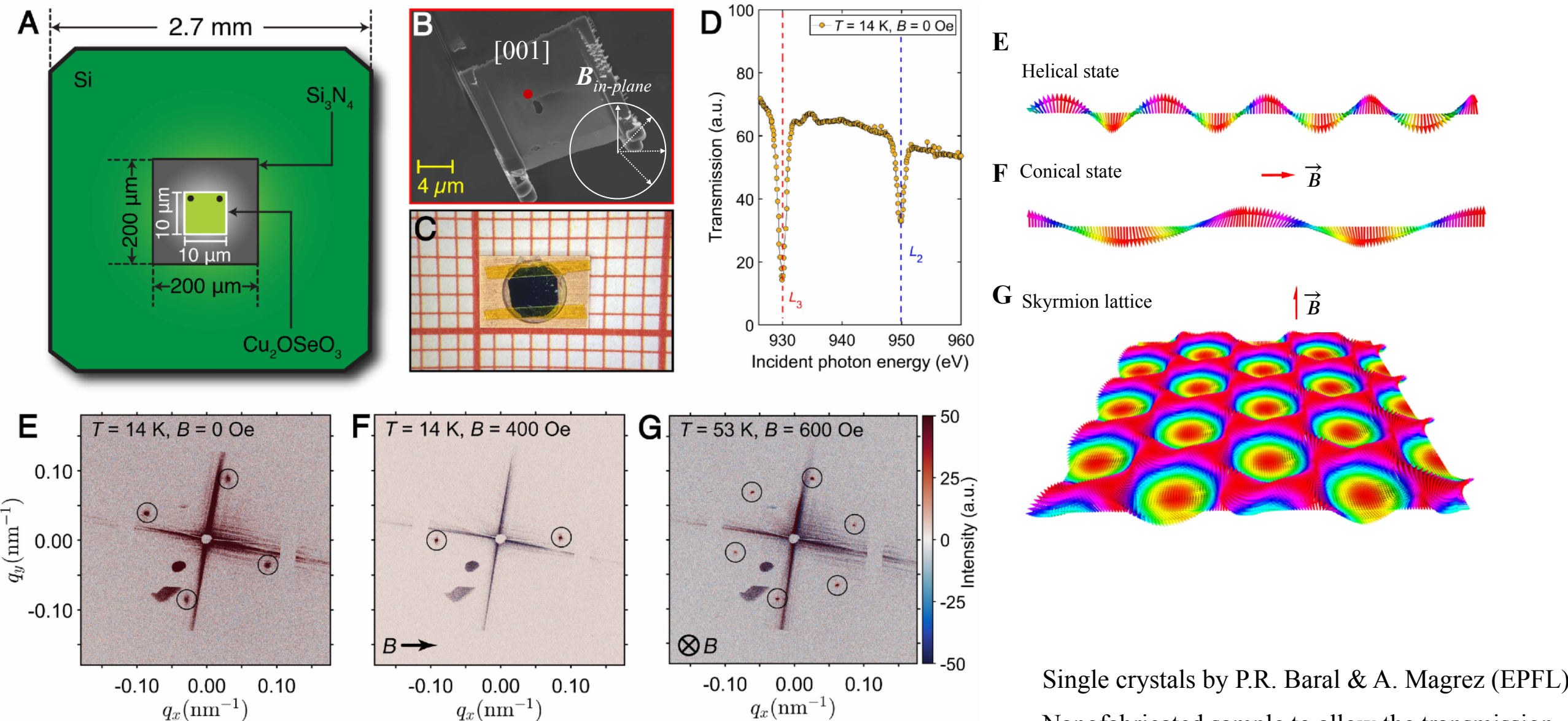


S. H. Moody, et al., Phys. Rev. Research 3, 043149 (2021)

Zn-doped Cu_2OSeO_3

High-resolution SAXS for capturing weak effects in cubic helimagnets

