



Magnetic imaging on the nanometer scale

Quantum Functionalities of Nanomagnets

Ingelheim, Germany

18.06.2025

Prof. Martino Poggio

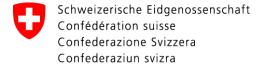


poggiolab.unibas.ch











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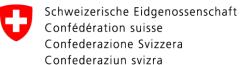






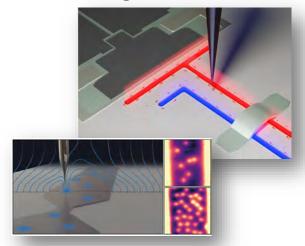




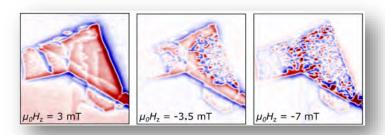




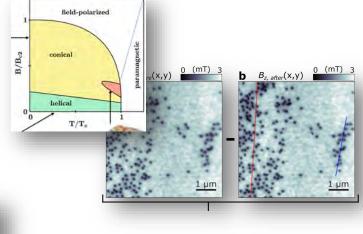
Superconducting thin films & circuits



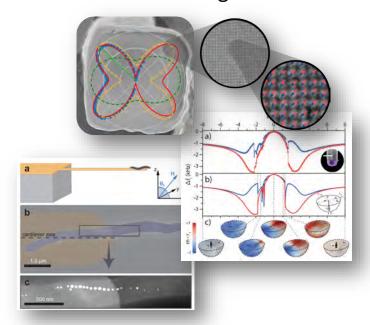
2D magnetic materials



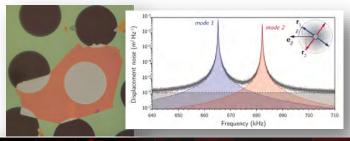
Chiral magnets



Nanomagnets

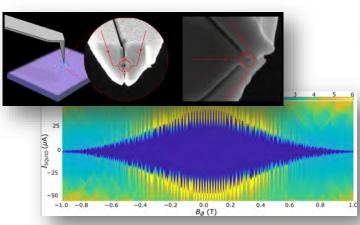


Nano & Opto-mechanics



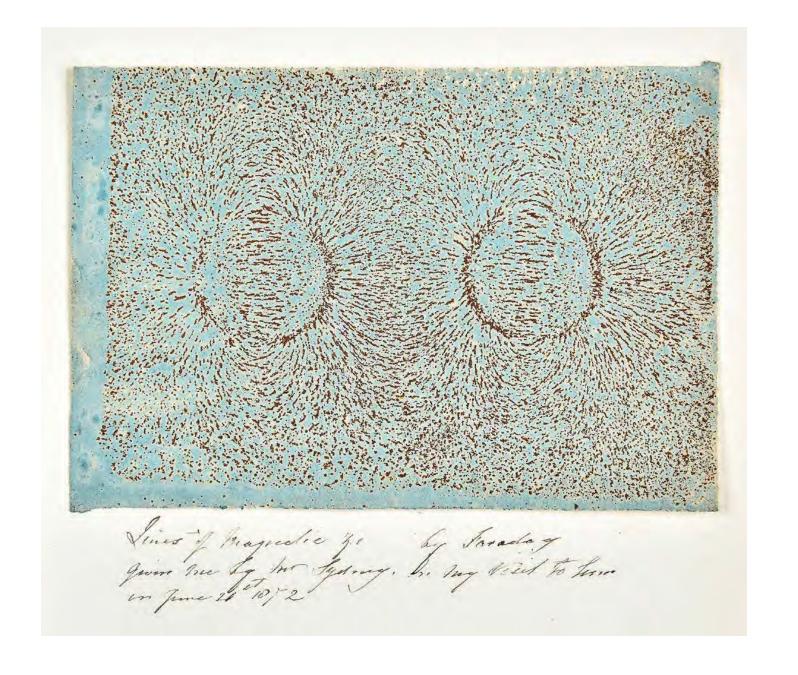


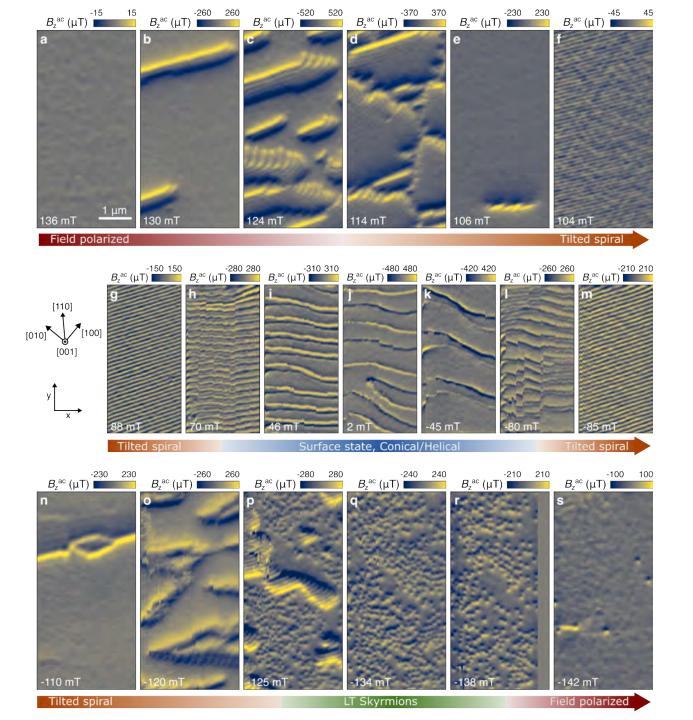
Superconducting sensors

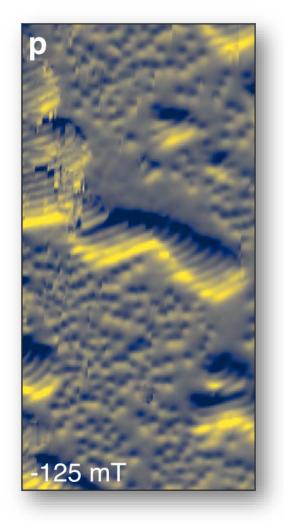


Outline

- Introduction
- Imaging magnetism via magnetic fields
- Magnetic force microscopy (MFM)
 - Bilayer EuGe₂
- Scanning SQUID microscopy (SSM)
 - Few-layer CrSBr
 - Monolayer CPS₄
- Conclusion







Surface of Cu₂OSeO₃, Marchiori et al., *Commun. Mater.* **5**, 202 (2024)

Energence of 2D materials and vdW heterostructures

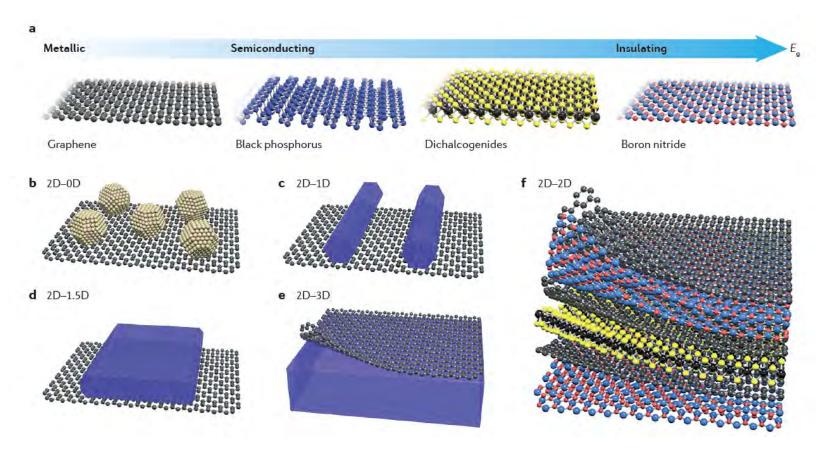


Figure 1 | **Two-dimensional layered materials and van der Waals heterostructures. a** | A broad library of two-dimensional layered materials (2DLMs) with varying chemical composition, atomic structures and electronic properties, with an increasing bandgap from left to right. $\mathbf{b} - \mathbf{f}$ | Van der Waals heterostructures formed by integrating the dangling-bond-free 2DLMs with 0D nanoparticles or quantum dots (panel \mathbf{b}), 1D nanowires (panel \mathbf{c}), 1.5D nanoribbons (panel \mathbf{d}), 3D bulk materials (panel \mathbf{e}) and 2D nanosheets (panel \mathbf{f}).

Energence of 2D materials and vdW heterostructures

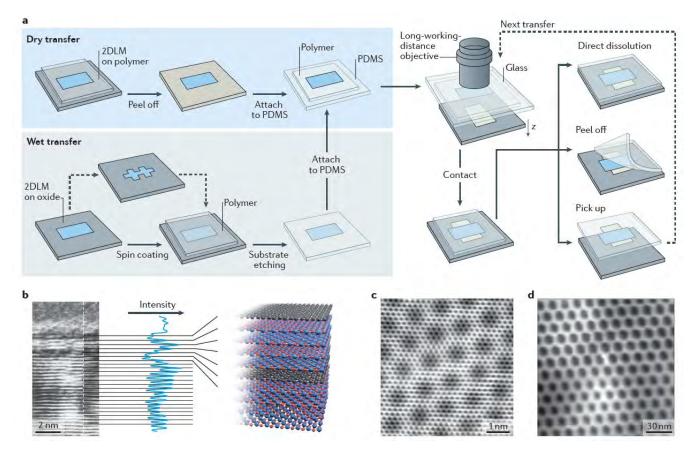


Figure 2 | Assembly and characterization of 2D–2D vdWHs. a | Schematic illustration of state-of-the-art alignment transfer processes for van der Waals heterostructure (vdWH) integration. Wet and dry transfer techniques are used to attach the target sheet to the stamp material. The stamp is then attached to a glass slide and placed in a transfer microscope. Micromanipulators allow for the precise alignment of sheets using a long-working-distance objective lens. The polymer transfer stamp can either be chemically dissolved away, mechanically peeled off or used to pick up the entire stack for further transfer steps. b | False-coloured high-resolution cross-sectional scanning tunnelling electron microscopy image of the BN-graphene–BN-graphene stack (left) and a corresponding schematic representation (right). c,d | Moiré pattern of graphene on BN (panel c) and a much larger moiré pattern of the commensurate-incommensurate transition of graphene on BN (panel d). 2DLM, two-dimensional layered material; BN, boron nitride; PDMS, poly(dimethyl siloxane). Panel b is from REF. 71, Nature Publishing Group. Panel c is courtesy of Brian LeRoy, University of Arizona, USA. Panel d is from REF. 73, Nature Publishing Group.

