

# How doubly-magic is the nucleus ${}^{78}_{28}\text{Ni}_{50}$ ?

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Doubly-magic nuclei play an important role in nuclear structure theory as testbeds for shell model calculations. Data on neutron-rich nuclei in the regions of doubly-magic nuclei as  ${}^{78}_{28}\text{Ni}_{50}$  and  ${}^{132}_{50}\text{Sn}_{82}$  have a decisive influence on nucleosynthesis calculations. They comprise the longer-lived waiting-point nuclei determining the duration of the r-process as well as the matter flow through the abundance maxima related to the magic neutron numbers [1].

Remaining deficiencies in the calculation of isotopic abundances have been interpreted by our group as signatures of nuclear structure near the neutron drip-line. Pronounced abundance troughs prior to the maxima have their origin in an overestimation of the neutron shell strength far from stability in global mass models such as FRDM and ETFSI-1.

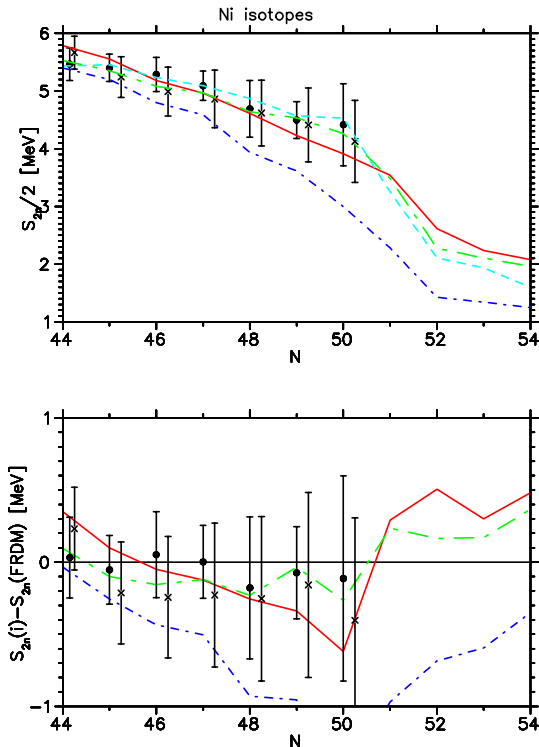


Figure 1: The two-neutron separation energies ( $S_{2n}$ ) across the shell gap at  $N=50$  for Ni-isotopes (upper part) and the differences to the FRDM model (lower part). Experimental values: crosses from 1995 mass evaluation [2] and circles from 2003 evaluation [3]. Theoretical masses: FRDM: cyan, ETFSI-1: green, ETFSI-Q: blue, HFB-2: red

A weakening ("quenching") of spherical shells with increasing isospin, resulting in a gradual setting in of collectivity has been predicted by HFB calculations for the spherical shells at  $N=50$ , 82 and 126, and is well established in the meantime for the lower neutron-magic numbers.

Only recently, first decay data on the very neutron-rich

doubly-magic nuclei  ${}^{78}_{28}\text{Ni}_{50}$  could be obtained at MSU [4]. It is of interest to determine the mutual influence of the proton and neutron magic numbers far from stability. The shell strength can be derived from the difference of the two-neutron separation energies  $S_{2n}$  prior and behind a magic neutron number. As an example, Fig. 1 displays the  $S_{2n}$  values for Ni isotopes.

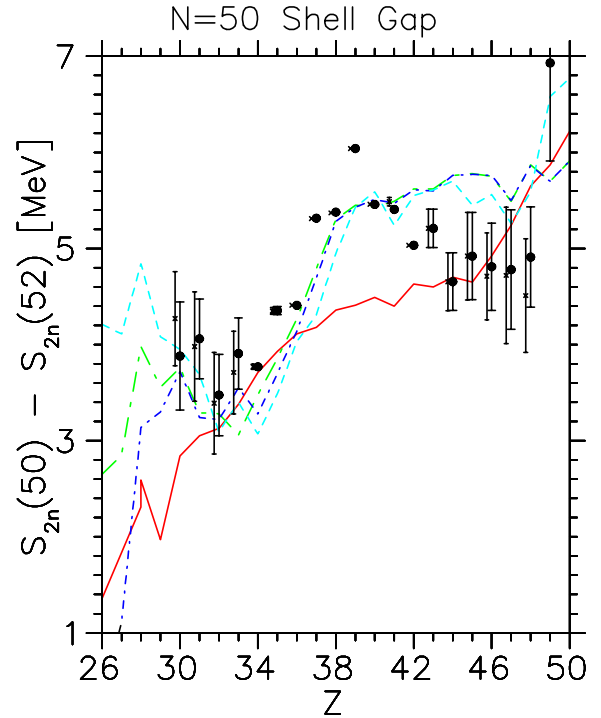


Figure 2: The  $N=50$  shell gap as a function of  $Z$ . Experimental values from the 2003 mass evaluation [2] (black and red circles) are compared to the 1995 ones [3] (crosses and black circles). [Same colour coding for the mass models as in Fig. 1.]

In Fig. 2 the  $N=50$  shell gap over a wide  $Z$  range is shown. The maximum corresponds to  ${}^{90}_{40}\text{Zr}_{50}$ , where the subshell closure at  $Z=40$  reinforces the  $N=50$  magic number. Most global mass models predict a local maximum for  ${}^{78}\text{Ni}$ . However, so far below  $Z=30$ , i.e.  ${}^{80}\text{Zn}$ , experimental masses have not been determined, so that this prediction cannot be verified. Direct mass measurements at MSU or the FRS-ESR of the GSI will in future extend the range of experimental masses in this region to more neutron-rich isotopes.

## References

- [1] K.-L. Kratz et al., *Ap. J.* **403**, 216 (1993).
- [2] G. Audi et al., *Nucl. Phys.* **A595**, 409 (1995).
- [3] G. Audi et al., *Nucl. Phys.* **A729**, 3 (2003).
- [4] P. Hosmer et al., 1<sup>st</sup> JINA-Workshop, Seattle (2004).