Cu_2OSeO_3 is a $P2_13$ cubic chiral magnet with the spiral period of 60 nm



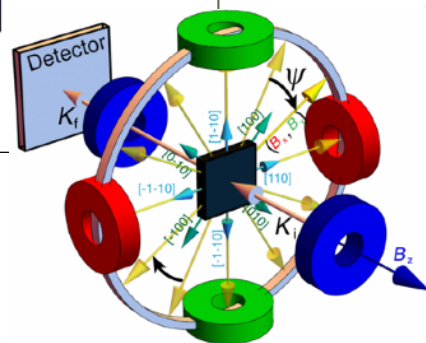
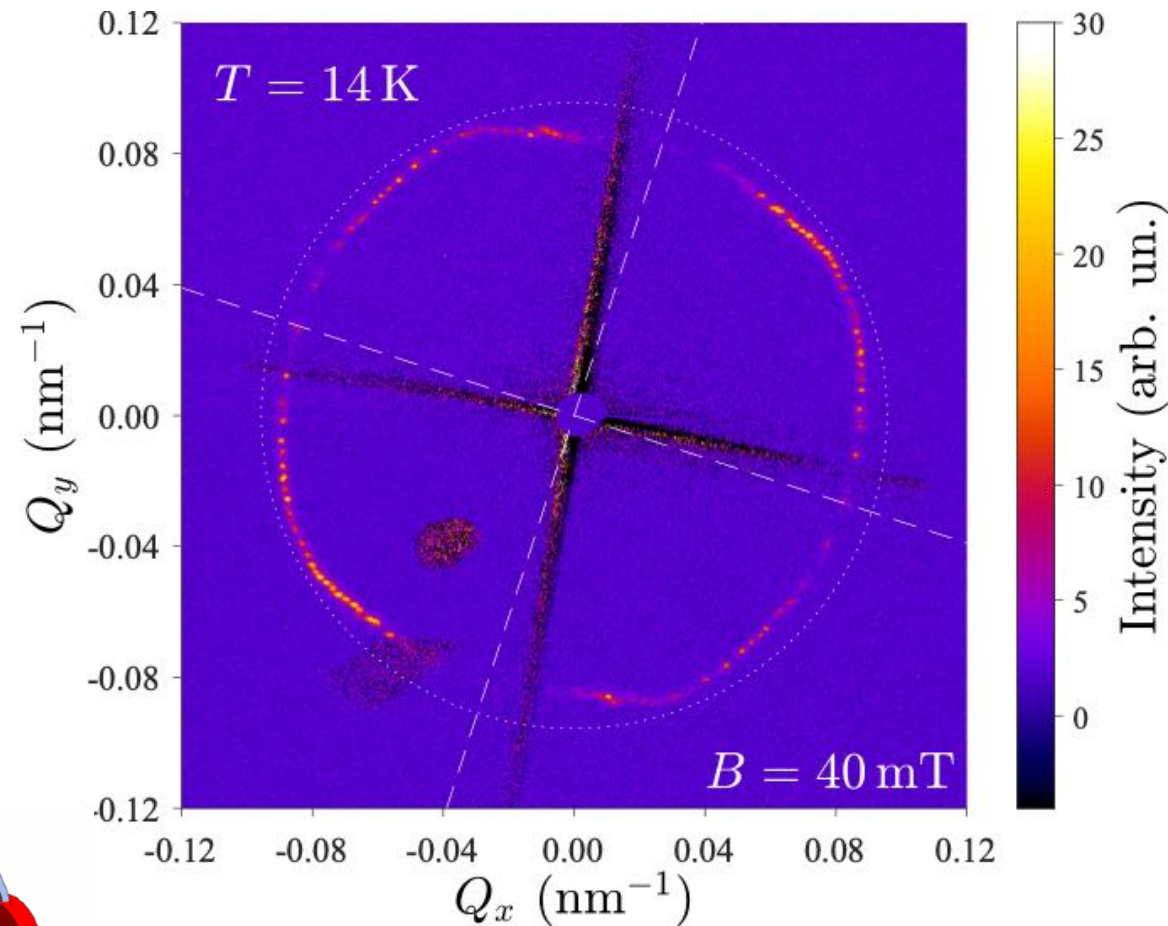
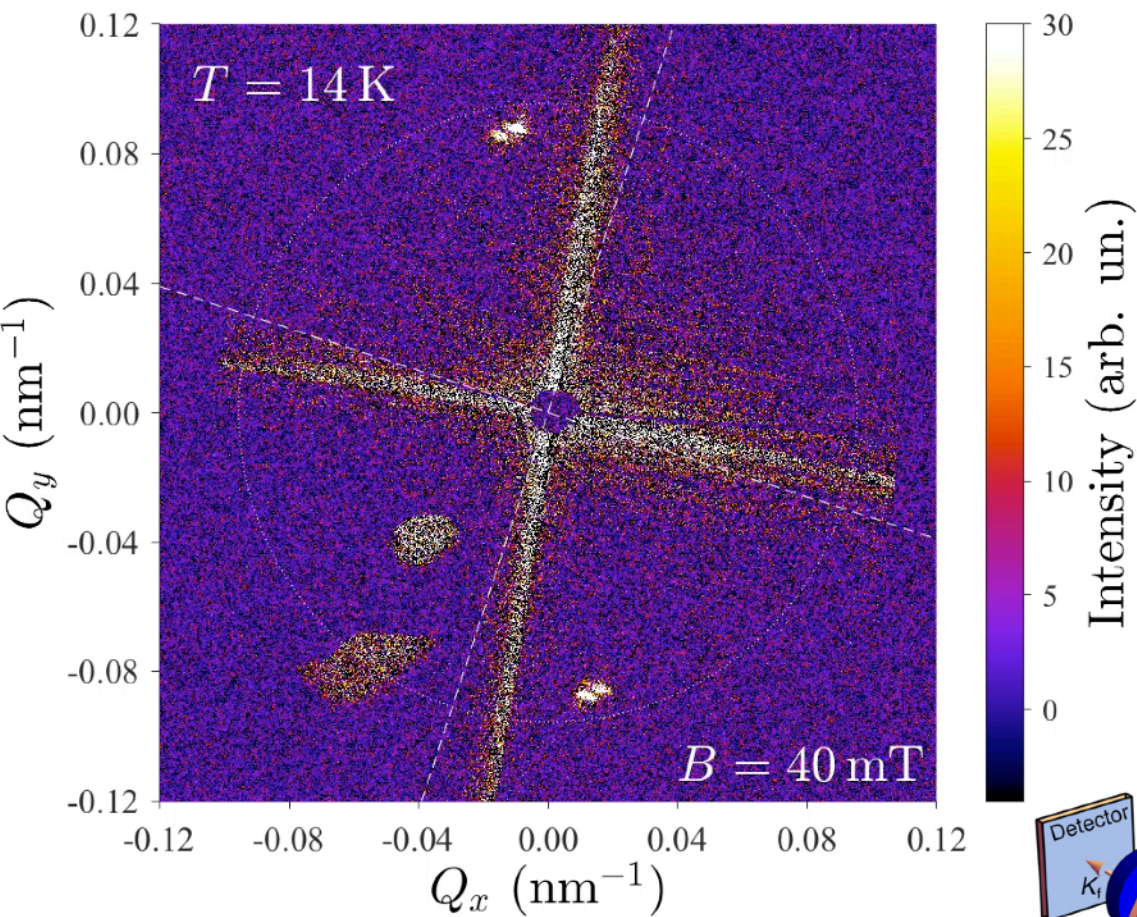
Single crystals by P.R. Baral & A. Magrez (EPFL)