Correlated states in atomically layered materials

ARTICLE

doi:10.1038/nature26160

Unconventional superconductivity in magic-angle graphene superlattices

Yuan Cao¹, Valla Fatemi¹, Shiang Fang², Kenji Watanabe³, Takashi Taniguchi³, Efthimios Kaxiras^{2,4} & Pablo Jarillo-Herrero¹

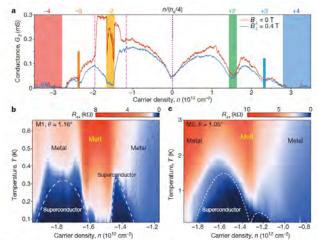


Figure 2 | Gate-tunable superconductivity in magic-angle TBG. a, Two-probe conductance $G_2 = l T V_{\rm bias}$ of device M1 $(\theta = 1.16^\circ)$ measured in zero magnetic field (red) and at a perpendicular field of $B_1 = 0.4$ T (blue). The curves exhibit the typical V-shaped conductance near charge neutrality (n = 0), vertical purple dotted line) and insulating states at the superlattice bandgaps $n = \pm n_o$, which correspond to filling ± 4 electrons in each motife unit cell (blue and red bars). They also exhibit reduced conductance at intermediate integer fillings of the superlattice owing to Coulomb interactions (other coloured bars). Near a filling of -2 electrons per unit cell, there is considerable conductance enhancement at zero field that is superseed in $B_1 = 0.4$ T. This enhancement sizenals the onset of

superconductivity. Measurements were conducted at $70\,\mathrm{mK}$; $V_{bias} = 10\,\mu\mathrm{V}$. b, Four-probe resistance R_{ss} measured at densities corresponding to the region bounded by pink dashed lines in a, versus temperature. Two superconducting domes are observed next to the half-filling state, which is labelled 'Mott' and centred around $-n_t/2 = -1.58 \times 10^{12}\,\mathrm{cm}^{-2}$. The remaining regions in the diagram are labelled as 'metal' owing to the metallic temperature dependence. The highest critical temperature observed in device M1 is $T_c = 0.5\,\mathrm{K}$ (at 50% of the normal-state resistance). c, As in b, but for device M2, showing two asymmetric and overlapping domes. The highest critical temperature in this device is $T_c = 1.7\,\mathrm{K}$.

LETTER

doi:10.1038/nature22391

Layer-dependent ferromagnetism in a van der Waals crystal down to the monolayer limit

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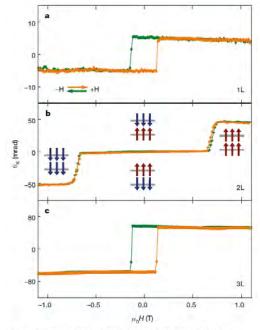
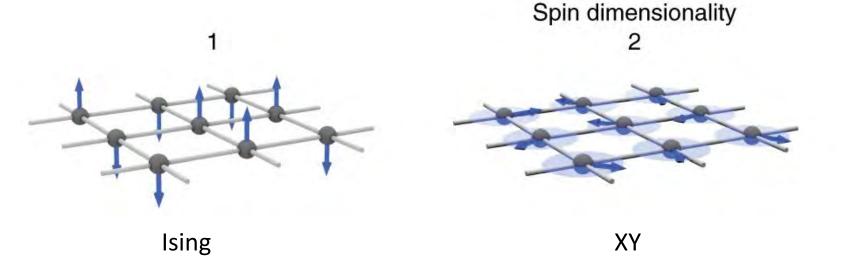
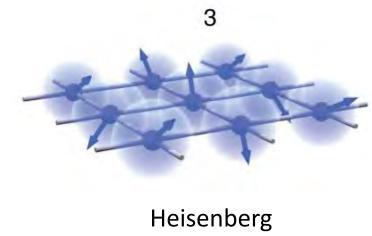


Figure 3 | Layer-dependent magnetic ordering in atomically-thin Crl_3 , a, MOKE signal on a monolayer (1L) Crl_3 flake, showing hysteresis in the Kerr rotation as a function of applied magnetic field, indicative of ferromagnetic behaviour. b, MOKE signal from a bilayer Crl_3 showing vanishing Kerr rotation for applied fields ± 0.65 T, suggesting antiferromagnetic behaviour. Insets depict bilayer (2L) magnetic ground states for different applied fields. ϵ , MOKE signal on a trilayer (3L) flake, showing a return to ferromagnetic behaviour.

Magnetism in two-dimensions





Model phase transitions

Unique topology

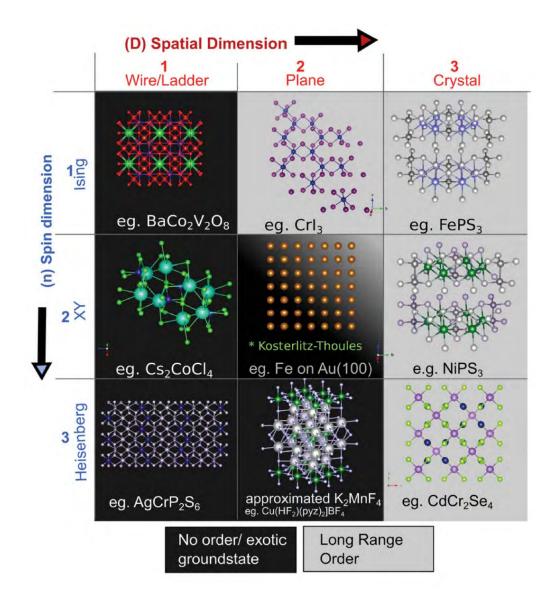
Phases governed by the Berezinski-Kosterlitz-Thouless

(BKT) transition

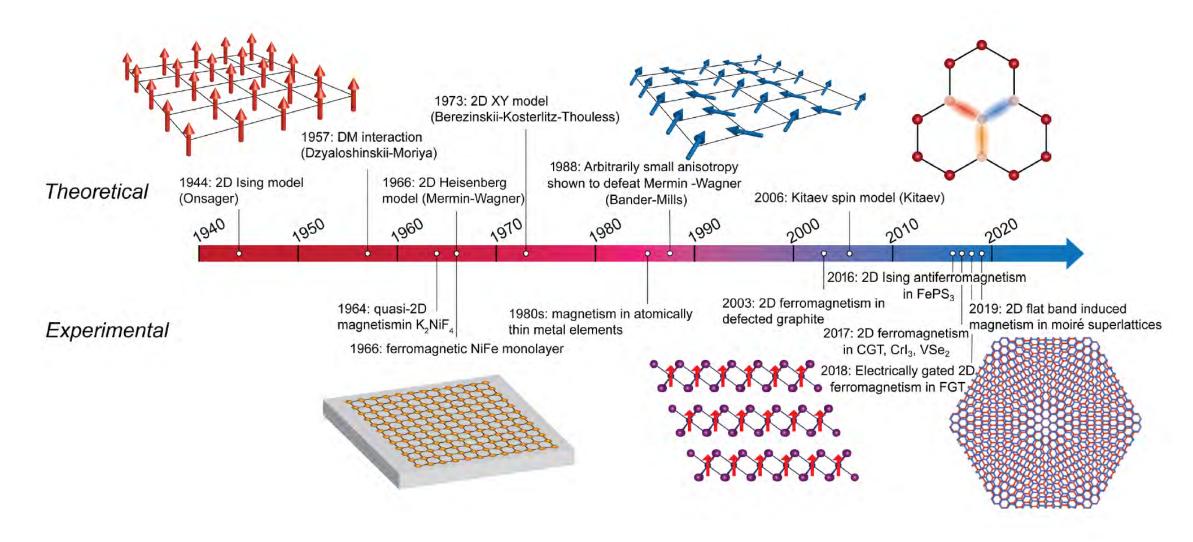
Mermin Wagner theorem – long range order?

Existence of magnetic anisotropy becomes crucial for ordering

Magnetic order



Two-dimensional magnetism timeline



How should we measure 2D magnets?

 Measurement techniques: neutron scattering, muon spin rotation, NMR, ESR, x-ray scattering, magnetization and susceptibility, STM, heat capacity & thermal transport

- 2D magnets contain few magnetic moments
- Correlation lengths are short

Need techniques combining high magnetic sensitivity and spatial resolution

How to measure the magnetism of 2D systems?

Map weak magnetic field patterns with high spatial resolution

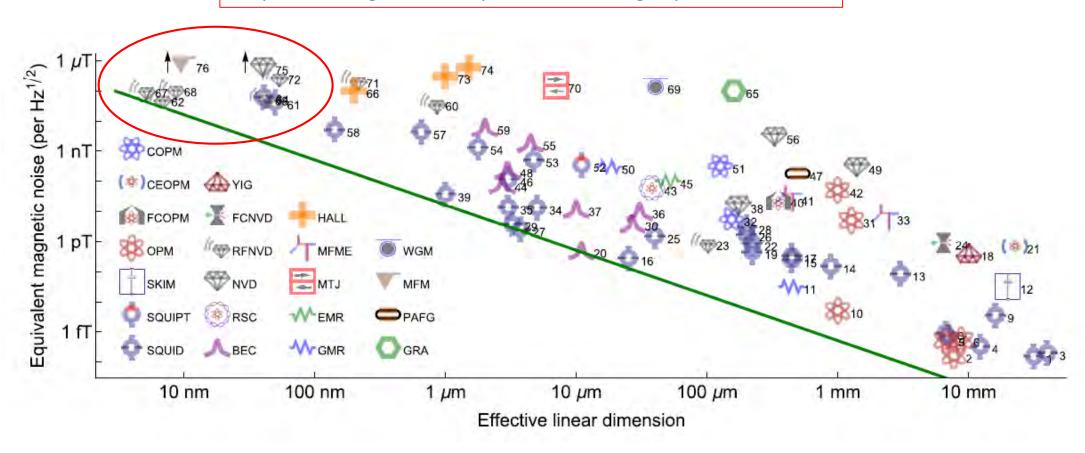
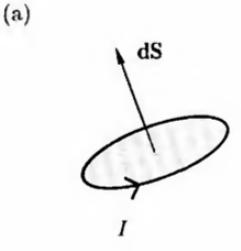


Fig. 1.1 (a) An elementary magnetic moment. $d\mu = IdS$, due to an elementary current loop. (b) A magnetic moment $\mu = I \int dS$ (now viewed from above the plane of the current loop) associated with a loop of current I can be considered by summing up the magnetic moments of lots of infinitesimal current loops.



