

Beta-decay half-lives of $^{129g,m}\text{Cd}$ and ^{133}Cd

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Beta-decay half-lives of very neutron-rich cadmium isotopes were determined at CERN-ISOLDE using the isobaric selectivity of the RILIS ion source. Beta-delayed neutron measurements resulted in an improved half-life of $^{129g,m}\text{Cd}$ and the first measurement of ^{133}Cd [1].

The region around the doubly magic ^{132}Sn has been the subject of several experimental investigations in the last years. The main reason for this big interest is the $N=82$ “waiting-point” area, where the r-process matter flow climbs up the shell closure from ^{125}Tc ($Z=43$) to ^{130}Cd ($Z=48$). The process is able to break out of the shell closure in the indium chain ($Z=49$) and rushes from there towards the next neutron-magic “waiting-point” area at $A \approx 195$.

Production and ionization of Cd isotopes

Neutron-rich medium-mass isotopes are normally produced at ISOLDE by high energy (1–1.4 GeV) proton-induced spallation of ^{238}U . This method produces also large amounts of proton-rich Rb and Cs isotopes. To avoid this problem, a “neutron-converter” was used.

The Cd isotopes were ionized by a laser ion source (RILIS) in three steps. The frequency of the first step lies in the ultra-violet of the spectrum and is done by frequency tripling with two BBO crystals. The ionization efficiency is about 10%. The disadvantage of this method is the unavoidable background of surface ionized species, in this case mostly In. This problem is not solved yet.

Half-life measurements

After ionization and isobaric separation, the nuclides were implanted on an aluminum tape of the Mainz tape station. Beta-delayed neutron (βdn) data were taken with our Mainz 4π neutron detector, consisting of 64 ^3He proportional-counter tubes arranged in three concentric rings in a polyethylene matrix.

Table 1: Comparison of experimental β -decay half-lives, literature values and QRPA model predictions

Nuclide		$T_{1/2}(\text{exp.})$	$T_{1/2}(\text{Lit.})$	$T_{1/2}(\text{QRPA})$	$T_{1/2}(\text{QRPA})$
		[ms]	[ms]	(FY) [ms]	(Nilsson) [ms]
^{129}Cd	g	242 (8)	270	457	308
	m	104 (6)	-	369	252
^{133}Cd		57 (10)	-	140	46

The results of these measurements are summarized in Table 1 and Figures 1 and 2. The results are compared with literature values and QRPA predictions from a “folded Yukawa”-potential (FY) or from the Nilsson model.

In the data of ^{129}Cd an interesting effect occurs: the delayed neutrons from the isomeric $\nu d_{3/2}$ state are mostly detected in the inner ring of the neutron coun-

ter, whereas the neutrons from the $\nu h_{11/2}$ ground state (g. s.) decay are detected mostly in the outer ring.

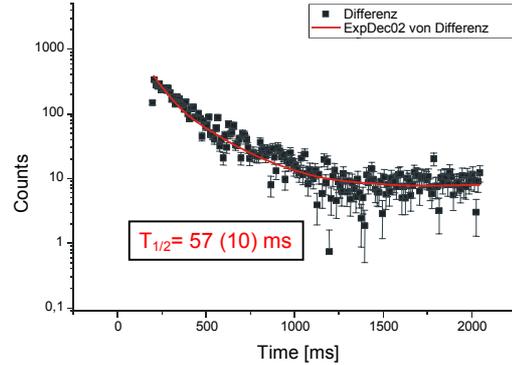


Fig. 1: Delayed neutron spectrum of ^{133}Cd with the half-life of $T_{1/2}=57$ (10) ms. The background is mainly due to the β -delayed neutron daughter ^{132}In with $T_{1/2} = 206$ (5) ms.

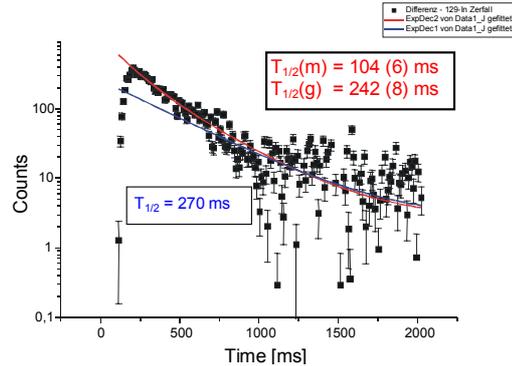


Fig. 2: Delayed neutron spectrum of ^{129}Cd . The blue line shows the simulated spectrum of ^{129}Cd with the half-life of 270 ms given in the literature.

The explanation of this effect are the different angular-momentum barriers for the emitted βdn 's. Neutrons emitted after β -decay of the $\nu d_{3/2}$ state have low angular momenta. This means, the average energy of these βdn 's is relatively low. Due to the thermalisation in the matrix of the neutron counter, these neutrons are predominantly detected in the inner ring. Neutrons emitted after the β -decay of the $\nu h_{11/2}$ g. s. have, due to their higher angular momenta and the corresponding γ -competition, higher energies. Therefore, more matrix material is required for thermalisation and the βdn 's are detected in the outer ring. This effect is clearly visible in the experimental data. In the outer ring the measured half-life is close to the literature value, in the inner ring clearly two half-lives are observed.

References

[1] O. Arndt, *Diplomarbeit* (2003)