Nanofabricated sample to allow the transmission

High-resolution SAXS for capturing weak effects in cubic helimagnets

Individual images of the conical states with 3 deg. step

Summed image at 14 K

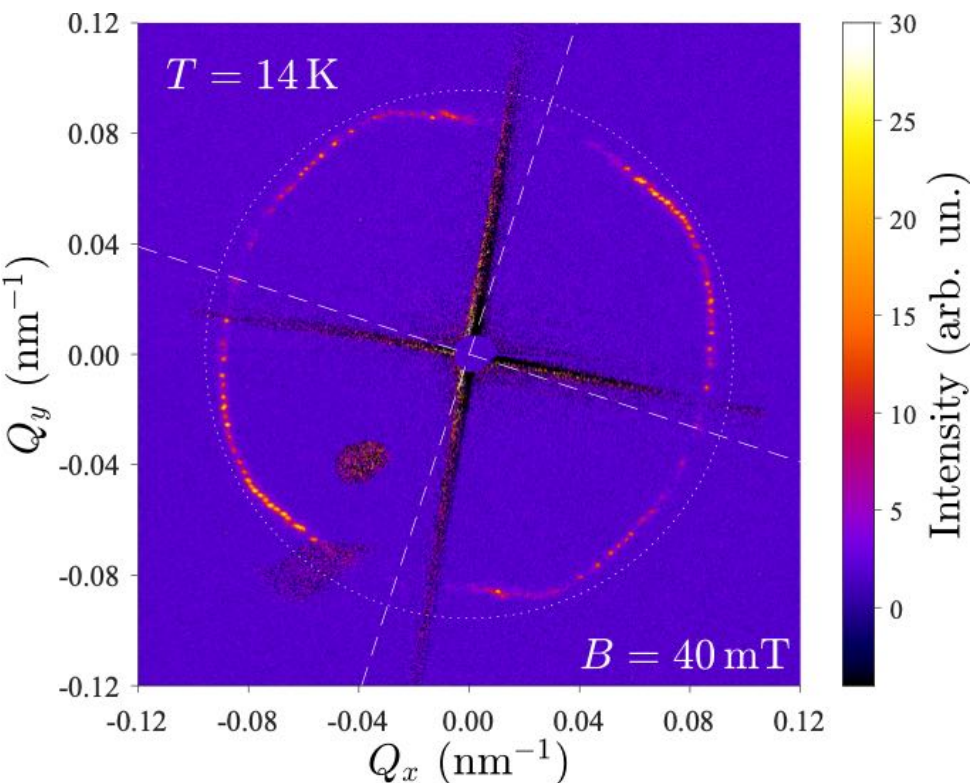


Peaks are not lying at the same $|Q|$!