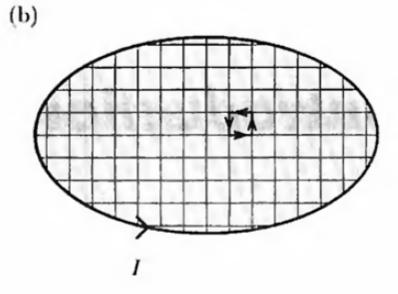
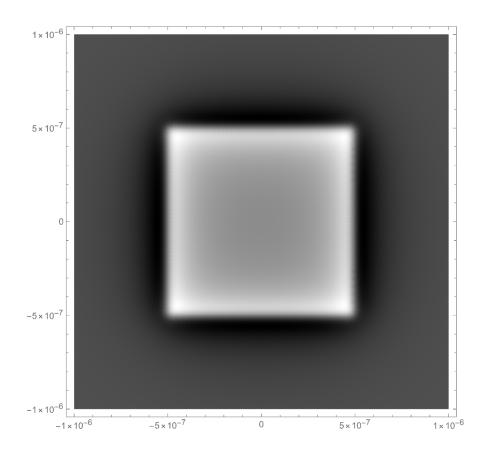
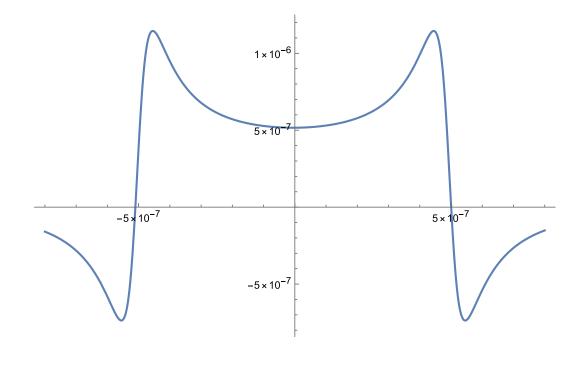


Image the edges

Out-of-plane stray field from a uniformly polarized 2D magnet









Nanoscale magnetic field imaging for 2D materials

Estefani Marchiori¹, Lorenzo Ceccarelli¹, Nicola Rossi¹, Luca Lorenzelli², Christian L. Degen 6 and Martino Poggio 6 1,3 🖂

Abstract | Nanoscale magnetic imaging can provide microscopic information about length scales, inhomogeneity and interactions of materials systems. As such, it is a powerful tool to probe phenomena such as superconductivity, Mott insulating states and magnetically ordered states in 2D materials, which are sensitive to the local environment. This Technical Review provides an analysis of weak magnetic field imaging techniques that are most promising for the study of 2D materials: magnetic force microscopy, scanning superconducting quantum interference device microscopy and scanning nitrogen-vacancy centre microscopy.

Magnetic imaging by "force microscopy" with 1000 Å resolution

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(Received 19 December 1986; accepted for publication 19 March 1987)

We describe a new method for imaging magnetic fields with 1000 Å resolution. The technique is based on using a force microscope to measure the magnetic force between a magnetized tip and the scanned surface. The method shows promise for the high-resolution mapping of both static and dynamic magnetic fields.

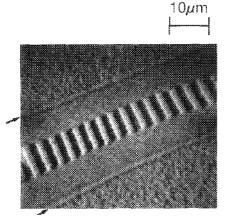
Appl. Phys. Lett. 50 (20), 18 May 1987

Magnetic force microscopy: General principles and application to longitudinal recording media

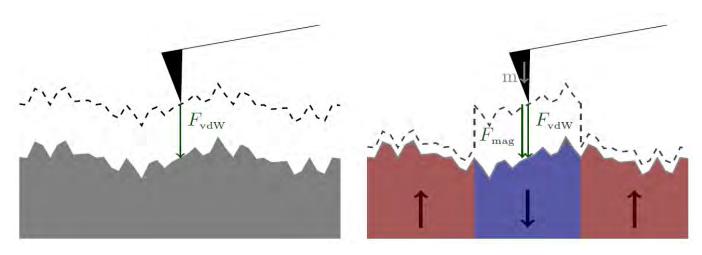
D. Rugar, H. J. Mamin, P. Guethner, S. E. Lambert, J. E. Stern, J. E. Stern, McFadyen, and T. Yogi^{b)}

IBM Research Division, Almaden Research Center, 650 Harry Road, San Jose, California 95120-6099

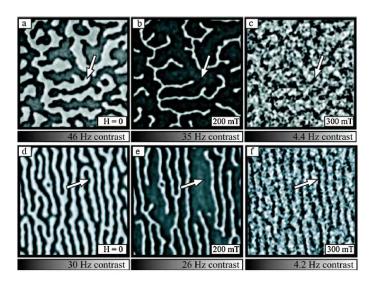
(Received 15 January 1990; accepted for publication 13 April 1990)



MFM achieves down to 10 nm resolution

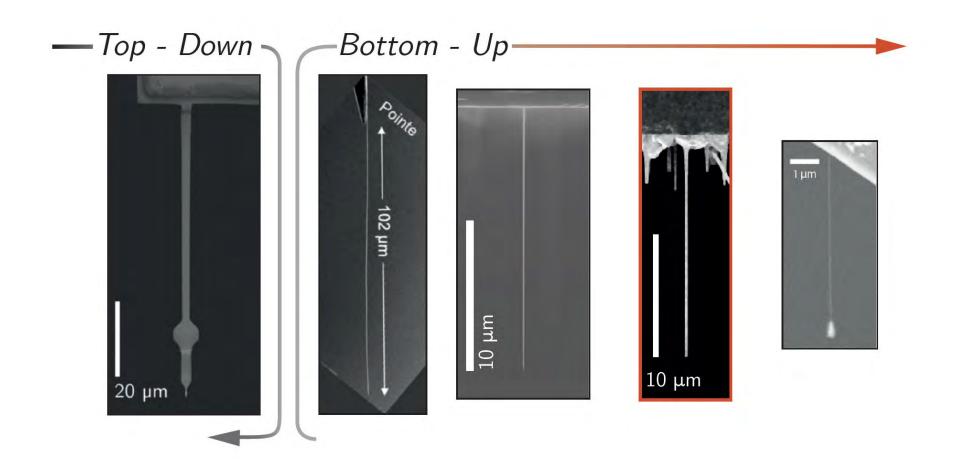


Schwenk, Ph.D. Thesis in Physics, University of Basel (2016).



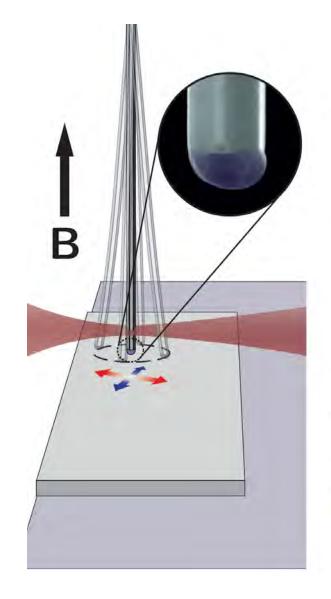
Schmid et al., Phys. Rev. Lett. 105, 197201 (2010).

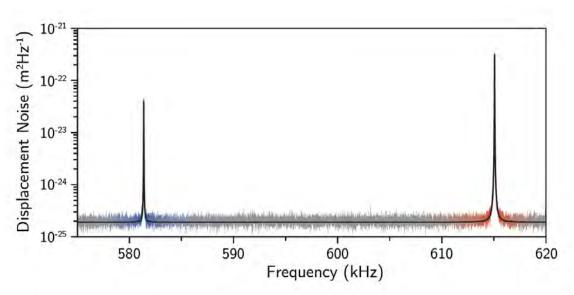
Nanowires as force sensors and scanning probes



Rossi, Ph.D. Thesis in Physics, University of Basel (2019).

