

# Decay-Scheme of $^{130}\text{Cd}$

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With the technical improvements and results described in the preceding annual report about “First  $\gamma$ -Spectroscopic Study of the N=82 r-Process “Waiting-Point” Nucleus  $^{130}\text{Cd}$ ”, we were able to collect  $\beta\gamma$ -coincidence data and to determine the  $Q_\beta$ -value for  $^{130}\text{Cd}$ . This result completed our earlier investigations and allowed to propose a first  $\beta$ -decay scheme of  $^{130}\text{Cd}$ , that can be compared to shell-model predictions.

## Decay scheme

The most important nuclear physics quantities in this context are the excitation energy of the lowest  $1^+$  2QP state and its log ft value for Gamow-Teller (GT) feeding. The input data needed to calculate the log ft are the GT-decay energy, which is deduced from the difference between the  $Q_\beta$ -value and the  $E(1^+)$ , the  $\beta$ -decay half-life of 162 ms and the  $\beta$ -feeding of 69%. The predictions for the  $Q_\beta$ -value of  $^{130}\text{Cd}$  lie between 8.93 MeV from mass models with shell-quenching (e.g. HFB-SkP and ETFSI-Q) and 7.43 MeV from the unquenched FRDM, while the Audi and Wapstra mass evaluation [1] leads to a value of 8.5 MeV. Our experimental value of 8.32 MeV [2] was deduced from a simplified Fermi-Kurie plot (see Fig.1) and confirms those high values from quenched models. This  $Q_\beta$  value is needed in order to obtain a physically consistent picture of  $^{130}\text{Cd}$   $\beta$ -decay.

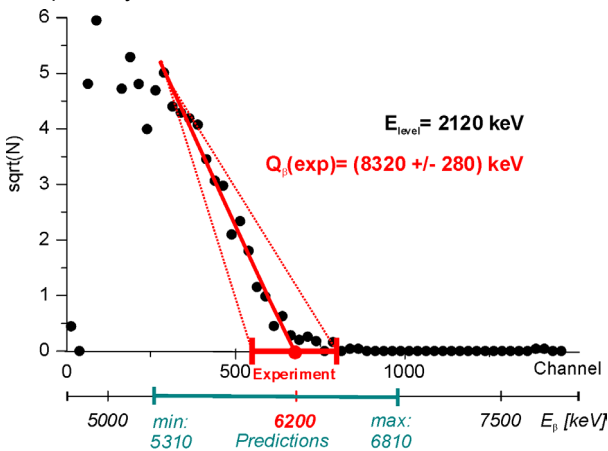


Fig. 1: Simplified Fermi-Kurie plot for  $^{130}\text{Cd}$ . Shown is the summation over all five transitions, that come from the  $1^+$  level at 2120 keV. From its x-axis intercept, a  $\beta$ -endpoint energy of 6200 keV was derived.

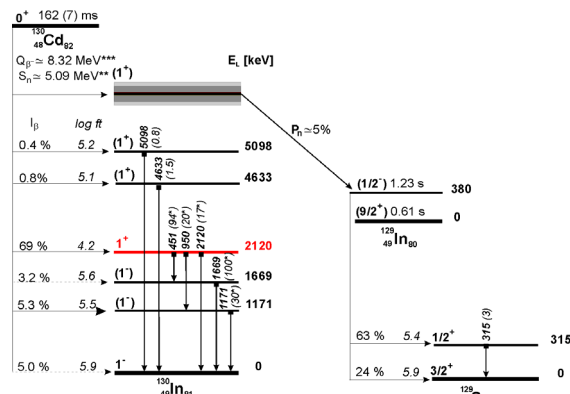


Fig. 2: Decay scheme of  $^{130}\text{Cd}$ . The  $S_n$ -value is taken from [1].

With these data, a log ft value of 4.2 for the  $0^+ \rightarrow 1^+$  GT-transition was deduced. The  $\beta$ -feeding of the ground-state suggests a first-forbidden (ff)-transition, while the log ft's of the levels at 1669 keV and 1171 keV lie between the usual values for GT- and ff-decays.

## OXBASH model predictions

Fig.3 shows a comparison of the experimental level scheme of  $^{130}\text{In}$  with two calculations of the OXBASH model [3,4]. Prediction (A) was made before knowing our results, and the position of the  $[\pi g_{9/2} \otimes \nu g_{7/2}]$  2QP  $1^+$ -state was calculated to 1382 keV. This theoretical  $E(1^+)$  was considerably lower than our experimental value of 2120 keV. In the recalculation labelled (B), the  $1^+$  level could be lifted to 1874 keV by weakening the interaction of the  $[\pi g_{9/2} \otimes \nu g_{7/2}]$  two-body matrix element. Apart from the above state, the three low-lying  $1^-$  levels are also of interest, because they may be populated by ff-transitions. The two high-lying  $1^+$ -states of 4QP nature at 4633 keV and 5098 keV are beyond the present model calculations, since they require core-breaking.

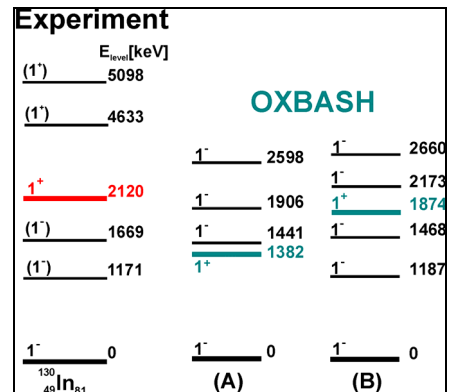


Fig. 3: Comparison of low-lying  $1^+$  and  $1^-$  levels in  $^{130}\text{In}$  with two calculations of the OXBASH shell model. Marked in red and blue are the positions of the lowest  $1^+$  level.

## Astrophysical implications

The high experimental  $Q_\beta$ -value of 8.32 MeV is another evidence for a “quenching” of the N=82 shell below  $^{132}\text{Sn}$  [5]. As a consequence of the weakening of the  $\pi\nu$ -monopole interaction, the half-lives for GT-decay of the so far unknown N=82 waiting-point nuclei  $^{128}\text{Pd}$ ,  $^{127}\text{Rh}$ ,  $^{126}\text{Ru}$  and  $^{125}\text{Tc}$  will become longer than predicted by recent shell models. This leads to a more reliable description of the build-up of the  $A \approx 130$   $N_{r,0}$  peak in r-process abundance calculations. Hence, a better understanding of the r-process matter flow through this “bottle-neck” region is derived, which also determines to a large extent the total duration of a classical r-process.

## References

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