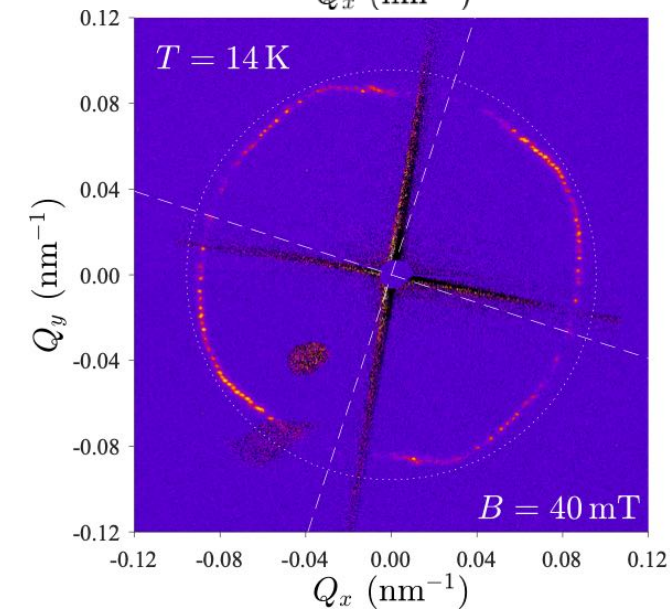
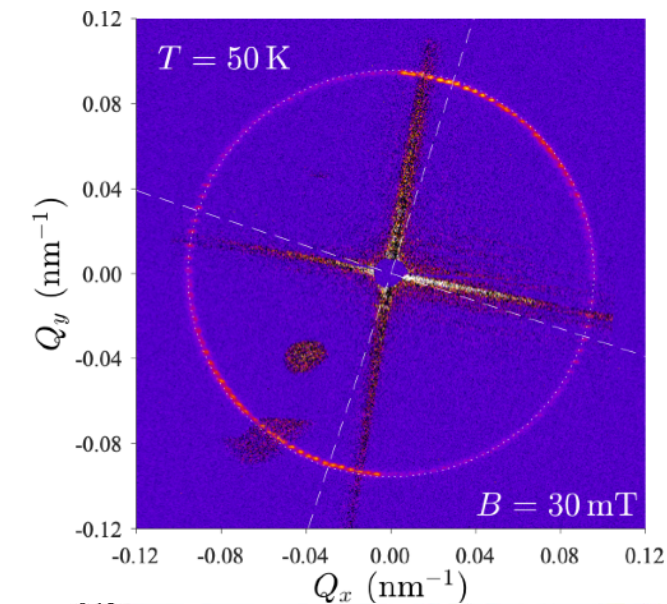
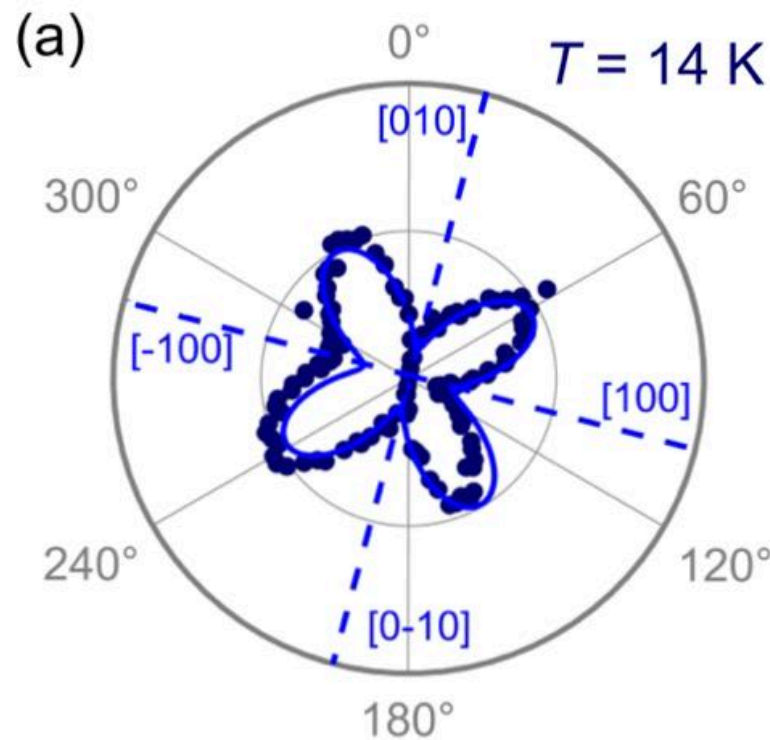
High-resolution SAXS for capturing weak effects in cubic helimagnets

50 K vs. 14 K

Summed pattern for all vector field angles



Fitted peak position



Strongly anisotropic periodicity of the conical texture!

Q is minimum for $\langle 100 \rangle$ and maximum for $\langle 110 \rangle \rightarrow F > 0$