NWs with magnetic tips

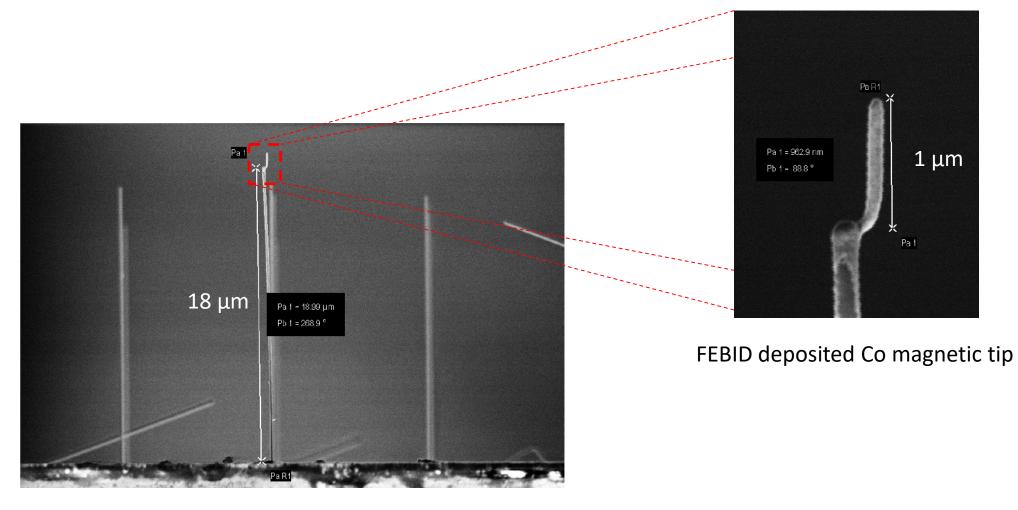




Length \sim 17 \pm 1 μ m with diameter \sim 225 \pm 15 nm

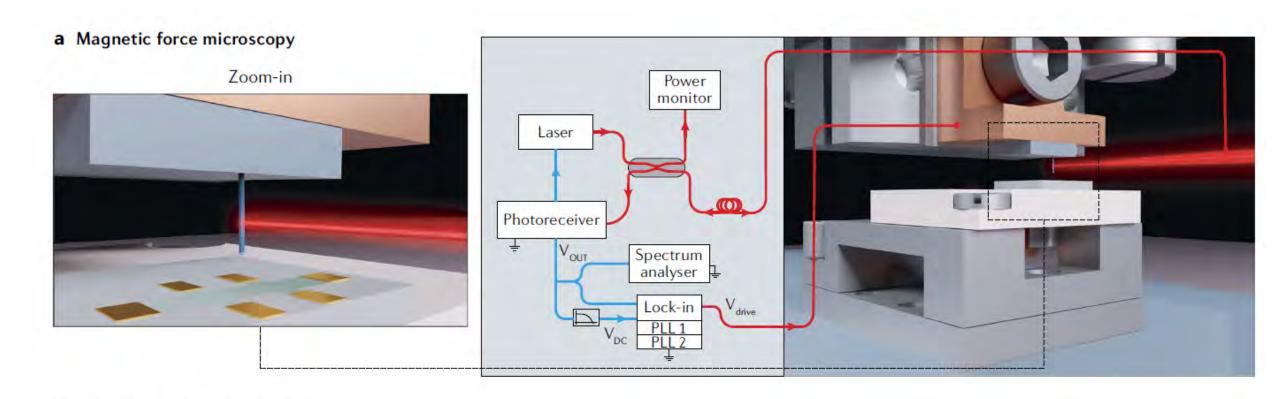
- $f \sim [500 700] \text{ kHz}$ $M \sim 700 \text{ fg}$
- $Q \sim [30 50] \times 10^3$ $\Gamma \sim 50 \times 10^{-15} \text{ kg/s}$
- $k \sim [1-10] \text{ mN/m}$ $F_{min} \sim 4 \text{ aN/}\sqrt{\text{Hz}}$

Magnet-tipped nanowires as MFM probes



Si NWs (Budakian Group, Waterloo)

MFM



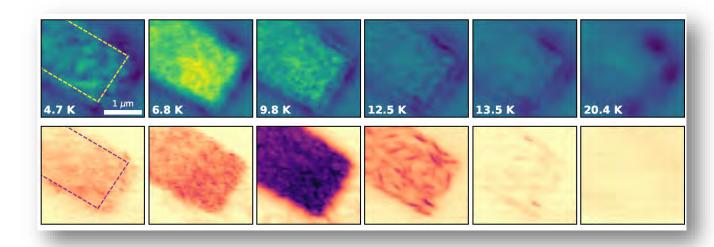
Magnetic Force Microscopy of bilayer EuGe₂



Hinrich Mattiat
Ph.D. Student

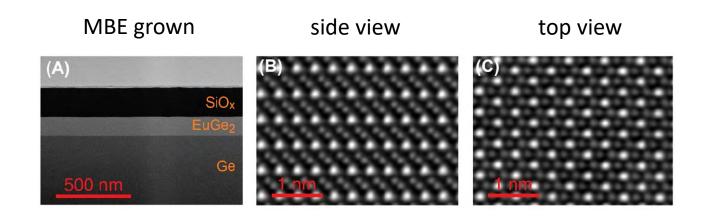


Lukas Schneider Ph.D. Student

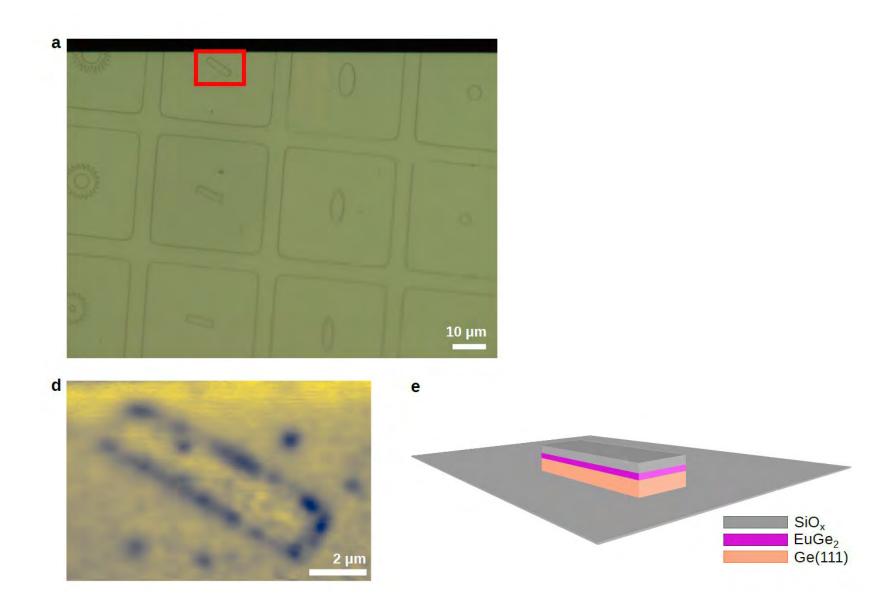


Magnetism in 2D EuGe₂

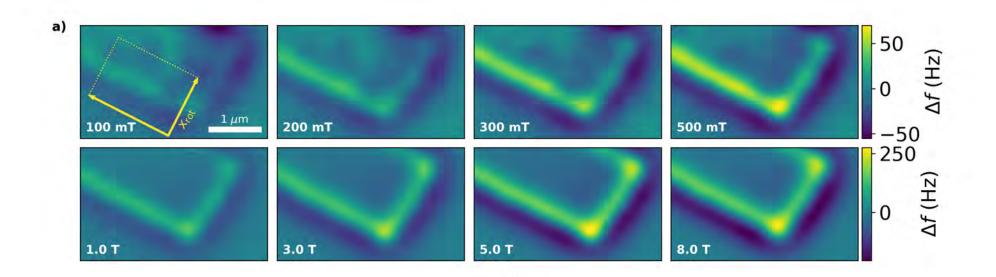
- Monolayer exhibits in-plane ferromagnetism (FM)
- Multi-layer (bulk) is stacked antiferromagnetically (AFM)
- layer dependent transition from AFM to FM evolves gradually from bulk to the monolayer, with evidence of a coexistence of both AFM and FM orders (exchange bias)



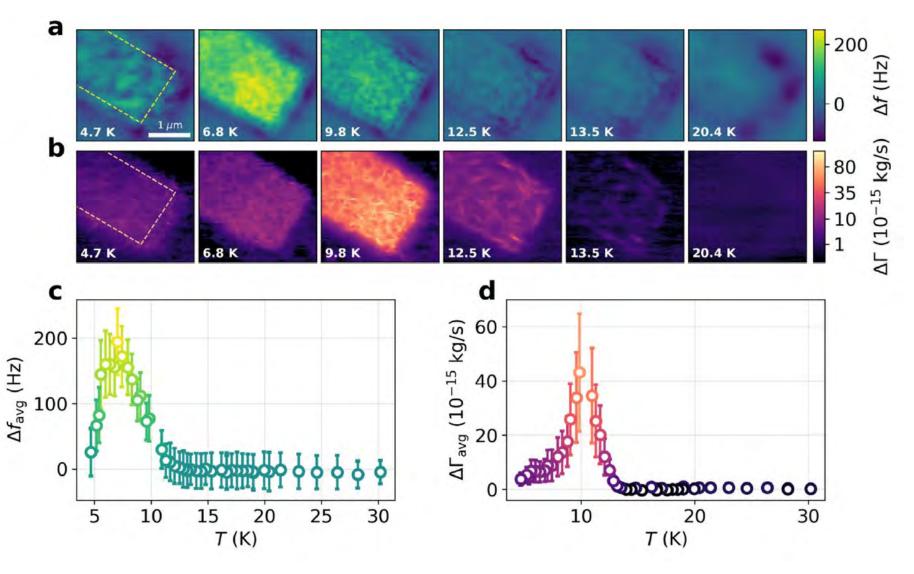
Bilayer EuGe₂: Samples



Bilayer EuGe₂: Applying out-of-plane field

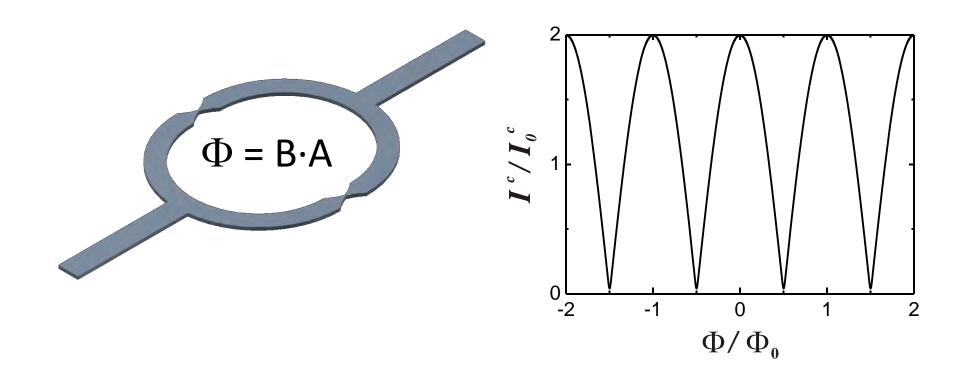


Bilayer EuGe₂: Temperature dependence



Mattiat et al., Nanoscale 16, 5302 (2024).

Superconducting quantum interference device (SQUID)



SQUID critical current: $I^c(\Phi) = 2I_0^C \cos|\Box|/\Phi_0|$

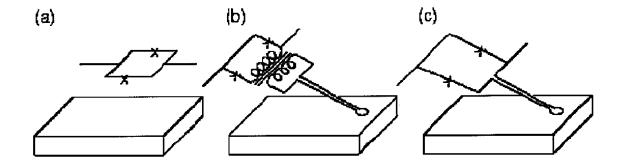
SCANNING SQUID MICROSCOPY

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John P. Wikswo, Jr.

Department of Physics and Astronomy, Vanderbilt University, Nashville, Tennessee 37235; e-mail: wikswojp@ctrvax.vanderbilt.edu



Annu. Rev. Mater. Sci. 1999. 29:117-48

2024 roadmap on magnetic microscopy techniques and their applications in materials science

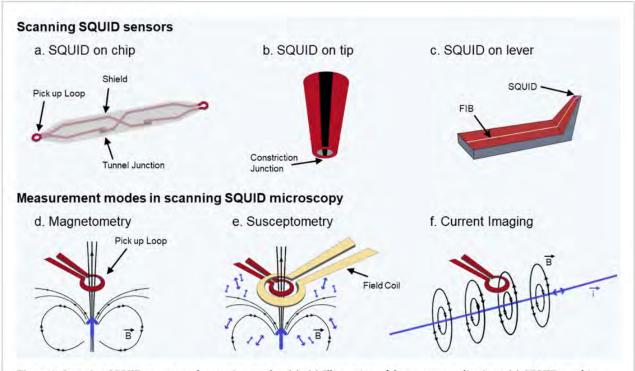
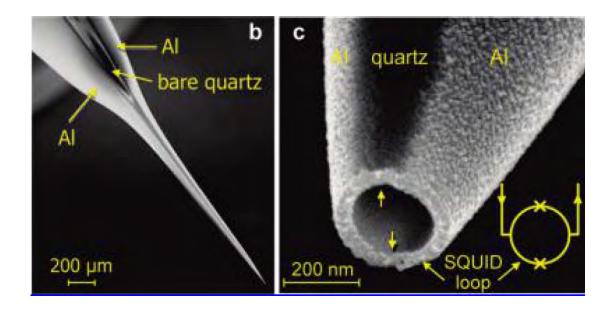


Figure 1. Scanning SQUID sensors and operation modes. (a)—(c) Illustration of three sensor realisations: (a) SQUID on chip (SOC) (b) SQUID on tip (SOT) (c) SQUID on lever (SOL). (d)—(f) Measurement modes in SSM. In (d) magnetometry, the SQUID's pick-up loop captures the magnetic field lines penetrating its area. By rastering the pick-up loop above the sample's surface we generate a map of the local magnetic flux. In (e) susceptometry, current flowing through a single turn field coil concentric with the pick-up loop creates a magnetic field at the sample's surface. The sample's response to this field is sensed by the pick-up loop. By modulating the field coil current we lock into the response, allowing for simultaneous measurement of magnetometry and susceptometry. In (f) current imaging, the magnetic fields generated by the current flow are imaged, allowing for reconstruction of the current density map.

Self-Aligned Nanoscale SQUID on a Tip

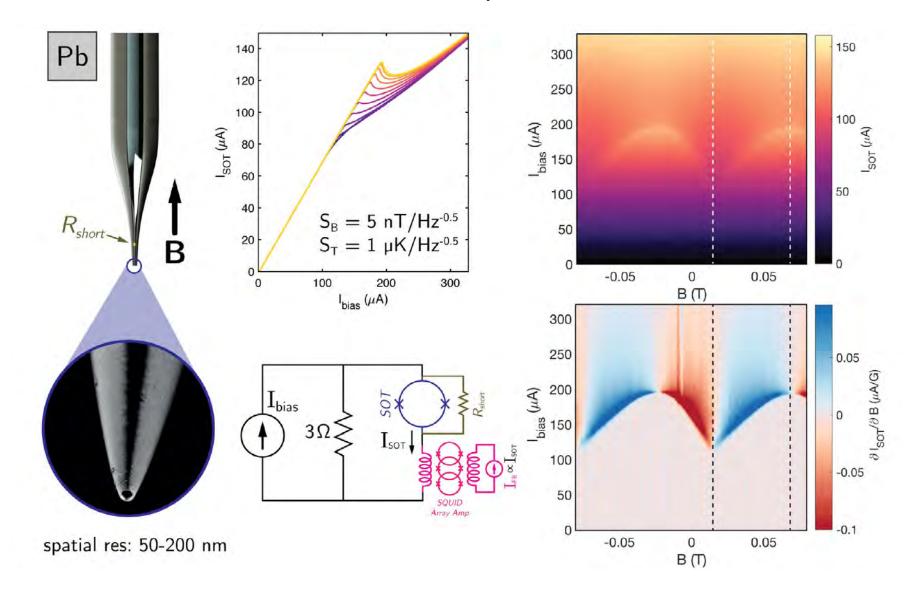
Amit Finkler,*',† Yehonathan Segev,† Yuri Myasoedov,† Michael L. Rappaport,† Lior Ne'eman,† Denis Vasyukov,† Eli Zeldov,† Martin E. Huber,† Jens Martin,§ and Amir Yacoby§

†Department of Condensed Matter Physics, Weizmann Institute of Science, Rehovot 76100, Israel, †Departments of Physics and Electrical Engineering, University of Colorado, Denver, Colorado 80217, and §Department of Physics, Harvard University, Cambridge, Massachusetts 02138



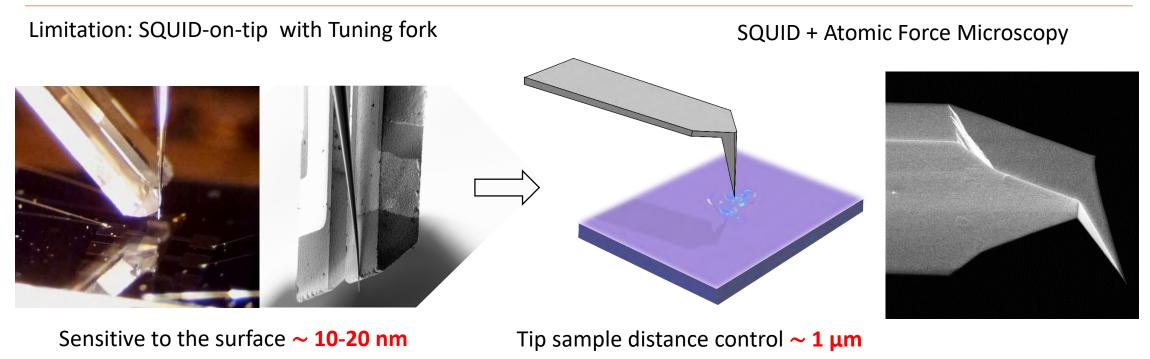
Nano Lett. **2010**, 10, 1046-1049

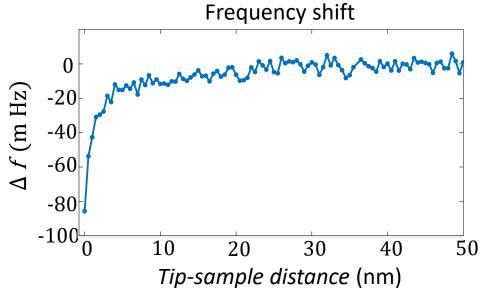
SQUID-on-tip sensor

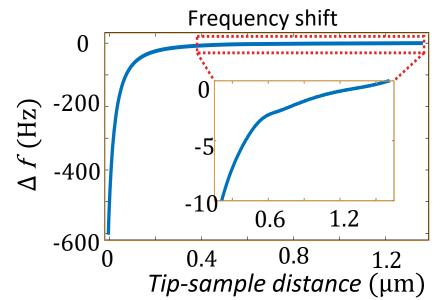


Vasyukov et al., Nat. Nanotechnol. 8, 639 (2013)

SQUID-on-Cantilever



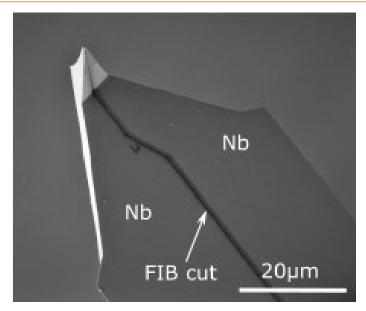


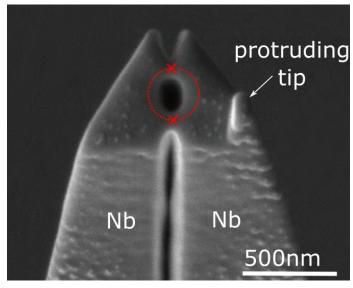


Fabrication of SQUID-on-lever



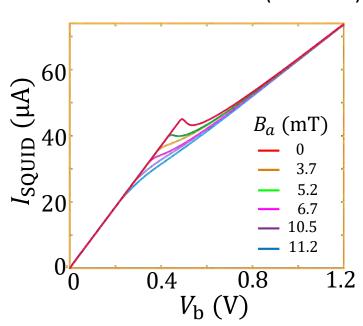
Effective diameter ~ 150 - 350 nm

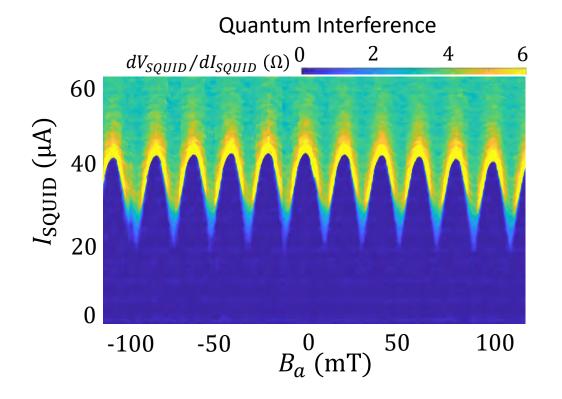




Magnetic and Thermal Sensitivity



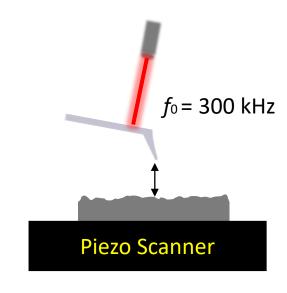




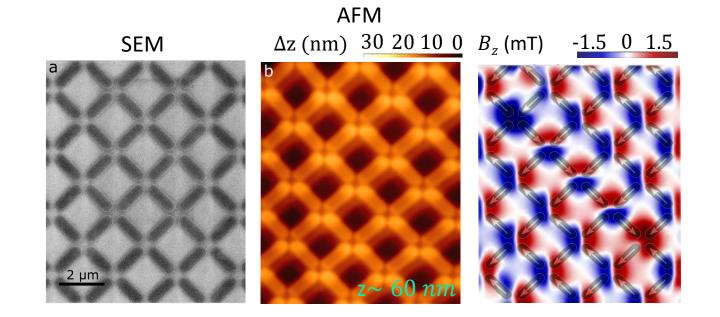
Flux Sensitivity, $S_{\Phi}^{1/2} \sim 0.2 \, \mu \Phi_0 / \mathrm{Hz}^{1/2}$