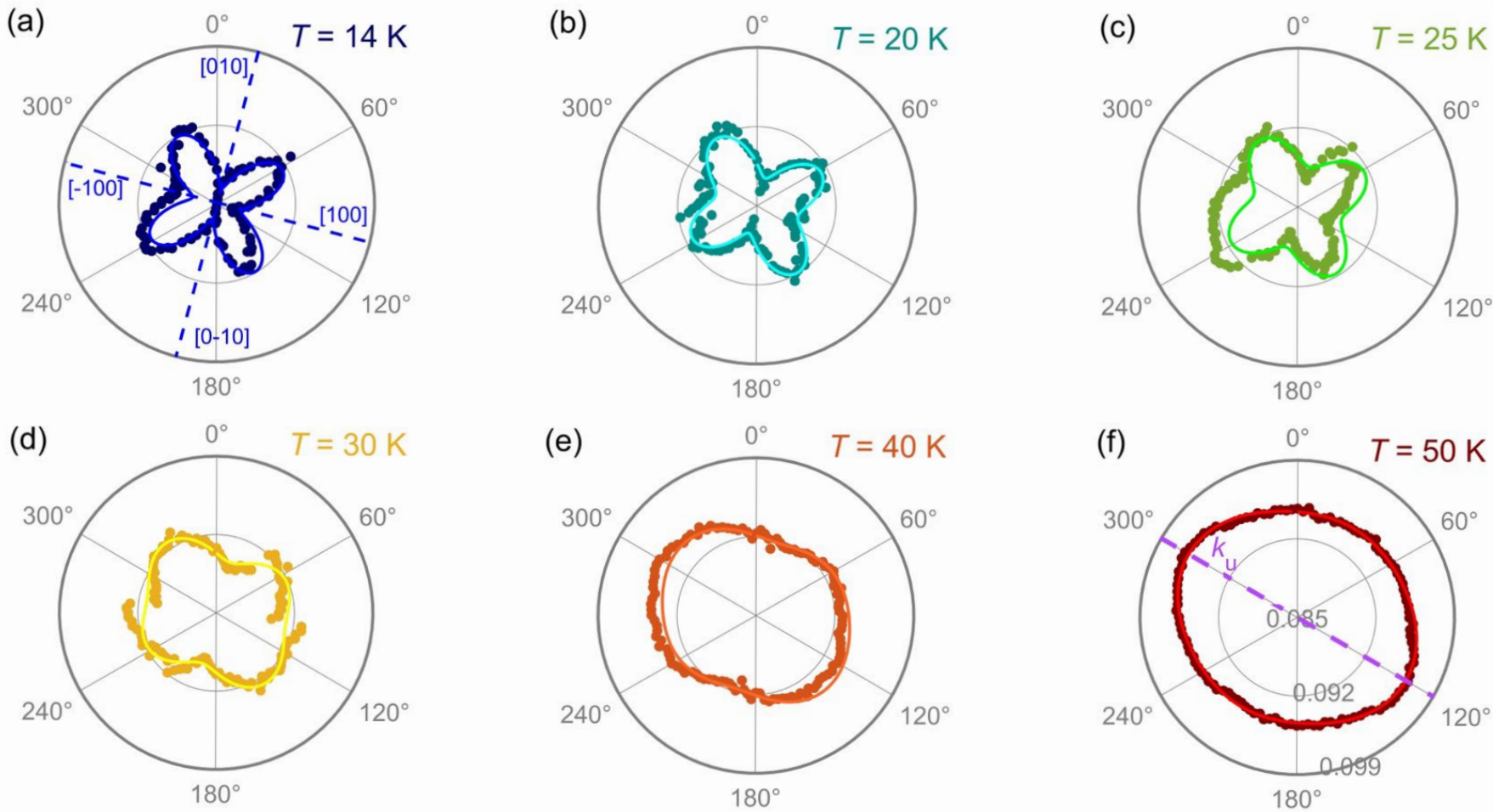
$$Q = Q_0 \left\{ 1 - \frac{F_{AEI} \sin^2 2\psi}{4J} \right\} - \frac{JZ^2 \cot^2 \alpha}{2D^3} \sin^2 2(\psi - \phi)$$



Theory by O. Utesov (IBS, Korea Rep.)

EXAMPLE I

High-resolution SAXS for capturing weak effects in cubic helimagnets

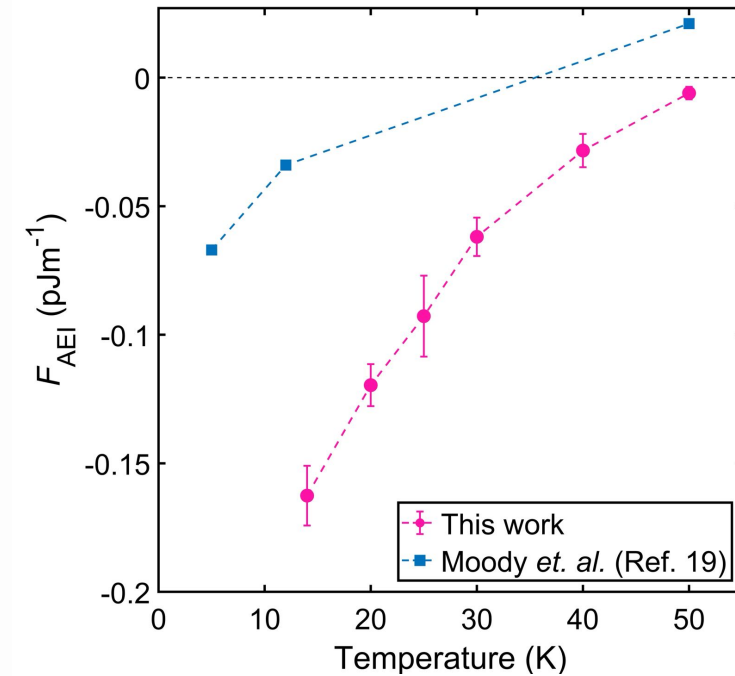


$$Q = Q_0 \left\{ 1 - \frac{F_{AEI} \sin^2 2\psi}{4J} \right\} \left(\frac{JZ^2 \cot^2 \alpha}{2D^3} \sin^2 2(\psi - \phi) \right)$$

Anisotropy gradually vanishes on warming to 50 K ($T_c = 55$ K), but a weak uniaxial distortion is present due to the tensile strain