Temperature Sensitivity, $S_{\rm T}^{1/2} \sim 600 \, {\rm nK/Hz^{1/2}}$

Scanning SQUID on Cantilever Microscopy

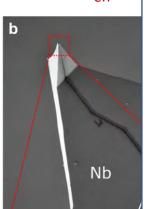


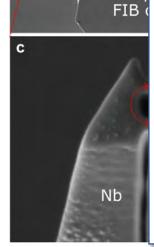
Home-built scanning probe microscope



Daducing the city of COLLID on lavor









José María De Teresa (Zaragoza)



Dieter Kölle (Tübingen)



Armin Knoll (IBM Zürich)



FiBsuperProbes

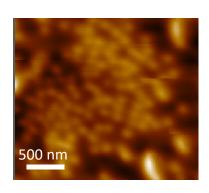


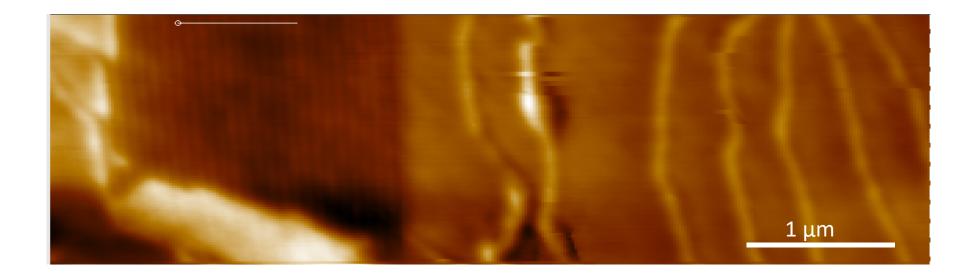
fibsuperprobes.com

Sub 100-nm spatial resolution

Cu₂OSeO₃

- Helical, conical and skyrmion phases
- Lengthscale of helimagnetic order: 62 nm





In preparation.

Scanning SQUID microscopy of strained 2D CrSBr



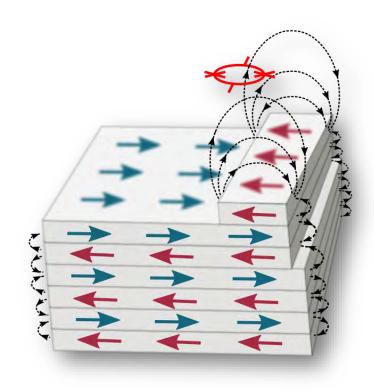
Dr. Kousik Bagani Post-doctoral Researcher



Andriani Vervelaki Ph.D. Student



Daniel Jetter
Ph.D. Student



Strain response of 2D CrSBr

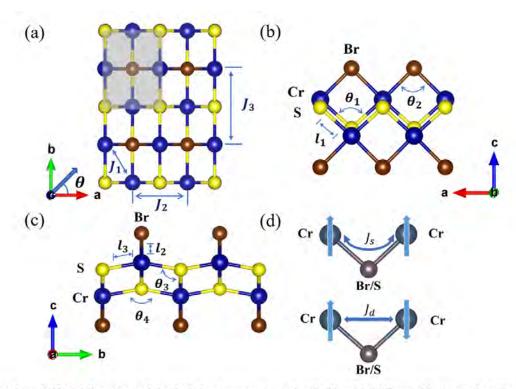
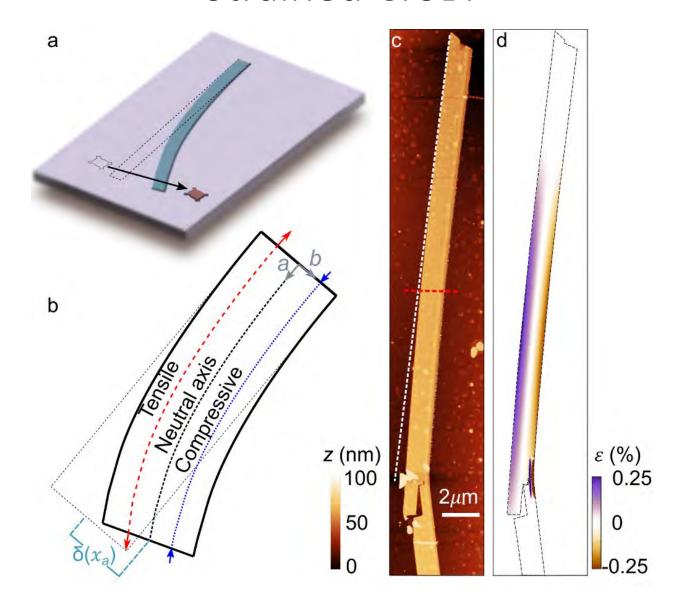
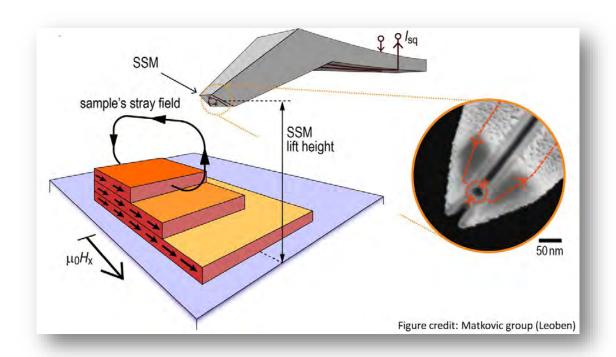


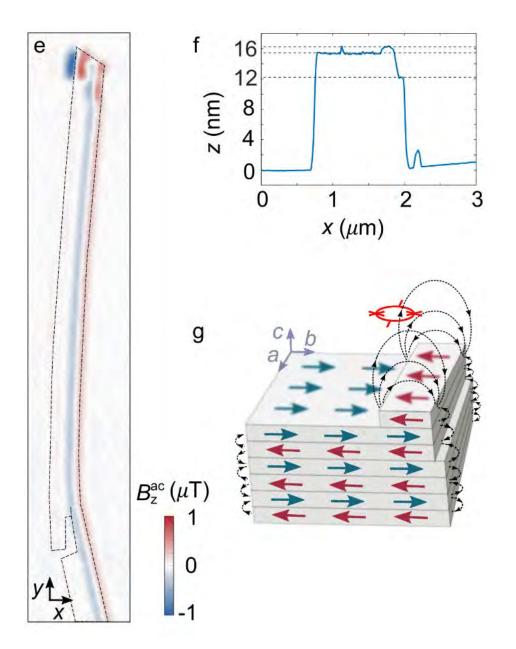
Fig. 1. Structure of monolayer CrSBr. (a) Top view, and the shaded part represents a unit cell, (b) and (c) different side views, where the distributions of atoms in three different directions are shown, and there are three different types of bond lengths l_i (i = 1, 2, 3) and four different bond angles θ_i (i = 1, ..., 4), respectively.

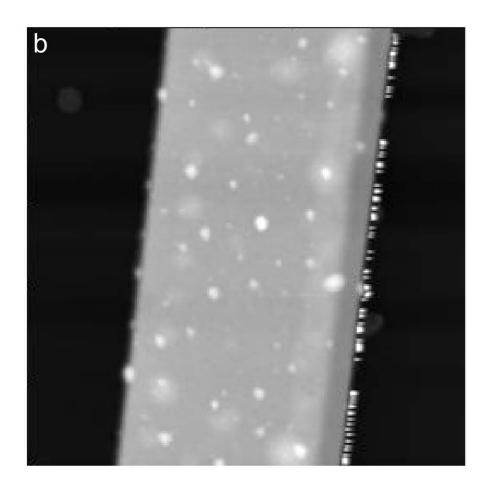
- A-type antiferromagnet
- $T_N = 132 \text{ K in bulk}$
- Magnetic anisotropy is biaxial with easy axis along a-axis
- Interlayer exchange responds to strain along a-axis
 - Compressive enhance AF
 - Tensile reduce AF