F vs T

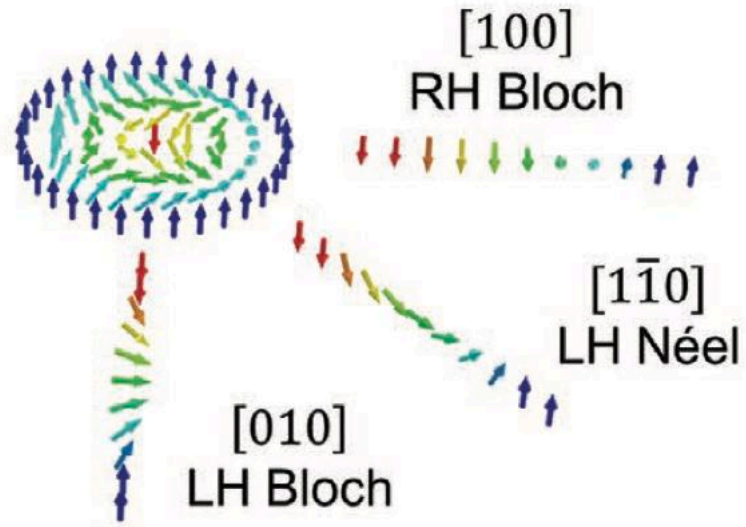
Dramatic increase compared to the Zn-doped COSO



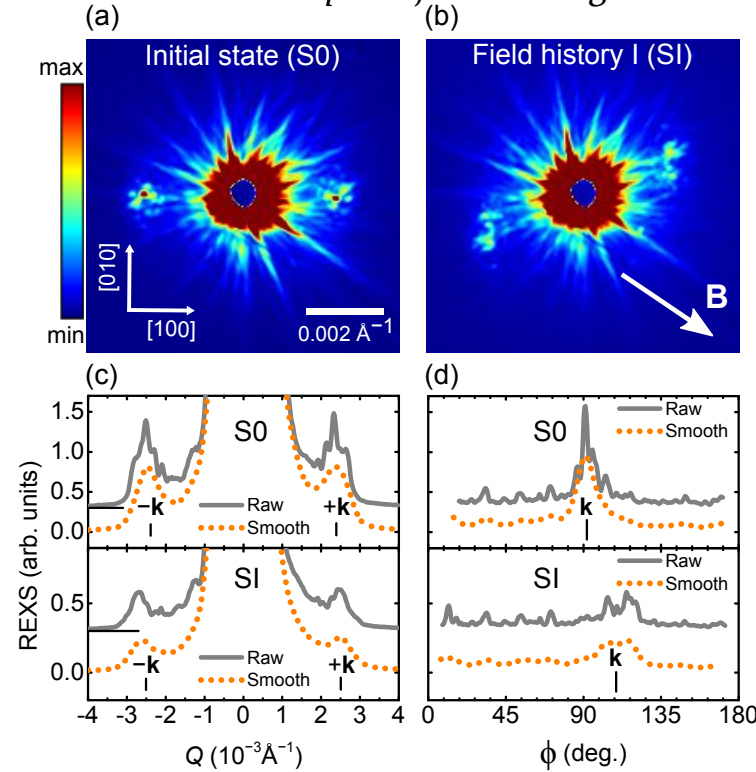
Reciprocal space resolution and acquisition time are superior compared to the neutron scattering

Tetragonal (D_{2d}) anti-skyrmion Heusler alloy $Mn_{1.4}PtSn$

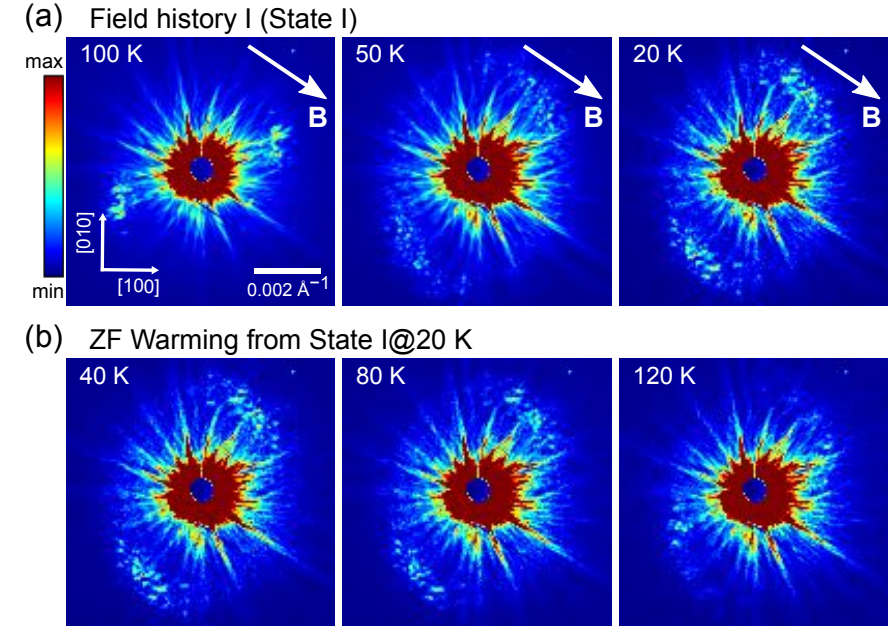
tilting of the spiral plane by the in-plane field



Ex-situ in-plane field training:



Actually depends on temperature



T. Ma, et al., Adv. Mater. 2020, 32, 2002043

- Flexibility in the field orientation compared to Lorentz TEM: ex-situ in-plane field on D_{2d} $Mn_{1.4}PtSn$ at BOREAS (ALBA)
- The spiral moves towards orthogonal-to-the-field axis - typical to cycloid/stripe domain but it is unexpected for this symmetry group
- At low T s gets more ‘cycloidal’

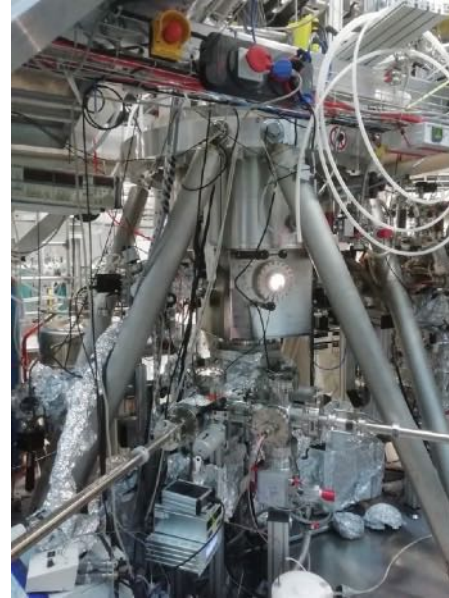
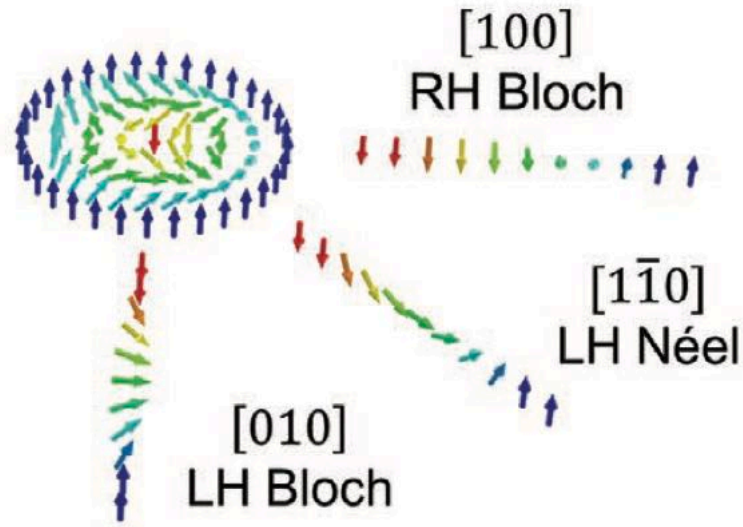
Dipolar interactions!

SAXS on single-Q and multi-Q textures in tetragonal Heusler alloy

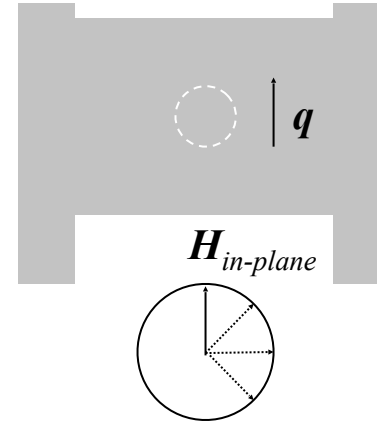
In collaboration with TU Dresden and MPI CPfS

Tetragonal (D_{2d}) anti-skyrmion Heusler alloy $Mn_{1.4}PtSn$

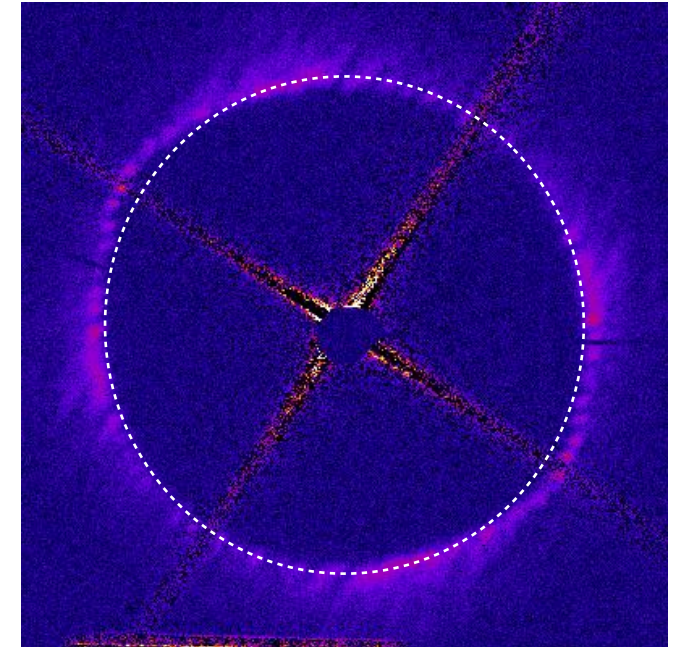
Full rotation of the spiral plane by the in-plane field!



VEKMAG @ BESSY-II



Summed over all azimuthal angles



T. Ma, et al., Adv. Mater. 2020, 32, 2002043

- Full control over the in-plane spiral propagation plane via vector magnetic field at VEKMAG (BESSY-II)
- At 300 K the spiral can be rotated freely within the tetragonal plane
- The Q -vector anisotropy depends on the sample shape



Dipolar interactions!