Strained CrSBr



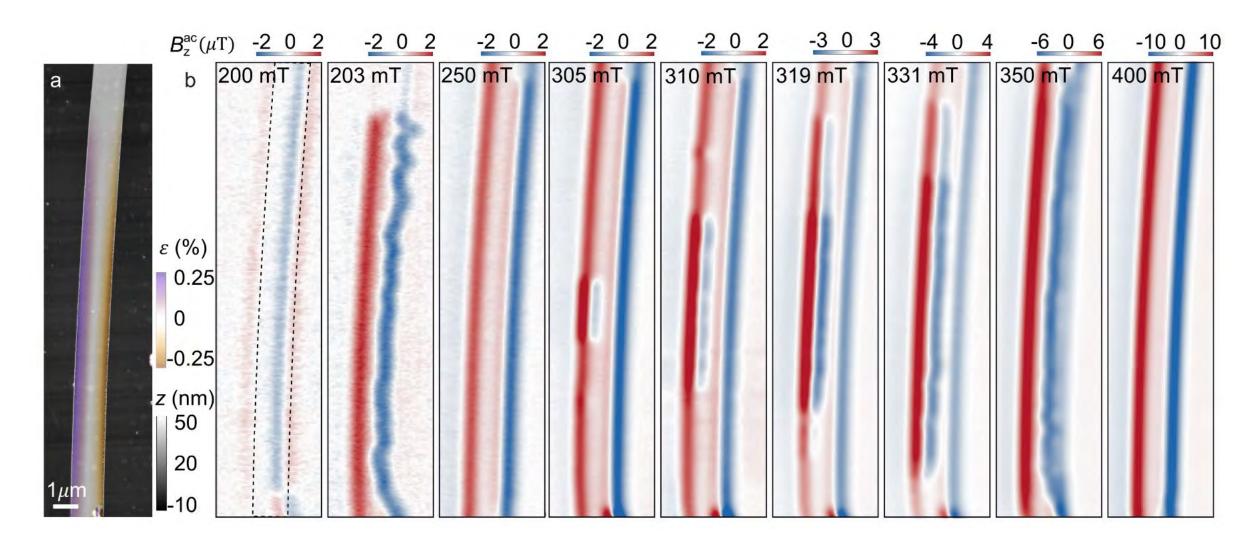


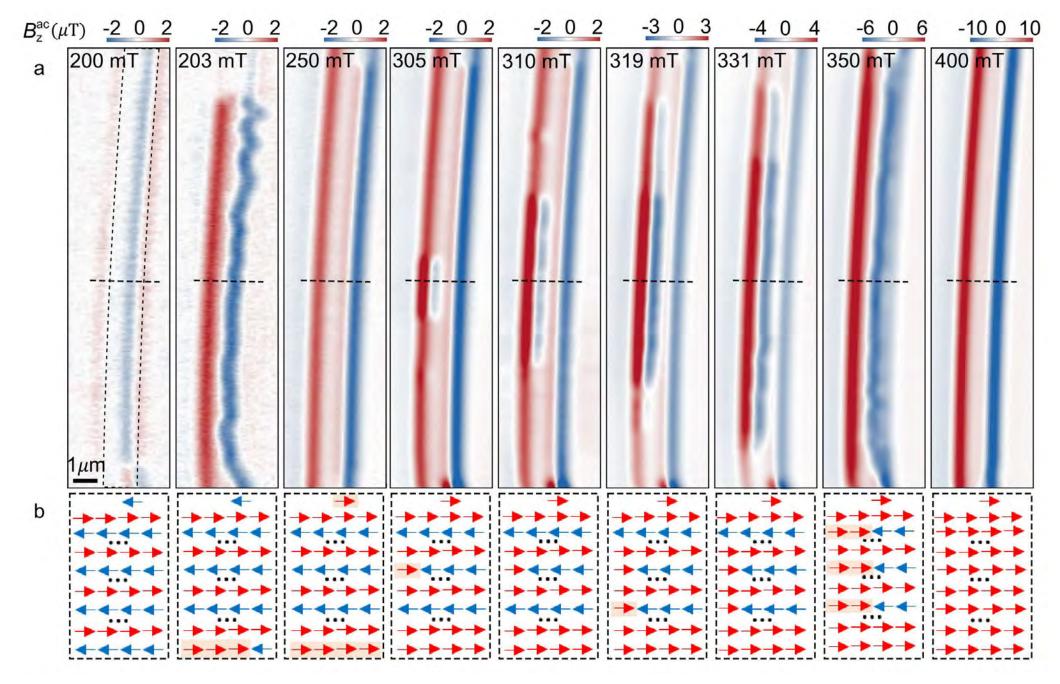




Bagani et al., Nano Lett. 24, 13068 (2024).

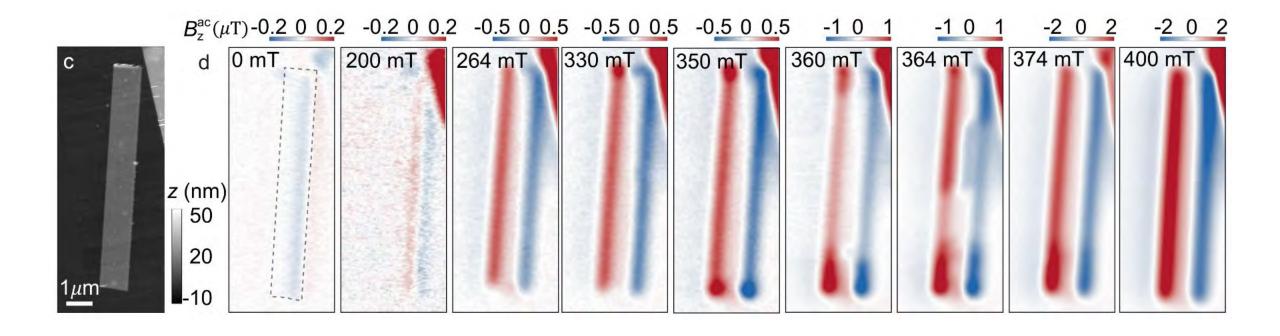
Magnetic field dependence: strained sample

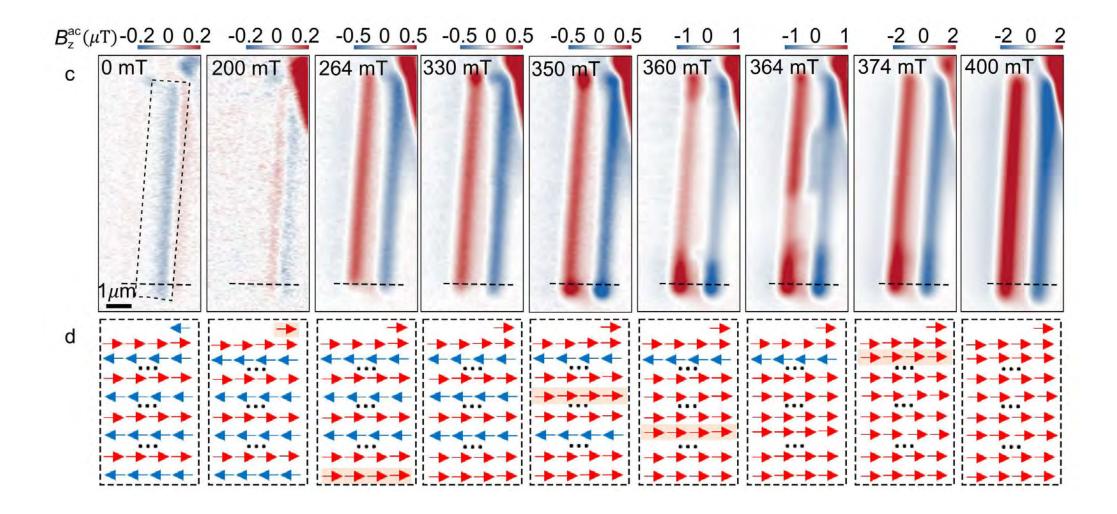




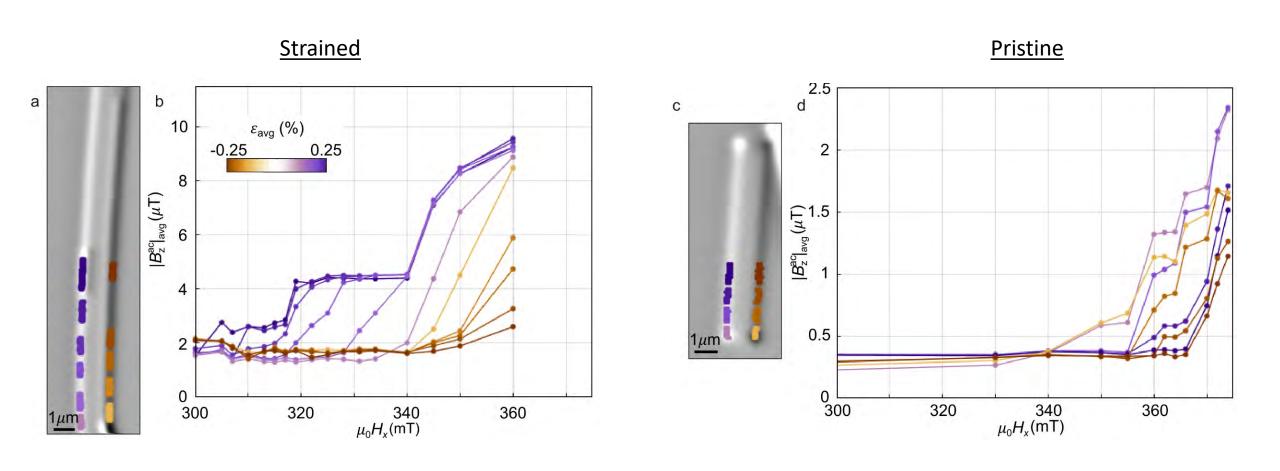
Bagani et al., Nano Lett. 24, 13068 (2024).

Magnetic field dependence: pristine sample

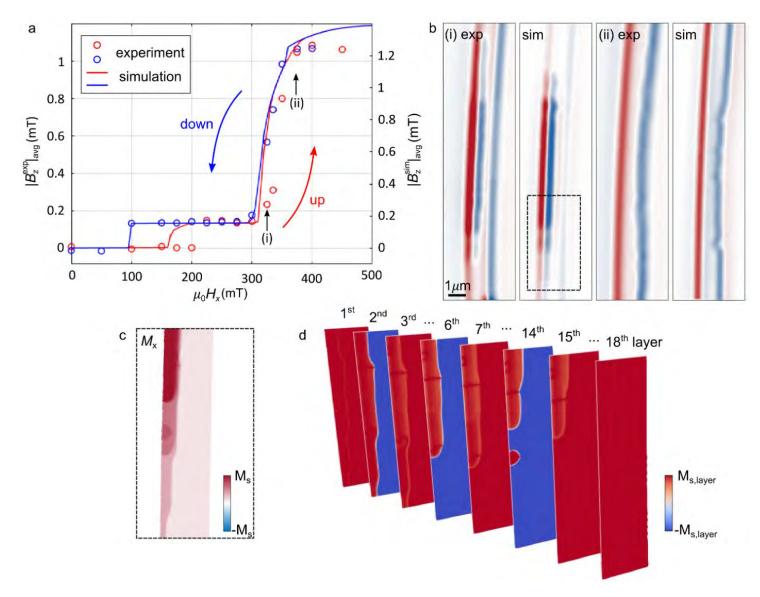




Magnetic field dependence



Comparison to micromagnetic simulations



Scanning SQUID microscopy of monolayer CrPS₄



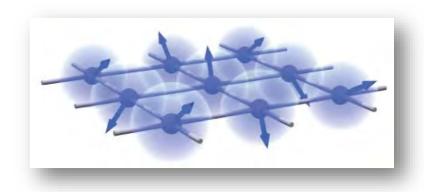
Dr. Paritosh Karnatak
Research Scientist



Dr. Andriani Vervelaki
Post-doctoral Researcher

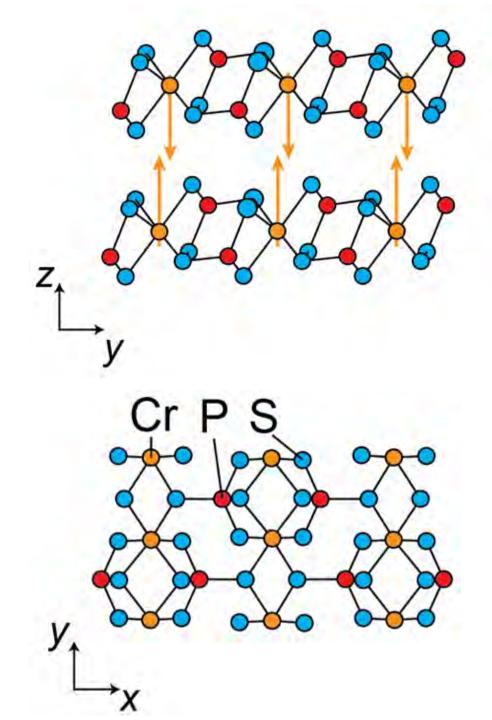


Katharina Kress
Ph.D. Student



CrPS₄

- Bulk CrPS₄ is a van der Waals (vdW) A-type semiconducting antiferromagnet
- Weak out-of-plane anisotropy
- Neel temperature in bulk $T_N = 38 \text{ K}$
- Monolayer's Curie temperature: $T_C = 23 \text{ K}$

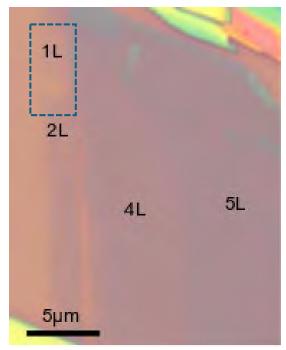


CrPS₄

- Exfoliated $CrPS_4$ (CPS) with various thicknesses: 1 10s of layers
- CPS encapsulated with h-BN

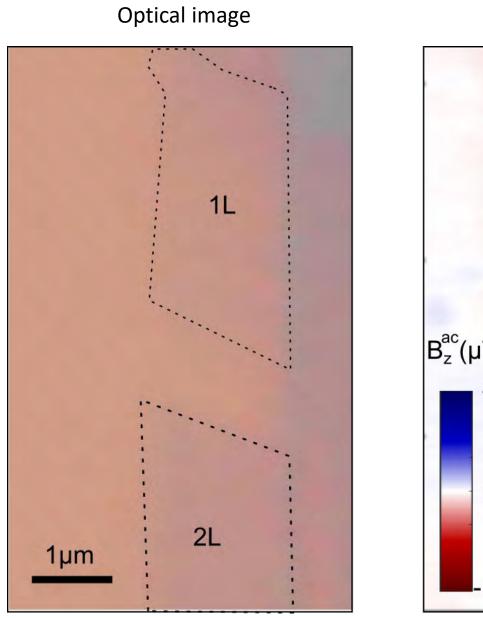
Fabrication:
Dr. Menghan Liao, Prof. Alberto F. Morpurgo
University of Geneva

Optical image before encapsulation

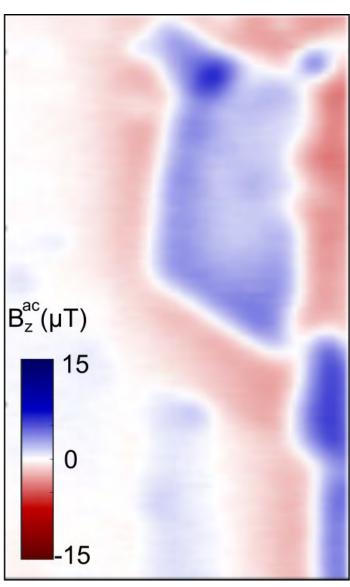


Study: Imaging of the monolayer's (1L) magnetic reversal and comparison of the behavior with thicker areas

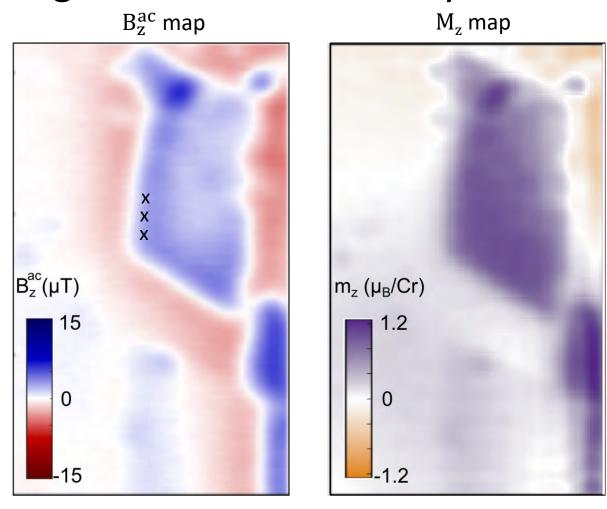
Stray field measurements

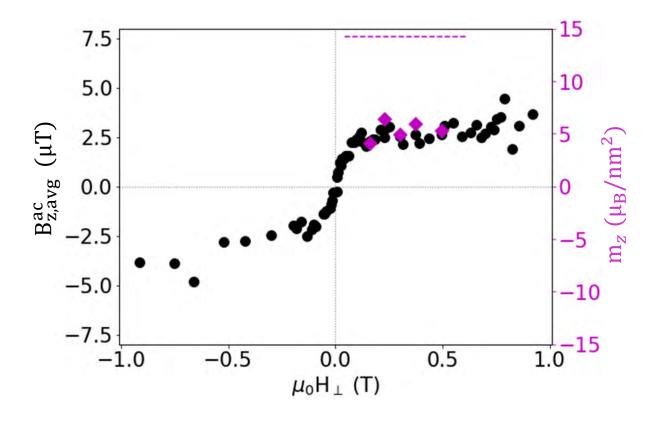


 B_{z}^{ac} map



Magnetism in the monolayer

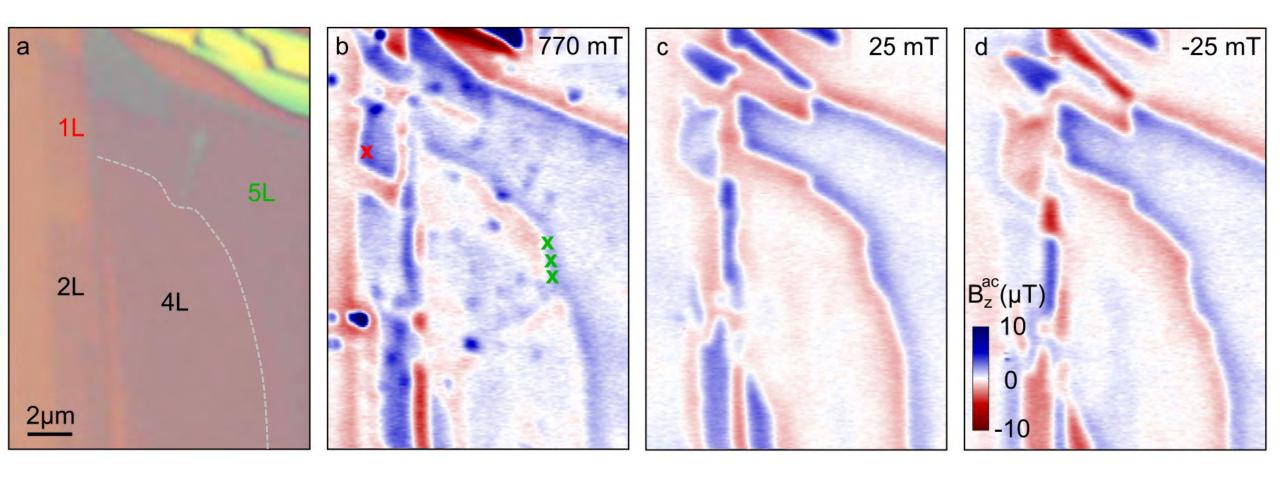




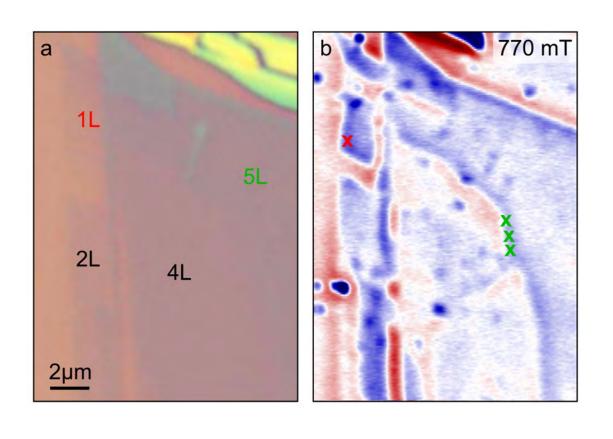
No remanence observed

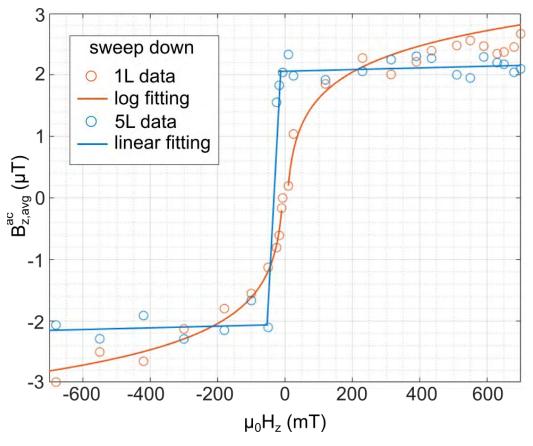
At ~ 0.5 T, $m \leq 40\%$ of the expected saturation magnetisation*

Magnetisation reversal – few layers



Magnetisation reversal – few layers





1L

- No retentivity
- Coercive field is ~ 0

5L

- $\sim 100\%$ retentivity
- 25 mT 50 mT coercive field

Recent related references from our group

Reviews

- Magnetic microscopy for materials science: J. Phys. Mater. 7, 032501 (2024)
- Magnetic field imaging for 2D materials: Nat. Rev. Phys. 4, 49 (2022)

Magnetic microscopy probes

- Nb and MoGe SQUID-on-tips: Appl. Phys. Lett. 122, 192603 (2023)
- SQUID-on-lever scanning probes: Phys. Rev. Appl. 17, 034002 (2022)
- Nanowire MFM fabricated by FEBID: Phys. Rev. Appl. 13, 044043 (2020)

Chiral Magnets

Imaging the surface of bulk Cu₂OSeO₃: Commun. Mater. 5, 202 (2024)

2D Magnets

- Imaging 2D strained CrSBr: Nano Lett. 24, 13068 (2024)
- Imaging of 2D Cr₂Ge₂Te₆: Commun. Mater. 5, 40 (2024)
- MFM of bilayer EuGe₂: Nanoscale 16, 5302 (2024)