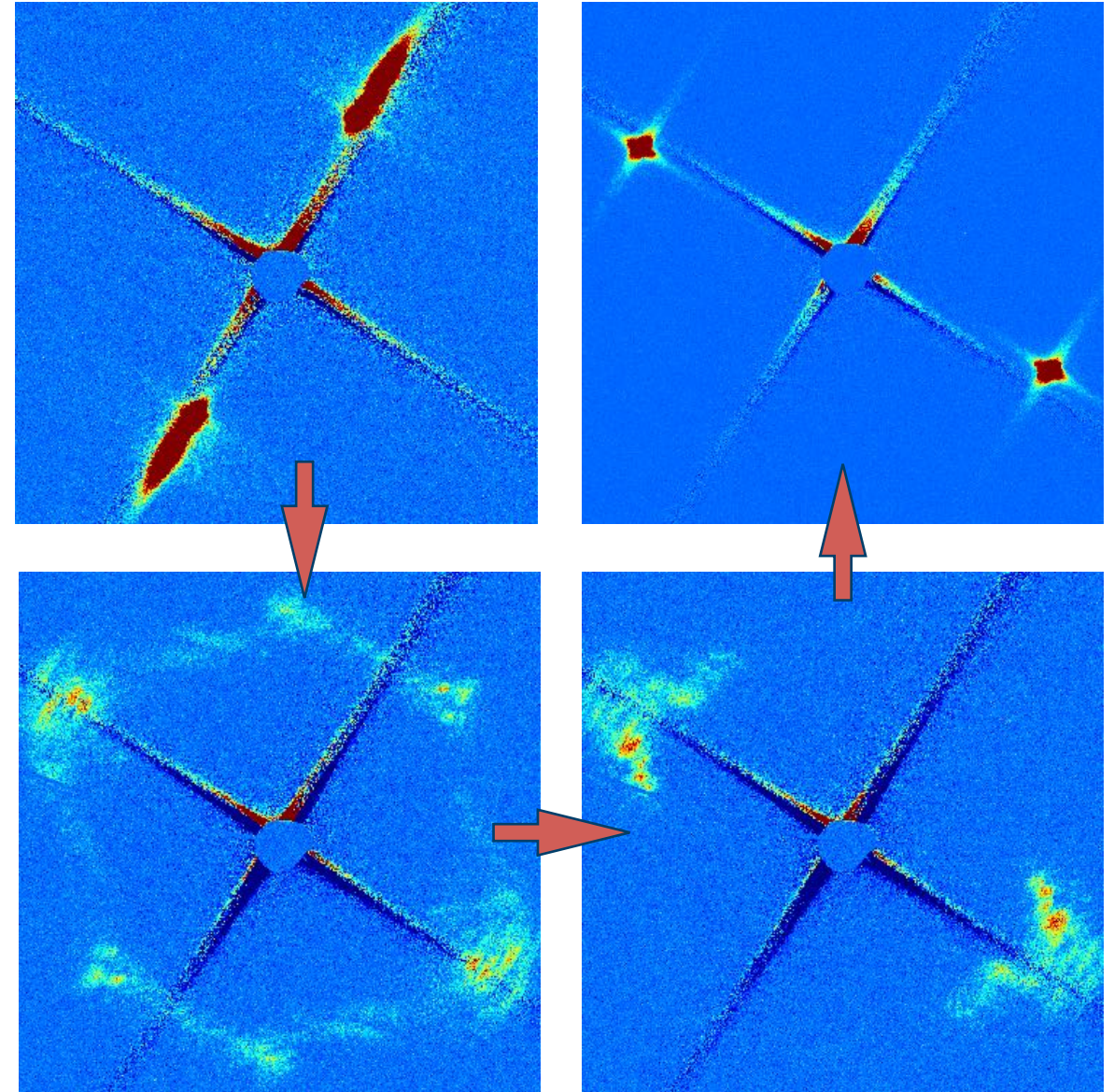
History dependence of the spiral switching:

- (1) Prepare a single domain by the corresponding in-plane field
- (2) Apply a finite field in the direction of k
- (3) The selected domain re-orientates to the orthogonal direction

Orthogonal helices in a D_{2d} system have opposite chiralities (in principle)



To change the helical propagation vector from $[100]$ to $[010]$ the system goes through an intermediate triple- Q state



CONCLUSIONS

- ❑ Resonant SAXS is an excellent tool study long-periodic spin textures
- ❑ Element selectivity is naturally provided
- ❑ Flexible sample environment
- ❑ Sensitive to very small sample volumes, e.g. $1 \times 1 \times 0.1 \mu\text{m}^3$
- ❑ Requires rather complex sample preparation
- ❑ Complementary to neutron scattering and electron microscopy

- ❑ Allows to extract weak parameters of the spin Hamiltonian (AEI) in cubic chiral magnets (FeGe, Cu_2OSeO_3 , CoZnMn, FeCoSi, etc.)
- ❑ Adds new knowledge of the spin spiral behaviours in (not quite) D_{2d} systems

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Thank you for your attention!



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